

The Magnetic Field Profile in Strongly Magnetized Neutron Stars

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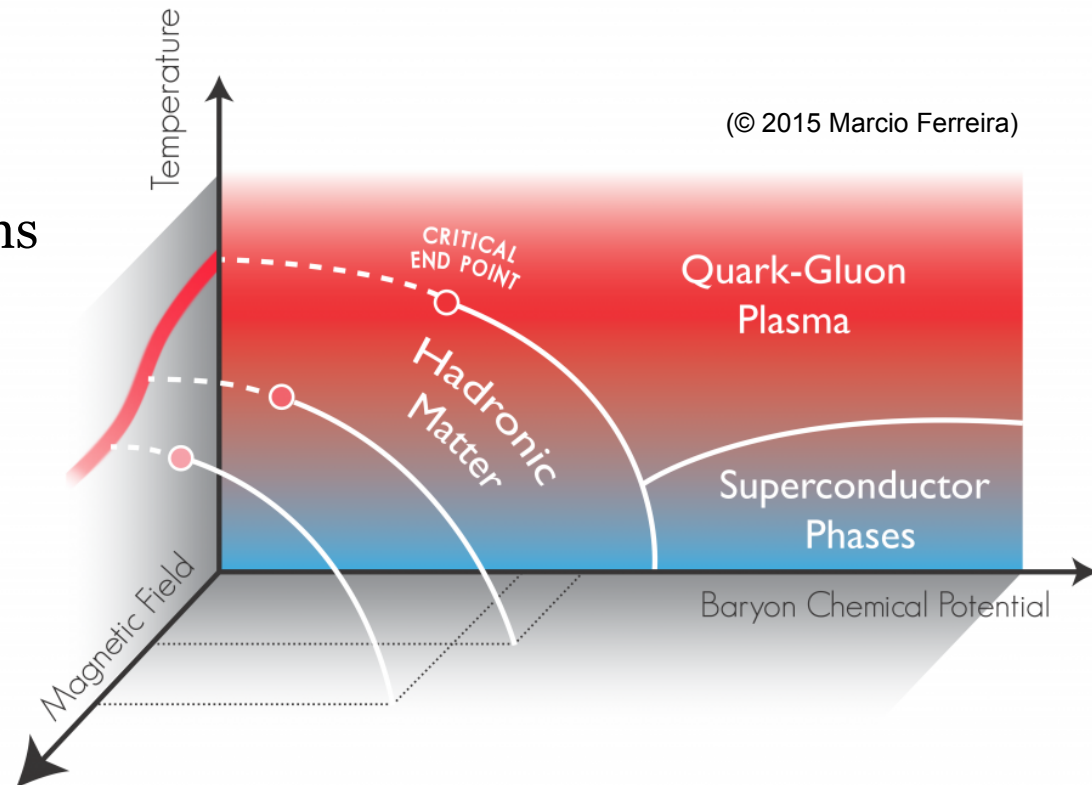
arXiv: 1612.05795



★ Motivation

- understand how B changes inside neutron stars to calculate

- changes in EoS stiffness
- changes in population
- possible phase transitions



- transport properties (thermal and electric conductivities), etc.
- temperature profiles (for fixed entropy per baryon)

Franzon, Dexheimer,
and Schramm
Phys. Rev. D (2016)

★ Ad hoc B profiles from literature

- as a function of baryon density:

$$B^*(n_B/n_0) = B_{surf} + B_0 \left[1 - e^{-\beta(n_B/n_0)^\gamma} \right]$$

Bandyopadhyay,
Chakrabarty and Pal
Phys. Rev. Let. (1997)

- as a function of chemical potential:

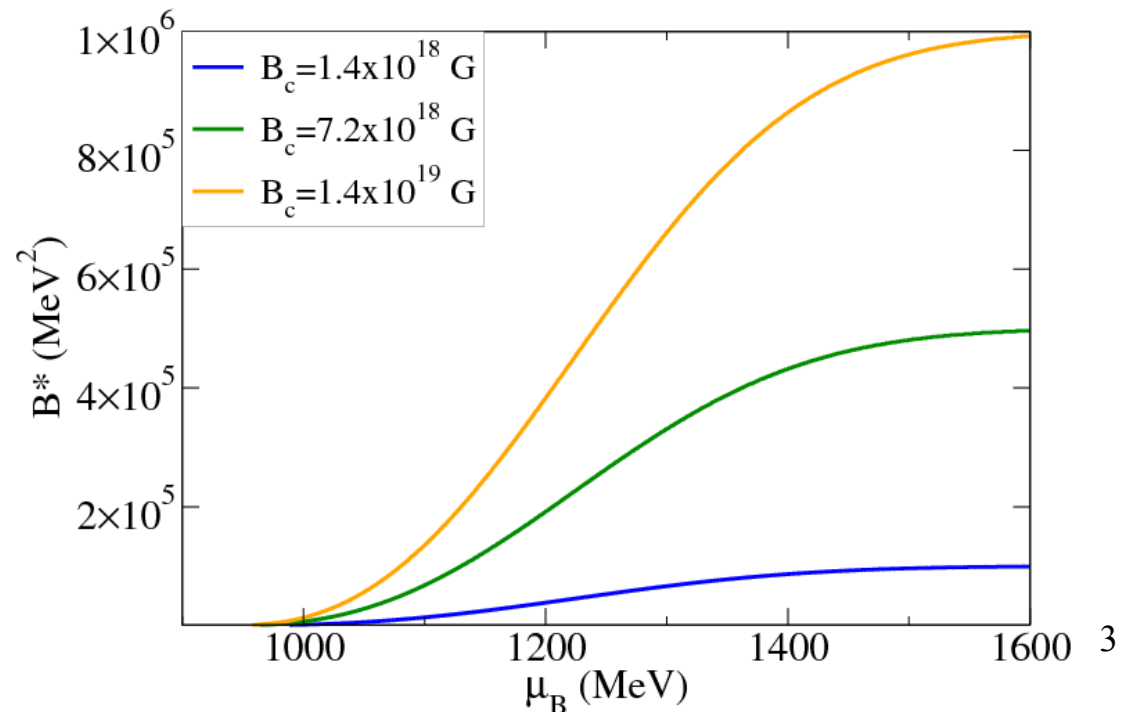
$$B^*(\mu_B) = B_{surf} + B_c \left[1 - e^{b \frac{(\mu_B - 938)^a}{938}} \right]$$

Dexheimer, Negreiros
and Schramm
Eur. Phys. J. A (2012)

and others

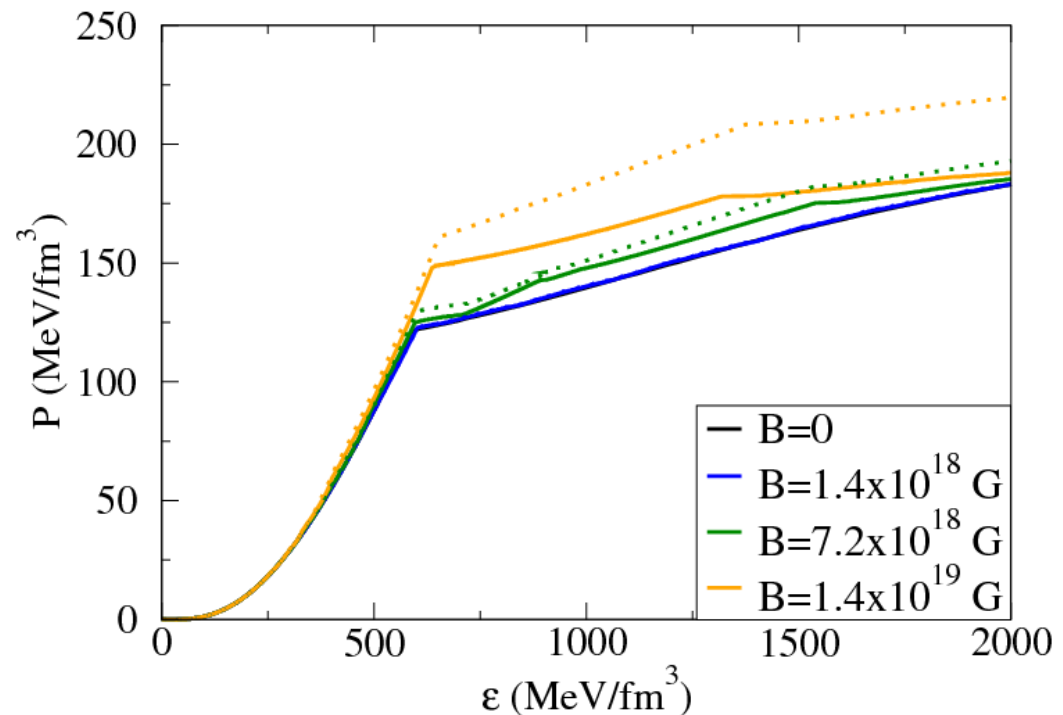
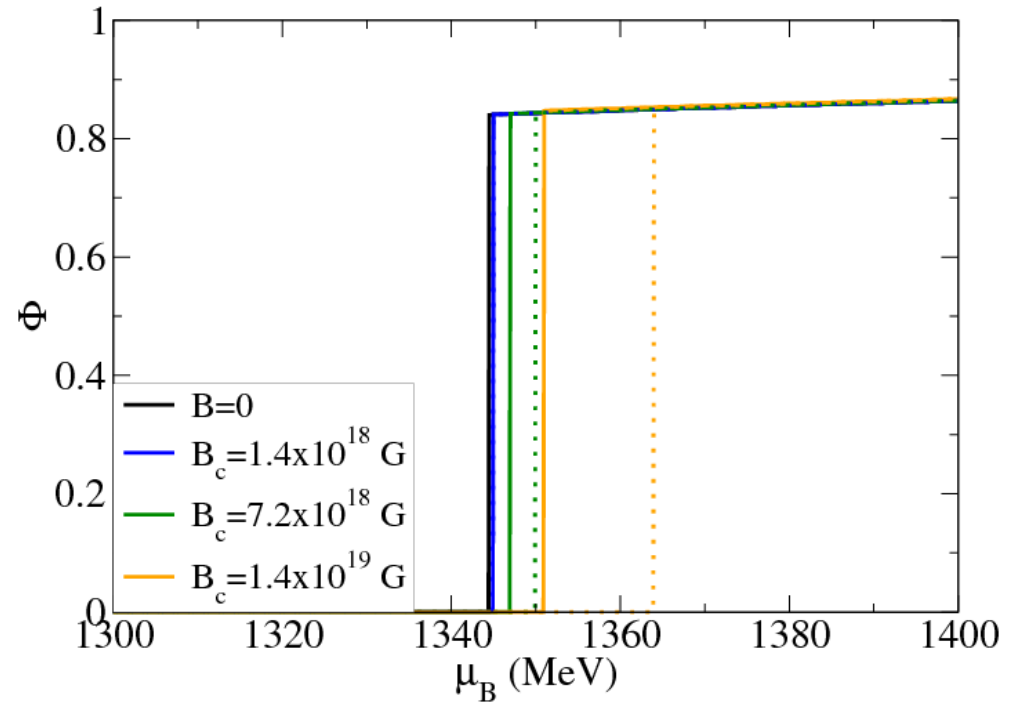
- do not fulfill
Maxwell's equations

Menezes and Alloy
arXiv:1607.07687 (2016)



★ Ad hoc B-profile effects on chiral (CMF) EOS

- Phi signals deconfinement which is pushed to larger chemical potentials by B
- effect increased by AMM (dotted line)
- equation of state gets stiffer with B



☆ Ingredients

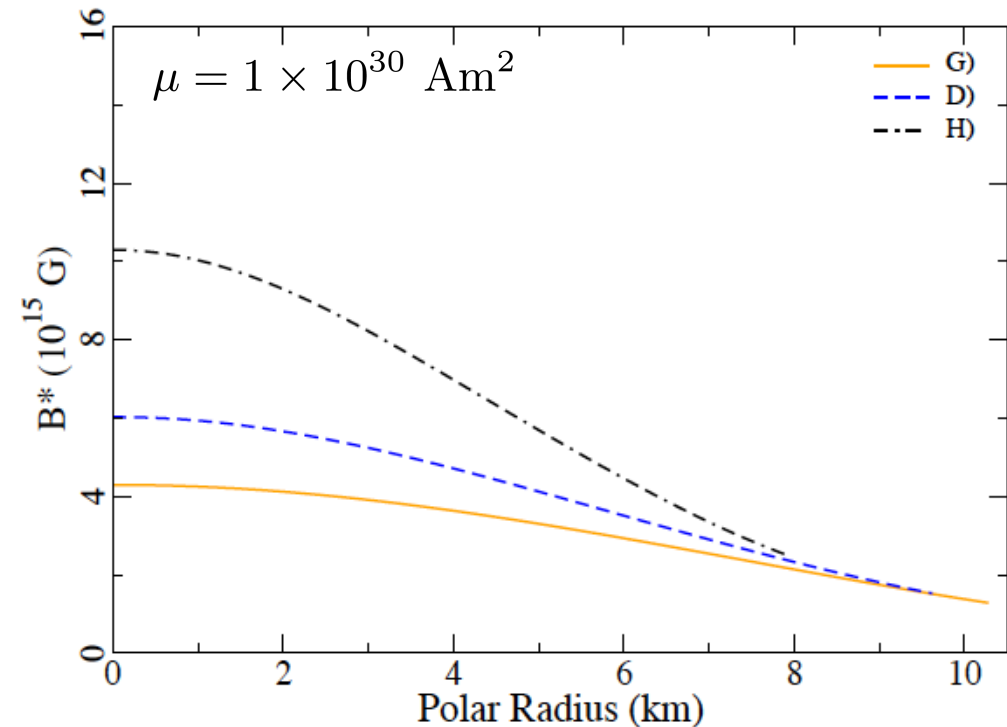
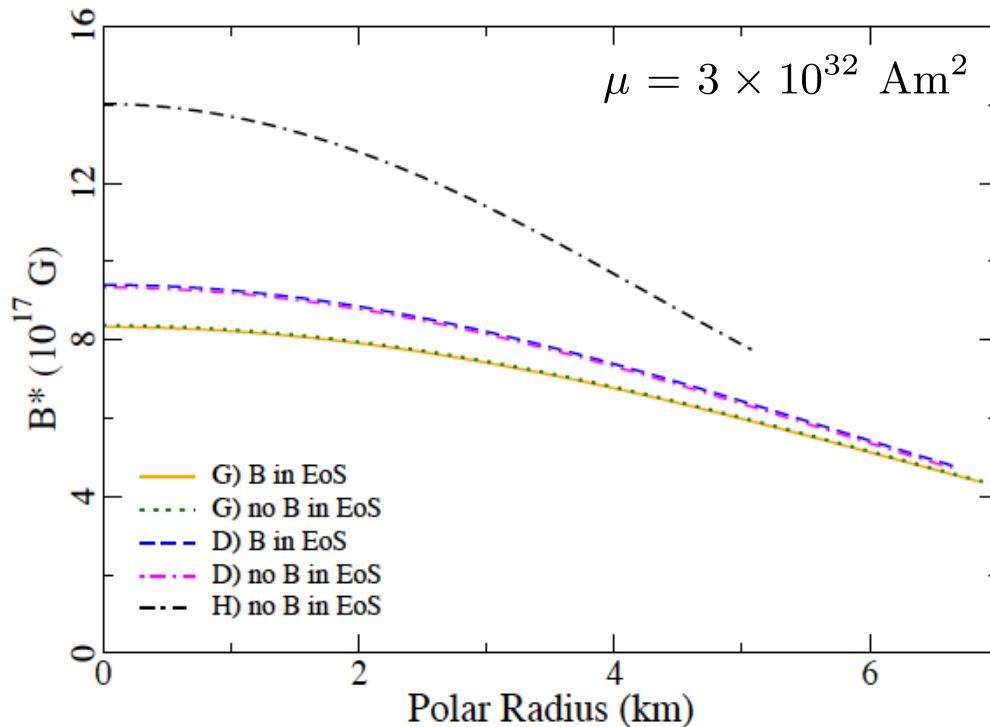
- EoS's:

- include magnetic field effects
- respect nuclear and astro constraints
- possess different degrees of freedom
 - 1) hadronic: **G-model** (many-body forces (MBF) among nucleons simulated by non-linear self-couplings)
 - 2) hybrid: **D-model** (chiral (CMF) model with nucleons, hyperons, quarks)
 - 3) quarks: **H-model** (3-flavor NJL model with vector-isoscalar interaction)

- General relativity:

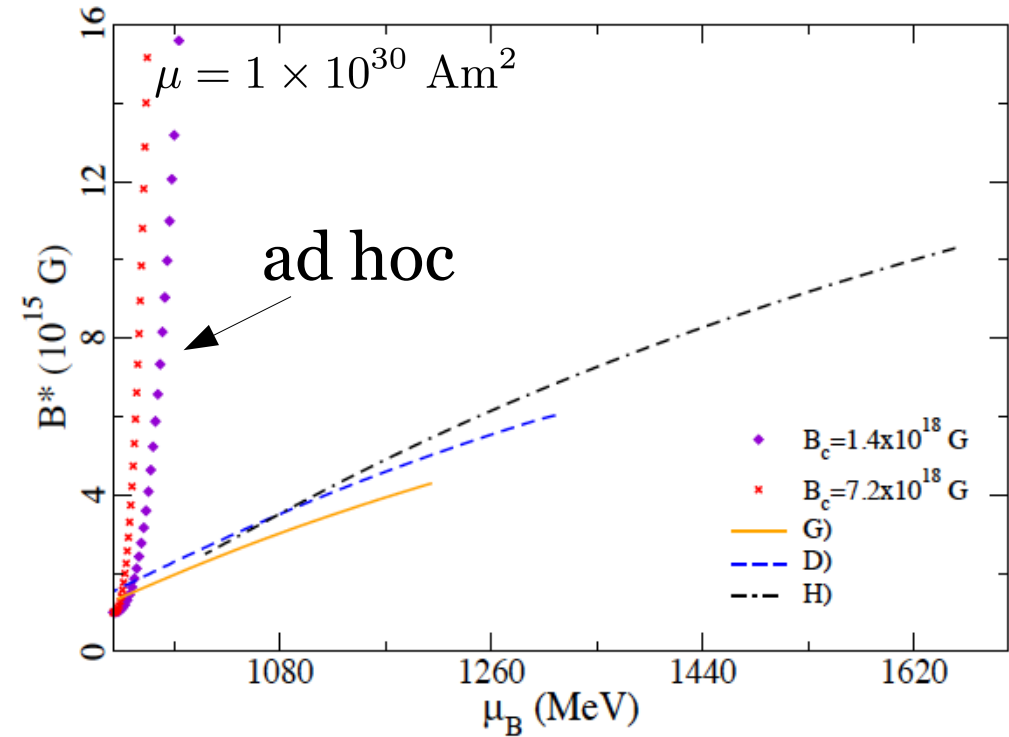
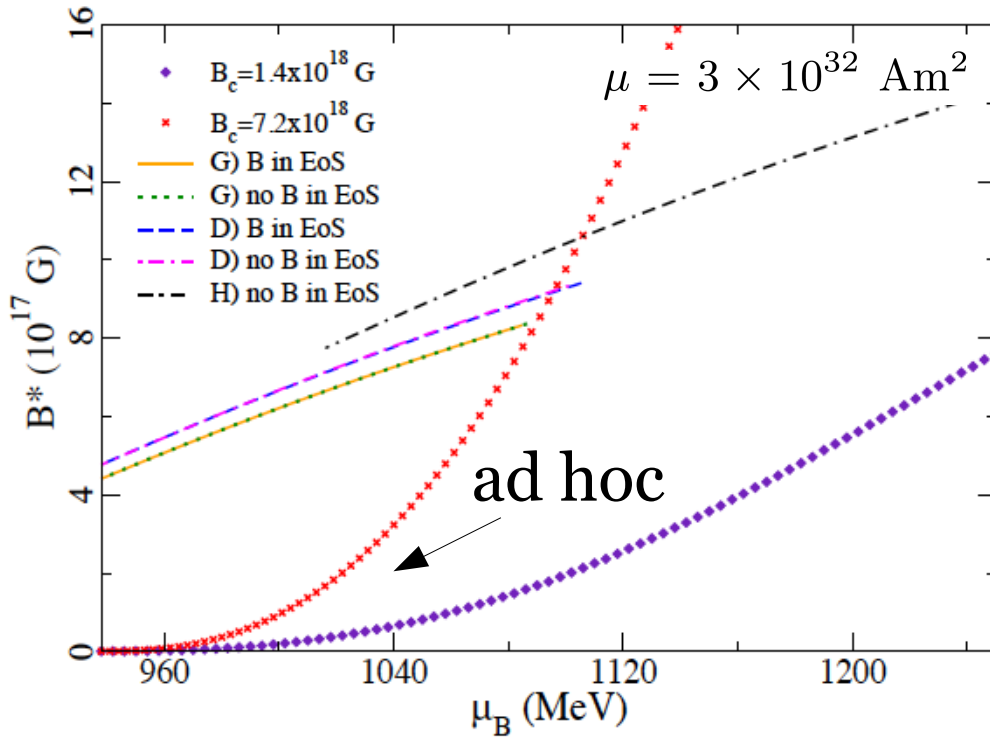
- equilibrium configurations from Einstein-Maxwell's field equations in spherical polar coordinates
- assumes a poloidal magnetic field configuration produced self-consistently by a macroscopic current (stellar radius, angle theta, and dipole magnetic moment μ)
- LORENE C++ class library for numerical relativity

★ B-profile versus macroscopic quantity in a massive star



- B in EoS makes very little or no difference in profile
- different EoS's show different magnetic field strengths, but have approximately the same profile shape (for any μ)

★ B-profile versus macroscopic quantity in a massive star



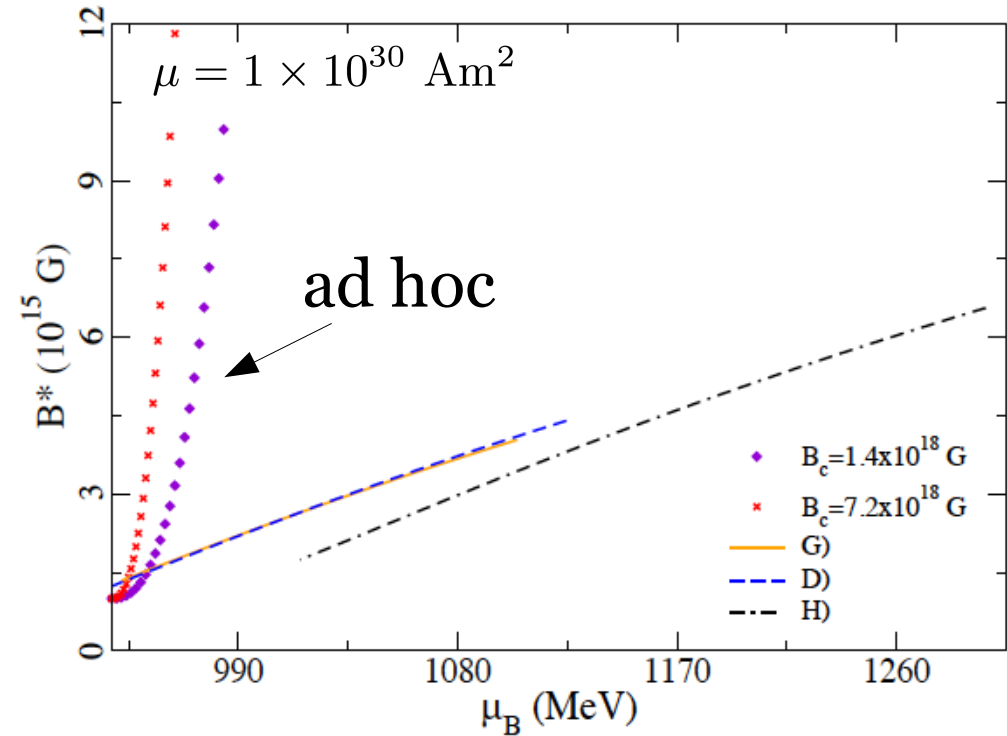
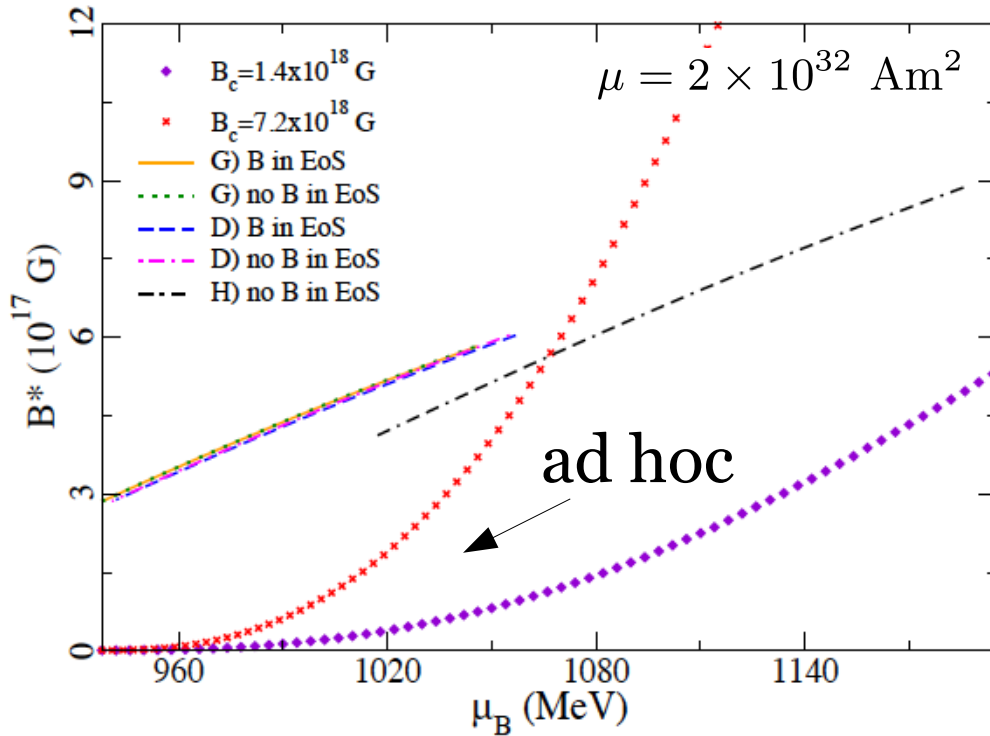
- the B profile is quadratic with respect to chemical potential

$$B^*(\mu_B) = (a + b\mu_B + c\mu_B^2) \mu, \text{ with } a = -1.68 \times 10^{-14},$$

$$b = 2.80 \times 10^{-17}, \quad c = -8.92 \times 10^{-21} \text{ and}$$

μ_B given in MeV and μ in A.m² to produce B* in G

★ B-profile versus macroscopic quantity in a lower mass star



- the B profile is still quadratic with respect to chemical potential but with coefficients $a = -2.60 \times 10^{-14}$, $b = 4.16 \times 10^{-17}$ and $c = -1.35 \times 10^{-20}$
- nuclear EoS's overlap

☆ Summary and Outlook

- we provide the first realistic magnetic-field profile for the EoS of magnetized neutron stars
- our profile is obtained from the solution of Einstein's equations and does not violate Maxwell's equations
- our results allow anyone to include a B profile in any neutron star EoS in a simple way to study changes in stiffness, population, phase transitions, temperature, transport properties, etc. due to B effects in their models
- we are currently analyzing the effects of the B profile in a phase diagram for neutron star matter



★ Magnetic Field in EOS at T=0

- B in the z-direction
- x, y energy levels quantized
- anomalous magnetic moment (AMM)

$$E_{i\nu s}^* = \sqrt{k_{z_i}^2 + \left(\sqrt{m_i^{*2} + 2\nu|q_i|B^*} - s_i\kappa_i B^* \right)^2}$$

$$E_{i_s}^* = \sqrt{k_i^2 + (m_i^{*2} - s_i\kappa_i B^*)^2}$$

$\kappa_i \rightarrow$ coupling strength of baryons
with electromagnetic field tensor

maximum
Landau level

$$\nu_{max} = \frac{E_{i_s}^{*2} + s_i\kappa_i B^* - m_i^{*2}}{2|q_i|B^*}$$