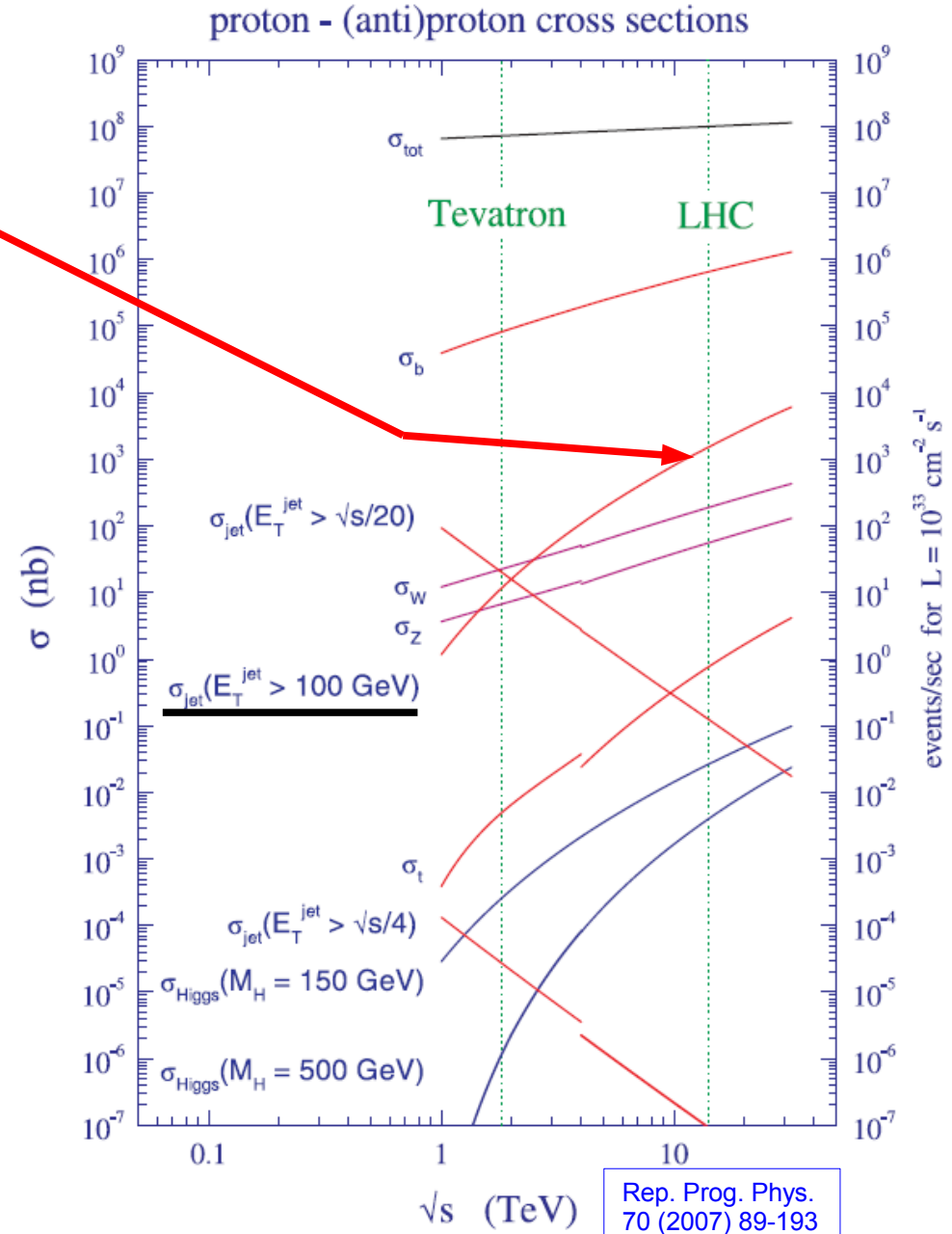
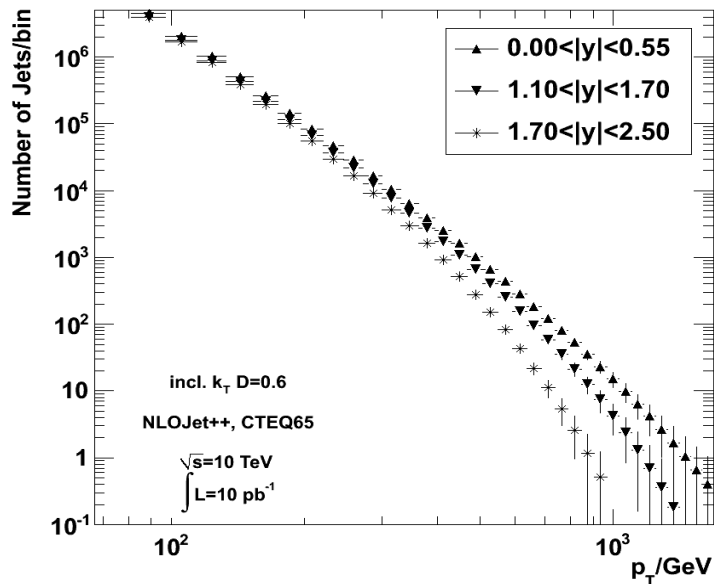


High p_T Jet Studies with CMS

On behalf of the CMS Collaboration
Andreas Oehler
University of Karlsruhe (KIT)

- **QCD Cross Section at the LHC**
- **Jet Measurement at CMS**
 - Jet Energy Calibration
 - Jet p_T Response
- **QCD & New Physics with High p_T Jets**
 - Inclusive Jets
 - Dijet Mass Cross Section
 - Dijet Angular Distributions
 - Jet Shapes
- **Outlook**

- QCD processes dominate LHC cross sections
- High p_T jet measurements will allow:
 - Detector commissioning
 - Sensitivity to new physics
 - Confrontation of pQCD at the TeV scale
 - Tests of current PDF extrapolations
 - Probe α_s
 - Understand multijet production (background to other searches)



- **Jets are sprays of particles**, collimated due to QCD interaction.
- **Jet algorithms** define an unambiguous way to **combine particles** to one fourvector:
 - Cone algorithms
 - Successive recombination algorithms
- **Input** to the algorithms are fourvectors from:
 - Calorimetric energy depositions
 - Tracks
 - Particle or energy flow reconstructed objects
 - Simulated or generated particles

Jet algorithms used in CMS

Iterative Cone, $R=0.5$ (in trigger)

fast, predictable reconstruction time

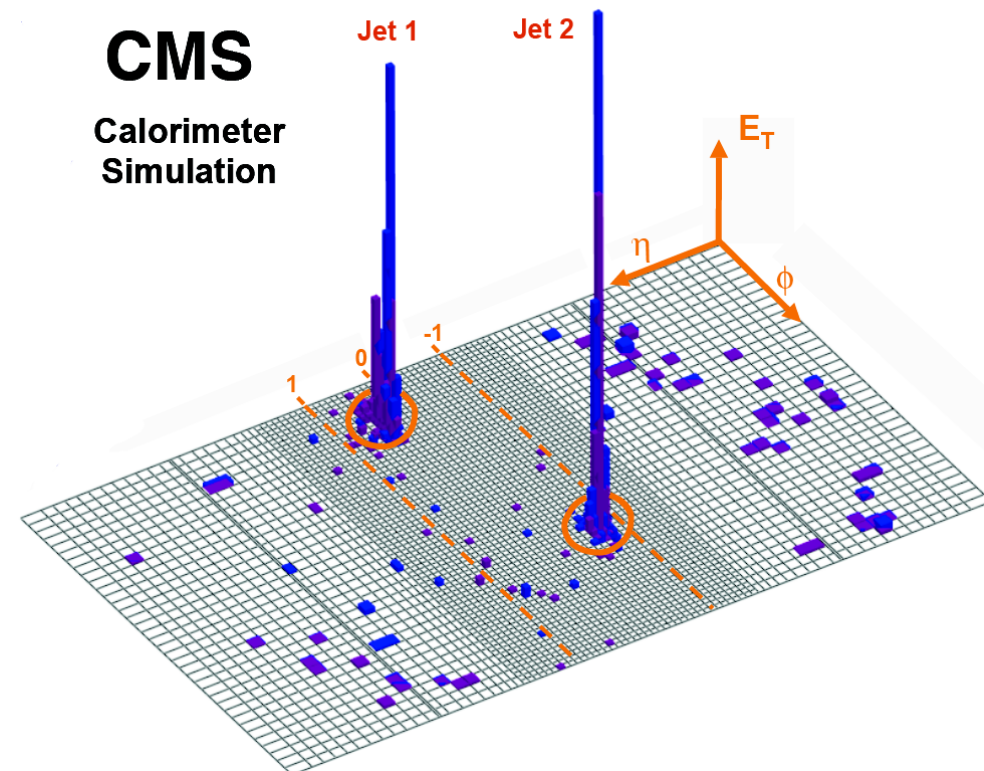
SISCone, $R=0.5, 0.7$

infrared and collinear safe,
FastJet implementation

incl. k_T , $D=0.4, 0.6$

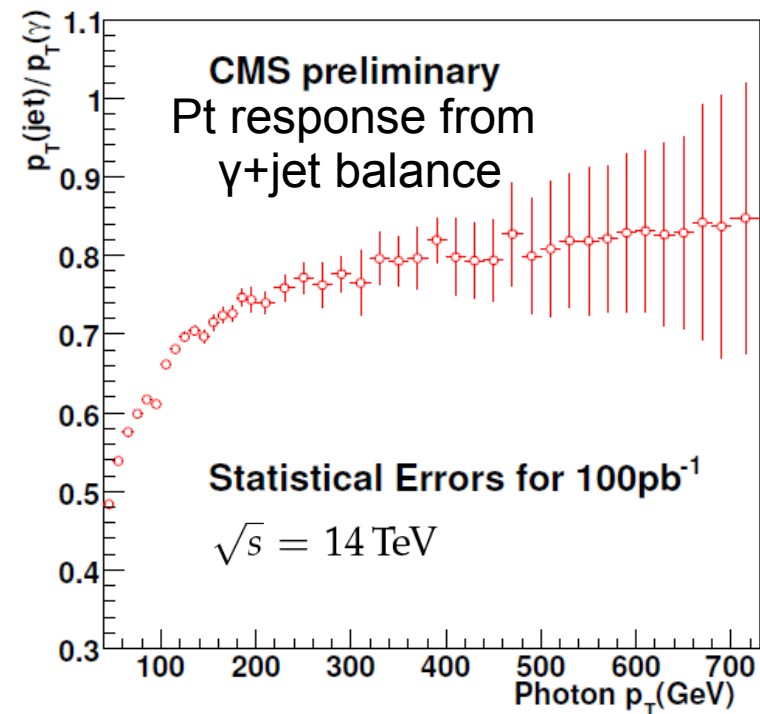
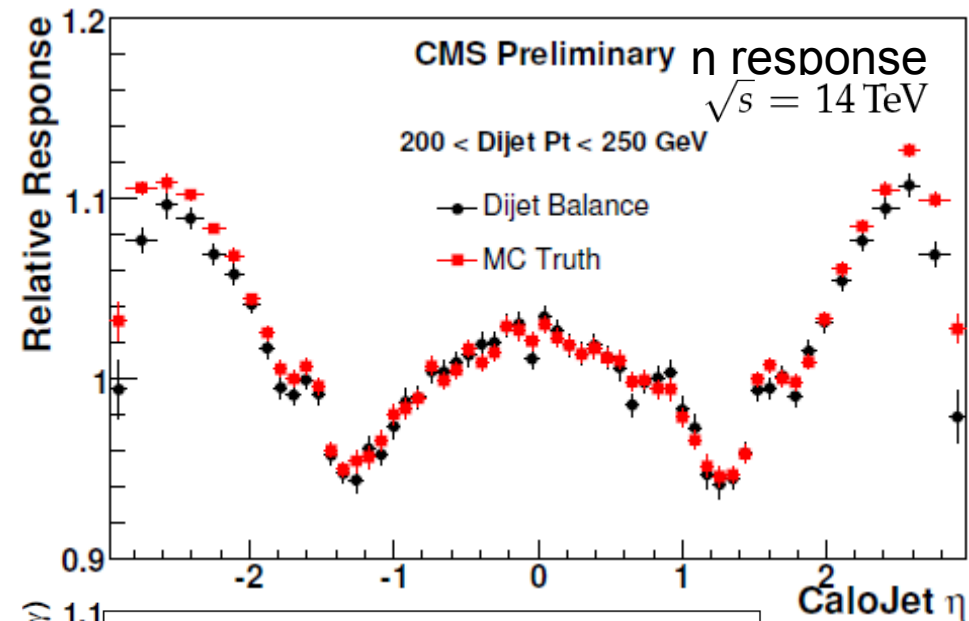
infrared and collinear safe,
use *FastJet* implementation

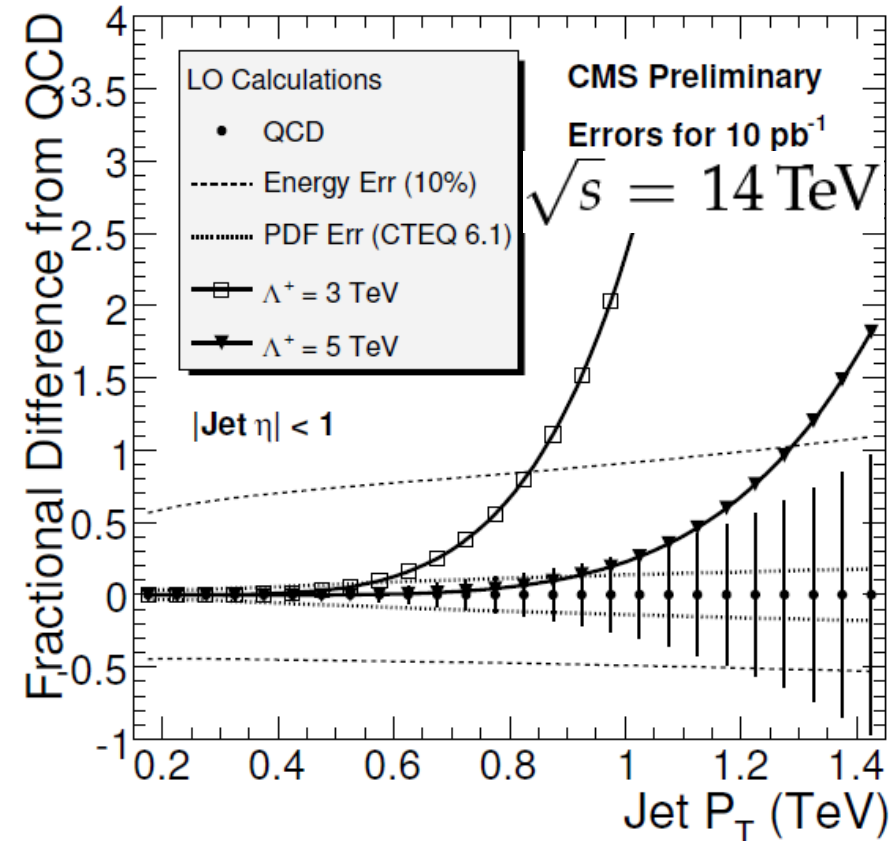
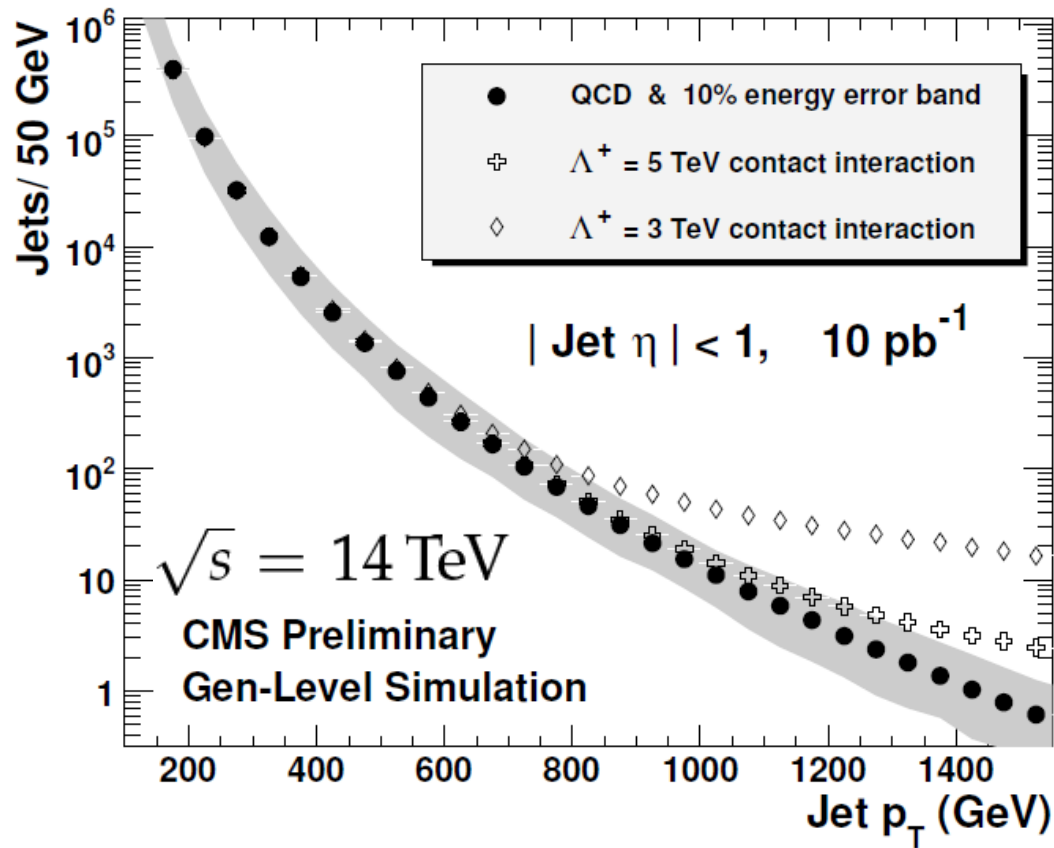
CMS
Calorimeter
Simulation



FastJet: <http://www.lpthe.jussieu.fr/~salam/fastjet/>

- **Default: correction to particle level**
- **Factorized approach** (like e.g. CDF)
 - **Offset correction** (remove pile-up and noise contribution)
 - **Relative correction** (flattens the jet response in η)
 - **Absolute correction** (flattens the jet response in p_T)
- **Data driven approach:**
 - Di-jet balancing (relative correction)
 - γ +jet, Z+jet balancing (absolute correction)
- **Optional corrections:**
 - Electromagnetic fraction dependence
 - Flavour dependence
 - Parton level
 - Underlying event
- **Systematic uncertainty $\sim 10\%$ at startup, improving to $\sim 5\%$ with first data.**



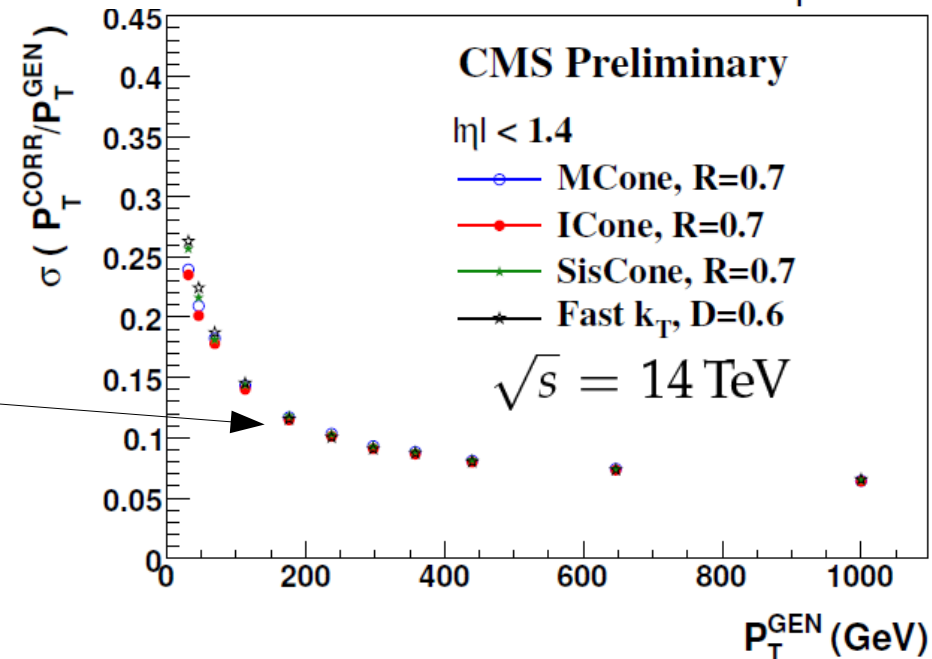
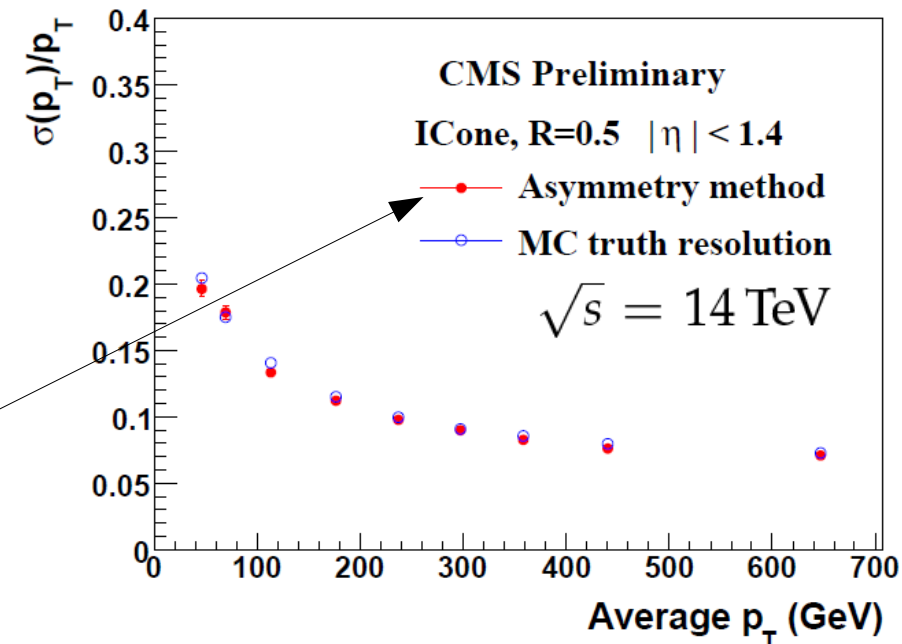


- Important jet commissioning measurement.
- Even a small amount of data is enough to exceed the Tevatron p_T reach (~700GeV).
- Sensitive to **contact interactions**: with 10pb⁻¹ @ 14TeV a discovery beyond the Tevatron limit (2.7 TeV) is possible.
- The jet energy scale is the dominating experimental systematic uncertainty.
- Could be used to constrain PDFs, which requires profound knowledge of systematic uncertainties.

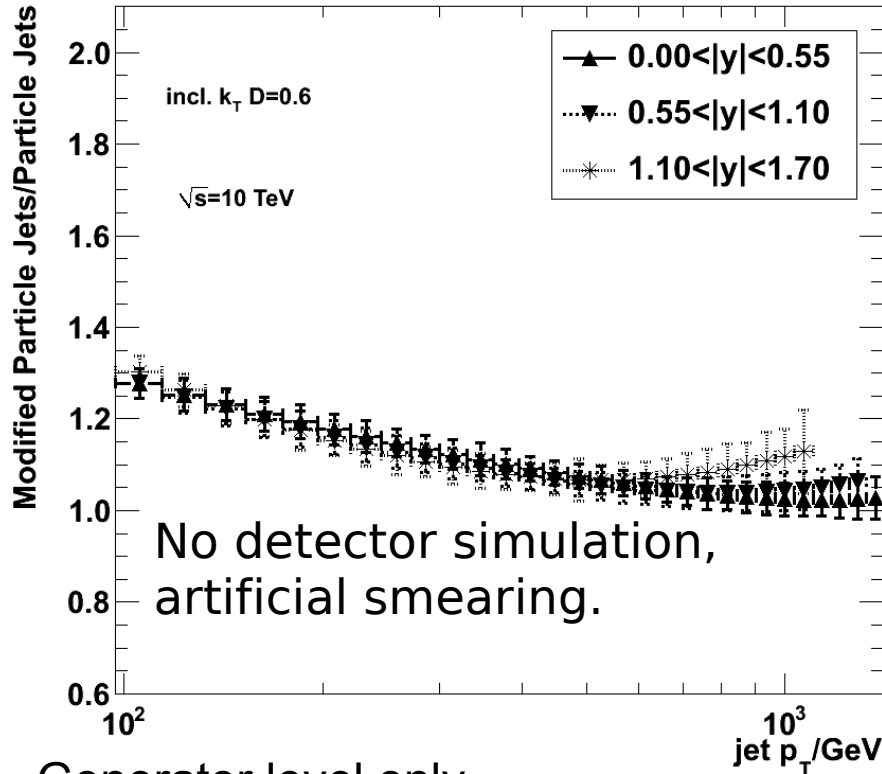
- Finite detector resolution affects measurements. The jet p_T resolution is a leading systematic effect for e.g. inclusive jets after jet energy scale and luminosity.
- Can be determined from MC truth matching (MC tuned to data), or by direct measurement from dijet events (Asymmetry Method).
- Jet p_T resolution usually parametrized by

$$\sigma(p_T) = p_T \cdot \sqrt{C^2 + \frac{S^2}{p_T} + \frac{N^2}{p_T^2}}$$

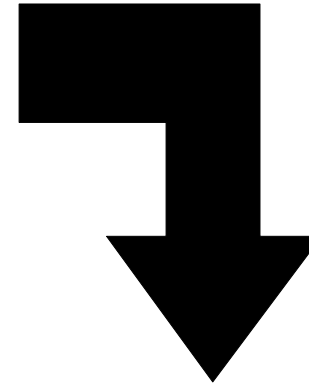
- Fairly independent of the specific jet algorithm.
- Can use combined calorimetric+track information (compensation correction) to improve the jet energy resolution



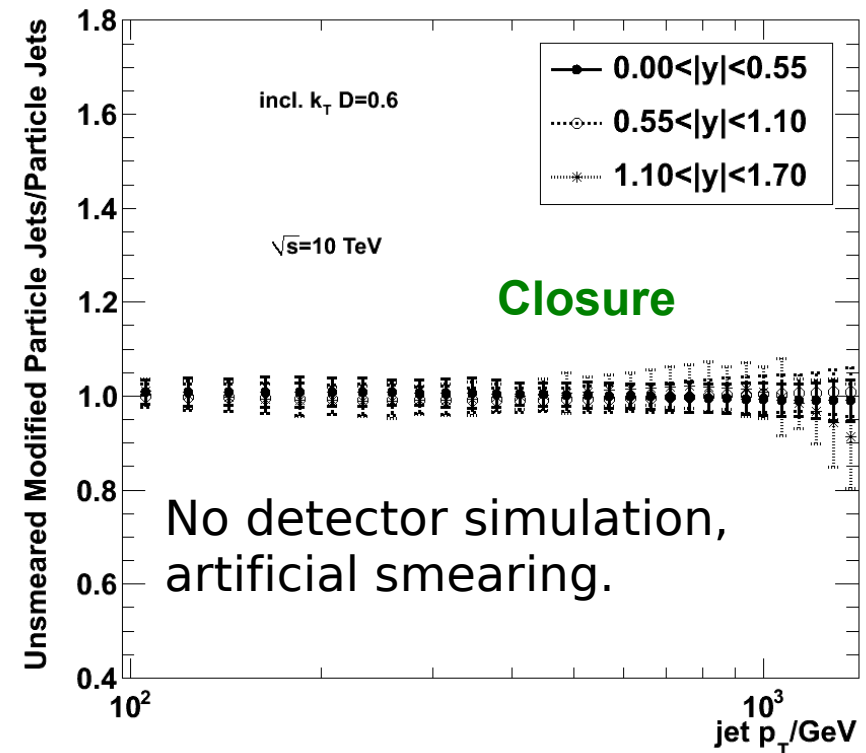
Inclusive jet spectrum needs a correction for the finite jet p_T resolution



- Generator level only.
- Artificially smear jets by Gaussian with an arbitrary (but reasonable) width.
- Apply Ansatz Method for unsmearing.
- Method corrects p_T smearing effects on steeply falling spectrum.



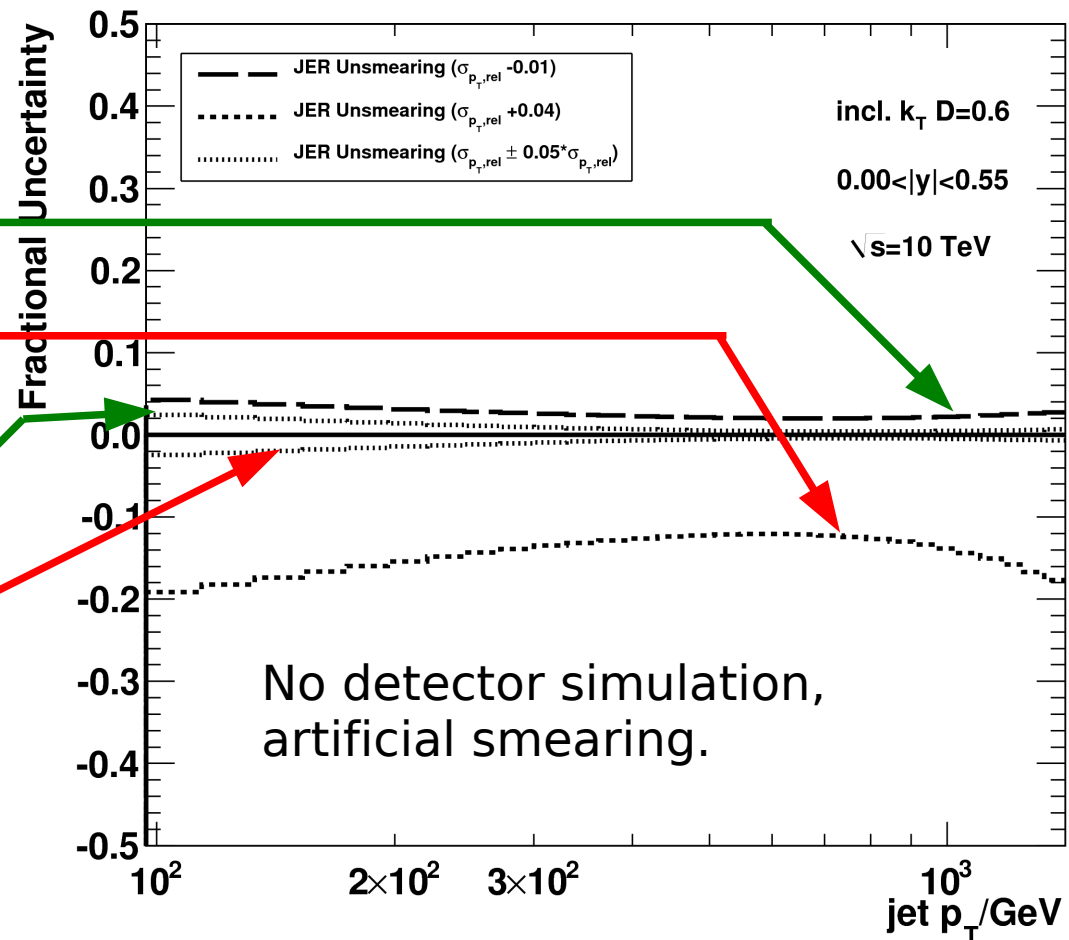
unsmearing by
"Ansatz Method"



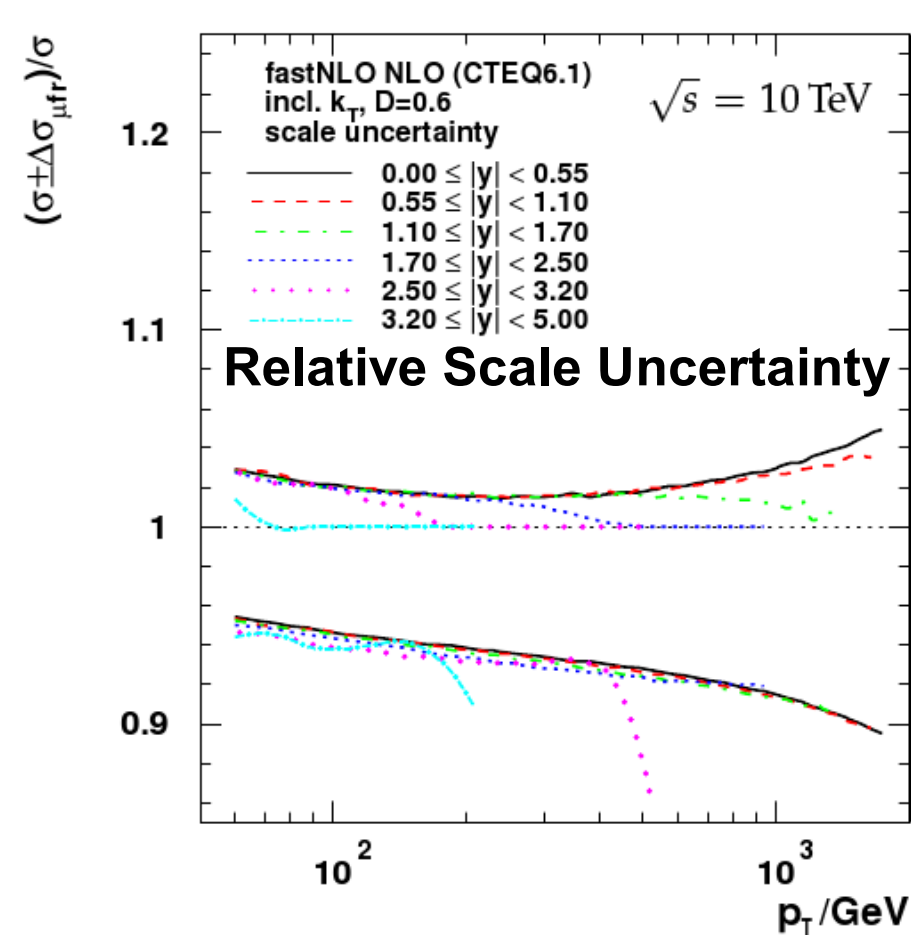
- Ansatz Method can correct finite resolution effects on the steeply falling spectrum
- However a good knowledge of the resolution is required
- A wrong assumption of the resolution can shift the final spectrum easily by some percent

Two unsmearing scenarios studied:

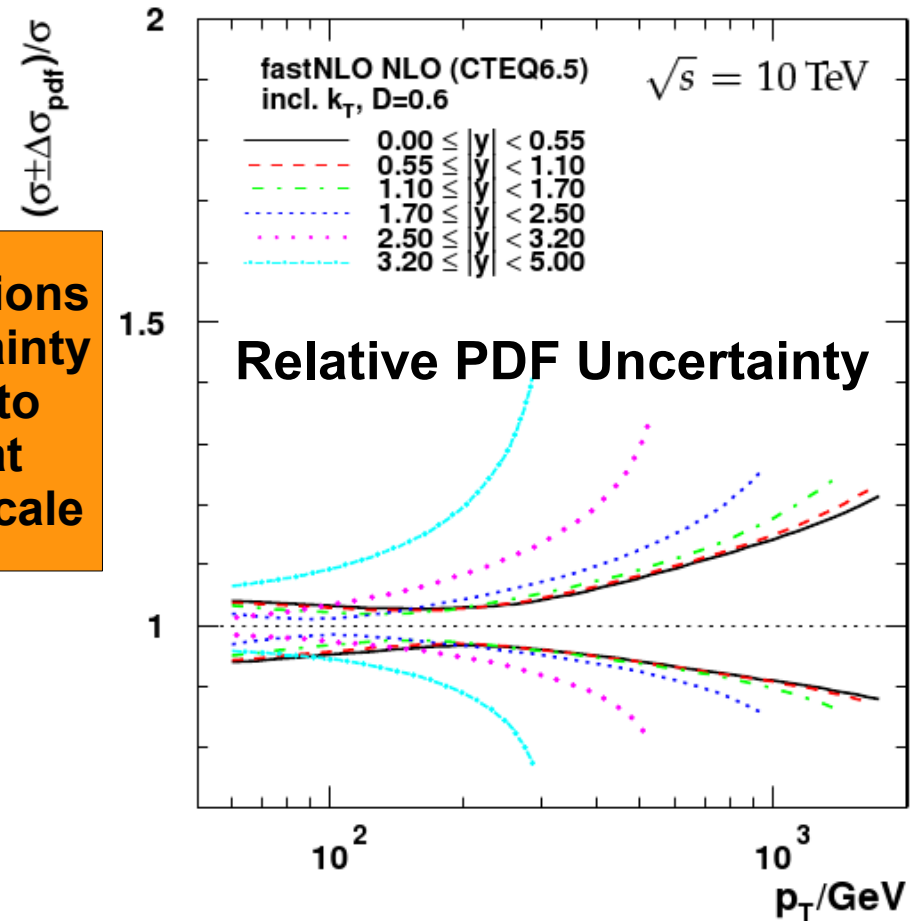
- Shift response used by *fraction of jet p_T*
 - assume **1% better** resolution in unsmearing than used for smearing
 - assume **4% worse** resolution in unsmearing than used in smearing
- Shift response used by *fraction of itself*
 - assume **5% better** resolution in unsmearing than used for smearing
 - assume **5% worse** resolution in unsmearing than used for smearing



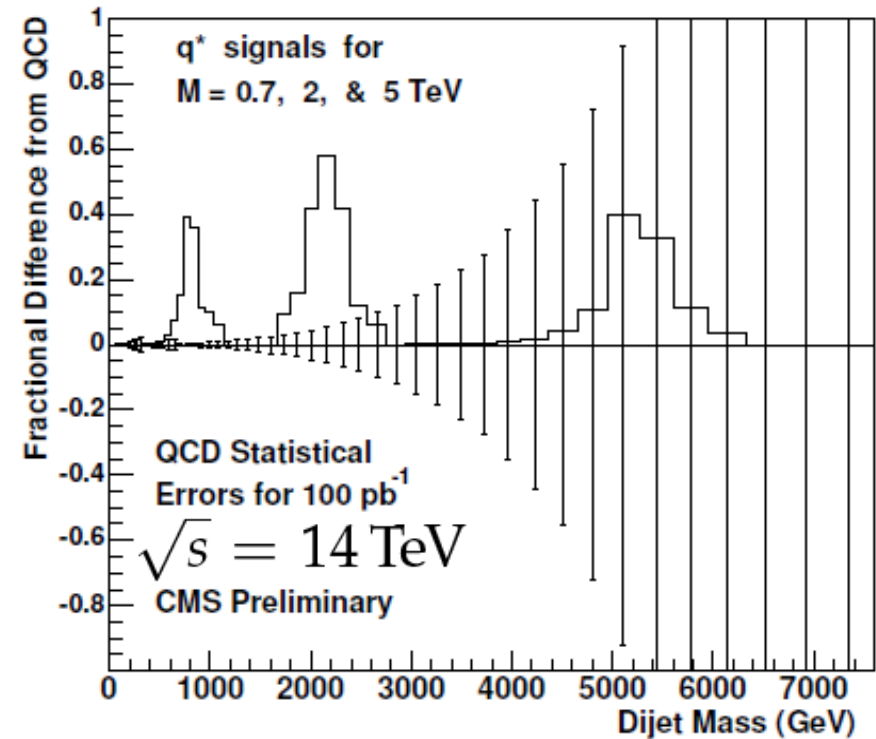
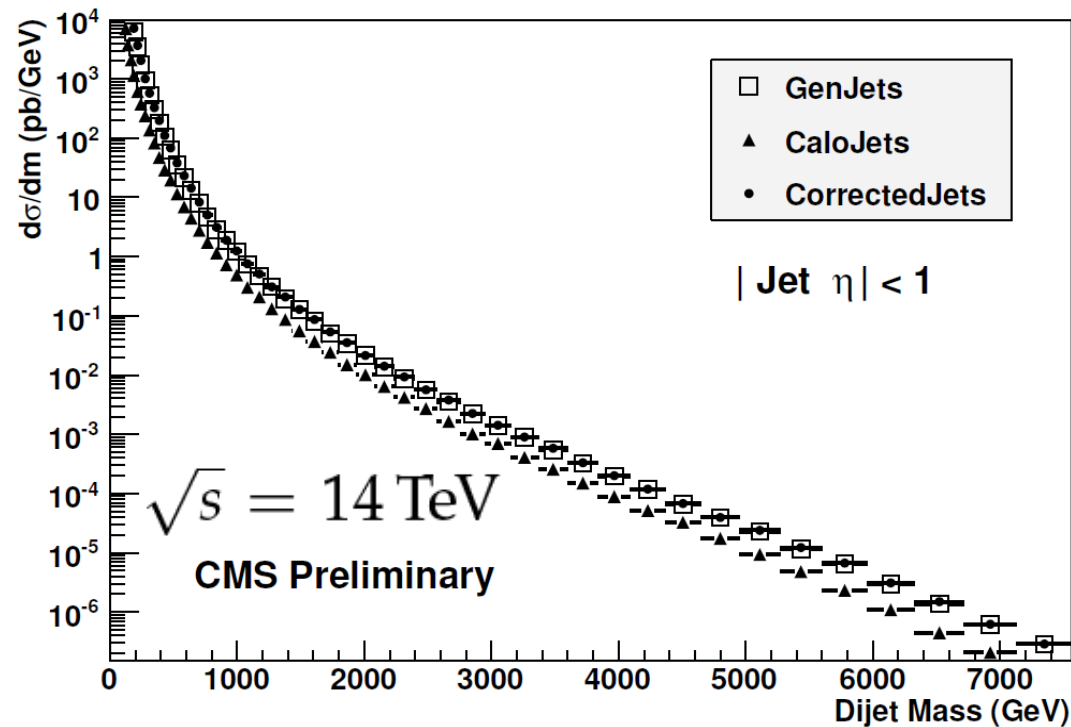
- Dominant theoretical uncertainties come from the uncertainties of the parton distribution functions and scale uncertainties in the calculations



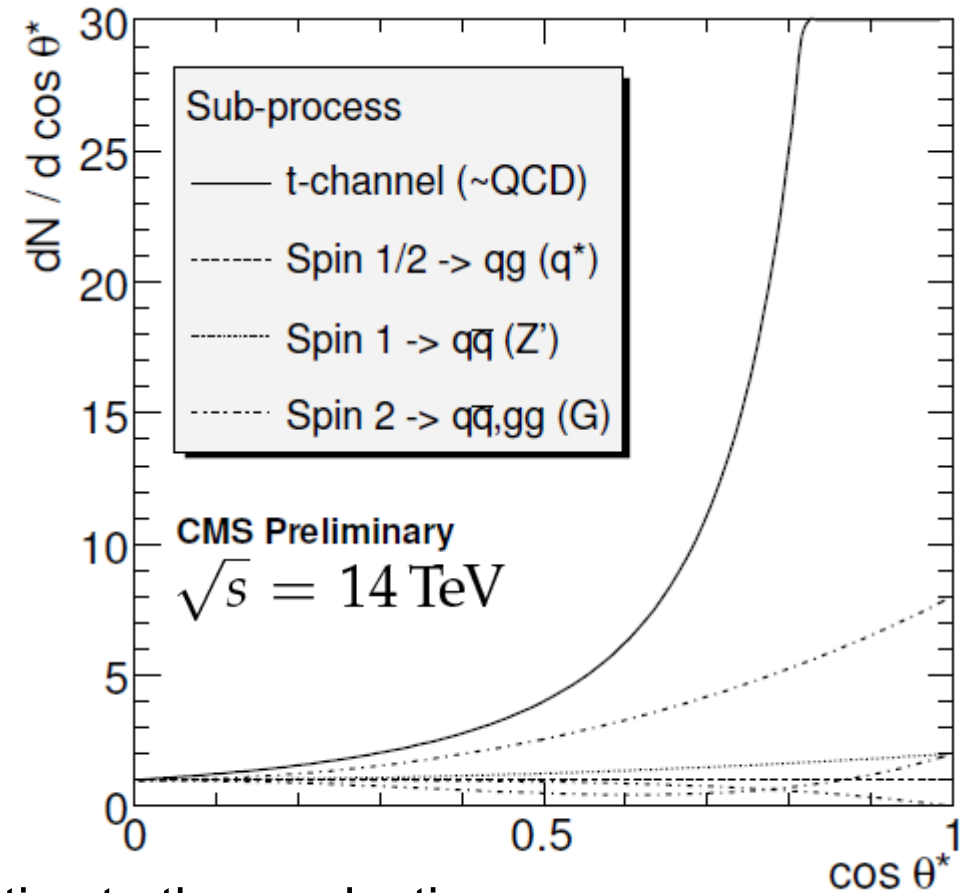
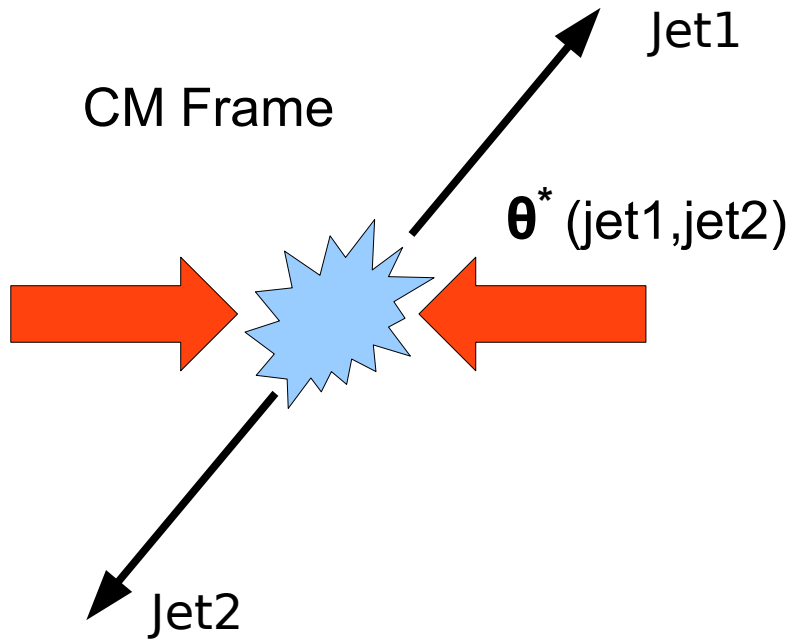
Contributions to uncertainty rise up to 5-20% at the TeV scale



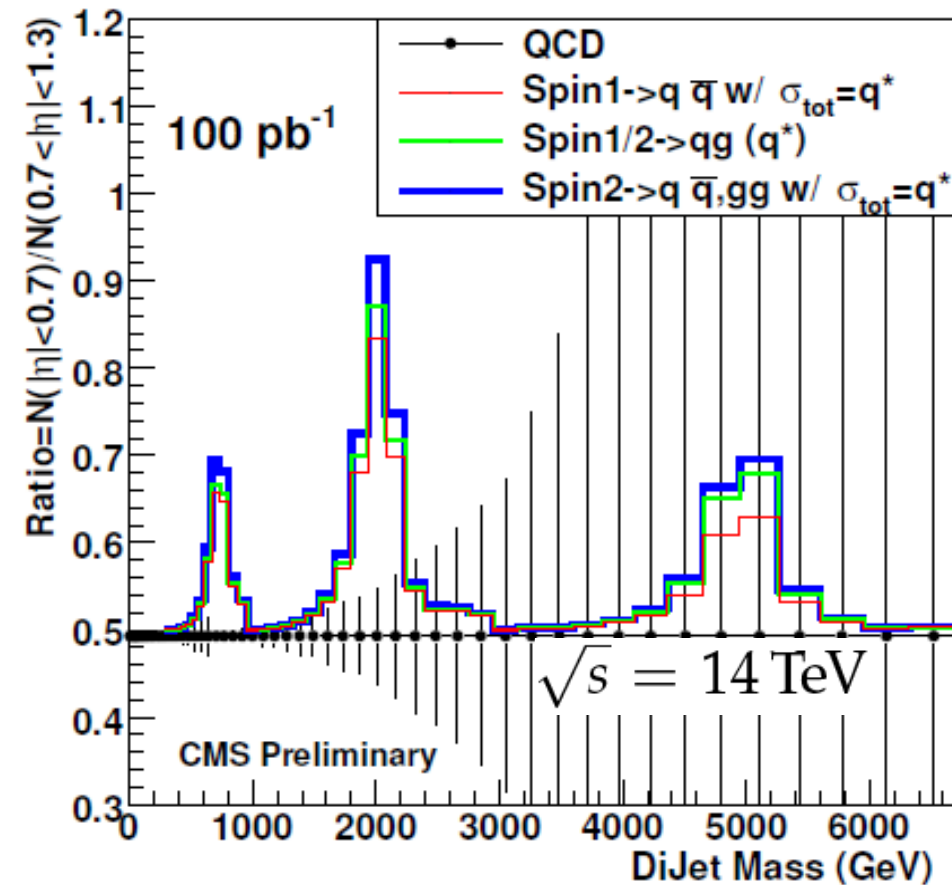
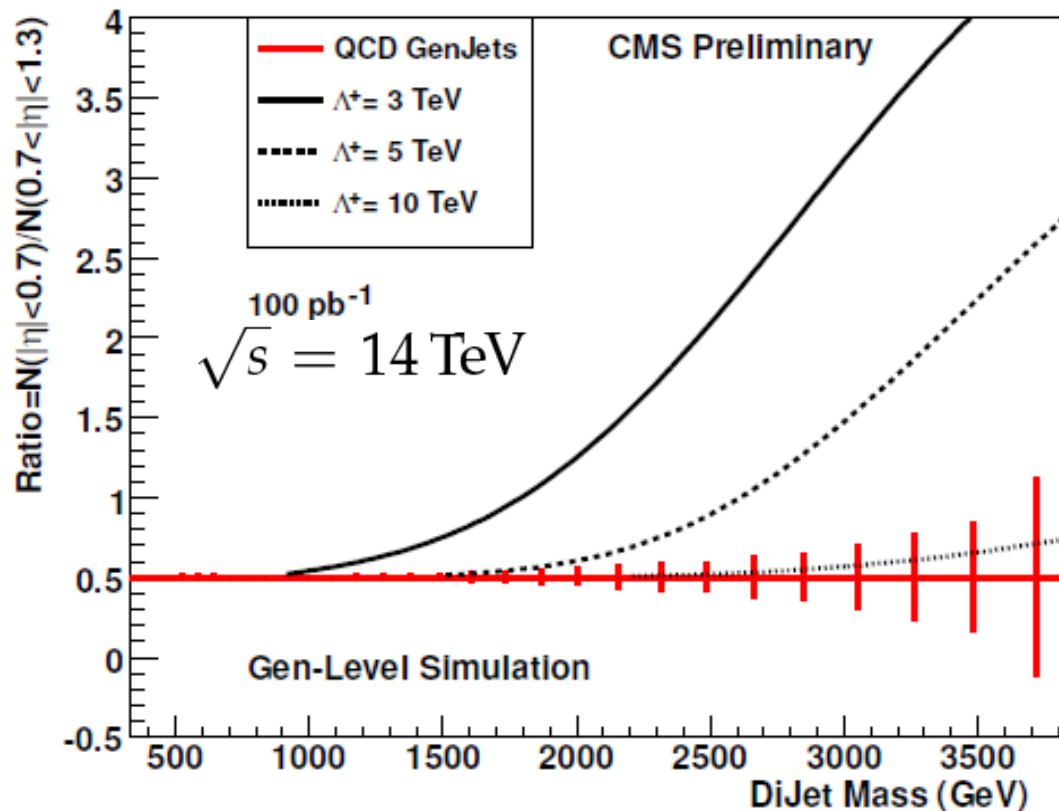
- Probing pQCD with high p_T jets means mainly comparing measurements to calculations in NLO. This procedure currently involves uncertainties from the non-perturbative corrections needed (uncertainty about 15% at low p_T , reduced to almost zero at 1 TeV).



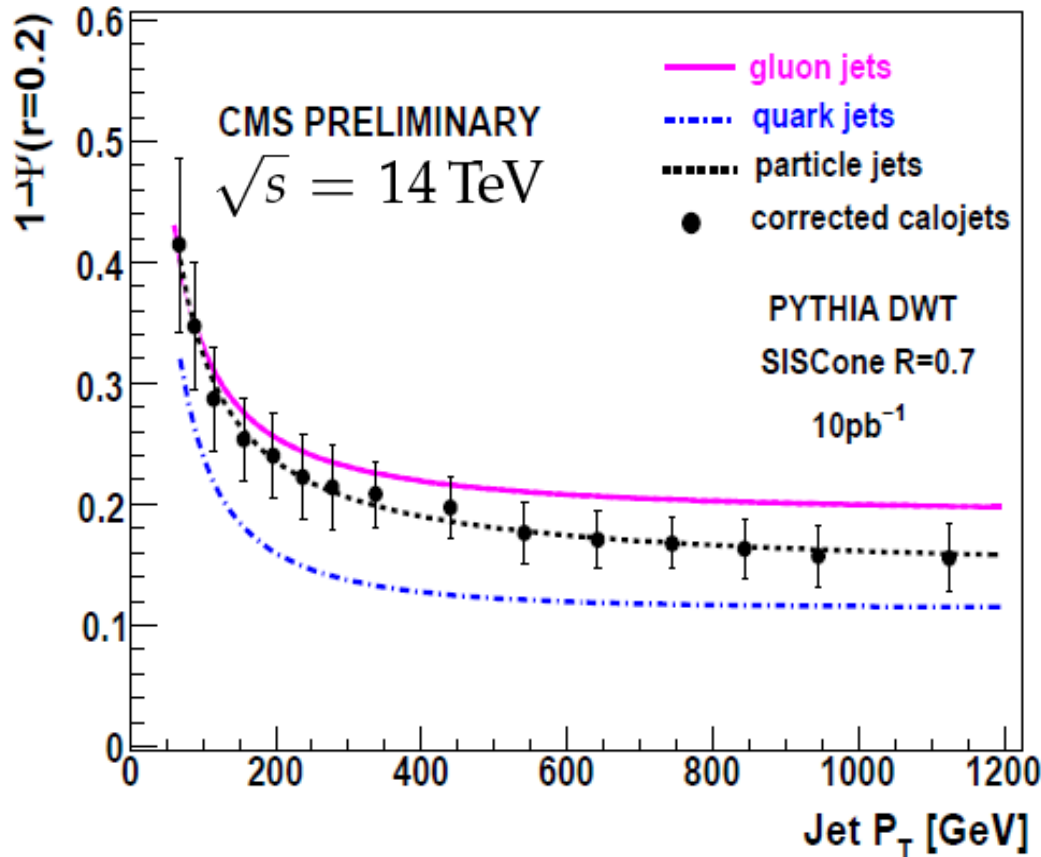
- Invariant mass of 2 leading jets.
- The jet energy scale is the dominant experimental systematic uncertainty
- Sensitive to objects decaying into 2 jets (**dijet resonances**: q^* , Z' , etc.)
- **With 100 pb^{-1} a 2 TeV excited quark can be clearly seen (Tevatron excludes: $M < 0.87$ TeV)**



- The dijet angular distributions are sensitive to the production process:
 - QCD dominated by t-channel scattering
 - Resonance decays and contact interactions tend to be more isotropic
- The dijet angular distributions will confirm any deviation from QCD observed in the dijet mass spectrum



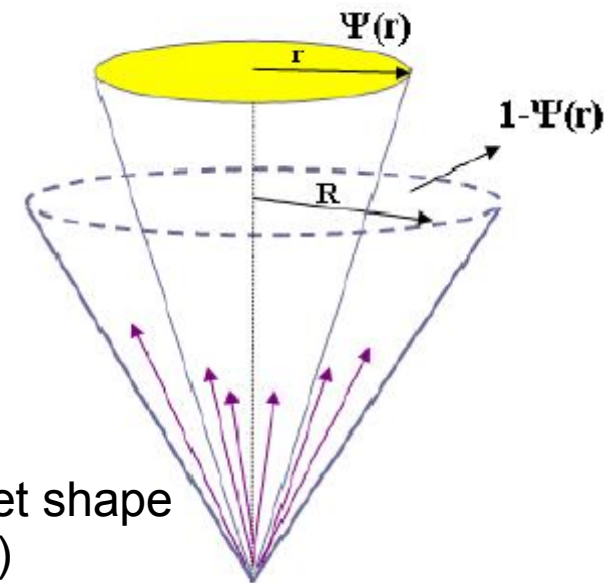
- The dijet ratio $N(|\eta| < 0.7) / N(0.7 < |\eta| < 1.3)$ vs mass is a robust way to identify deviations from QCD.
- Sensitivity to **contact interactions** (discovery up to $\Lambda = 6.8$ TeV with 100 pb⁻¹ @ 14 TeV)
- Sensitivity to **dijet resonances**. At higher luminosity spin measurement is also feasible.
- Ideal measurement for startup. JES and luminosity uncertainties are largely cancelled.



For Hadronic Event Shape studies see talk of M. Weber.

- Jet shape measurements can be used to discriminate between different underlying event models.
- Can be used to distinguish gluon initiated jets from quark jets.

- Measurement of the average integrated (differential) energy flow inside jets.
- Jet shape measurements can be used to test the showering models in the MC generators.



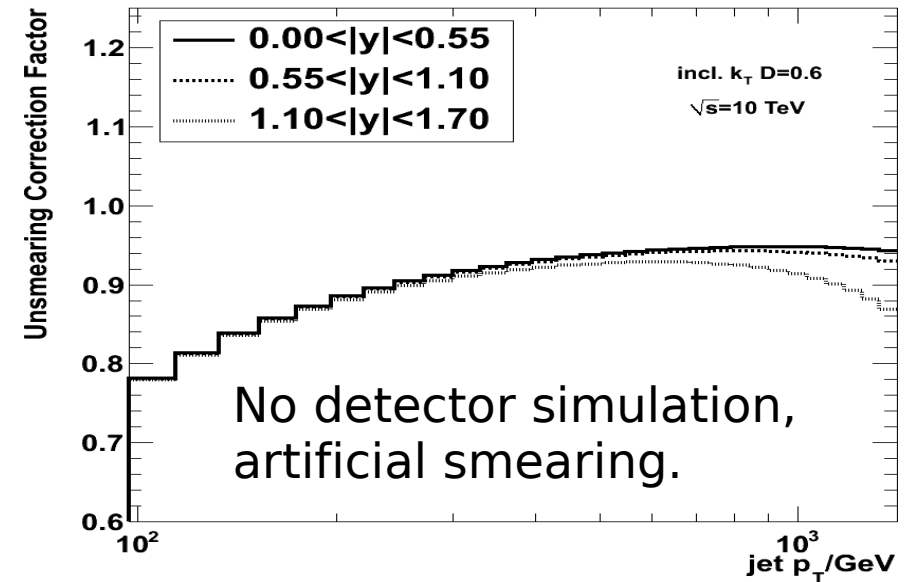
Integrated jet shape $\Psi(r)$

- Jet final states are critical both for studying the detector performance and for the “re-discovery” of the Standard Model at the LHC.
- LHC p-p collisions will probe a previously unexplored kinematic region and a **small amount of early data** will be enough for QCD measurements **exceeding the p_T reach of the Tevatron**.
- Jet measurements are sensitive to new physics. **Contact interactions and resonances decaying into dijets can be discovered early on**, even with 10% JES uncertainty at startup.
- Precision QCD measurements will require much more data and effort to constrain the systematic uncertainties.
- **Establishing the JES with a small systematic uncertainty ($\sim 1\%$) is a challenging task which will require a huge effort and a large integrated luminosity.**

BACKUP SLIDES

Motivation

The **observed** cross section is higher **than** the true one due to the falling shape of the spectrum and the finite p_T resolution. More events migrate into a bin of measured p_T than out of it.



Unsmearing steps:

- Analytical expression of the p_T resolution $\Rightarrow R(p'_T, p_T) = \frac{1}{\sqrt{2\pi}\sigma(p'_T)} \exp\left[-\frac{(p'_T - p_T)^2}{2\sigma^2(p'_T)}\right]$
- Ansatz function with free parameters to be determined by the data $\Rightarrow f(p_T) = N \cdot p_T^{-a} \cdot \left(1 - \frac{2 \cosh(y_{min}) p_T}{\sqrt{s}}\right)^b \exp(-\gamma p_T)$
- Fitting the data with the Ansatz function smeared with p_T resolution. $\Rightarrow F(p_T) = \int_0^\infty f(p'_T) R(p'_T, p_T) dp'_T$
- Unsmearing correction calculated bin by bin. $\Rightarrow C_{bin} = \frac{\int_{bin} f(p_T) dp_T}{\int_{bin} F(p_T) dp_T}$

- Correction factor for non perturbative effects, for which NLO calculations do not account for
- PYTHIA: V6.4, D6T, CTEQ6L1
- HERWIG++: default tune, MRST2001

