

# Implications of Direct Detection Blind Spots at the LHC

Peisi Huang

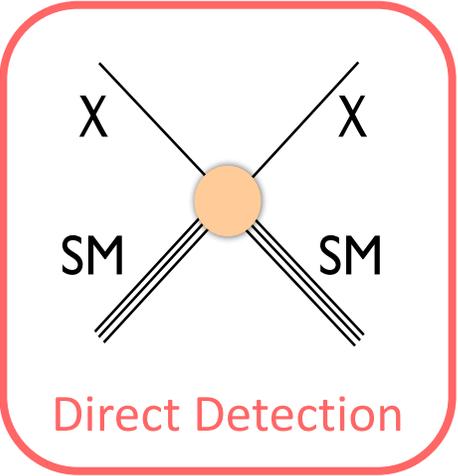
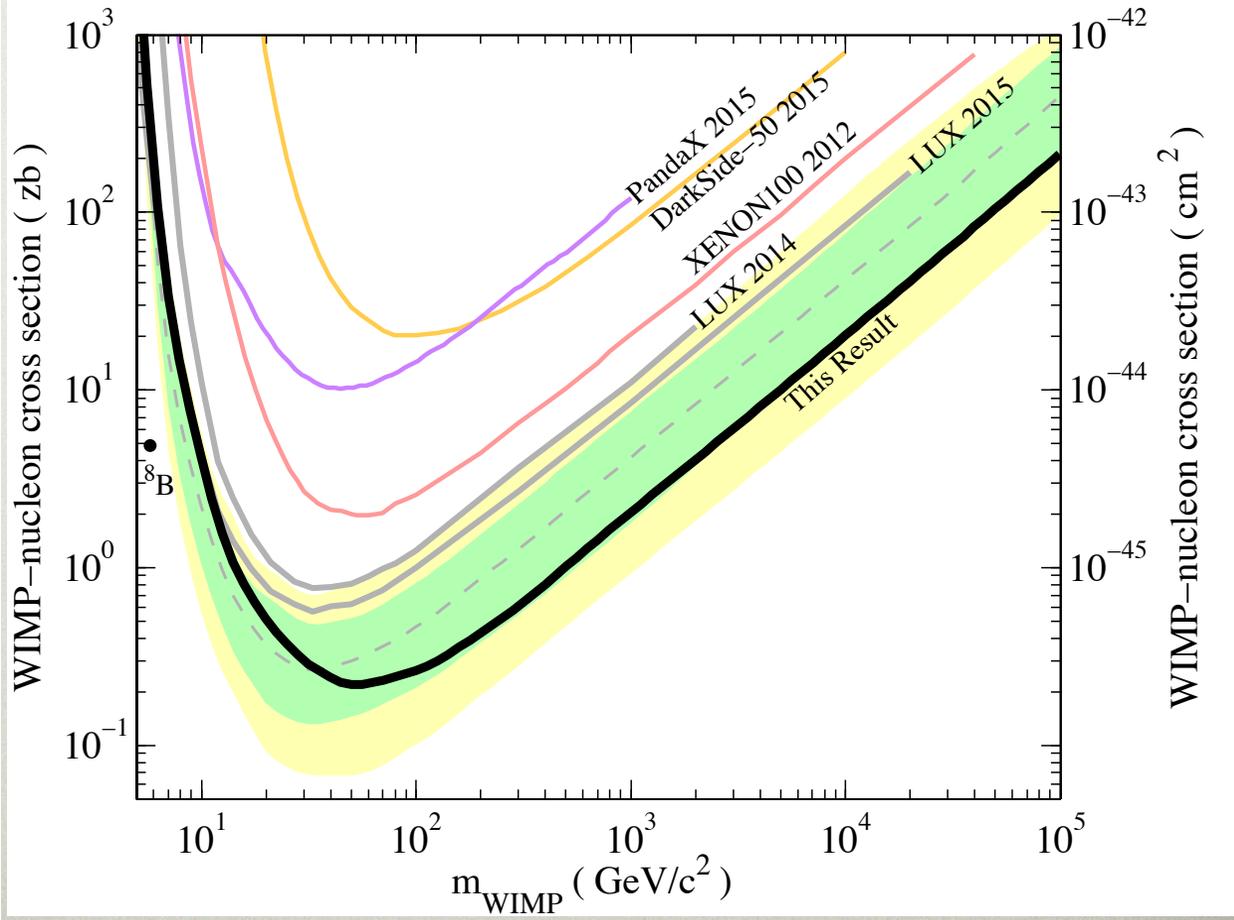
Texas A&M University

NP Interpretations at the LHC Workshop, Argonne

PH, C. Wagner arXiv:1404.0392

PH, R. Roglans, D. Spiegel, Y. Sun and C. Wagner arXiv:1701.02737

# Direct Detection



- The direct detection experiments are pushing rapidly into the region with a Higgs exchange
- Suppress the neutrino direct detection rate?

# outline

- How to suppress the neutralino direct detection rate?
  - Blind Spot scenarios
- Deviations from the Blind Spots
  - current constraints
  - future reaches
- How to probe the blind spot scenarios at the LHC?
  - heavy Higgs searches
  - Electroweakino searches

# outline

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# Suppress the neutralino direct detection rate

- Consider a neutralino scattering off a down-type quark.

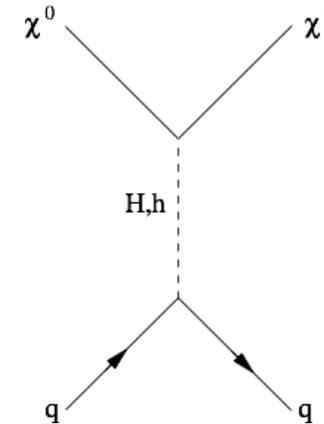
In gauge eigenstates

$$h = \frac{1}{\sqrt{2}}(\cos \alpha H_u - \sin \alpha H_d)$$

$$H = \frac{1}{\sqrt{2}}(\sin \alpha H_d + \cos \alpha H_u).$$

Amplitude

$$a_d \sim \frac{m_d}{\cos \beta} \left( \frac{-\sin \alpha g_{\chi\chi h}}{m_h^2} + \frac{\cos \alpha g_{\chi\chi H}}{m_H^2} \right).$$



$$\tilde{\chi} = N_{i1} \tilde{B} + N_{i2} \tilde{W} + N_{i3} \tilde{H}_d + N_{i4} \tilde{H}_u \quad L \supset -\sqrt{2}g'Y_{H_u}\tilde{B}\tilde{H}_uH_u^* - \sqrt{2}g\tilde{W}^a\tilde{H}_u t^a H_u^* + (u \leftrightarrow d)$$

Couplings

$$g_{\chi\chi h} \sim (g_1 N_{i1} - g_2 N_{i2})(-\cos \alpha N_{i4} - \sin \alpha N_{i3})$$

$$g_{\chi\chi H} \sim (g_1 N_{i1} - g_2 N_{i2})(-\sin \alpha N_{i4} + \cos \alpha N_{i3}).$$

# Suppress the neutralino direct detection rate

$$a_d \sim \frac{m_d(g_1 N_{i1} - g_2 N_{i2})}{\cos \beta} \left[ N_{i4} \sin \alpha \cos \alpha \left( \frac{1}{m_h^2} - \frac{1}{m_H^2} \right) + N_{i3} \left( \frac{\sin^2 \alpha}{m_h^2} + \frac{\cos^2 \alpha}{m_H^2} \right) \right]$$

## Loop Effects

$$L = f_d \bar{d}_L d_R H_d^0 + \epsilon_d f_d \bar{d}_L d_R H_u^{0*} + h.c.,$$

$$a_d \sim \frac{\bar{m}_d(g_1 N_{i1} - g_2 N_{i2})}{\cos \beta} \left[ N_{i4} \sin \alpha \cos \alpha \left( \frac{1 - \epsilon_d / \tan \alpha}{m_h^2} - \frac{1 + \epsilon_d \tan \alpha}{m_H^2} \right) + N_{i3} \left( \frac{\sin^2 \alpha (1 - \epsilon_d / \tan \alpha)}{m_h^2} + \frac{\cos^2 \alpha (1 + \epsilon_d \tan \alpha)}{m_H^2} \right) \right]$$

$$\bar{m}_d \equiv \frac{m_d}{1 + \epsilon_d \tan \beta} \quad \epsilon_d \approx \frac{2\alpha_s}{3\pi} M_3 \mu C_0(m_0^2, m_R^2, |M_3|^2)$$

$$C_0(X, Y, Z) = \frac{y}{(x-y)(z-y)} \log(y/x) + \frac{z}{(x-z)(y-z)} \log(z/x).$$

When 1st&2nd gen squarks are heavy,  $\epsilon_d$  is suppressed

# Suppress the neutralino direct detection rate

$$a_d \sim \frac{m_d(g_1 N_{i1} - g_2 N_{i2})}{\cos \beta} \left[ N_{i4} \sin \alpha \cos \alpha \left( \frac{1}{m_h^2} - \frac{1}{m_H^2} \right) + N_{i3} \left( \frac{\sin^2 \alpha}{m_h^2} + \frac{\cos^2 \alpha}{m_H^2} \right) \right]$$

$$N_{i3} \sim (m_\chi \cos \beta + \mu \sin \beta) \quad \sin \alpha \approx -\cos \beta$$

$$N_{i4} \sim (m_\chi \sin \beta + \mu \cos \beta).$$

Pierce, Shah, and Freese. arXiv:1309.7351

$$a_d \sim \frac{m_d}{\cos \beta} \left[ \cos \beta (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} - \mu \sin \beta \cos 2\beta \frac{1}{m_H^2} \right]$$

$$a_u \sim \frac{m_u}{\sin \beta} \left[ \sin \beta (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} + \mu \cos \beta \cos 2\beta \frac{1}{m_H^2} \right]$$

$$a_p = \left( \sum_{q=u,d,s} f_{Tq}^{(p)} \frac{a_q}{m_q} + \frac{2}{27} f_{TG}^{(p)} \sum_{q=c,b,t} \frac{a_q}{m_q} \right) m_p$$

# Blind Spot in Dark Matter Direct Detection

$$\sigma_p^{SI} \sim \left[ \underbrace{(F_d^{(p)} + F_u^{(p)})(m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2}}_{\text{125 GeV Higgs Exchange}} + \underbrace{\mu \tan \beta \cos 2\beta (-F_d^{(p)} + F_u^{(p)}/\tan^2 \beta) \frac{1}{m_H^2}}_{\text{Heavy Higgs Exchange}} \right]^2$$

$F_d^{(p)} \approx 0.15, F_u^{(p)} \approx 0.14$

Blind Spot at  $(m_{\chi_1^0} + \mu \sin 2\beta) = 0$

$$2 (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} \simeq - \mu \tan \beta \frac{1}{m_H^2}$$

See Cheung, Hall, Pinner, and Ruderman. arXiv:1211.4873  
 Han, Kling, Su, and Wu. arXiv:1612.02387

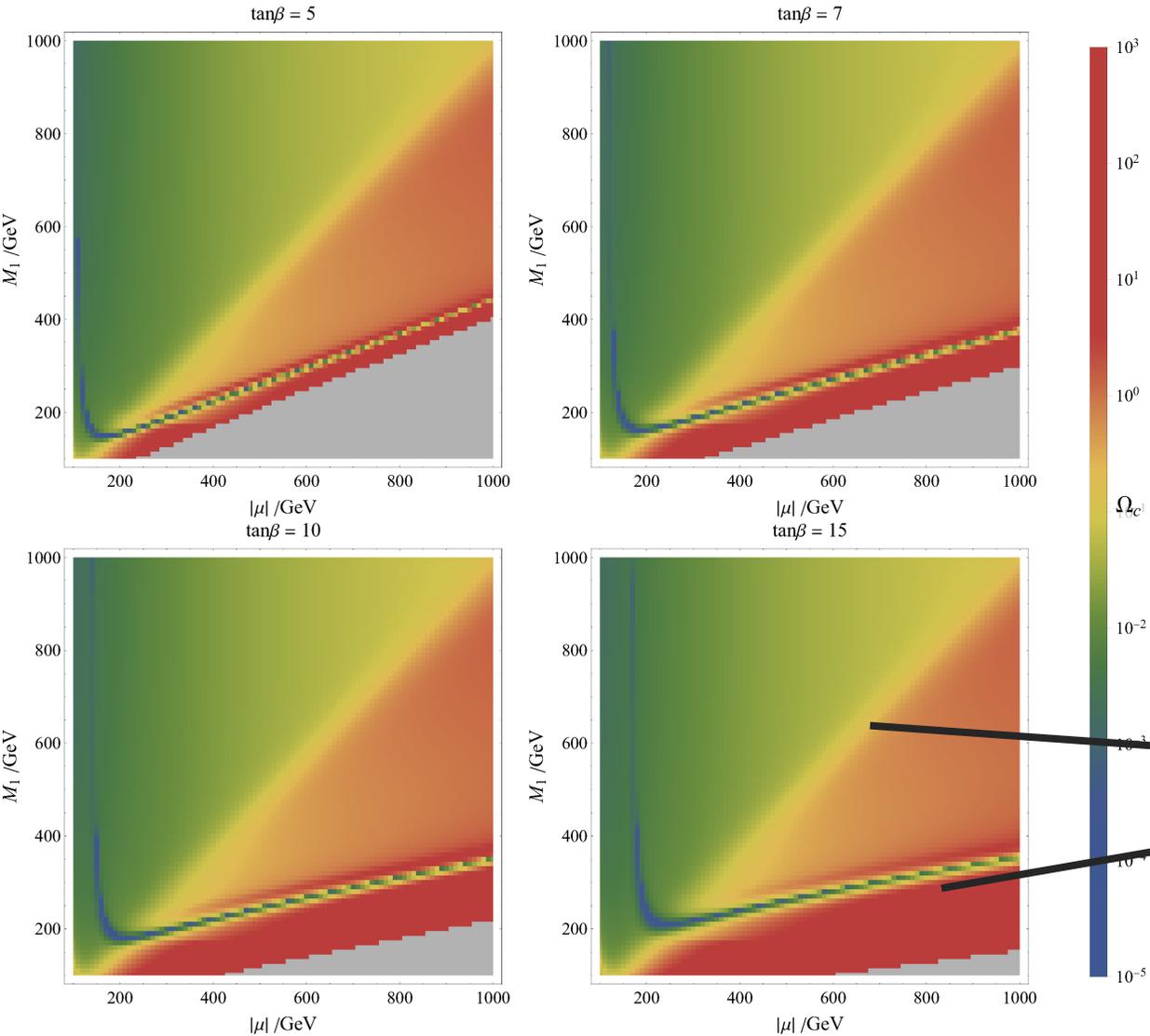
PH, C. Wagner, 14

Higgsino mass  $\mu < 0$

- Suppress the lightest neutralino coupling to 125 GeV Higgs
- Destructive interference between the 125 GeV Higgs and the heavy Higgs exchange

Reduce the pMSSM parameter space to  $M_1, \mu, \tan \beta,$  and  $m_A$   
 For a full pMSSM study, see Cahill-Rowley et al, 1405.6716. For a loop level analysis, see Berlin, Hooper, and McDermott, arXiv:1508.05390

# Blind Spot and the Relic Density



Relevant parameters:  
 $M_1$ ,  $\mu$ ,  $\tan\beta$ , and  $m_A$

Under Abundant Over Abundant

$m_A$  is chosen to minimize the DD rate

$$2 (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} \simeq - \mu \tan \beta \frac{1}{m_H^2}$$

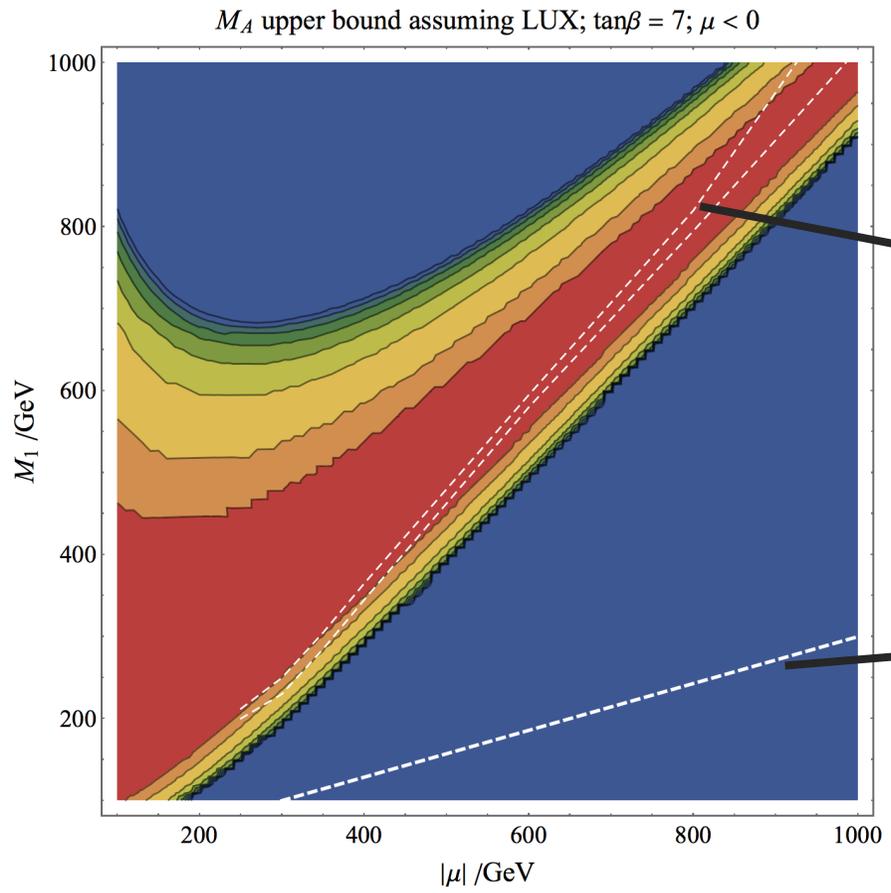
Two branches  
 well tempered  $M_1 \sim \mu$   
 A-funnel  $m_\chi \sim m_A/2$

# outline

- How to suppress the neutralino direct detection rate?
  - Blind Spot scenarios
- Deviations from the Blind Spots
  - current constraints
  - future reach
- How to probe the blind spot scenarios at the LHC?
  - heavy Higgs searches
  - Electroweakino searches

# Current Direct Detection Constraints

$$\sigma_p^{SI} \sim \left[ (F_d^{(p)} + F_u^{(p)})(m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} + \mu \tan \beta \cos 2\beta (-F_d^{(p)} + F_u^{(p)}/\tan^2 \beta) \frac{1}{m_H^2} \right]^2 \quad \rightarrow \quad m_A \text{ upper bound}$$



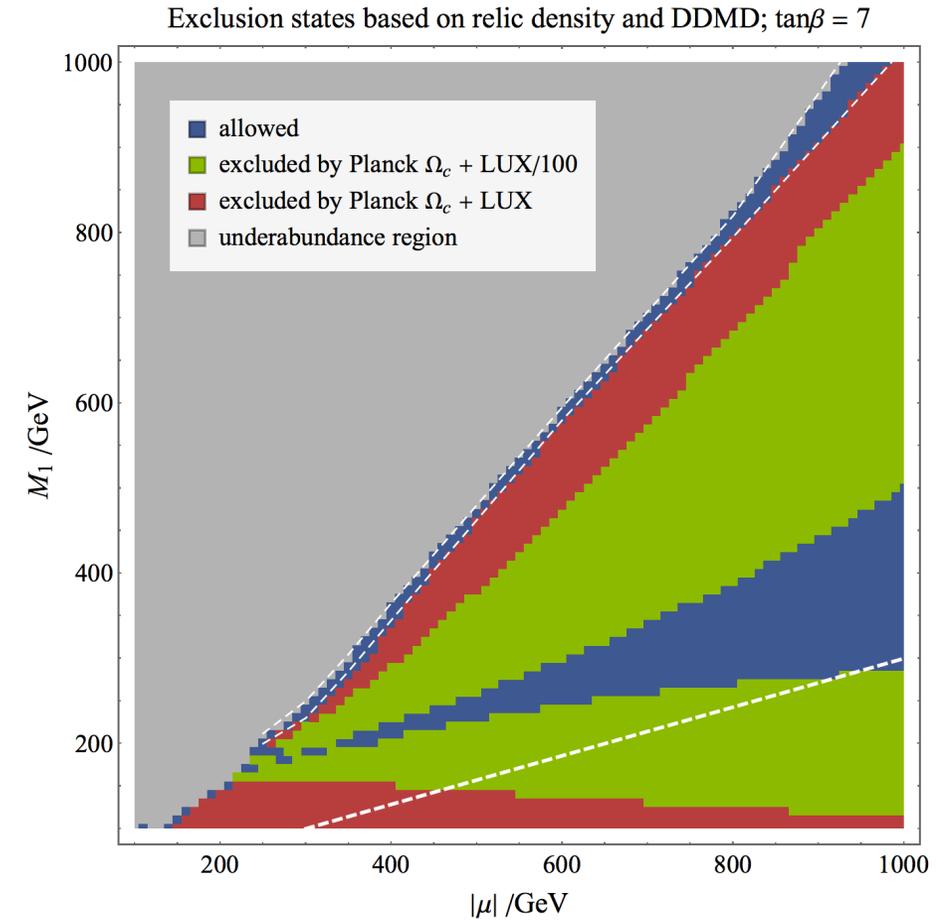
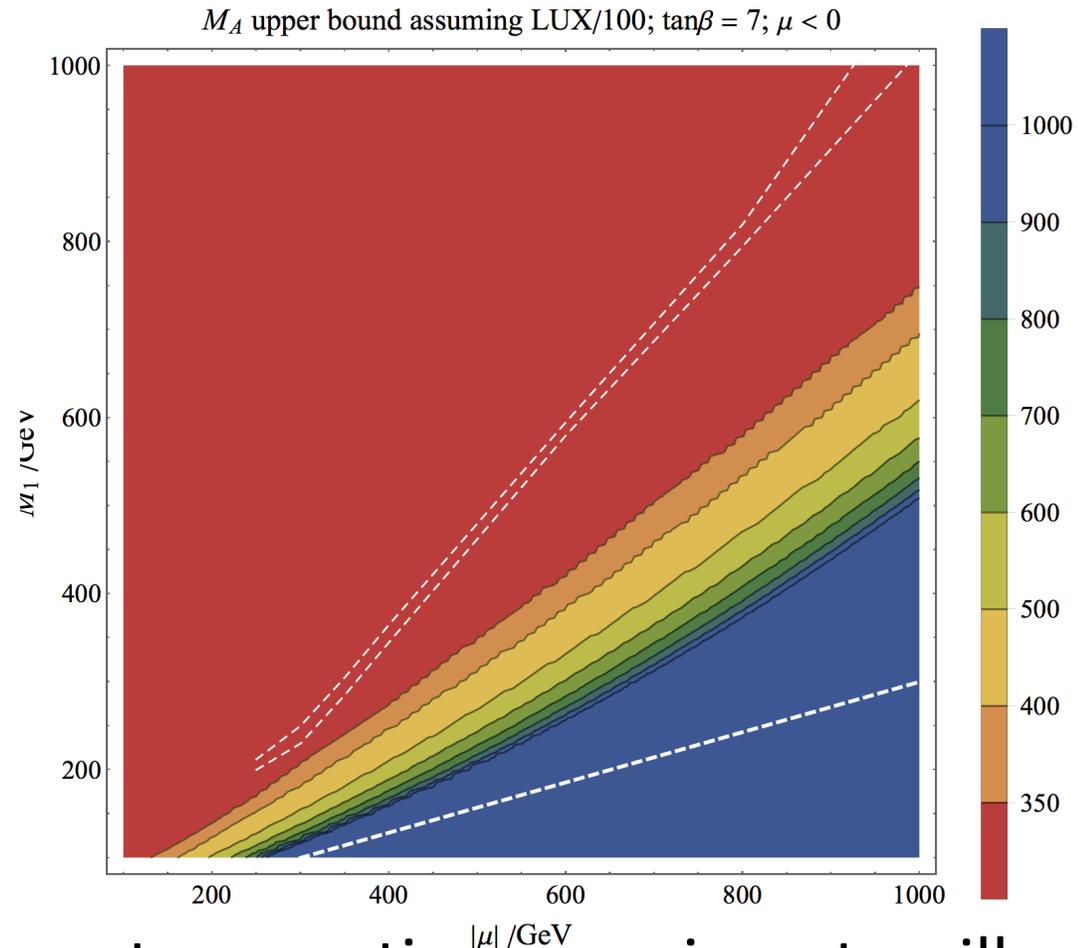
Assume the right relic density

In the well-tempered region,  $m_A < 350 - 400$  GeV

The A-funnel region is allowed by LUX and relic density considerations

neutralino higgs coupling vanishes

# Future reach



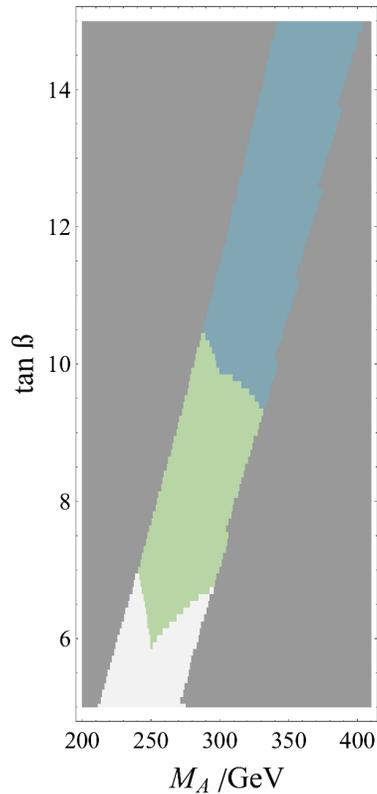
next generation experiments will push  $m_A$  to be smaller than 300 GeV  
in the well-tempered region

# outline

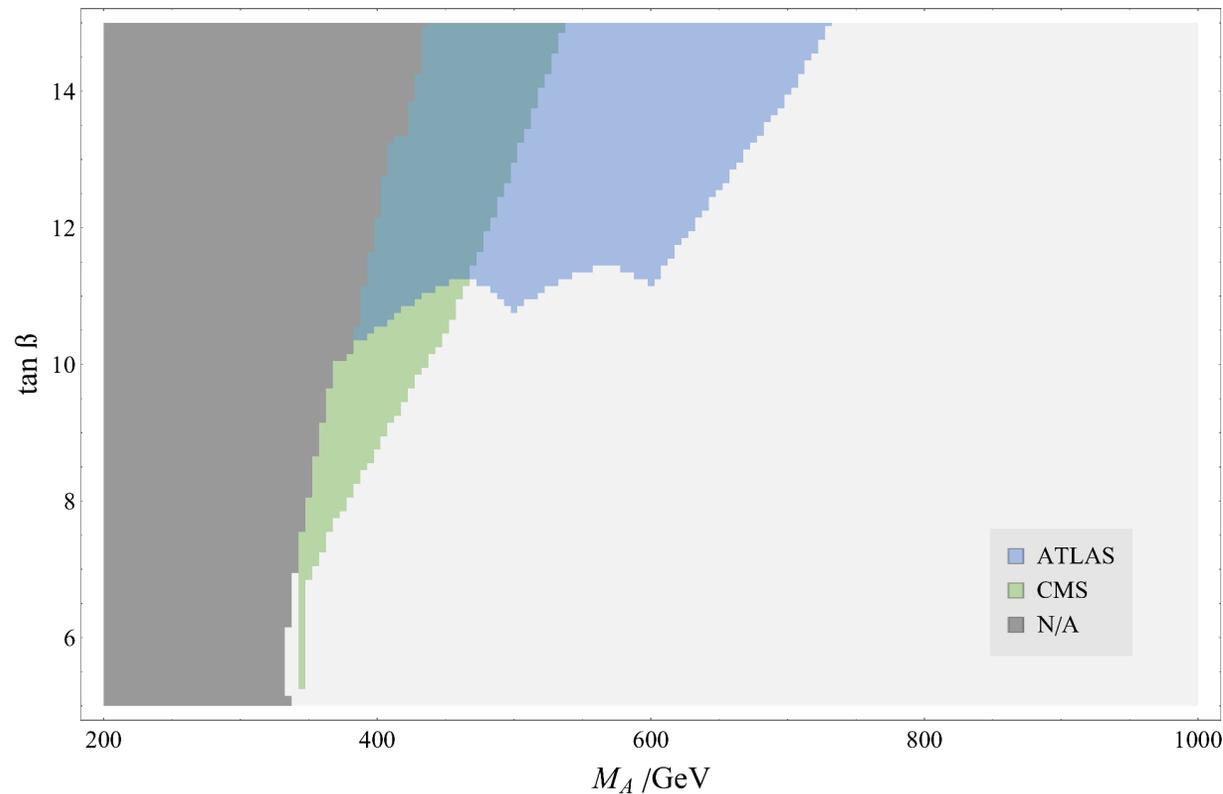
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# Probe the Blind Spots scenarios – collider searches

$$2 (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} \simeq - \mu \tan \beta \frac{1}{m_H^2}$$



Well-tempered



A-funnel

Well-tempered region :  
open for  $\tan \beta < 6$   
completely ruled out for  
 $\tan \beta \gtrsim 7$

A-funnel region starts to  
get excluded as  $\tan \beta$   
increases

# Probe the Blind Spots scenarios – collider searches

- hbb coupling

$$\frac{g_{hbb}}{g_{hbb}^{SM}} = -\frac{\sin \alpha}{\cos \beta} = \sin(\beta - \alpha) - \tan \beta \cos(\beta - \alpha)$$

hVV coupling  $\sim 1$

controls the modification

- One loop level, sizable  $\tan \beta$

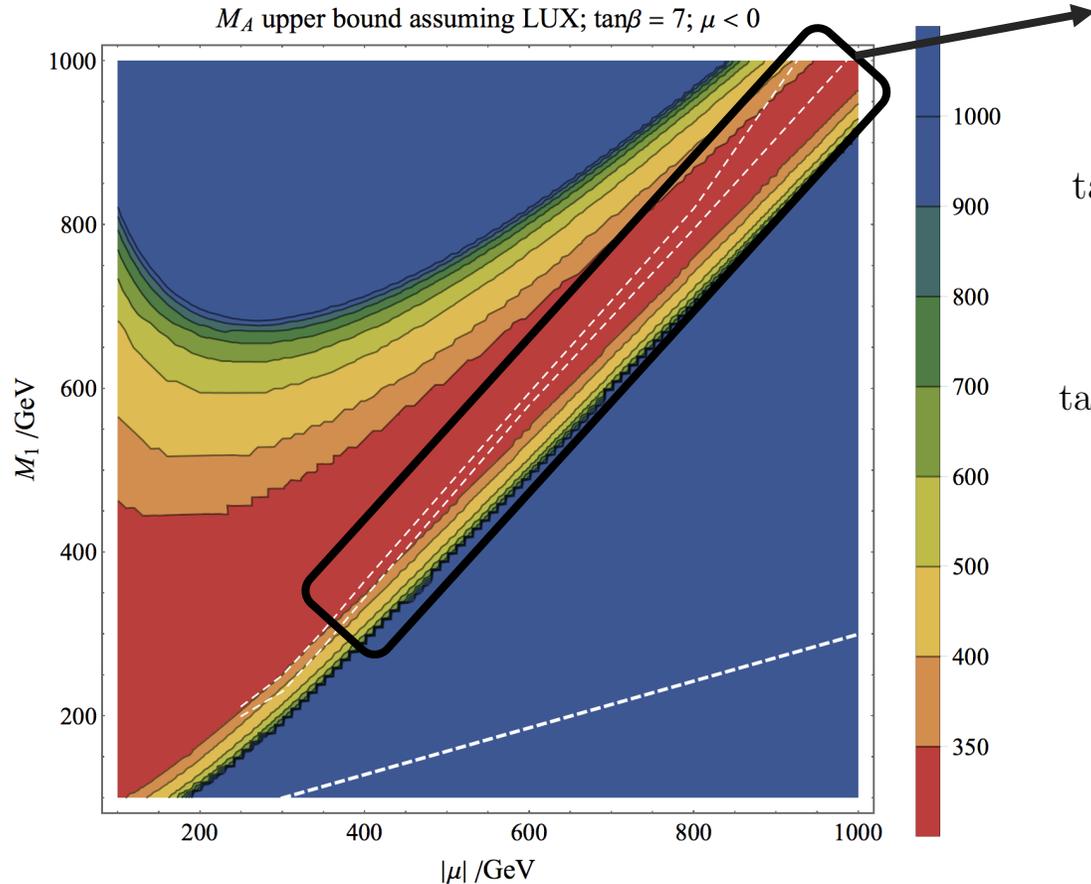
$A_t \lesssim 3 M_s$ , vacuum stability

$$\tan \beta \cos(\beta - \alpha) \simeq \frac{-1}{m_H^2 - m_h^2} \left[ m_h^2 + m_Z^2 + \frac{3m_t^4}{4\pi^2 v^2 M_S^2} A_t \mu \tan \beta \left( 1 - \frac{A_t^2}{6M_S^2} \right) \right]$$

$M_s \gtrsim \text{TeV}$ , proper Higgs mass

$m_A \gtrsim 350 \text{ GeV}$  to be consistent with the current Higgs data.

# Probe the Blind Spots scenarios – collider searches



In tension with the current Higgs data  
Extended the Higgs sector

$$\tan\beta \cos(\beta - \alpha) \simeq \frac{-1}{m_H^2 - m_h^2} \left[ m_h^2 + m_Z^2 + \frac{3m_t^4}{4\pi^2 v^2 M_S^2} A_t \mu \tan\beta \left( 1 - \frac{A_t^2}{6M_S^2} \right) \right]$$

$$\tan\beta \cos(\beta - \alpha) \simeq \frac{-1}{m_H^2 - m_h^2} \left[ m_h^2 + m_Z^2 - \lambda^2 v^2 + \frac{3m_t^4}{4\pi^2 v^2 M_S^2} A_t \mu \tan\beta \left( 1 - \frac{A_t^2}{6M_S^2} \right) \right]$$

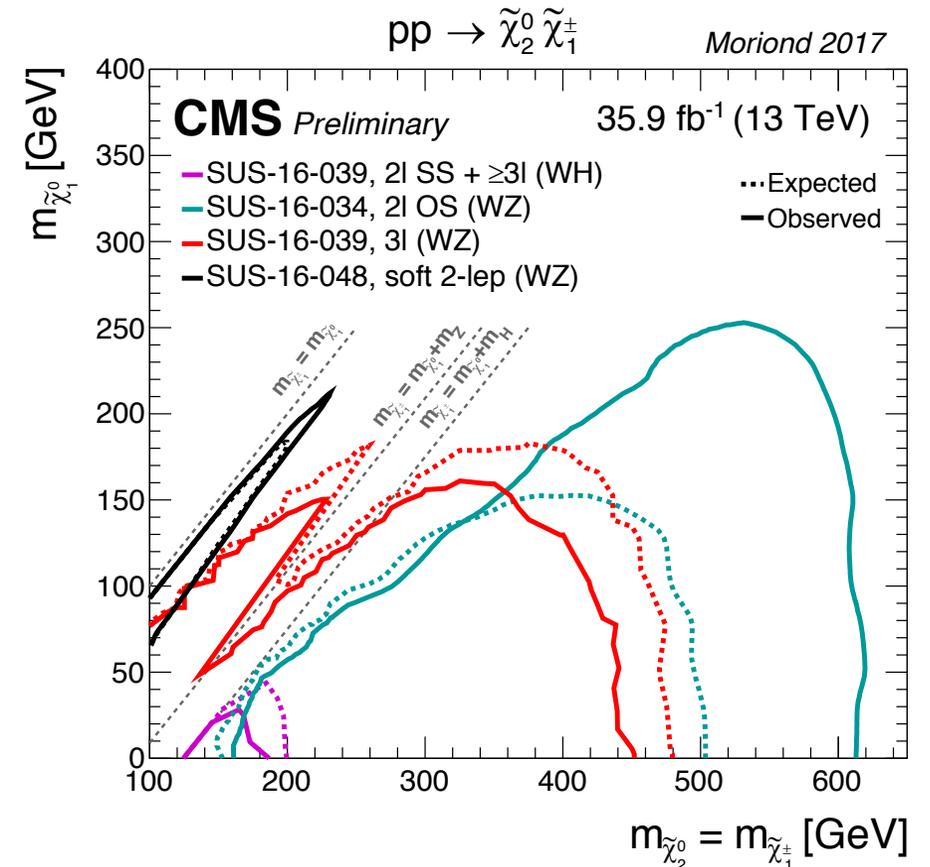
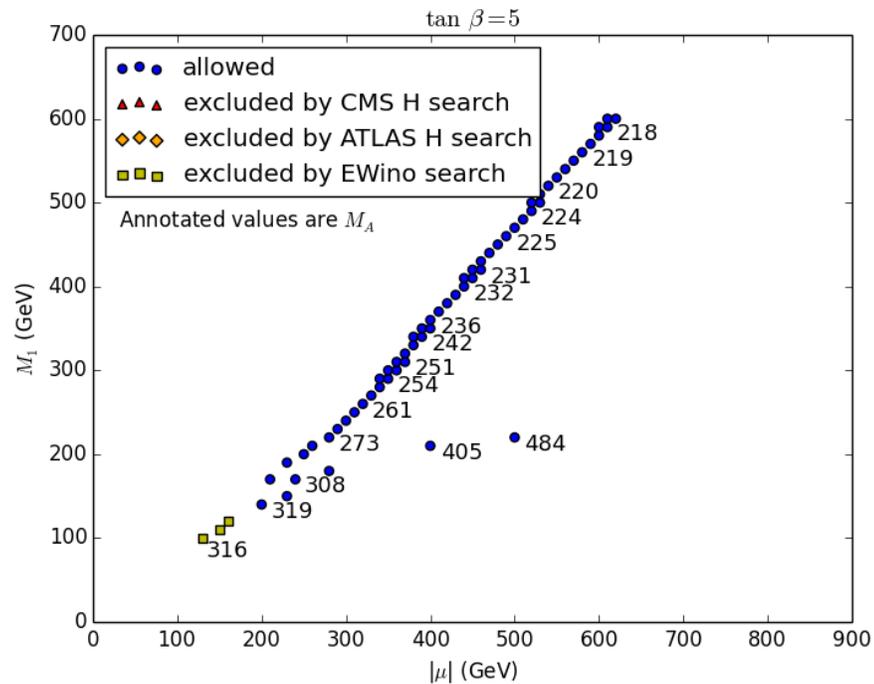
$\lambda \simeq 0.6-07$ , even for  $m_H \simeq 200$  GeV, the modification in SM Higgs coupling is small enough.

Blind spot scenarios in NMSSM, see Badziak, Olechowski, and Szczerbiak, arXiv:1512.02472

# Probe the Blind Spots scenarios – collider searches

- 3 neutralinos, (mixtures of bino and Higgsinos), the lightest chargino(Higgsino-like)

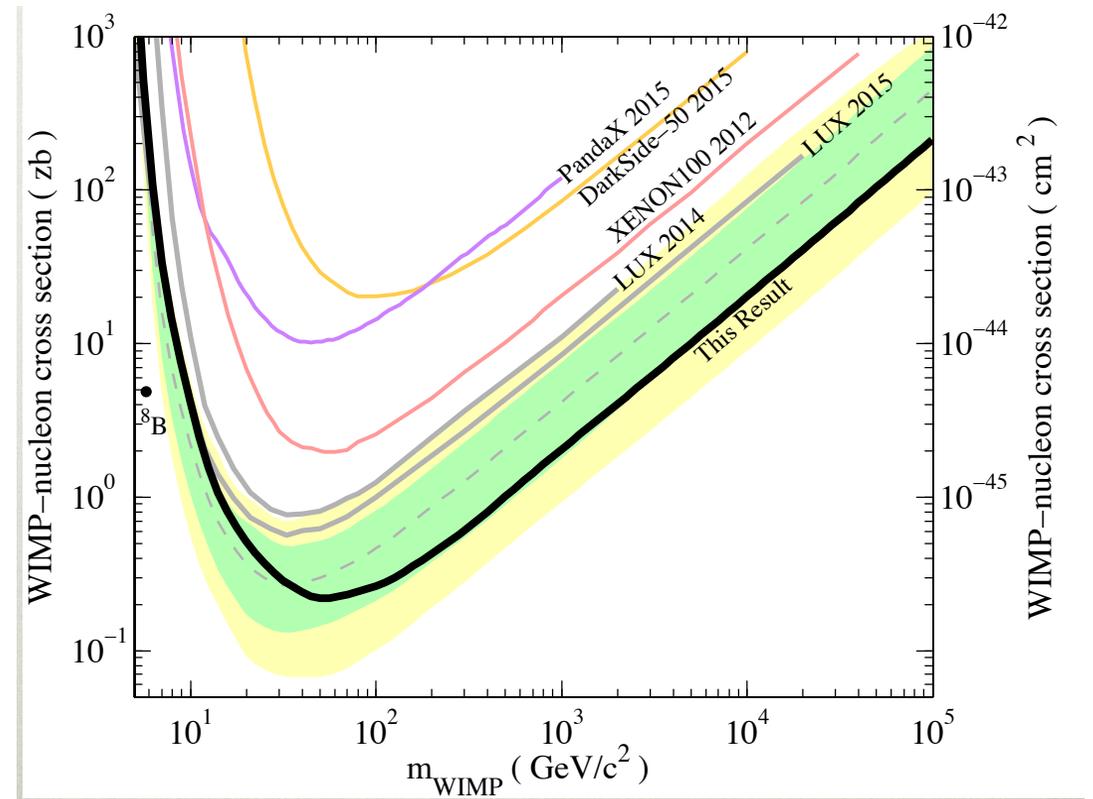
$$\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow W Z \tilde{\chi}_1^0 \tilde{\chi}_1^0 \quad \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow W h \tilde{\chi}_1^0 \tilde{\chi}_1^0$$



# Conclusion

- Identify a blind spot scenario for neutralino dark matter
- possible probes at the LHC
  - heavy Higgs searches, precision Higgs, and Electroweakinos searches

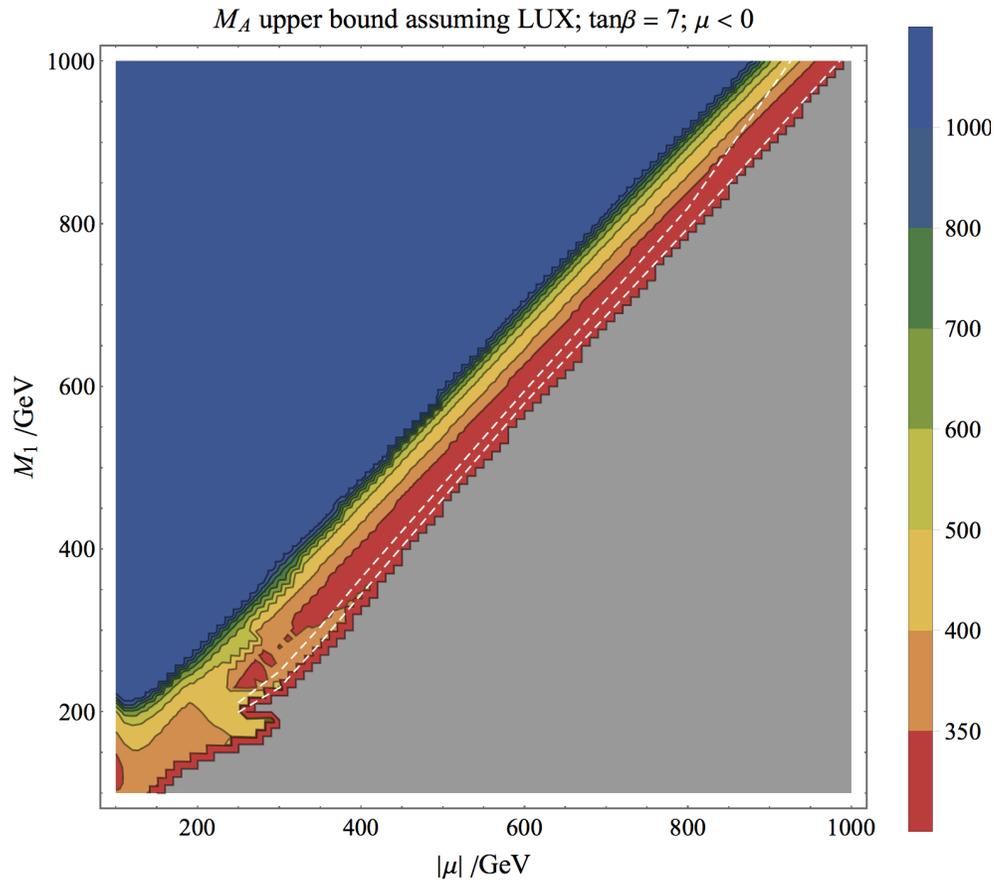
$$2 (m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} \simeq - \mu \tan \beta \frac{1}{m_H^2}$$



backup

# Current Direct Detection Constraints

$$\sigma_p^{SI} \sim \left[ (F_d^{(p)} + F_u^{(p)})(m_\chi + \mu \sin 2\beta) \frac{1}{m_h^2} + \mu \tan \beta \cos 2\beta (-F_d^{(p)} + F_u^{(p)}/\tan^2 \beta) \frac{1}{m_H^2} \right]^2 \rightarrow m_A \text{ upper bound}$$



Rescale according to the local density

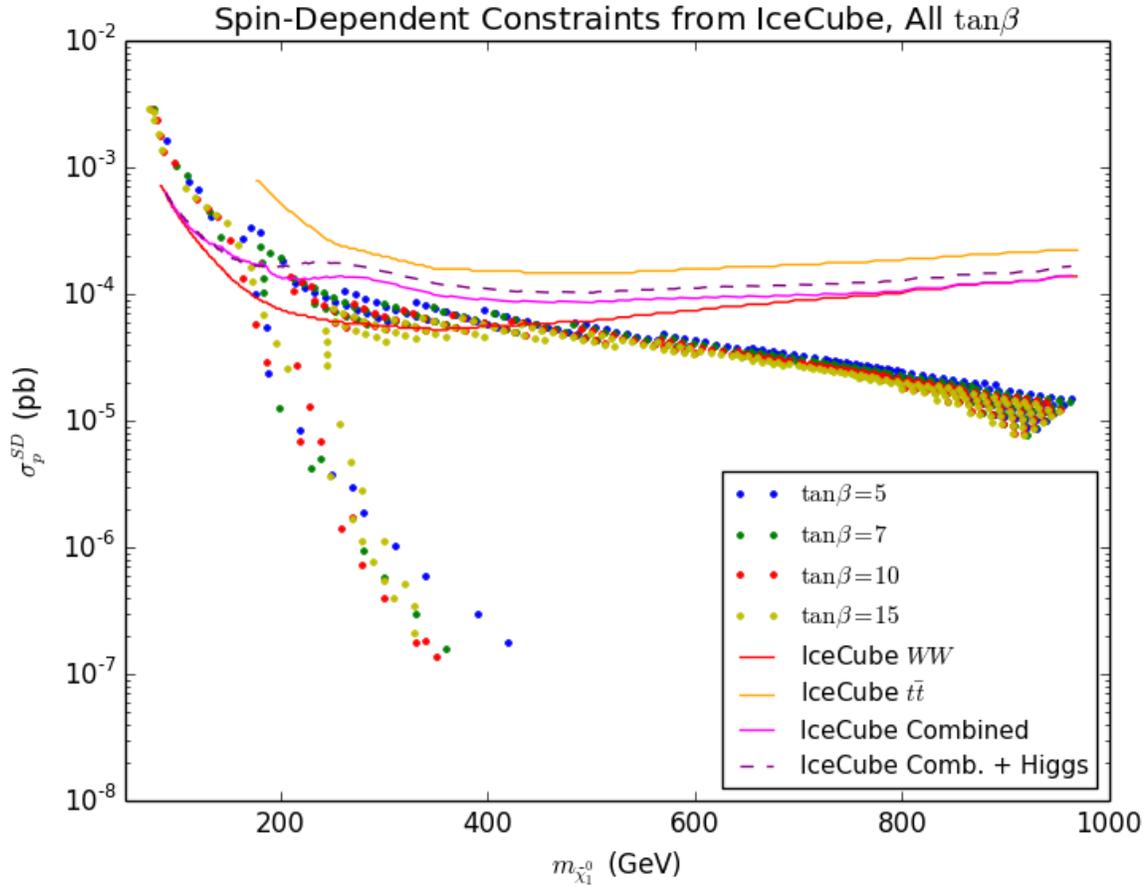
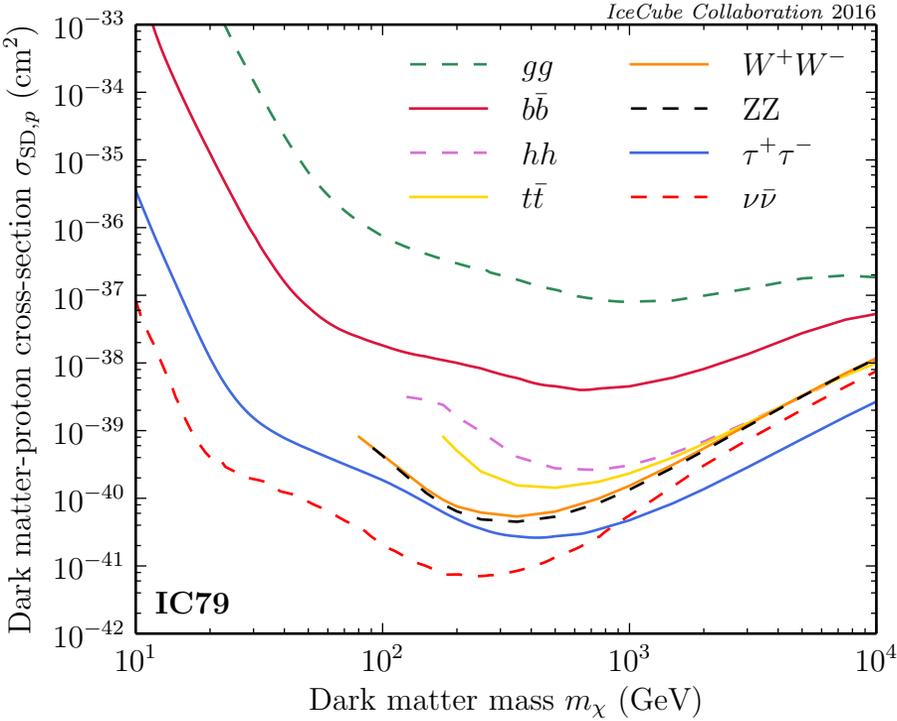
In the well-tempered region,  $m_A < 350 - 400$  GeV

relaxed  $m_A$  bound to the left of the well-tempered region

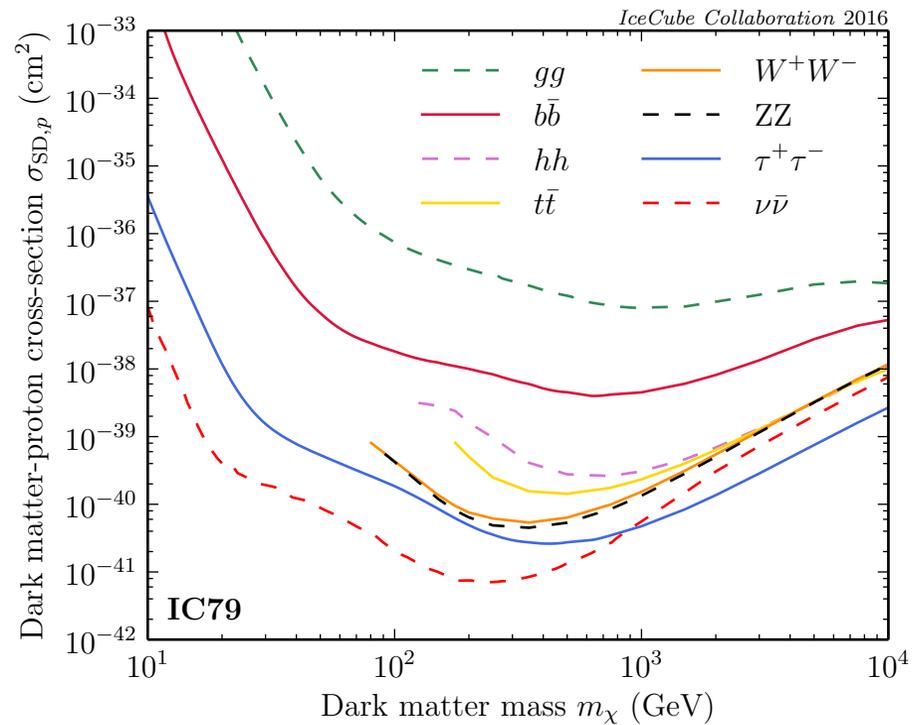
Will discuss direct heavy Higgs search and precision Higgs later

# Probe the Blind Spots scenarios – IceCube

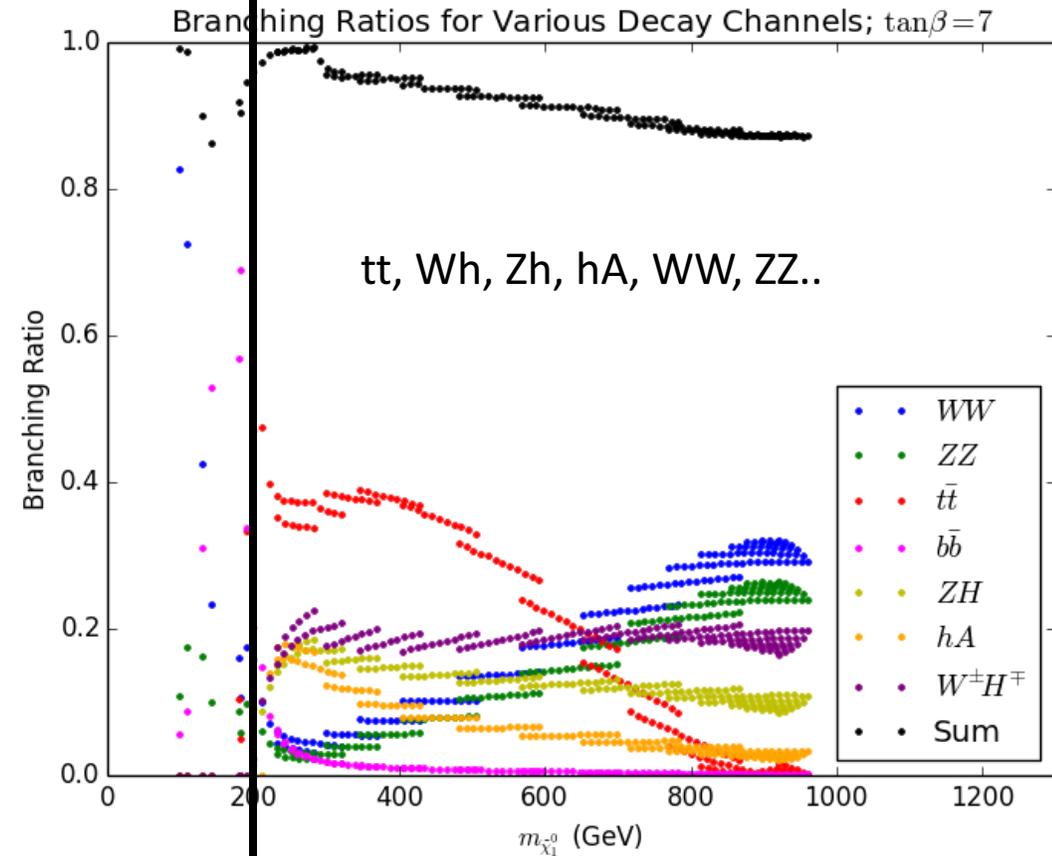
exclude the well-tempered region for  $m_\chi \lesssim 200$  GeV



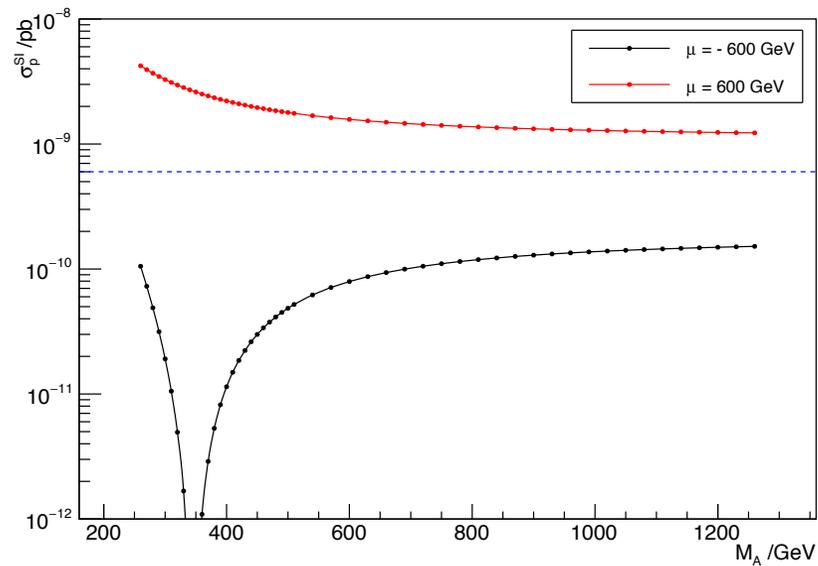
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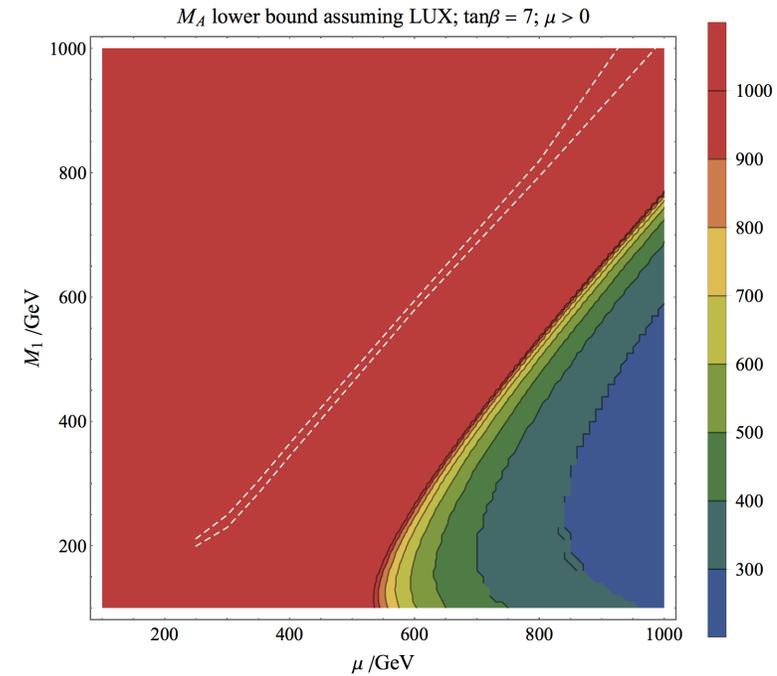
bb, weak



# Positive $\mu$

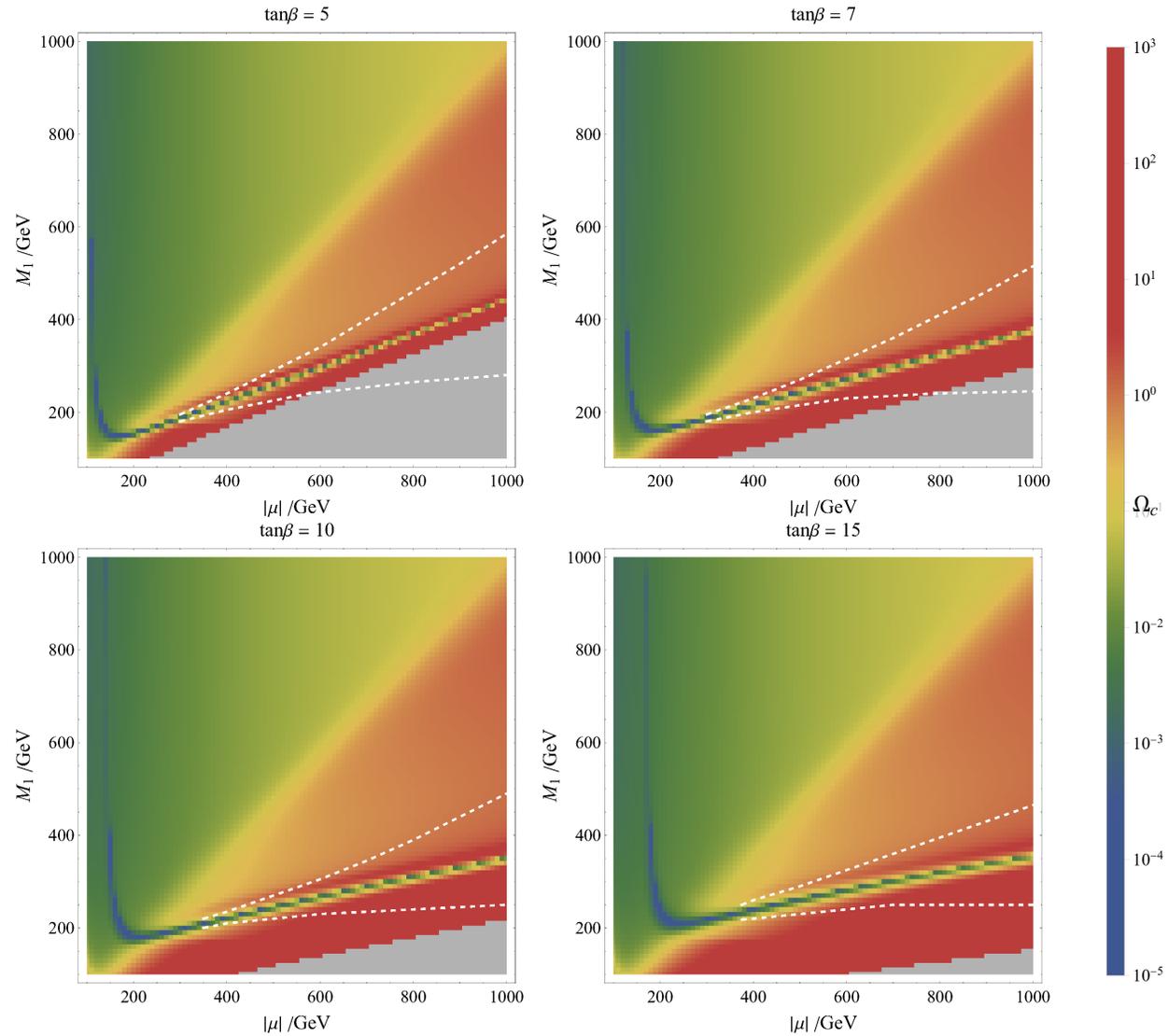


Larger Higgs Neutralino coupling  
Constructive interference

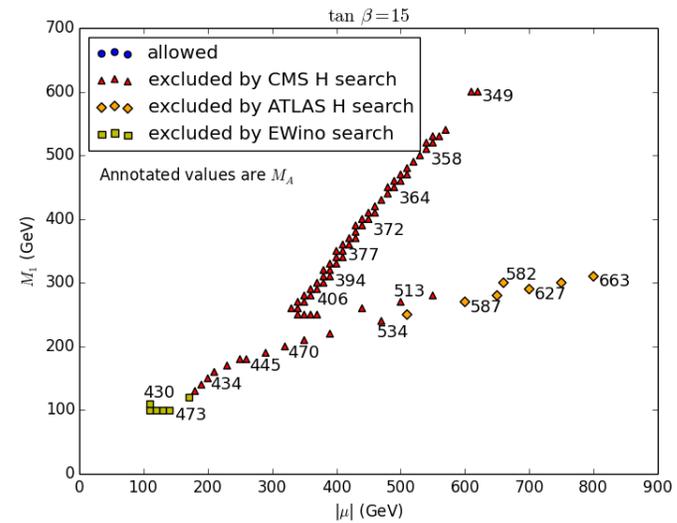
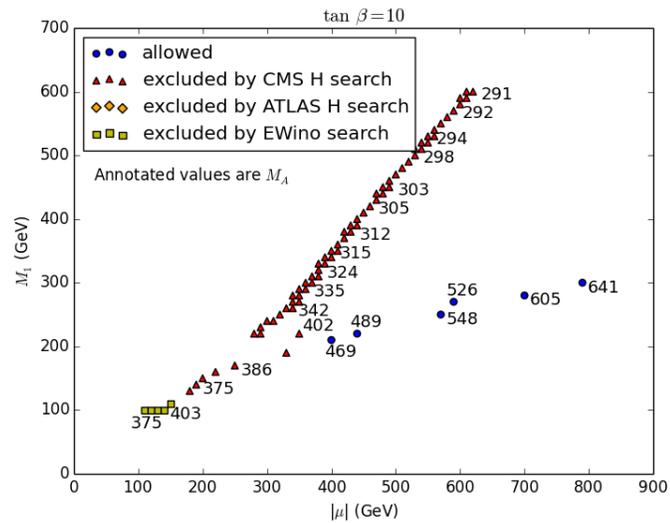
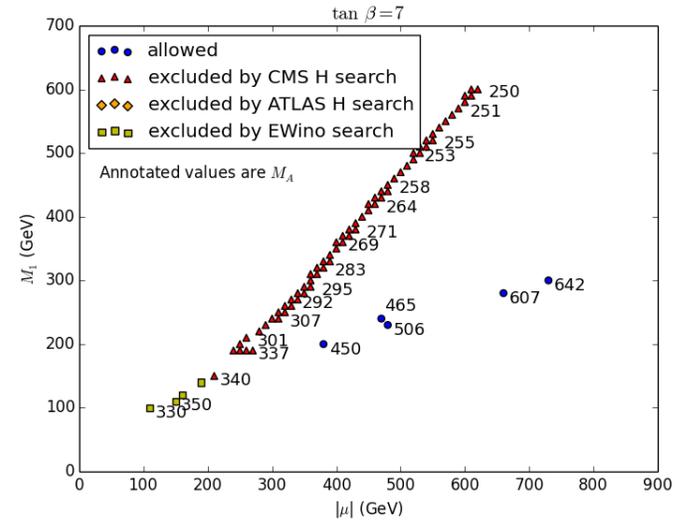
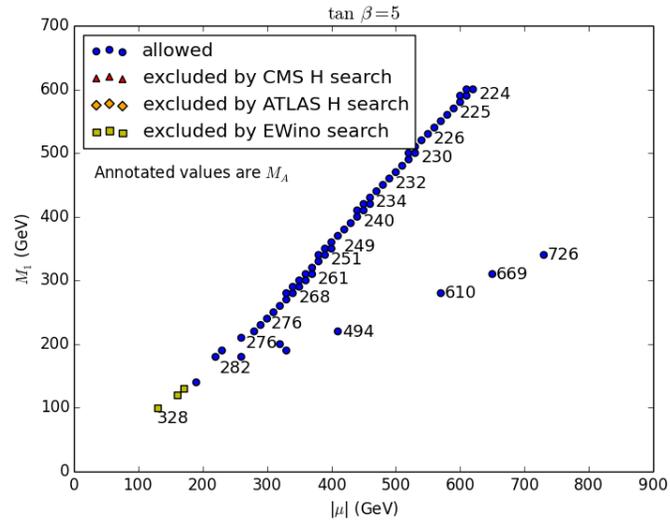


lower bound on  $m_A$

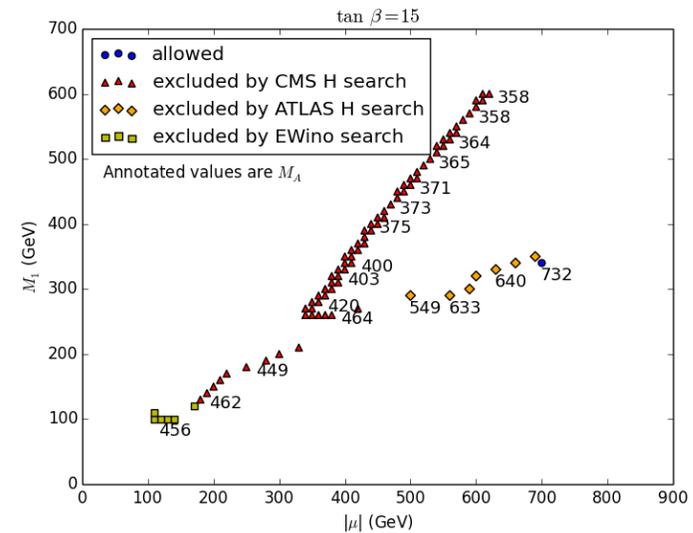
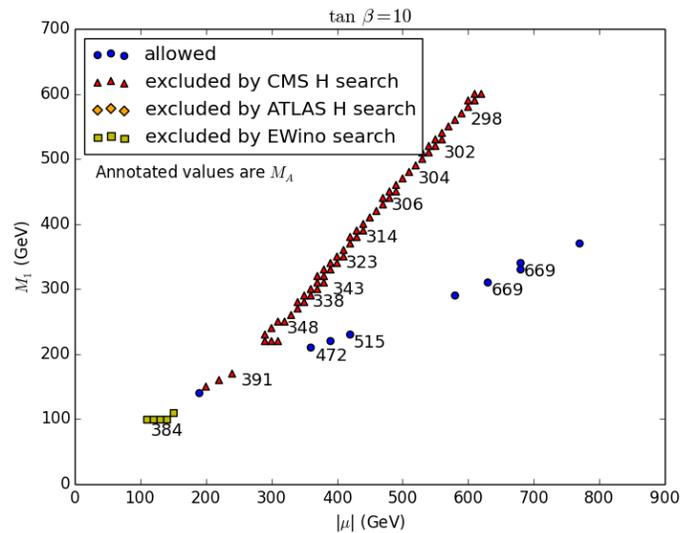
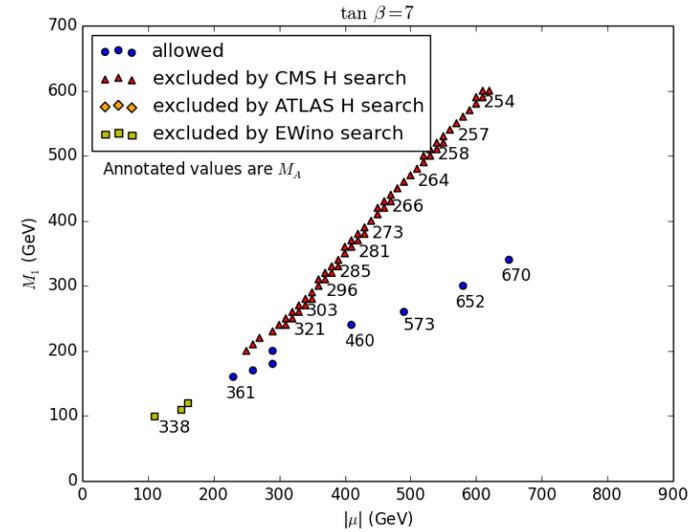
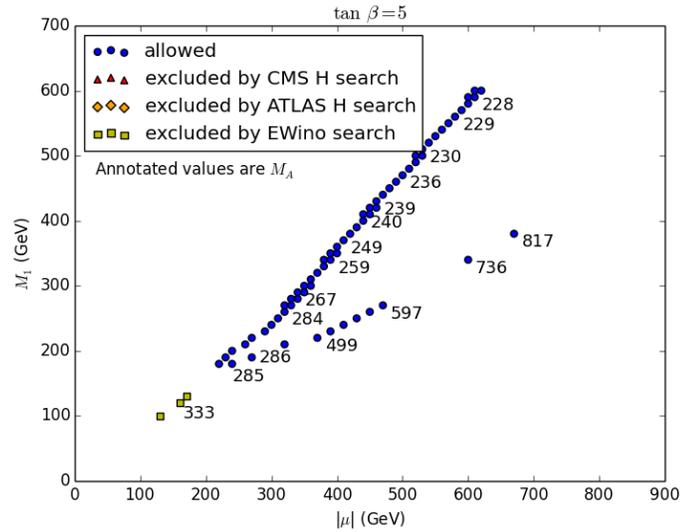
# A –funnel and the LUX limit



# Exclusions



# Exclusions, mA upper bound



# EW-ino production cross section

