

# Measurement of charged particle spectra in pp collisions at CMS

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*DIS09: XVII International Workshop  
on Deep-Inelastic Scattering and Related Subjects*

# Outline

- Motivation of interest
- State of knowledge
- The CMS Silicon Tracker: layout and performances
- Minimum Bias trigger performances
- CMS proposed analysis
- Conclusions

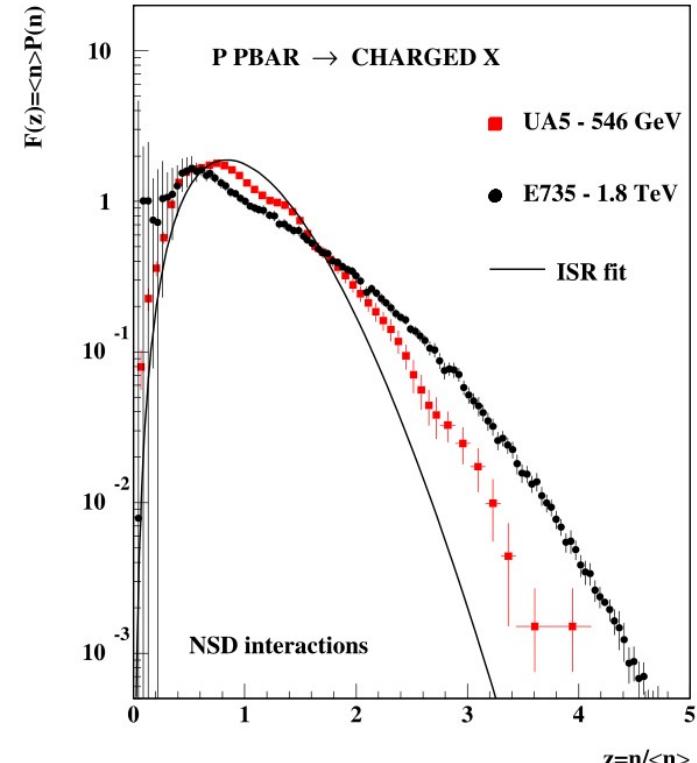
# Motivation

Minimum Bias (mostly non-single diffractive inelastic interactions)  
~65% of the total LHC scattering process:  $\sigma_{NSD} = \sigma_{tot} - \sigma_{elas} - \sigma_{SD}$

- analyse global properties of hadron production:  
higher energies → new opportunity for comprehension, sensitivity to fragmentation models
- predict background levels associated to many physics processes:  
at design luminosity of  $\mathcal{L}=1.0 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$  and for the final expected beam bunch pattern of  $2808 \times 2808$  there will be on average 20 collisions per recorded event: 1 signal + ~19 min-bias pile-up
- understand the complex nature of the radiation environment in which LHC's detector systems will operate (particular attention for Si tracker)

# History

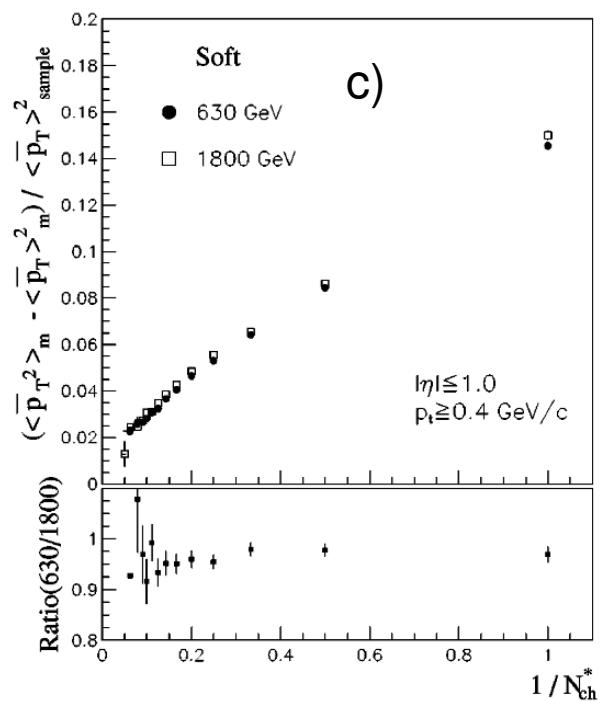
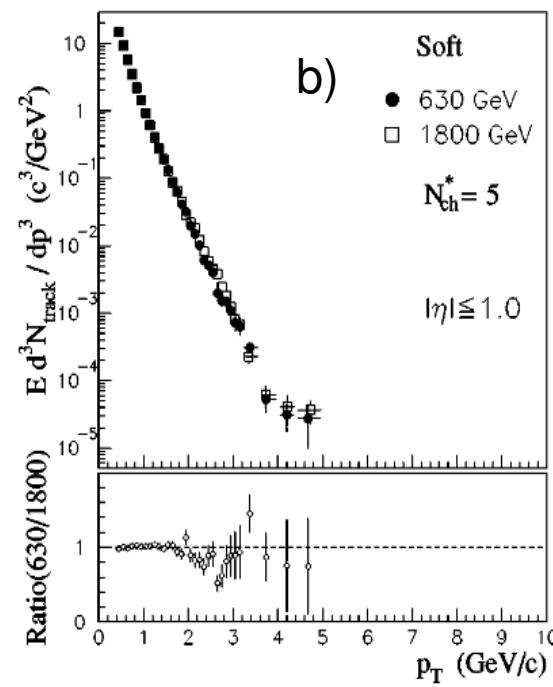
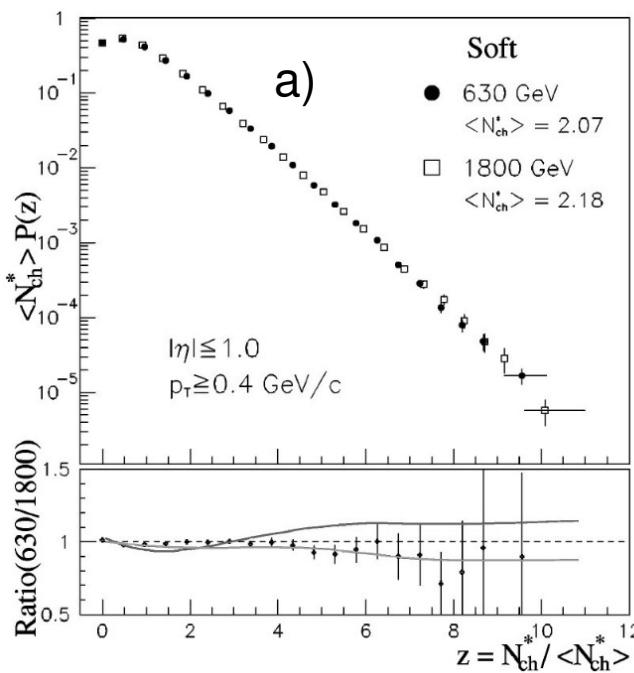
- asymptotic scaling of multiplicity distribution in hadron-hadron collisions (KNO scaling) predicted in 1972;
- early confirmation at 10-30-60 GeV;
- first observations of violation at higher energies from the UA5 collaboration (c.m. energies of 200-546-900 GeV) → notion of multiple parton interaction [T. Sjöstrand prd 36 (1987) 2019] ;
- CDF at Tevatron:
  - minimum bias trigger events at 1800 and 630 GeV
  - novel approach dividing dataset into “soft” and “hard” substitution samples (based on presence of calorimetric activity in the transverse region)



# Tevatron studies

Evidence of energy invariance in “soft” sample:

- a) KNO scaling
- b) energy invariant  $p_T$  distributions at fixed multiplicities
  - in purely “soft” interactions the number of produced charged particles is the only global event variable changing with  $\sqrt{s}$
  - ← consequence/confirmation: mean  $p_T(N)$  invariance
- c)  $\langle p_T \rangle_{\text{event}}$  dispersion scaling and compatibility with no dynamical fluctuations

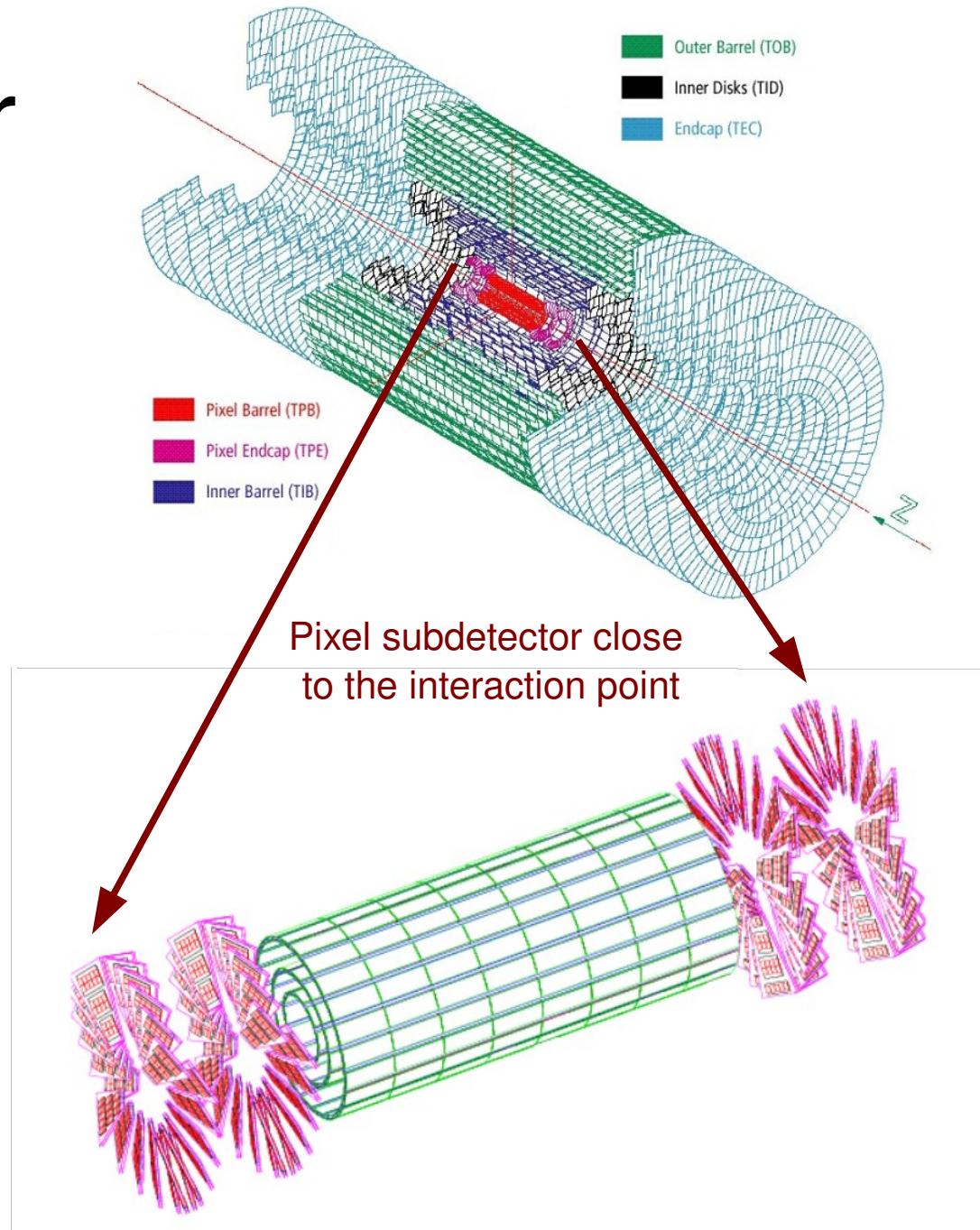


# The Pixel Tracker

- ~1 m total length
- 3 barrel layers at radii of 4 , 7 and 11 cm
- 2 disks on each side at 34 and 46 cm
- about 60 million of pixels with shape of  $100 \times 150 \mu\text{m}^2$
- design resolution of about  $10 \mu\text{m}$  for the  $r\text{-}\phi$  measurement and about  $20 \mu\text{m}$  for the  $z$  measurement

High resolution three dimensional space points:

- primary vertex reconstruction  
→ track seeding  
→ fast HLT track reconstruction



# Performances

## TRACKER:

November cosmic ray data acquisition with magnetic field (CRAFT)

### • Operational Fraction:

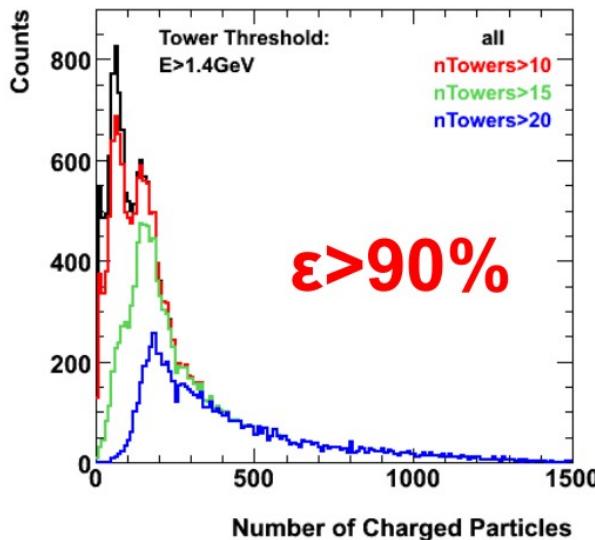
- Barrel pixels: 99.1%
- Forward pixels: 94.0%  
(readout chips without bias voltage,  
reparation in course)

### • Cosmic Track Finding Efficiency:

- Tag&Probe technique:
  - tag = cosmic tracks in muon systems
  - probe = tracker tracks CKF track algorithm
- Combinatorial Track Finder, standard for collision  
(Kalman filter)

Efficiency (%)	Data	MC
CKF	$99.5 \pm 0.1$	$99.9 \pm 0.1$

## MB trigger: hadronic forward based



### • Backgrounds:

- **Beam-gas** collisions: at start up ( $43 \times 43$  bunch pattern) expected rate of 5 mHz; for the full luminosity ( $2808 \times 2808$  bunch pattern) expected rate of 35 kHz. Comparing to the beam bunch crossing-clock collision rate (32 MHz): 1 collision per 1000 would be a beam gas interaction;
- **Beam-halo** collisions: no events triggered out of  $10^4$  MC events;
- **Cosmics**: measured negligible during cosmic data acquisition with magnetic field;

Charged particle spectra at CMS

# CMS proposed analysis

- Charged spectra with reconstructed particle and low- $p_T$  adapted algorithm
- Pseudorapidity distribution with hit-counting
- {Tracklets: work in progress}

# Low $p_T$ tracks reconstruction (I)

## Pixel subdetector permits:

- specific low  $p_T$  algorithm
- same analysis for pp and PbPb interaction (thanks to low occupancy)

## Basic logic (both standard and low pT tracks):

make hit pairs then predict range of possible coordinates of the third hit

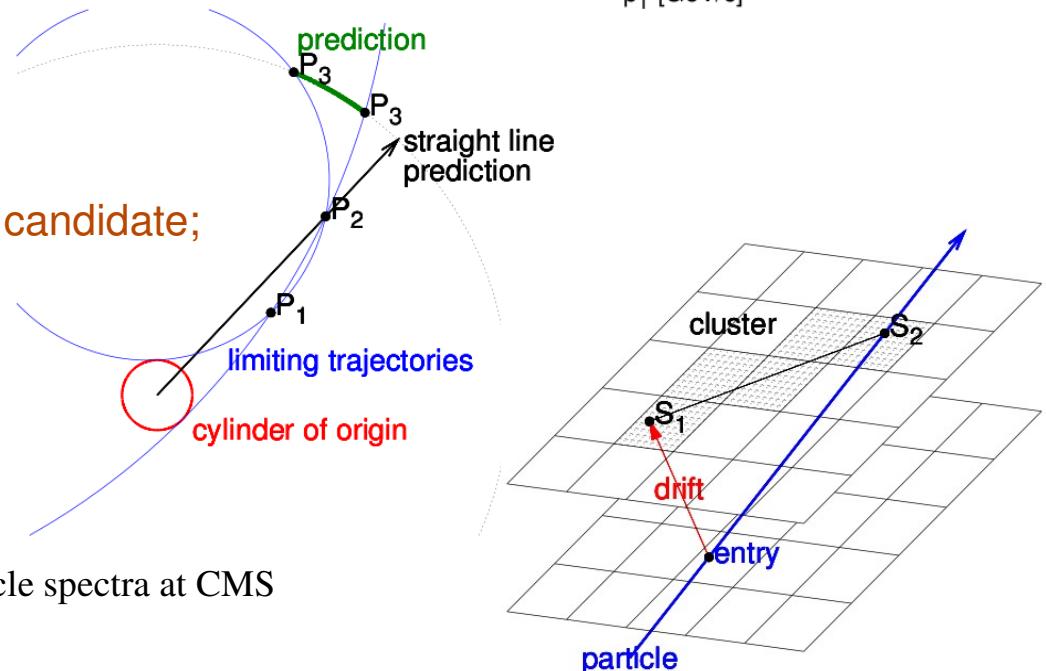
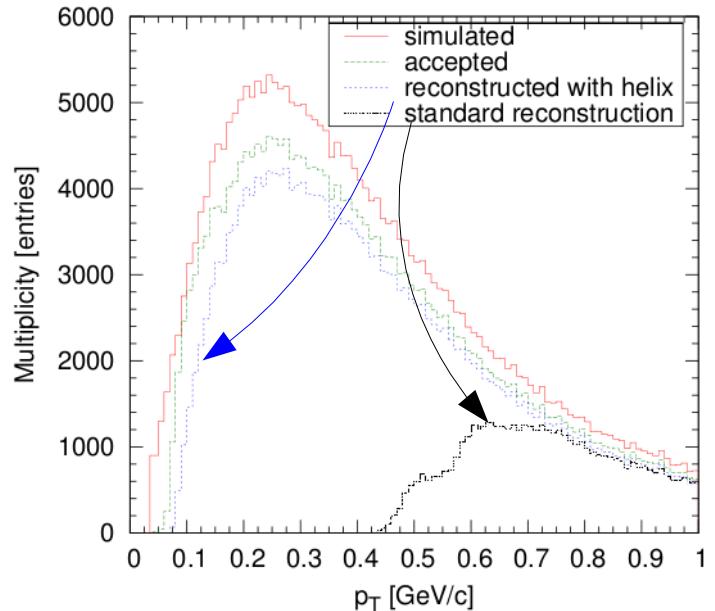
## Criteria specific for low pT:

tracks coming from a cylinder of origin; minimum  $p_T$  requirement ( $0.075\text{GeV}/c$ );

the track must be able to reach the third hit layer candidate; helix propagation;

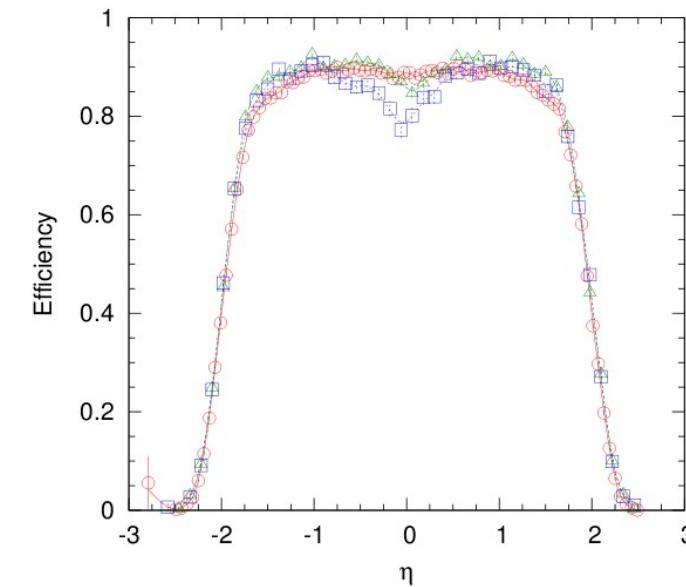
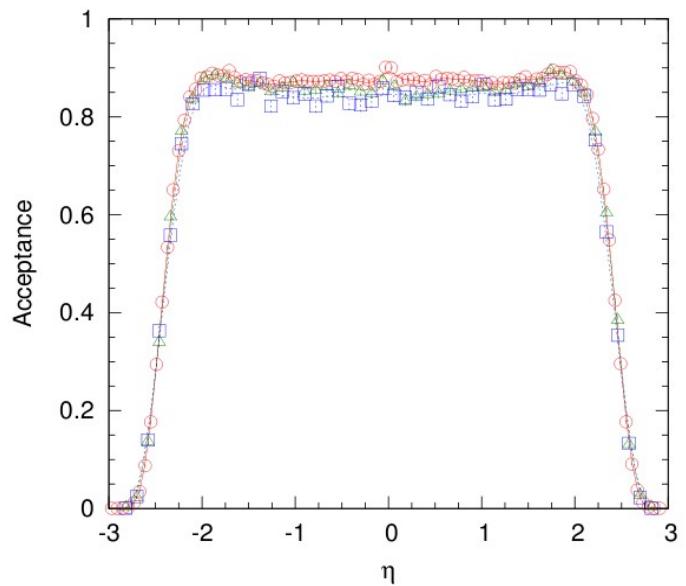
## Triplet cleaning:

false rejection checking compatibility between predicted incidence angle and cluster shape

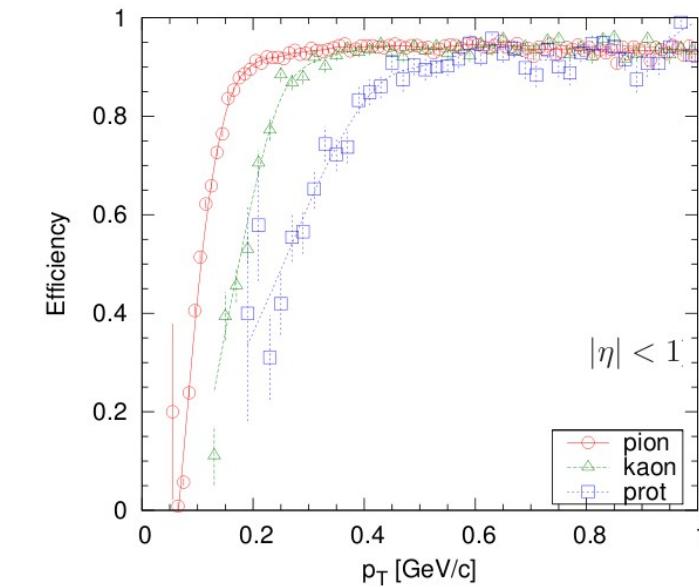
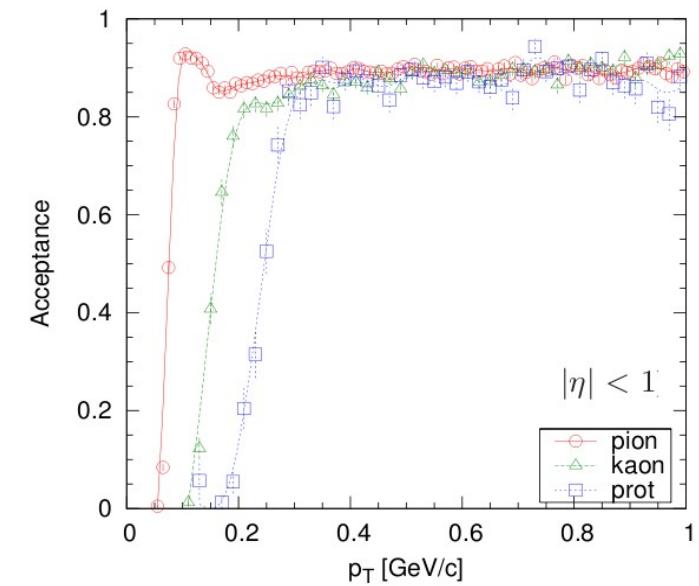


# Low $p_T$ tracks reconstruction (II)

Essentially flat acceptance above  $p_T$  of 0.1, 0.2 and 0.3 GeV/c for pions, kaons and protons



Flat efficiency in the range  $|\eta| < 1.5$ , except for protons

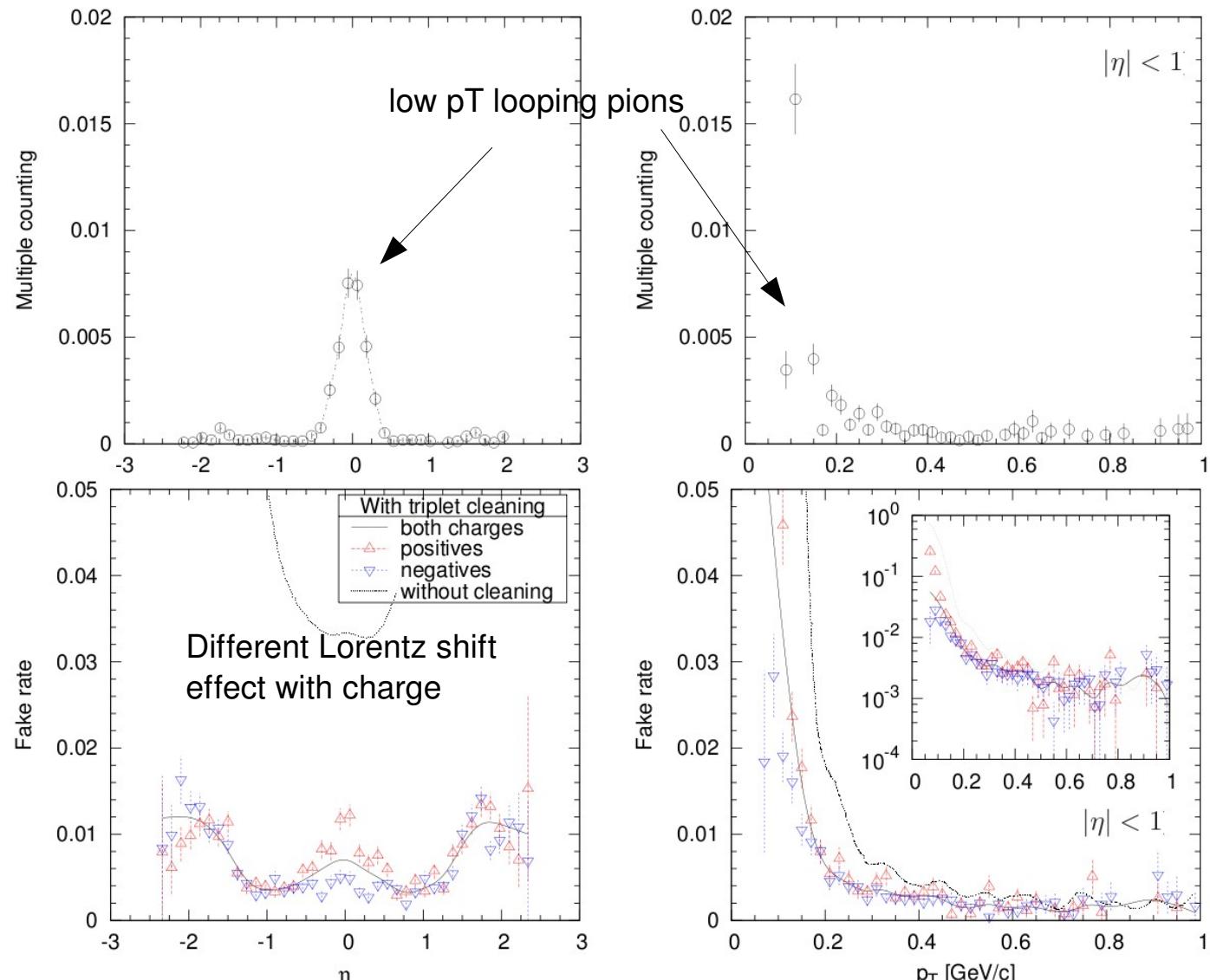


# Low $p_T$ tracks reconstruction (III)

Multiple counting  
 (essentially looper) below  
 1% around  $\eta=0$ , quickly  
 to zero for  $|\eta|>0.5$

Clean up imposing  
 vertexing +  
 cluster-shape/incident-angle  
 relation:

low fake rate at 0.5%(2%)  
 for  $\eta=0$  ( $|\eta|=2$ ), below 1%  
 for  $p_T>0.16\text{GeV}/c$



# Charged hadron spectra with low $p_T$ tracks

Binning measured tracks in  $\Delta N_{measured}(\eta, p_T)$  and correcting for all the considered effects:

$$\Delta N_{corrected} = \frac{(1 - fakeRate)(1 - feedDown)}{geomAccep \cdot algoEffic \cdot (1 - multiCount)} \cdot \Delta N_{measured}$$

→ Triple differential invariant yield:  $E \frac{d^3 N}{dp^3} = \frac{d^3 N}{d\phi dy dp_T} = \frac{1}{2\pi p_T} \frac{d^2 N}{dy dp_T} = \frac{1}{2\pi p_T} \frac{E}{p} \frac{d^2 N}{d\eta dp_T}$

Fitted by a reparametrization  
of a Tsallis function:

$$E \frac{d^3 N}{dp^3} = \frac{dN}{dy} \frac{(n-1)(n-2)}{2\pi nT [nT + (n-2)m]} \left[ 1 + \frac{E_T(p_T)}{nT} \right]^{-n}$$

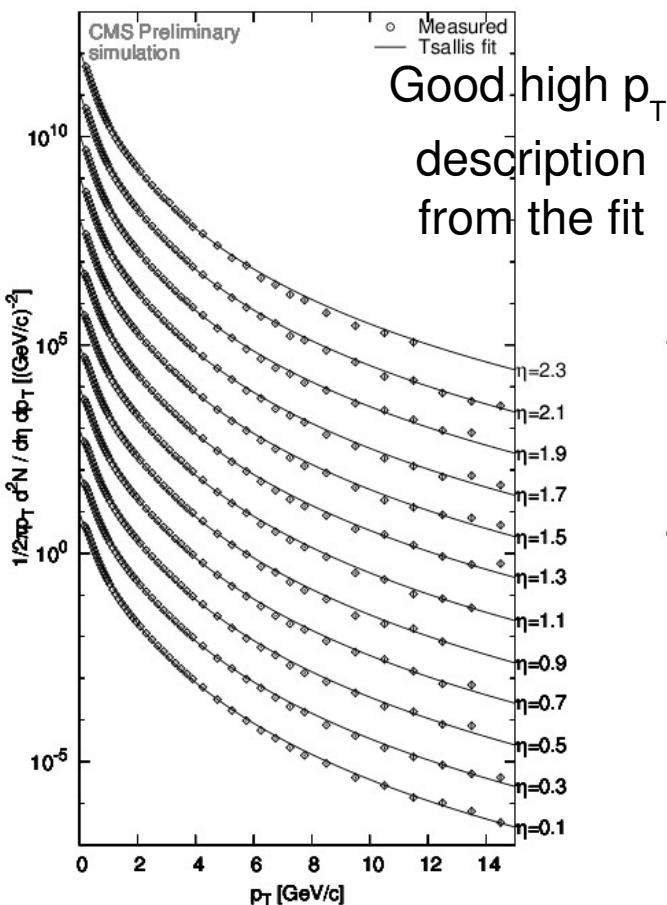
Combining both {

- thermal low  $p_T$ :  $E \frac{d^3 N}{dp^3} = A \exp \left[ \frac{-E(p_T)}{T} \right]$
- high  $p_T$  empirical power-law tail:  $E \frac{d^3 N}{dp^3} = \frac{A}{(1 + p_T/p_0)^n}$

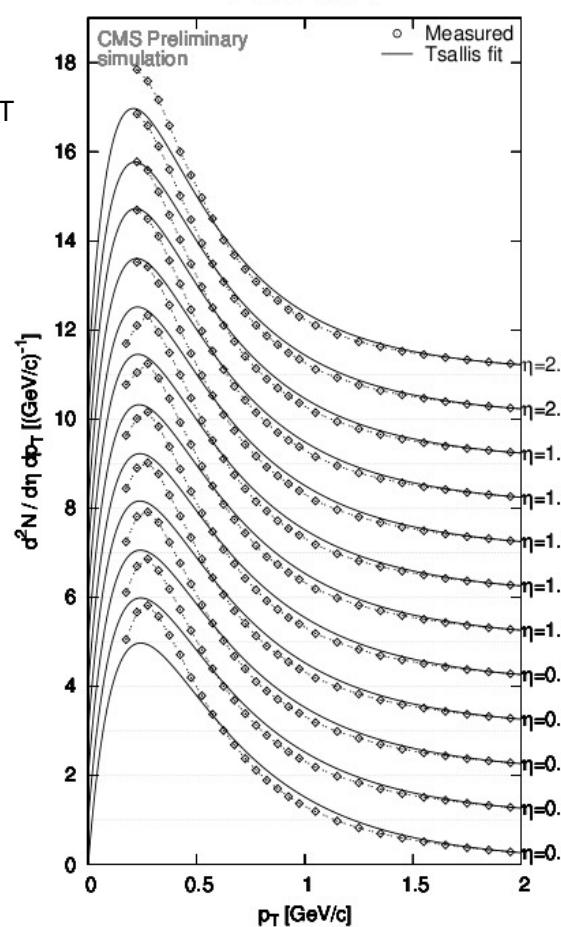
obtaining the  $p_T$  integrated  $dN/dy$  distribution from fit !

# Charged hadron spectra with low $p_T$ tracks

Charged hadrons



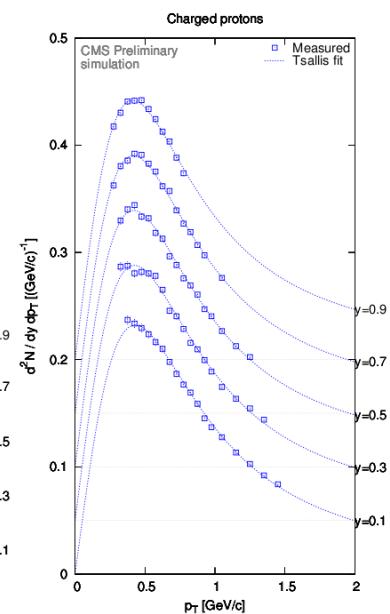
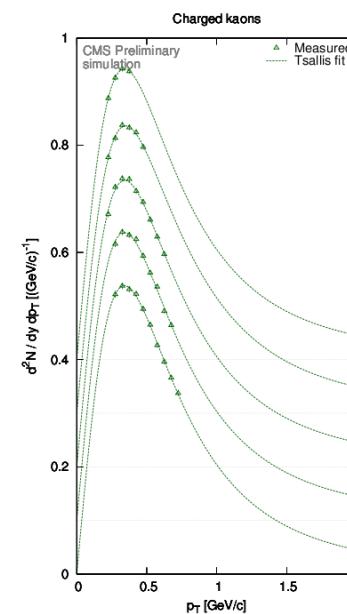
Charged hadrons



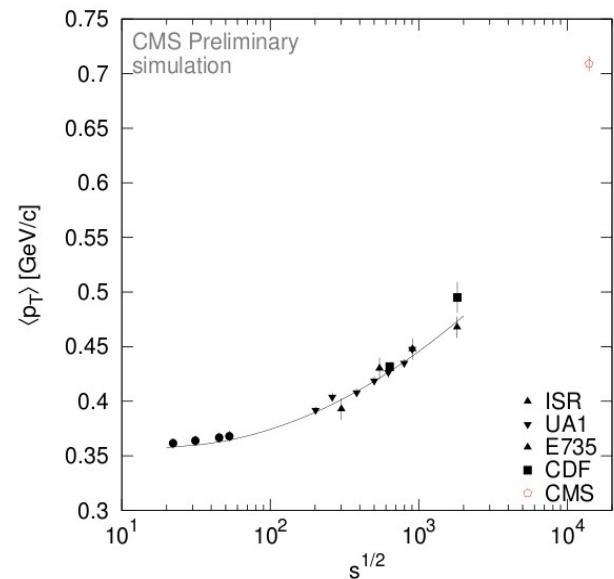
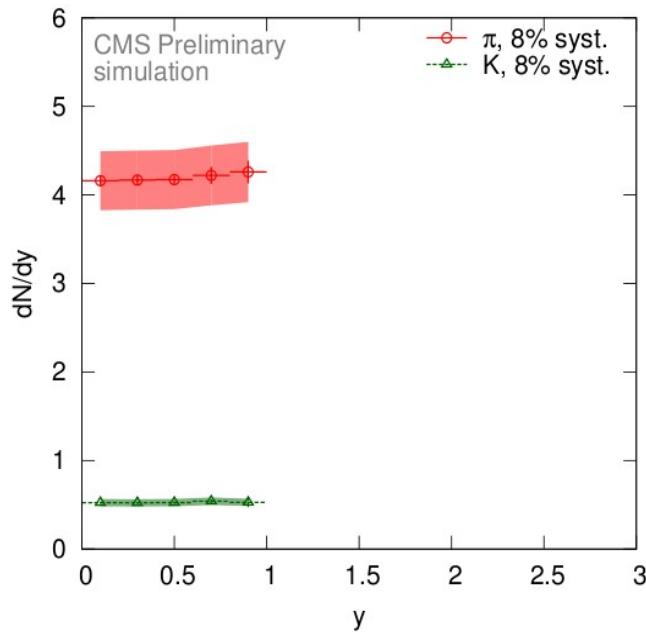
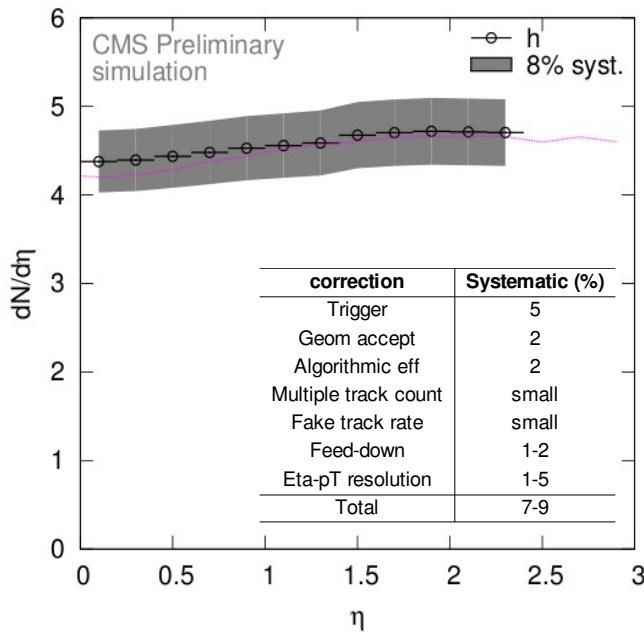
Bad low  $p_T$  description  
for unidentified hadrons

- sum the measured differential yields for  $0.2 < p_T < 2 \text{ GeV}/c$
- thermal fit extrapolation down to  $0 \text{ GeV}/c$

Better situation for  
identified hadrons:



# Charged hadron spectra with low $p_T$ tracks



- Differential cross-section:

$$d\sigma_{inelastic} = \frac{dN}{\Delta N_{events}} \cdot \sigma_{inelastic} = \frac{dN}{\Delta N_{events}} \cdot \frac{\Delta N_{events}}{Lumi \Delta t} = \frac{dN}{Lumi \Delta t}$$

(measuring time  $\Delta t$  and Luminosity  $Lumi$  known and independently measured)

# Hit counting method

- Early startup:

low luminosity scenario → no pile-up

(needed only  $10^4$  events here for the analysis)

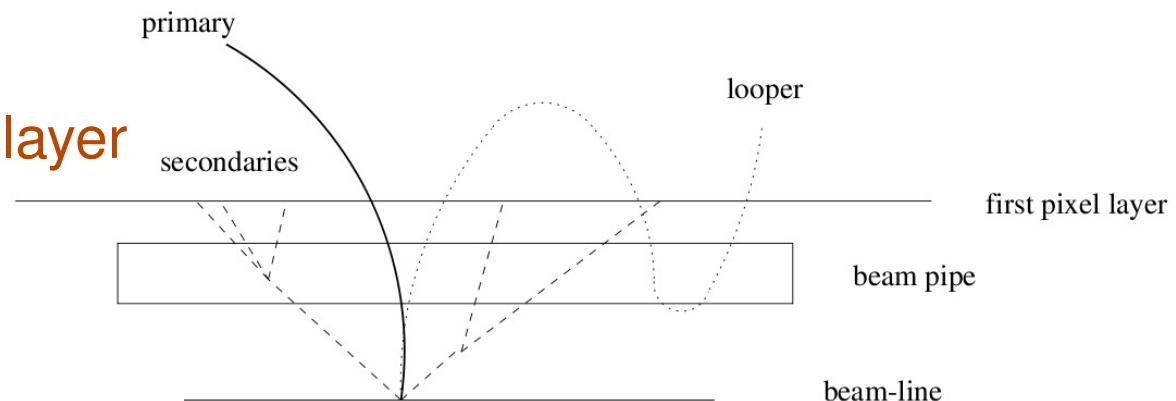
- Method:

(as PHOBOS@RHIC)

counting hits on one pixel barrel layer

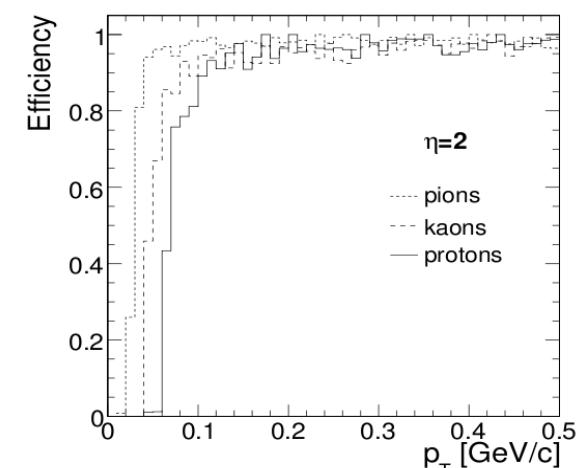
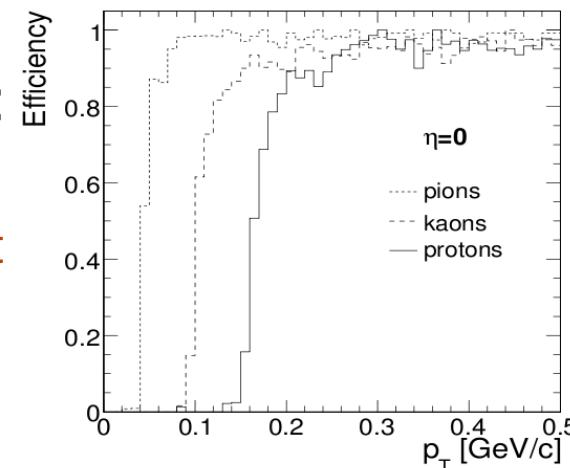
→ correcting for:

- Looper particle
- Decays from primary
- Particles from beam-pipe interaction



- Small minimum  $p_T$  requirement:

~30 MeV/c to leave hit in the first layer  
(to be compared with 75 MeV/c request from full tracking)



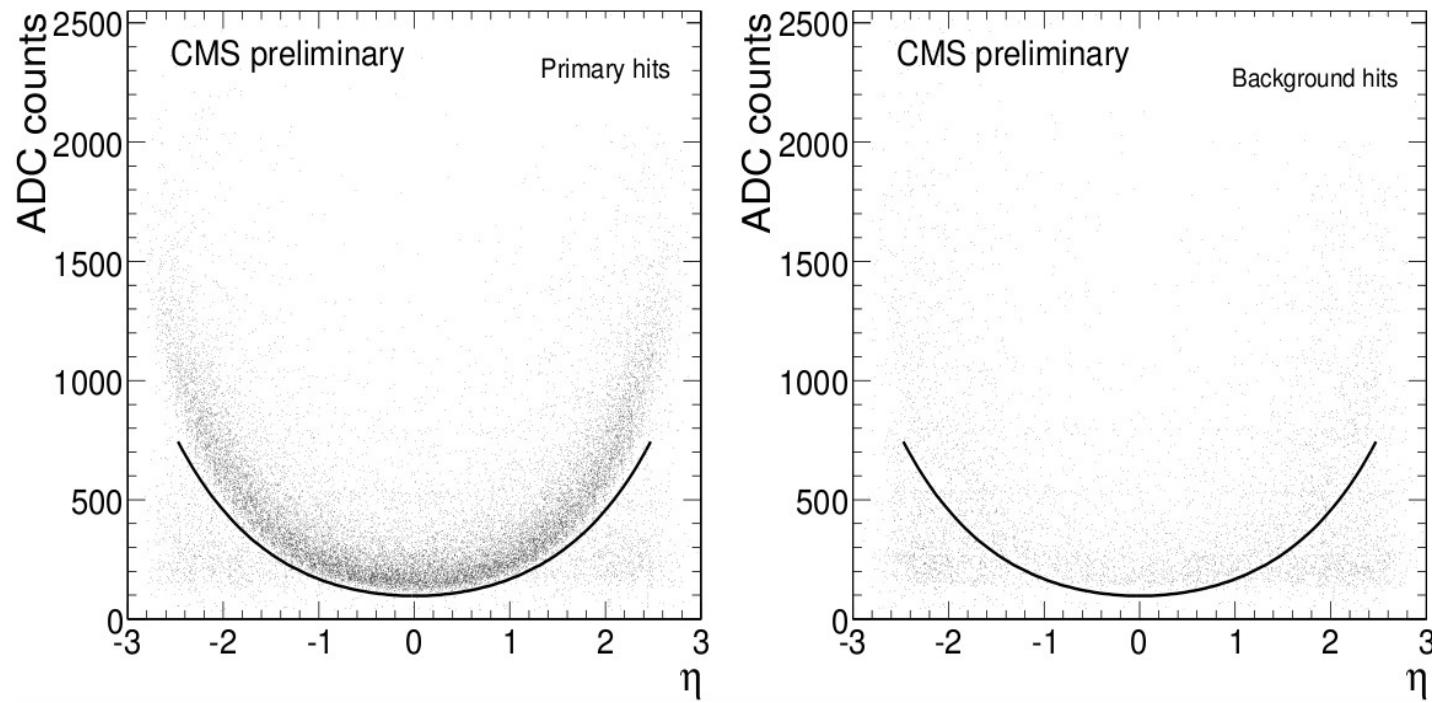
Charged particle spectra at CMS

# Hit counting method

Hit selection criteria: release of energy  $dE/dx$  in the Silicon

**Primary:** average energy deposit proportional to  $\cosh(\eta)$  and Landau-distributed around the mean

**Background:** no strict relation crossing-angle/pseudorapidity



# Hit counting method

Ingredients<sup>\*</sup> for calculating the pseudorapidity distribution of charge hadrons:

$$\frac{dN}{d\eta} = \frac{\sum_M E_t^{real} \cdot \frac{1}{\epsilon(M)} \cdot \frac{dN}{d\eta}(M)}{\sum_M E_t^{real} \cdot \frac{1}{\epsilon(M)}}$$

(\* more in backup)

## Systematics:

### • MODEL DEPENDENCES from corrections

the analysis uses 2 MC dataset: one giving the correction, one acting as the “real” measure  
→ need to evaluate the response in case of momentum and multiplicity distribution extreme differences between MC and “real” data;

simulated scenarios:

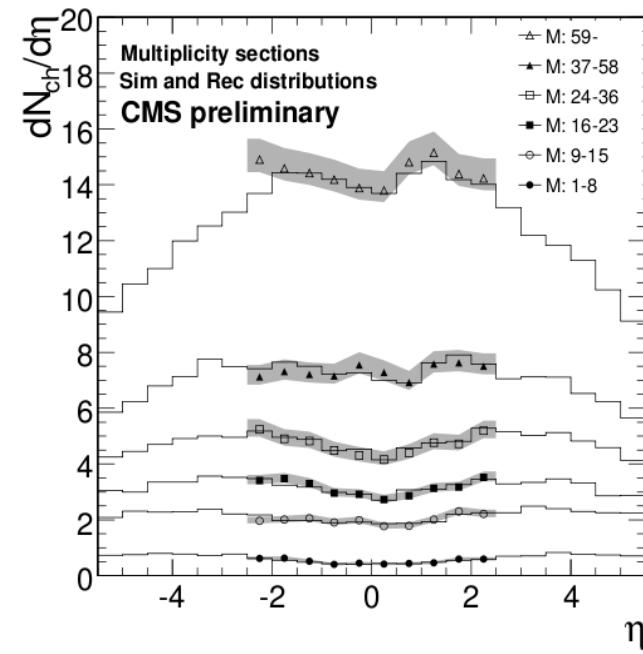
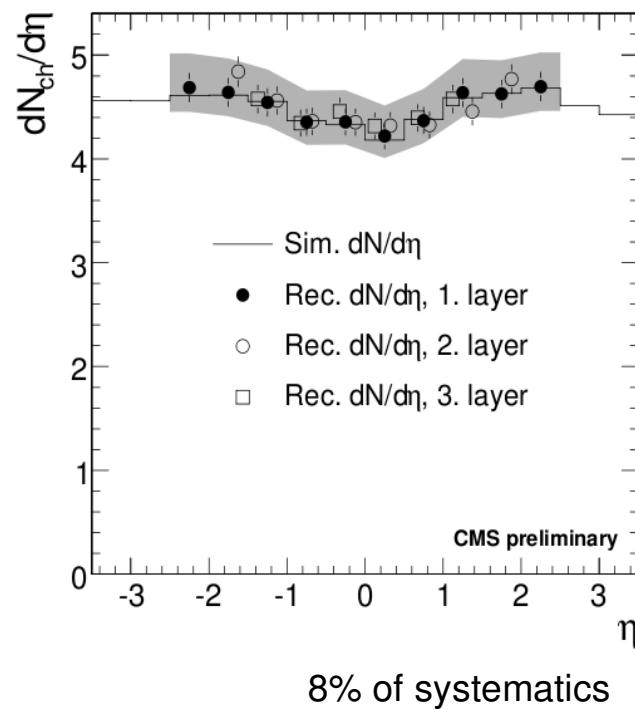
- a) multiplicity-distribution\*0.5 (random removal), multiplicity-distribution\*2 (merging dataset), unchanged pT-distribution;
- b) p-distribution\*0.5, p-distribution\*1.5, unchanged multiplicity

### • VERTEXING

estimated effect of uncertainty on actual vertexing efficiency ( $\text{eff}_V$ ):

analysis without vertex reconstruction → calculated the vertex bias →  $\text{sys}_V = (1 - \text{eff}_V) * \text{vertex-bias}$

# Hit counting method



# Conclusion

- CMS is ready for first collisions analysis
- Results feed through to detector + physics community
- Interesting results available after 1 month of collisions (1<sup>st</sup> analysis) or even few days (2<sup>nd</sup> analysis)

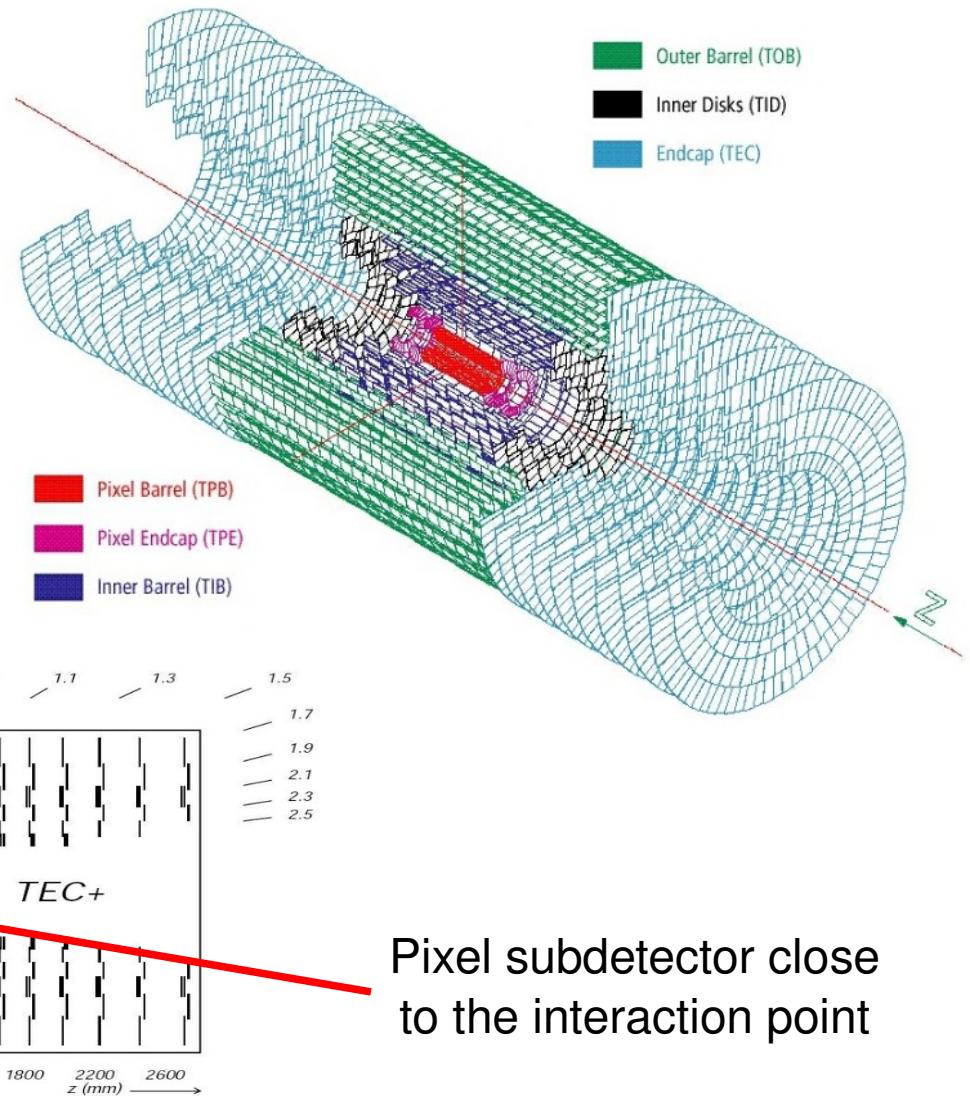
# Bibliography

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- *Measurement of charged hadron spectra in proton-proton collisions at sqrt(s)=14TeV, F. Sikler, K.Krajczar, CMS AN-2007/021*
- *Pseudorapidity distributions of charged hadrons in minimum bias p-p collisions at sqrt(s)=14TeV, K.Krajczar, G. Veres, CMS AN 2008/018*

# Start of Back up slides

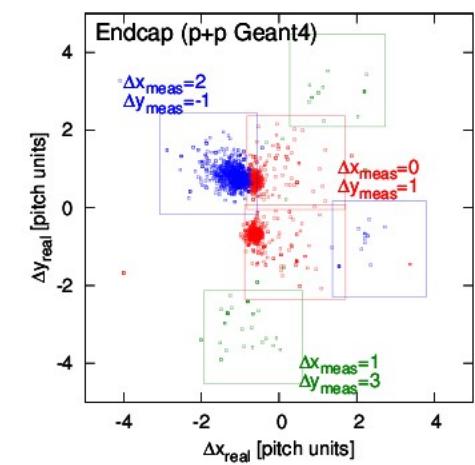
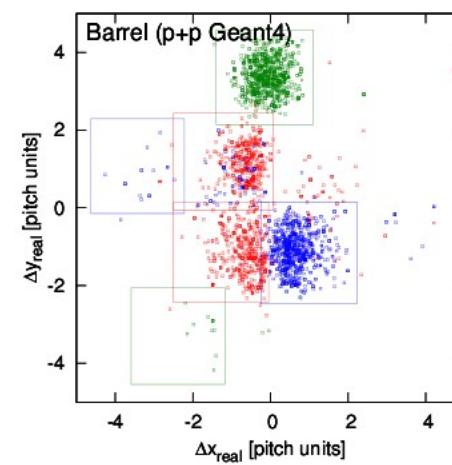
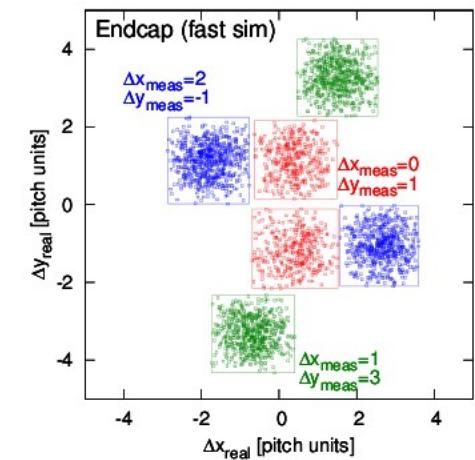
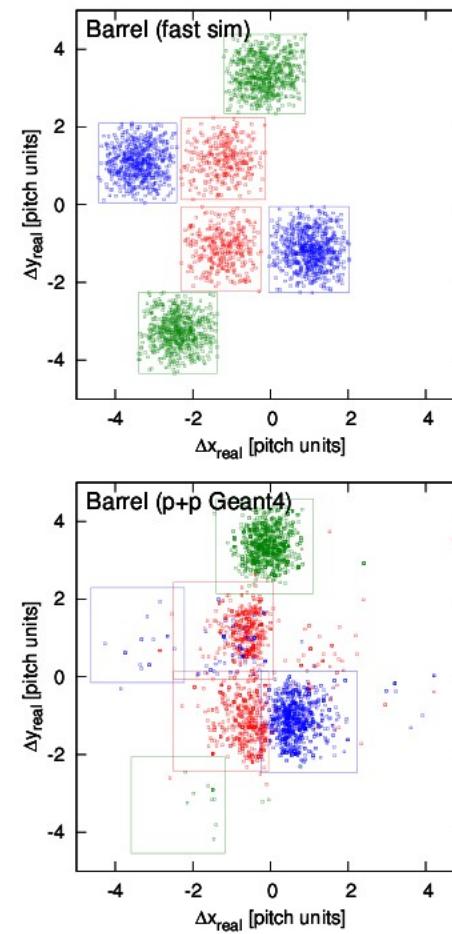
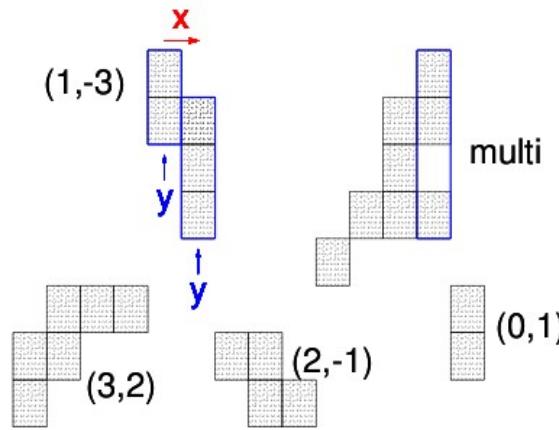
# The CMS Silicon Tracker

- Total area:  $\approx 1 \text{ m}^2$  (Pixel) +  $200 \text{ m}^2$  (Strip)
- Coverage up to  $|\eta| < 2.4$
- Active elements:  
66 million pixels + 9.6 million silicon strips

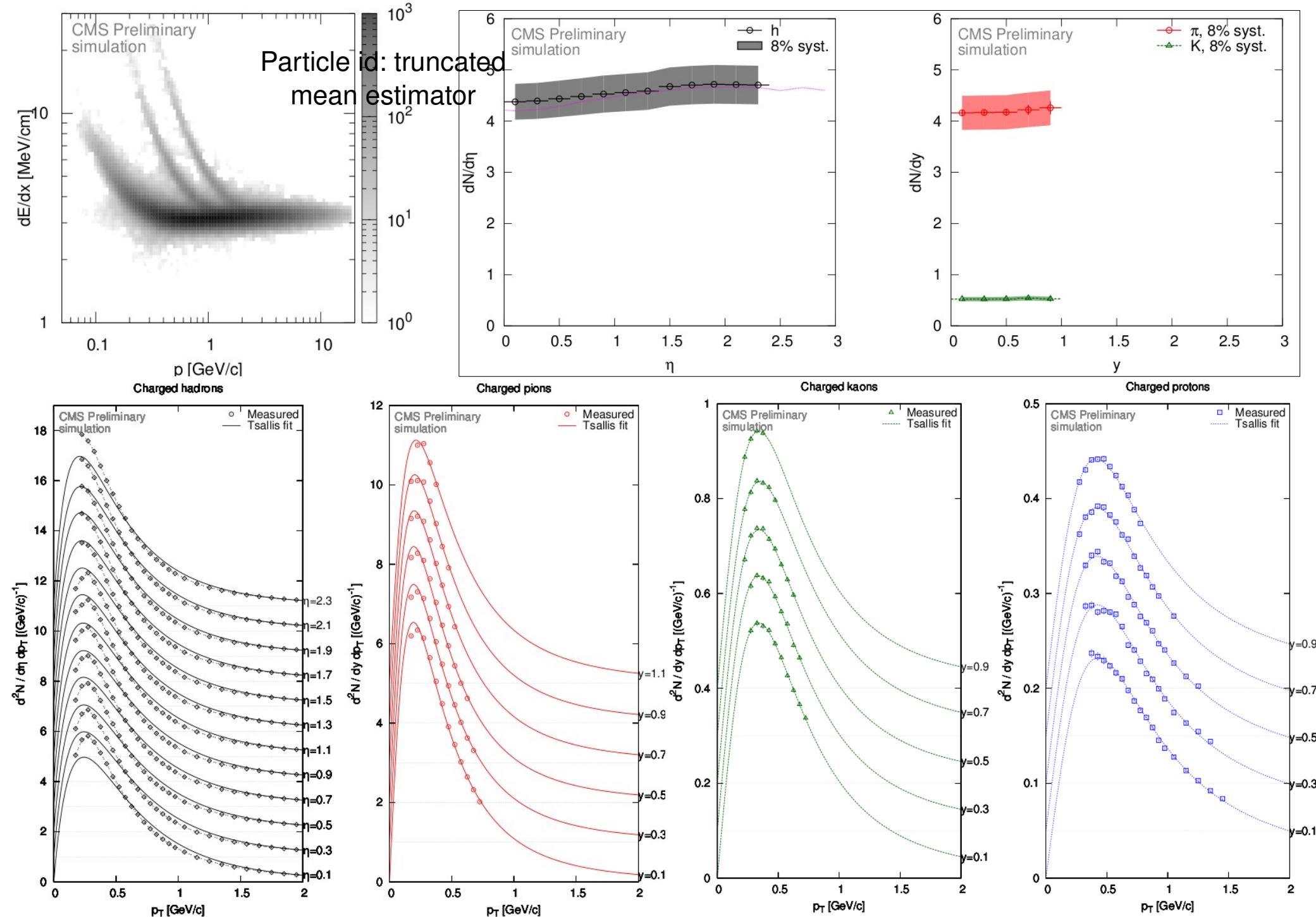


# Track cleaning from shape

Low  $p_T$  track reconstruction



# Charged hadron spectra with low pT track



# Hit counting method

## Corrections

### •Event selection:

ratio between MC events passing selection cut and total events with  $M$  pixel hits:

$$\epsilon(M) = E_t^{MC}(M)/E^{MC}(M)$$

dependent on the number of hits on the layer (“multiplicity”  $M$ ).

→ to correct the  $M$ -dependent measured multiplicity distribution of charged particle  $\frac{dN}{d\eta}(M)$

### •hit/tracks:

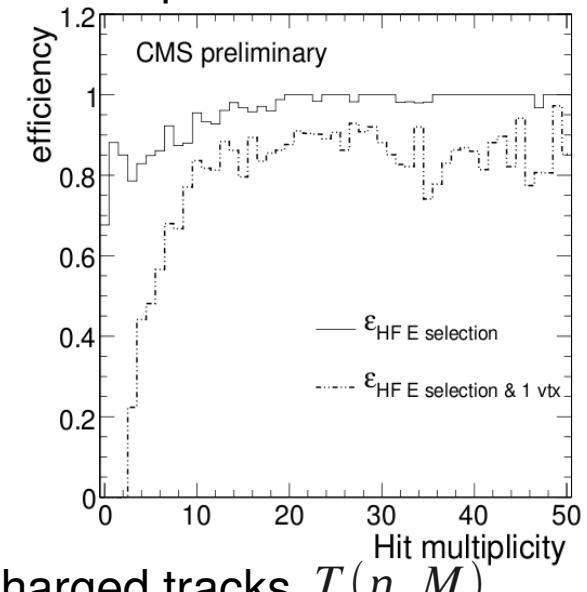
convert the measured number of hits  $H(\eta, M)$  into number of charged tracks  $T(\eta, M)$  through simulation prediction

$$\chi(\eta, M) = \frac{H^{MC}(\eta, M)}{T^{MC}(\eta, M)}$$

← small model-dependence  
(only detector geometry)

and finally measure the  $M$ -dependent

multiplicity distribution of charged hadrons:  $\frac{dN}{d\eta}(\eta, M) = \frac{1}{\chi(\eta, M)} \frac{H^{real}(\eta, M)}{E_t(M)}$



## Corrections

# Hit counting method

- **Looper:**

particle with  $p_T < 800 \text{ MeV}/c$  not reaching calorimeters apparatus;

both primary and secondary;

properly rejected for high eta ( $|\eta| > 2$ ) with the  $dE/dx$  hit selection

→ estimate the hits from looper passing  $dE/dx$  cut from high-eta rejected hits:

$$looperHit_{passing}^{real} = \left( \frac{looperHit_{passing}^{MC}}{totalHit_{notpassing}^{MC}} \right) \cdot totalHit_{notpassing}^{real}$$