

Holography - hints from  
classical gravity

The gravitational phase  
space as a tool to construct  
holographic dualities

Conclusions

# Holography and Applications

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# An unorthodox 'derivation' of AdS/CFT

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- Let us imagine we are back in 1996 and know nothing about AdS/CFT.
- We can go a long way in 'deriving' the AdS/CFT correspondence by studying classical AdS gravity:

$$R_{\mu\nu} - \frac{1}{2}(R - 2\Lambda)g_{\mu\nu} = 0$$

- A mathematically interesting question to ask is what is the space of solutions, i.e. the classical phase space, of such an equation.
- Finding the space of non-singular solutions is an extremely hard problem in dimensions higher than 3. In 3 dimensions it is given by the Teichmüller space of Riemann surfaces [K. Krasnov, ...]
- The problem remains hard if one allows for singular solutions, but progress can be made...

# The Fefferman-Graham theorem

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## Fefferman-Graham theorem

- In an open neighborhood of the conformal boundary at  $r = \infty$  the most general solution of Einstein's equation with a negative cosmological constant takes the form

$$ds^2 = dr^2 + \gamma_{ij}(r, x)dx^i dx^j, \quad i, j = 1, \dots, d$$

where

$$\gamma_{ij} = e^{2r} \left( g_{(0)ij} + e^{-2r} g_{(2)ij} + \dots + g_{(d)ij} e^{-dr} + h_{(d)ij} e^{-dr} r + \dots \right)$$

- The tensors  $g_{(0)ij}(x)$  and the transverse traceless part of  $g_{(d)ij}(x)$  can be independently prescribed and they *uniquely* determine the bulk metric in an open neighborhood of the conformal boundary.
- In particular,  $g_{(2)}, \dots, g_{(d-2)}$  and  $h_{(d)}$  are algebraically determined in terms of  $g_{(0)}$  and its derivatives, while all higher order terms are determined in terms of  $g_{(0)}$  and  $g_{(d)}$ .

# Classical phase space

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- The space parameterized by the tensors  $g_{(0)ij}$  and  $g_{(d)}$  is a symplectic manifold equipped with the symplectic form

$$\Omega = \int d^d x \delta g_{(0)ij} \wedge \delta g_{(d)}^{ij}$$

- This classical phase space contains singular solutions as a generic assignment of  $g_{(0)}$  and  $g_{(d)}$  leads to a singularity at some point far enough from the conformal boundary.
- There is nevertheless enough structure to allow us to 'derive' the kinematics of AdS/CFT...

- Einstein's equation implies that  $g_{(d)}$  must satisfy the constraints:

$$g_{(0)}{}^{ij}g_{(d)ij} = \mathcal{A}[g_{(0)}], \quad D_{(0)}{}^i \left( g_{(d)ij} - g_{(0)}{}^{kl}g_{(d)kl}g_{(0)ij} \right) = 0$$

where  $\mathcal{A}[g_{(0)}]$  is a local functional of  $g_{(0)}$  that coincides with the conformal anomaly of a  $d$ -dimensional conformal field theory! (with certain values for the anomaly coefficients.)

- Given a conformal Killing vector of the metric  $g_{(0)ij}$ , the above identities imply that the following charges are conserved:

$$Q_\xi = \int d\sigma^i \left( g_{(d)ij} - g_{(0)}{}^{kl}g_{(d)kl}g_{(0)ij} \right) \xi^j$$

- The Poisson algebra of these charges provides a representation of the asymptotic symmetry algebra  $so(2, d)$  on phase space. In three dimensions ( $d = 2$ ), this turns out to be two copies of a centrally extended Virasoro algebra [Brown-Henneaux]!

- We have seen that the quantity

$$\mathcal{T}_{ij}[g_{(0)}] = \left( g_{(d)ij} - g_{(0)}{}^{kl} g_{(d)kl} g_{(0)ij} \right)$$

has the kinematic properties required of the stress tensor of a conformal field theory coupled to the metric  $g_{(0)ij}$ . What about the dynamics?

- To compute the corresponding 'correlation functions' we need to evaluate the functional derivatives

$$\langle T(x_1) \dots T(x_n) \rangle \sim \frac{\delta \mathcal{T}[g_{(0)}(x_1)]}{\delta g_{(0)}(x_2) \dots \delta g_{(0)}(x_n)}$$

- Determining  $g_{(d)}[g_{(0)}]$  in full generality corresponds to finding the subspace of the above phase space that includes only *non-singular* solutions. We now see that it also corresponds to computing *all* correlation functions of the dual QFT!

- To compute the  $n$ -point function we need to compute  $g_{(d)}[g_{(0)}]$  to order  $n - 1$  in  $g_{(0)}$  around a background solution.
- Remarkably, the dynamics encoded in such correlation functions is consistent with that of a strongly coupled CFT!

#### Recent examples of CFT dynamics from AdS gravity

- 2-point function in the background of a black hole: quark-gluon plasma
- $g_{(d)}[g_{(0)}]$  in a derivative expansion around a black hole: relativistic Navier-Stokes equation [Bhattacharyya, Hubeny, Minwalla, Rangamani]

## To summarize...

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The classical phase space of AdS gravity has the structure of a CFT!



# Holographic dualities

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- Ultimately, an holographic duality must involve a quantum theory of gravity, such as string theory, on one side, and a non-gravitational quantum theory on the other.
- In systems that admit a semiclassical limit, such as the large  $N$  limit, one can expect that gravity captures some corner of the dynamics, around which systematic corrections can be computed.
- Within this class of systems, the systematic construction of the classical phase space of gravitational theories provides a method for 'constructing' holographic dualities in both directions:

# From QFT to gravity and back

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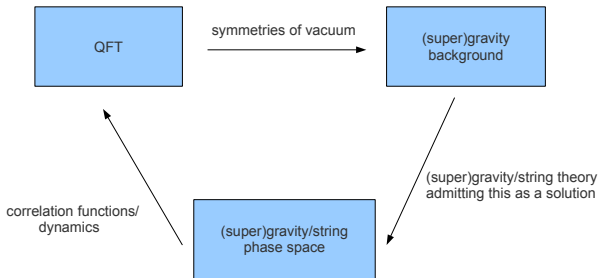
Conclusions

- There are many physically important strongly coupled QFTs that we would like to try to understand via holography.
- Prominent examples are QCD and non-relativistic CFTs that arise at quantum critical points in condensed matter physics.
- Starting with some basic knowledge, such as the symmetries of the vacuum, of such QFTs, one can try to construct a gravitational background as a starting point in finding the holographic dual.
- Constructing the gravitational phase space in such a background then will lead to dynamical information about the QFT...

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# From gravity to exotic QFTs

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- However, systematically constructing the classical phase space of gravity can lead to exotic QFTs, non-local quantum theories, or exotic UV completions.
- The Klebanov-Strassler and the Maldacena-Núñez backgrounds provide such examples of exotic UV completions to versions of  $\mathcal{N} = 1$  supersymmetric Yang-Mills theory.
- Moreover, one may hope that such an approach may shed light on the holographic dual of gravity in asymptotically flat spacetimes.

# Summary & Conclusions

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Classical gravity contains QFT!