

Axions, supersymmetry, and all that: Emergent phenomena in condensed matter physics

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Lake Louise Winter Institute
February 22, 2018



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What is more important, particle physics or condensed matter?

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“Die Festkörperphysik ist eine Schmutzphysik”

(Solid state physics is dirt physics)

“One shouldn’t wallow in dirt” (letter to Peierls)

“squalid-state physics”



Why study “dirt physics”?



“The behavior of large and complex aggregates of elementary particles, it turns out, is not to be understood in terms of a simple extrapolation of the properties of a few particles”

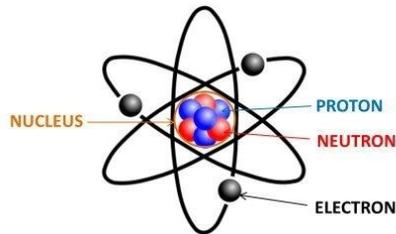
(P. W. Anderson, “More is Different”, Science ‘72)

Collective phenomena

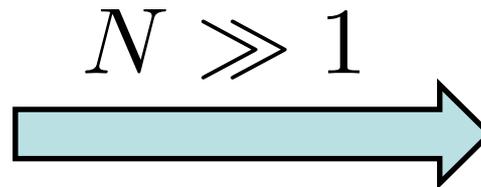


“The behavior of large and complex aggregates of elementary particles, it turns out, is not to be understood in terms of a simple extrapolation of the properties of a few particles”

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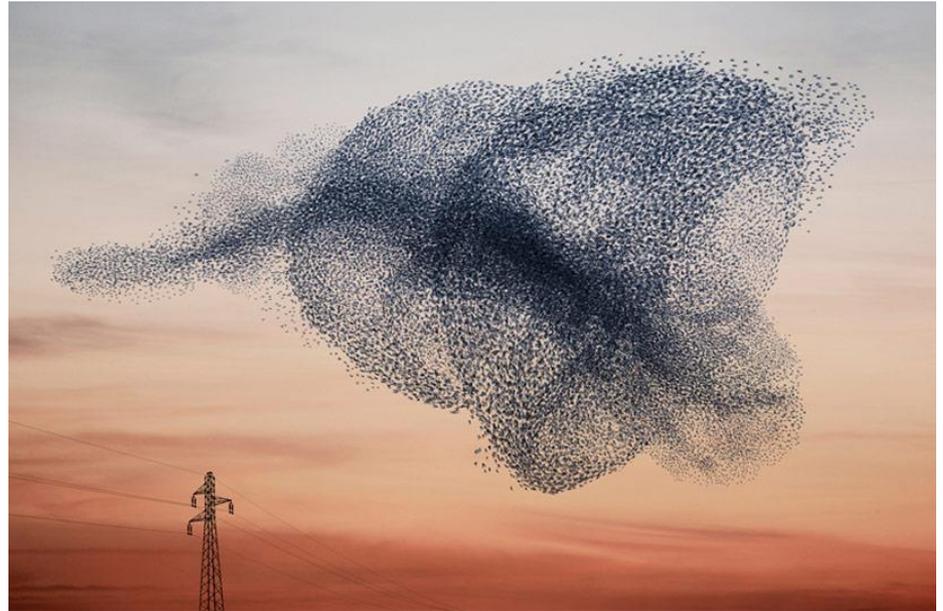
charge, spin, mass



rigidity, resistance,
specific heat...

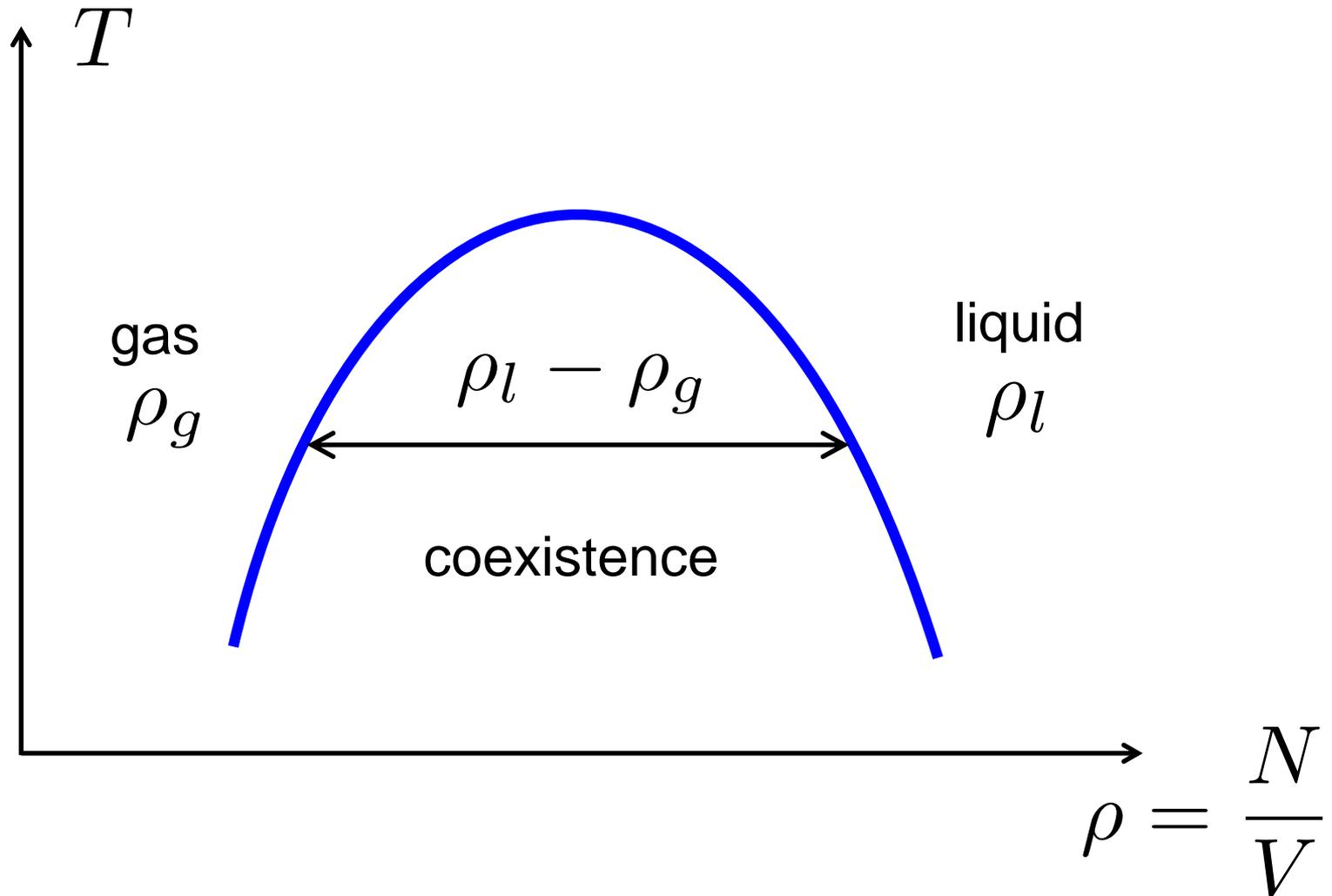
Emergence and universality

- Collective behavior can be qualitatively distinct from microscopic behavior (**emergence**)

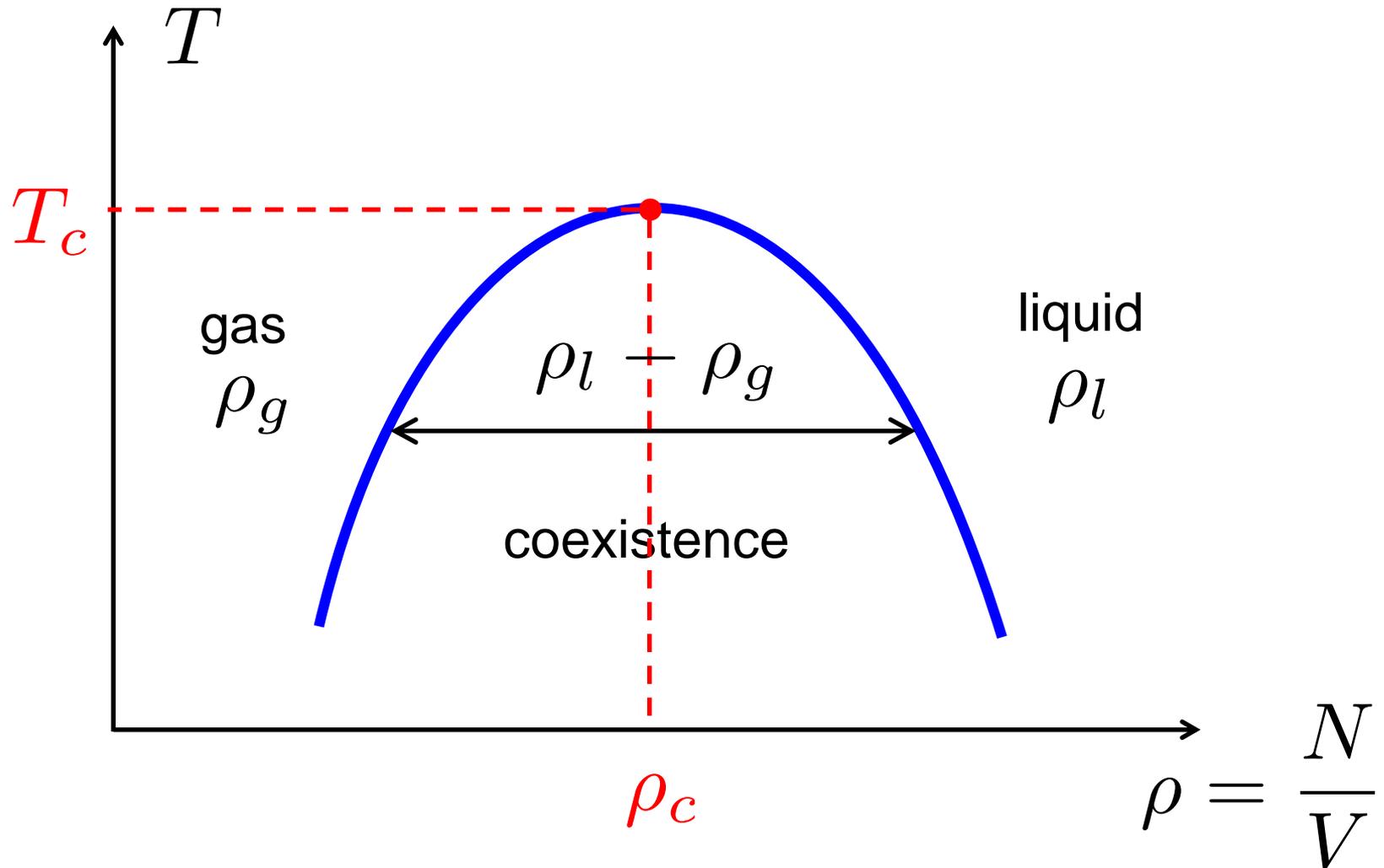


- The laws of physics that describe collective phenomena on long length/time scales are **universal**

Liquid-gas critical point

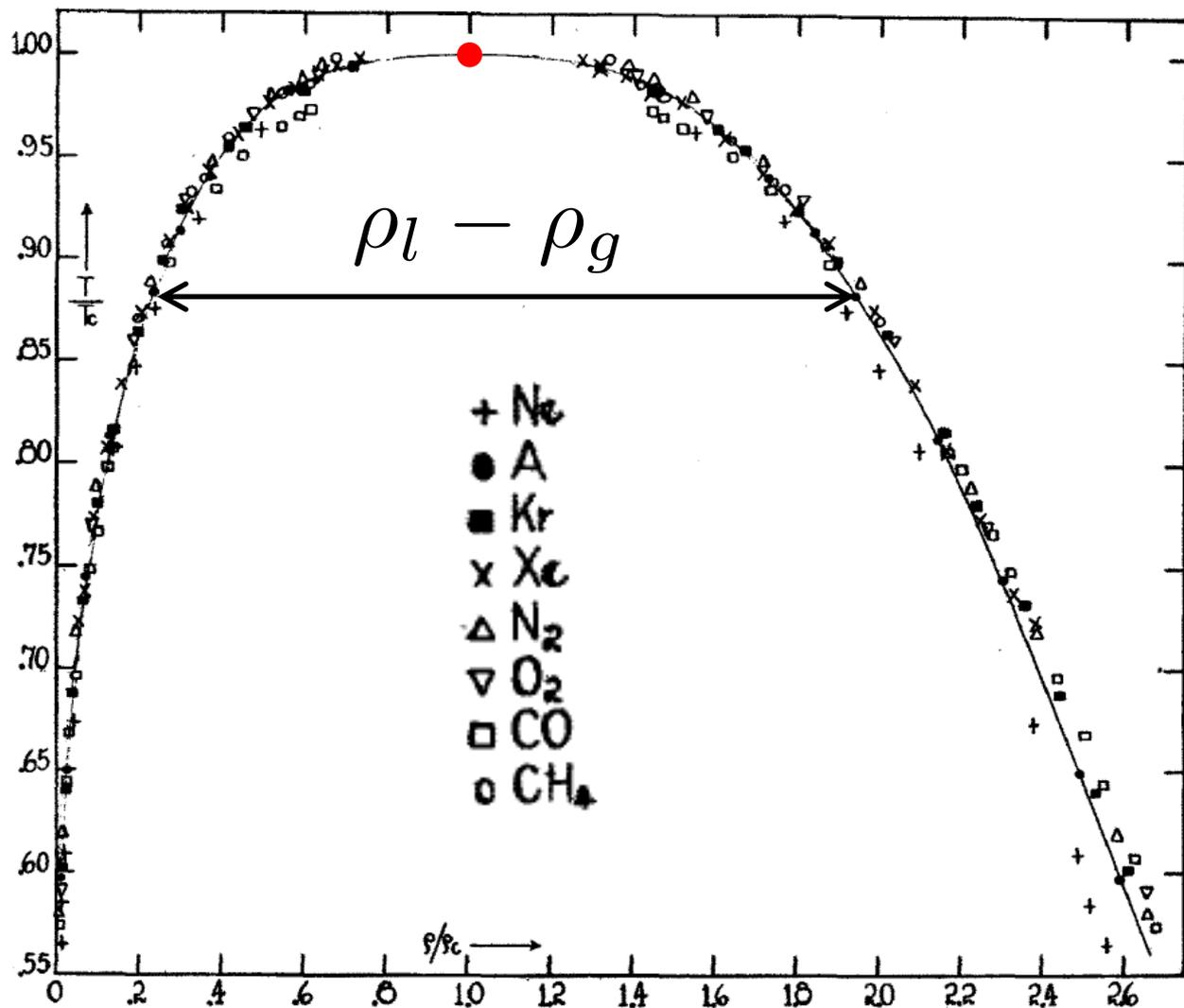


Liquid-gas critical point



$$T/T_c$$

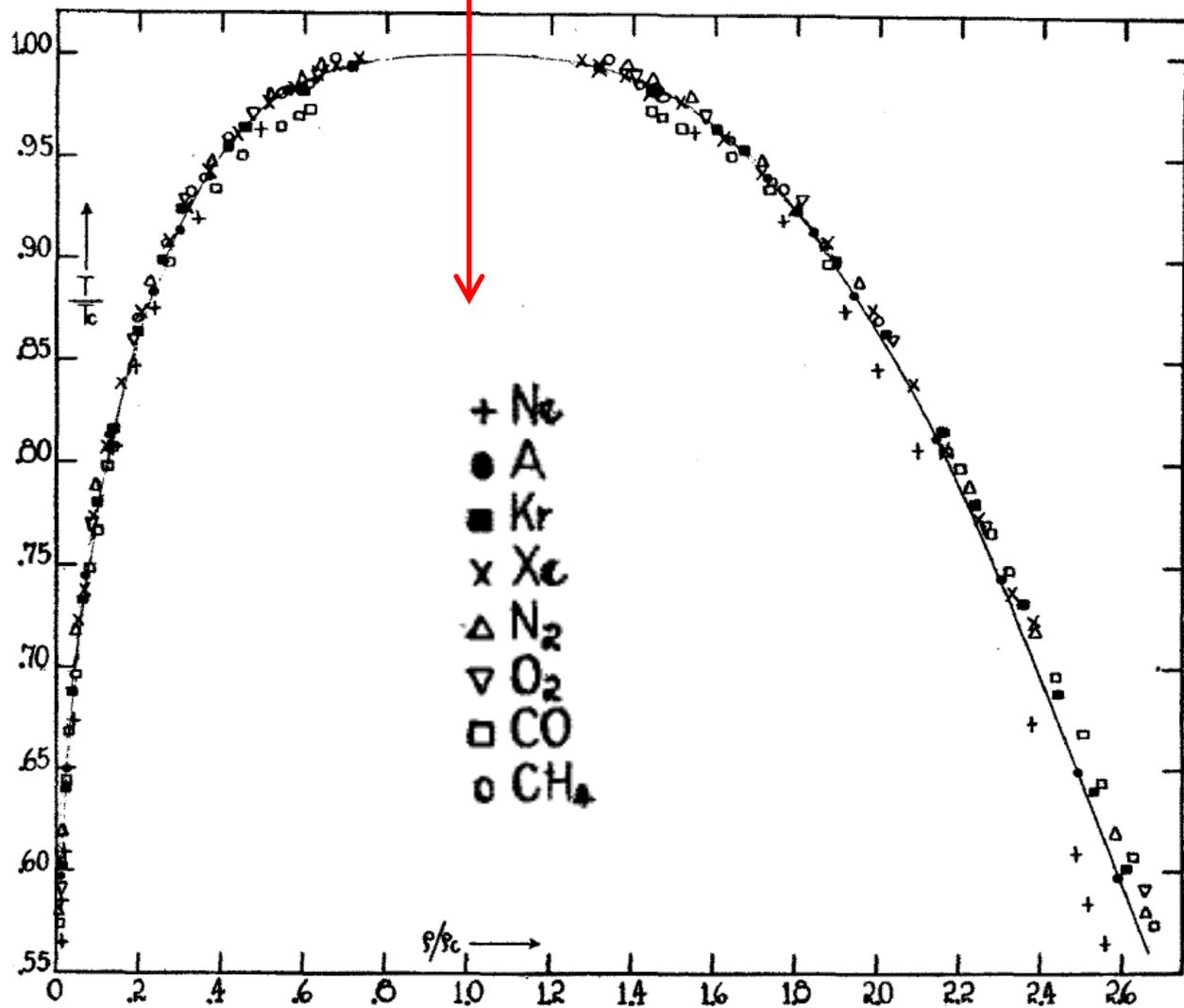
$$|\rho_l - \rho_g| \propto |T - T_c|^\beta$$

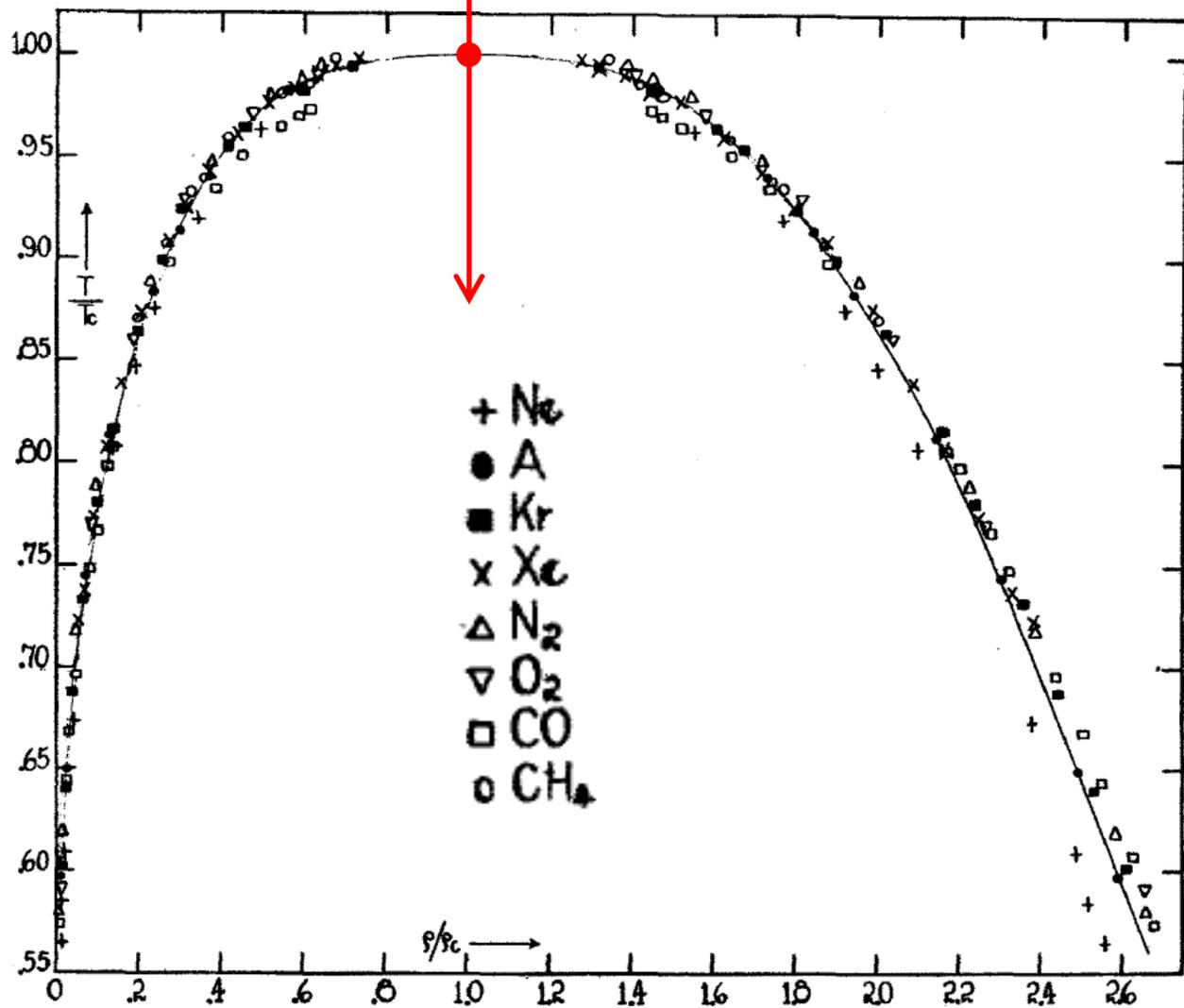


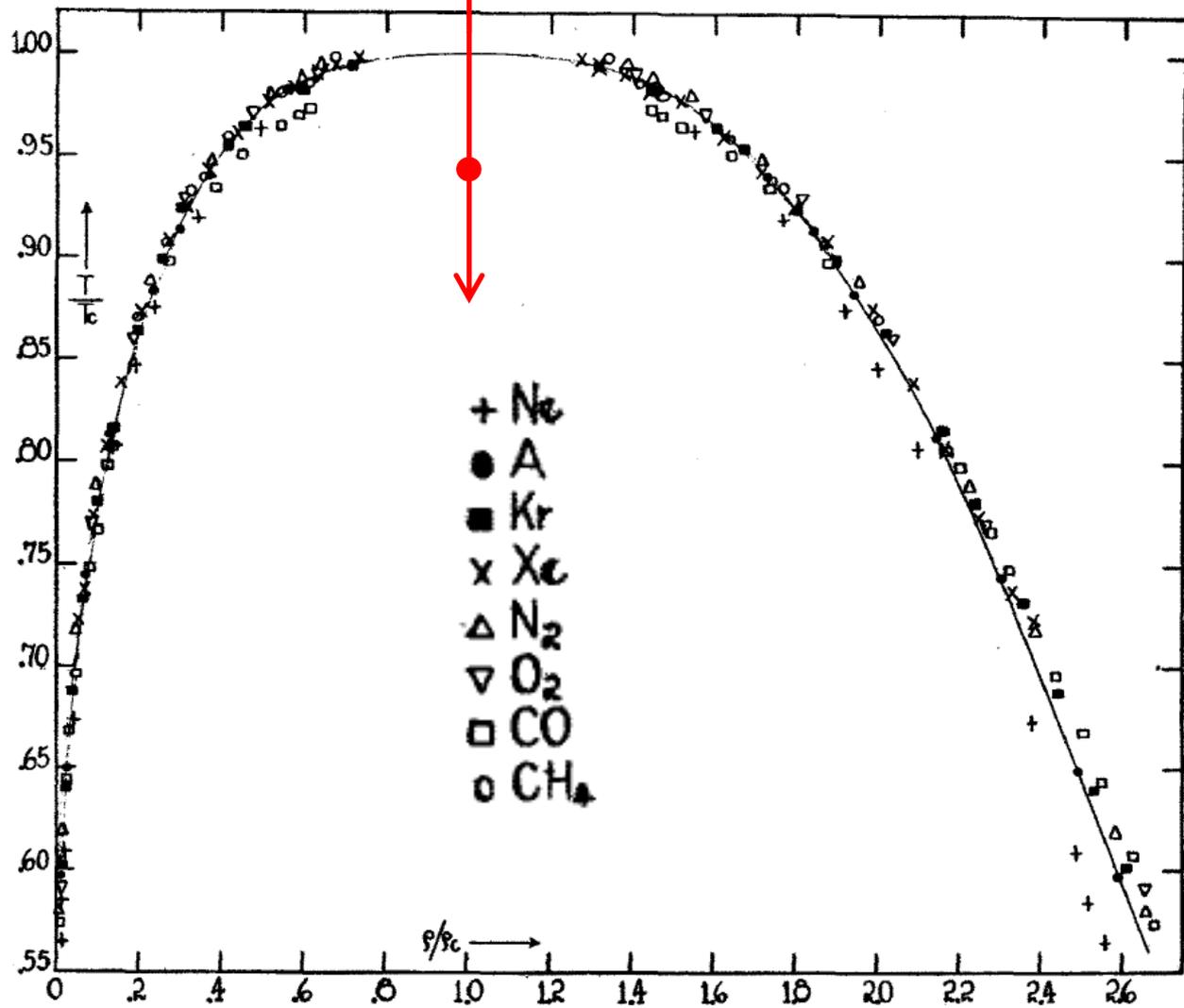
$$\beta \approx 0.32$$

for all
substances!

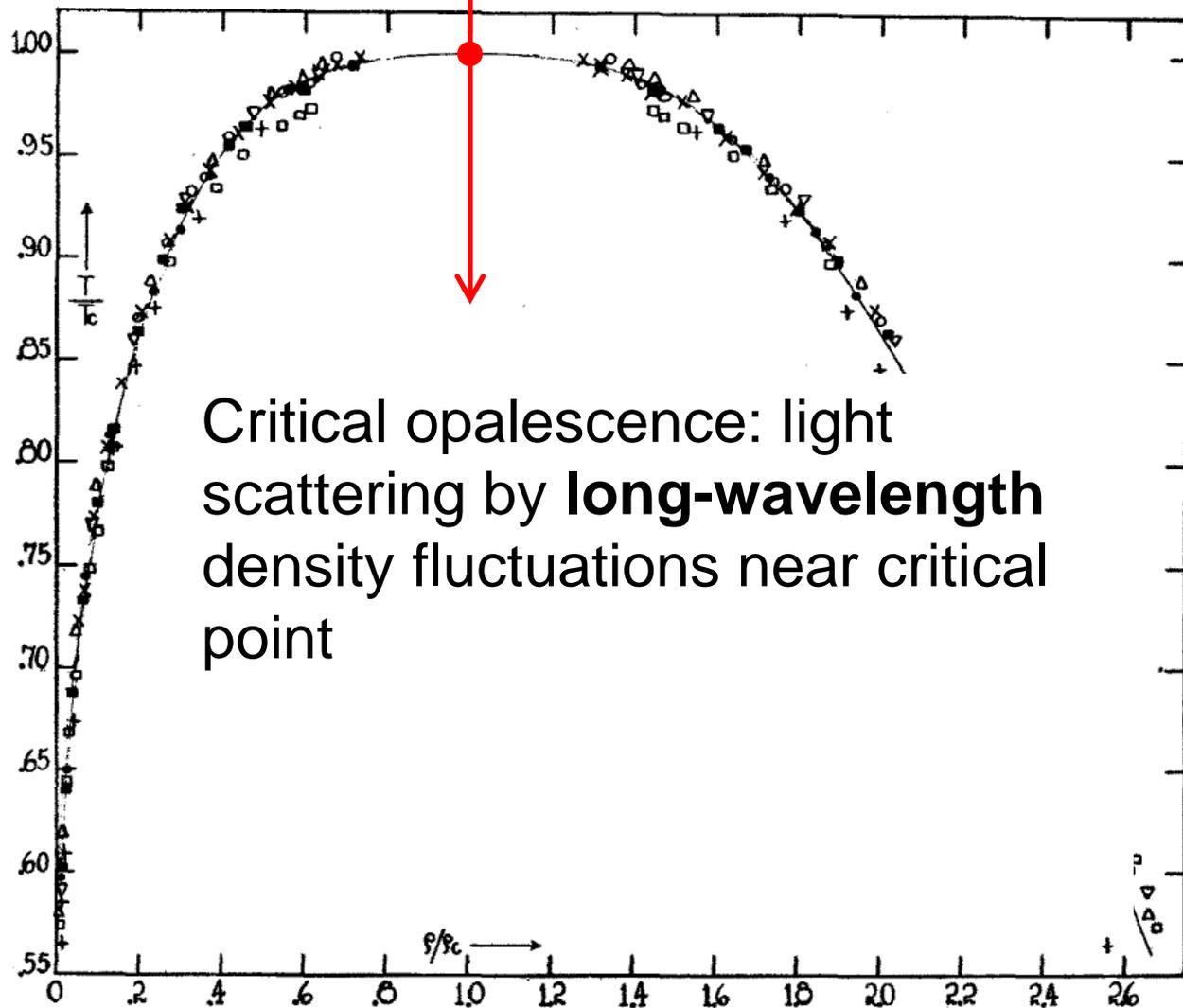
$$\rho/\rho_c$$

T/T_c  ρ/ρ_c

T/T_c  ρ/ρ_c

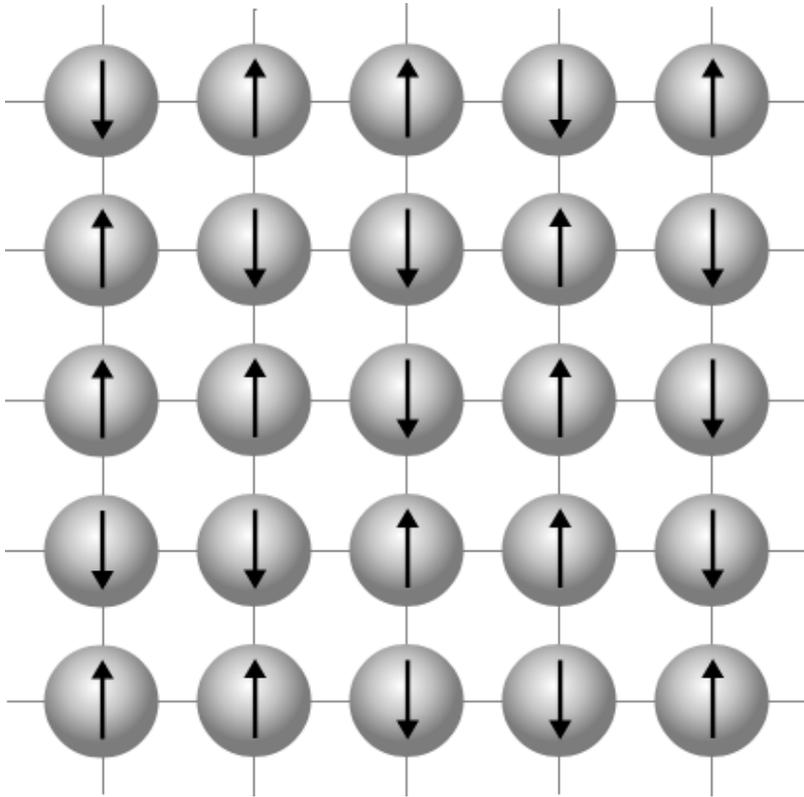
T/T_c  ρ/ρ_c

$$T/T_c$$



$$\rho/\rho_c$$

Ising model



$$H = -J \sum_{\langle ij \rangle} s_i s_j,$$

$$s_i = \pm 1$$

T

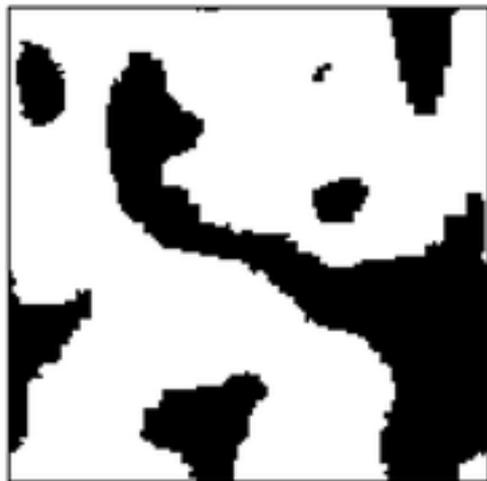


paramagnet

● $T_c \approx 2.3J/k_B$

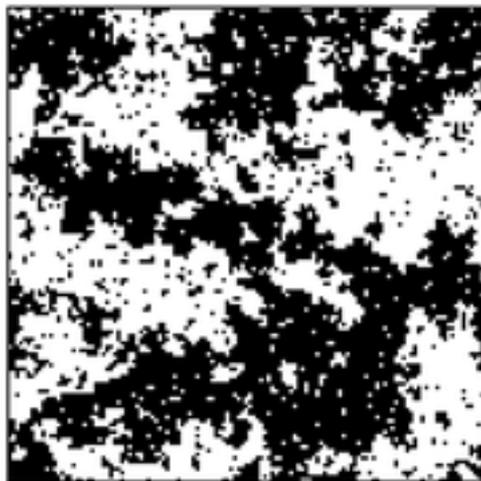
ferromagnet

$$T \ll T_c$$



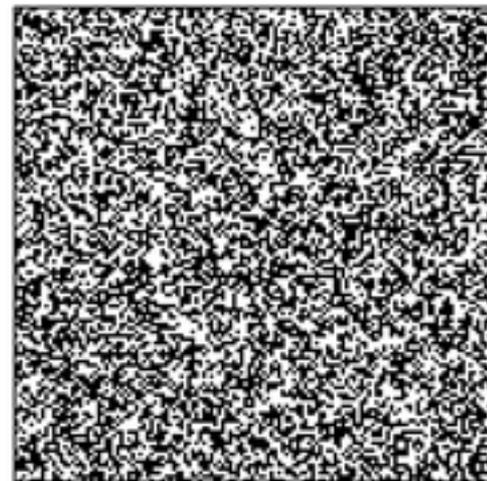
$$\xi \sim a$$

$$T \approx T_c$$

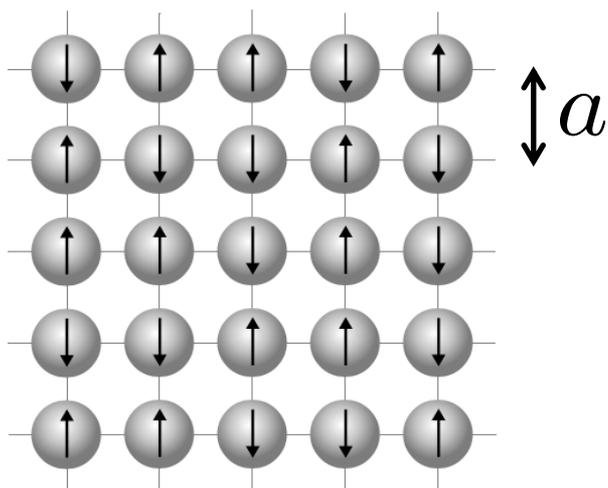


$$\xi \gg a$$

$$T \gg T_c$$



$$\xi \sim a$$



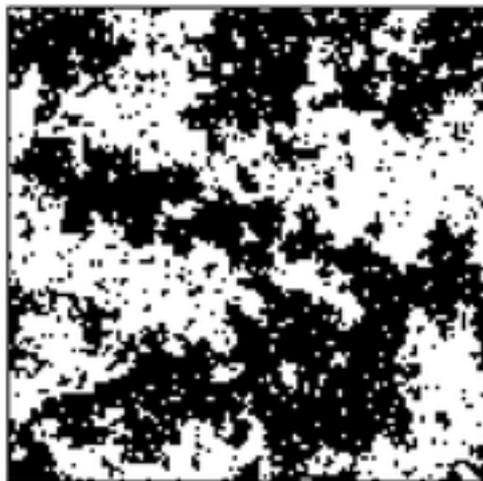
- Correlation length ξ for fluctuations **diverges** at criticality

$$T \ll T_c$$



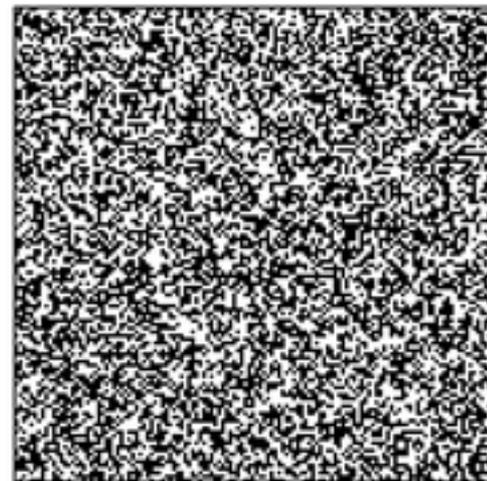
$$\xi \sim a$$

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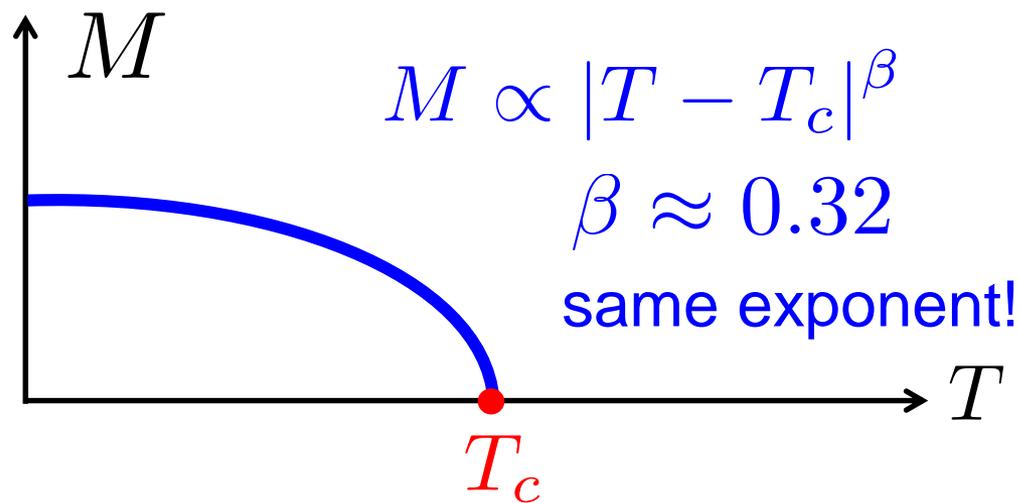
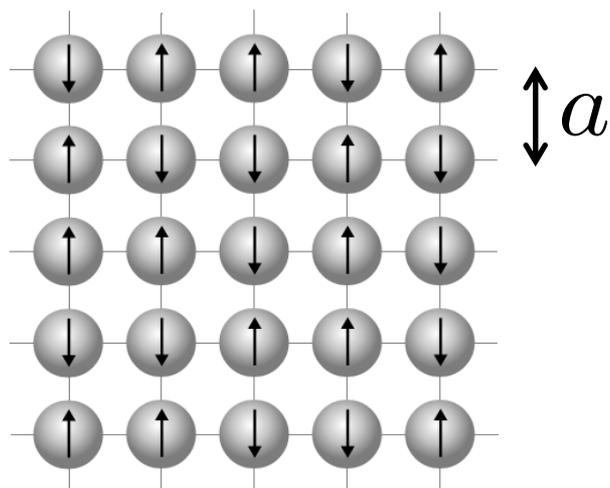


$$\xi \gg a$$

$$T \gg T_c$$

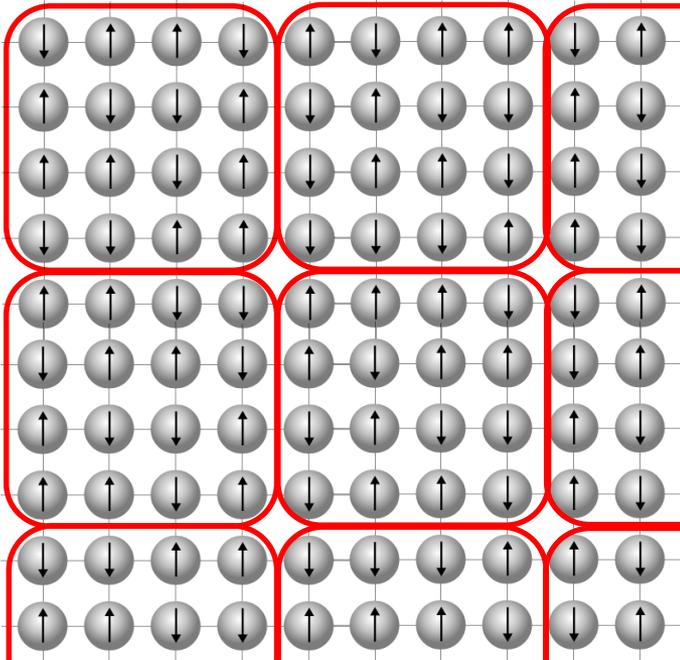


$$\xi \sim a$$



From spins to fields

$\phi(\mathbf{x})$

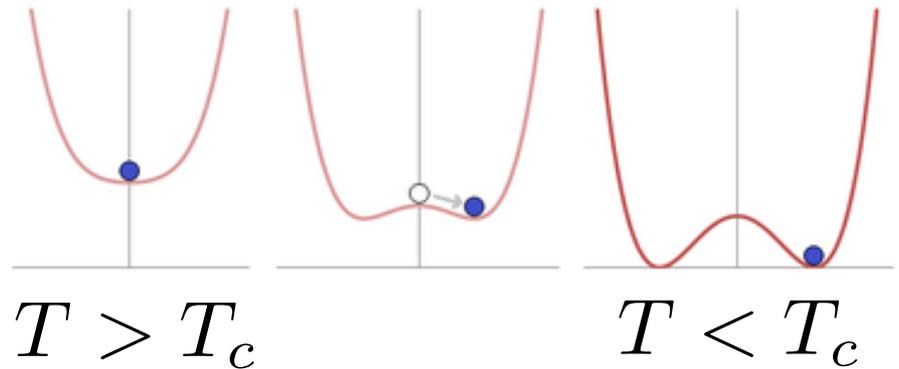


- Critical properties governed by long-wavelength (IR) fluctuations, for which microscopic (UV) details are unimportant: **universality**
- Coarse-grained description is sufficient: statistical **field** theory

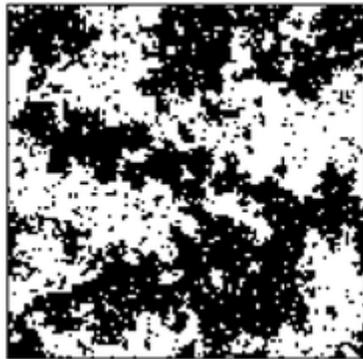
$$Z = \sum_{\{s_i\}} e^{-\beta H} \approx \int \mathcal{D}\phi(\mathbf{x}) e^{-S}$$

$$S = \int d\mathbf{x} \left(\frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{4!} \phi^4 \right)$$

$$m^2 \propto T - T_c$$



- At long wavelengths, emergent **rotational** invariance
- At critical point, also emergent **scale** invariance: universal critical exponents (*Wilson, Fisher, Polyakov...*)



$$S = \int dx \left(\frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{4!} \phi^4 \right)$$

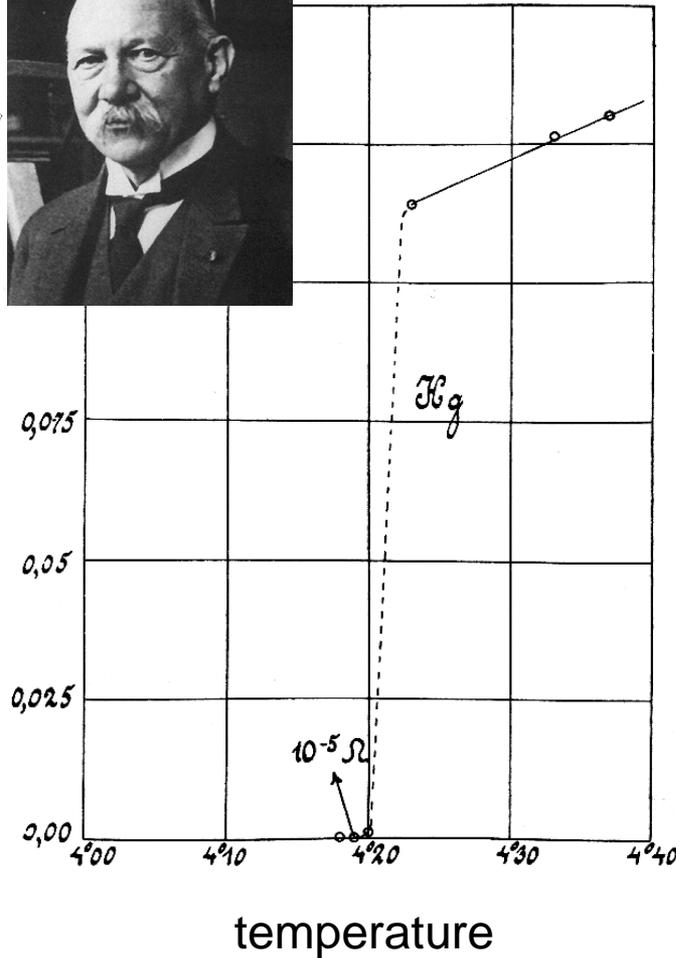
- Looks like “relativistic” QFT for scalar particle... but **static!**
- Describes spatial correlations, no time direction
- No time = no particles
- What about **quantum** statistical mechanics?



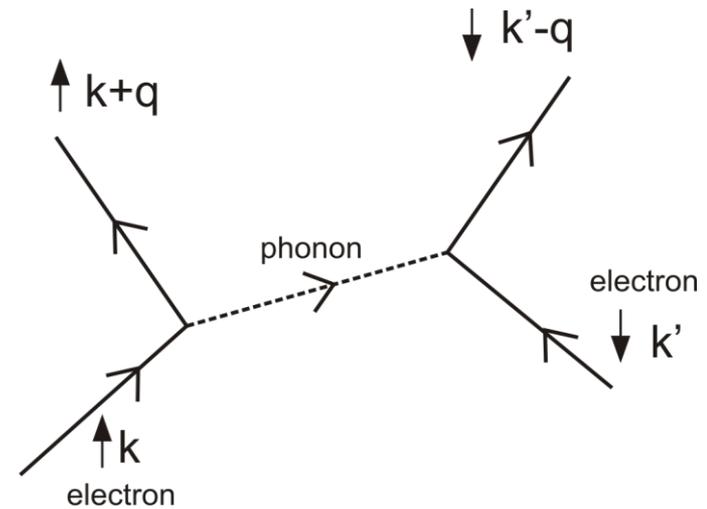
Superconductivity



electrical resistance



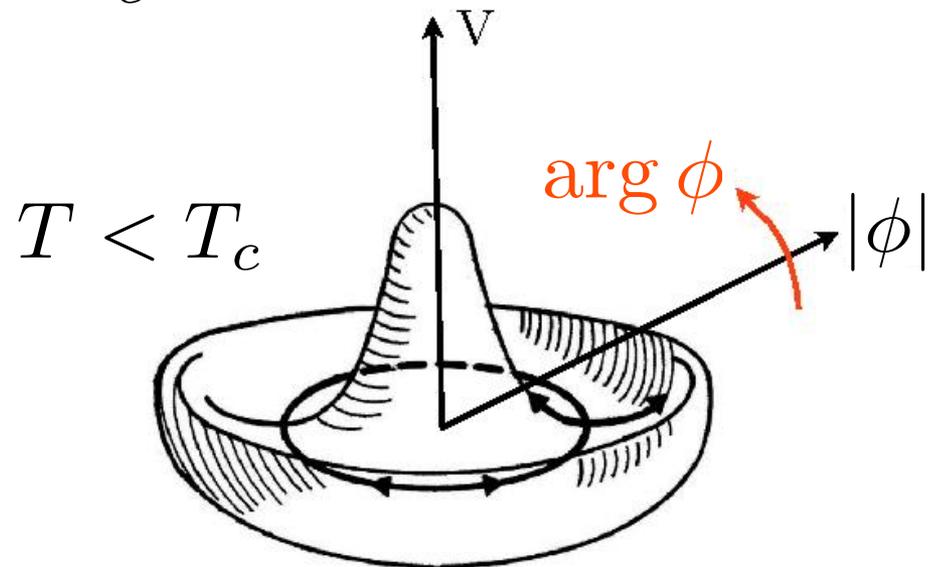
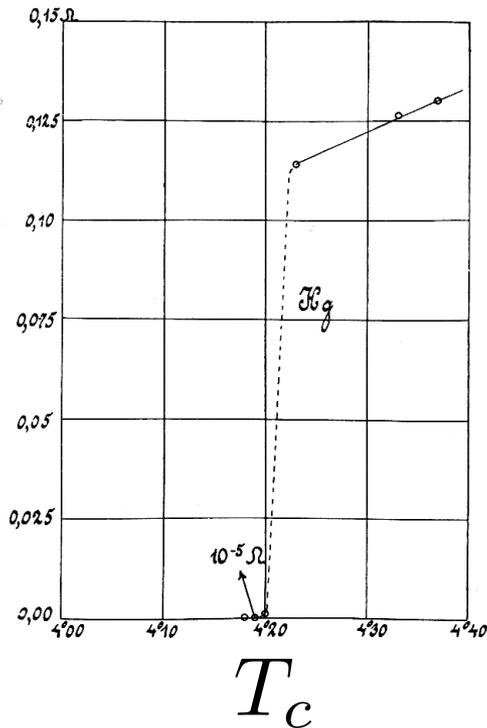
- Quantum phenomenon: Cooper pairing mediated by phonon exchange (BCS '57)



- Near superconducting transition, coarse-grained Cooper pair field = complex scalar ϕ
- Long-wavelength description is again a **static** field theory!
(Ginzburg & Landau, '50)

$$S = \int d\mathbf{x} \left(\frac{1}{2m} |(-i\hbar\nabla - 2e\mathbf{A})\phi|^2 + \alpha|\phi|^2 + \beta|\phi|^4 \right)$$

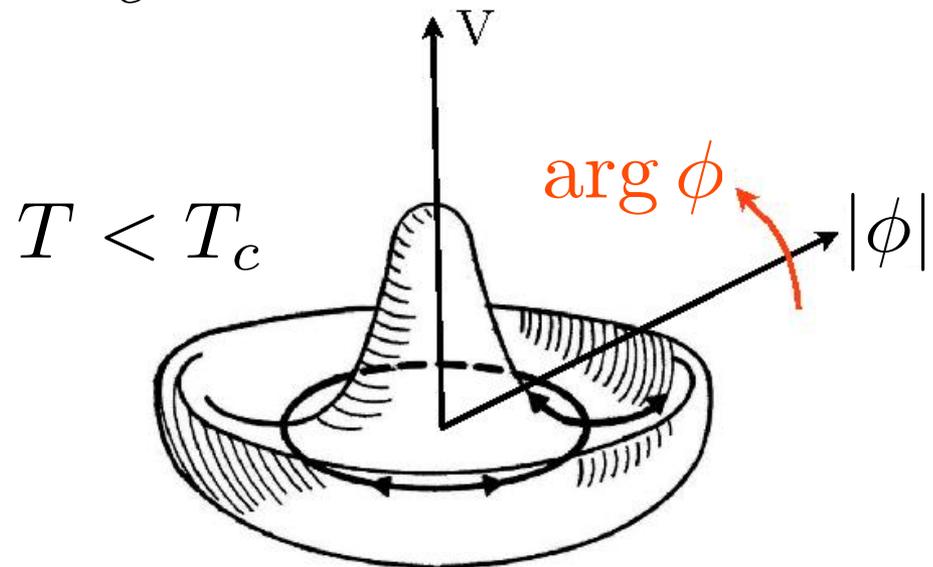
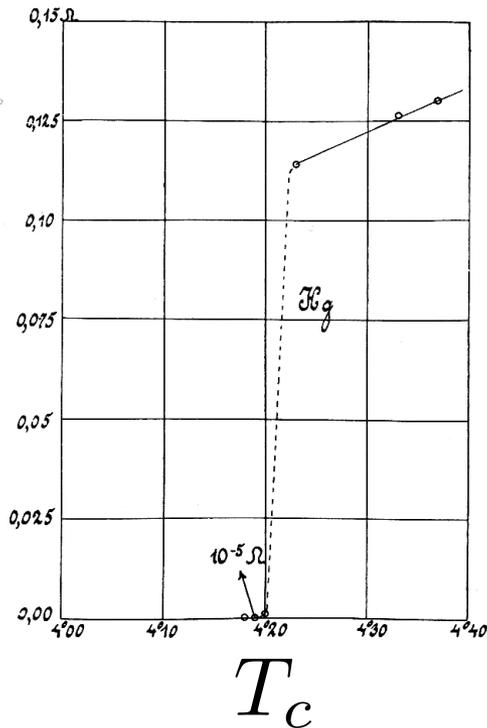
$$\alpha \propto T - T_c$$



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From Euclid to Minkowski

Can “real” QFTs emerge in condensed matter?

From Euclid to Minkowski

Can “real” QFTs emerge in condensed matter? Yes!

- Relativistic QFTs: Weyl/Dirac fermions and axion electrodynamics in CMP
- Supersymmetric QFTs: Higgsinos and gauginos in CMP?



Banff International Research Station

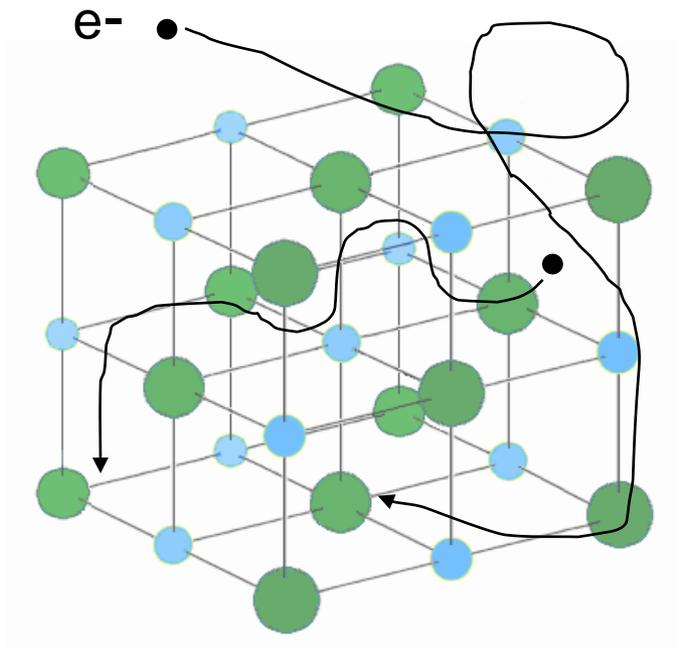
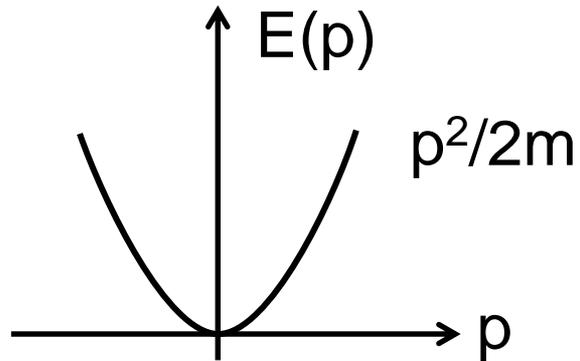
for Mathematical Innovation and Discovery

Relativistic Fermions and Nodal Semimetals from Topology

February 11-16, 2018

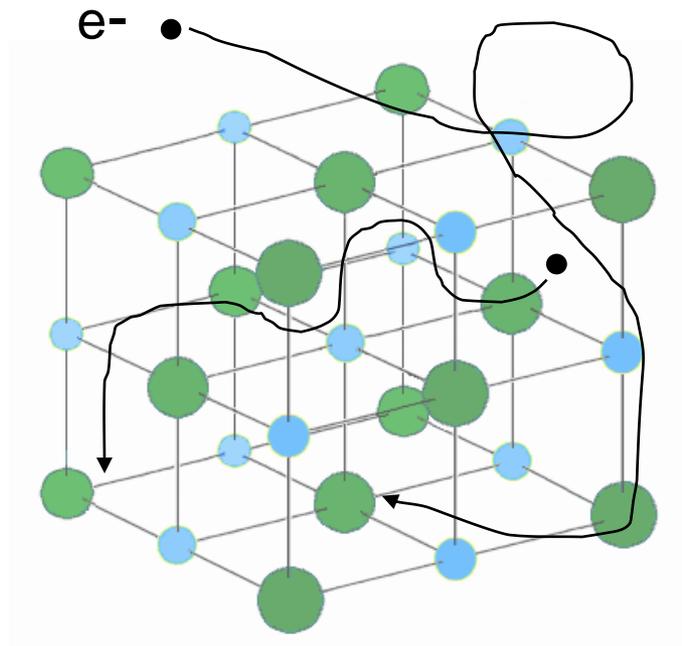
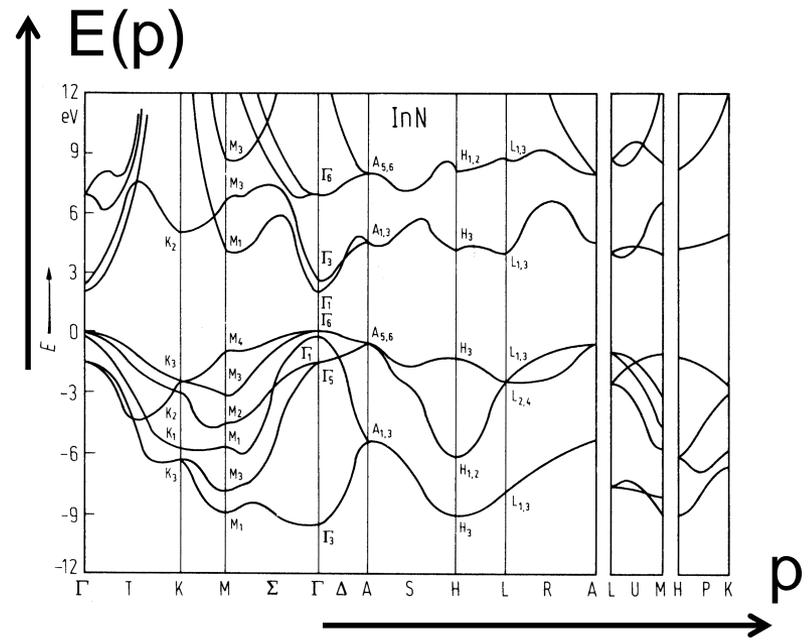
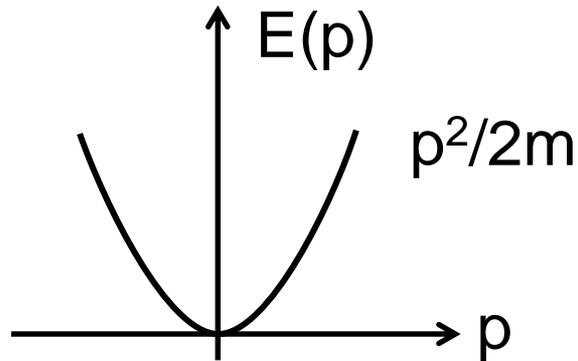


Band theory of solids

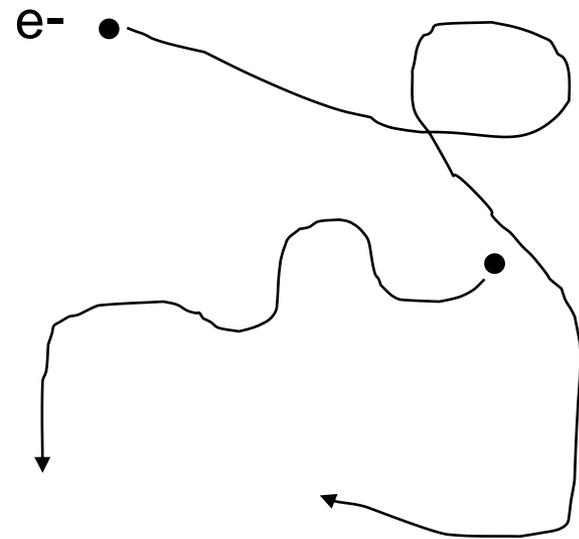


Bloch, Wilson, Peierls, Wigner... ('20-'30s)

Band theory of solids

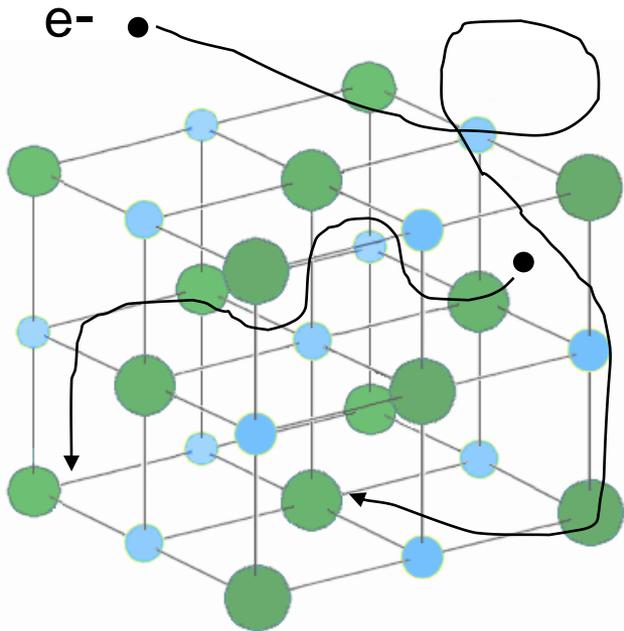
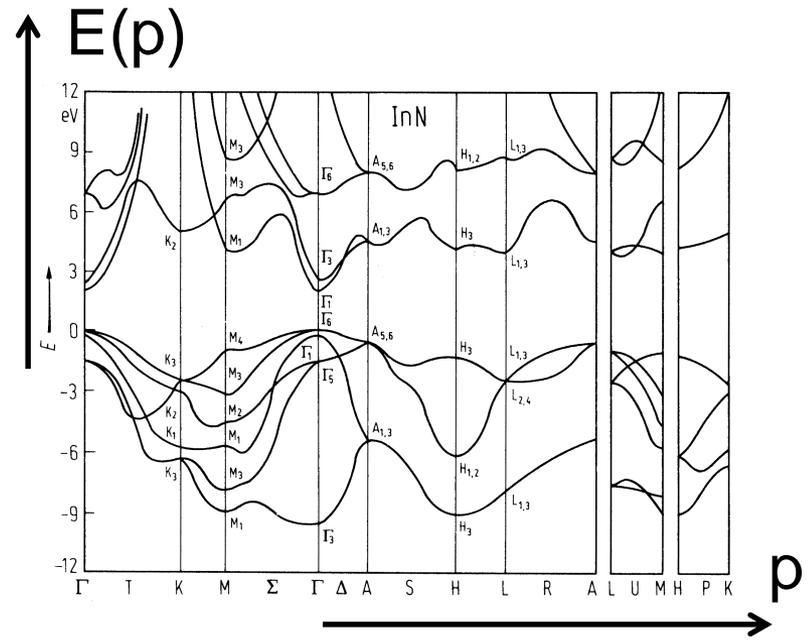
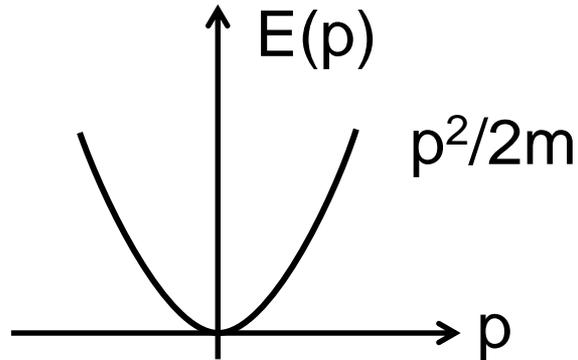


\approx

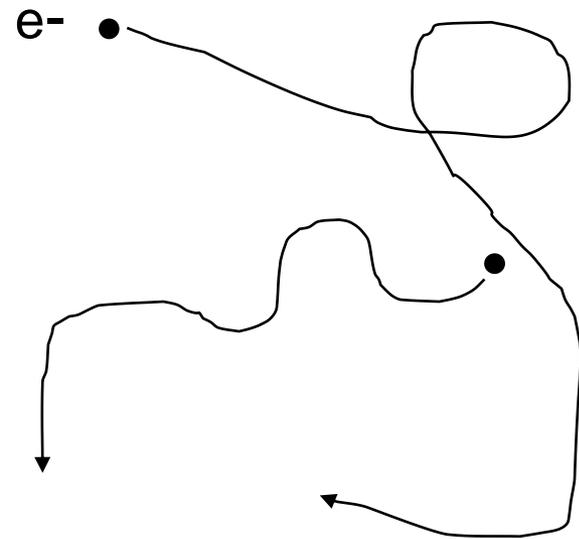


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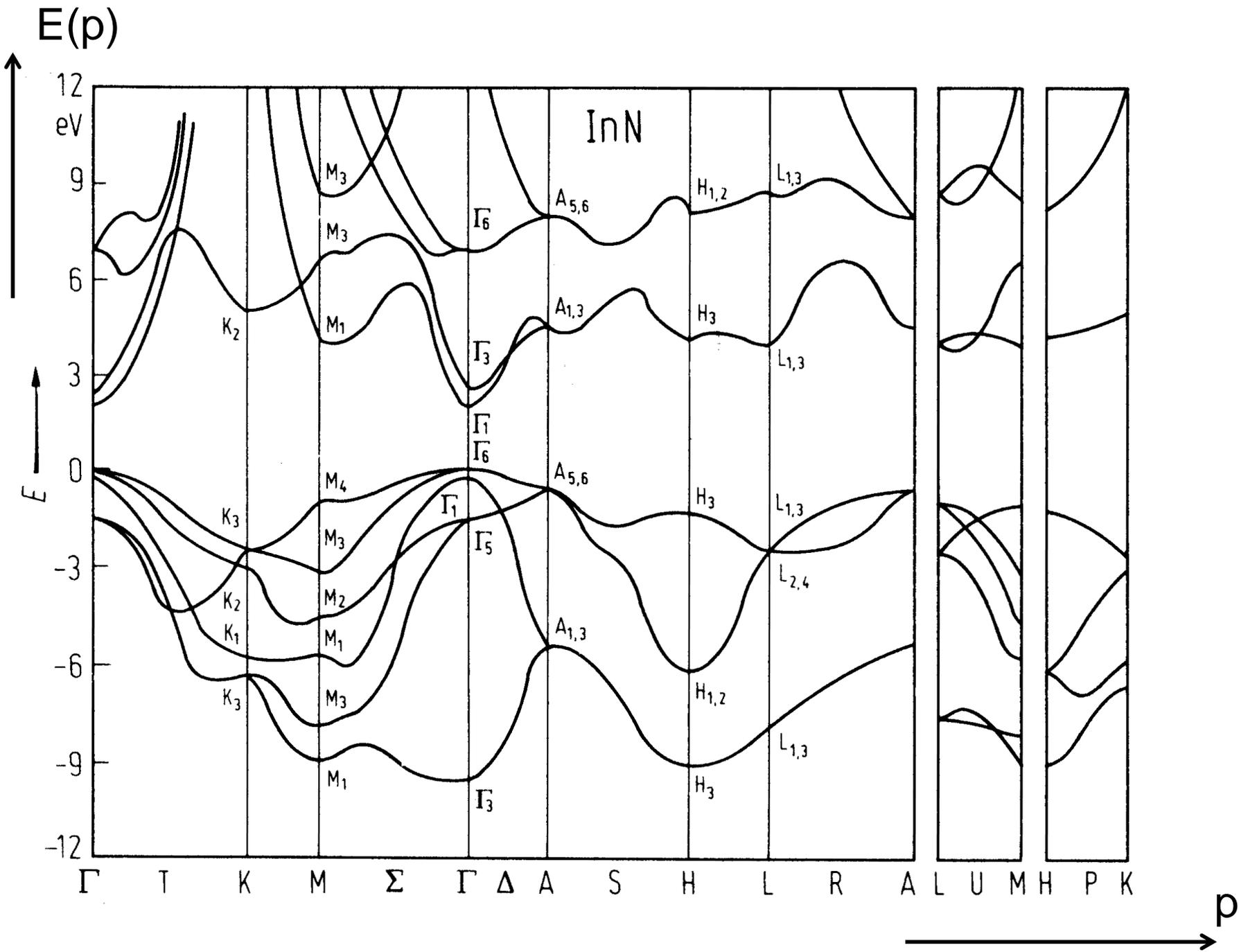
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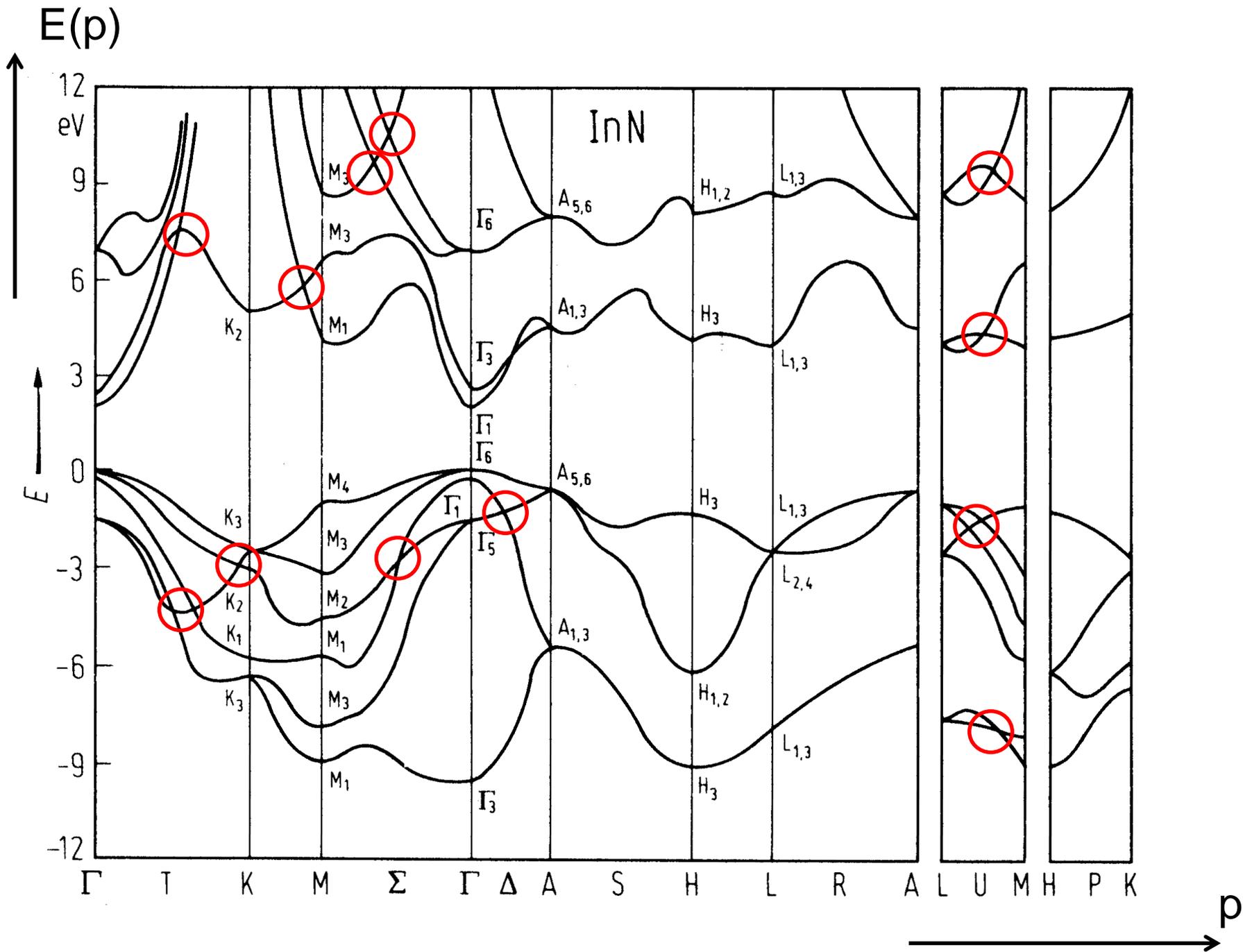


\approx

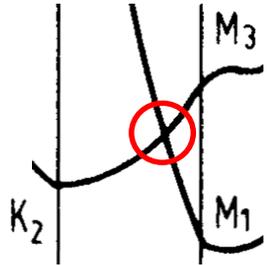


Each crystal is a new “vacuum” with different properties





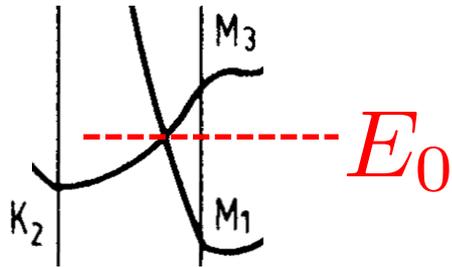
von Neumann-Wigner theorem



- “Level repulsion” in QM... why so many accidental degeneracies?



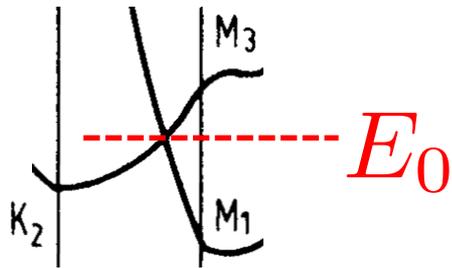
von Neumann-Wigner theorem



- “Level repulsion” in QM... why so many accidental degeneracies?

$$H = \begin{pmatrix} E_0 & 0 \\ 0 & E_0 \end{pmatrix}$$

von Neumann-Wigner theorem

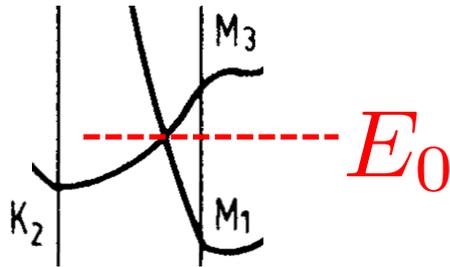


- “Level repulsion” in QM... why so many accidental degeneracies?

$$H = \begin{pmatrix} E_0 + \Delta_3 & \Delta_1 - i\Delta_2 \\ \Delta_1 + i\Delta_2 & E_0 - \Delta_3 \end{pmatrix}$$

$$E_{\pm} = E_0 \pm \sqrt{\Delta_1^2 + \Delta_2^2 + \Delta_3^2}$$

von Neumann-Wigner theorem



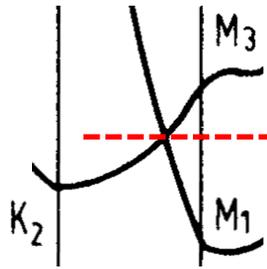
- “Level repulsion” in QM... why so many accidental degeneracies?

$$H = \begin{pmatrix} E_0 + \Delta_3 & \Delta_1 - i\Delta_2 \\ \Delta_1 + i\Delta_2 & E_0 - \Delta_3 \end{pmatrix}$$

$$E_{\pm} = E_0 \pm \sqrt{\Delta_1^2 + \Delta_2^2 + \Delta_3^2}$$

- In 3D, can adjust p_x, p_y, p_z to set $\Delta_1 = \Delta_2 = \Delta_3 = 0$
- 2-fold degenerate **points** with **linear** dispersion are generic (“topologically stable”)

Weyl fermions in CMP

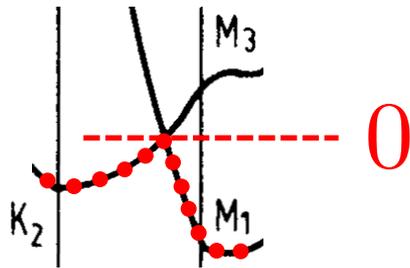


E_0

- Weyl fermion with “speed of light” = $|v|$ and chirality = $\text{sign}(v)$

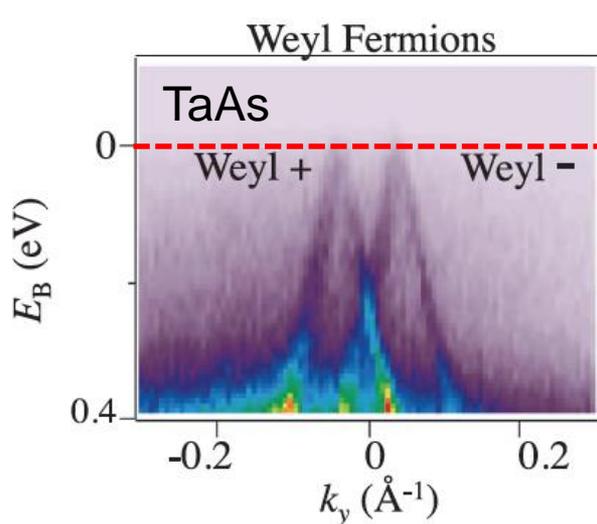
$$H = E_0 + v\boldsymbol{\sigma} \cdot \boldsymbol{p}$$

Weyl fermions in CMP

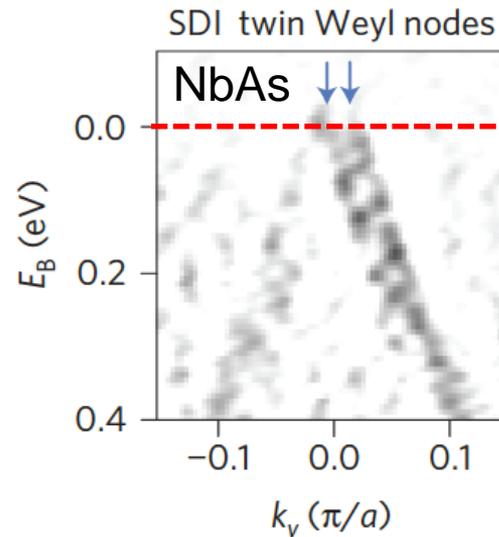


$$H = v\sigma \cdot p$$

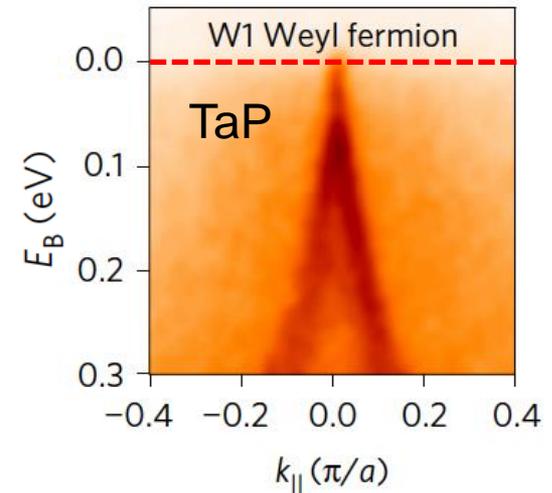
- Weyl fermion with “speed of light” = $|v|$ and chirality = $\text{sign}(v)$
- To mimic relativistic vacuum, Fermi level must be at the Weyl point: “Weyl semimetal”



Xu et al., Science '15



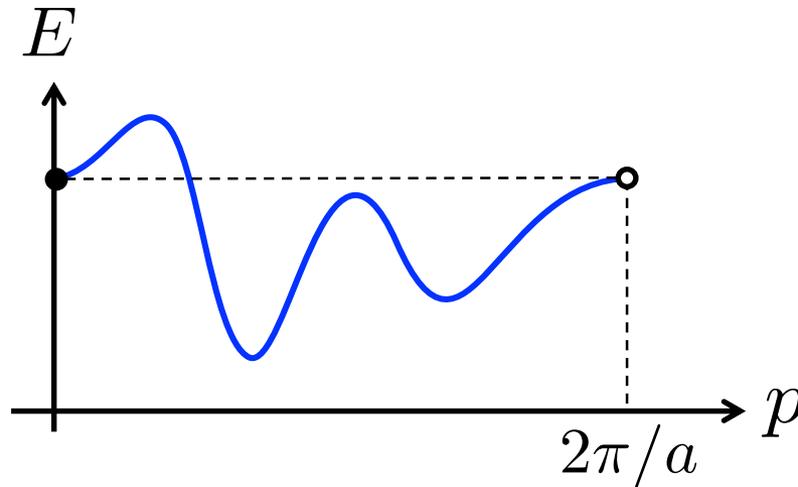
Xu et al., Nature Phys. '15



Zhang et al., Nature Phys. '17

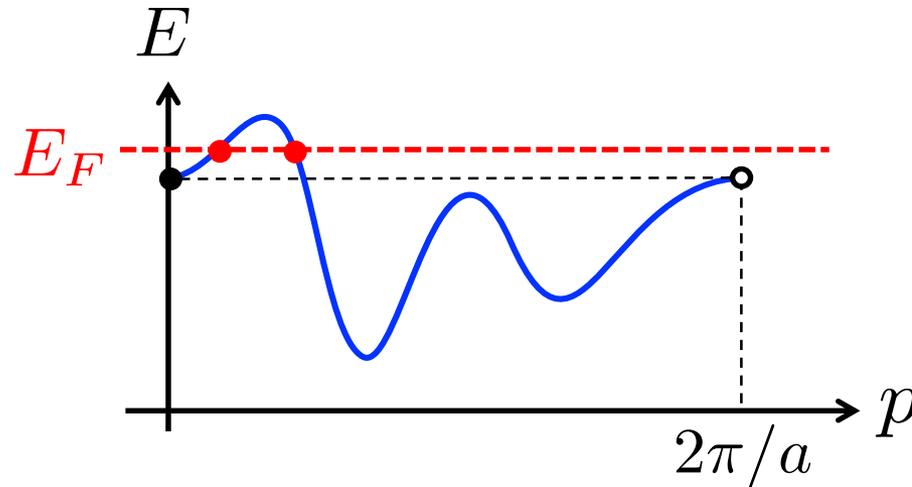
Chiral fermions on a lattice?

- In CMP, Weyl fermions always appear in R/L pairs
- Well-known from lattice gauge theory! “Fermion doubling problem”
- $E(p)$ must be **periodic** on a lattice, e.g. in 1D:



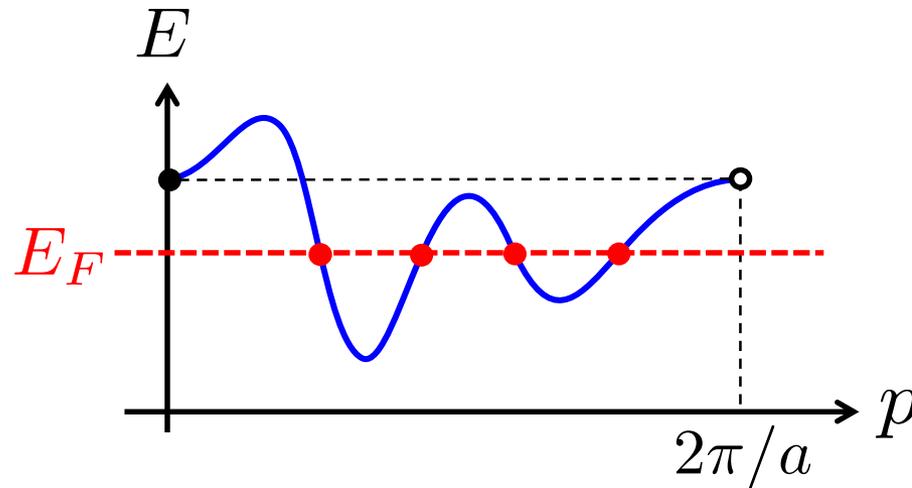
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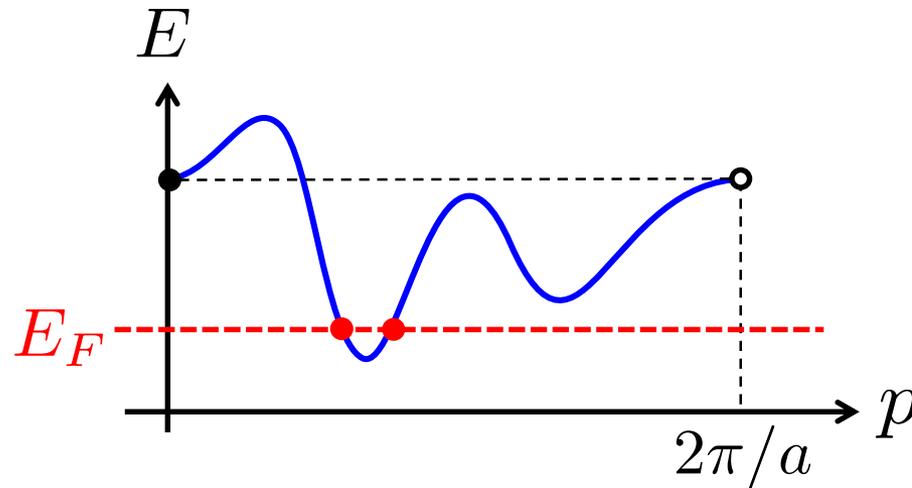
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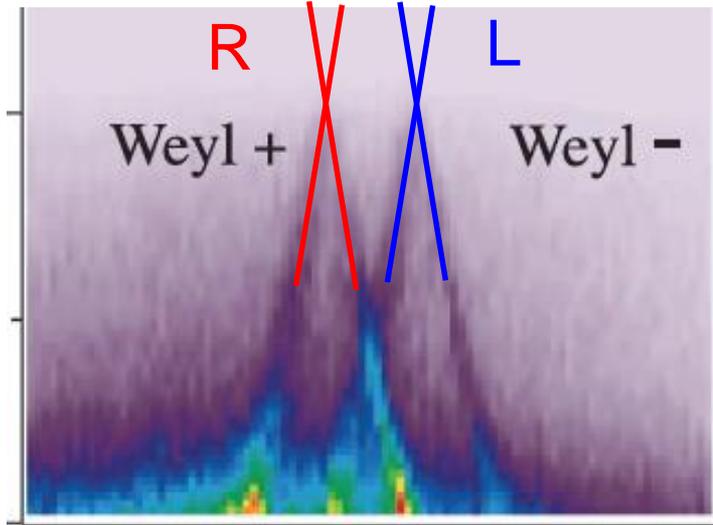
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- Analogous in 3D (*Nielsen & Ninomiya, NPB '81*)

Emergent chiral symmetry

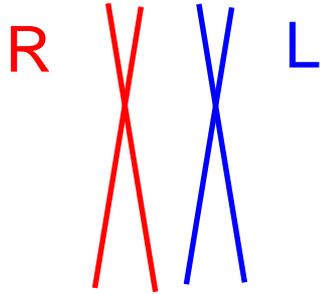


- For energies near E_F , **emergent** chiral symmetry protected by vN-W theorem

$$\psi_R \rightarrow e^{i\alpha} \psi_R$$

$$\psi_L \rightarrow e^{-i\alpha} \psi_L$$

Emergent chiral symmetry



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- What if R/L Weyl fermions at the same momentum point?

Broken chiral symmetry



- For energies near E_F , **emergent** chiral symmetry protected by vN-W theorem

$$\psi_R \rightarrow e^{i\alpha} \psi_R$$

$$\psi_L \rightarrow e^{-i\alpha} \psi_L$$

- What if R/L Weyl fermions at the same momentum point?
- Fourfold degeneracy, not stable! Dirac mass generally allowed, (explicitly) broken chiral symmetry

$$m \bar{\Psi} e^{i\theta \gamma^5} \Psi = m \cos \theta \bar{\Psi} \Psi + i m \sin \theta \bar{\Psi} \gamma^5 \Psi$$

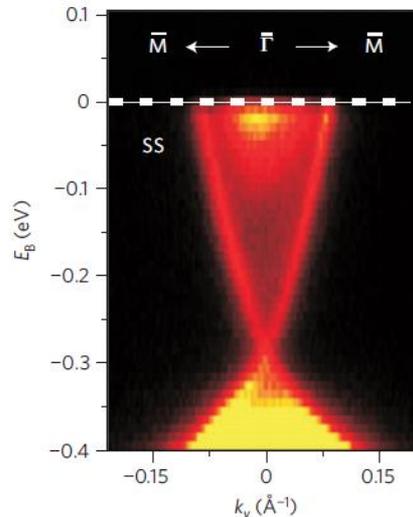
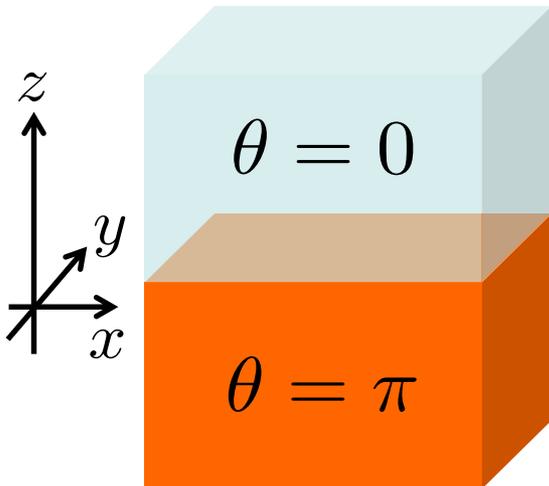
normal mass

axial mass

θ and T

$$m\bar{\Psi}e^{i\theta\gamma^5}\Psi = m\cos\theta\bar{\Psi}\Psi + im\sin\theta\bar{\Psi}\gamma^5\Psi$$

- Most materials have T symmetry: only $\theta=0$ or $\theta=\pi$
- Interface between $\theta=0$ and $\theta=\pi$: mass domain wall, traps fermion zero mode (*Jackiw & Rebbi, PRD '76*)

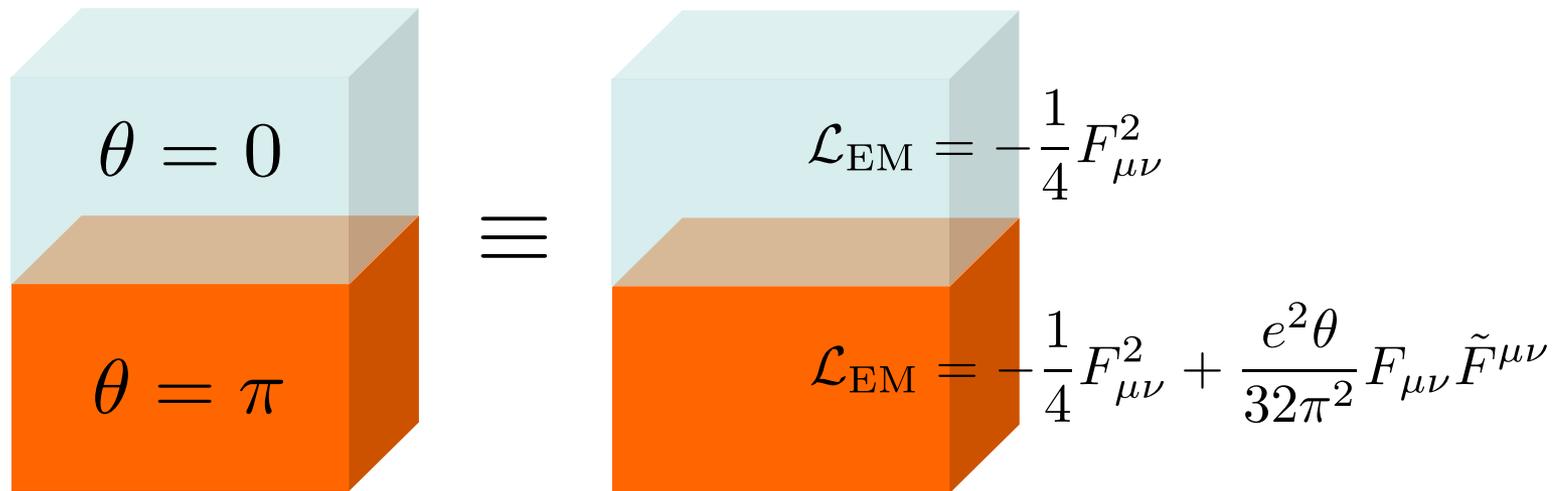


- Zero mode = massless **(2+1)D** Dirac fermion!
- $\theta=\pi$ materials = “topological insulators”

Xia et al., Nature Phys. '09

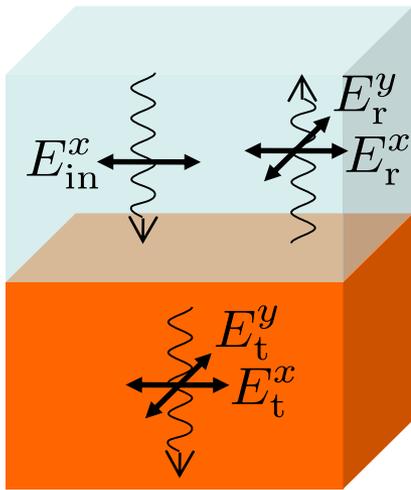
Axion electrodynamics

- Coupling to E&M: chiral rotation converts the mass phase to an electromagnetic **theta angle** $\theta=\pi$



Axion electrodynamics

- Theta angle can be measured via Kerr/Faraday effect

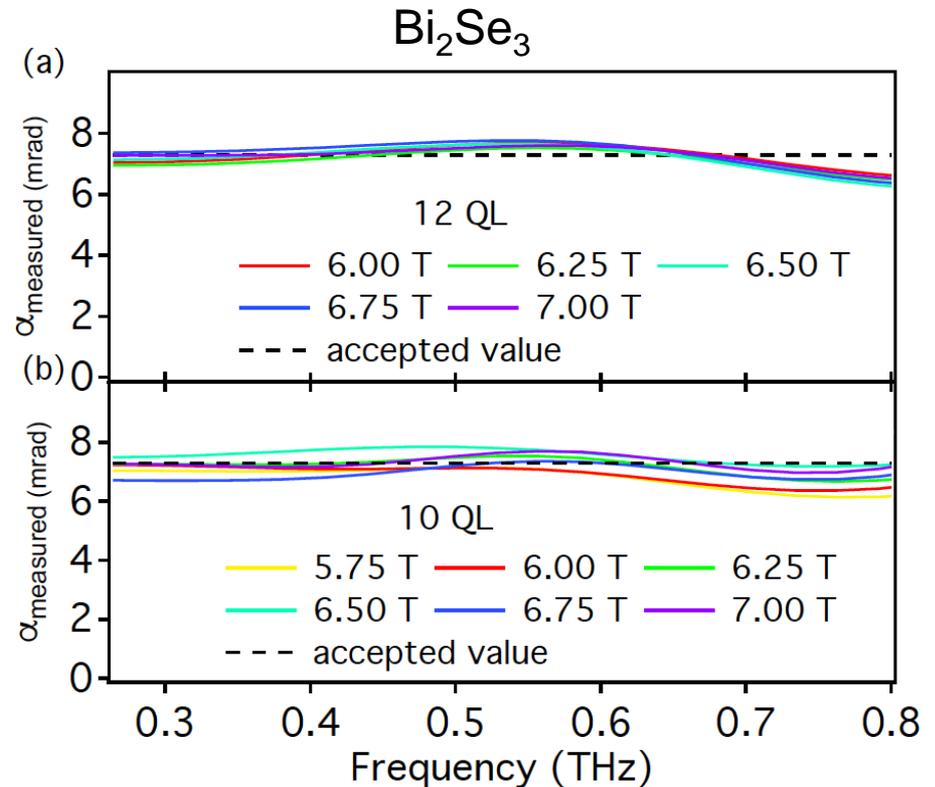


$$\tan \theta_K = \frac{E_r^y}{E_r^x}$$

$$\tan \theta_F = \frac{E_t^y}{E_t^x}$$

$$\frac{\cot \theta_F + \cot \theta_K}{1 + \cot^2 \theta_F} = \alpha$$

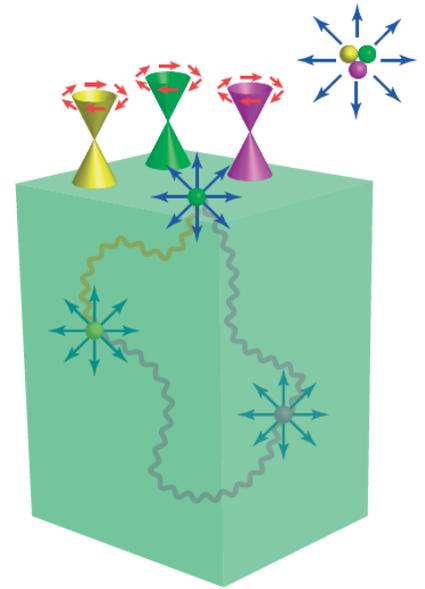
JM et al., PRL '10



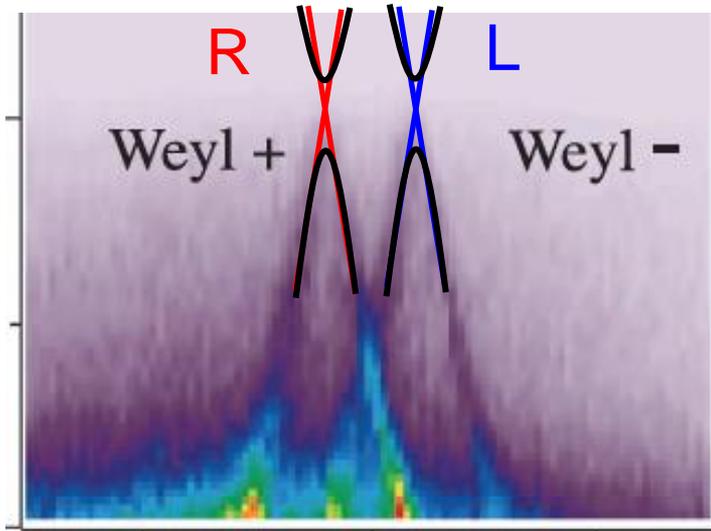
Wu et al., Science '16

θ and T, revisited

- EM θ angle different from 0, π can be consistent with T, if there exist **deconfined, fractionally charged** collective excitations
- Example: $e/3$ quasiparticles in the fractional quantum Hall effect (*Laughlin, PRL '83*)
- Quasiparticles with charge e/q : $\theta = \pi \frac{p}{q}$, p, q odd
- Reason: modifies Dirac quantization condition = changes periodicity of θ angle
- Ongoing search for “fractional” topological insulators



Dark matter on a chip?

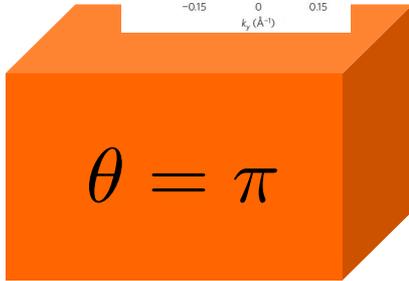
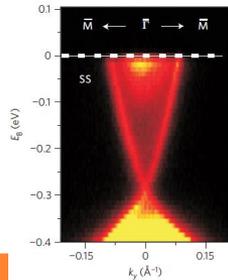


- Chiral symmetry in Weyl semimetals can in principle be broken **spontaneously** as a result of electron-electron interactions

- Pseudoscalar Goldstone mode = “dynamical axion field”
- Topological defects in the broken symmetry phase = “axion strings”, trap chiral fermion zero modes (*Callan & Harvey, NPB '85*)
- Search for interacting Weyl materials (CeRu_4Sn_6 , CeSb , ...)

SUSY in condensed matter?

Superconducting Dirac fermions



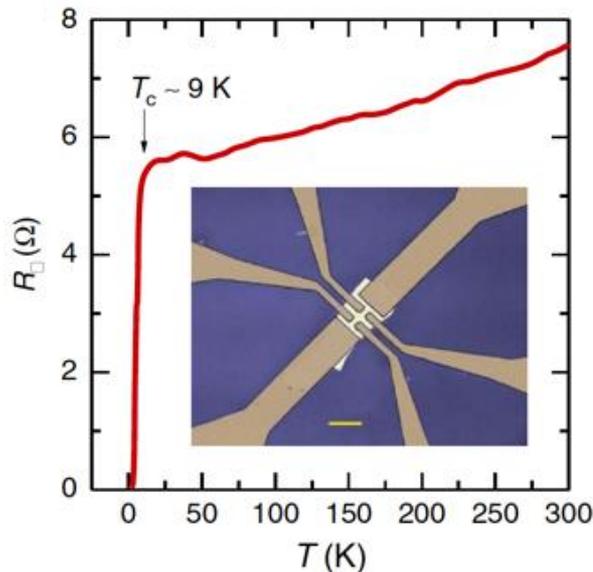
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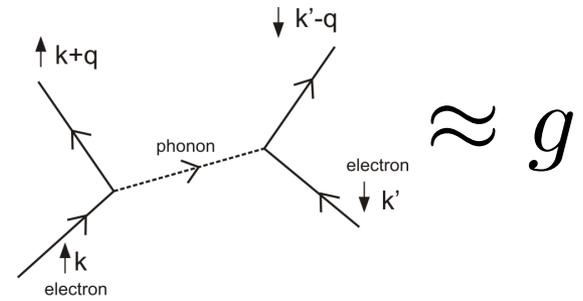
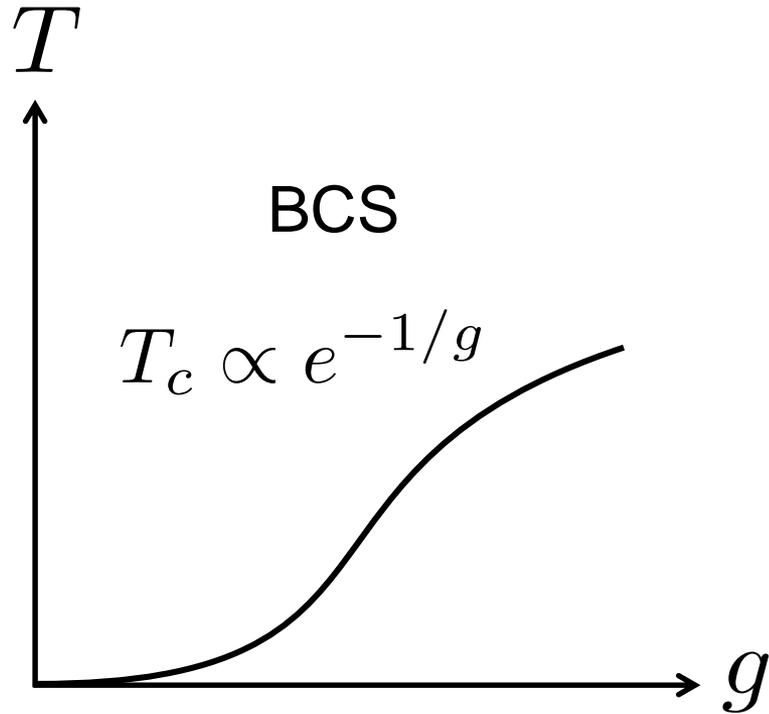
Emergent surface superconductivity in the topological insulator Sb_2Te_3

Lukas Zhao¹, Haiming Deng¹, Inna Korzhovska¹, Milan Begliarbekov¹, Zhiyi Chen¹, Erick Andrade², Ethan Rosenthal², Abhay Pasupathy², Vadim Oganesyan^{3,4} & Lia Krusin-Elbaum^{1,4}

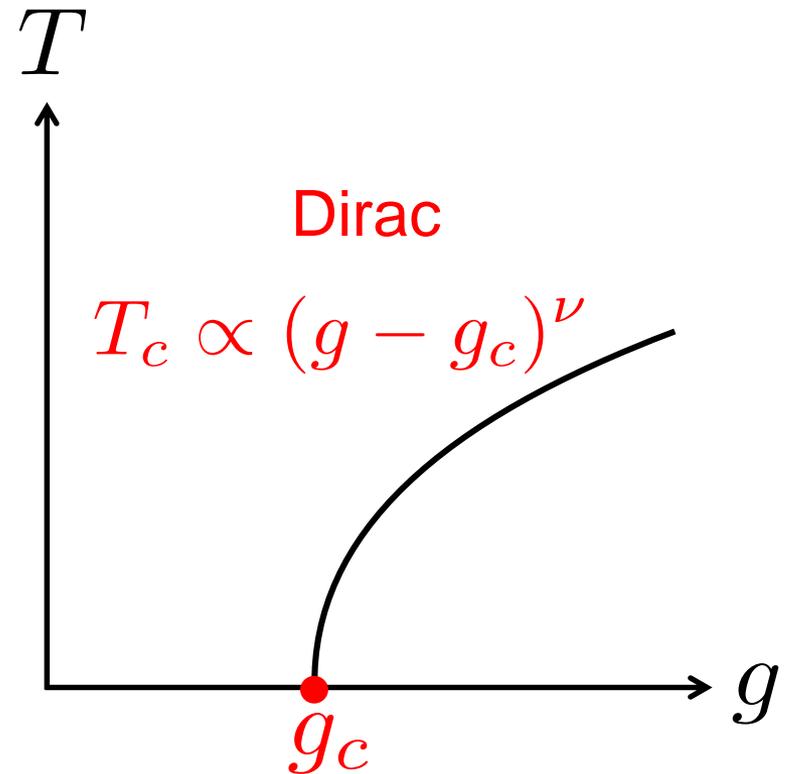
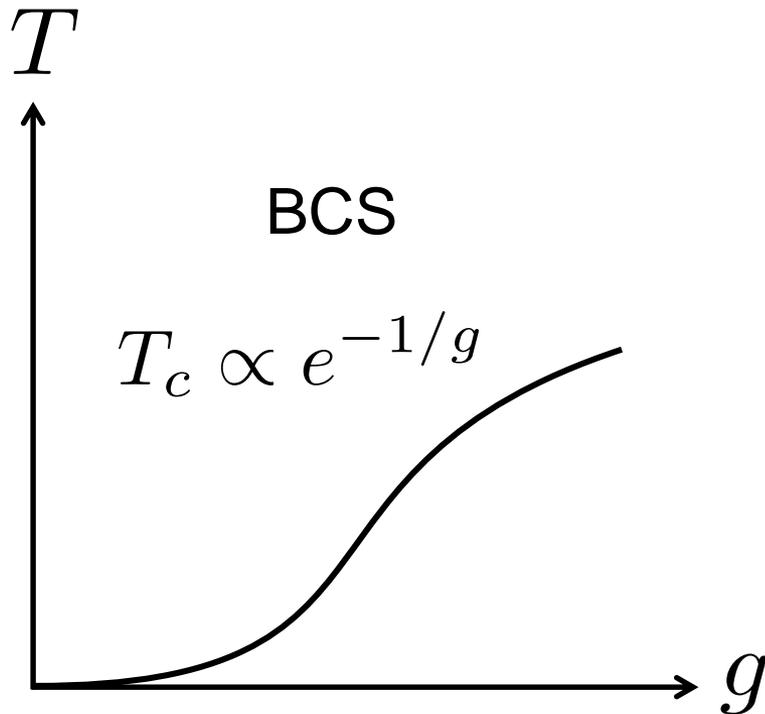


- What is special about superconducting Dirac fermions in 2+1 dimensions?

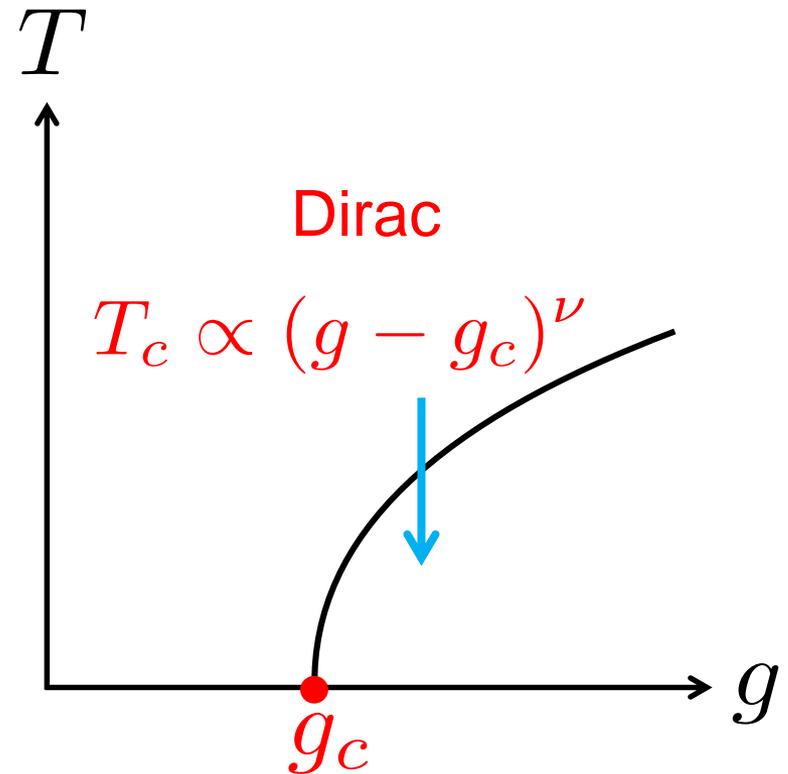
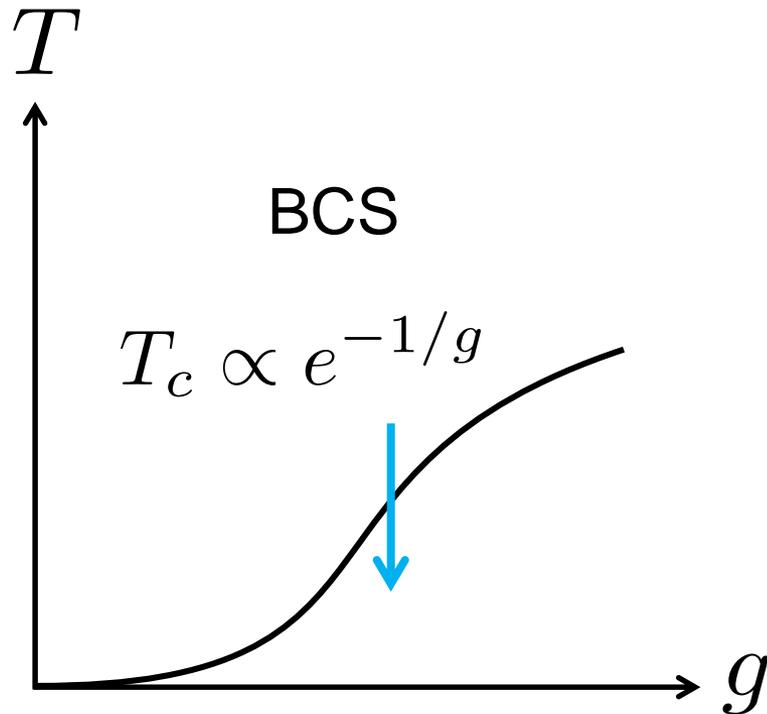
BCS vs Dirac



BCS vs Dirac

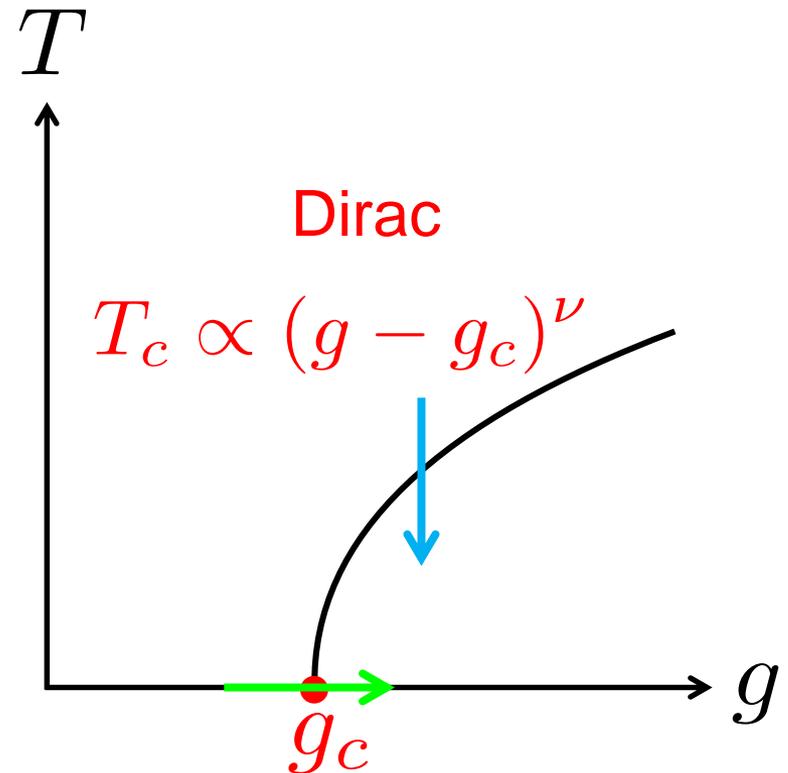
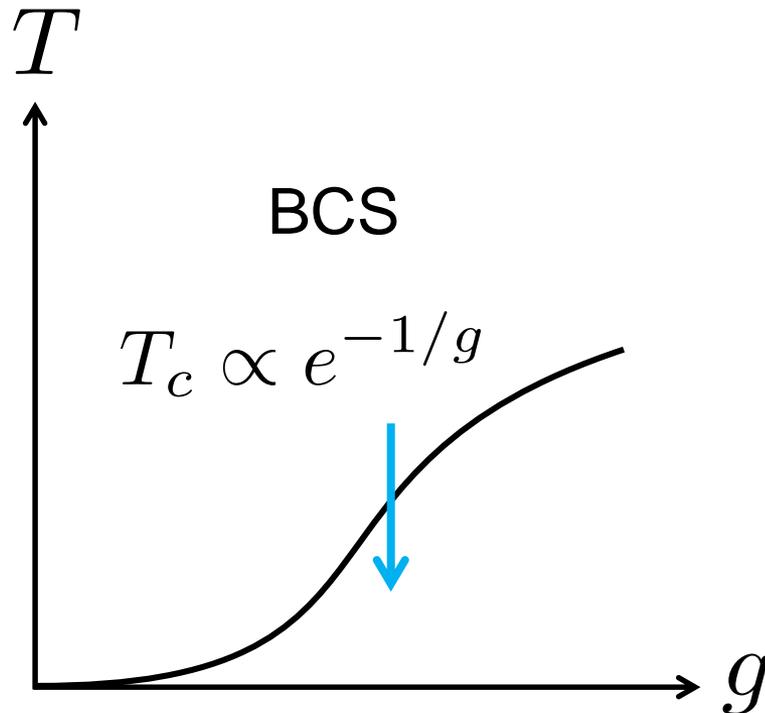


BCS vs Dirac



- Thermal phase transition similar to BCS (static GL)

BCS vs Dirac



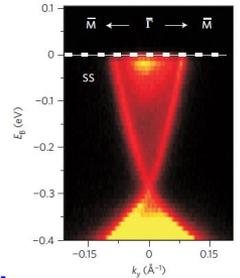
- Thermal phase transition similar to BCS (static GL)
- Dirac: additional **quantum phase transition** ($T=0$)!

Quantum Ginzburg-Landau theory

- Quantum phase transition: driven by **quantum** rather than thermal fluctuations
- Described by **time-dependent** GL action:

$$S = \int dx dt \left[i\bar{\psi}(\gamma^0 \partial_t + c_f \boldsymbol{\gamma} \cdot \nabla) \psi \right]$$

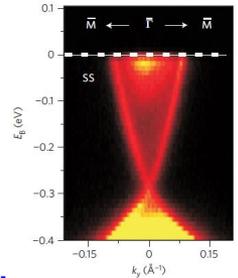
Dirac fermion ψ



Quantum Ginzburg-Landau theory

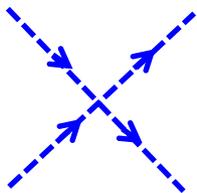
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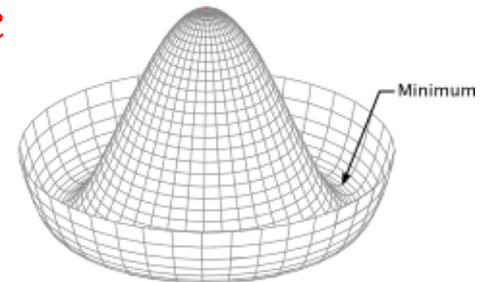
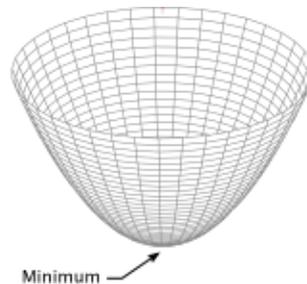
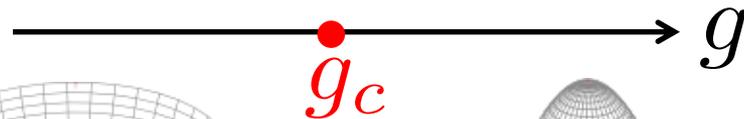


$$+ |\partial_t \phi|^2 - c_b^2 |\nabla \phi|^2 - \alpha |\phi|^2 - \beta |\phi|^4$$

Cooper pair field ϕ ----->



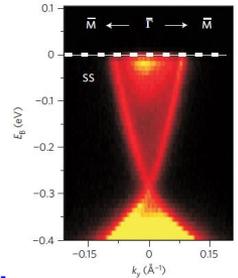
$$\alpha \propto g_c - g$$



Quantum Ginzburg-Landau theory

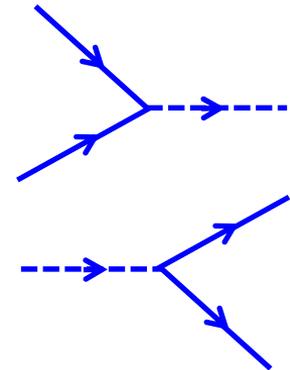
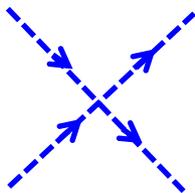
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$$S = \int dx dt \left[i\bar{\psi}(\gamma^0 \partial_t + c_f \boldsymbol{\gamma} \cdot \nabla) \psi + |\partial_t \phi|^2 - c_b^2 |\nabla \phi|^2 - \alpha |\phi|^2 - \beta |\phi|^4 + h(\phi^* \psi_\uparrow \psi_\downarrow + \text{c.c.}) \right] \quad \text{pair formation/breaking}$$

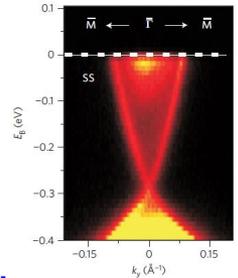
Cooper pair field ϕ



Quantum Ginzburg-Landau theory

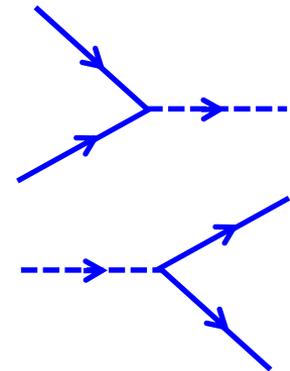
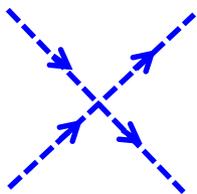
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$$+ |\partial_t \phi|^2 - c_b^2 |\nabla \phi|^2 - \alpha |\phi|^2 - \beta |\phi|^4 + h(\phi^* \psi_\uparrow \psi_\downarrow + \text{c.c.}) \quad \text{Cooper pair field } \phi$$

pair formation/breaking



- **True QFT**, but not (yet) relativistic... two velocities

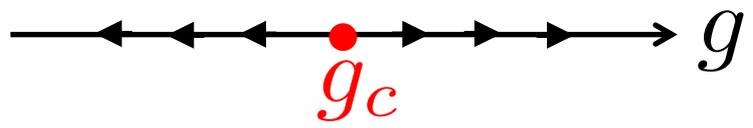
Quantum critical point

- Quantum critical point =
(IR unstable) RG fixed point,
scale invariance in space **and** time



- 1-loop RG: fermion/boson velocities flow to a common value = **emergent Lorentz invariance**

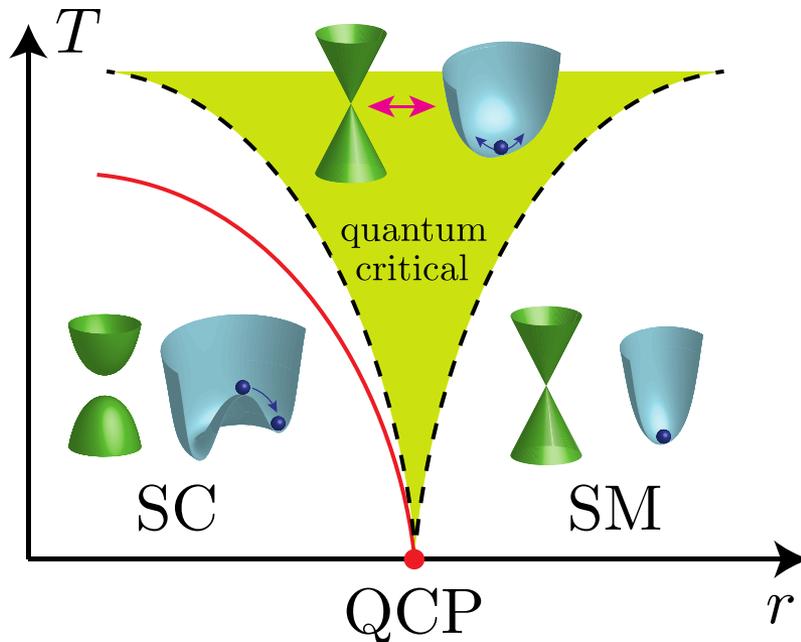
Quantum critical point

- Quantum critical point = 
(IR unstable) RG fixed point,
scale invariance in space **and** time
- 1-loop RG: fermion/boson velocities flow to a common value = **emergent Lorentz invariance**
- In fact, $|\phi|^4$ and Yukawa couplings also flow to common value: **emergent N=2 SUSY!** (*Grover, Sheng, Vishwanath, Science '14; Ponte & Lee, NJP '14*)
- Remains true at 3-loop level (*Zerf, Lin, JM, PRB '16*)

A new universality class

$$\mathcal{L} = i\bar{\psi}\gamma_{\mu}\partial_{\mu}\psi + |\partial_{\mu}\phi|^2 + r|\phi|^2 + h^2|\phi|^4 + h(\phi^*\psi^T i\sigma^y\psi + \text{h.c.})$$

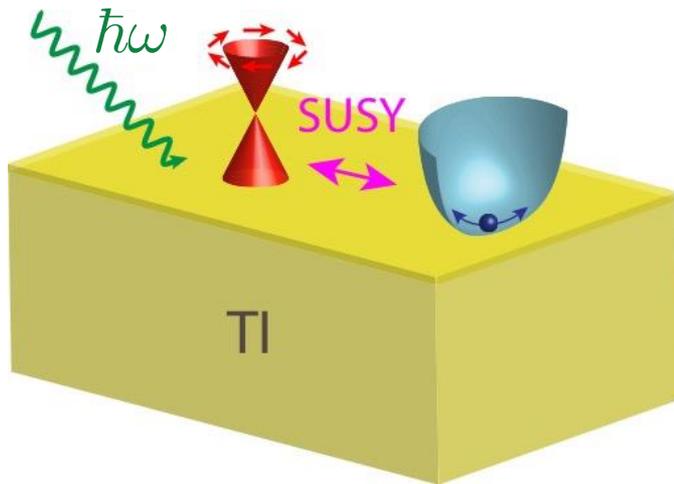
- N=2 Wess-Zumino model = 1 chiral multiplet X with $W=X^3$. Dirac fermion and Cooper pair are superpartners!



- “SUSY version” of usual Ginzburg-Landau theory
- Thermodynamic critical exponents known at 4 loops (Zerf, Lin, JM, PRB ‘16; Zerf et al., PRD ‘17)

A new quantum universality class

- **Quantum** phase transition: besides (static) critical exponents, also universal **dynamical** properties
- In (2+1)D, AC conductivity $\sigma(\omega)$ is a **universal constant** at quantum critical points (*Damle & Sachdev, PRB '97*)

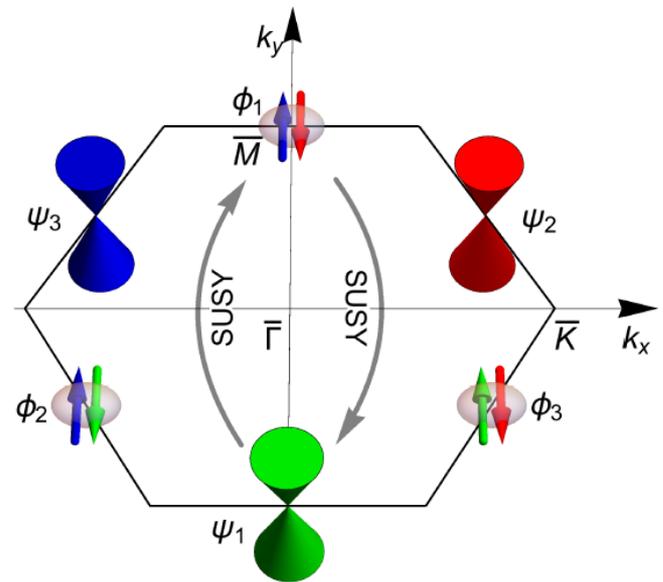


- Can be computed exactly for this transition despite strong coupling, unprecedented in $D > 1+1$

$$\sigma(\omega) = \frac{5(16\pi - 9\sqrt{3})}{243\pi} \frac{e^2}{\hbar} \approx 0.2271 \frac{e^2}{\hbar}$$

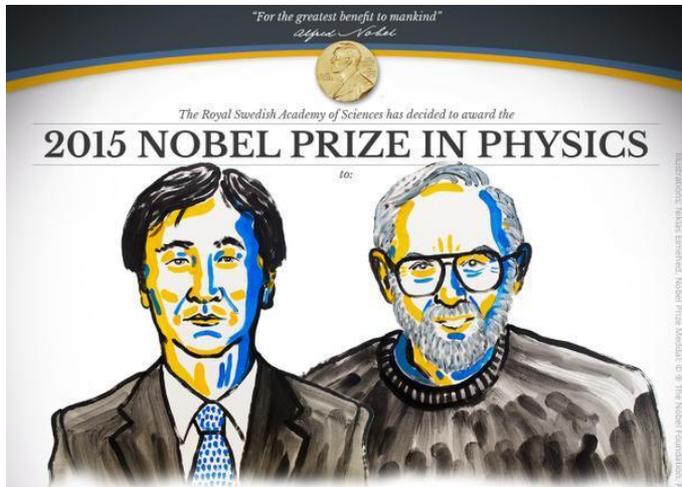
SUSY and mirror symmetry

- Topological insulators have **odd #** of surface Dirac fermions
- **3** Dirac fermions: different universality class = XYZ model (*Aharony et al., NPB '97*)
- Infrared dual to $N=2$ SQED₃ with one “flavor” of matter fields
- Bosonic Cooper pairs & Dirac fermions are dual to U(1) vortices, “photons” and “photinos”

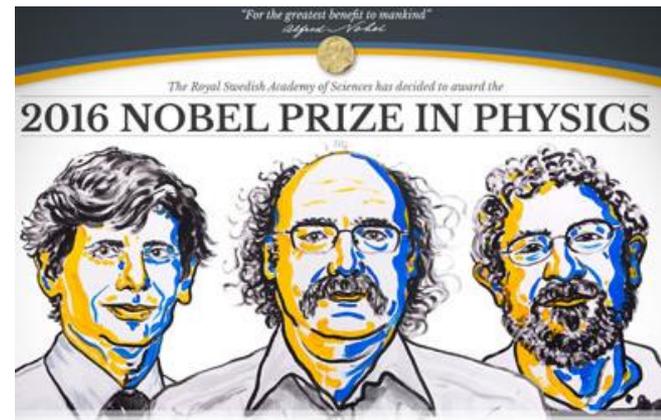


Conclusion

- An exciting time for HEP and CMP alike!
- Theory predictions inspired by HEP have led to “topological revolution” in CMP
- “Dirt” as potential platform to observe exotic (emergent) HEP phenomenology?



neutrino oscillations



topological phases of matter

*To see a World in a Grain of Sand
And a Heaven in a Wild Flower
Hold Infinity in the palm of your hand
And Eternity in an hour*

(William Blake, *Auguries of Innocence*, 1863)