

Heavy Flavours Working Group: News from Theory

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Theoretical talks

I. General-Mass Variable-Flavor-Number Scheme (GM-VFNS)

- Inclusive photoproduction of D^* mesons at NLO (Hubert Spiesberger)
- Charm hadroproduction and the charm content of the proton (Bernd Kniehl)
- Analysis of heavy-quark electroproduction data in VFNS (Sergey Alekhin)

II. SFs, PDFs and HF Cross-sections

- 3-loop corrections to the moments of $F_2^{(c,b)}(x, Q^2)$ for $Q^2 \gg m_Q^2$ (Sebastian Klein)
- Global PDF analysis in an intermediate-mass scheme (Pavel Nadolsky)
- MC @ NLO for heavy quarks at HERA (Tobias Toll)
- Higher-order soft corrections to $t\bar{t}$ production (Ulrich Langenfeld)
- Heavy flavour effects in the virtual photon SF to NLO in QCD (Tsuneo Uematsu)

III. Studies of (J/ψ , Υ) Data

- Theoretical status of J/ψ production at HERA (Pierre Artoisenet)
- J/ψ polarization from fixed-target to collider energies (Pietro Faccioli)
- Production and polarization of Υ in the K_T -factorization approach (Nikolai Zotov)

IV. Spectroscopy of the states X,Y,Z and the b baryons (Antonio Polosa; Marek Karliner)

$ep \rightarrow D^{*\pm} X$ at low Q^2 in GM-VFNS

General-Mass Variable-Flavor-Number Scheme

H. Spiesberger

for 1-particle inclusive heavy meson production

- The problem:
Conventionally, PDFs and FFs are defined in the $\overline{\text{MS}}$ scheme
 $\overline{\text{MS}}$ scheme is based on a massless calculation
Can not use $\overline{\text{MS}}$ PDFs and FFs in a massive calculation?
- The solution:
Match massless and massive calculations:

$$d\sigma_{\text{sub}} = \lim_{m \rightarrow 0} d\tilde{\sigma}(m) - d\hat{\sigma}_{\overline{\text{MS}}}$$

The **subtracted cross section** (in a massive calculation)

$$d\hat{\sigma}(m) = d\tilde{\sigma}(m) - d\sigma_{\text{sub}}$$

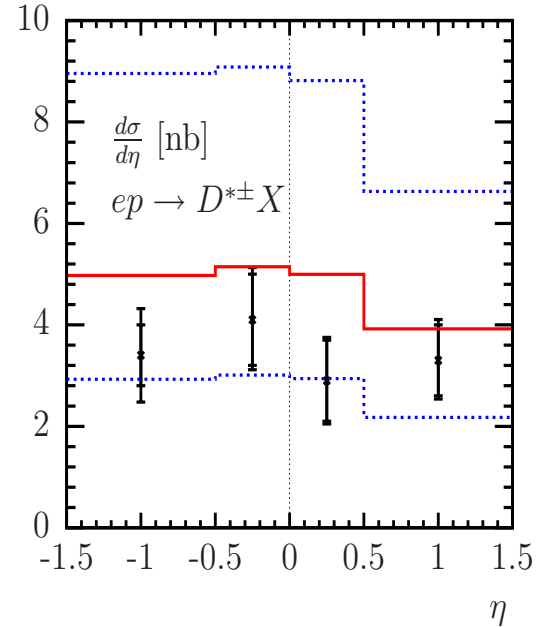
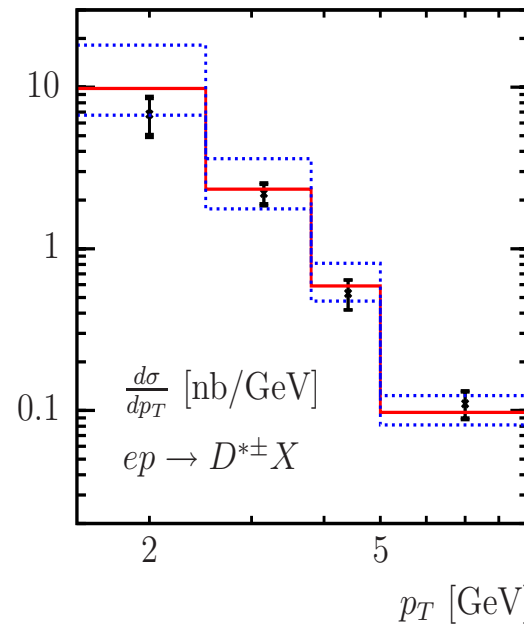
can be used with $\overline{\text{MS}}$ parton distribution and fragmentation functions

- The **GM-VFNS** (general-mass variable flavor number scheme) combines: **resummation of large logarithms** $\log(\mu/m)$ in PDFs and FFs with **finite mass terms** in the hard cross section

Work in collaboration with B. A. Kniehl, G. Kramer, I. Schienbein

$ep \rightarrow D^{*\pm} X$ at low Q^2 in GM-VFNS

$e^\pm p \rightarrow D^* + X$ at low Q^2 : $0.05 < Q^2 < 0.7 \text{ GeV}^2$



scale variations by factor 2 around $\mu_{R,F} = \sqrt{p_T^2 + m^2}$, data: [ZEUS PLB649](#)

More applications available for: $\gamma + \gamma \rightarrow D^{*\pm} + X$: [EPJC22](#), [EPJC28](#);
 $\gamma^* + p \rightarrow D^{*\pm} + X$: [EPJC38](#), [arXiv:0902.3166](#); $p + \bar{p} \rightarrow (D^0, D^{*\pm}, D^\pm, D_s^\pm, \Lambda_c^\pm) + X$:
[PRD71](#), [PRL96](#), [arXiv:0901.4130](#); $p + \bar{p} \rightarrow B + X$: [PRD77](#)

Open charm hadroproduction (Kniehl)

OUR THEORETICAL BASIS FOR $p\bar{p} \rightarrow D^*X$

Factorization Formula:

$$d\sigma(p\bar{p} \rightarrow D^*X) = \sum_{i,j,k} \int dx_1 dx_2 dz f_i^p(x_1) f_j^{\bar{p}}(x_2) \times \\ d\hat{\sigma}(ij \rightarrow kX) D_k^{D^*}(z) + \mathcal{O}(\alpha_s^{n+1}, (\frac{\Lambda}{Q})^p)$$

Q : hard scale, $p = 1, 2$

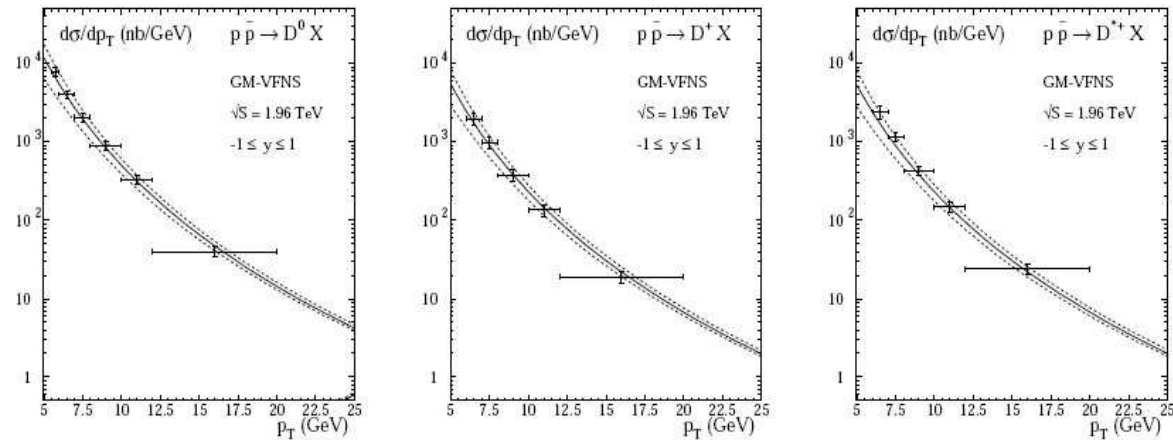
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- $d\hat{\sigma}(\mu_F, \mu'_F, \alpha_s(\mu_R), \frac{m_h}{p_T})$: hard scattering cross sections free of long-distance physics $\rightarrow m_h$ kept
 - PDFs $f_i^p(x_1, \mu_F), f_j^{\bar{p}}(x_2, \mu_F)$: $i, j = g, q, c$ [$q = u, d, s$]
 - FFs $D_k^{D^*}(z, \mu'_F)$: $k = g, q, c$

\Rightarrow need short distance coefficients including heavy quark masses

[1] J. Collins, 'Hard-scattering factorization with heavy quarks: A general treatment', PRD58(1998)094002

Open charm hadroproduction (Kniehl)

HADROPRODUCTION OF D^0, D^+, D^{*+}, D_S^+ GM-VFNS RESULTS W/ KKKSc FFs [1]



- $d\sigma/dp_T$ [nb/GeV] $|y| \leq 1$ prompt charm
- Uncertainty band: $1/2 \leq \mu_R/m_T, \mu_F/m_T \leq 2$ ($m_T = \sqrt{p_T^2 + m_c^2}$)
- CDF data from run II [2]
- GM-VFNS describes data within errors

[1] B.K.,Kramer,Schienbein,Spiesberger, arXiv:0901.4130[hep-ph], PRD(to appear)

[2] Acosta et al., PRL91(2003)241804

HF contr. in $F_2(x, Q^2)$ for $Q^2 \gg m_Q^2$ (Klein)

Moments of the 3-Loop Corrections to the Heavy Flavor Contribution to $F_2(x, Q^2)$ for $Q^2 \gg m^2$

1

Sebastian Klein (DESY) in collaboration with I. Bierenbaum and J. Blümlein

[0904.3536 [hep-ph]]

- QCD precision analyses require the description of the heavy quark contributions to 3-loop order to line up with the accuracy reached in the massless case.
 \Rightarrow QCD analysis and determination of $\alpha_s(M_Z^2)$ from DIS data: $\delta\alpha_s/\alpha_s < 1\%$.

- Factorization of heavy flavor Wilson coefficients $H_{j,i}$ for $Q^2 \gg m^2$ [F_2 : $Q^2/m^2 \geq 10$]

$$H_{(2,L),i}^{S,NS}\left(\frac{Q^2}{\mu^2}, \frac{m^2}{\mu^2}\right) = \underbrace{A_{ki}^{S,NS}\left(\frac{m^2}{\mu^2}\right)}_{\text{massive OMEs}} \otimes \underbrace{C_{(2,L),k}^{S,NS}\left(\frac{Q^2}{\mu^2}\right)}_{\text{light Wilson coefficients}}.$$

[van Neerven, Smith, et. al. 1996]

- Analytic results are known for $Q^2 \gg m^2$ at NLO for $F_2(x, Q^2)$. [$F_L(x, Q^2)$ at NNLO]
- light Wilson coefficients known at NNLO [Vermaseren, Vogt, Moch, 2005.]
 \Rightarrow massive OMEs needed at $O(a_s^3)$.
- This calculation: We obtain fixed Mellin moments of all massive OMEs at $O(a_s^3)$

3-loop corrections to HF contr. in $F_2(x, Q^2)$ for $Q^2 \gg m_Q^2$

- The calculation of fixed moments of the massive operator matrix elements at $O(a_s^3)$ has been finished up to $N = 10, 12, 14$ (≈ 250 days of computer time)

$$A_{Qq}^{(3),PS} : (2, 4, \dots, 12); \quad A_{qq,Q}^{(3),PS}, A_{gg,Q}^{(3)} : (2, 4, \dots, 14);$$
$$A_{qq,Q}^{(3),NS\pm} : (2, 3, \dots, 14); \quad A_{Q(q)g}^{(3)}, A_{gg,Q}^{(3)} : (2, 4, \dots, 10);$$

- Massive OMEs occur as well as transition functions to define a variable flavor number scheme starting from a fixed flavor number scheme.

- The complete OMEs but the constant terms are derived in terms of renormalization constants known for all N.

- The calculation provides first independent checks on the 3-loop anomalous dimensions $\gamma_{qq}^{(2)}$, $\gamma_{qq}^{(2),PS}$ and on respective color projections of $\gamma_{qq}^{(2),NS\pm}$, $\gamma_{gg}^{(2)}$ and $\gamma_{gq}^{(2)}$.

\Rightarrow Agreement for the terms $\propto T_F$ with [Larin, Ritbergen, Vermaseren, 1994;

Larin, Nogueira, Ritbergen, Vermaseren, 1997; Retey, Vermaseren, 2001; Vogt, Moch, Vermaseren, 2004.]

- We corrected some left outs even at NLO in the literature.
- First phenomenological parameterization to come up soon.

Heavy flavour effects in the virtual photon SF to NLO (Uematsu)

The heavy quark mass effects in parton picture and OPE

$$F_2^\gamma(x, Q^2, P^2) = \vec{q}^\gamma(y, Q^2, P^2, m^2) \otimes \vec{C} \left(\frac{x}{y}, \frac{\bar{m}^2}{Q^2}, \bar{g}(Q^2) \right)$$

Photon structure function

PDF

Coefficient function

←
mass dependence

Parton interpretation of twist-2 operators \vec{O}_n

$$\int_0^1 dx x^{n-1} \vec{q}^\gamma(x, Q^2, P^2, m^2)$$

$$= \vec{A}_n \left(1, \frac{\bar{m}^2(P^2)}{P^2}, \bar{g}(P^2) \right) T \exp \left[\int_{\bar{g}(Q^2)}^{\bar{g}(P^2)} dg \frac{\gamma_n(g, \alpha)}{\beta(g)} \right]$$

no mass dependence

↙

where

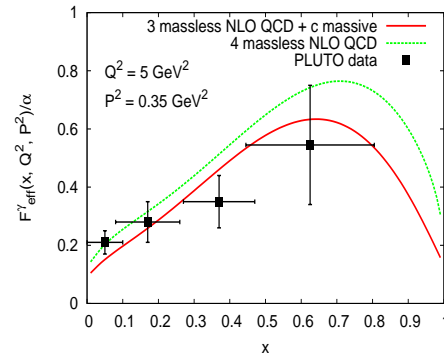
$$\langle \gamma(P^2) | \vec{O}_n(\mu^2) | \gamma(P^2) \rangle = \vec{A}_n \left(\frac{P^2}{\mu^2}, \frac{\bar{m}^2(\mu^2)}{\mu^2}, \bar{g}(\mu^2) \right)$$

Perturbatively calculable!

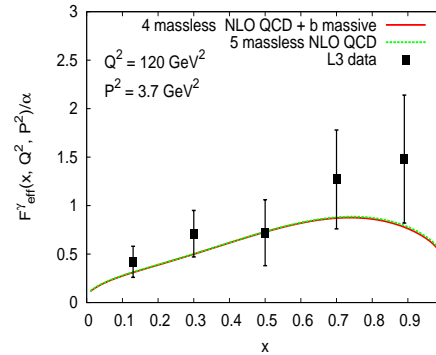
←
mass dependence

Heavy flavour effects in the virtual photon SF to NLO (Uematsu)

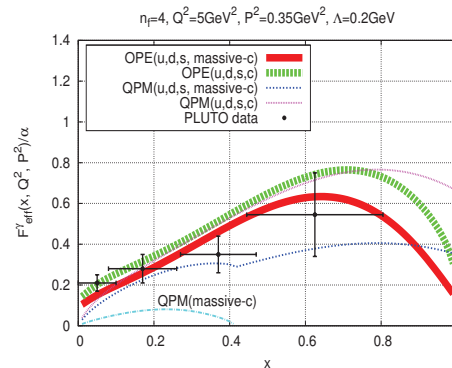
Heavy flavour effects vs. experimental data for F_{eff}^γ



QCD prediction vs. PLUTO data



QCD prediction vs. L3 data



QCD prediction with PM Box contribution

In the massive quark limit
 $\Lambda_{\text{QCD}}^2 \ll P^2 \ll m^2 \ll Q^2$

Heavy quark mass
 threshold effects
 illustrated by PM Box

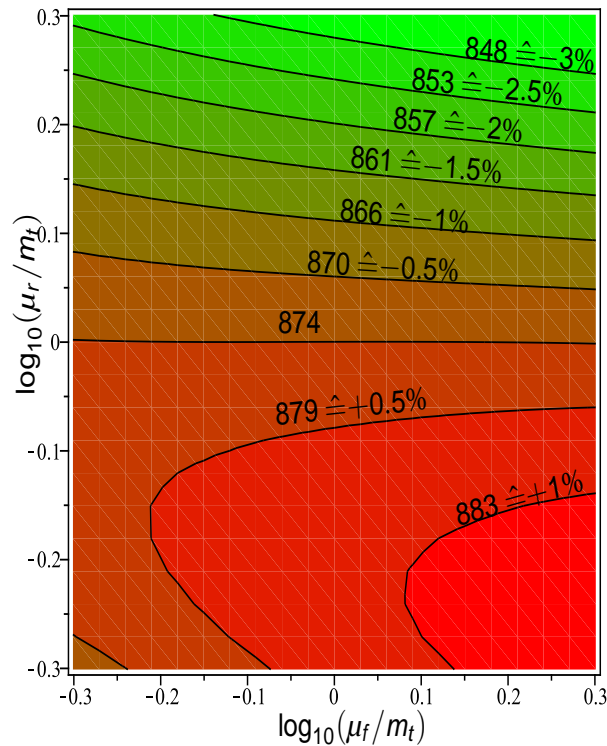


$$x_{\text{max}} = \frac{1}{1 + \frac{4m^2}{Q^2}}$$

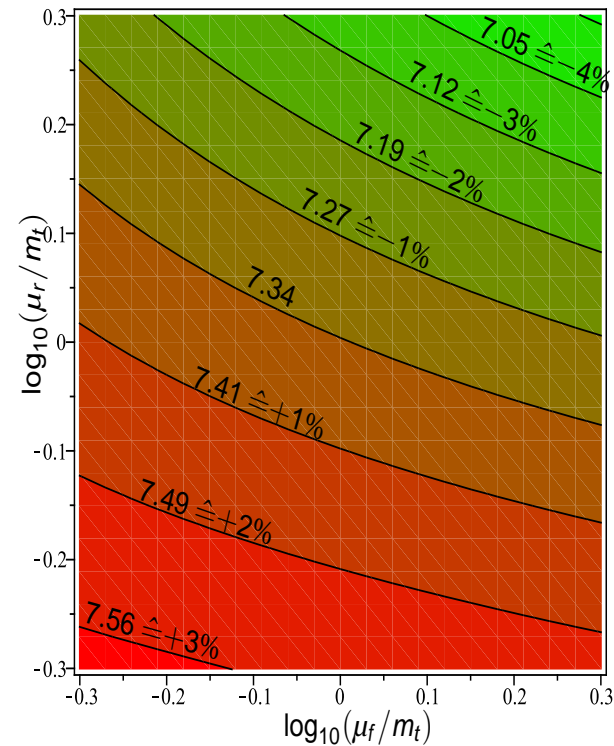
$t\bar{t}$ production cross sections (Langenfeld)

Scale dependence of the $t\bar{t}$ production cross section

Contour lines of the total hadronic cross section for top pair production in the $\mu_f - \mu_r$ - plane in pb (\log_{10} - scale)



LHC



Tevatron



1/2

$t\bar{t}$ production cross sections (Langenfeld)

Cross section Predictions

- Total cross section prediction for $m_t = 173 \text{ GeV}$ and PDF set Cteq6.6:

$$\sigma(p\bar{p} \rightarrow t\bar{t}) = 7.34_{-0.38}^{+0.23} \text{ pb @ Tevatron}$$

$$\sigma(pp \rightarrow t\bar{t}) = 874_{-33}^{+14} \text{ pb @ LHC 14 TeV}$$

- Fit functions to the hadronic cross section for different scales ($\mu = 1/2m_t, m_t, 2m_t$) for the two colliders Tevatron and LHC (10/14 TeV) and for two different PDF sets (Cteq6.6, MSTW 2008 NNLO), w/o PDF errors.

$$\text{Ansatz: } \sigma(m_t) = a + bx + cx^2 + dx^3 + ex^4 + fx^5 + gx^6, \\ x = (m_t/\text{GeV} - 173).$$

- Total cross section with combined theoretical and PDF error:

$$\sigma(p\bar{p} \rightarrow t\bar{t}) = 7.34_{-0.76}^{+0.66} \text{ pb @ Tevatron}$$

$$\sigma(pp \rightarrow t\bar{t}) = 874_{-50}^{+26} \text{ pb @ LHC 14 TeV}$$

Theoretical status of J/ψ production at HERA (Artoisenet)

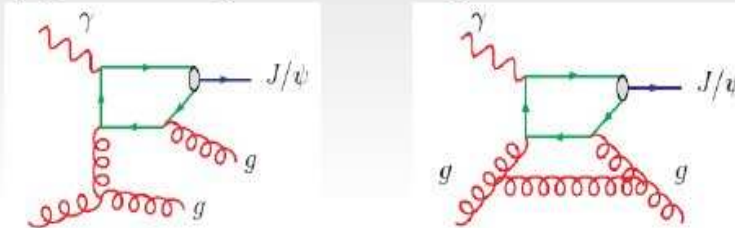
Colour-singlet yield at NLO

M. Kramer (1995)

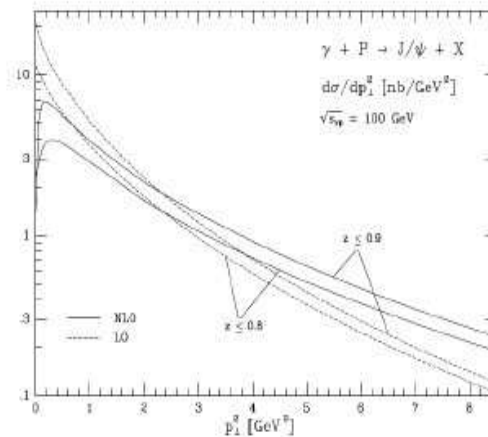
P. A., J. Campbell, F. Maltoni, F. Tramontano (2009)

C.-H. Chang, R. Li, J.-X. Wang (2009)

- Typical Feynman diagrams:



- Differential cross section:



substantial enhancement at
high transverse momentum
(new kinematic contributions)

Kramer, Nucl.Phys.B459 :3-50,1996.

Theoretical status of J/ψ production at HERA (Artoisenet)

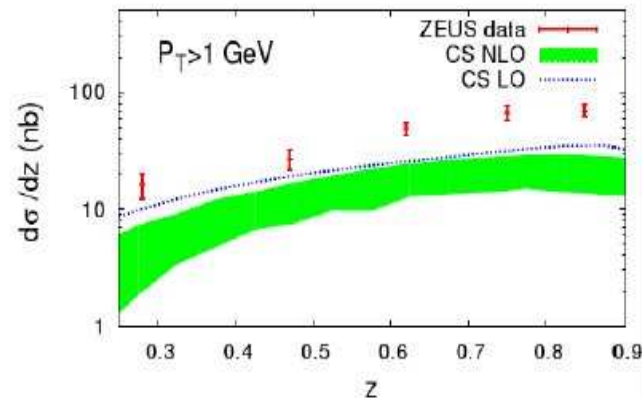
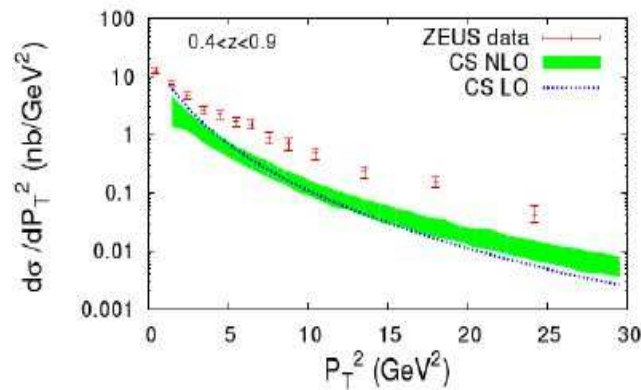
Comparison with the data

- Colour singlet yield at NLO undershoots the data

mass uncertainty: $1.4 \text{ GeV} < m_c < 1.6 \text{ GeV}$

scale uncertainty:

$$\mu_0 = 4m_c, \quad 0.5\mu_0 < \mu_r, \mu_f < 2\mu_0, \quad 0.5 < \frac{\mu_r}{\mu_f} < 2$$



Theoretical status of J/ψ production at HERA (Artoisenet)

J/Ψ Polarization

- J/Ψ polarization deduced from the angular distribution of the produced leptons
- spin quantization axis: $\hat{z} = -\frac{\mathbf{P}_N}{|\mathbf{P}_N|}$ (target frame)
- angular distr. of the leptons in the J/Ψ rest frame

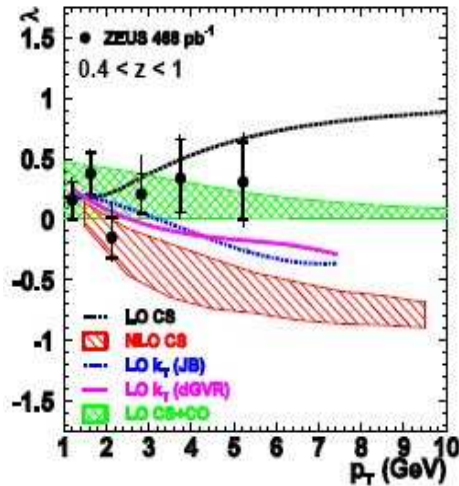
$$\frac{d\sigma}{d\Omega du} \propto 1 + \lambda(y) \cos^2 \theta + \mu(y) \sin 2\theta \cos \phi + \frac{\nu(y)}{2} \sin^2 \theta \cos 2\phi$$

with respect to the spin quantization axis, and

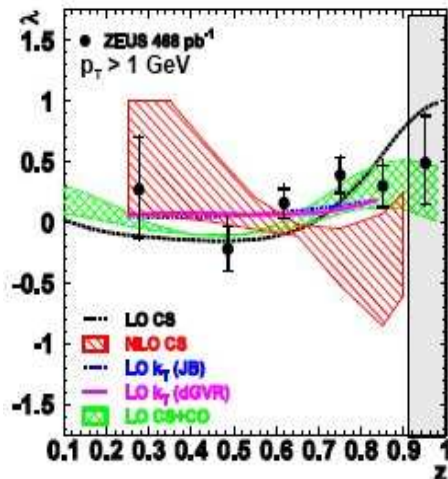
$$\lambda = \frac{\rho_{11} - \rho_{00}}{\rho_{11} + \rho_{00}}, \quad \mu = \frac{\sqrt{2} \text{Re} \rho_{10}}{\rho_{11} + \rho_{00}}, \quad \nu = \frac{2\rho_{1,-1}}{\rho_{11} + \rho_{00}}$$

Theory vs. Expt: J/ψ polarization at HERA (Bertolin)

J/ψ helicity at HERA



- LO CS and NLO CS predictions have opposite sign ... we initially thought NLO corrections would be small ...
- LO k_T CS has the same sign of NLO, parton transverse momentum, k_T , mimics NLO terms
- LO CS+CO is flat
- data are consistent with being flat in the probed p_T range
- proton dissociative background mostly at low p_T
- analysis redone for $z < 0.9$, effects in the sys. errors

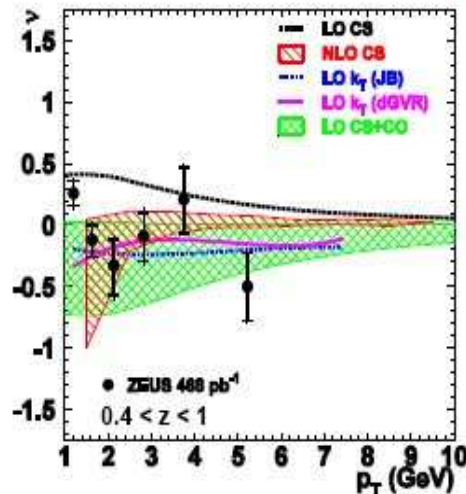


- LO CS describe the data well
- NLO CS has large uncertainties ... $p_T > 1$ GeV may be not enough ...
- LO k_T CS not too different from LO
- LO CS+CO is pretty much the same as LO CS
- proton dissociative is at the 60 – 70 % level for $0.9 < z < 1$, $\ll 5$ % elsewhere

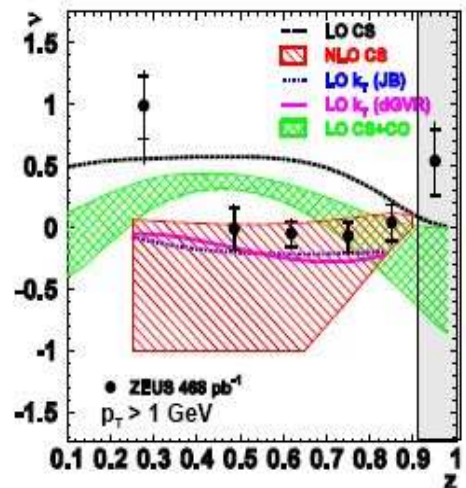
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Theory vs. Expt: J/ψ polarization at HERA (Bertolin)

J/ψ helicity at HERA (cont.)



- LO CS is positive ... all other predictions are negative ... and in better agreement with the data
- LO k_T CS is pretty much as NLO CS
- LO CS+CO is flat
- data are consistent with being flat in the probed p_T range
- proton dissociative background mostly at low p_T
- analysis redone for z < 0.9, effects in the sys. errors



- LO CS does not describe the data, positive
- NLO CS has large uncertainties ... negative ... p_T > 1 GeV may be not enough ...
- LO k_T CS fine ... except at low z
- LO CS+CO does not describe the data, positive
- proton dissociative is at the 60 – 70 % level for 0.9 < z < 1, << 5 % elsewhere

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Conclusion

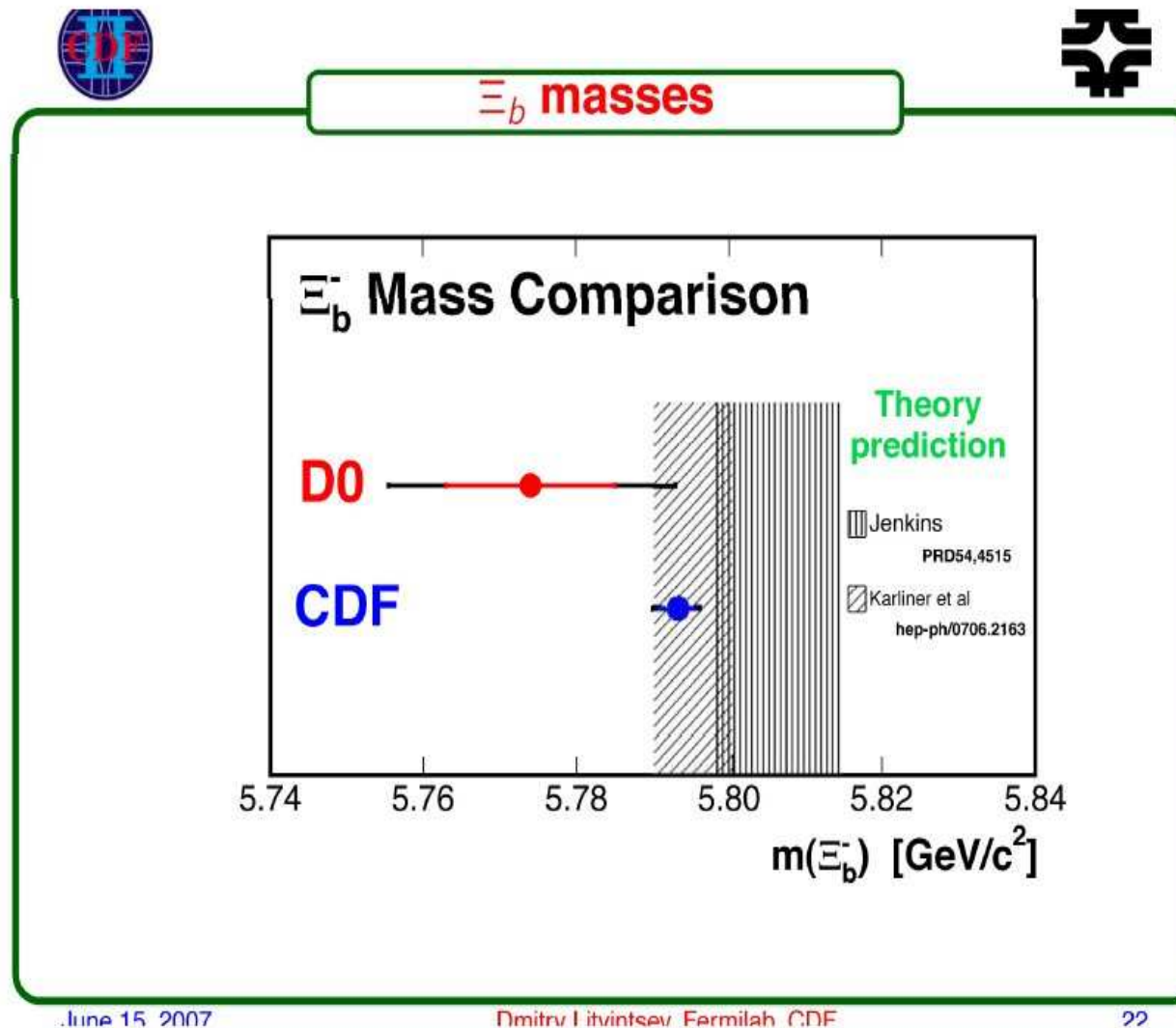
- New results on the polarization of the J/ψ in photoproduction (colour-singlet at NLO)
- Colour-singlet prediction alone does not describe all features of the data collected at HERA
- NLO computation of the polarization associated to the colour-octet P-wave transition may solve the current discrepancy

Quark Model and spectroscopy of the b -baryons (Karliner)

Summary

- Constituent quark model with color HF interaction gives highly accurate predictions for heavy baryon masses
- a challenge for theory: derivation from QCD
- constituent quark masses depend on the spectator quarks
- $M_{\Sigma_b} - M_{\Lambda_b} = 194 \text{ MeV}$ vs 192 in EXP (CDF)
- $M(\Sigma_b^*) - M(\Sigma_b) = 22 \text{ MeV}$ vs 21 MeV in EXP (CDF)
- $\mu_{\Lambda_c} = 0.43 \text{ n.m.}$ $\mu_{\Lambda_b} = -0.067 \text{ n.m.}$
- meson-baryon effective supersymmetry
- meson/baryon HF splitting confirms Cornell potential
- Ξ_b mass prediction: 5795_{-5} MeV vs $5793 \pm 2.4 \pm 1.7 \text{ MeV}$
- puzzle in $Y(5S)$ decays: $\bar{b}b u \bar{d}$ candidates?

Quark Model and spectroscopy of the b -baryons (Karlner)



X, Y, Z, ... spectroscopy (Polosa)

X, Y, Z, ... Spectroscopy

... is exotic for sure. The first of this series of narrow resonances, the X(3872), is not a charmonium and challenges both its possible competitive explanations: molecule and tetraquark

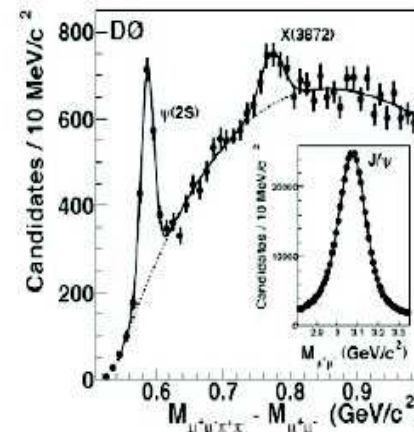
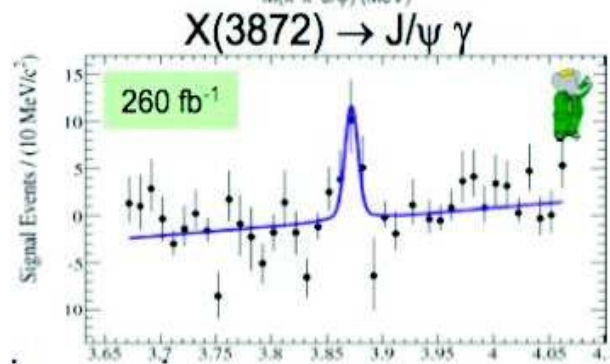
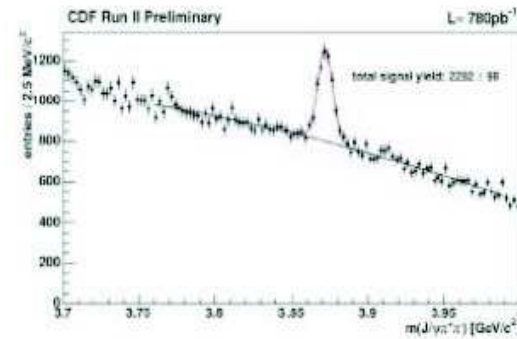
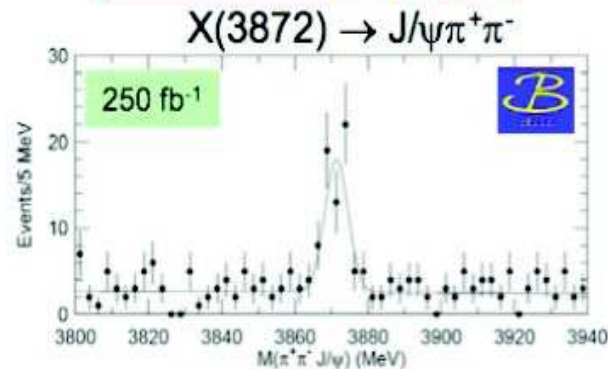
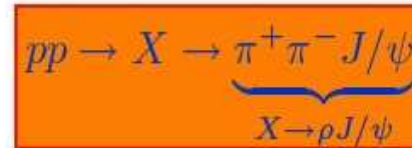
The tetraquark picture has clear predictions: one of them, the existence of charged particles decaying into charmonium + pions, is still debated by experiments.

The molecule picture seems to be at odds with the large prompt production cross sections at hadron colliders.

Other options, like the hybrid, are still possible for particles like the Y(4260).

X, Y, Z, ... spectroscopy (Polosa)

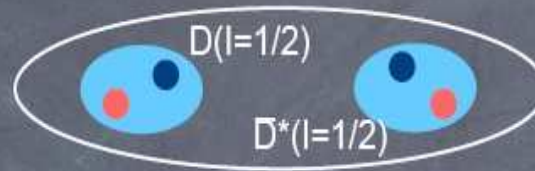
Experimental Facts



$$X \rightarrow \gamma J/\psi \mapsto C = +1 \quad \text{and} \quad X \rightarrow \rho^0 J/\psi \rightarrow (\pi^+ \pi^-)_S J/\psi \mapsto P = 1$$

X, Y, Z, ... spectroscopy (Polosa)

I: The molecule interpretation



$$\mathcal{E}_0 \sim -0.25 \pm 0.40 \text{ MeV}$$

$$\frac{\hbar^2}{\mu r_0^2} - \frac{g^2}{4\pi} \frac{e^{-\frac{m_\pi r_0}{\hbar}}}{r_0} = \mathcal{E}_0$$

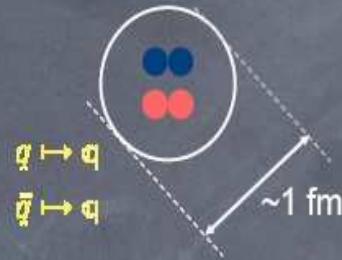
$$r_0 \sim 8 \text{ fm}$$

This is a large (with respect to strong interaction characteristic range) loosely bound state: can it be formed prompt in pp collisions at high energy?

CDF measures a prompt xsect of about 3 nb - a calculation of prompt production cross section (Bignami, Grinstein, Piccinini, Polosa w.i.p.) points at a much smaller value.

X, Y, Z, ... spectroscopy (Polosa)

II: The tetraquark interpretation



$$\mathfrak{q} = [q \uparrow q \downarrow]_{\mathfrak{s}_c, \mathfrak{s}_s} \text{ or, more precisely, } \mathfrak{q}_{i\alpha} = \epsilon_{ijk} \epsilon_{\alpha\beta\gamma} \bar{q}_C^{j\beta} \gamma_5 q^{k\gamma}$$

-This is a compact, tightly bound object.

Drawbacks: Many states are predicted. X^+ and X^- yet to be found. (Some indication by Belle of charged particles decaying into charmonium + charged pions).

-Two neutral X , with a difference in mass of few MeV and with different decay modes are expected in order to explain the observed isospin violation in $X \rightarrow \psi\rho$, $\psi\omega$ (which seem to occur at the same rate).

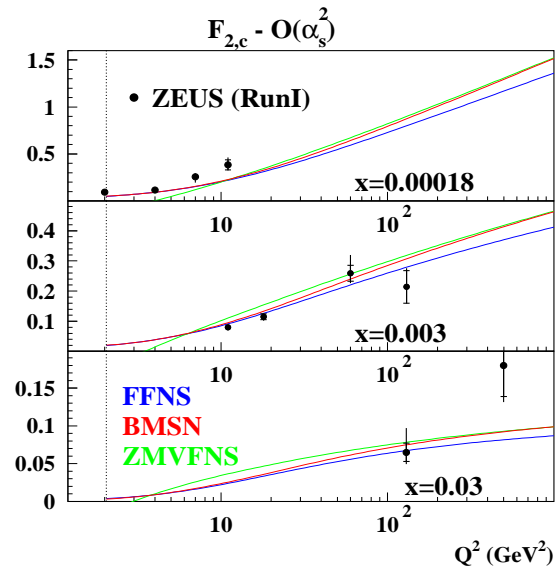
-A doubly charged state $W^{++}=[cu][db\ sb]$ could exist (very unlikely for a molecule: at 10 fm electromagnetic repulsion starts to compete with strong interaction)

Backups

- VFNS and heavy quark electroproduction data
Sergey Alekhin
- MC @ NLO for heavy quarks at HERA
Tobias Toll
- Global PDF analysis in an intermediate-mass scheme
Pavel Nadolsky
- J/ψ polarization from fixed-target to collider energies
Pietro Faccioli
- Production and polarization of Upsilon mesons in the K_T -factorization approach
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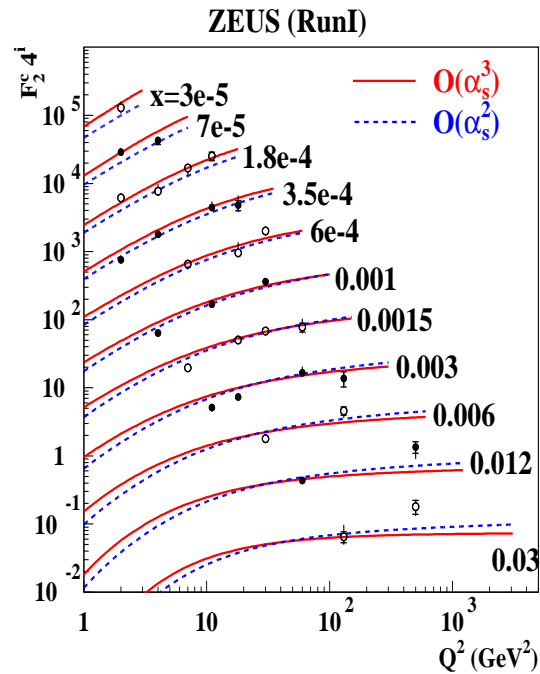
VFNS and heavy quark electroproduction data (Alekhin)

(S. Alekhin, J. Blümlein, S. Klein, S. Moch)



- The Buza-Matiounine-Smith-van Neerven (BMSN) prescription for the variable-flavour-number scheme (VFNS), is not too far from the fixed-flavour-number scheme (FFNS) for the realistic HERA kinematics; it seems to be the case for any smooth matching.
- The remaining discrepancies with the data cannot be cured by the smooth VFNS.

VFNS and heavy quark electroproduction data (Alekhin)



- The FFNS with account of the partial $O(\alpha_s^3)$ corrections provides a good description of the HERA data at small/moderate Q^2 .
- At large Q^2 it undershoots the data due to negative contribution from $c_{2,g}^{(2,1)}$, the missing contribution from $c_{2,g}^{(2,0)}$ at large η must be positive in order to improve the agreement.

(c.f. S.Klein's talk)

Conclusions

- MC@NLO constructed for HERA
- At parton level it behaves as expected
- Where pQCD is expected to be important MC@NLO is doing well

χ^2/ndf combined for all beauty plots shown :

	MC@NLO	FMNR	HERWIG (rescaled)	ndf
χ^2	0.86	4.51	1.31	57

Charm Issues

- NLO calculation unable to describe Charm data
- Shapes are well described
- One analysis too large, one too small
- Charm mass is low $m_t=1.5$ GeV, also Q^2 is low
- Probing the edge of pQCD
- This leads to large sensitivity to the scales:

χ^2/ndf H1 2006 D^* + other jet

Scale	m_t	$2 \cdot m_t$	$4 \cdot m_t$	ndf
$p_t(D^*)$	25.87	20.22	16.34	4
$\eta(D^*)$	3.97	2.09	1.47	5
$\Delta\eta(D^*, jet)$	2.99	1.64	1.07	6
$p_t(jet)$	4.78	2.29	1.53	4
$\eta(jet)$	3.48	1.56	1.12	5
$\Delta\phi(D^*, jet)$	4.20	2.79	2.20	6
Total	6.77	4.50	3.47	30

By increasing the scale, the description improves by > 3 units in χ^2/ndf

Global PDF analysis in an intermediate-mass scheme (Nadolsky)

Rescaling variables in heavy-flavor DIS

P. Nadolsky, W.-K. Tung, arXiv:0903:2667

- **General-mass scheme:** leading heavy-quark mass effects are captured by a **rescaling variable** χ in HQ PDF's:

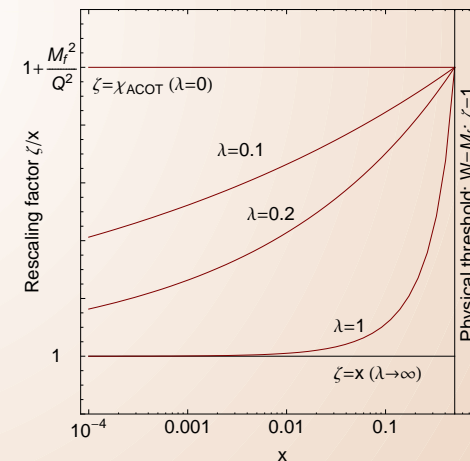
$$x = \chi / \left(1 + M_f^2 / Q^2 \right), \text{ with } M_f^2 = 4m_c^2 \text{ (} m_c^2 \text{) in NC (CC) DIS}$$

Barnett, Haber, Soper; Tung, Kretzer, Schmidt

- **Zero-mass scheme at NLO:** realistic threshold behavior is reproduced by a **generalized rescaling variable** $\zeta(\lambda)$:

$$x = \zeta / \left(1 + \zeta^\lambda M_f^2 / Q^2 \right), \text{ with } 0 \leq \lambda \lesssim 1$$

- ▶ $\lambda = 0$: $\zeta = \chi$ (the ACOT- χ variable)
- ▶ $\lambda \gtrsim 1$: $\zeta \approx x$ (no rescaling)
- ▶ $\zeta \approx x$ for $Q^2 \gg M_f^2$

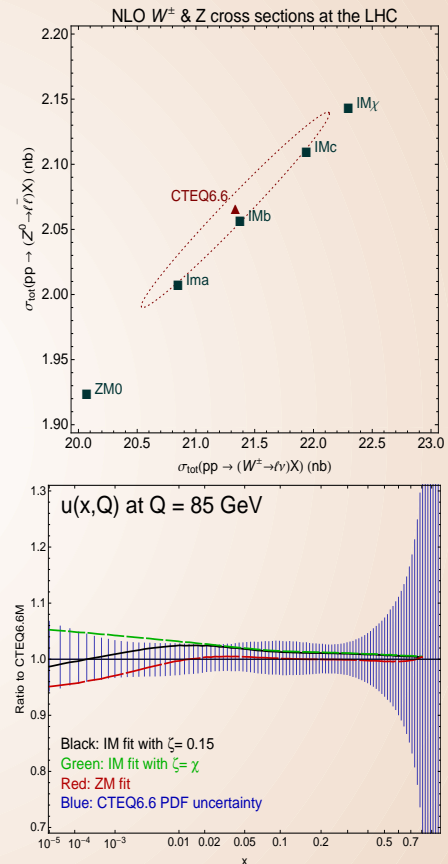


Global PDF analysis in an intermediate-mass scheme (Nadolsky)

Intermediate-mass scheme

PDF fits with ZM Wilson coefficients evaluated at $\zeta(x, M_f/Q, \lambda)$ closely reproduce GM fits for some λ values

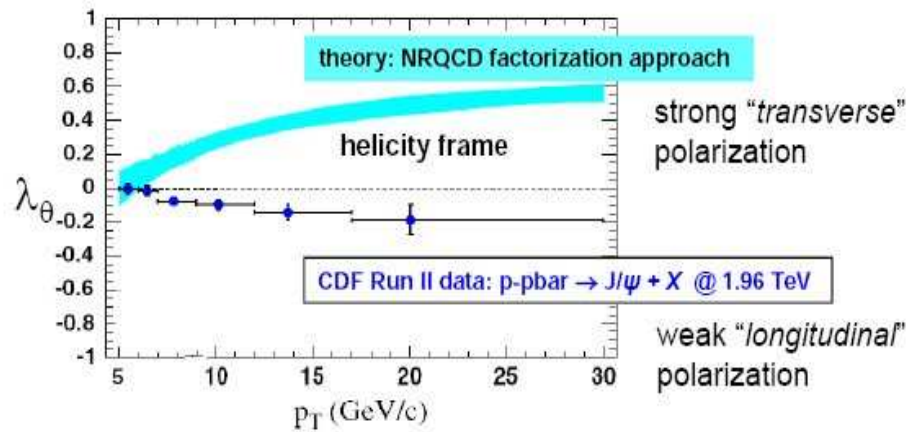
This **intermediate-mass** formulation improves the NLO ZM scheme and approximates (more fundamental) GM predictions by simple means



J/ψ polarization (Faccioli)

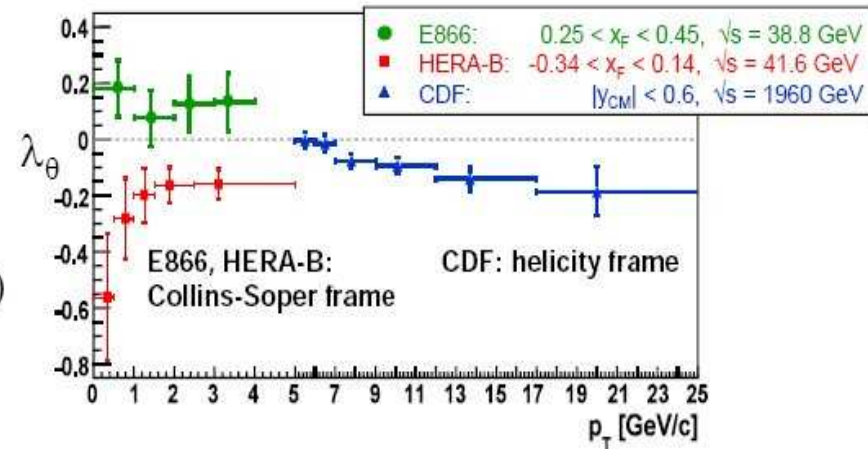
The “ J/ψ polarization puzzle”

theory
vs
experiment:



experiment
vs
experiment:

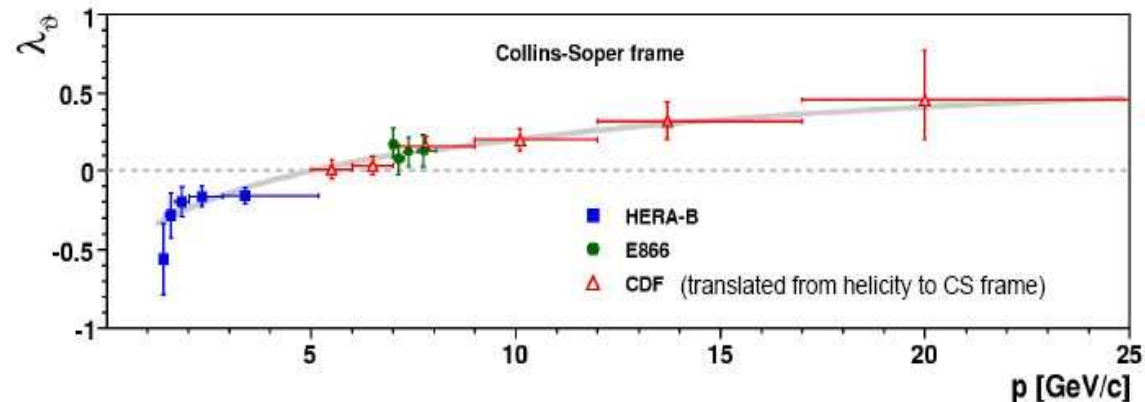
(we address *this*
part of the problem)



The “ J/ψ polarization puzzle”

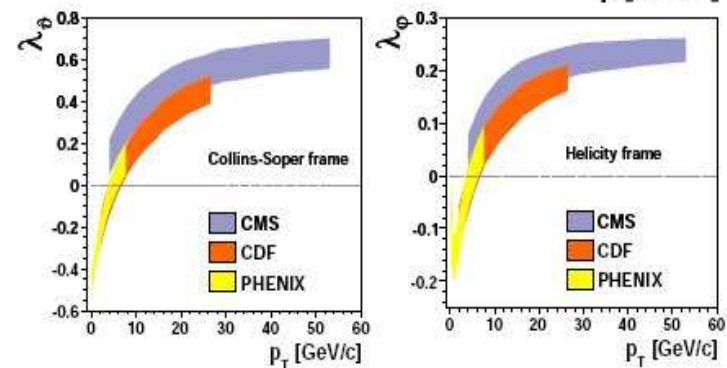
Towards a clearer view of the experimental picture

The polarization *depends* on the reference frame. In the Collins-Soper frame (assumed to be “the natural frame”), the E866, HERA-B and CDF J/ψ polarization results nicely overlap as a function of the J/ψ cms total momentum, showing that the polarization changes from longitudinal at small momentum to transverse at high momentum:



According to this scenario, at very high- p_T the experiments should observe:

- large transverse polarization in the CS frame
- large azimuthal anisotropy in the helicity frame



Υ polarization in K_T factorization

- Goal: Theoretical predictions on the polarization of the Υ at the Tevatron & LHC
- $p\bar{p}$ and pp X-sections calculated as a convolution of u.g.d. $\mathcal{F}_g(x, k_T^2, \mu^2)$ and partonic X-section $\hat{\sigma}$

$$\sigma_{pp} = \int \mathcal{F}_g(x_1, k_{1T}^2, \mu^2) \cdot \mathcal{F}_g(x_2, k_{2T}^2, \mu^2) \hat{\sigma}_{gg}(x_1, x_2, k_{1T}^2, k_{2T}^2) dx_1 dx_2 dk_{1T}^2 dk_{2T}^2$$

- BFKL prescription for the off-shell gluon spin density matrix:

$$\overline{\epsilon_g^\mu \epsilon_g^{*\nu}} = p_p^\mu p_p^\nu x_g^2 / |k_T|^2 = k_T^\mu k_T^\nu / |k_T|^2$$

- Choice of u.g.d.:

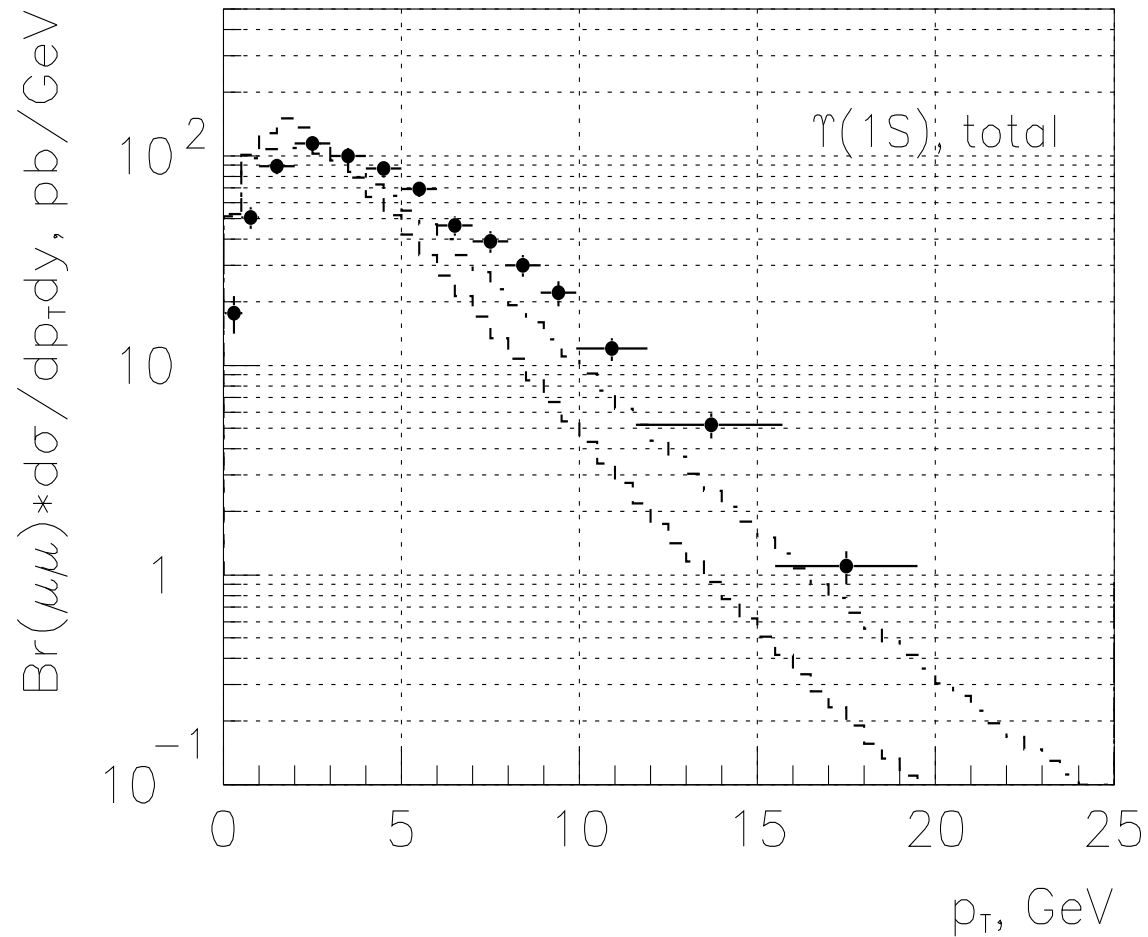
- dGRV, derived from the collinear gluon density $G(x, \mu^2)$, differentiating it w.r.t. μ^2 and setting $\mu^2 = k_T^2$
- JB parametrization [J. Blümlein; J. Phys. G 19 (1993) 1623], solution of leading order BFKL equation in the double logarithmic approximation

- Production Mechanisms:

- Direct gluon-gluon fusion: $g + g \rightarrow \Upsilon + g$
- Production of the P-wave states $g + g \rightarrow \chi_b \rightarrow \Upsilon + \gamma$
- Polarization state characterized by the spin alignment parameter α :

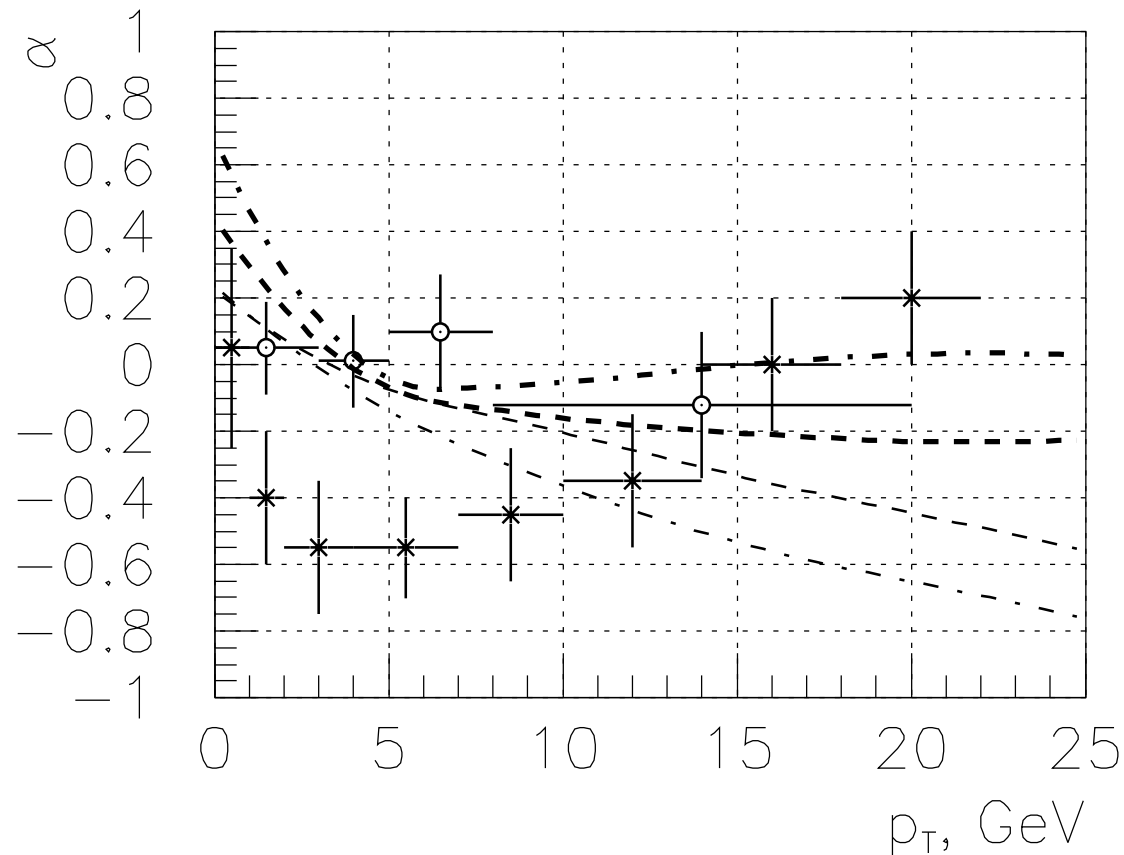
$$\alpha(\mathcal{P}) = (d\sigma/d\mathcal{P} - 3d\sigma_L/d\mathcal{P}) / (d\sigma/d\mathcal{P} + d\sigma_L/d\mathcal{P})$$

Υ X-section at the Tevatron & Theory (Zotov)



- Dashed-dotted (upper) histograms: JB Parametrization
- Dashed (lower) dGRV parametrization

Spin alignment parameter $\alpha(\Upsilon)$ vs. p_T (Zotov)



- Upper data points (D0 Run I: Direct Υ and $\chi_b \rightarrow \Upsilon + \gamma$) & Theoretical curves (upper: JB, lower: dGRV)
- Lower data points (D0 Run II: Direct Υ) & Theoretical curves (upper: dGRV, lower: JB)
- Conclusion: Need more theoretical studies