

Gaps between jets and soft gluon resummation

Simone Marzani
University of Manchester

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In collaboration with Jeff Forshaw and James Keates

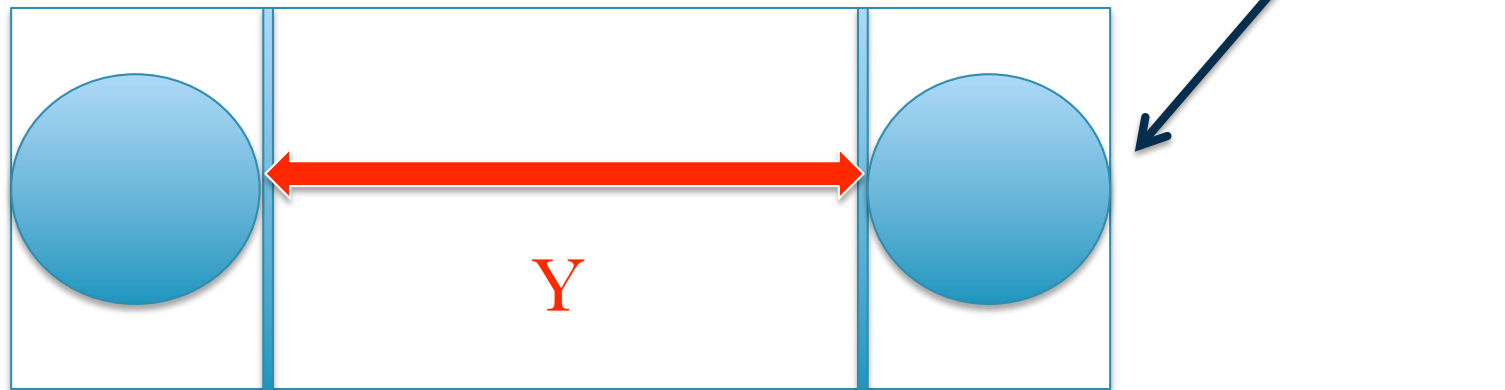
Outline

- Importance of rapidity gaps for LHC physics
- Soft gluons
- Discovery of super-leading logarithms
- Phenomenological studies and comparison to HERWIG++
- Conclusions and Outlook

The observable

Production of two jets with

- transverse momentum Q
- rapidity separation Y



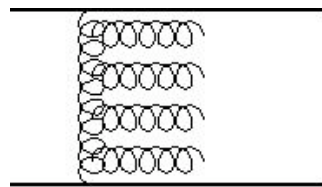
- Emission with $k_T > Q_0$
forbidden in the inter-jet region

Plenty of QCD effects

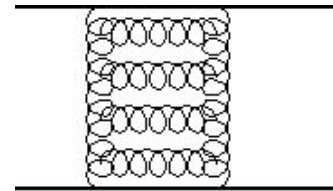
“wider” gaps

Y

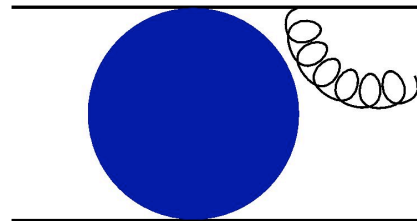
Forward BFKL
(Mueller-Navelet jets)



Non-forward BFKL
(Mueller-Tang jets)



Wide-angle soft
radiation

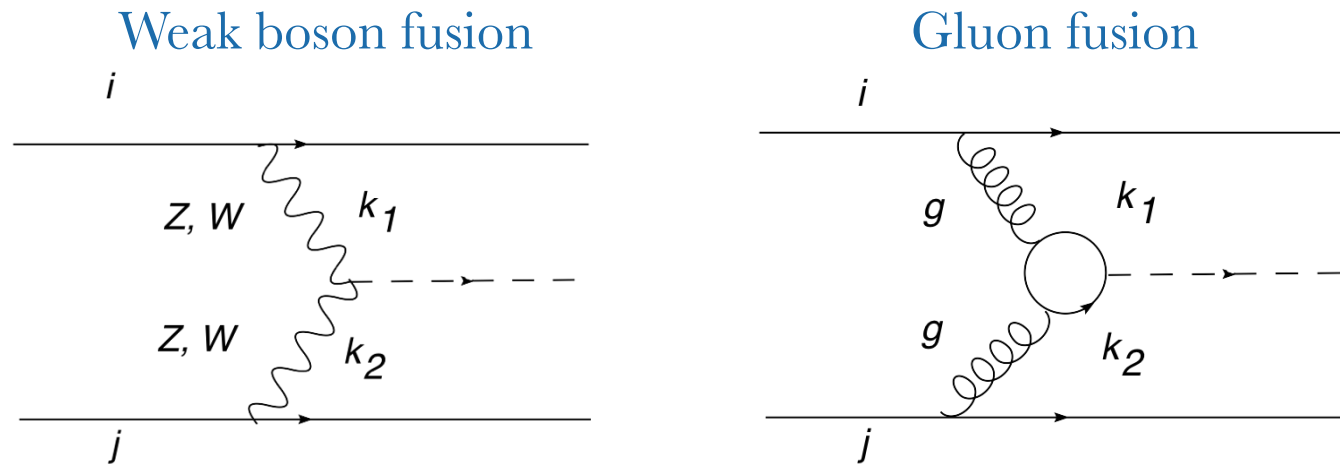


Fixed order

Super-leading
logs

$$L = \ln \frac{Q}{Q_0} \quad \text{“emptier” gaps}$$

Higgs + 2 jets

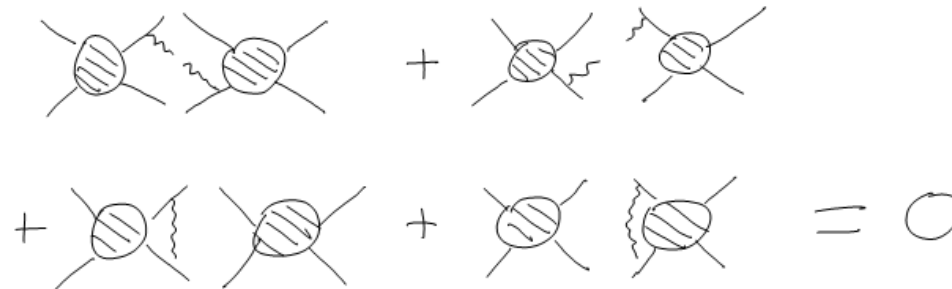


- Different QCD radiation in the inter-jet region
- To enhance the WBF channel, one can make a veto Q_0 on additional radiation between the tagged jets
- QCD radiation as in dijet production
- Important in order to extract the VVH coupling

Forshaw and Sjödal
arXiv:0705.1504 [hep-ph]

Soft gluons in QCD

- What happens if we dress a hard scattering with soft gluons?
- Sufficiently inclusive observables are not affected: real and virtual cancel via Bloch-Nordsieck theorem



- Soft gluon corrections are important if the real radiation is constrained into a small region of phase-space
- In such cases BN fails and miscancellation between real and virtual induces large logarithms

$$-\alpha_s \int_0^{Q_0} \frac{dE}{E} \Big|_{\text{real}} + \alpha_s \int_0^Q \frac{dE}{E} \Big|_{\text{virtual}} = \alpha_s \int_{Q_0}^Q \frac{dE}{E} \Big|_{\text{virtual}} = \alpha_s \ln \frac{Q}{Q_0}$$

Soft gluons in gaps between jets

- Naive application of BN:
real and virtual contributions cancel everywhere
except within the gap region for $k_T > Q_0$
- One only needs to consider **virtual corrections** with
$$Q_0 < k_T < Q$$
- Leading logs (LL) are resummed by iterating the one-loop result:

$$\mathcal{M} = e^{-\alpha_s L \Gamma} \mathcal{M}_0$$

soft anomalous dimension



Born



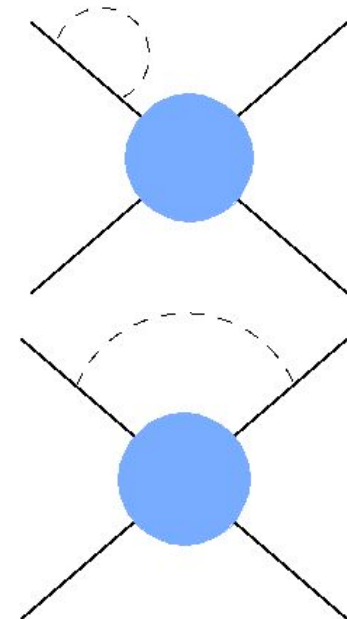
Oderda and Sterman
hep-ph/9806530

Colour evolution (I)

- The anomalous dimension can be written as

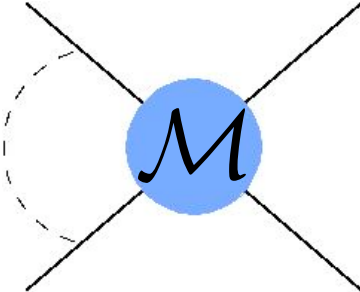
$$\Gamma = \frac{1}{2} Y T_t^2 + i\pi T_1 \cdot T_2 + \frac{1}{4} \rho (T_3^2 + T_4^2)$$

- T_i is the colour charge of parton i
- T_i^2 is a Casimir
- $T_t^2 = (T_1^2 + T_3^2 + 2T_1 \cdot T_3)$
is the colour exchange in the t -channel



Colour evolution (II)

- The $i\pi$ term is due to Coulomb gluon exchange

$$i\pi T_1 \cdot T_2 \mathcal{M} = \text{diagram}$$


- It doesn't play any role for processes with less than 4 coloured particles (e.g. DIS or DY)

$$T_1 + T_2 + T_3 = 0 \Rightarrow T_1 \cdot T_2 = \frac{1}{2}(T_3^2 - T_1^2 - T_2^2)$$

leading to an unimportant overall phase

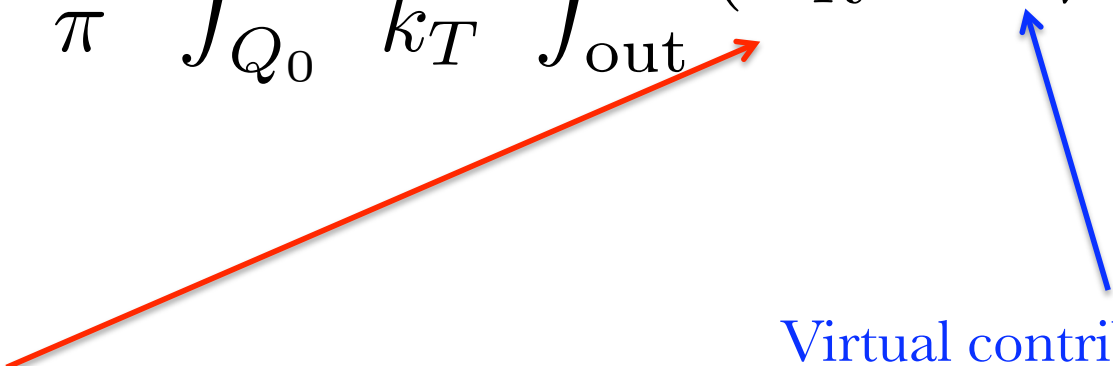
- Coulomb gluon contributions are *not* implemented in parton showers

Non-global effects

- However this naive approach completely ignores a whole tower of LL
- Virtual contributions are not the whole story because real emissions out of the gap are forbidden to remit back into the gap
- The full LL result is obtained by dressing the $2 \rightarrow n$ (i.e. $n-2$ out of gap gluons) scattering with virtual gluons (and not just $2 \rightarrow 2$)
- Resummation can be done (so far) only in the large N_c limit

One gluon outside the gap

- As a first step we compute the tower of logs coming from only one out-of-gap gluon:

$$\sigma^{(1)} = -\frac{2\alpha_s}{\pi} \int_{Q_0}^Q \frac{dk_T}{k_T} \int_{\text{out}} (\Omega_R + \Omega_V)$$


Real contribution:

- real emission vertex D^μ
- 5 - parton anomalous dimension Λ

Sjödahl
arXiv:0807.0555 [hep-ph]

Virtual contribution:

- virtual eikonal emission γ
- 4-parton anomalous dimension Γ

A big surprise

Conventional wisdom (“plus prescription” of DGLAP)

when the out-of-gap gluon becomes collinear with one of the external partons the real and virtual contributions should cancel

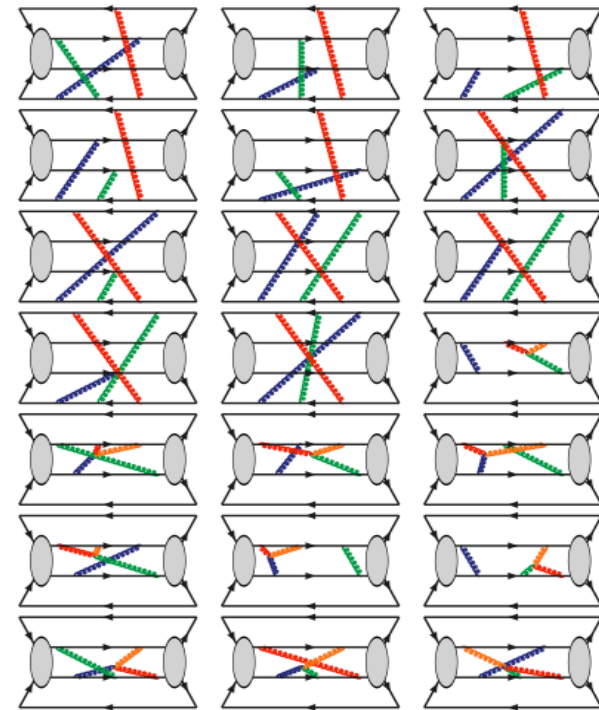
- It works when the out-of-gap gluon is collinear to one of the outgoing partons ✓
- But it fails for initial state collinear emission ✗
- Cancellation *does* occur for up to 3rd order relative to the Born, but fails at 4th order
- The problem is entirely due to the emission of Coulomb gluons
- As result we are left with super-leading logarithms (SLL):

$$\sigma^{(1)} \sim -\alpha_s^4 L^5 \pi^2 + \dots$$

Forshaw Kyrieleis Seymour
hep-ph/0604094

Fixed order calculation

- Gluons are added in all possible ways to trace diagrams and colour factors calculated using COLOUR
- Diagrams are then cut in all ways consistent with strong ordering
- At fourth order there are 10,529 diagrams and 1,746,272 after cutting.
- SLL terms are confirmed at fourth order and **computed for the first time at 5th order**

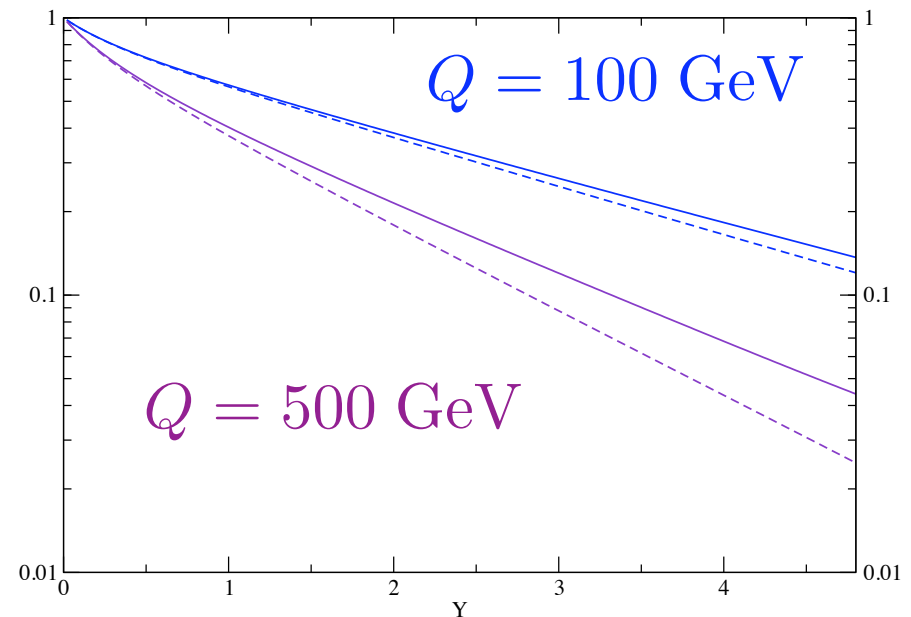
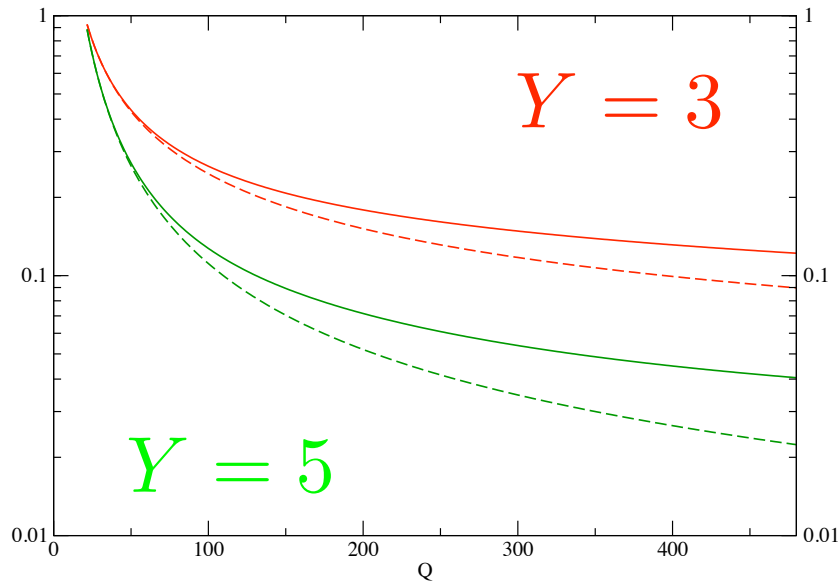


Keates and Seymour
arXiv:0902.0477 [hep-ph]

LHC results: no gluon outside the gap

$$f_{\text{gap}} = \sigma^{(0)} / \sigma^{\text{born}}$$

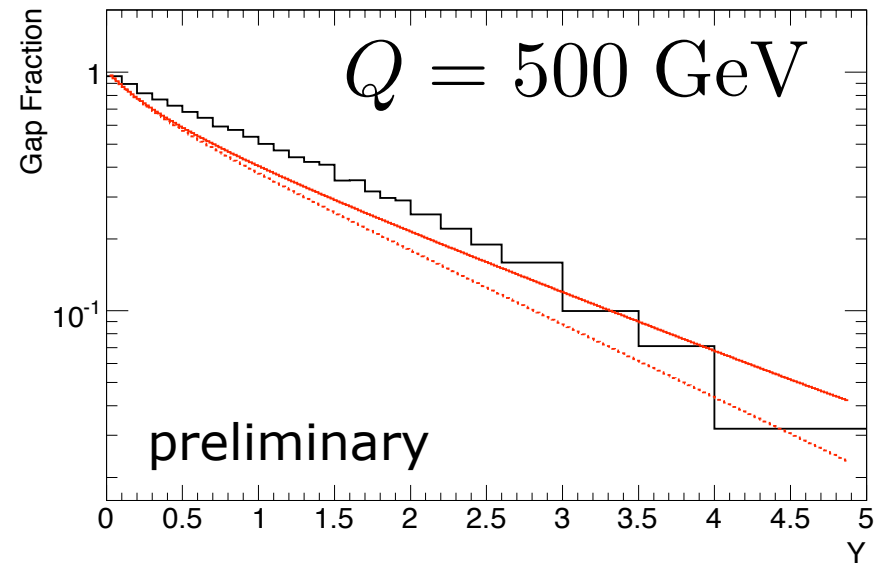
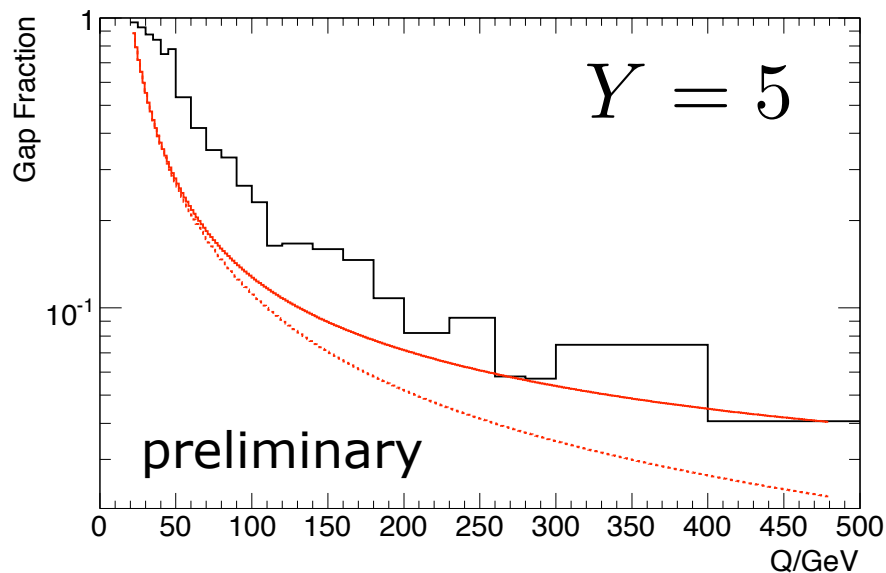
$$\begin{aligned}\sqrt{s} &= 14 \text{ TeV} \\ Q_0 &= 20 \text{ GeV} \\ R &= 0.4 \\ \eta_{\text{cut}} &= 4.5\end{aligned}$$



- solid lines: full resummation
- dashed lines: ignoring $i\pi$'s

Comparison to HERWIG++

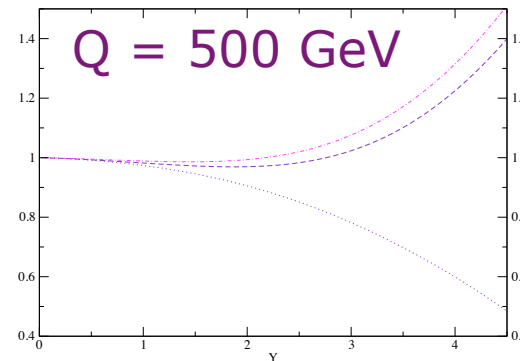
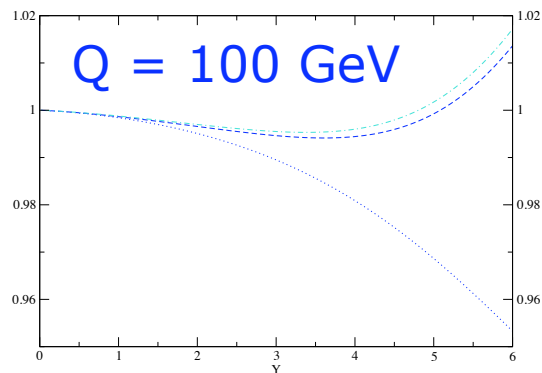
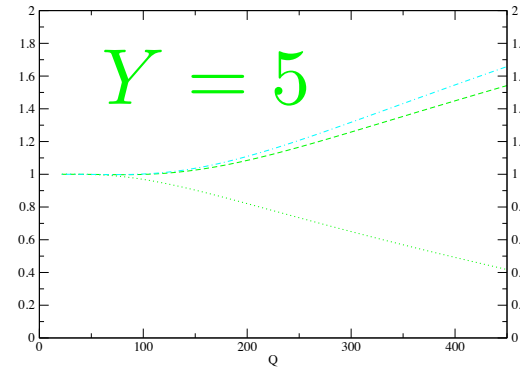
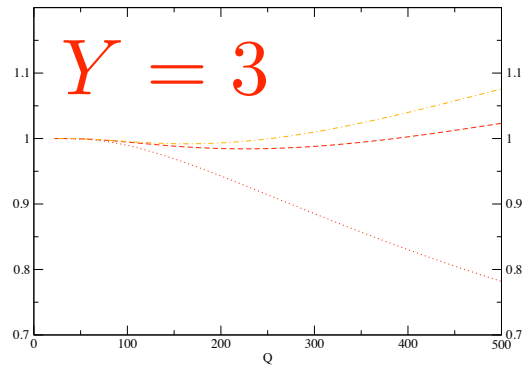
- We compare our results to HERWIG++
- LO scattering + parton shower (no hadronisation)
- Jet algorithm SIScone



- Different results at low Q and Y : energy-momentum conservation
- We need matching to NLO
- Other differences: large N_c limit and non-global effects

Phenomenology of SLL (I)

$$(\sigma^{(0)} + \sigma^{(1)} + \sigma^{(2)})/\sigma^{(0)}$$

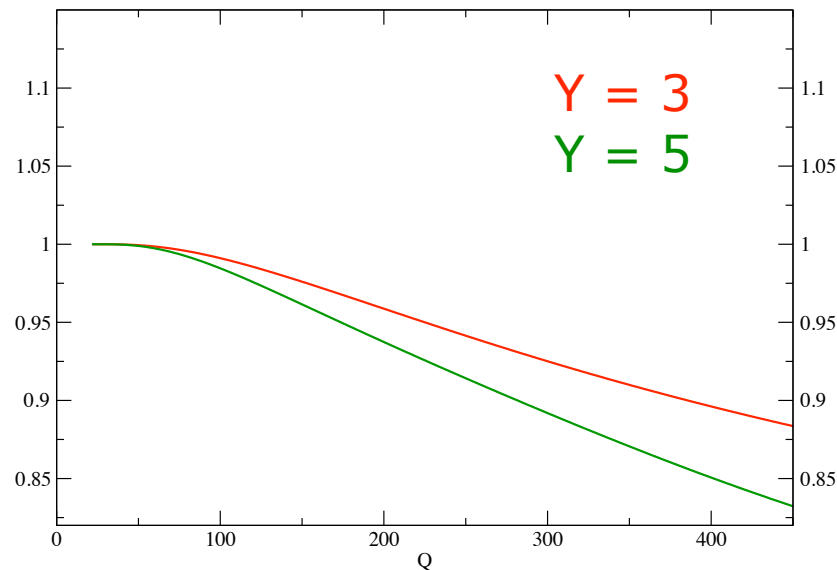


- dotted, one gluon, α_s^4
- dashed: one gluon, up to α_s^5
- dash-dotted: one+two gluons, up to α_s^5

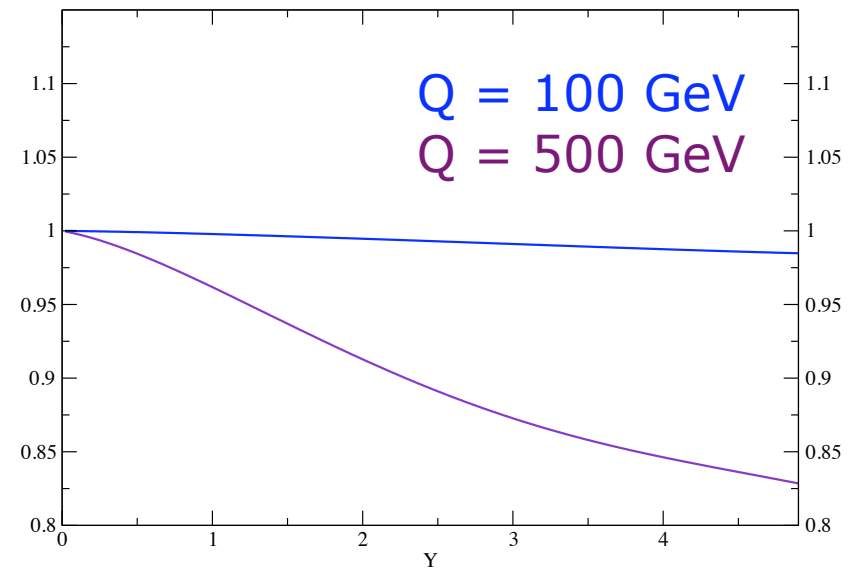
instability:
need of resummation

Phenomenology of SLL (II)

Resummed results (one out-of-gap gluon)



- $Y = 3$, $\sim 5\%$
- $Y = 5$, $\sim 10-15\%$



- $Q = 100$ GeV, $\sim 2\%$
- $Q = 500$, $\sim 10-15\%$

- SLL could have an effect as big as 10-15 % in quite extreme dijet configurations
- There are no SLL effect on Higgs+ jj, unless $Q_0 < 10$ GeV

Conclusions

- Early data: there is plenty of interesting QCD physics in gaps between jets
- More data: Higgs coupling to weak bosons
- Coulomb gluons play an important role
- Dijet cross-section could be sensitive to SLL at large Y and L (e.g. 300 GeV and $Y = 5$, $\sim 15\%$)

Outlook (phenomenology)

- Compute the best theory prediction for gaps between jets at the LHC:
 - matching with NLO
 - complete one gluon outside the gap
 - non-global (large N_c)
 - jet algorithm dependence
 - BFKL resummation

Outlook (theory)

- There is an interesting link between non-global logs and BK equation

Banfi, Marchesini and Smye
hep-ph/0206076

Avsar, Hatta and Matsuo
arXiv:0903.4285 [hep-ph]

- Understanding the origin of super-leading logs
 - k_t ordering ?
 - interaction with the remnants ?

on-going projects in Manchester