

Ion acceleration in nanostructured targets

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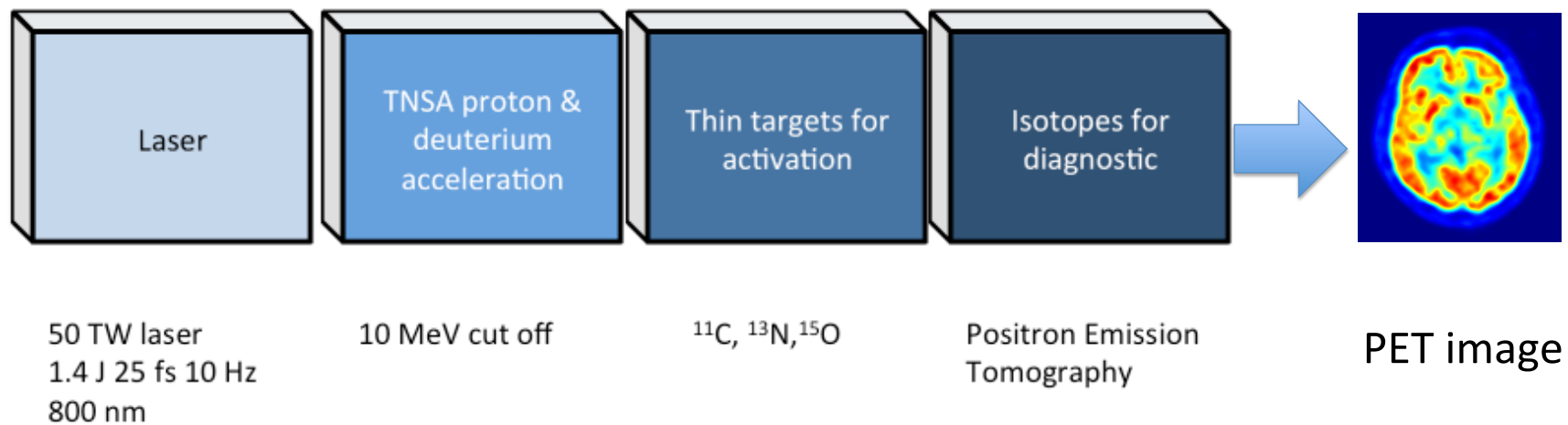


Outline

- Motivation
- Initial question
- How to estimate the prospects for PET.
- Simulations on ion acceleration in nano structured targets
- Conclusions

L2A2: Laser Laboratory for Acceleration and Applications at the USC

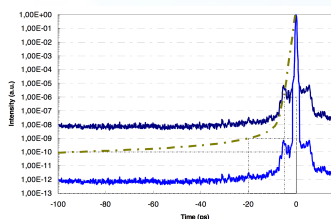
- Objectives:
 - Establish the new lab
 - Do basic science in laser plasma acceleration
 - Try to push new technology for radioisotope production from laser driven ions.
 - One of the important applications of laser driven plasma accelerators is the use of protons to produce radioisotopes.
 - The production of short-lived isotopes such as ^{11}C or ^{18}F is important in medicine for positron-emission tomography (PET).



Laser system @ L2A2

Peak power	45 TW	30 GW
Nominal Energy	1.4 J	1 mJ
Pulse duration	<25 fs	<35 fs
Central wavelength	800 nm +/- 10 nm	800 nm +/- 10 nm
Repetition rate	10 Hz	1 KHz

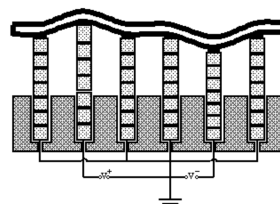
XPW: 10^{11}
ps contrast



Dazzler: Spectral
control



High power
Adaptive Mirror



Beam pointing
stability for
underdense
experiments

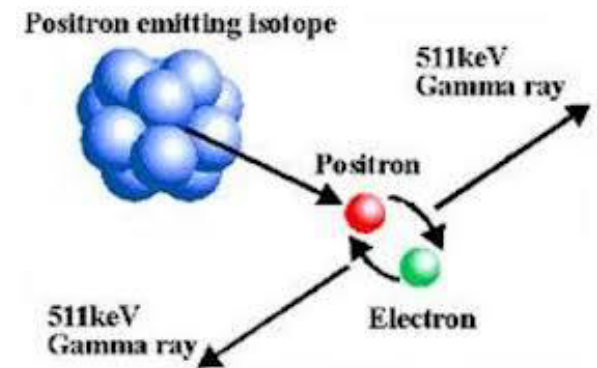
5μrad rms 500shots

Energy stability <
1.5 % rms
Strehl ratio >0.6
Divergence 15
urad

Motivation: Positron Emission Tomography (PET)

PET is considered the most sensitive 3D imaging technique:

- Positrons emitted by specific β^+ radionuclides annihilate producing two photons
- The back-to-back emission of pairs of photons from a given volume made it possible its reconstruction using appropriate image reconstruction algorithms



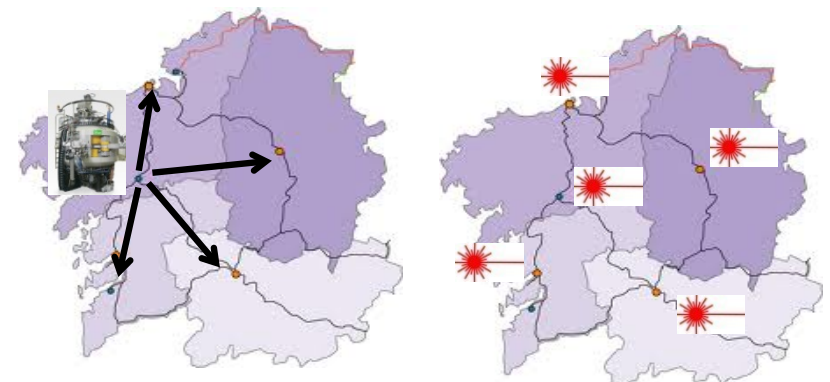
The production of radionuclides requires an accelerator with a complex infrastructure :

- The production strategy is based on a regional production center and distribution.
- Short-lived emitters can only be used close to the production center

Laser accelerators may represent an option:

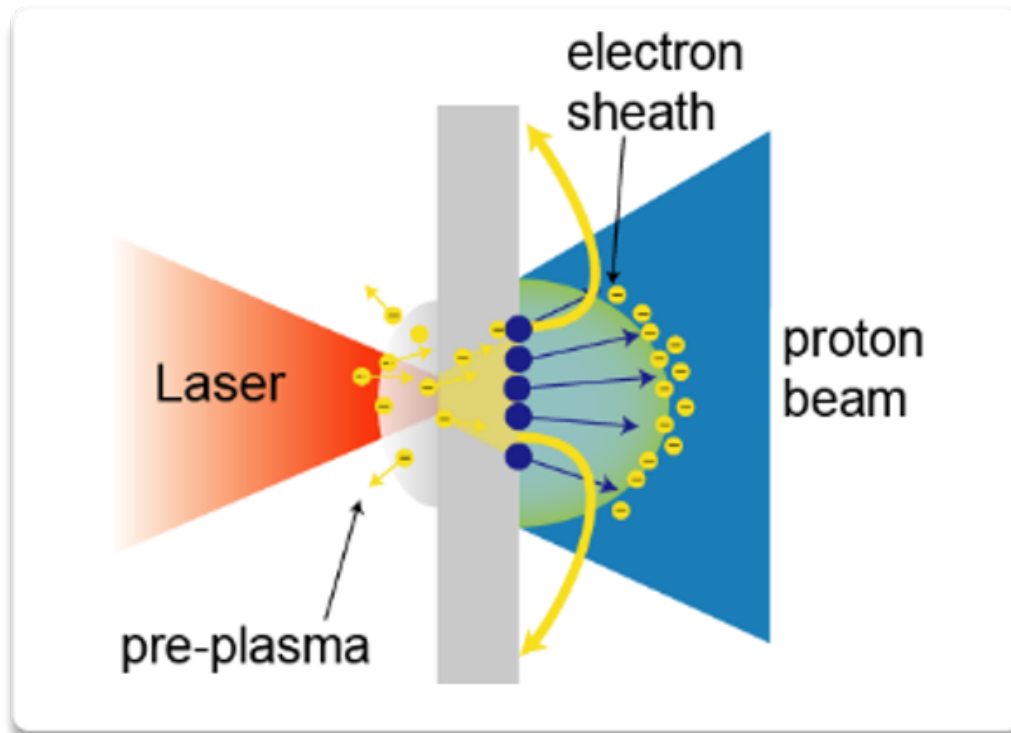
Simpler and cheaper infrastructure.

Compact device that could be installed in any hospital for the production of short-lived radioisotopes



^{11}C	^{13}N	^{15}O	^{18}F
20 min.	10 min.	2 min.	110 min.

Target Normal Sheath Acceleration (TNSA)



Characteristics

- Broad spectra
- High number of protons: 10^{10} - 10^{13}
- Low emittance: 4×10^{-3} mm/mrad
- Divergence 10-20 degrees
- Ultrashort 0.1-10 ps

Snively et al, PRL **85** (2000) 2945

Other observations:

Clark et al, PRL **84** (2000) 670

Maksimchuk et al, PRL **84** (2000) 4108

Ion acceleration by superintense laser-plasma interaction

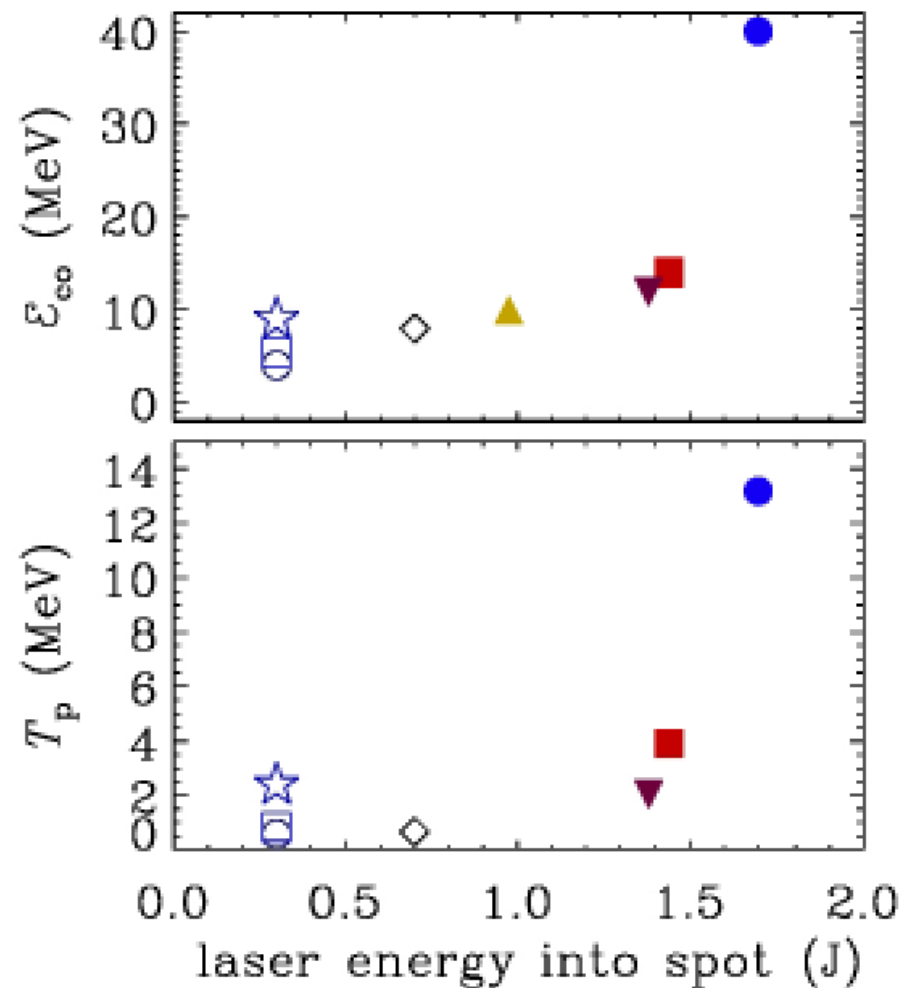
Andrea Macchi, Marco Borghesi, and Matteo Passoni

Rev. Mod. Phys. 85, 751

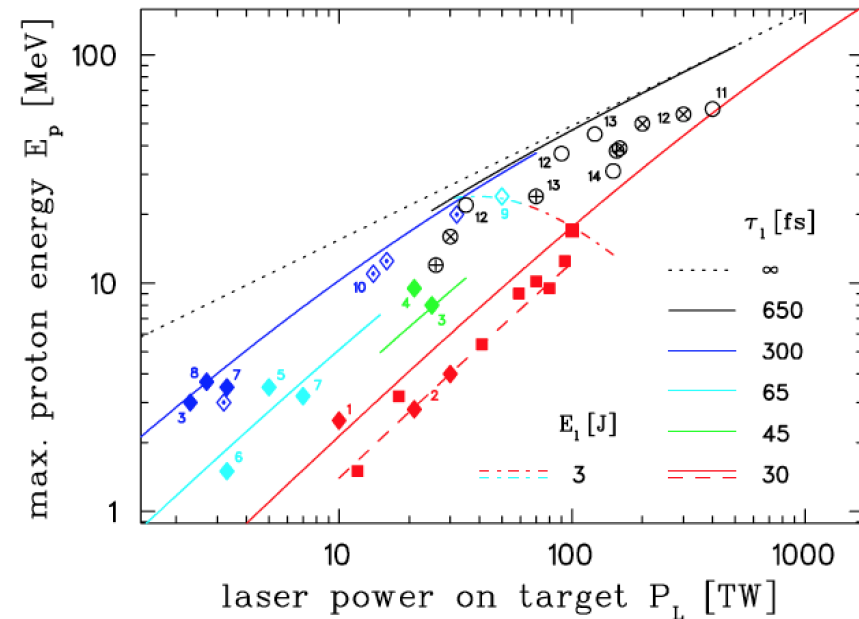
Why TNSA?

- The most studied phenomena
- Achievable with $<2\text{J}$ per pulse
- Micron thick targets instead of nm thick targets
- Several routes for optimization described
- A broad spectrum is useful depending on the cross section
- ▶ progress should be monitored on small-scale lasers with potential for high repetition rate and cost-effective applications

TNSA Protons: Laser energy below 2J



A. Macchi



K Zeil, et al, NJP 12 (2010) 045015 (16pp)

- **Overdense.**
- **TNSA**
- **10 um Al foils**
- **Single shot**
- **10 MeV protons**

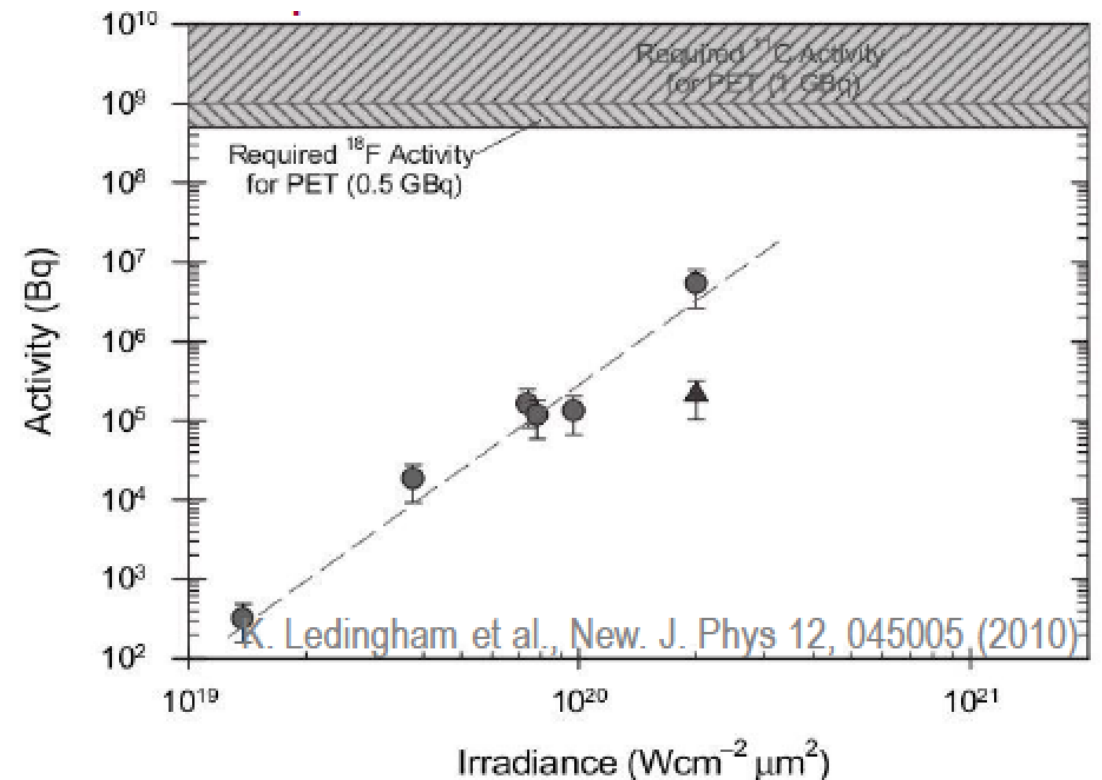
Initial question.

- Would it be possible to produce radioisotopes for medical imaging such as PET from laser driven ions?

The initial question:

Activation by laser produced protons

- 120 J, ~1 ps (10^{20} W/cm²), CR $\sim 10^6$
 ^{11}C activity/shot ~ 200 kBq
I. Spencer et al., NIMB 183, 449 (2001)
- 20-30 J, 0.3 – 0.8 ps ($1-6 \cdot 10^{19}$ W/cm²), CR $< 10^6$
 ^{11}C activity/shot ~ 1 MBq
J. Fuchs et al., PRL 94, 045004 (2005)
- 0.8 J, 40 fs ($6 \cdot 10^{19}$ W/cm²), CR $< 10^6$
 ^{11}C activity/shot ~ 1.2 kBq
S. Fritzler et al., App. Phys. Lett. 83, 3039 (2003)



I. Spencer, et al, Nuclear Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms 183 (2001) 449.
S. Fritzler, et al, Appl. Phys. Lett. 83 (2003) 3039.
K. W. D. Ledingham, et al, J. Phys. D: Appl. Phys. 37 (2004) 2341.
L. Robson, et al, High-power laser production of pet isotopes, Lecture Note in Physics 694 (2006) 191.

One estimation from 2006

JOURNAL OF APPLIED PHYSICS **100**, 113308 (2006)

Numerical simulation of isotope production for positron emission tomography with laser-accelerated ions

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*Département de Physique Théorique et Appliquée, CEA/DAM Ile-de-France, BP 12,
91680 Bruyères-le-Châtel, France*

Emmanuel d'Humières

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University of Nevada, Reno, Nevada 89557*

Sven Fritzler

Siemens Medical Solutions, Vacuum Technology, 91052 Erlangen, Germany

Victor Malka

*Laboratoire d'Optique Appliquée, ENSTA, UMR 7639 CNRS/Ecole Polytechnique,
91761 Palaiseau, France*

$$I = 10^{20} \text{ W/cm}^2$$

^{11}C	1.1×10^8	per shot	94.8 Gbq	1Khz	1hr
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^{18}F	3.1×10^7	per shot	9.7 Gbq	1Khz	1hr
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$$I = 4 \times 10^{20} \text{ W/cm}^2$$

^{18}F	104 Gb (2.8 Ci)	1Khz	1hr
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New estimation 2016

Study of the production yields of ^{18}F , ^{11}C , ^{13}N and ^{15}O positron emitters from plasma-laser proton sources at ELI-Beamlines for labeling of PET radiopharmaceuticals

Ernesto Amato^a, Antonio Italiano^{b,*}, Daniele Margarone^c, Benedetta Pagano^d, Sergio Baldari^{a,d}, Georg Korn^c

^a Section of Radiological Sciences, Department of Biomedical and Dental Sciences and of Morphologic and Functional Imaging, University of Messina, Italy

^b Istituto Nazionale di Fisica Nucleare, Gruppo Collegato di Messina, Italy

^c Institute of Physics ASCR, v.v.i. (FZU), ELI-Beamlines Project, 182 21 Prague, Czech Republic

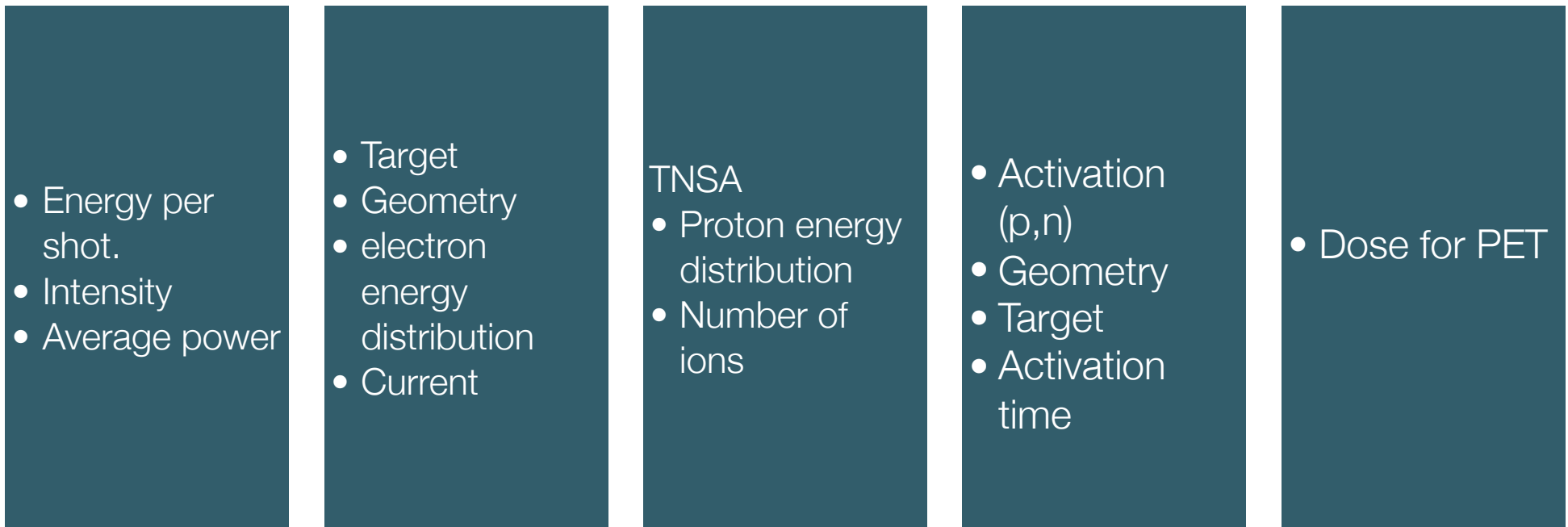
^d Nuclear Medicine Unit, University Hospital "G. Martino", Messina, Italy

Table 2

Total radionuclide activities produced at the end of bombardment (EOB), for different laser repetition rates and irradiation times, for the three proton spectra, according to Table 1.

Nuclide	$T_{1/2}$ (min)	A^{CYCL} (MBq)	$T_{\text{irr}}^{\text{CYCL}}$ (min)	Rep. rate (Hz)	1			5			10		
					T_{irr} (min)								
					15	60	120	15	60	120	15	60	120
^{18}F	110	52000 (F^-)	60	spectrum	Total Activity at EOB (MBq)								
				low	1.8	6.4	10.8	9.1	31.9	53.8	18.3	63.8	107.5
				medium	11.4	39.6	66.8	56.8	198.2	333.9	113.5	396.3	667.8
				high	54.4	189.8	319.9	271.9	949.2	1599.6	543.9	1898.4	3199.1
^{11}C	20	36000 (CO)	50	spectrum	Total Activity at EOB (MBq)								
				low	2.0	4.4	5.0	10.2	22.1	24.8	20.4	44.1	49.6
				medium	32.1	69.4	78.1	160.7	346.9	390.3	321.4	693.8	780.5
				high	198.6	428.6	482.2	993.0	2143.2	2411.1	1985.9	4286.4	4822.2
^{13}N	10	4000 (NH_3)	10	spectrum	Total Activity at EOB (MBq)								
				low	0.1	0.1	0.1	0.4	0.6	0.6	0.7	1.1	1.1
				medium	44.3	67.5	68.6	221.6	337.5	342.8	443.3	675.0	685.5
				high	502.8	765.7	777.7	2514.2	3828.5	3888.3	5028.4	7657.0	7776.6
^{15}O	2	74000 (O_2)	10	spectrum	Total Activity at EOB (MBq)								
				low	5.7	5.7	5.7	28.5	28.7	28.7	57.0	57.3	57.3
				medium	7.5	7.6	7.6	37.5	37.8	37.8	75.1	75.5	75.5
				high	43.3	43.5	43.5	216.3	217.5	217.5	432.6	435.0	435.0

Diagram of the estimation



Laser

Electrons

Protons

Activation

Dose

Diagram of the estimation.

- The estimation compare against the production of a cyclotrons.
- 2006. In order to compete with a regular cyclotron it is necessary to have **1PW, 1Khz**.
- 2016. *“In principle, the feasibility to produce **clinically relevant amounts of positron emitters, to be used for inline preparation of single doses of radiopharmaceuticals**,..., peak power (**PW**) and high repetition rate (**10 Hz**)”*.
- Consider a broader perspective.

Broader considerations for the estimation.

- Energy per shot.
- Intensity
- Average power

- Target
- Geometry
- electron energy dist
- Current

TNSA

- Proton energy dist
- Number of ions

- Activation (p,n)
- Geometry
- Target
- Activation time

- Dose for PET

Broader considerations for the estimation.

- Installation cost
- Complexity

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- Complex targets

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- Advance proton diag
- Online beam monitoring

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- Dose for PET

- Radiopharmacy cost

- Advance imaging technique

Broader considerations for the estimation.

Overall cost & complexity

- Installation cost
- Complexity

- Energy per shot.
- Intensity
- Average power

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Broader considerations for the estimation.

Overall cost & complexity

New paradigm for radioisotope production

- Installation cost
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- Dose for PET

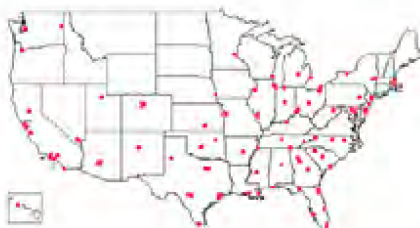
- Radiopharmacy cost

- Advance imaging technique

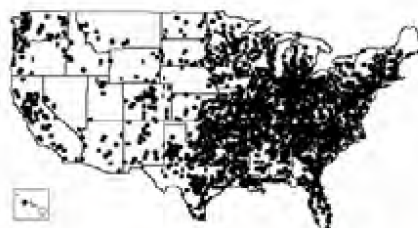
High repetition rate

New paradigm for radio isotope production

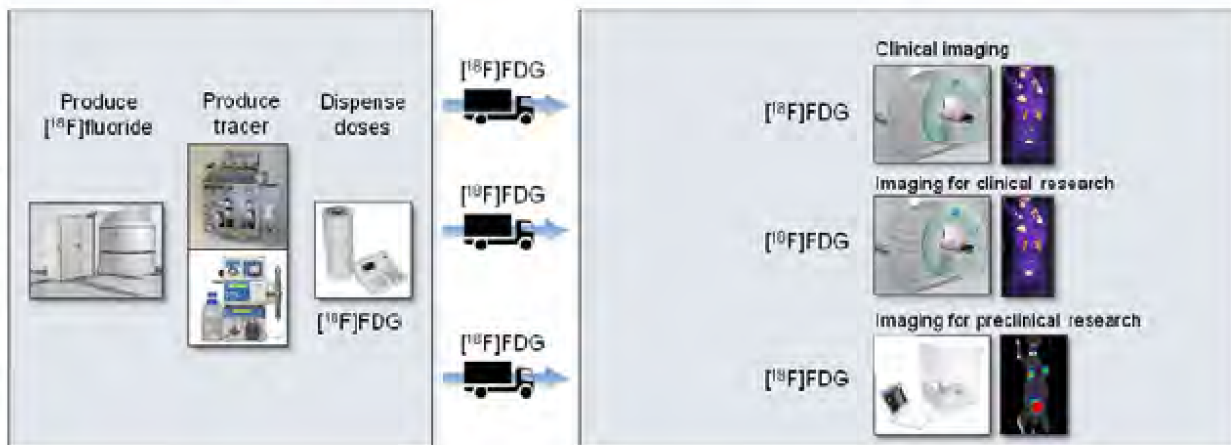
PET Radiopharmacies



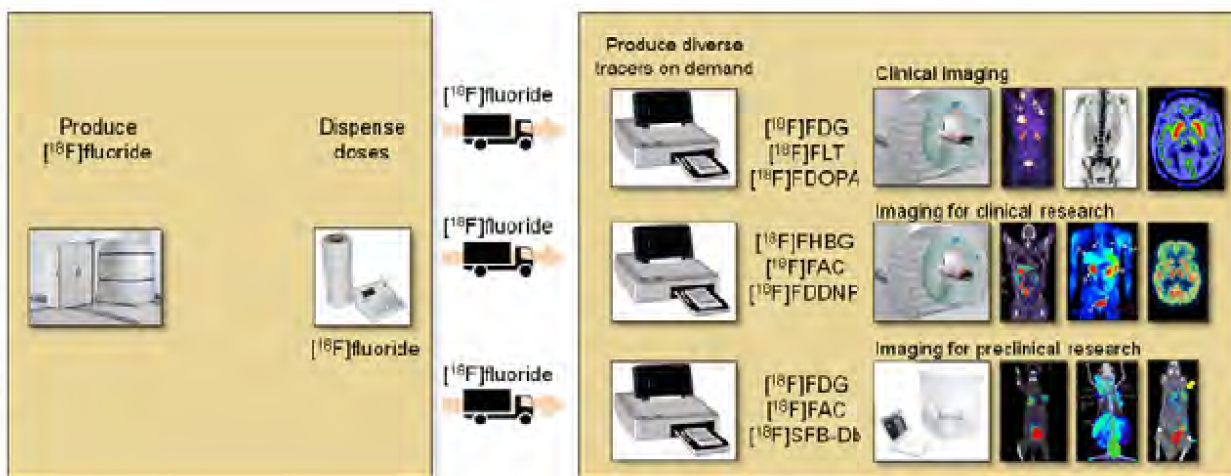
PET Centers



Centralized Production



Decentralized Production



Emerging Technologies for Decentralized Production of PET Tracers

Pei Yuin Keng¹, Melissa Esterby^{1,2} and R. Michael van Dam¹
¹Crumpp Institute for Molecular Imaging, Department of Molecular & Medical Pharmacology, David Geffen School of Medicine, University of California, Los Angeles
²Sofie Biosciences, Inc.
 USA

Overall cost & complexity

- **Large Cyclotrons**

- Commercial installations are large facilities
- A special building
- Designed to produce a large number of doses
- Cost effective use depend on number of PET scanners
- Strong restriction on research related to new pharmaceuticals in preclinical
- Short lifetime (110 min) implies dose will undergo multiple half life of decay before use
- Careful planning
- Only FDG

- **Laser**

- Not a special building, limited shielding
- Lowest possible energy
- Adjusted to single dose production
- Minimal down time
- Radiopharmacy tuned for single dose (Lab on chip)
- Fully automated extraction

Overall cost & complexity

- **Large Cyclotrons**

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- Only FDG

- **Laser**
- Not a special building, limited shielding



fully automated
extraction

New diagram for the estimation.

Overall cost & complexity

New paradigm for radioisotope production

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- Dose for PET

- Radiopharmacy cost

- Advance imaging technique

High repetition rate

Nano and micro structured target

State of the art

Microstructuring the front of the target surface enhances energy absorption and improves ion acceleration [1-7].

The increased absorption creates more energy available for the particles to be accelerated.

- [1] Phys. Plasmas 18(10) (2011)
- [2] Contrib. Phys. Plasmas 53(2), 173-178 (2013)
- [3] Plasma Phys. and Control. Fusion 58, 014038 (2016)
- [4] New J. Phys. 13, 053028 (2011)
- [5] Phys. Rev. Lett. 109, 234801 (2012)
- [6] Phys. Rev. Lett., 110, 215004 (2013)
- [7] Nature Lett. 439(26), 445-448 (2006)

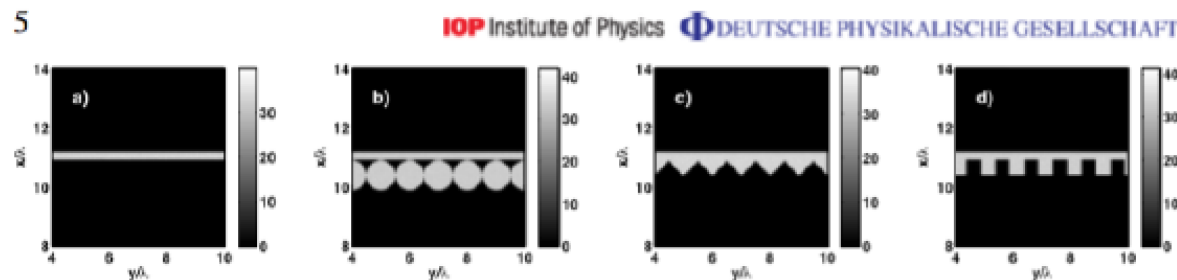


Image from [4]

Target (figure 1)	Electron temperature (MeV)	Electron divergence ($^{\circ}$)	Laser absorption (%)
(a)	0.10	14.8	3.8
(b)	0.40	39.7	55.2
(c)	0.42	41.8	80.5
(d)	0.37	40.9	43.9

Triangular shapes

Why a triangular shape?

Proton cutoff and energy absorption are enhanced compared to others [4].

There are prospects of using easier-to-make tilted triangles: possibly better for oblique incidence experiments.

Experimental limitations

Structure height below a wavelength and below bulk thickness

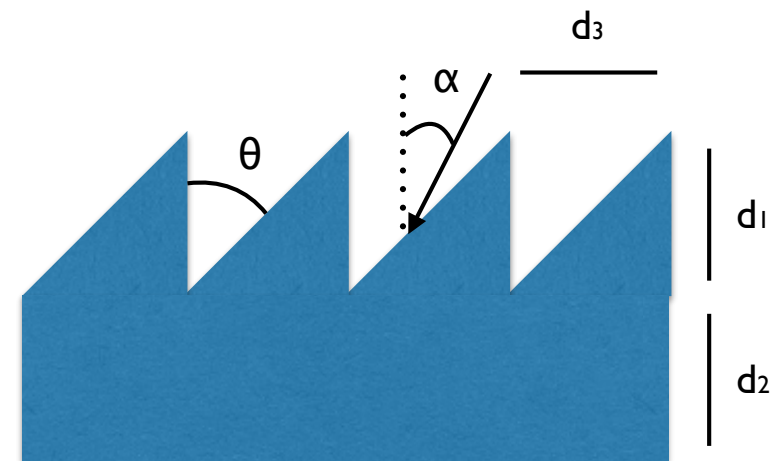
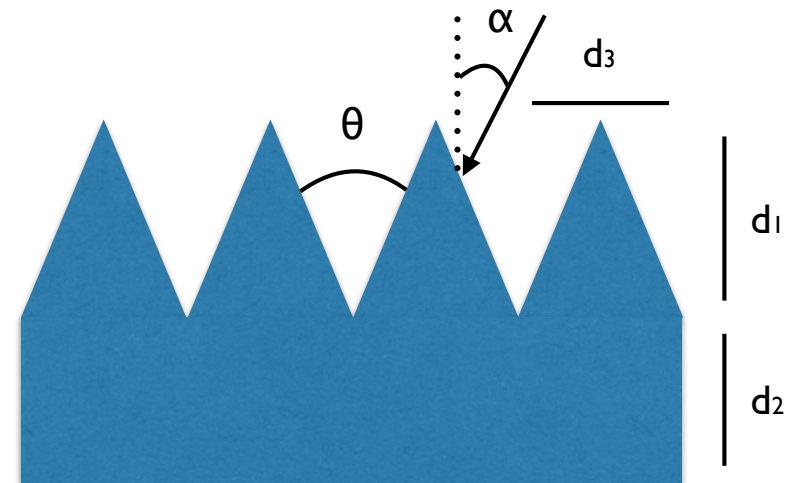
Parameters to vary

- The structure height
- The triangle width
- The bulk thickness
- The angle of incidence

[4] New J. Phys. 13, 053028 (2011)

Shapes

The angle of incidence of the laser is represented by α .



Triangular shapes

Why a triangular shape?

Proton cutoff and energy absorption are enhanced compared to others [4].

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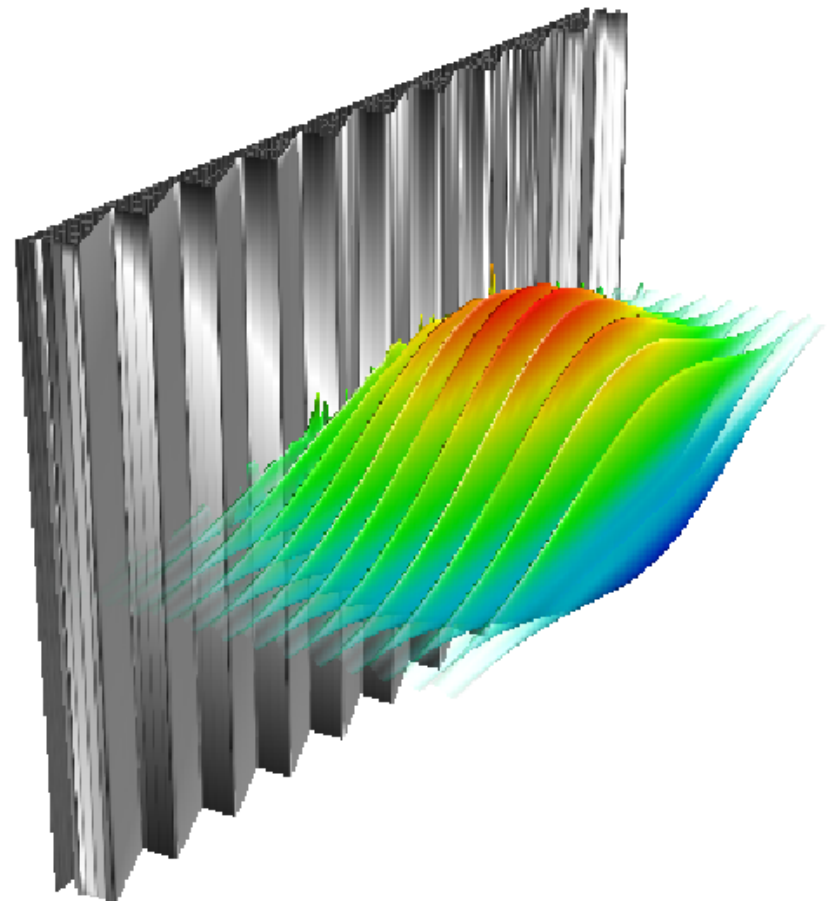
Structure height below a wavelength and below bulk thickness

Parameters to vary

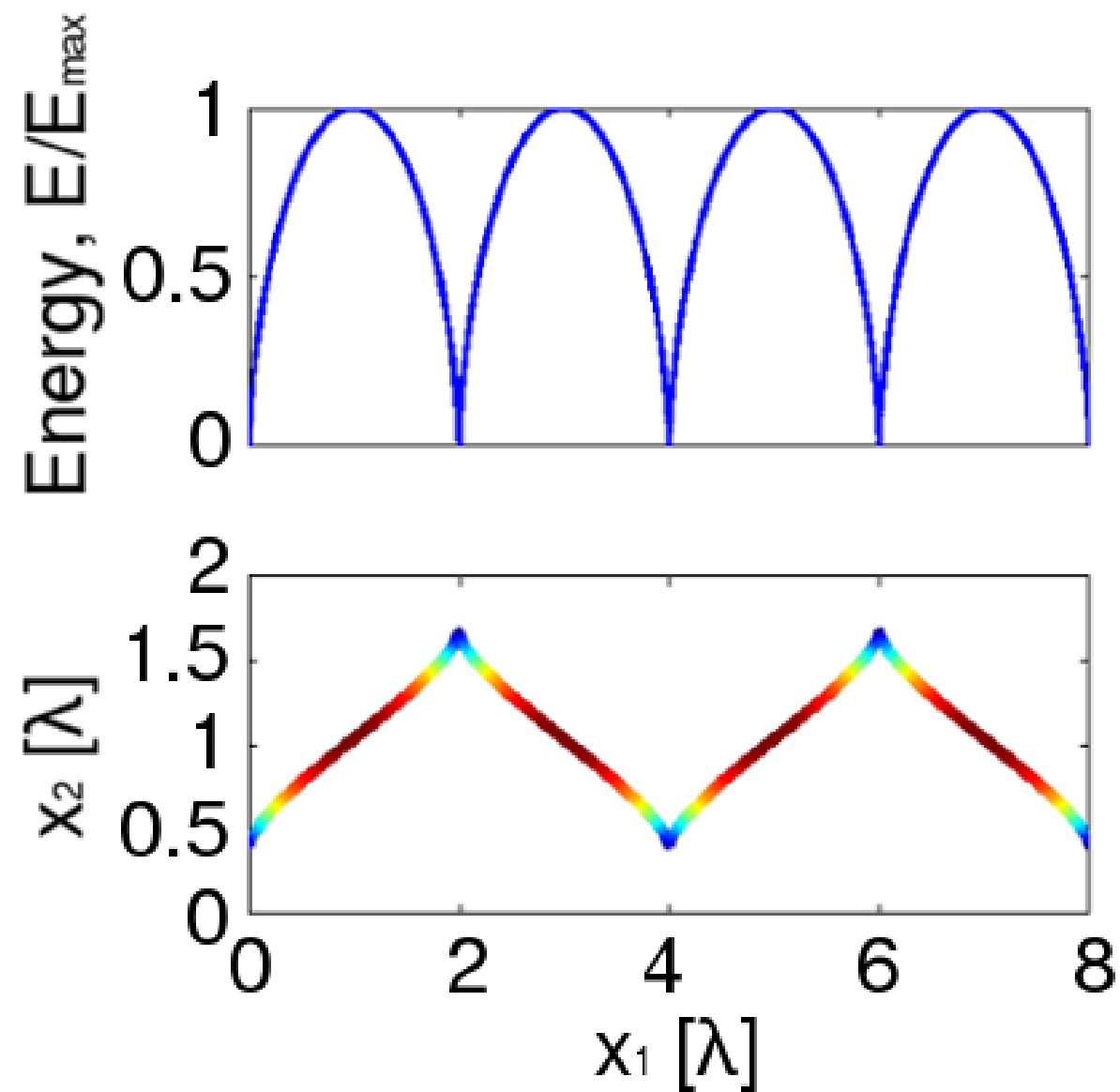
- The structure height
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Shapes

The angle of incidence of the laser is represented by α .

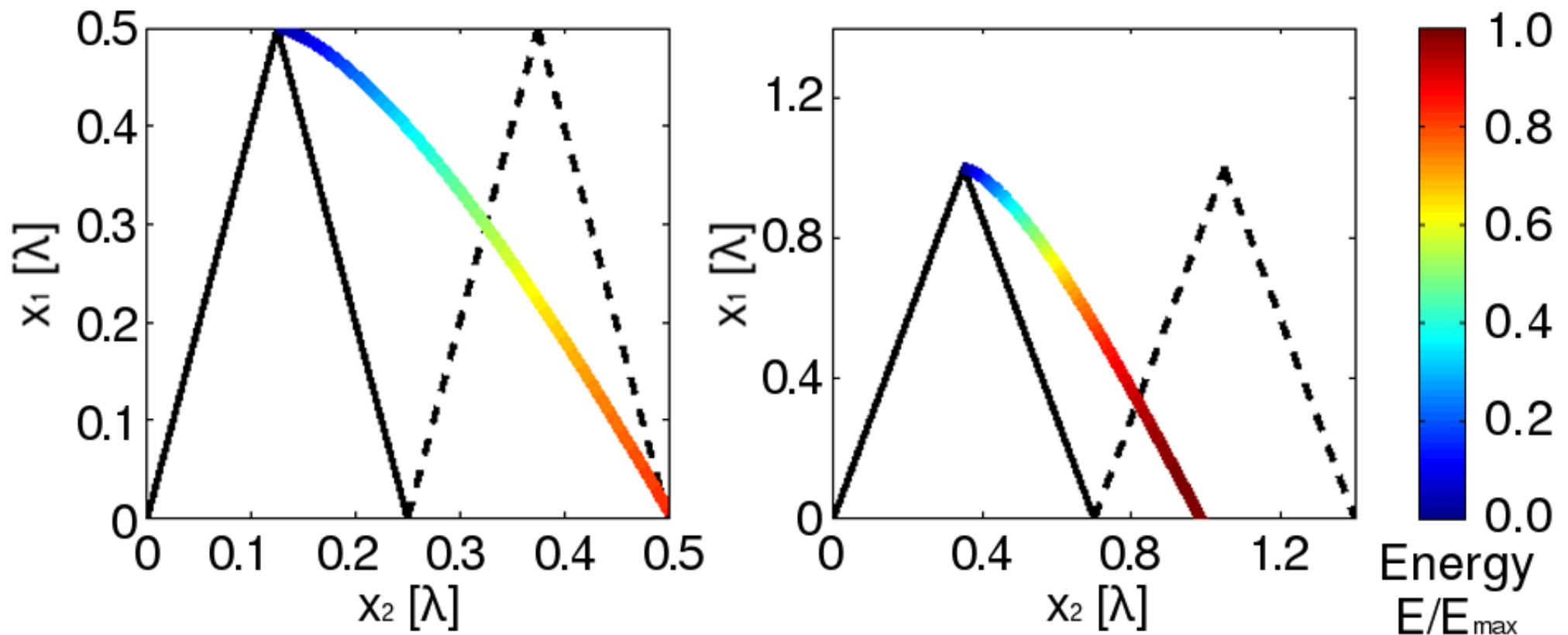


Simple idea



Simple idea

- Use the geometry of the target to optimize the energy delivered by the electron to the target

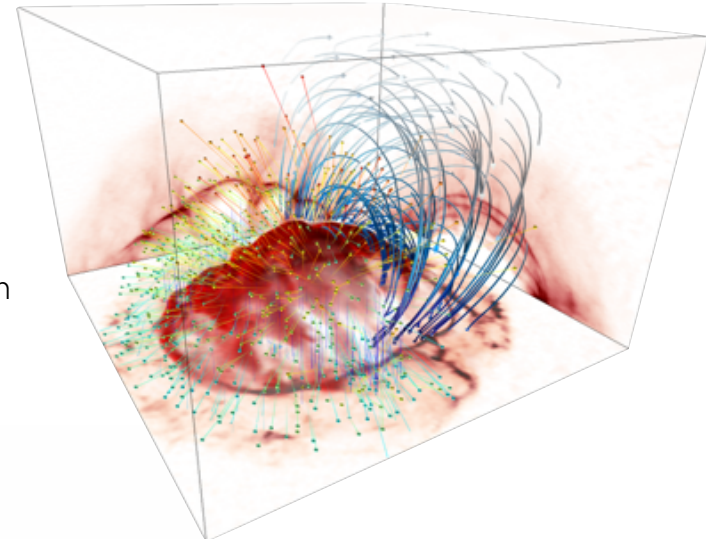


OSIRIS PIC code



osiris framework

- Massively Parallel, Fully Relativistic Particle-in-Cell (PIC) Code
- Visualization and Data Analysis Infrastructure
- Developed by the osiris.consortium
⇒ UCLA + IST



UCLA

Ricardo Fonseca

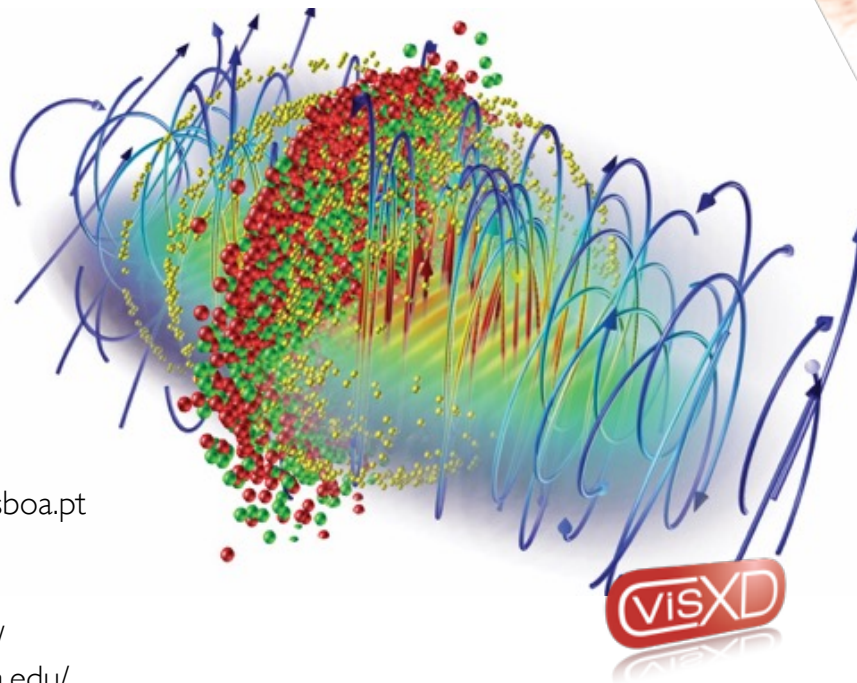
ricardo.fonseca@tecnico.ulisboa.pt

Frank Tsung

tsung@physics.ucla.edu

<http://epp.tecnico.ulisboa.pt/>

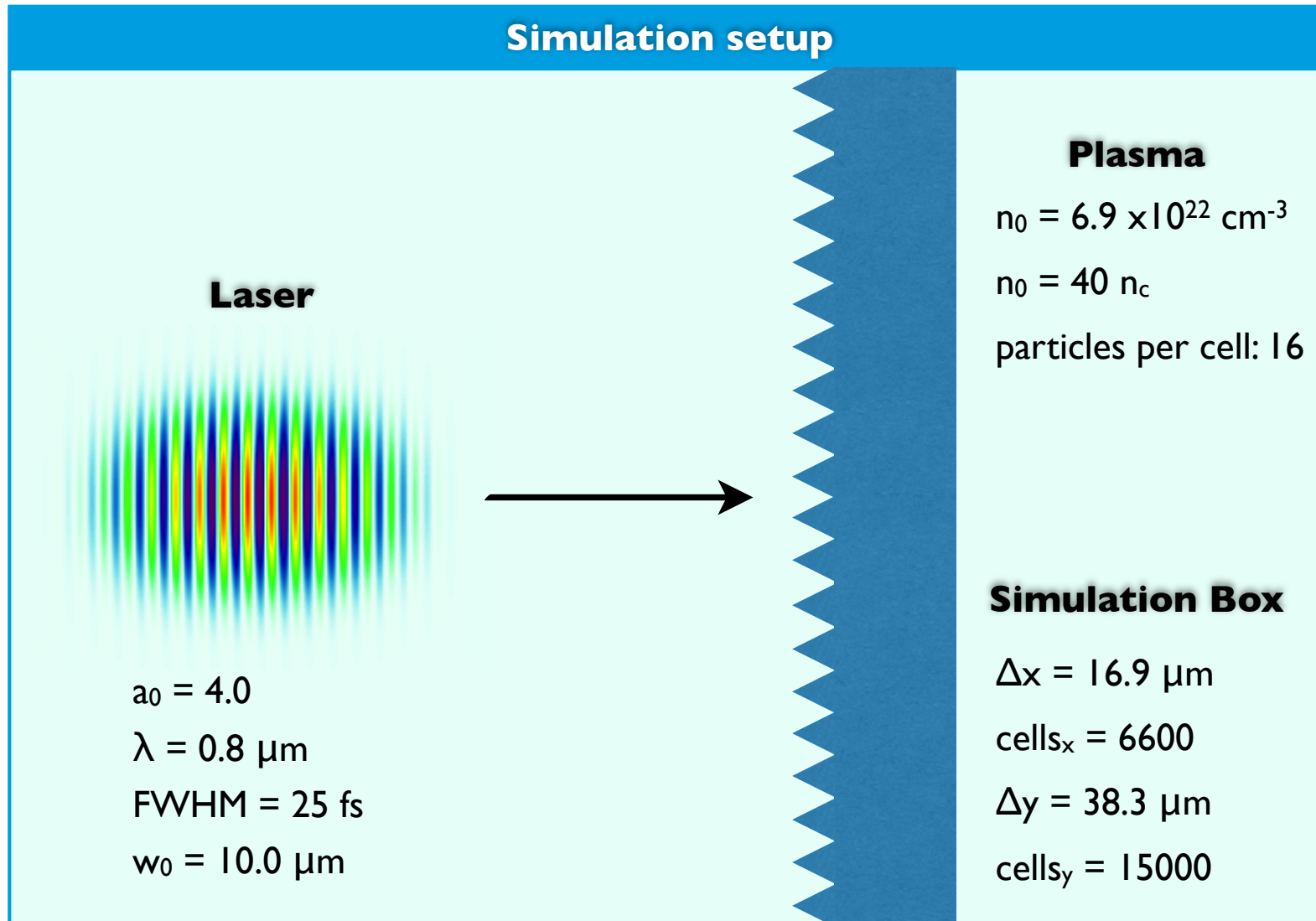
<http://plasmasim.physics.ucla.edu/>



code features

- Scalability to ~ 1.6 M cores
- SIMD hardware optimized
- Parallel I/O
- Dynamic Load Balancing
- Particle merging
- GPGPU support
- Xeon Phi support
- QED Module

Simulations in 2D for parameter scan



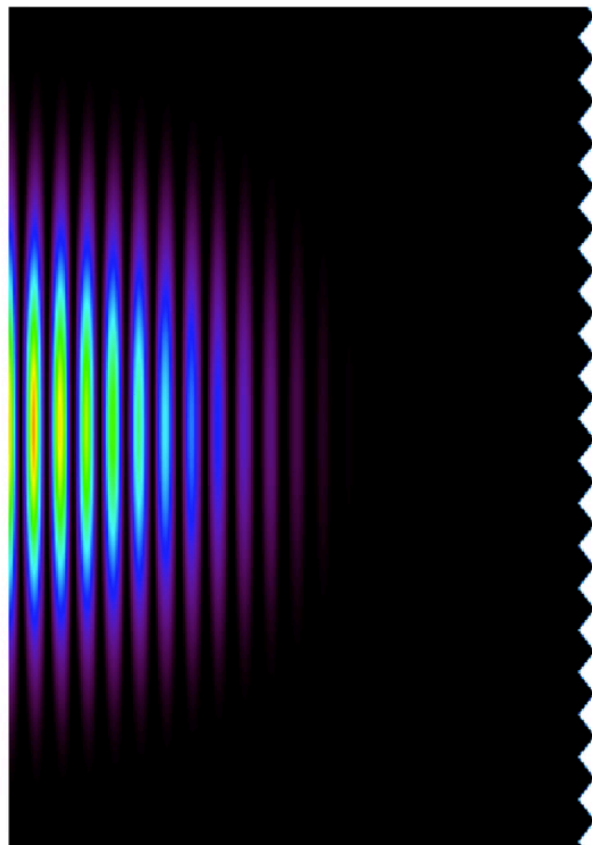
Different structures mean different absorption



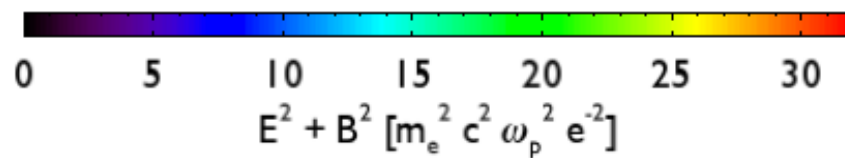
80.5% absorption

$d_1 = 0.25\lambda$

H



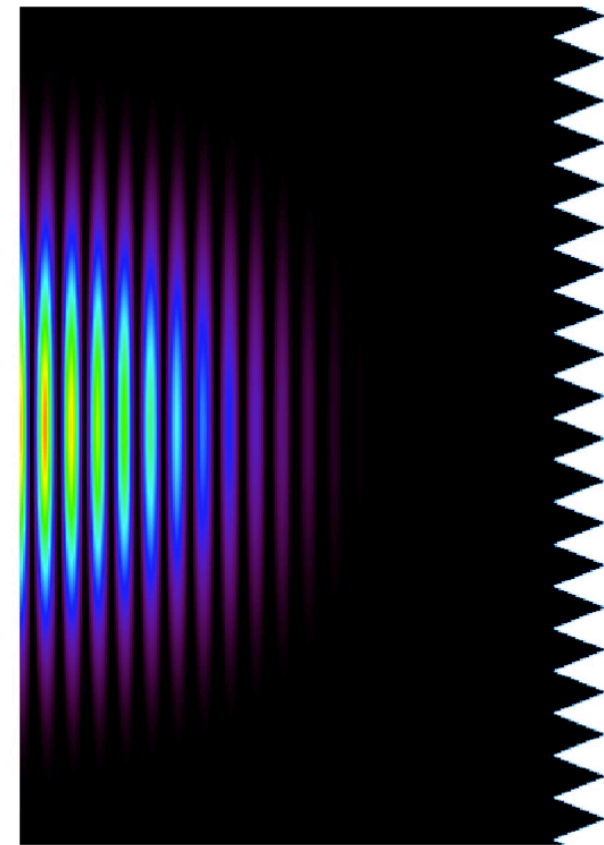
I $d_3 = 0.7\lambda$



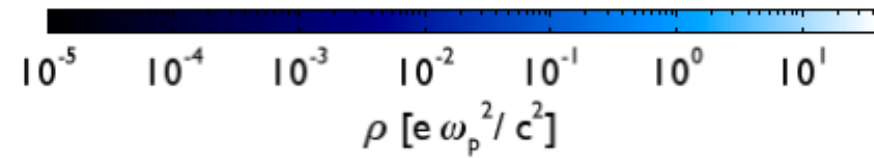
96.7% absorption

$d_1 = \lambda$

H



I $d_3 = 0.7\lambda$



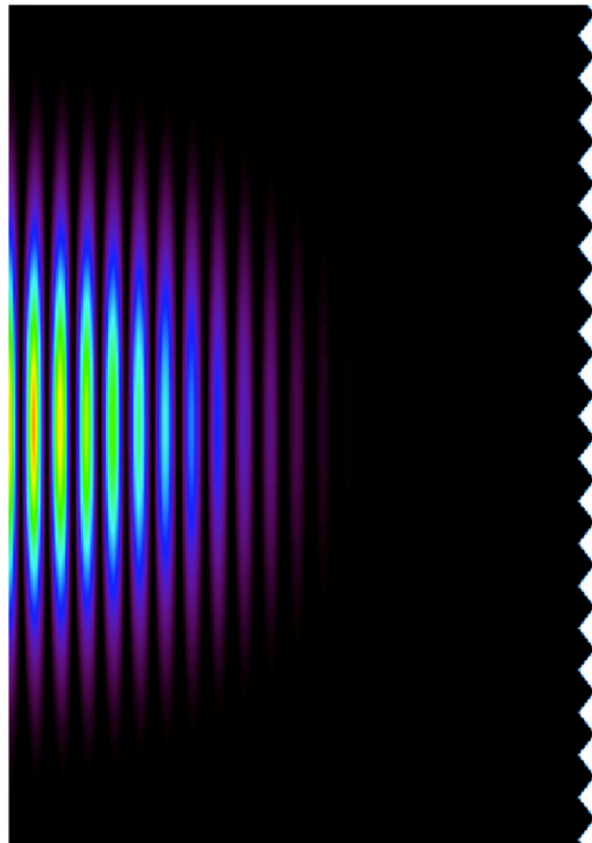
Different structures mean different absorption



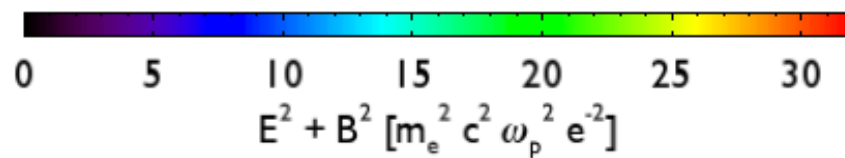
80.5% absorption

$d_1 = 0.25\lambda$

H



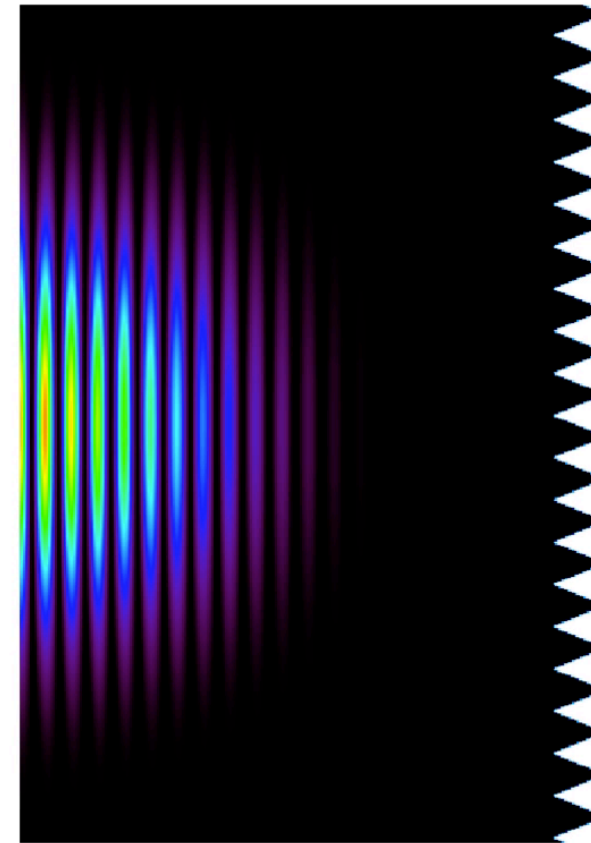
I $d_3 = 0.7\lambda$



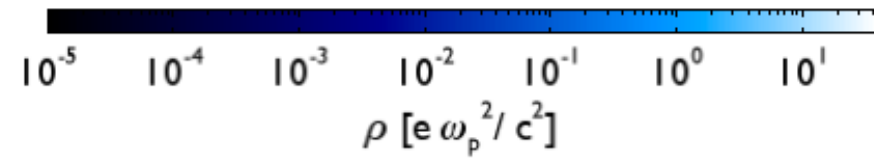
96.7% absorption

$d_1 = \lambda$

H

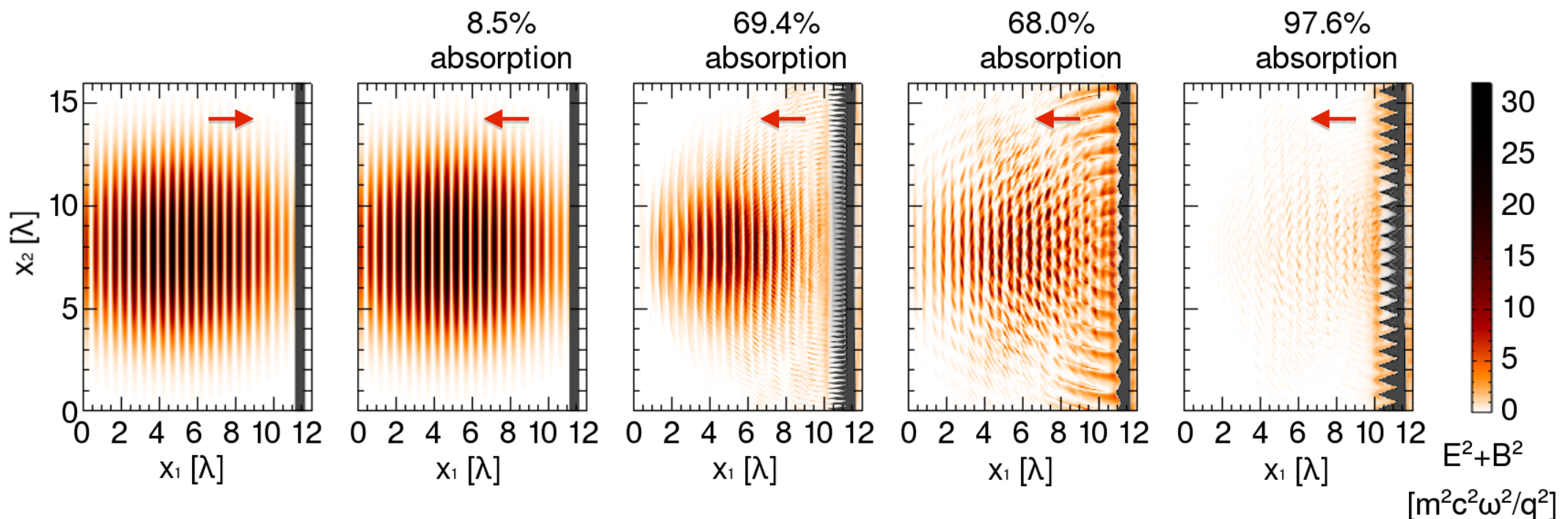


I $d_3 = 0.7\lambda$



Laser absorption

- Reflected field for different parameters of the structure



$$a_0 = 4.0$$

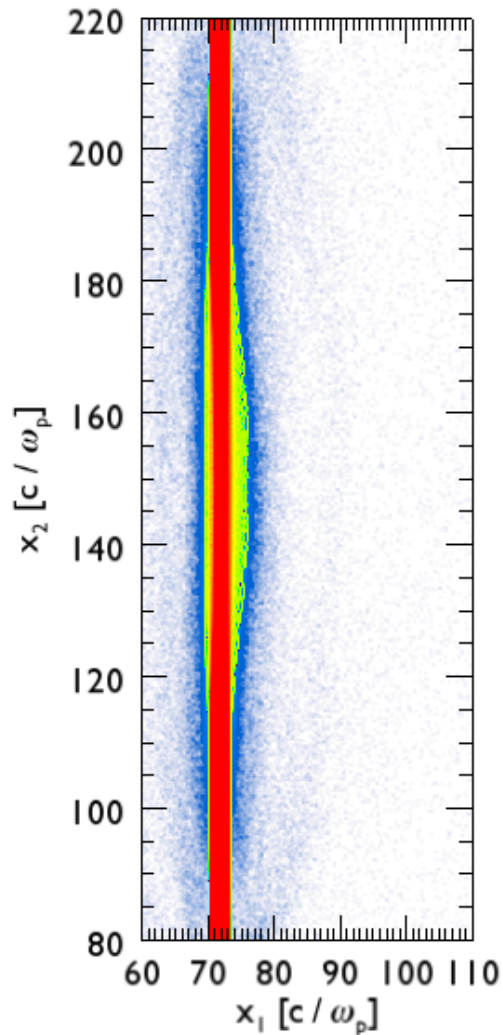
$$\lambda = 0.8 \mu\text{m}$$

$$\text{FWHM} = 25 \text{ fs}$$

$$w_0 = 10.0 \mu\text{m}$$

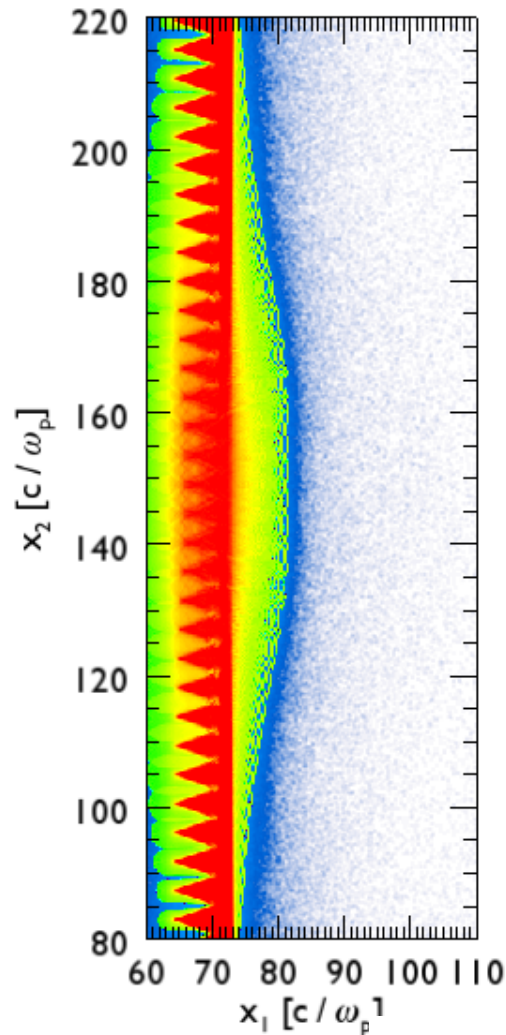
Mechanism for proton acceleration

Flat surface



$$d_2 = 0.5\lambda$$

Structured surface



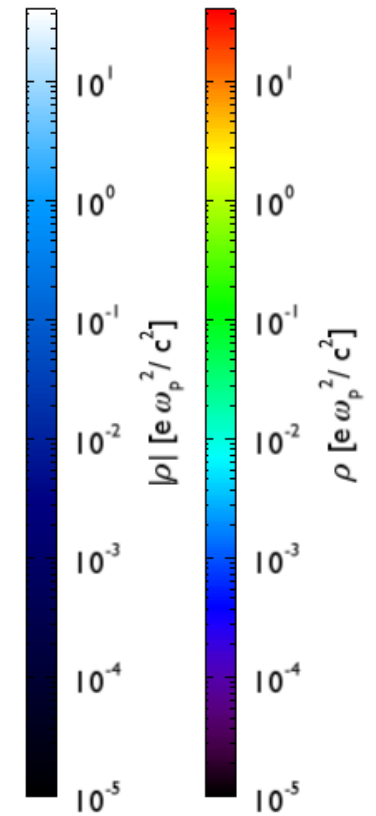
$$d_1 = \lambda$$

$$d_2 = 0.5\lambda$$

$$d_3 = 0.7\lambda$$

electron

proton



Effect of the width of the triangular structure

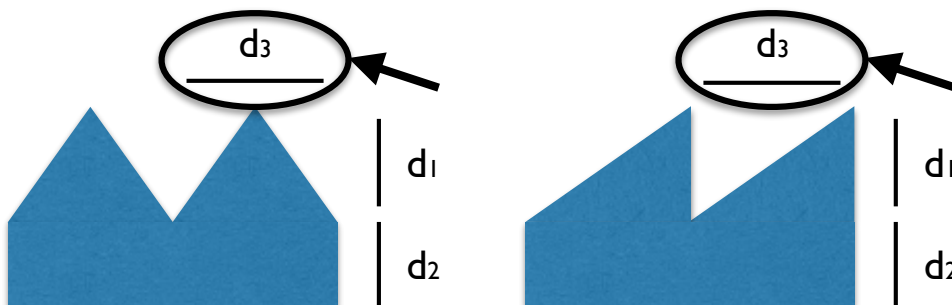
Effect of the width

- If $d_3 \ll \lambda$: The structures become invisible.
- If $d_3 \gg \lambda$: The surface becomes flat.

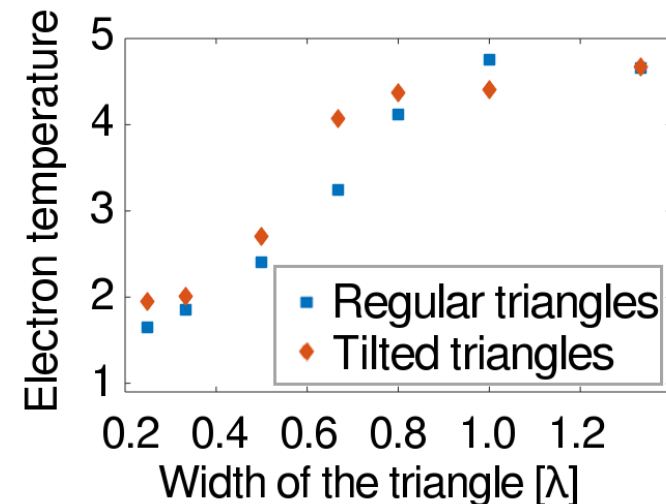
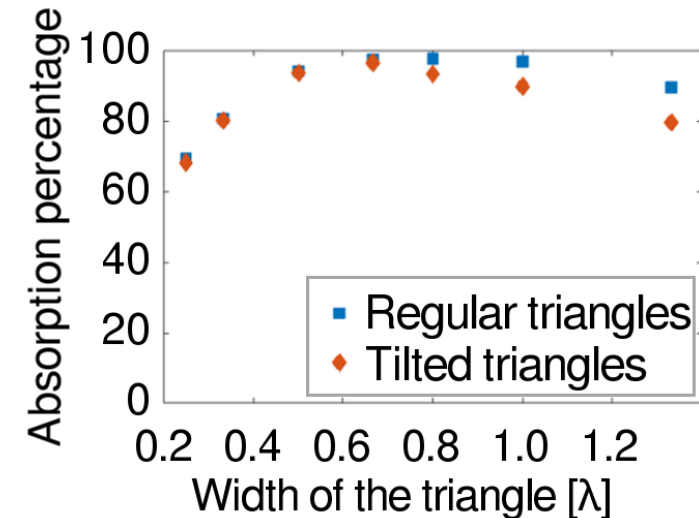
Model and data show maximum absorption for a width:

$$d_{3,\max} = 2a_0\lambda(\pi n_e)^{-1/2} = 0.7\lambda$$

$$d_2 = 0.5\lambda$$

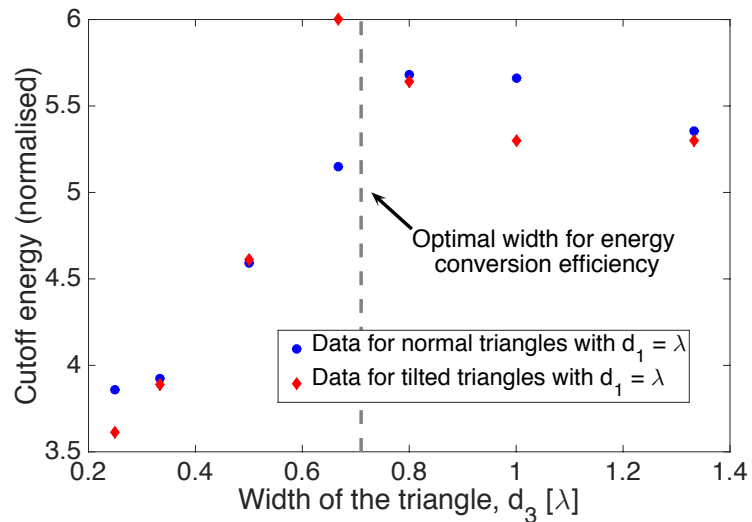


Absorption/Electron energy

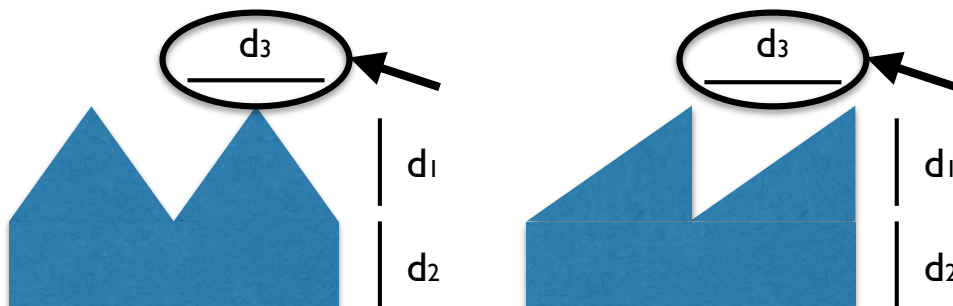


Effect of the width of the triangular structure

Proton cutoff energy ($t=405/w$)

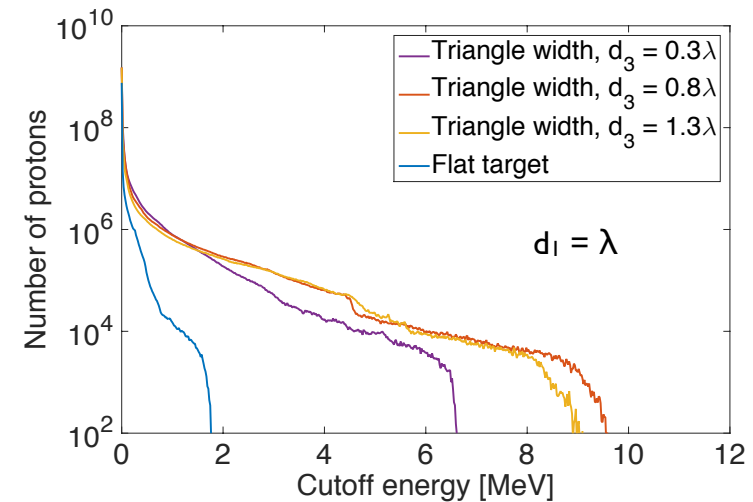


$$d_2 = 0.5\lambda$$

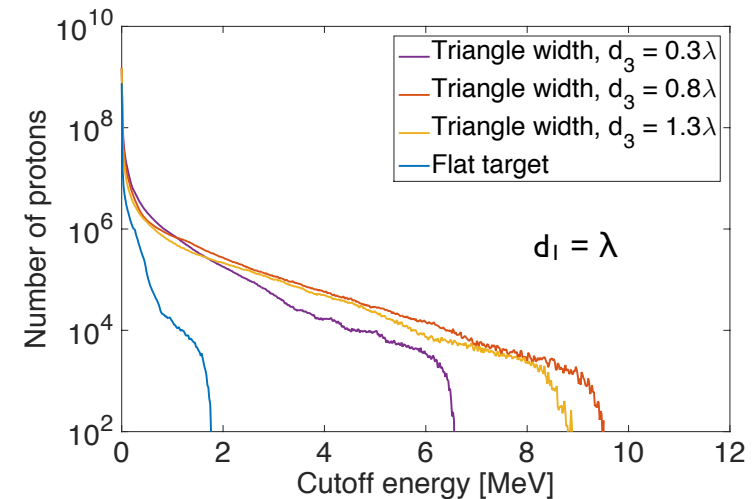


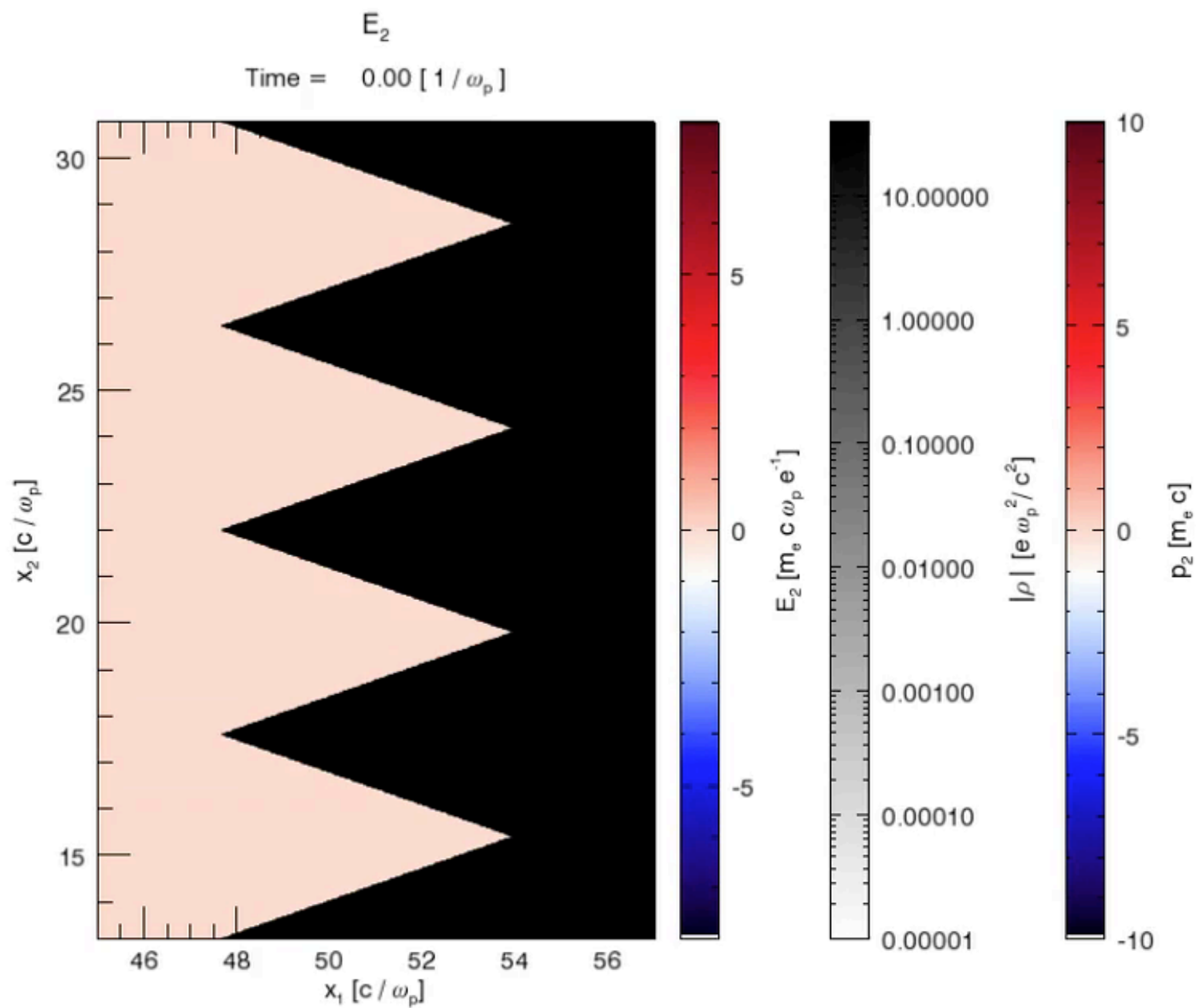
Proton spectra ($t=405/w$)

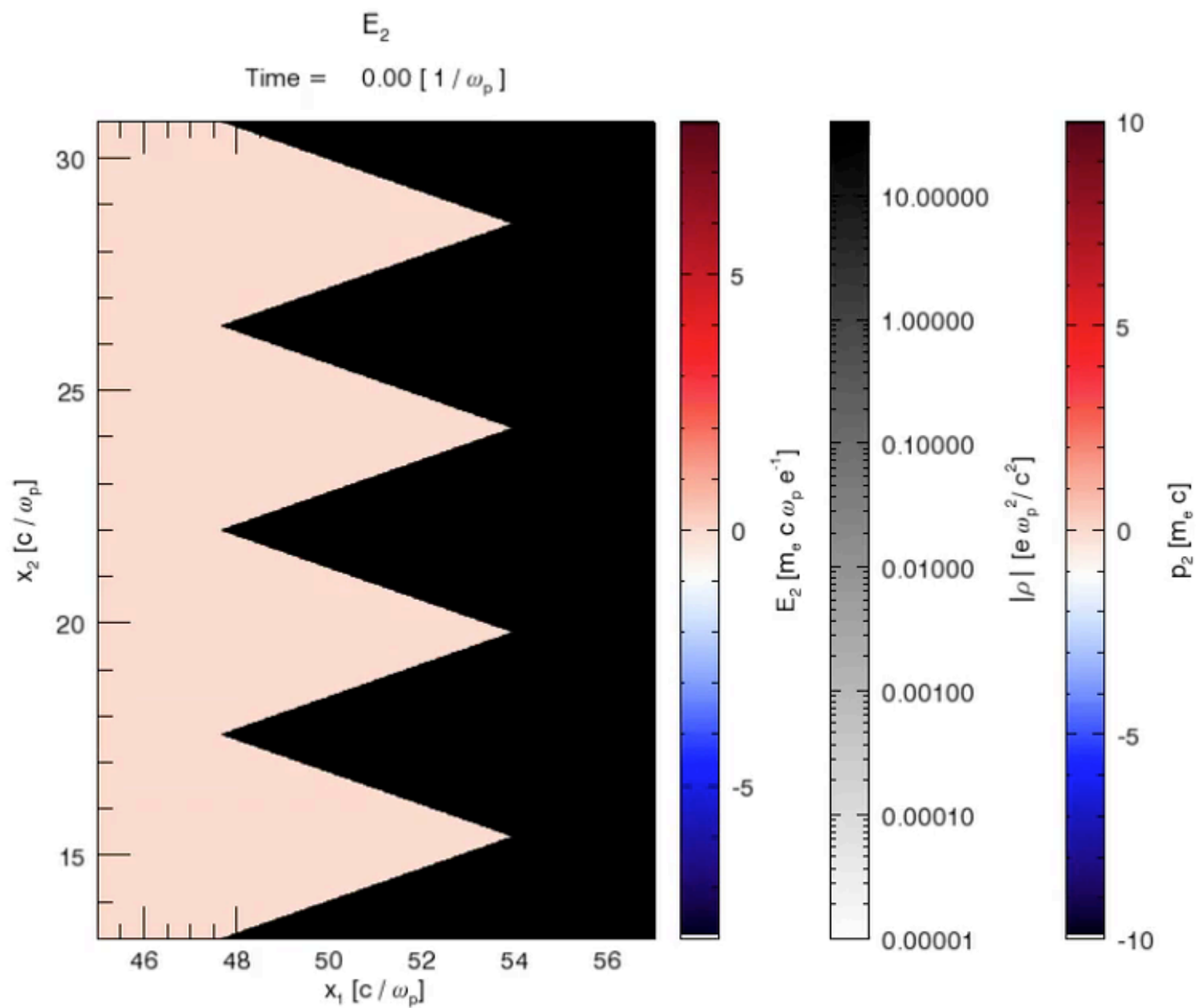
Normal triangles



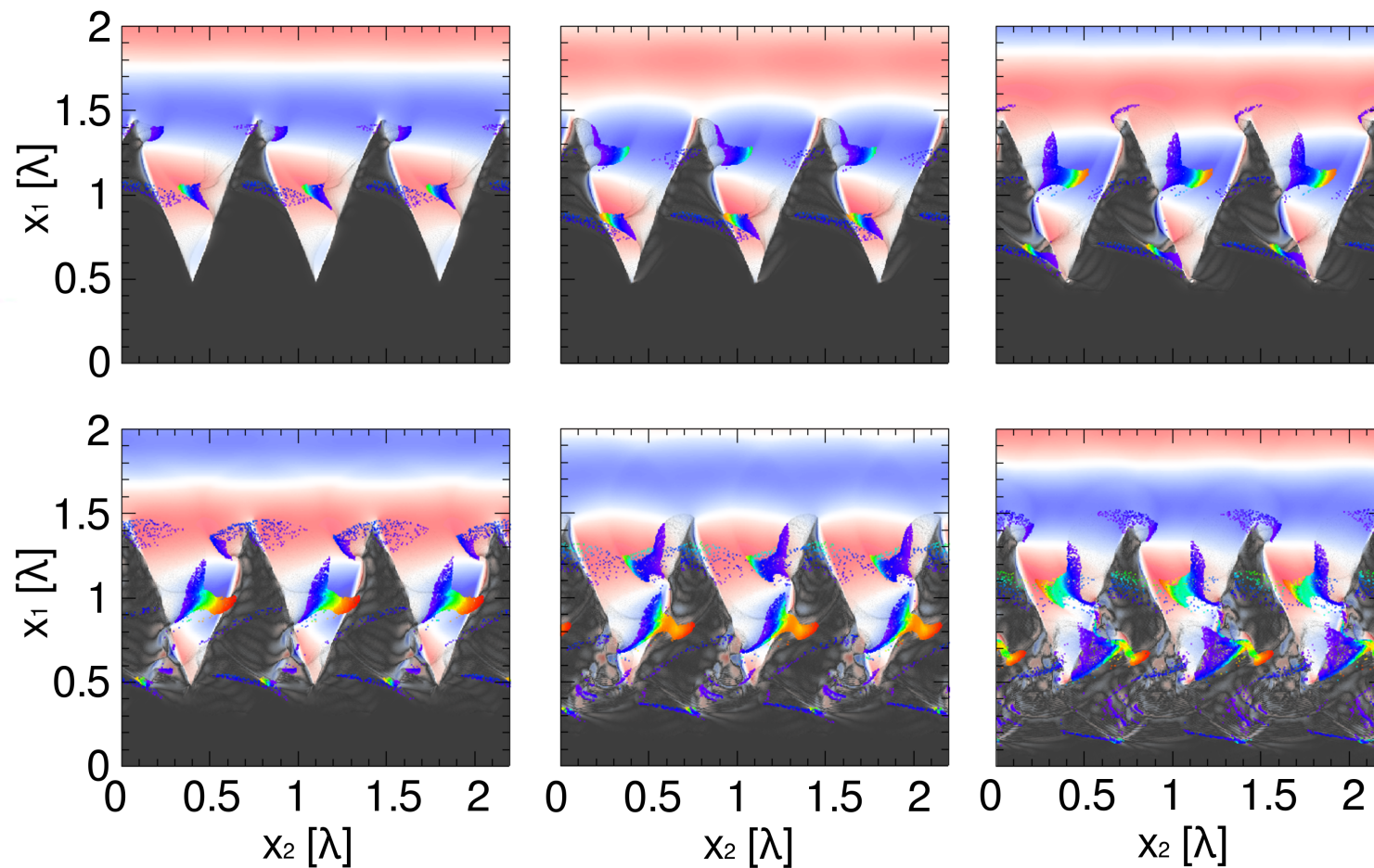
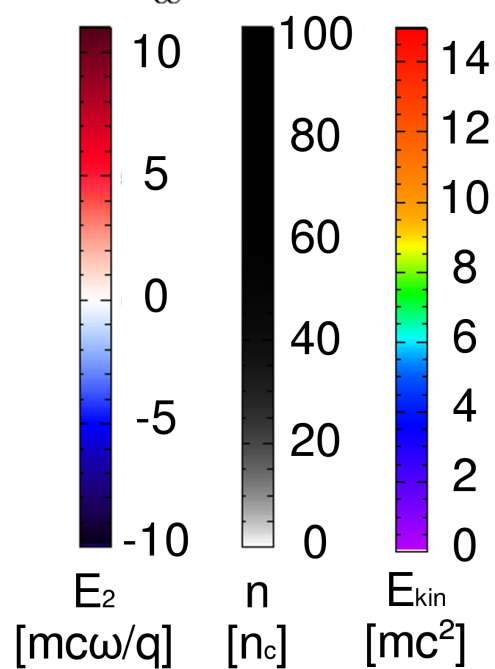
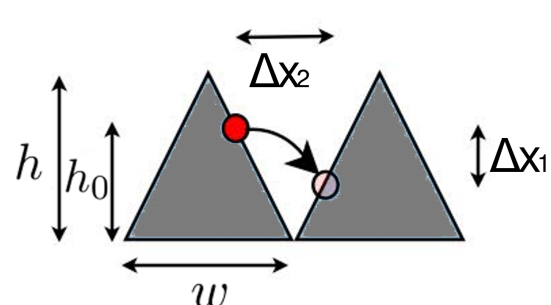
Tilted triangles



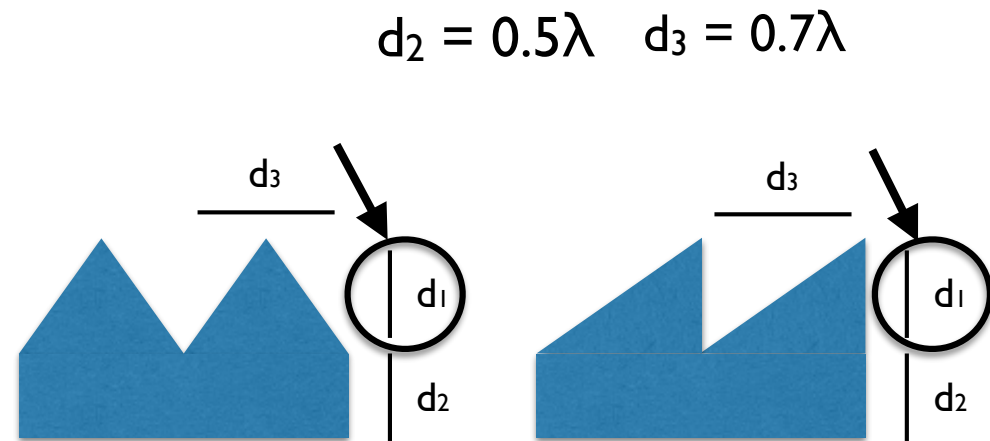
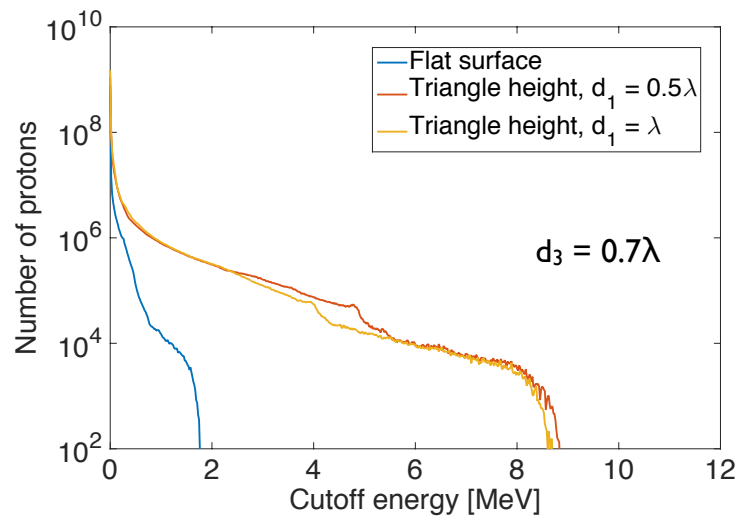




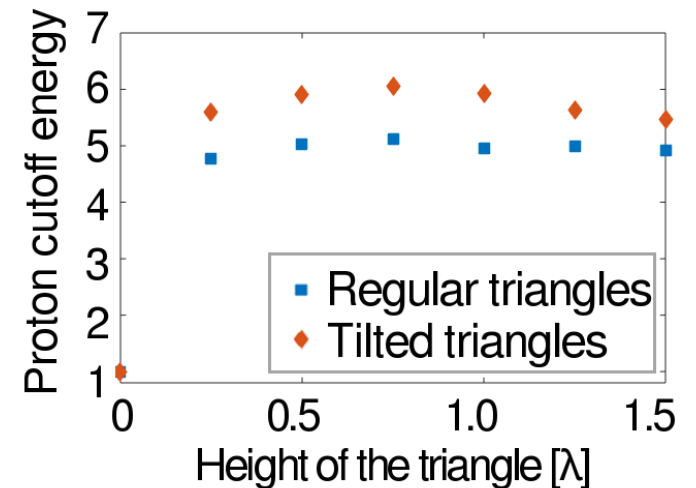
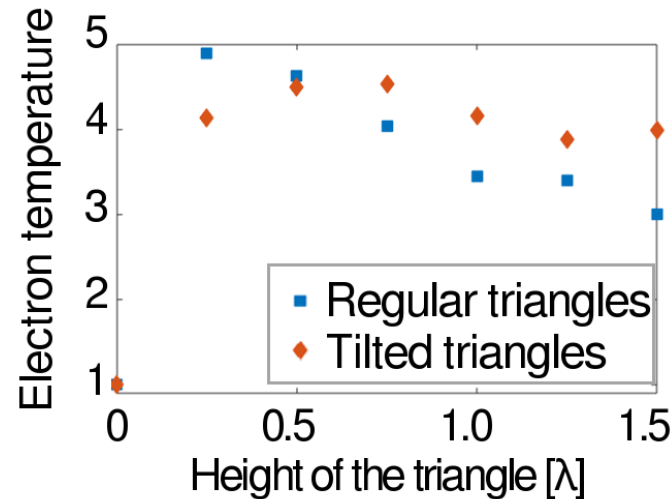
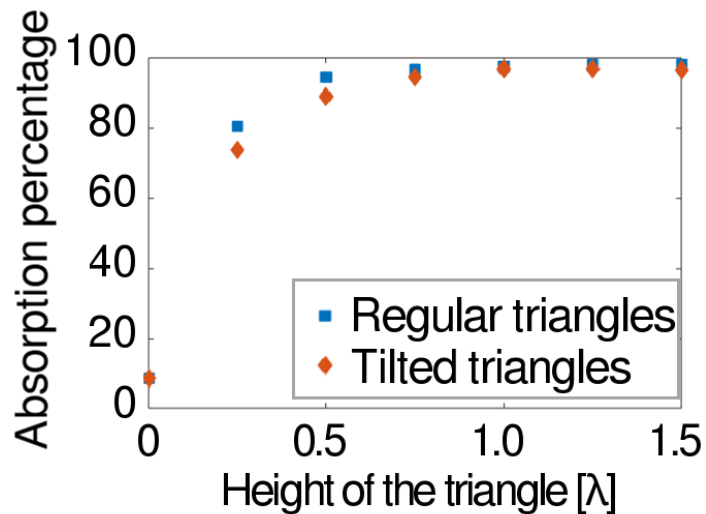
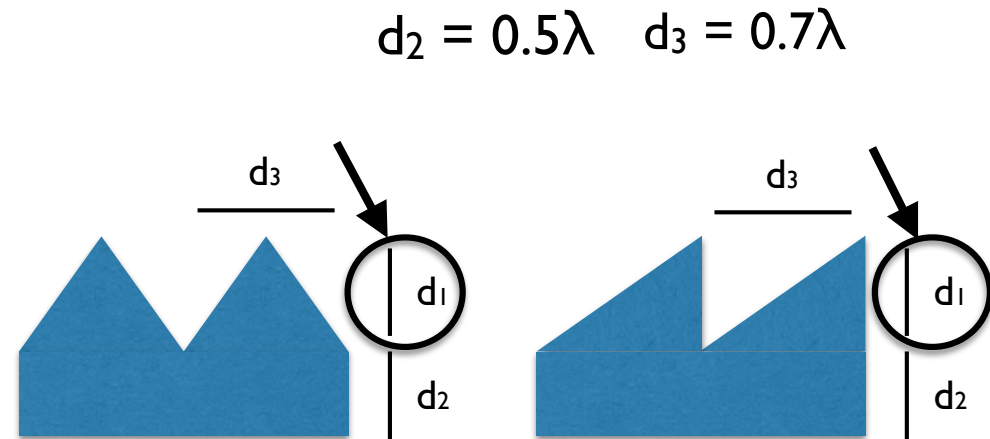
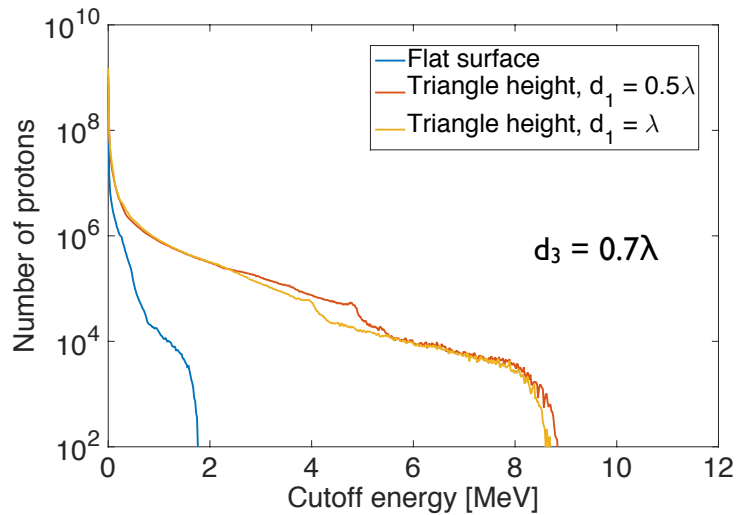
Mechanism for efficient absorption



Dependence on the height of the structure



Dependence on the height of the structure

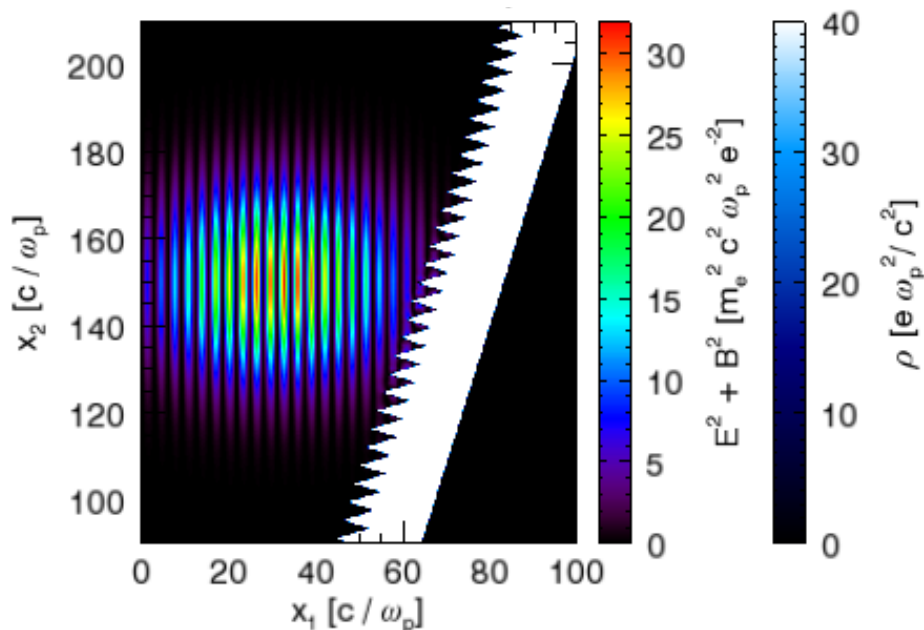


Angle of incidence

Main idea

Ion acceleration is optimized for a set of parameters, depending on the plasma density and laser peak intensity.

We want to obtain an estimate of the outcome in the most optimized situation for oblique incidence.

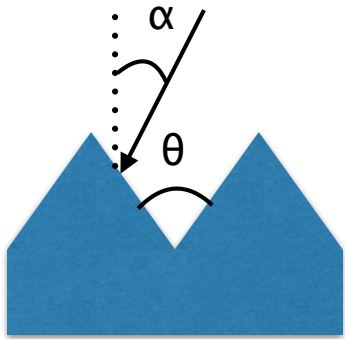


Optimal case

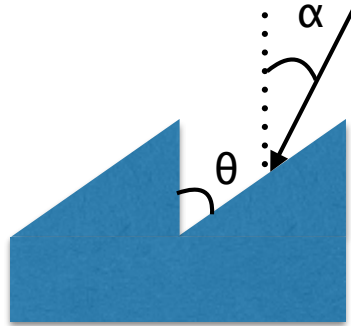
- Incidence: Oblique
- Structure: Tilted triangles
- Optimal width: $d_3 = 0.7\lambda$
- Height with saturation: $d_1 = \lambda$
- Realistic thickness: $d_2 = 2\lambda$

Dependence with the incident angle

$$d_2 = 0.5\lambda \quad d_3 = \lambda$$



$$d_1 = 1.2\lambda$$



$$d_1 = \lambda$$

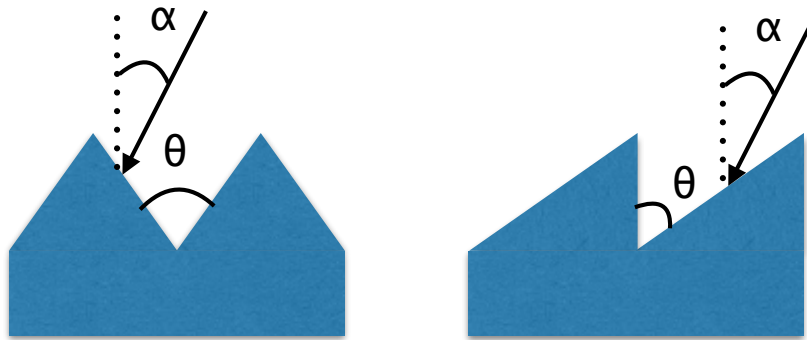
What we expect

We expect an asymmetry for the tilted triangles for negative or positive angles.

We choose the sizes of the structures to have the same angle $\theta=45^\circ$.

Dependence with the incident angle

$$d_2 = 0.5\lambda \quad d_3 = \lambda$$



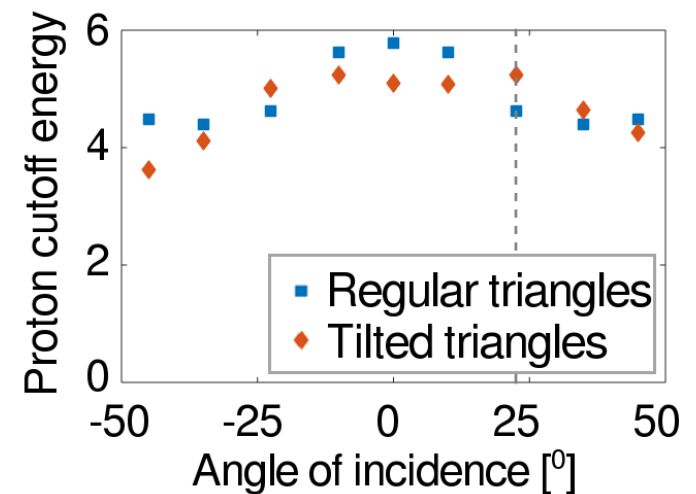
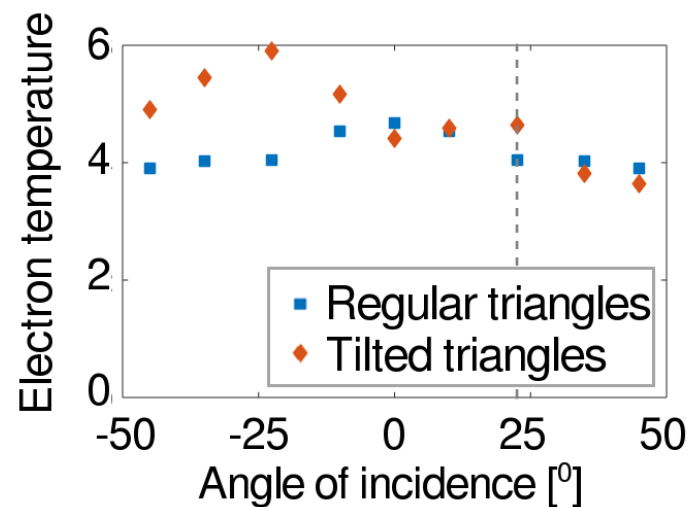
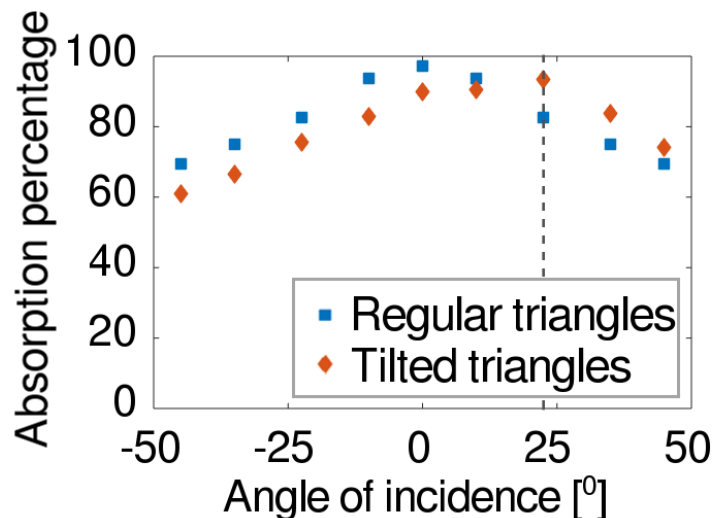
$$d_1 = 1.2\lambda$$

$$d_1 = \lambda$$

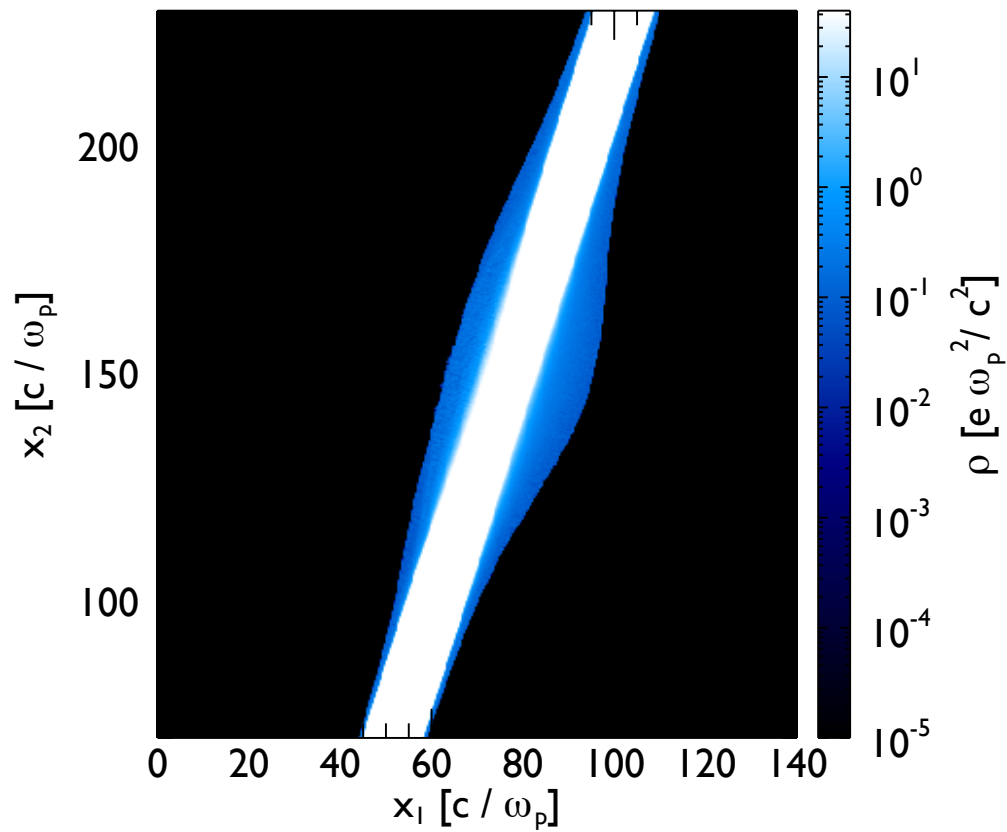
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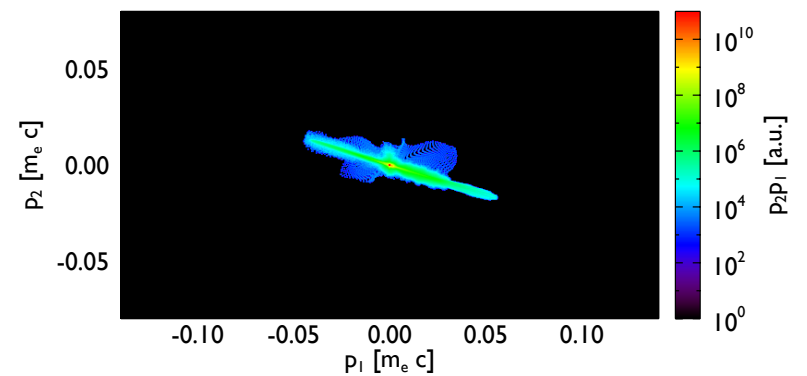
Proton charge density ($t=405/w$)



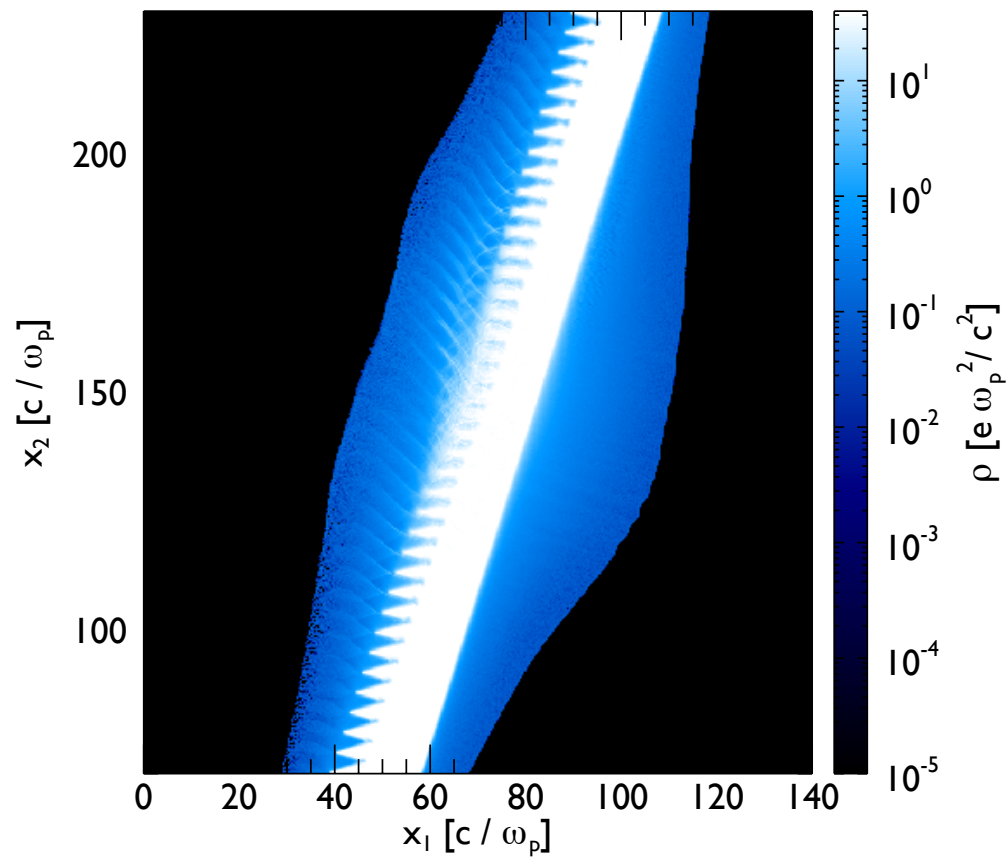
$$\alpha = 17.5^\circ \quad d_2 = 2\lambda$$

	Absorption	Proton cutoff
Flat	6,1 %	1,60 MeV
Structured	90,6 %	7,72 MeV
Ratio	14,9	4,8

Proton $p_2 p_1$ phasespace ($t=405/w$)



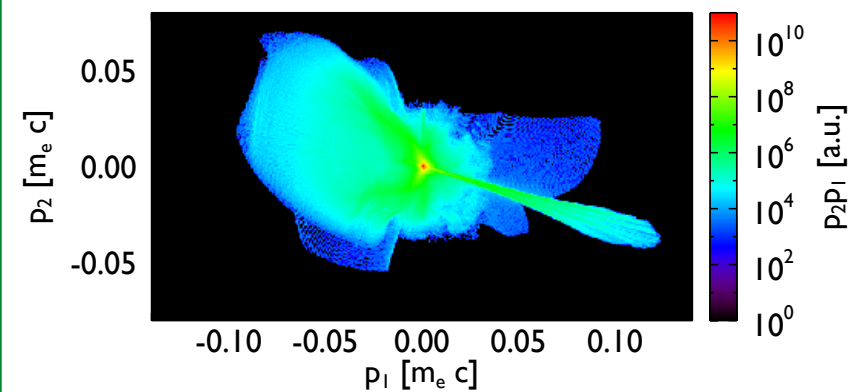
Proton charge density ($t=405/w$)



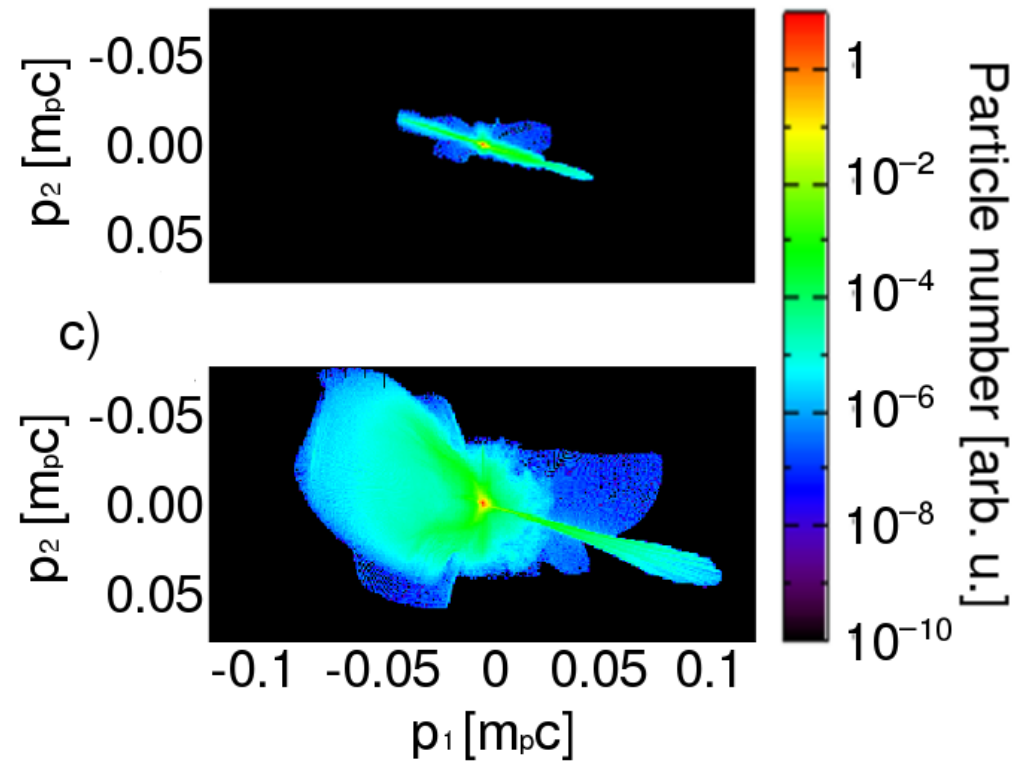
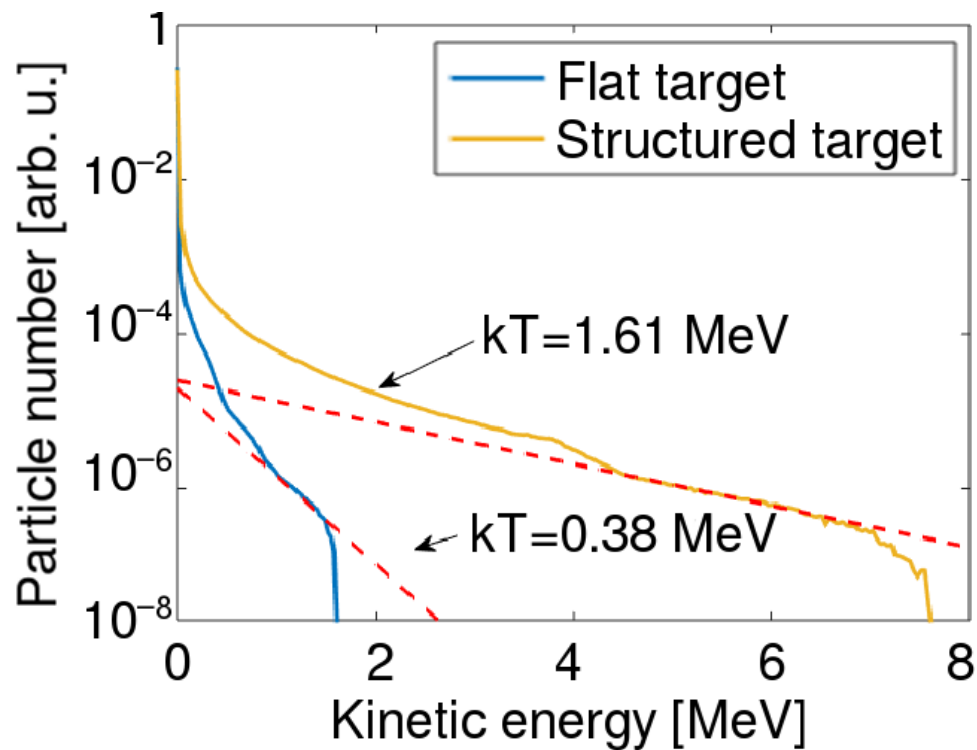
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Optimized proton beam



Why comparing

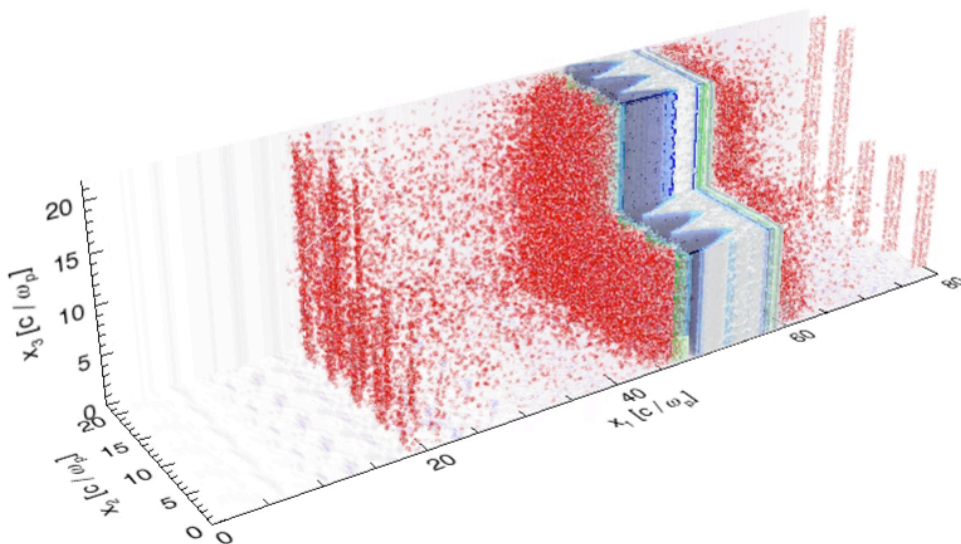
We need to know how the results scale from 2D to 3D to see how reliable they are.

We will use slab geometry to compare both cases.

Scaling of reflected energy and cutoff

	Regular triangles			Tilted triangles		
	2D	3D	3D/2D	2D	3D	3D/2D
Absorbed energy	95,4%			92,3%		
	No slab: 97,6 %	91,5 %	0,96	No slab: 97,0 %	89,5 %	0,97
Proton cutoff	4.29 MeV	4.28 MeV	0,998	4.50 MeV	4.21 MeV	0,936

Interaction



Simulation setup

Simulation:

x axis: 20.4 μm in 1600 cells

y/z axis: 2.8 μm in 220 cells

time step: 21.3 as

$d_1 = \lambda$

$d_2 = 0.5\lambda$

$d_3 = 0.7\lambda$

Laser:

$a_0 = 4.0$

$\lambda = 0.8 \mu\text{m}$

length = 18.8 fs

Plasma

$n_0 = 6.9 \times 10^{22} \text{ cm}^{-3}$

$n_0 = 40 n_c$

particles per cell: 1x2x2

Why comparing

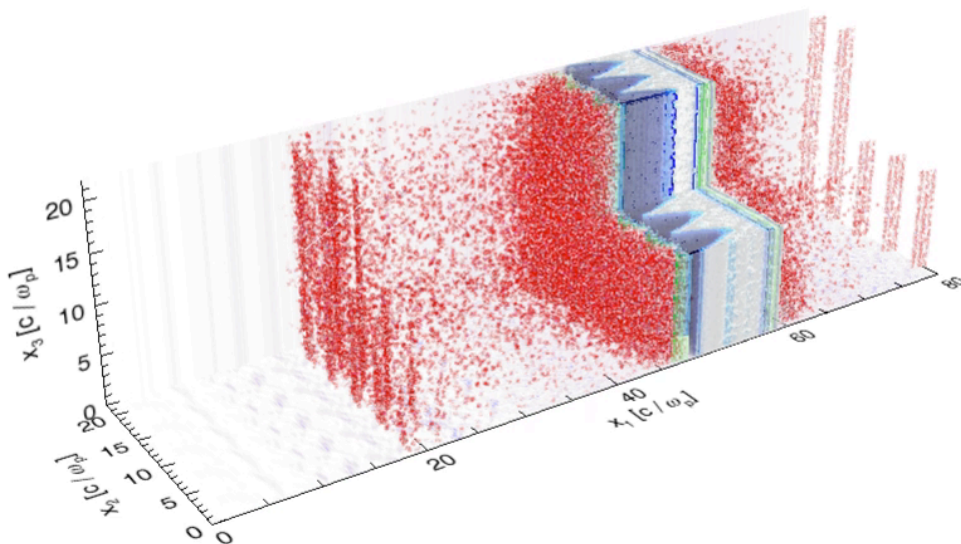
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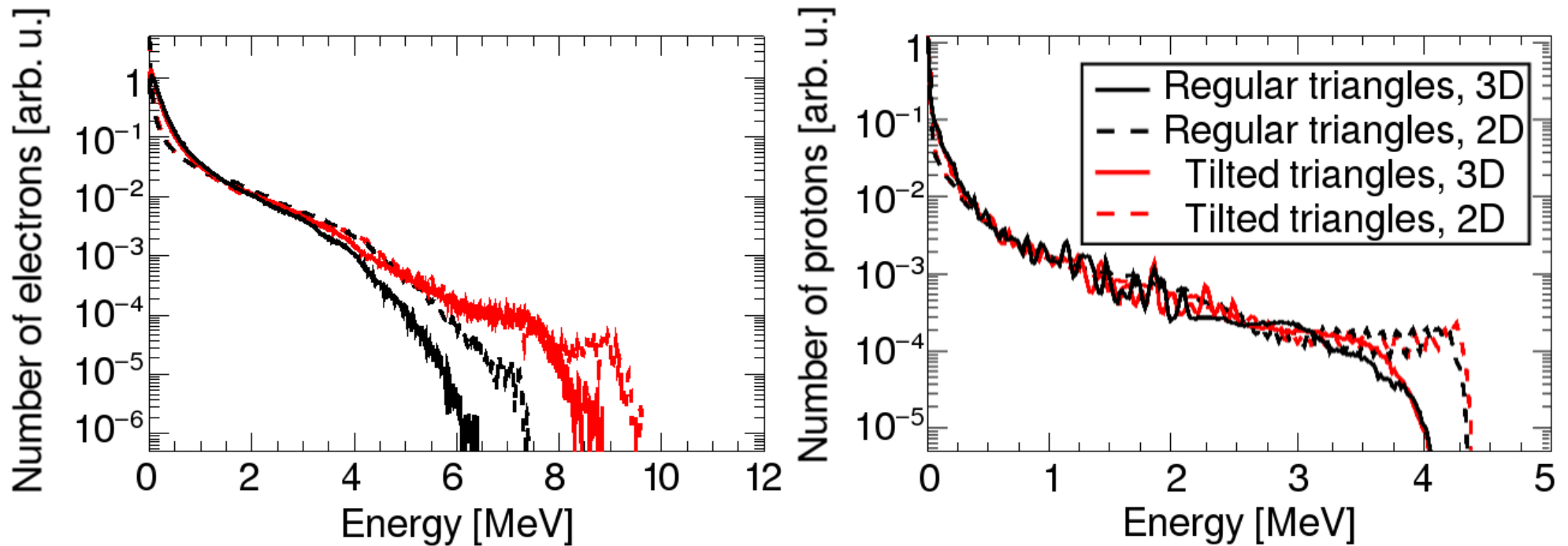
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$n_0 = 6.9 \times 10^{22} \text{ cm}^{-3}$

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particles per cell: 1x2x2

Comparison 2D vs 3D



3D: $3.1 \times 10^{11} p/shot$

New diagram for the estimation.

Overall cost & complexity

New paradigm for radioisotope production

- Installation cost
- Complexity

- Energy per shot.
- Intensity
- Average power

- Target
- Geometry
- electron energy dist
- Current

- Complex targets

- TNSA
- Proton energy dist
 - Number of ions

- Advance proton diag
- Online beam monitoring

- Activation (p,n)
- Geometry
- Target
- Activation time

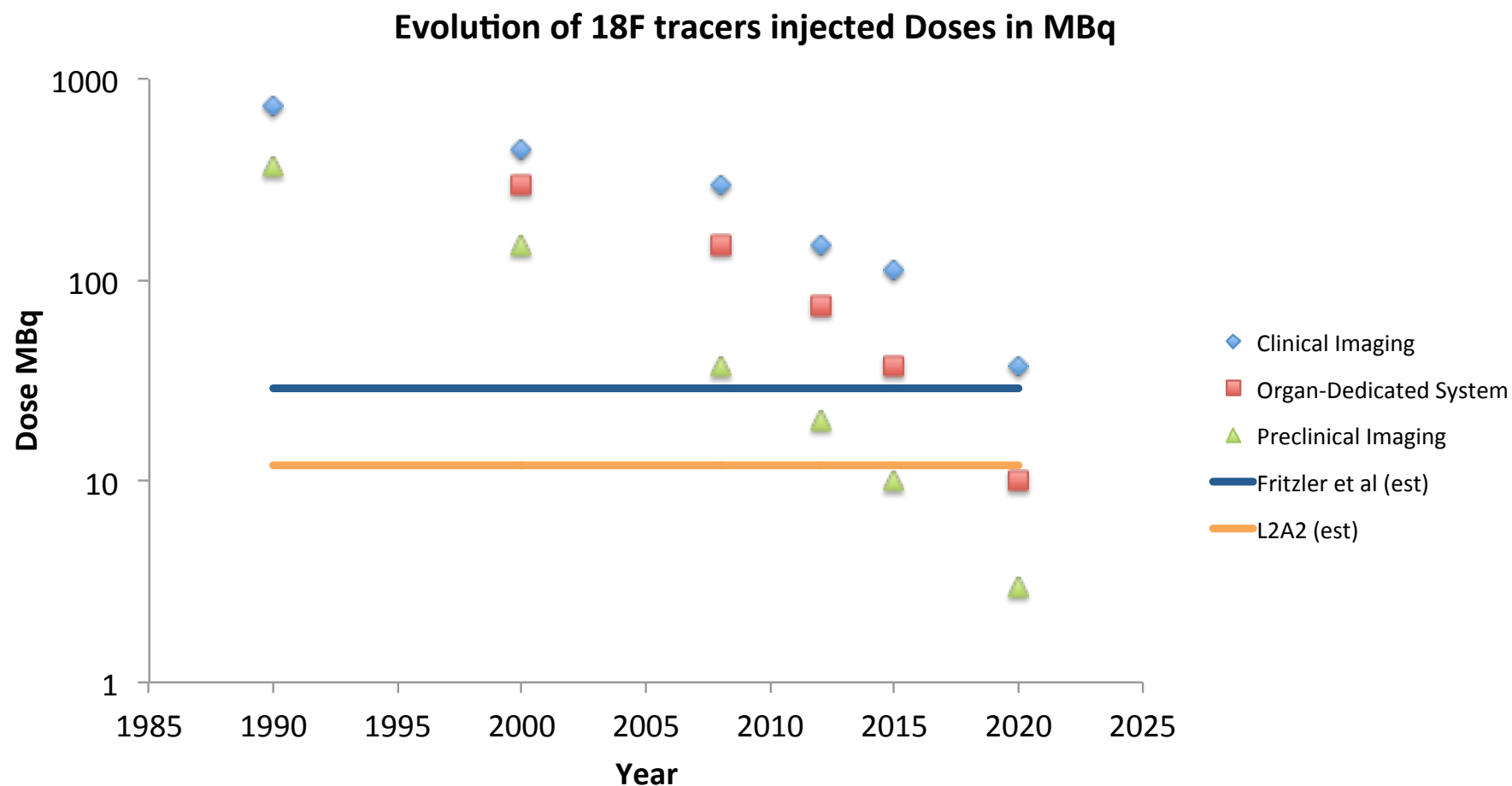
- Dose for PET

- Radiopharmacy cost

- Advance imaging technique

High repetition rate

The doses are continuously decreasing



New imaging techniques

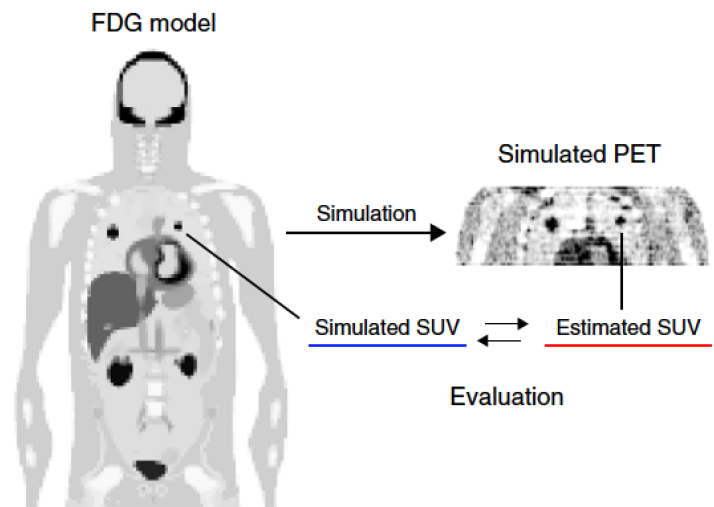


Fig. 1. Workflow was based on a realistic framework of multiple simulated FDG-PET studies.

Silva-Rodríguez J, et al. Simulated FDG-PET studies for the assessment of SUV quantification methods. Rev Esp Med Nucl Imagen Mol. 2014. <http://dx.doi.org/10.1016/j.remnm.2014.07.006>

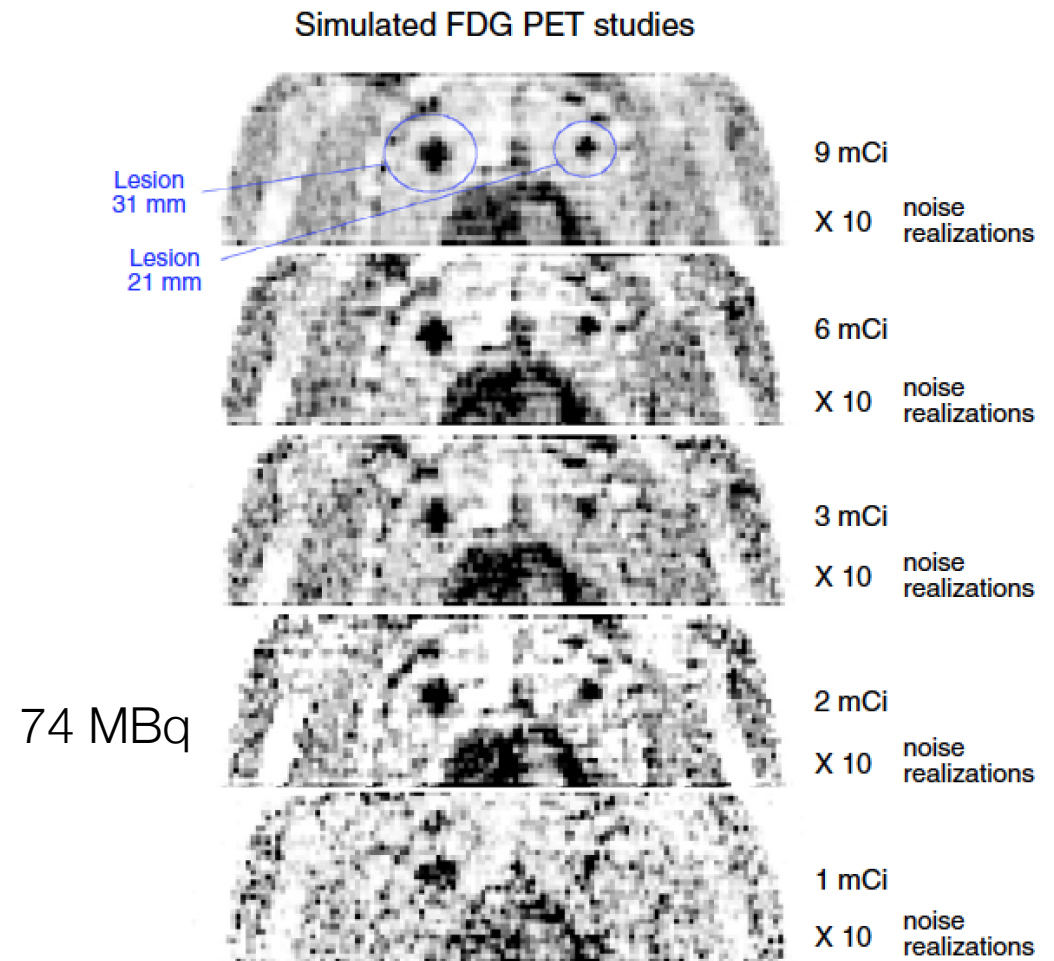


Fig. 2. Coronal views of simulated FDG-PET studies of patients with solitary pulmonary nodules for different injected FDG doses ranging from 9 mCi to 1 mCi.

Acknowledge.

Laser physics

- J. Arines, C. Bao, M. Flores (USC)
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- A. Paredes (UVi)



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D SALAMANCA

CAMPUS DE EXCELENCIA INTERNACIONAL



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Manuel Blanco

Marija Vranic



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