

# Small- $x$ resummed DIS phenomenology

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PDF and low- $x$  Session

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# MOTIVATION

# Small- $x$ resummation and phenomenology

- Any realistic small- $x$  resummation procedure should:
  - Include a stable solution of BFKL with running coupling at NLL
  - Match smoothly to DGLAP at moderate and large  $x$
  - Provide the complete set of splitting functions and (at least) DIS coefficient functions in a well defined factorization scheme
- Full set of small- $x$  resummed DIS splitting functions and coefficient functions recently available (Altarelli, Ball and Forte, NPB799:199-240,2008) where all above requirements met  $\rightarrow$  Next step is to study its phenomenological implications
- See also S. Marzani's talk for discussion on the small- $x$  resummation of the Drell-Yan process
- Important related work by the Ciafaloni-Colferai-Salam-Stasto (see A. Stasto's talk) and Thorne-White groups
- Not a review talk  $\rightarrow$  Apologies in advance for lack of references!



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# SMALL-X RESUMMATION & PRACTICAL IMPLEMENTATION



# ABF approach to small- $x$ resummation

- Duality between DGLAP kernel  $\gamma(N, \alpha_s)$  and BFKL kernel  $\chi(M, \alpha_s)$

$$\chi(\gamma(N, \alpha_s), \alpha_s) = N$$

Solution of BFKL and DGLAP equations coincide at leading twist (fixed  $\alpha_s$ )  
 $\rightarrow$  Double leading (logs of  $x$  and  $Q^2$ ) expansion

Automatic resummation of collinear poles of BFKL kernel  $\chi$  which makes its perturbative expansion not reliable

$$\gamma(1, \alpha_s) = 0, \quad \rightarrow \quad \chi(0, \alpha_s) = 1$$

- Symmetry of BFKL kernel upon gluon exchange: at fixed  $\alpha_s$

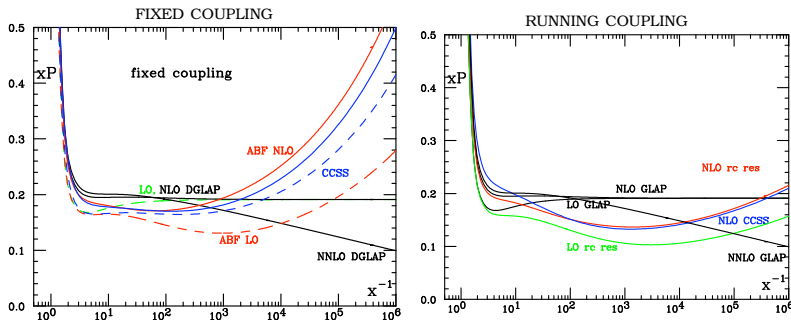
$$\chi(M) = \chi(1 - M)$$

Use collinear resummation for  $M \sim 0$  in DL expansion in anti-collinear region  
 $M \sim 1 \rightarrow$  Stable perturbative expansion at resummed level

- Running coupling resummation of BFKL solution
- Factorization scheme defined at the resummed level, whose relation with large- $x$  schemes is known exactly

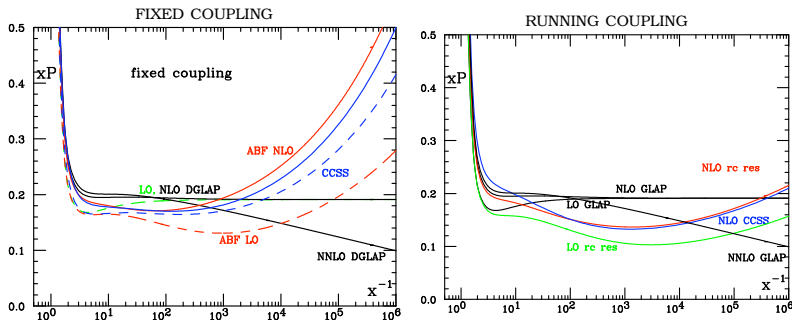


# Example: Resummed $xP_{gg}(x, \alpha_s)$ for $N_f = 0$



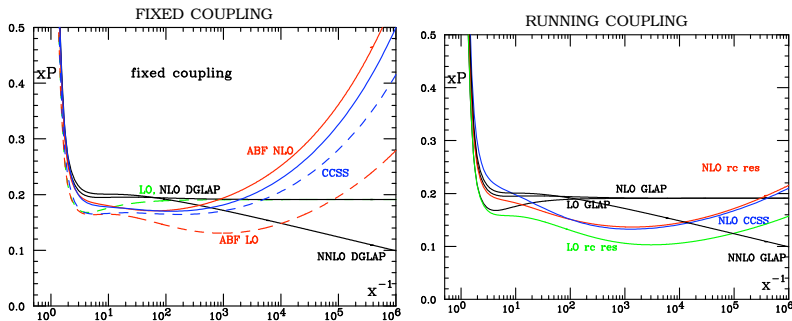
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- Fast convergence of resummed expansions
- Match to large- $x$  DGLAP expansion for  $x \leq 10^{-2}$
- Small- $x$  rise tamed by running coupling effects
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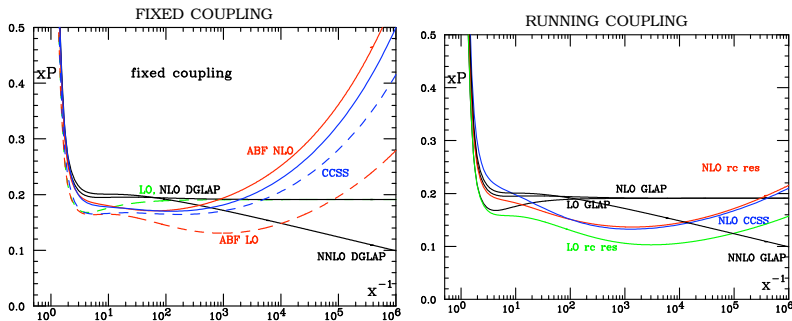
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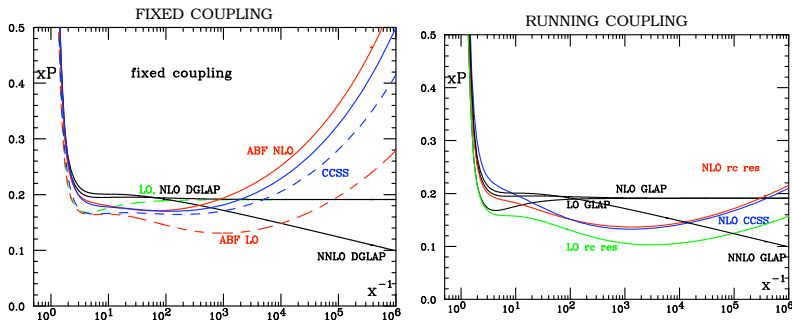


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# Resummed K-factors

For **realistic phenomenology**  $\rightarrow$  Compute **small- $x$  resummed K-factors** with **NNPDF1.0** **parton distributions** (see **my talk and M. Ubiali's one** this morning)

$$K_{F_2}^{\text{res}}(x, Q^2) \equiv F_2^{\Sigma}(x, Q^2) \Big|_{\text{NLOres}} / F_2^{\Sigma}(x, Q^2) \Big|_{\text{NLO}}$$

K-factors obtained **fixing structure functions** at **HERA** matching scale  $t_0 \equiv \ln Q_{\text{match}}^2$ :

$$F_{\text{NLO}}(N, t_0) = c_{\text{NLO}}(N, \alpha_s(t_0)) f_{\text{NLO}}(N, t_0).$$

**redefining the input PDFs**

$$f_{\text{res}}(N, t_0) = C_{\text{res}}^{-1}(N, \alpha_s(t_0)) F_{\text{NLO}}(N, t_0).$$

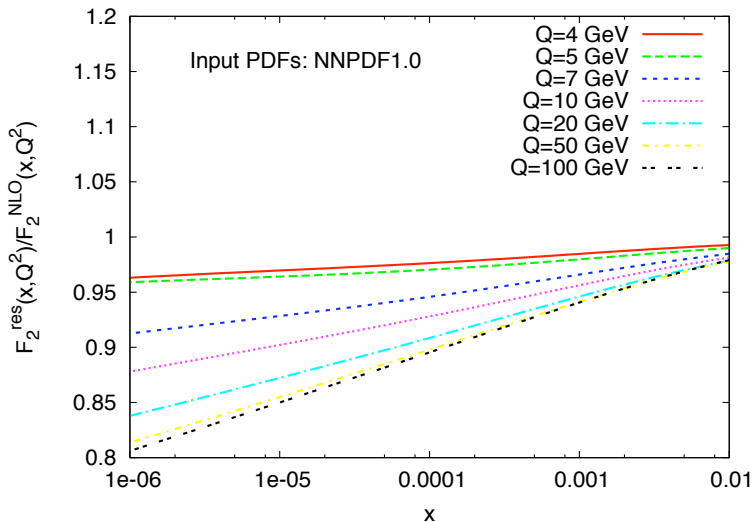
then **evolve with appropriate kernels** - closer to a real **PDF fit**!

$$\begin{aligned} F_{\text{res/nlo/nnlo}}(N, t) &= C_{\text{res/nlo/nnlo}}(N, \alpha_s(t)) \Gamma_{\text{res/nlo/nnlo}}(t, t_0) f_{\text{res/nlo/nnlo}}(N, t_0) \\ F_{\text{res}}(N, t_0) &= F_{\text{NNLO}}(N, t_0) = F_{\text{NLO}}(N, t_0) \end{aligned}$$

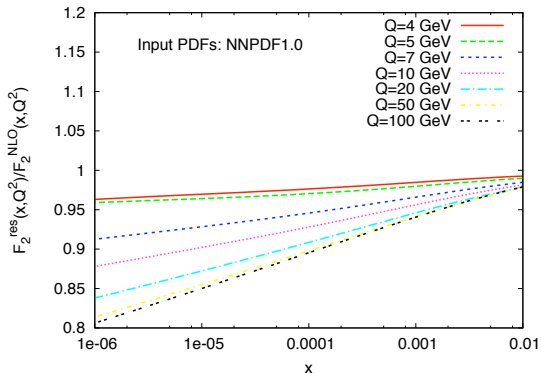
Results shown here assume  $Q_{\text{match}} = 3 \text{ GeV}$



# Resummed K-factors



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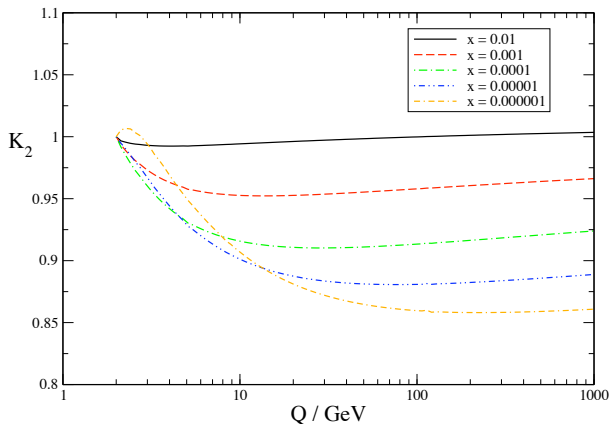


- Smooth matching with DGLAP region  $\lim_{x \rightarrow 1} K^{\text{res}}(x, Q^2) = 1$
- Resummed corrections similar in size to NNLO, but opposite direction
- Resummation effects small for  $Q^2 \sim Q_{\text{match}}^2$  (which fixes the PDFs), differences seen in  $Q^2$  evolution

# Resummed K-factors

Smooth large  $Q^2$  limit (Toy PDFs, [arxiv:0901.2504](https://arxiv.org/abs/0901.2504))

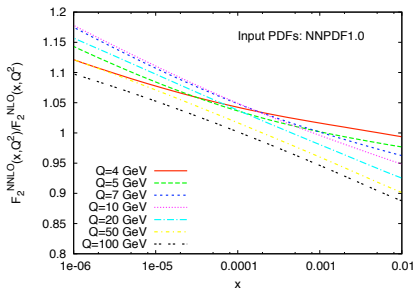
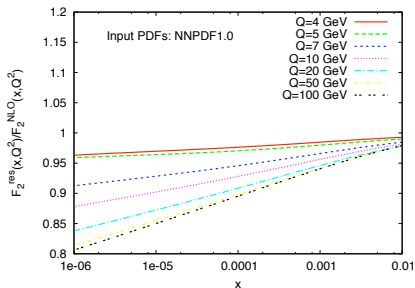
→ Asymptotic freedom at work!



# Resummed K-factors for HERA

Compare **NLL** and **NNLO** K-factors in HERA region

For  $F_2(x, Q^2)$



- NLL small- $x$  corrections have **similar size but opposite sign** as NNLO

# SMALL-X RESUMMATION DIS PHENOMENOLOGY

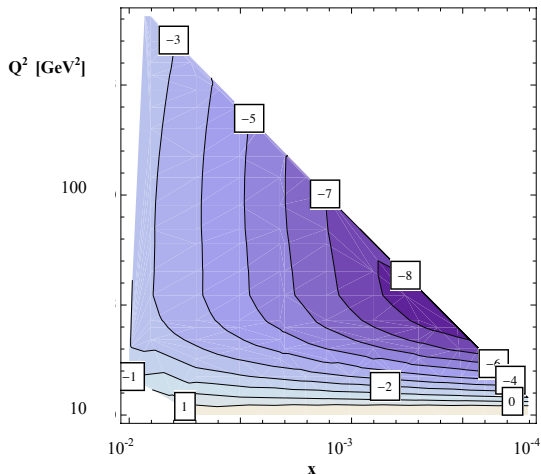




# Small- $x$ resummation at HERA

Size of **small- $x$  resummation effects** in HERA kinematics

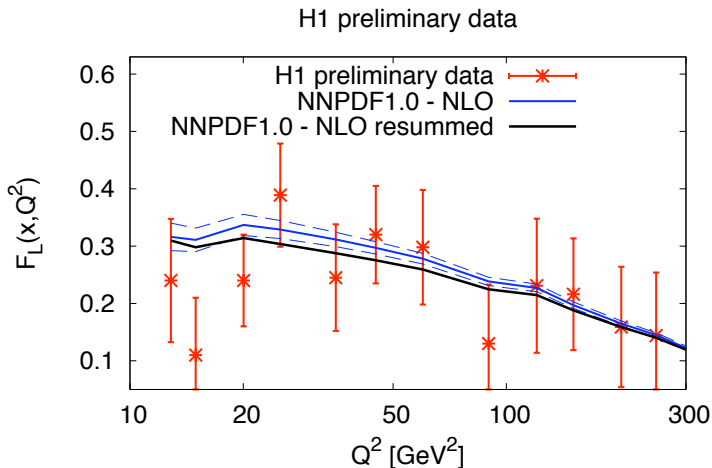
% contribution from NLOres to F2p – HERA kin



- Moderate size of corrections explain why DGLAP works so well at HERA ....
- ... but size large enough to be discriminated if ultimate HERA precision achieved!

# Small- $x$ resummed $F_L$

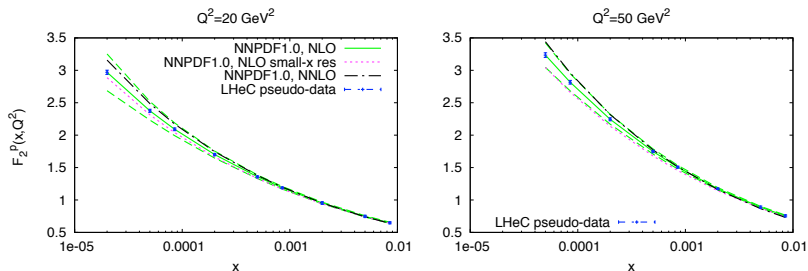
Small- $x$  resummation effects in  $F_L$  larger than PDF uncertainties, but too large experimental errors to discriminate scenarios



# Small- $x$ resummation at the LHeC

Small- $x$  resummation effects larger for the LHeC  
(see my talk at the *Future Facilities* session)

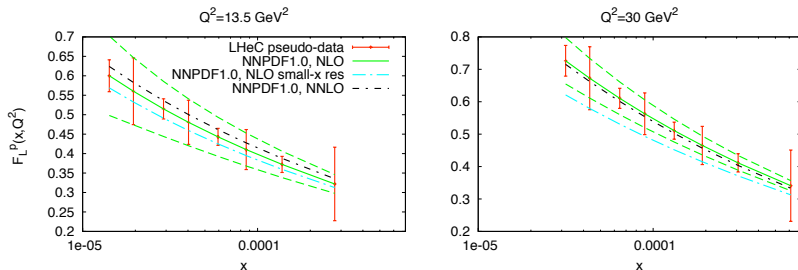
Again notice **opposite trend of NLOres and NNLO corrections**



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# ULTRA HIGH ENERGY NEUTRINO INTERACTIONS



# Ultra-High Energy (UHE) Neutrinos

- Neutrino telescopes (e.g. **ICECUBE**) and cosmic ray observatories (e.g. **AUGER**)  
→ potential for detecting **UHE** neutrinos up to  $E_\nu = 10^{12}$  GeV ( $\sqrt{s} \simeq 10^2 \sqrt{s}_{\text{LHC}}$ )
- Detection of cosmic neutrinos → **New window to the universe**
- Detection of cosmic neutrinos → **Probe the Standard Model** in a completely unexplored kinematical regime
- In this talk → Implications of UHE neutrinos for Quantum Chromodynamics (But also much to learn for **New Physics!**)



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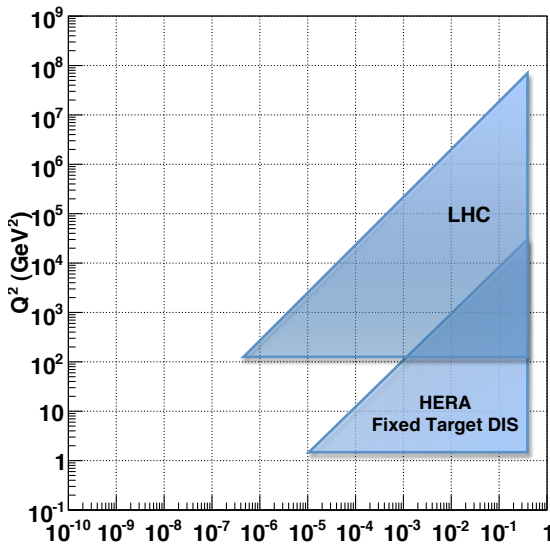


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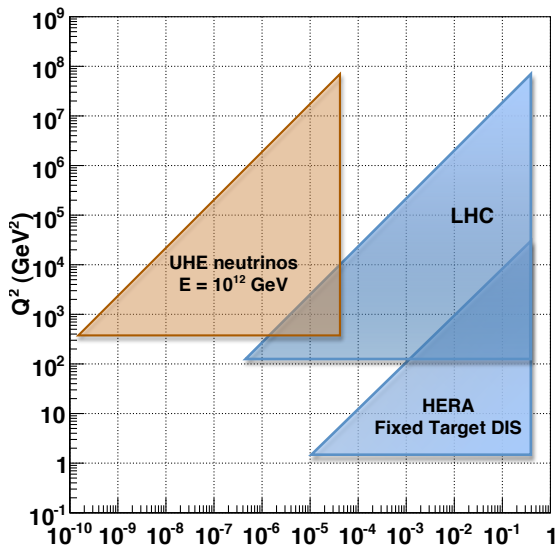
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# UHE Kinematics



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The total **neutrino-nucleus scattering** cross section:

$$\sigma^{\nu N}(E_\nu) = \int_{Q_0^2/s}^1 dx \int_{Q_0^2}^{xs} dQ^2 \frac{d^2 \sigma^{\nu N}}{dx dQ^2}(x, Q^2), \quad s = 2M_N E_\nu$$

can be decomposed in terms of **structure functions**

$$\begin{aligned} \frac{d^2 \sigma^{\nu(\bar{\nu})p}}{dx dQ^2}(x, y, Q^2) &= \frac{G_F^2}{4\pi x} \left( \frac{M_W^2}{Q^2 + M_W^2} \right)^2 \\ &\times \left[ Y_+ F_2^{\nu(\bar{\nu})p}(x, Q^2) - y^2 F_L^{\nu(\bar{\nu})p}(x, Q^2) \pm Y_- x F_3^{\nu(\bar{\nu})p}(x, Q^2) \right] \end{aligned}$$

Integration dominated by  $Q^2 \sim M_W^2$ ,  $x \ll 1$  region

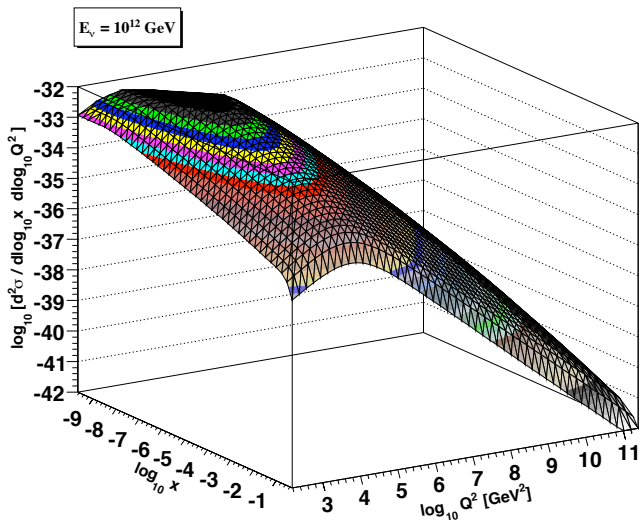
$$\langle x \rangle \sim \frac{M_W^2}{2M_N E_\nu} \ll 1$$

→ UHE interactions very sensitive to

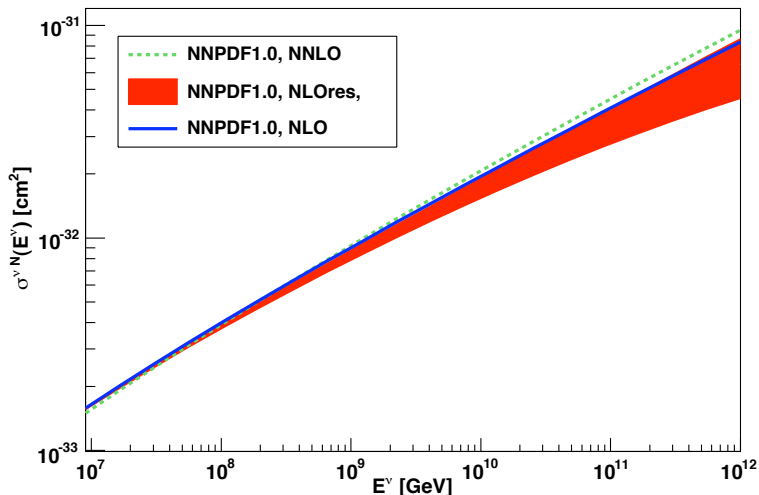
- Small- $x$  Parton Distributions
- Small- $x$  (BFKL) effects



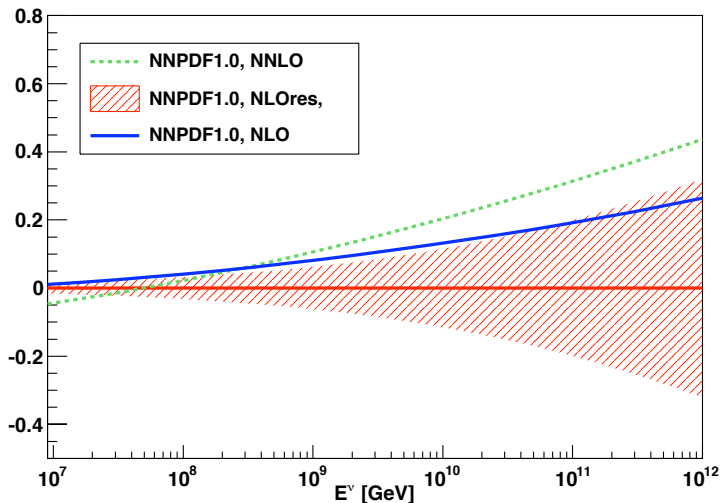
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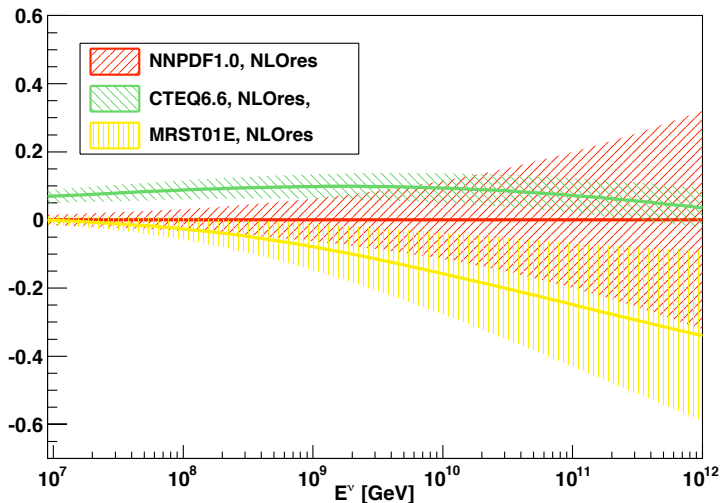
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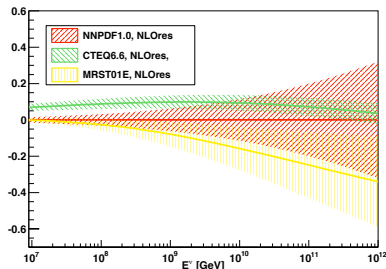
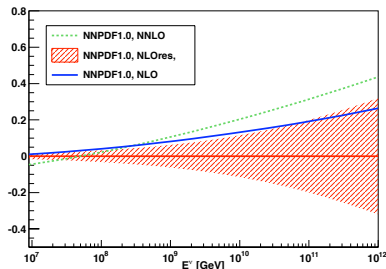


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# Small- $x$ resummation of UHE neutrinos



- Small- $x$  resummation decreases  $\sigma^\nu(E_\nu)$  with respect NNLO by a factor 2 at  $E_\nu = 10^{12}$  GeV (Similar trend in *saturation models*)
- NNLO-NLOres differences outside PDF uncertainty band
- PDF uncertainties sizable in *extrapolation region*  $\rightarrow$  Should be reduced in near future with final HERA and LHC data for small- $x$  PDFs

# OUTLOOK



# Outlook

- **BFKL small-x resummed DIS structure functions (with associated resummed K-factors)** available after several years of continuous efforts
- Small-x corrections have a moderate effect in HERA region, but could be discriminated if precision high enough, even more at the LHeC
- Cosmic neutrinos provide a useful testing ground for small-x QCD scenarios
- The K-factor approach provides realistic qualitative insights, but accuracy limited by **matching scale choice** → A fully resummed small-x PDF analysis is ultimately unavoidable
- Work in progress in this direction in the frameworks of the **CTEQ** (P. Nadolsky) and **NNPDF** global PDF analysis

Thanks for your attention!



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- The K-factor approach provides realistic qualitative insights, but accuracy limited by matching scale choice → A fully resummed small- $x$  PDF analysis is ultimately unavoidable
- Work in progress in this direction in the frameworks of the CTEQ (P. Nadolsky) and NNPDF global PDF analysis

**Thanks for your attention!**

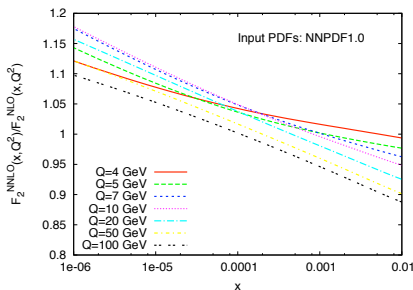
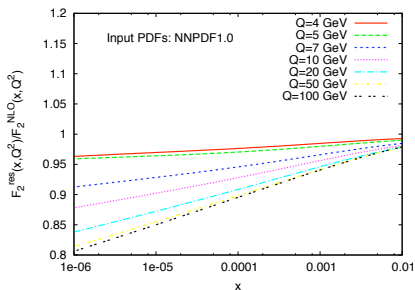


# EXTRA MATERIAL

# Resummed K-factors for HERA

Compare **NLL** and **NNLO** K-factors in HERA region

For  $F_2(x, Q^2)$

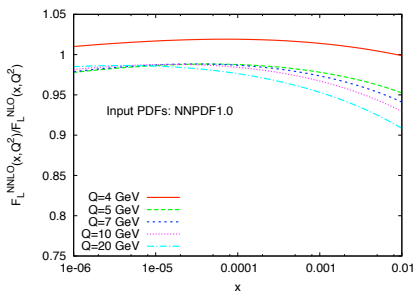
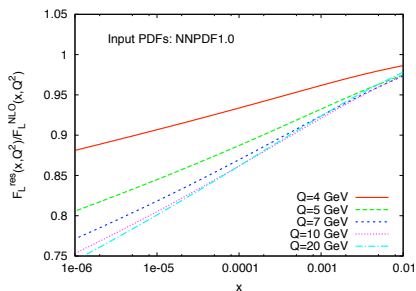


NLL small- $x$  corrections have **similar size but opposite sign** as NNLO

# Resummed K-factors for HERA

Compare **NLL** and **NNLO** K-factors in HERA region

For  $F_L(x, Q^2)$

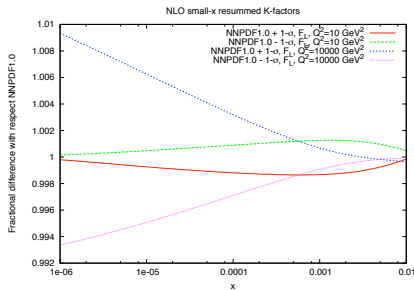
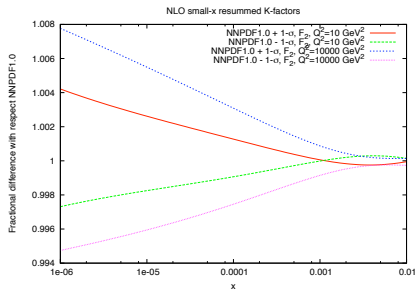


Note again that small- $x$  resummation effects **increase with  $Q^2$**  (as would be obtained from a resummed PDF analysis)

# Resummed K-factors for HERA

Compare **NLL** and **NNLO** K-factors in HERA region

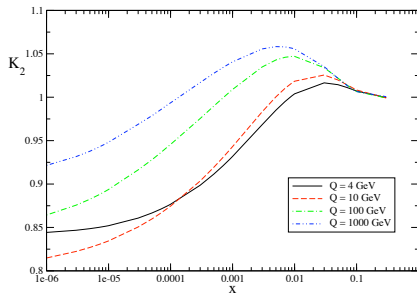
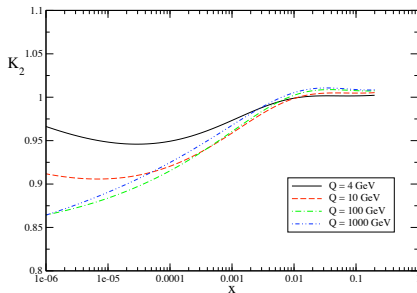
Accuracy of the K-factor approach not limited by choice of input PDF set



K-factors variation  $\leq 1\%$  if the  $1\text{-}\sigma$  NNPDF envelope instead of central PDFs used

# Comparison between ABF and TW

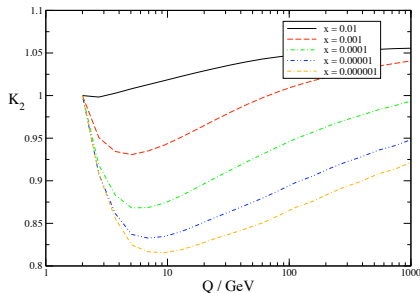
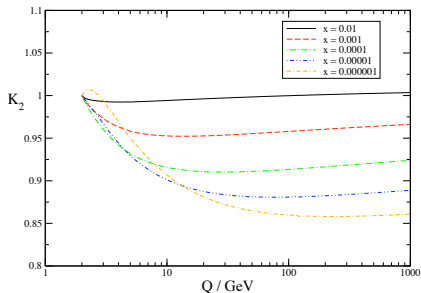
Compare  $F_2^{\text{NLL}}/F_2^{\text{NLO}}$  obtained from toy PDFs at  $Q_0^2 = 4 \text{ GeV}^2$



- Qualitative agreement:  $F_2^{\text{NLL}}$  smaller than  $F_2^{\text{NLO}}$
- Quantitative difference I: TW missing **smooth matching at large- $x$** , while for  $F_2^{\text{NLL}}/F_2^{\text{NLO}} \rightarrow 1$

# Comparison between ABF and TW

Compare  $F_2^{\text{NLL}}/F_2^{\text{NLO}}$  obtained from toy PDFs at  $Q_0^2 = 4 \text{ GeV}^2$

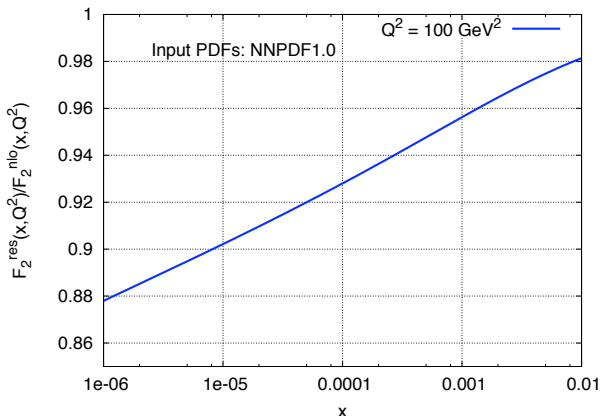


- Qualitative agreement:  $F_2^{\text{NLL}}$  smaller than  $F_2^{\text{NLO}}$
- Quantitative difference II: **ABF resummation stabilized at large- $Q^2$**  (asymptotic freedom) while TW not  $\rightarrow$  Collinear resummation of BFKL kernel missing in TW



# Dependence with matching scale

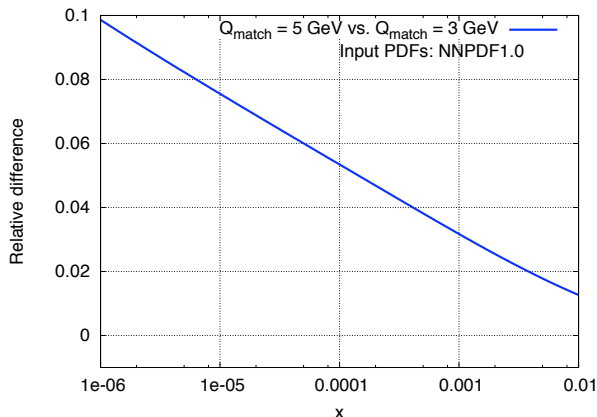
Compare  $Q_{\text{match}} = 3 \text{ GeV}^2$  with  $Q_{\text{match}} = 5 \text{ GeV}^2$





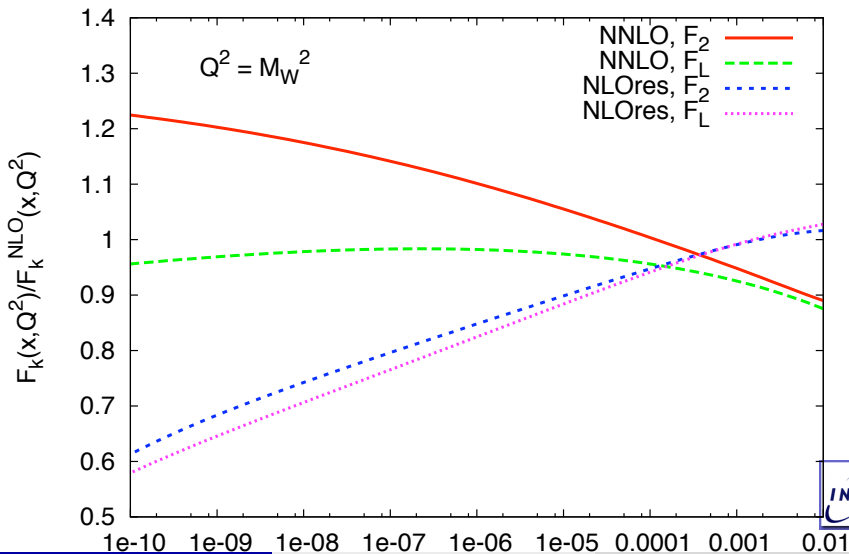
# Dependence with matching scale

Compare  $Q_{\text{match}} = 3 \text{ GeV}^2$  with  $Q_{\text{match}} = 5 \text{ GeV}^2$



Matching scale dominant uncertainty in the K-factor approach to small- $x$  resummation  
Need for a fully small- $x$  resummed global PDF fit

# K-factors for UHE neutrinos

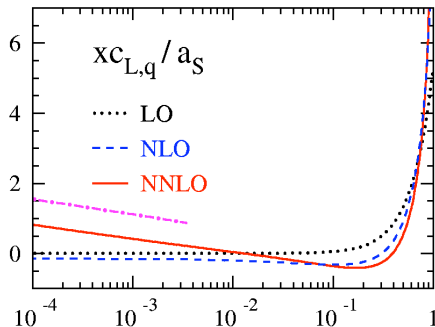


# Why to resum small- $x$ terms?

Instability of perturbation theory (I)

## THE COEFFICIENT FUNCTION $C_L$

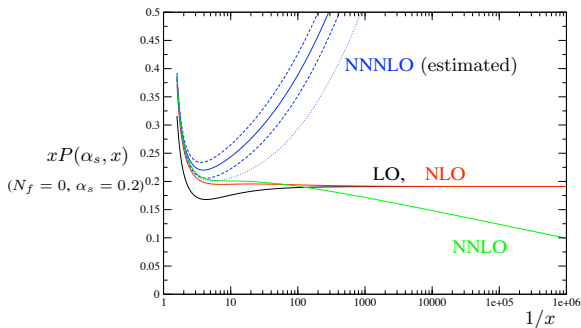
(Moch, Vermaseren, Vogt 2005)



# Why to resum small- $x$ terms?

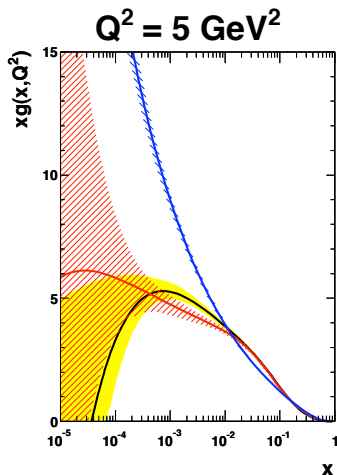
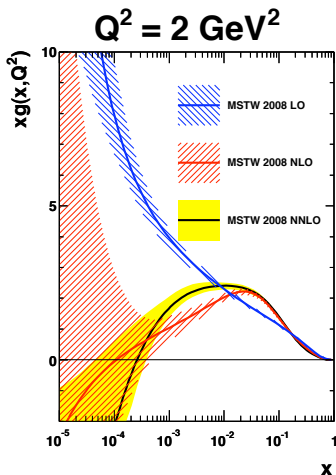
## Instability of perturbation theory (II)

$$xP(\alpha_s, x) \underset{x \rightarrow 0}{\sim} \alpha_s c_1^{(1)} + \alpha_s^2 c_2^{(1)} + \alpha_s^3 \left( c_3^{(2)} \ln x + c_3^{(1)} \right) + \alpha_s^4 \left( c_4^{(4)} \ln^3 x + c_4^{(3)} \ln^2 x + c_4^{(2)} \ln x + c_4^{(1)} \right) + \dots$$



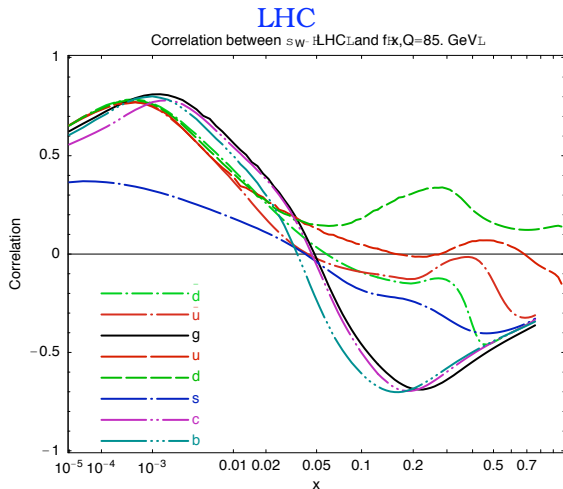
# Why to resum small- $x$ terms?

## Instability of NNLO global PDF determinations



# Why to resum small- $x$ terms?

Small- $x$  PDFs propagate to LHC inclusive observables (CTEQ6.6 study)



# Small-x resummation FAQ

- **Q:** can one resum large small  $x$  corrections to all orders?  
A: yes, both at the leading and subleading level for anomalous dimensions and deep-inelastic and HQ hard coefficients  
BFKL 75-76, FL 98; Ciafaloni, Catani and Hautmann, 91-94
- **Q:** can one combine small  $x$  resummation with standard perturbative evolution?  
A: yes, by performing a suitable double BFKL+GLAP expansion  
Ball, Forte 95, + Altarelli (ABF) 2000
- **Q:** can one obtain a stable perturbative expansion at the resummed level?  
A: yes, provided one exploits suitable physical constraints: momentum conservation & underlying gluon symmetry  
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- **Q:** can one understand the success of NLO perturbation theory despite large small  $x$  terms?  
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- **Q: does the resummation interfere with the choice of factorization scheme?**  
A: no: BFKL resummation is compatible with collinear factorization; resummation can be performed in any scheme  
ABF 01, CCSS 02, ABF 08
- **Q: can one estimate the ambiguities in the resummation and how large are they?**  
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→ LHC Small-x resummed phenomenology around the corner!



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