

**You can hide  
but you have to run:  
direct detection with vector mediators**

**Francesco D'Eramo**



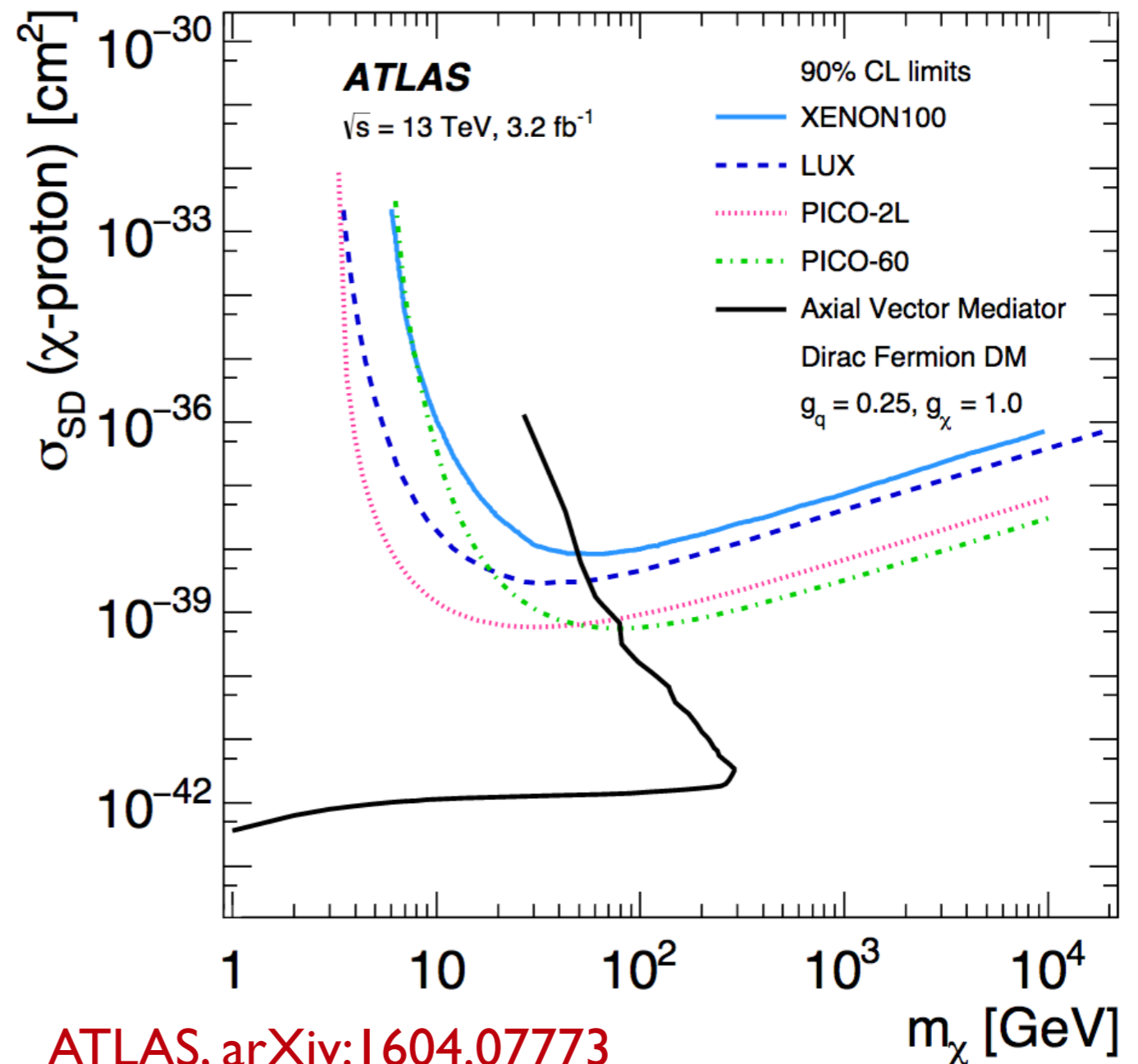
**DMWG - 22 June 2016**

# LHC vs Direct Detection

Two probes of DM couplings at vastly different energies

How are LHC limits translated into the  $(m_{\text{DM}}, \sigma_{\text{SD}})$  plane?

$$\mathcal{L} = g_\chi A_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi + g_q A_\mu \bar{q} \gamma^\mu \gamma^5 q$$



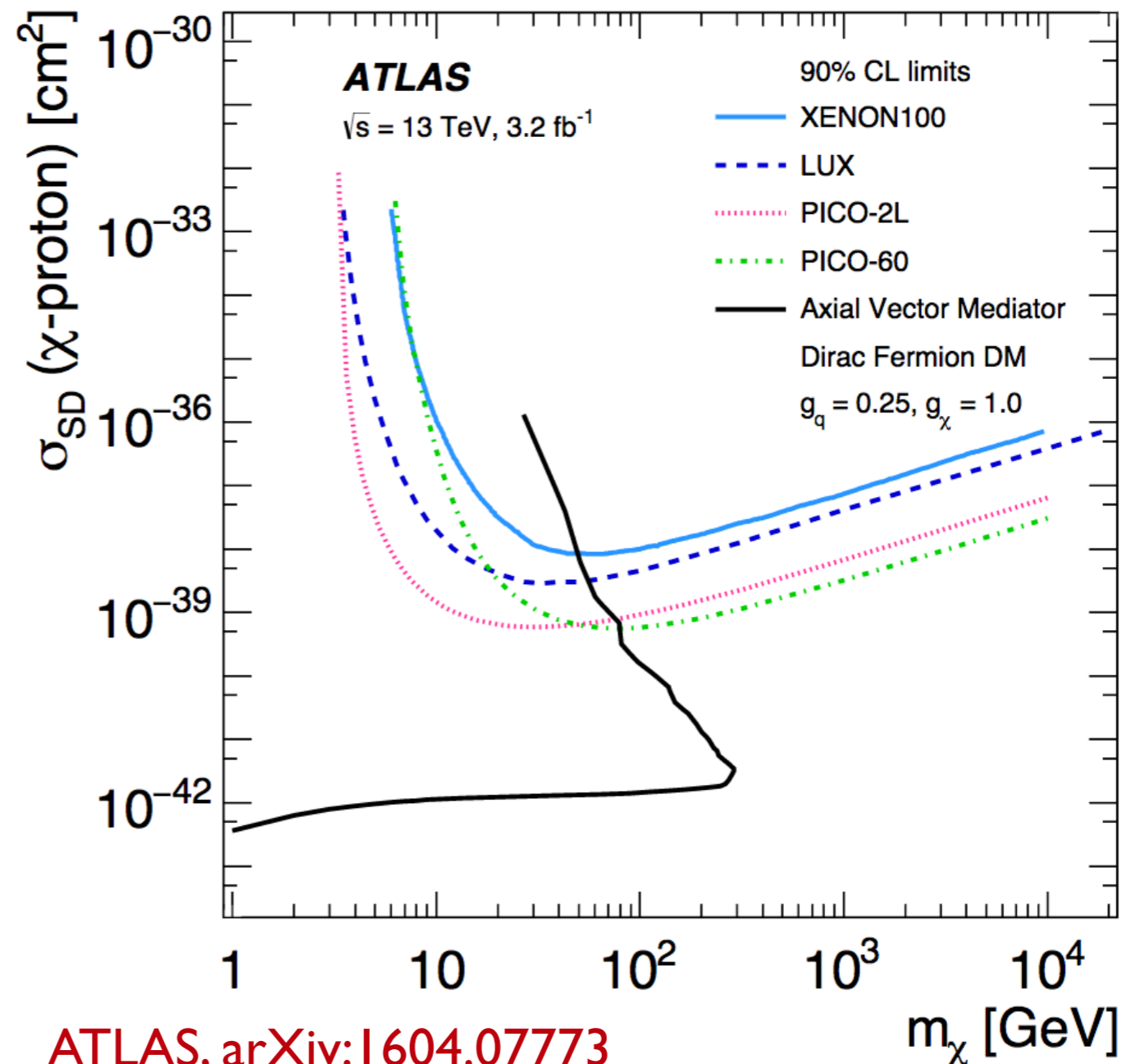
ATLAS, arXiv:1604.07773

# LHC vs Direct Detection

LHC bounds have to be evolved down to the direct detection scale

You have to run!  
(Renormalization Group Evolution — RGE)

$$\mathcal{L} = g_\chi A_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi + g_q A_\mu \bar{q} \gamma^\mu \gamma^5 q$$

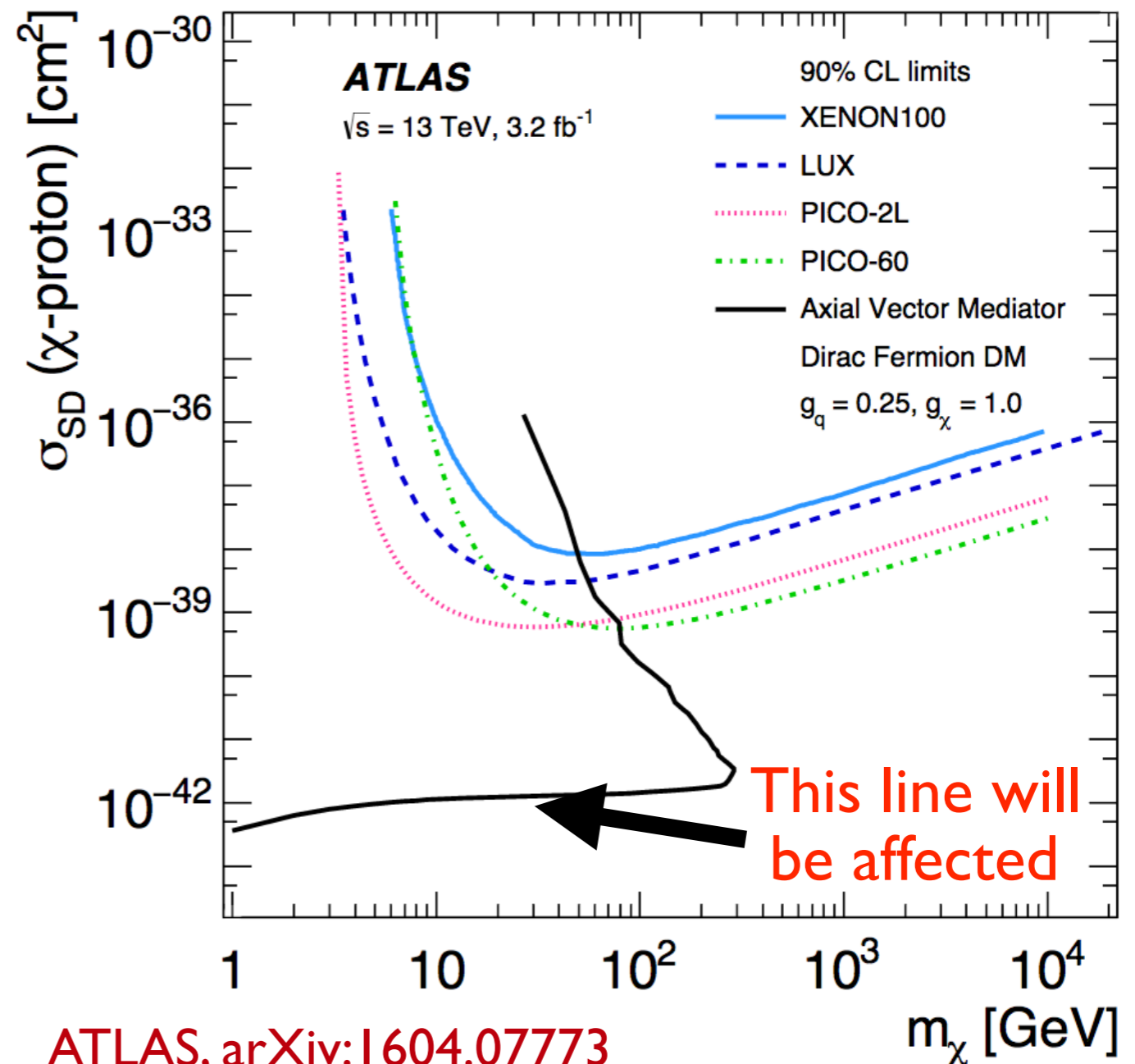


# LHC vs Direct Detection

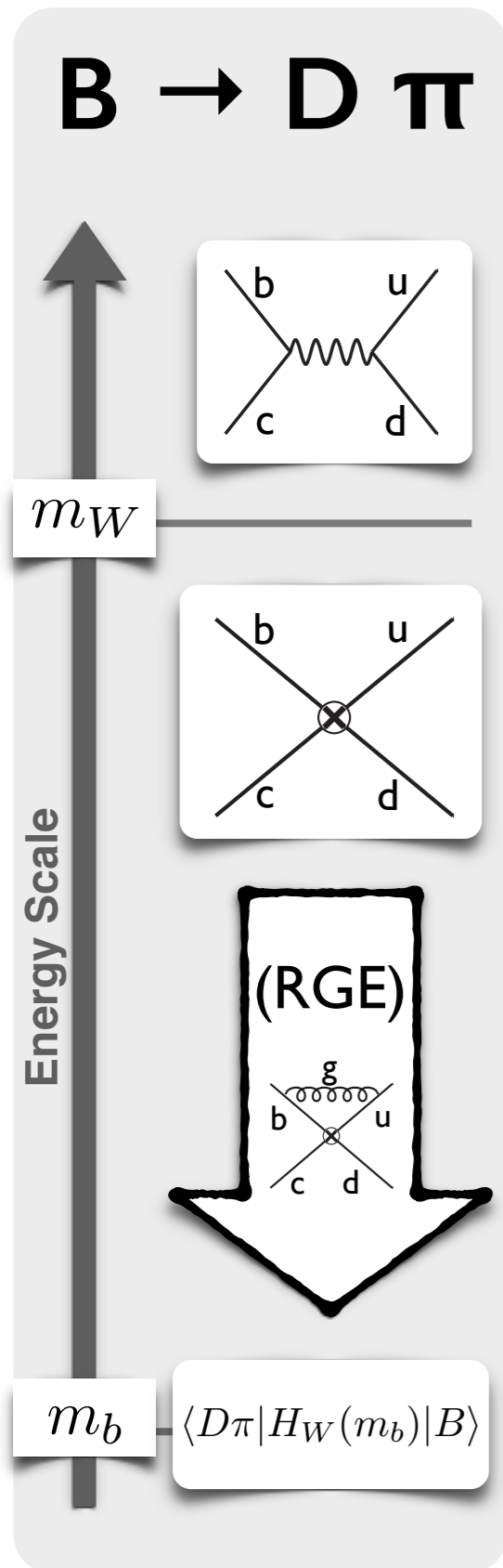
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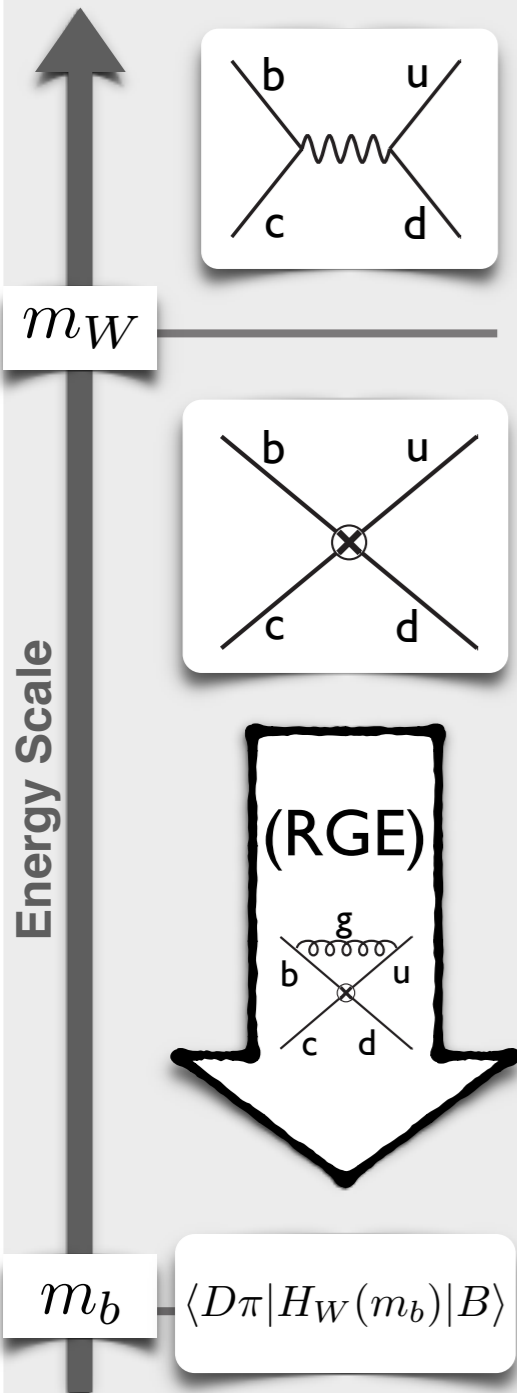


# DM Direct Detection

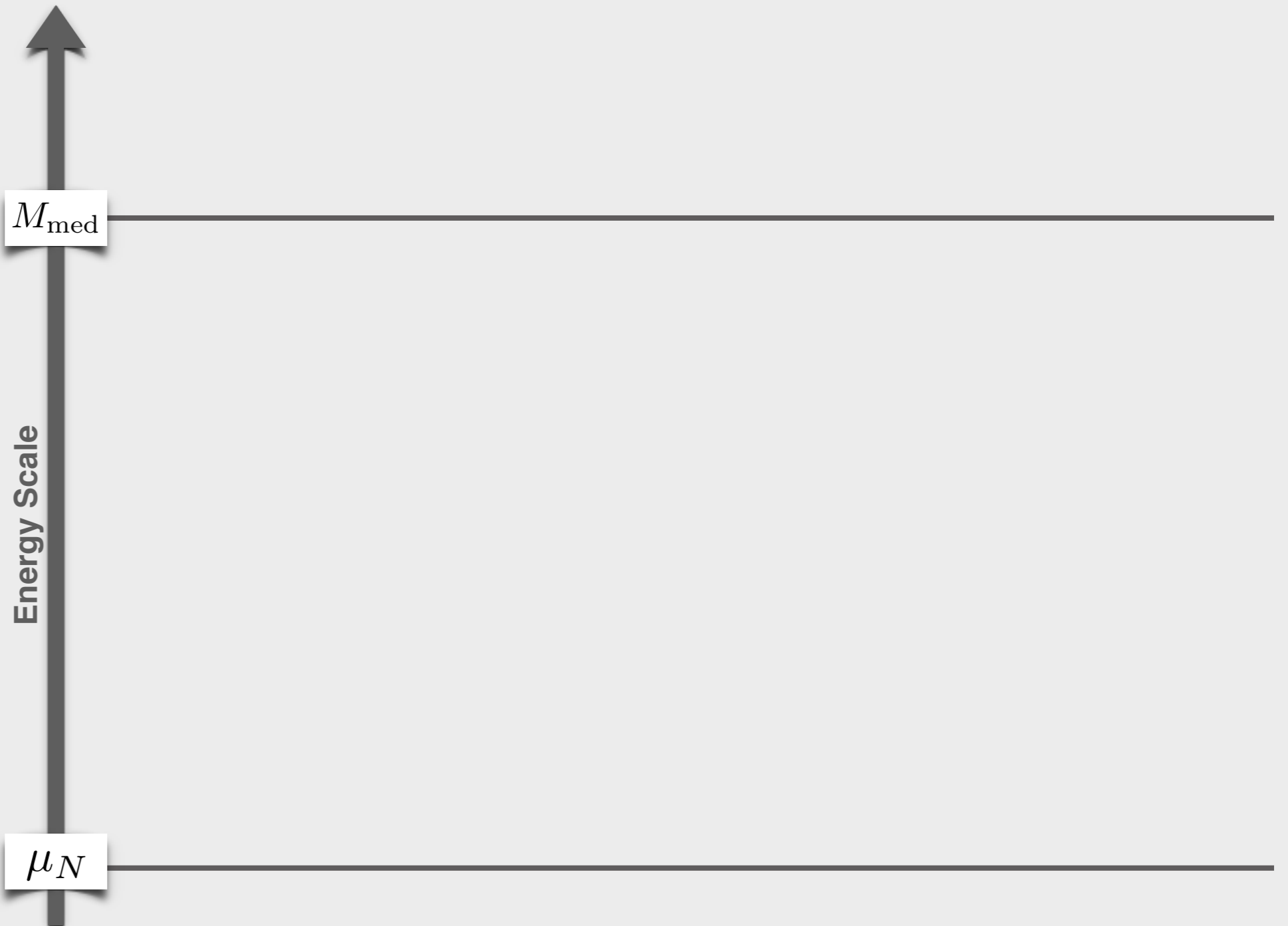


# DM Direct Detection

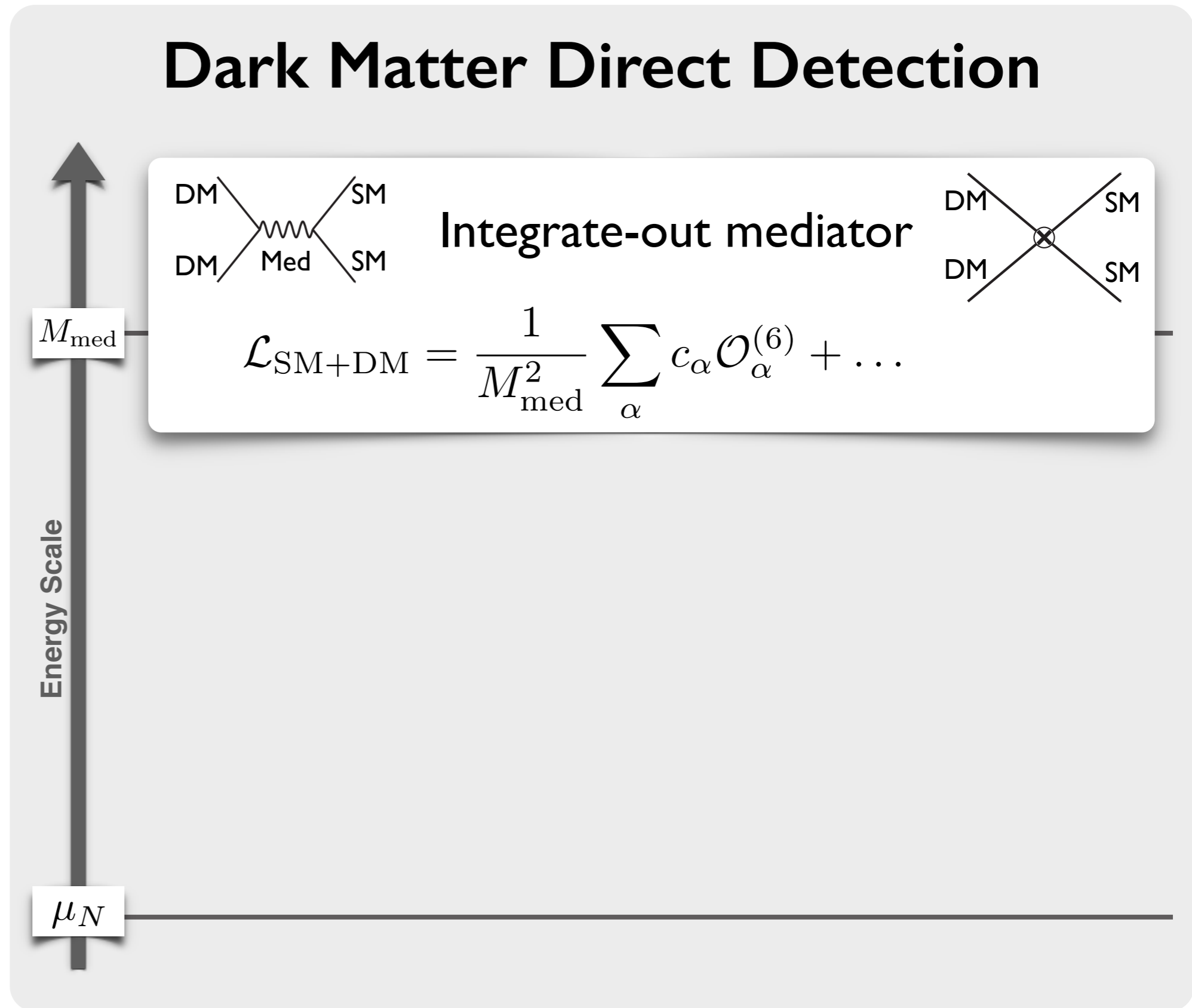
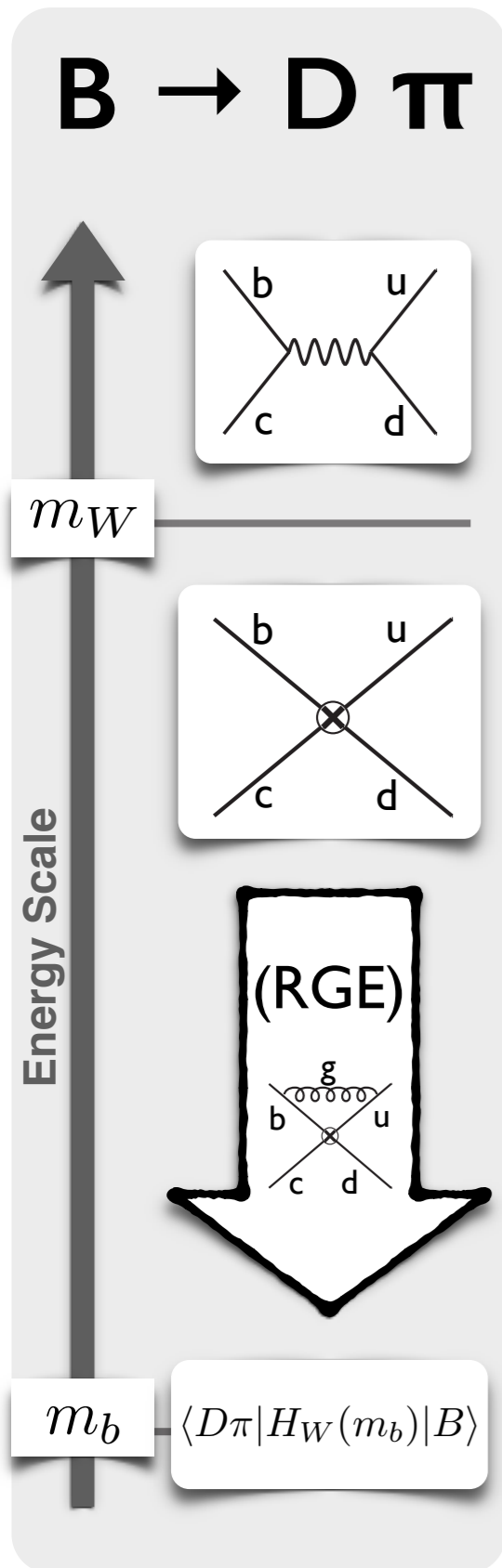
$B \rightarrow D \pi$



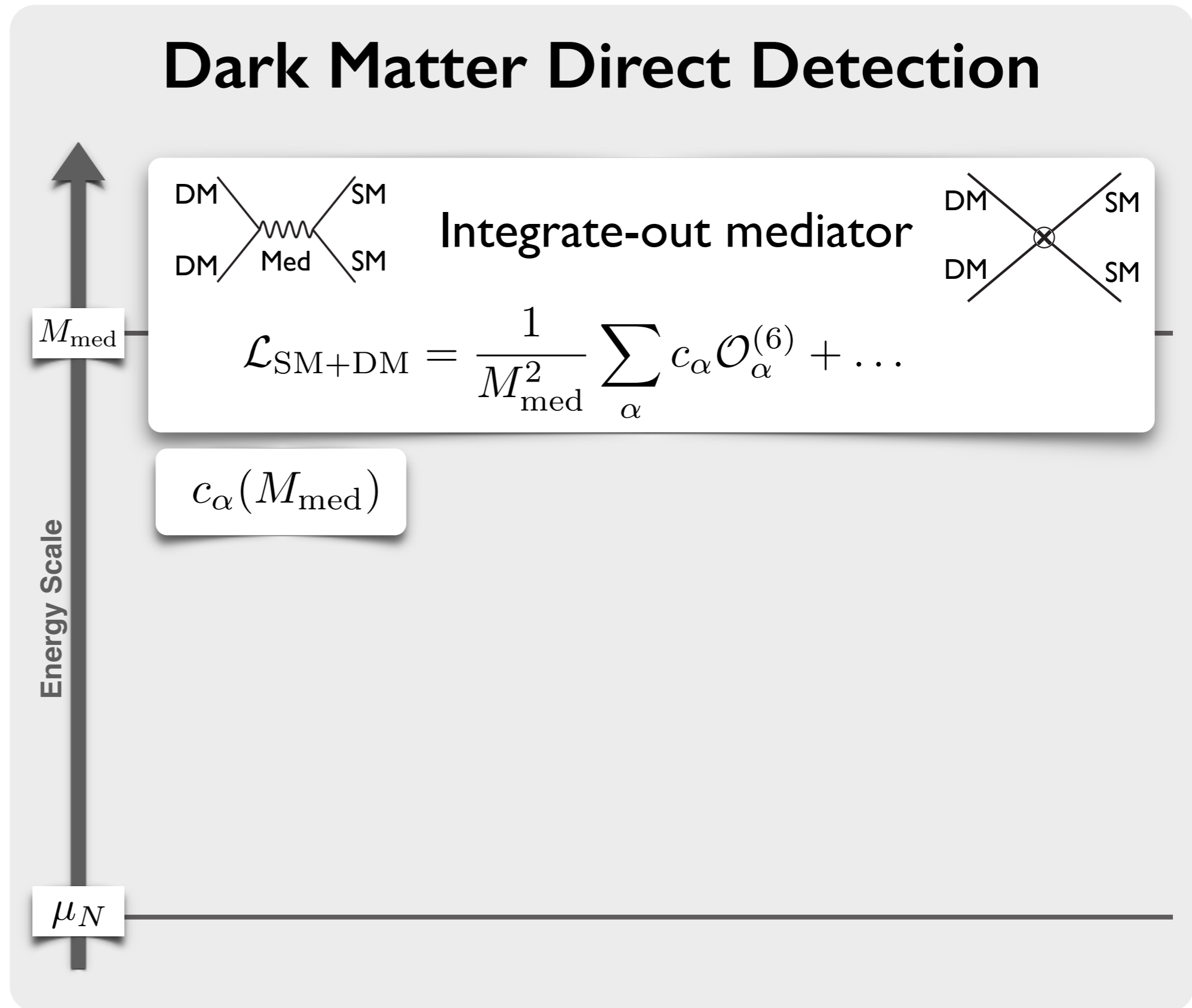
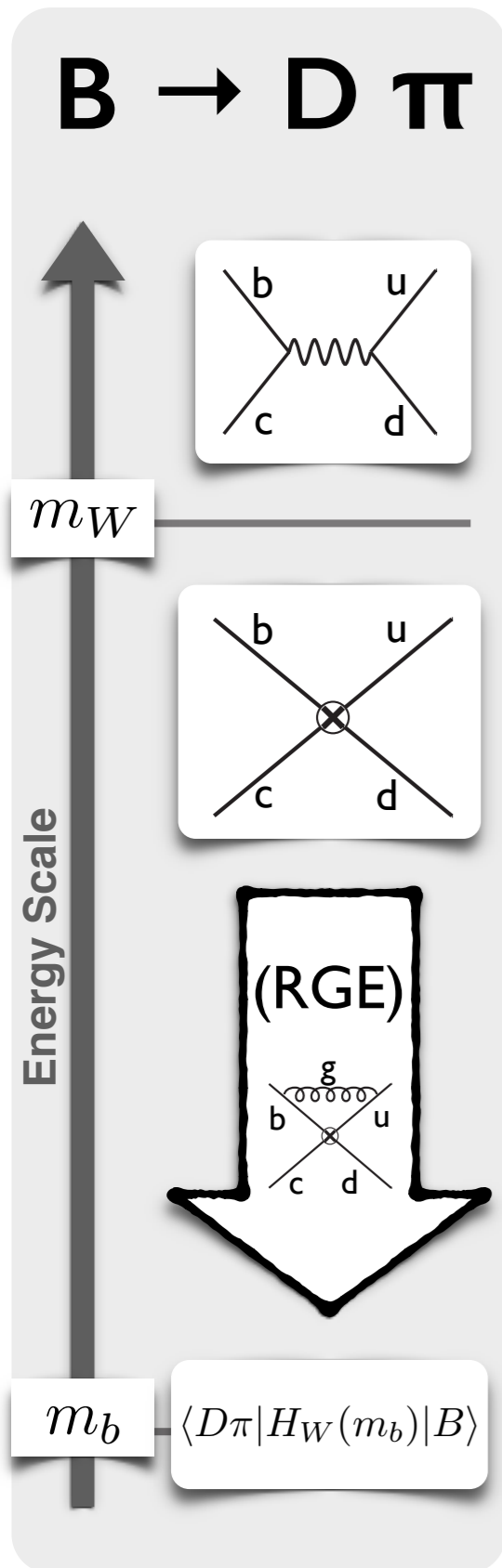
Dark Matter Direct Detection



# DM Direct Detection

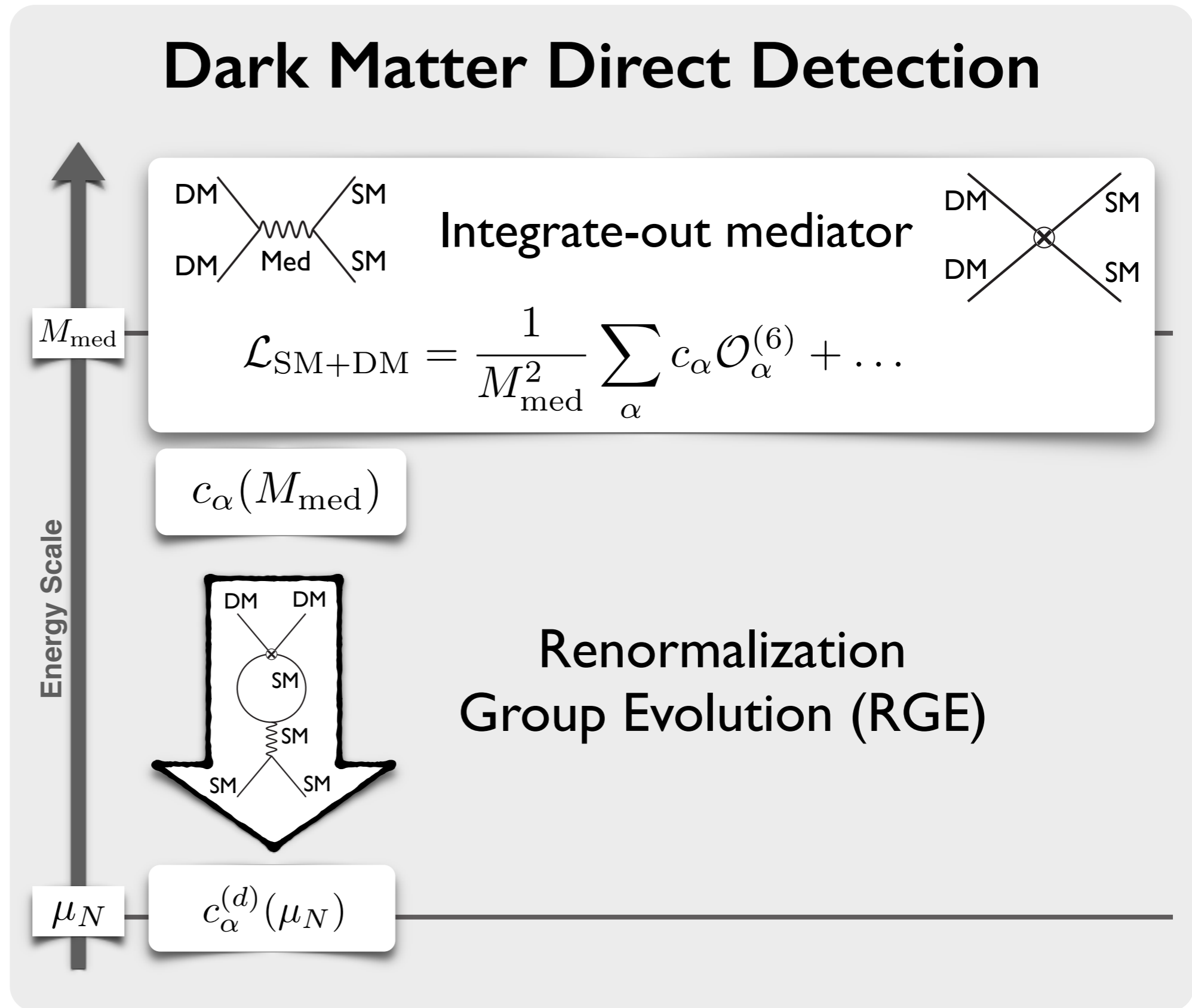
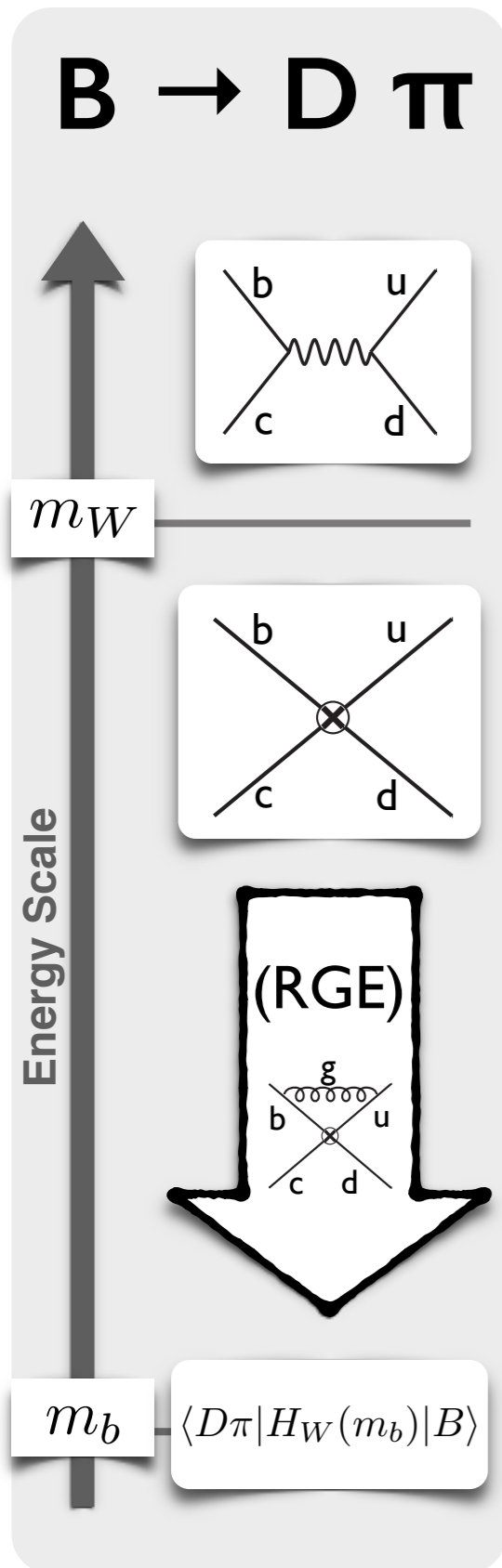


# DM Direct Detection

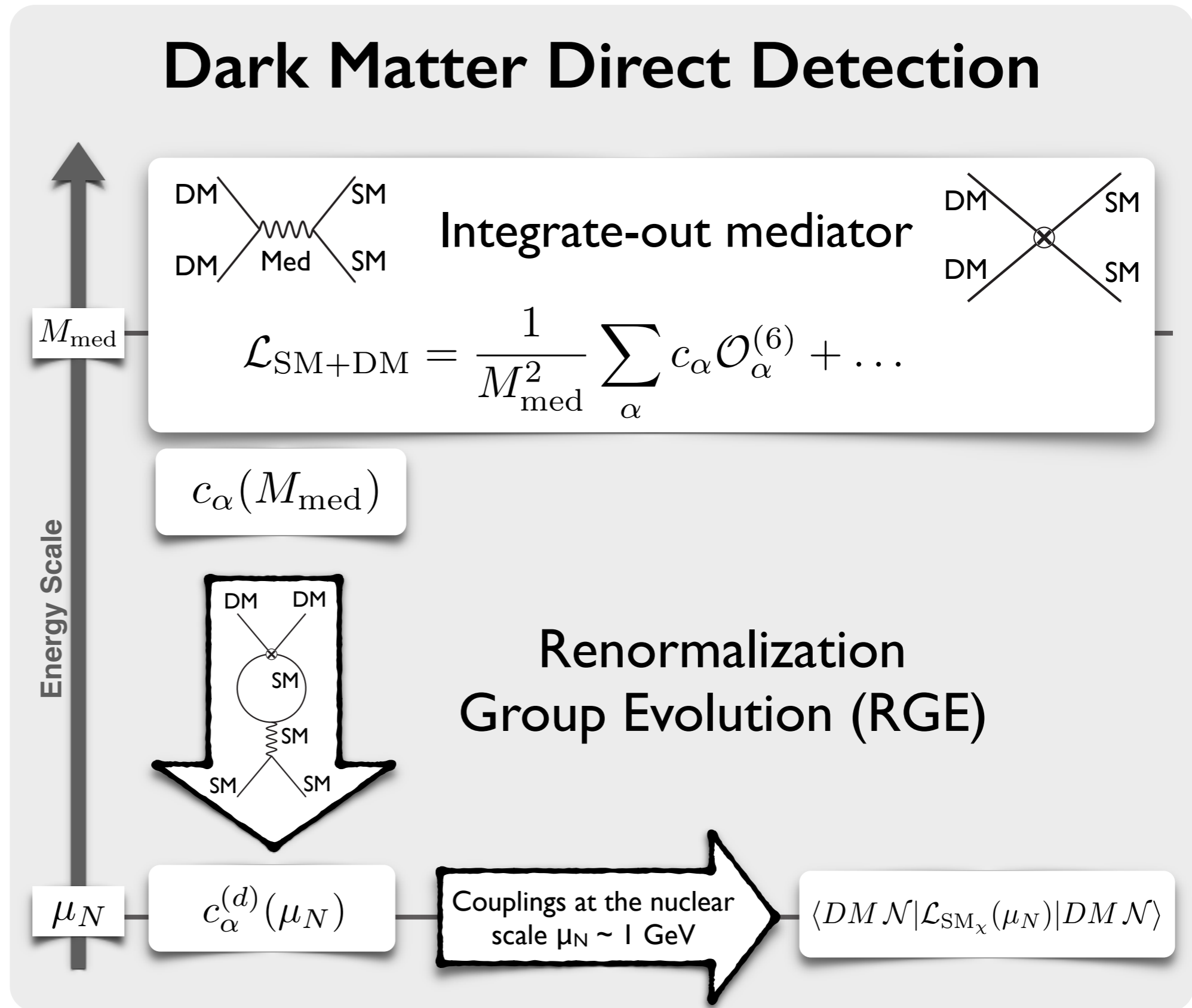
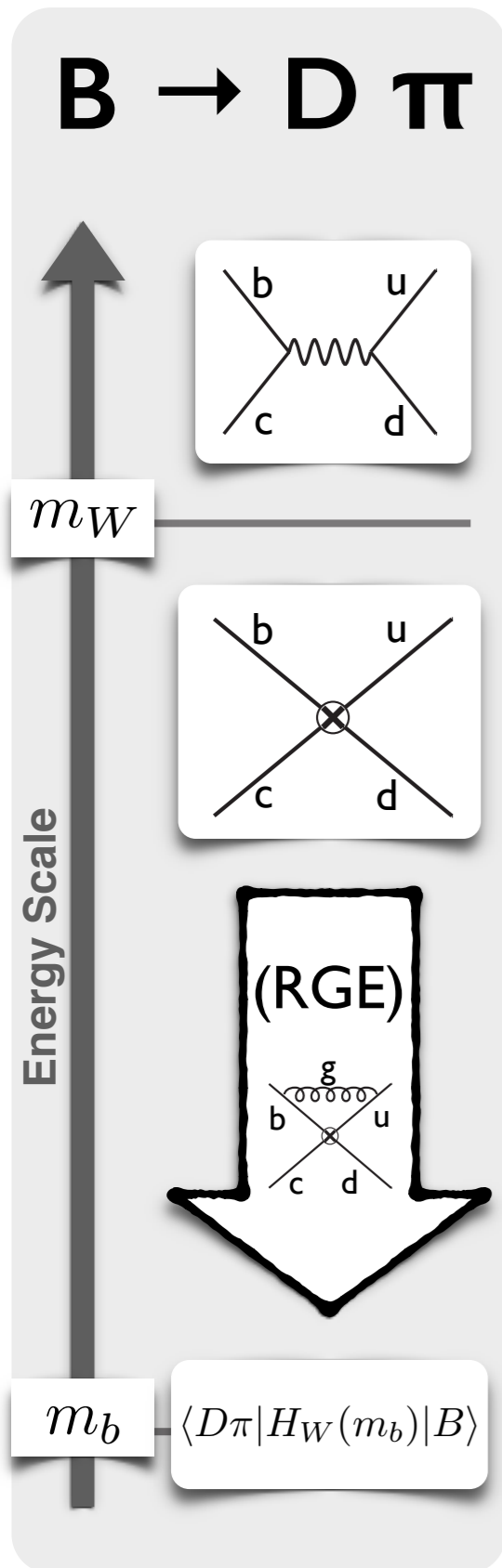




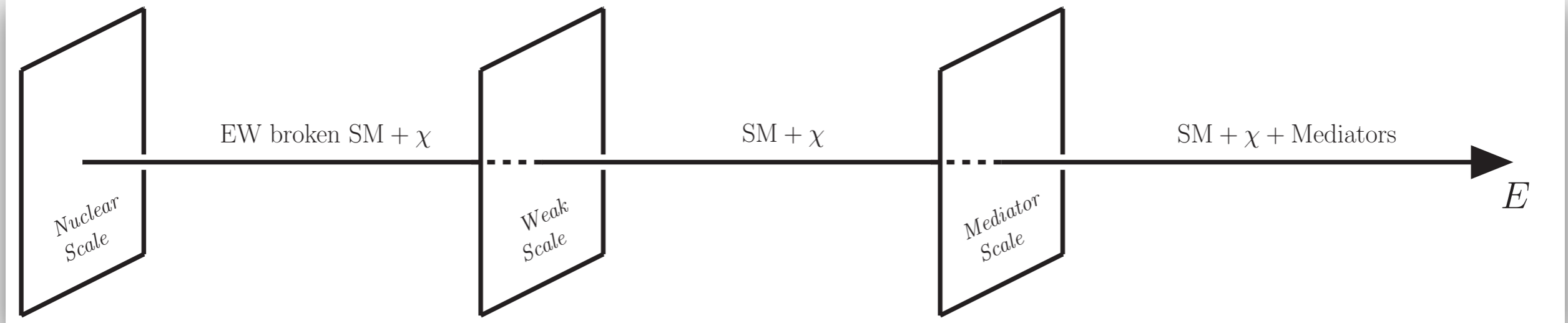
# DM Direct Detection



# DM Direct Detection



# Vector Mediators RGE



Crivellin, FD, Procura, Phys.Rev.Lett. 112 (2014), arXiv:1402.1173  
FD, Procura, JHEP1504 (2015), arXiv:1411.3342



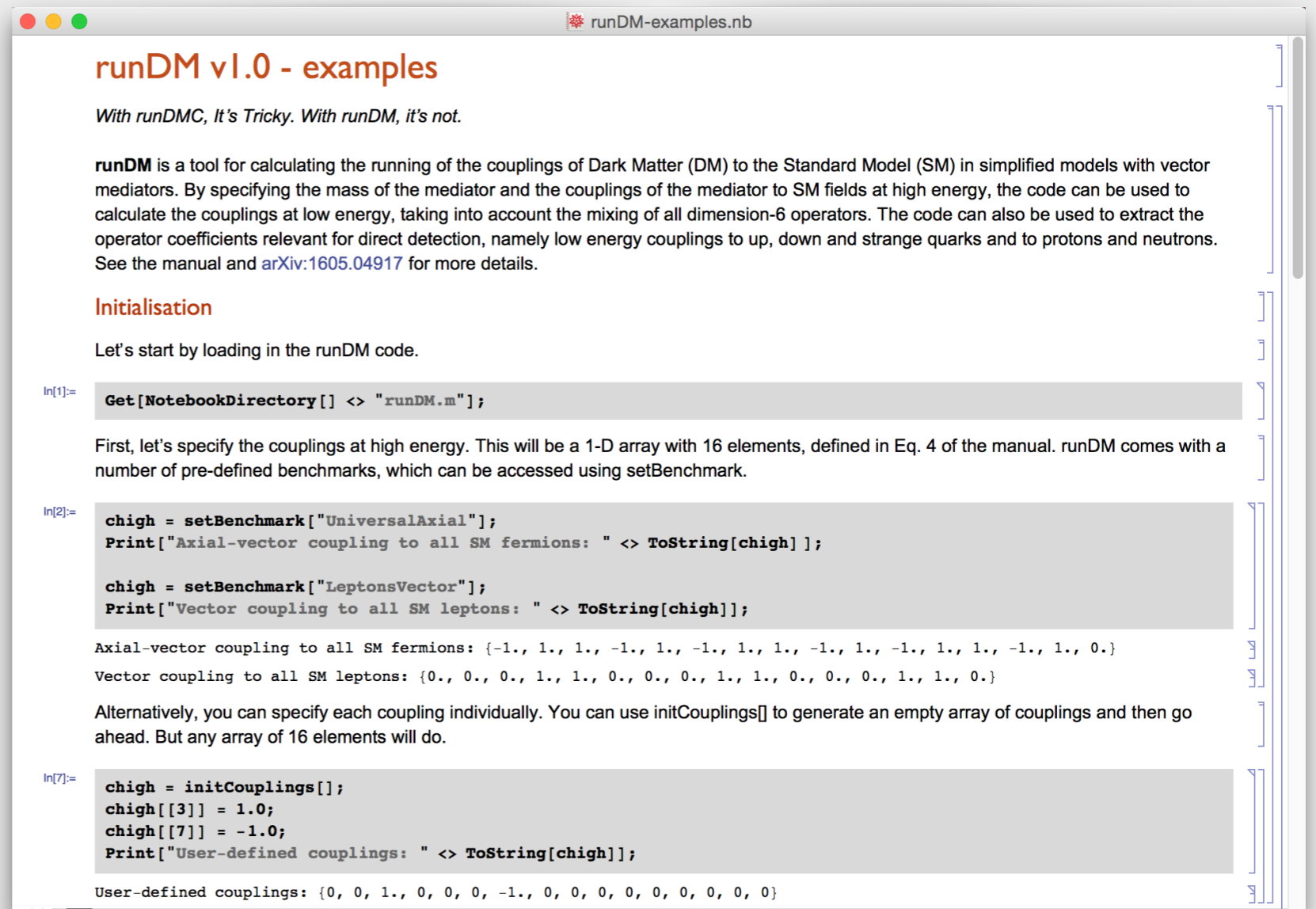


# runDM: code for RGE

## Inclusion of RGE effects automatic

### INPUT:

Effective couplings  
at an arbitrary  
energy scale



```
runDM v1.0 - examples

With runDMC, It's Tricky. With runDM, it's not.

runDM is a tool for calculating the running of the couplings of Dark Matter (DM) to the Standard Model (SM) in simplified models with vector mediators. By specifying the mass of the mediator and the couplings of the mediator to SM fields at high energy, the code can be used to calculate the couplings at low energy, taking into account the mixing of all dimension-6 operators. The code can also be used to extract the operator coefficients relevant for direct detection, namely low energy couplings to up, down and strange quarks and to protons and neutrons. See the manual and arXiv:1605.04917 for more details.

Initialisation

Let's start by loading in the runDM code.

In[1]:= Get[NotebookDirectory[] <> "runDM.m"];

First, let's specify the couplings at high energy. This will be a 1-D array with 16 elements, defined in Eq. 4 of the manual. runDM comes with a number of pre-defined benchmarks, which can be accessed using setBenchmark.

In[2]:= chigh = setBenchmark["UniversalAxial"];
Print["Axial-vector coupling to all SM fermions: " <> ToString[chigh]];

chigh = setBenchmark["LeptonsVector"];
Print["Vector coupling to all SM leptons: " <> ToString[chigh]];

Axial-vector coupling to all SM fermions: {-1., 1., 1., -1., 1., -1., 1., 1., -1., 1., -1., 1., 1., -1., 1., 0.}
Vector coupling to all SM leptons: {0., 0., 0., 1., 1., 0., 0., 0., 1., 1., 0., 0., 0., 1., 1., 0.}

Alternatively, you can specify each coupling individually. You can use initCouplings[] to generate an empty array of couplings and then go ahead. But any array of 16 elements will do.

In[7]:= chigh = initCouplings[];
chigh[[3]] = 1.0;
chigh[[7]] = -1.0;
Print["User-defined couplings: " <> ToString[chigh]];

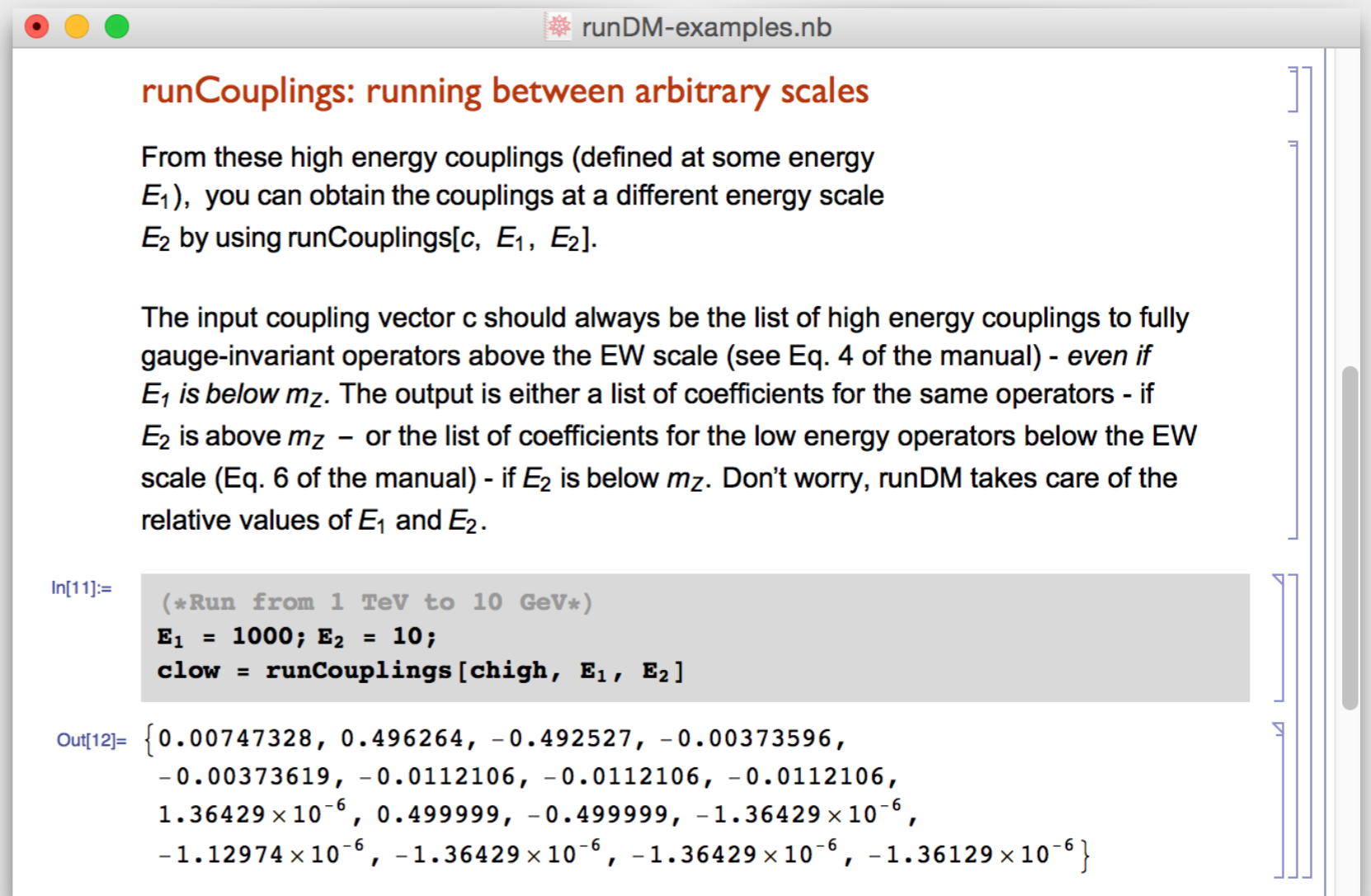
User-defined couplings: {0, 0, 1., 0, 0, 0, -1., 0, 0, 0, 0, 0, 0, 0, 0, 0}
```

# runDM: code for RGE

## Inclusion of RGE effects automatic

### OUTPUT I:

RG evolved couplings at a second arbitrary energy scale (useful for future ID studies)



**runCouplings: running between arbitrary scales**

From these high energy couplings (defined at some energy  $E_1$ ), you can obtain the couplings at a different energy scale  $E_2$  by using `runCouplings[c,  $E_1$ ,  $E_2$ ]`.

The input coupling vector `c` should always be the list of high energy couplings to fully gauge-invariant operators above the EW scale (see Eq. 4 of the manual) - *even if  $E_1$  is below  $m_Z$* . The output is either a list of coefficients for the same operators - if  $E_2$  is above  $m_Z$  - or the list of coefficients for the low energy operators below the EW scale (Eq. 6 of the manual) - if  $E_2$  is below  $m_Z$ . Don't worry, runDM takes care of the relative values of  $E_1$  and  $E_2$ .

```
In[11]:= (*Run from 1 TeV to 10 GeV*)
E1 = 1000; E2 = 10;
cLow = runCouplings[cHigh, E1, E2]
```

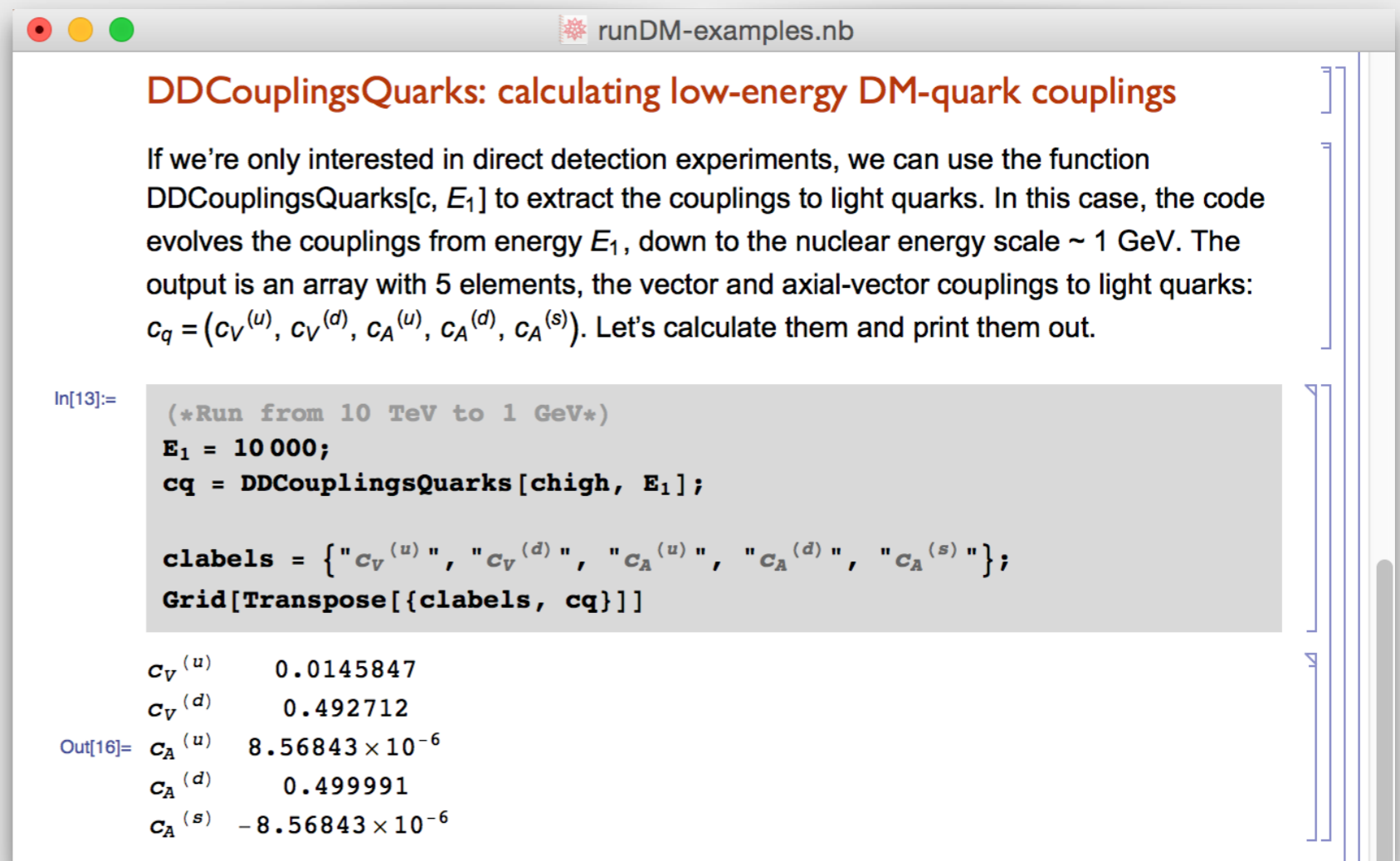
```
Out[12]= {0.00747328, 0.496264, -0.492527, -0.00373596,
-0.00373619, -0.0112106, -0.0112106, -0.0112106,
1.36429 × 10-6, 0.499999, -0.499999, -1.36429 × 10-6,
-1.12974 × 10-6, -1.36429 × 10-6, -1.36429 × 10-6, -1.36129 × 10-6}
```

# runDM: code for RGE

## Inclusion of RGE effects automatic

### OUTPUT II:

RG evolved  
couplings for  
direct detection



**DDCouplingsQuarks: calculating low-energy DM-quark couplings**

If we're only interested in direct detection experiments, we can use the function `DDCouplingsQuarks[c, E1]` to extract the couplings to light quarks. In this case, the code evolves the couplings from energy  $E_1$ , down to the nuclear energy scale  $\sim 1$  GeV. The output is an array with 5 elements, the vector and axial-vector couplings to light quarks:  $c_q = (c_V^{(u)}, c_V^{(d)}, c_A^{(u)}, c_A^{(d)}, c_A^{(s)})$ . Let's calculate them and print them out.

```
In[13]:= (*Run from 10 TeV to 1 GeV*)
E1 = 10 000;
cq = DDCouplingsQuarks[chigh, E1];

clabels = {"cV(u)", "cV(d)", "cA(u)", "cA(d)", "cA(s)"};
Grid[Transpose[{clabels, cq}]]
```

```
Out[16]= cV(u)    0.0145847
cV(d)    0.492712
cA(u)    8.56843 × 10-6
cA(d)    0.499991
cA(s)   -8.56843 × 10-6
```



# Results: quarks vector

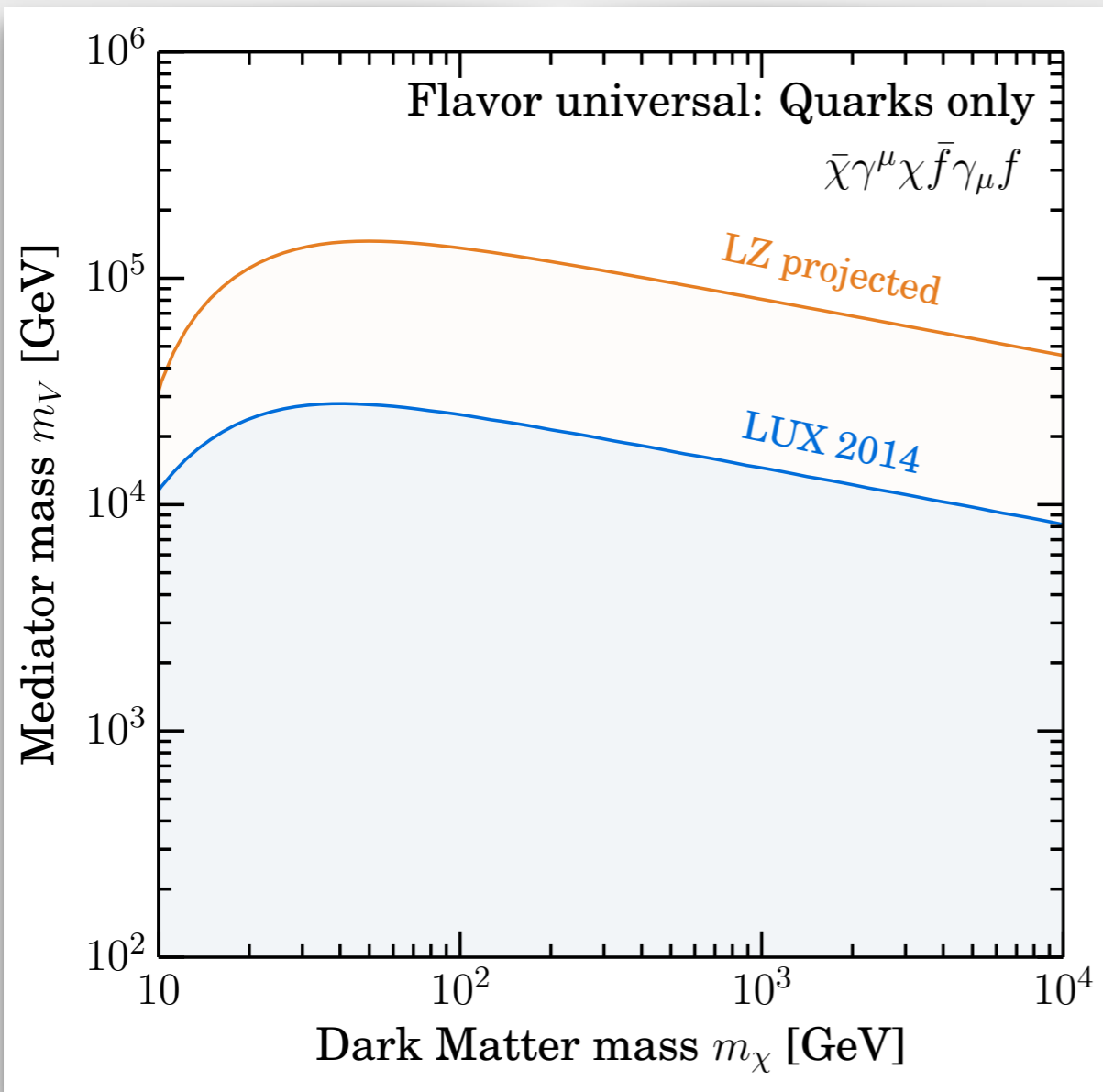
Flavor universal couplings  
to quark vector currents

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{m_V^2} J_{\text{DM} \mu} \sum_{i=1}^3 \left[ \bar{u}^i \gamma^\mu u^i + \bar{d}^i \gamma^\mu d^i \right]$$

# Results: quarks vector

Flavor universal couplings  
to quark vector currents

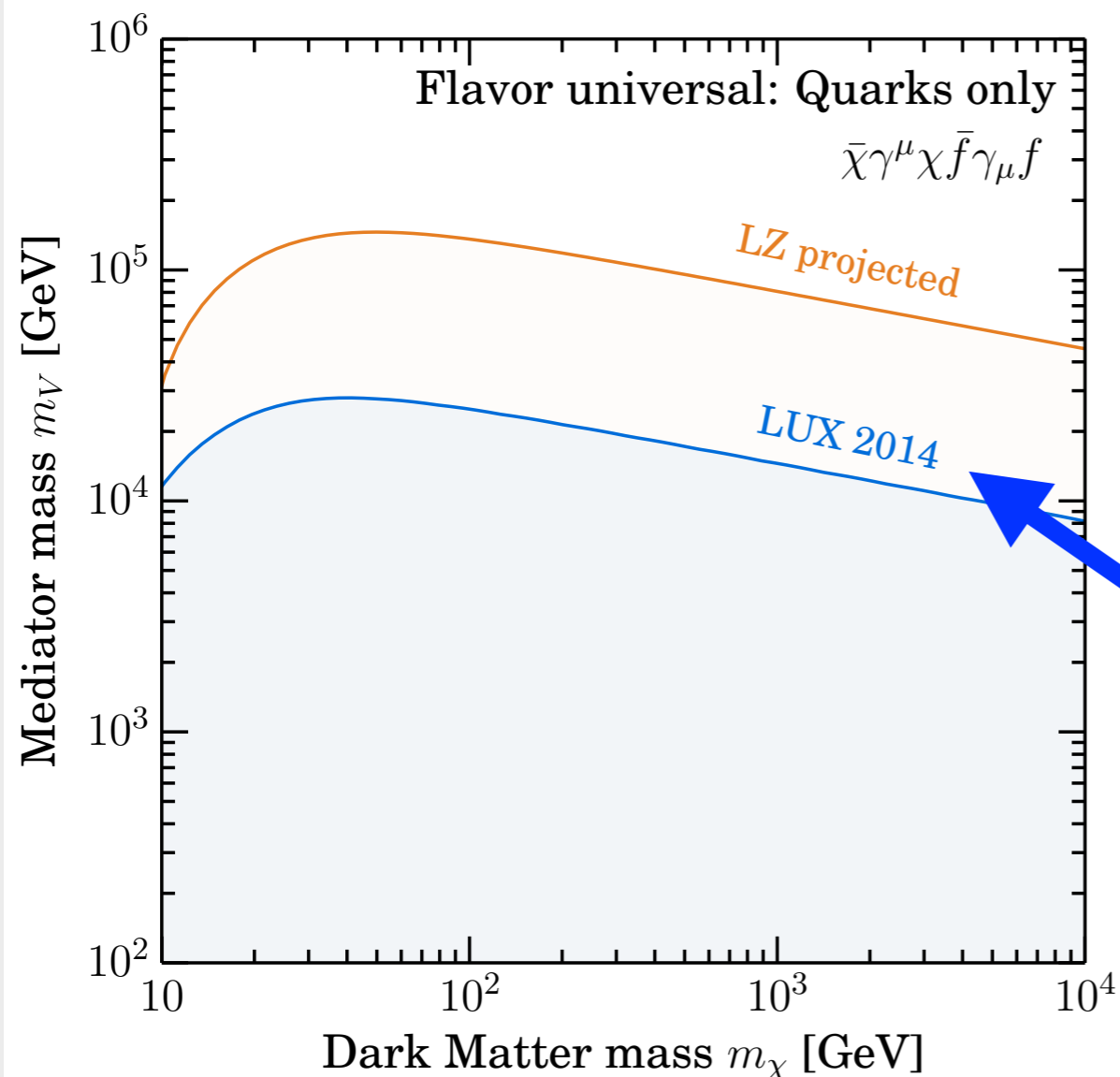
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RGE:

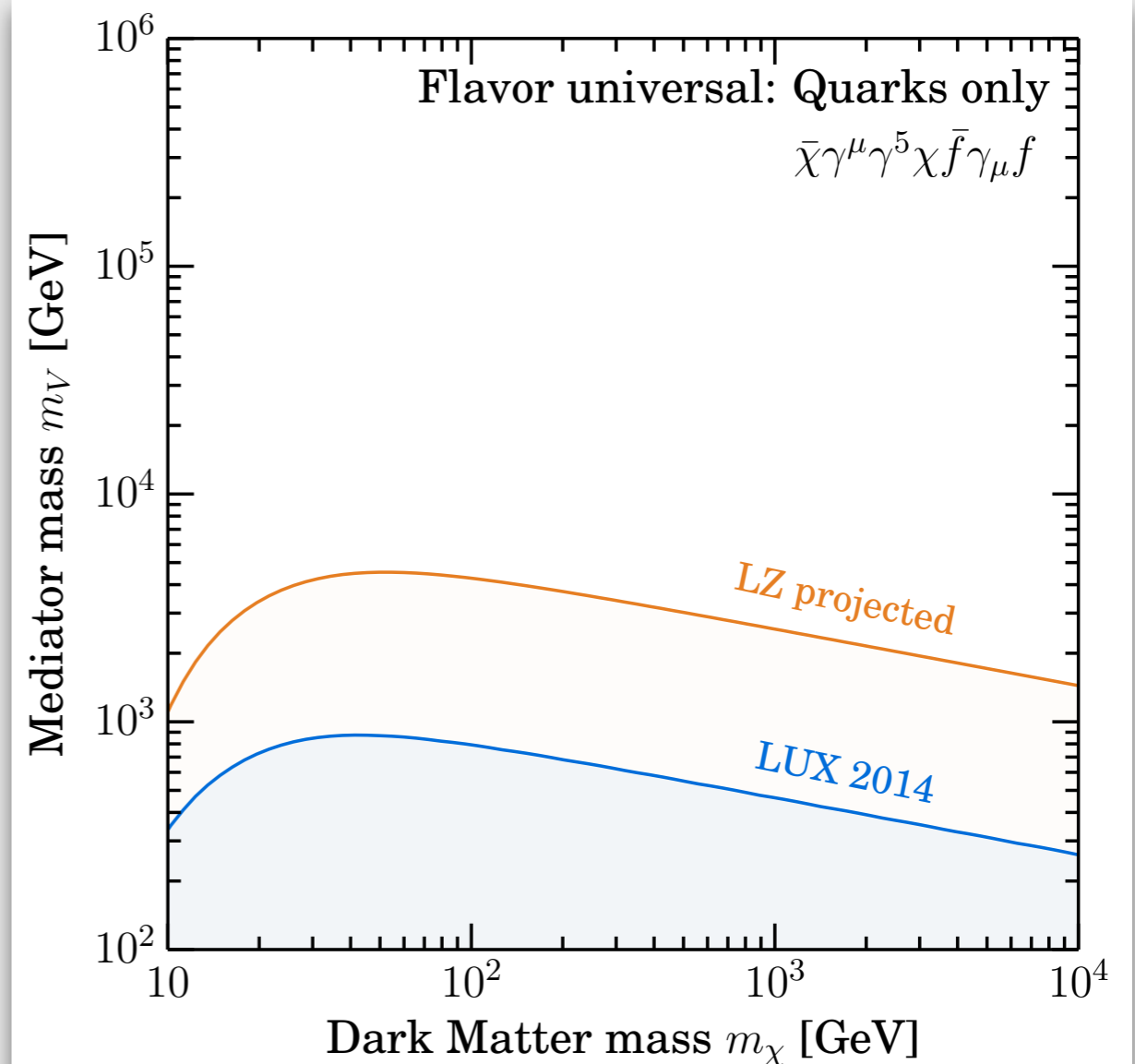
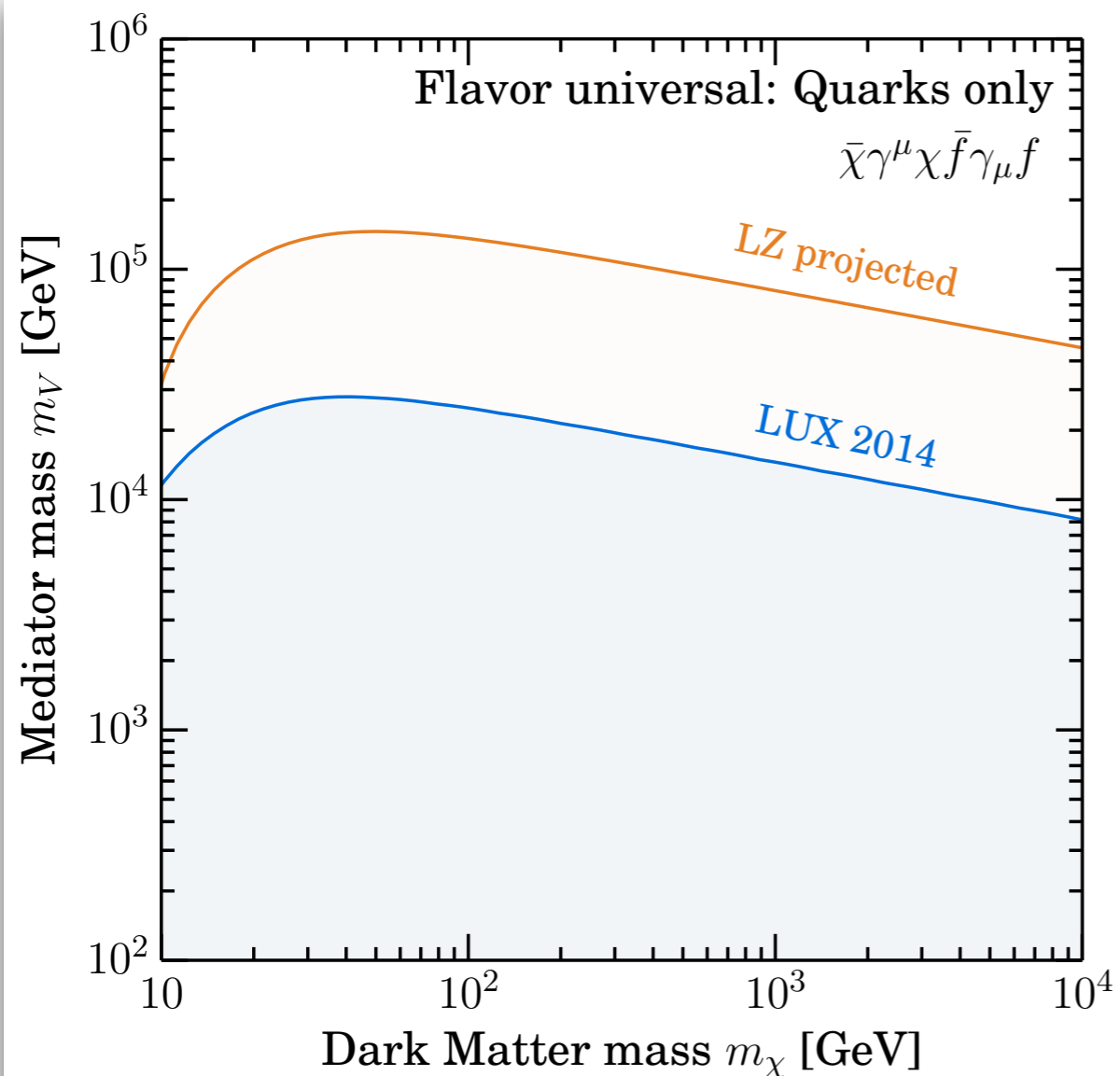
O(1%) correction  
to EFT couplings

Very strong bounds,  
meaningful results also  
for loop-induced rates

# Results: quarks vector

Flavor universal couplings  
to quark vector currents

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{m_V^2} J_{\text{DM} \mu} \sum_{i=1}^3 \left[ \bar{u}^i \gamma^\mu u^i + \bar{d}^i \gamma^\mu d^i \right]$$



# Results: quarks axial

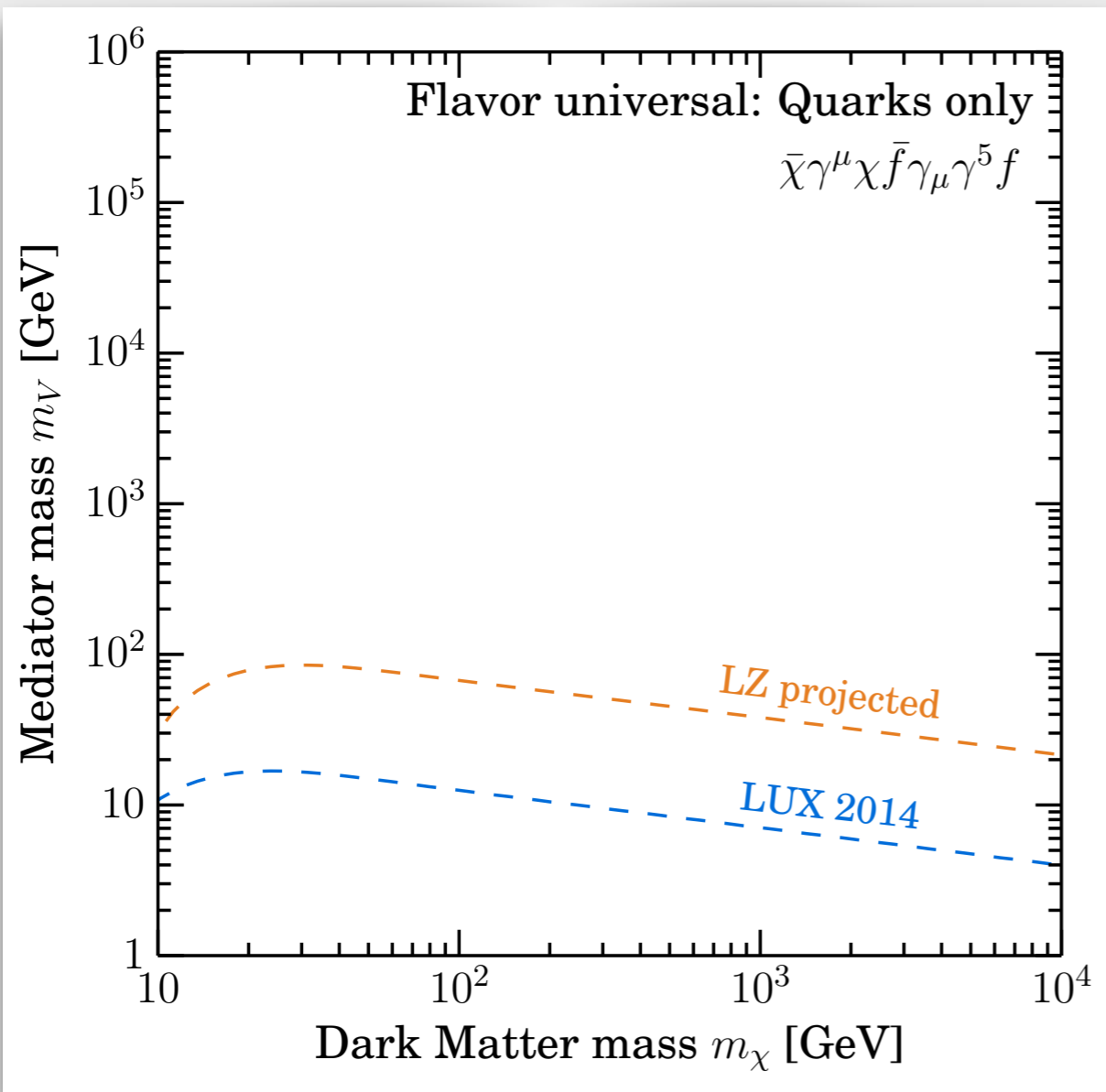
Flavor universal couplings  
to quark axial currents

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{m_V^2} J_{\text{DM}\mu} \sum_{i=1}^3 \left[ \bar{u}^i \gamma^\mu \gamma^5 u^i + \bar{d}^i \gamma^\mu \gamma^5 d^i \right]$$

# Results: quarks axial

Flavor universal couplings  
to quark axial currents

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{m_V^2} J_{\text{DM} \mu} \sum_{i=1}^3 \left[ \bar{u}^i \gamma^\mu \gamma^5 u^i + \bar{d}^i \gamma^\mu \gamma^5 d^i \right]$$



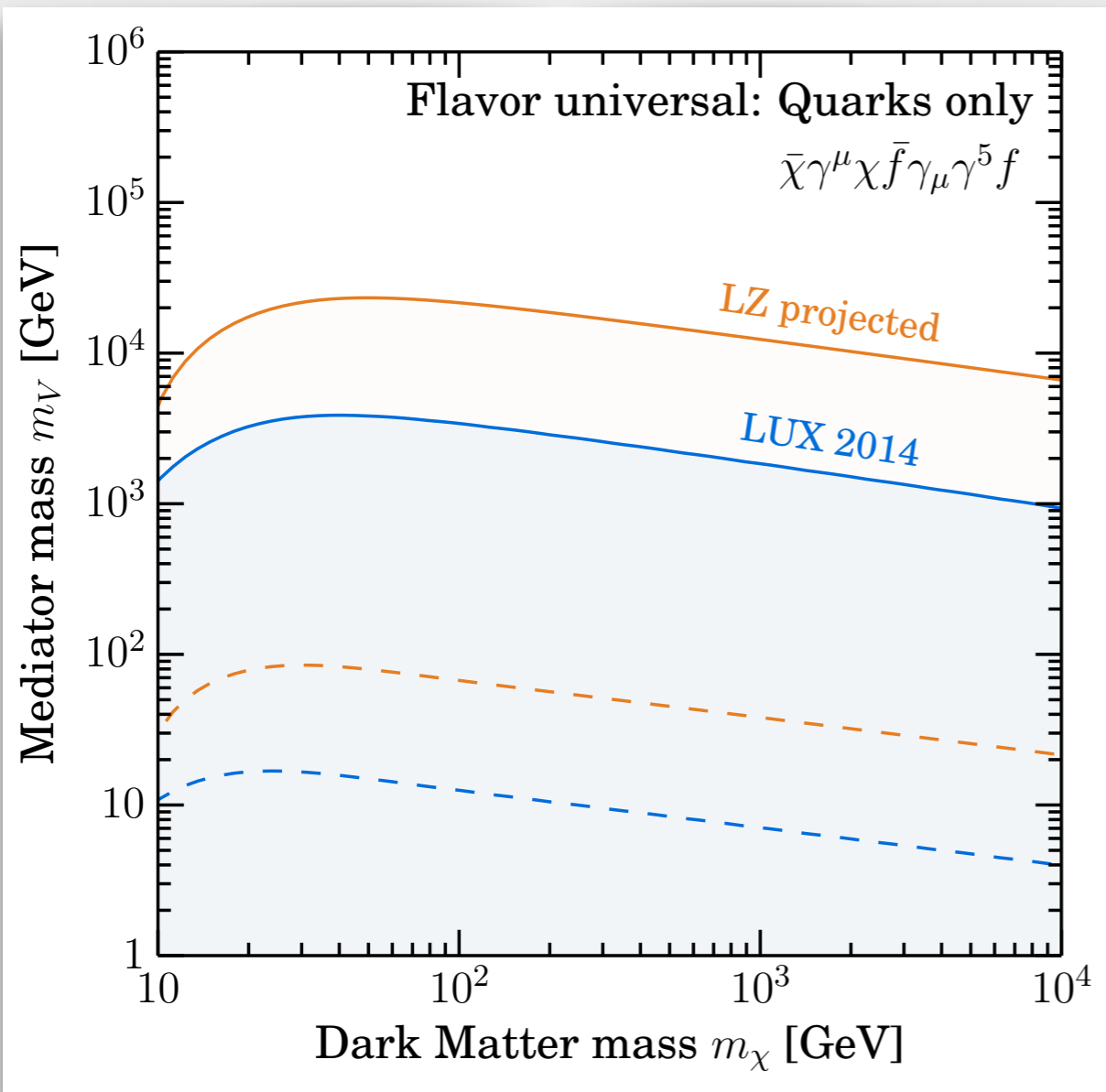
— without RGE

— with RGE

# Results: quarks axial

Flavor universal couplings  
to quark axial currents

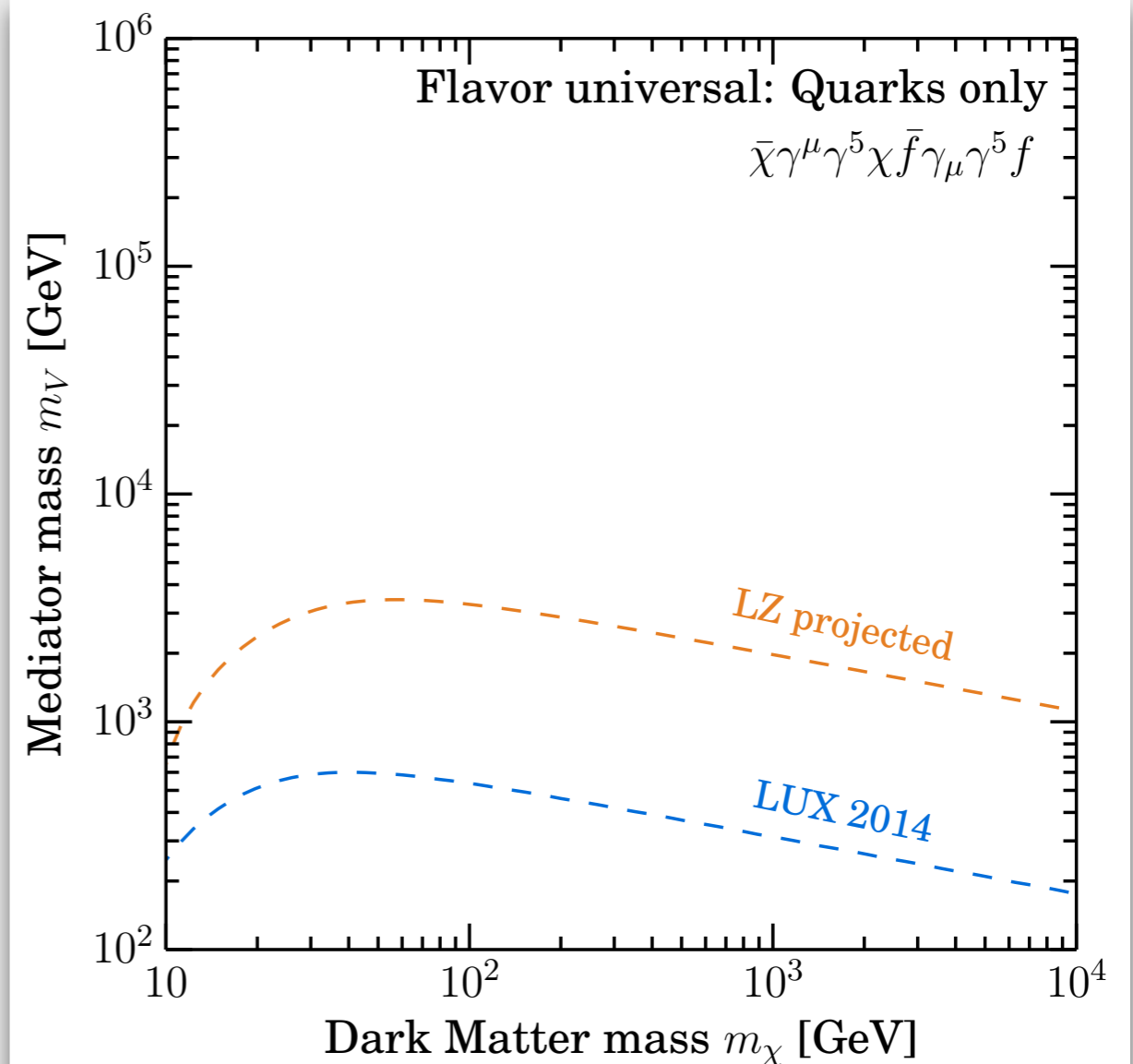
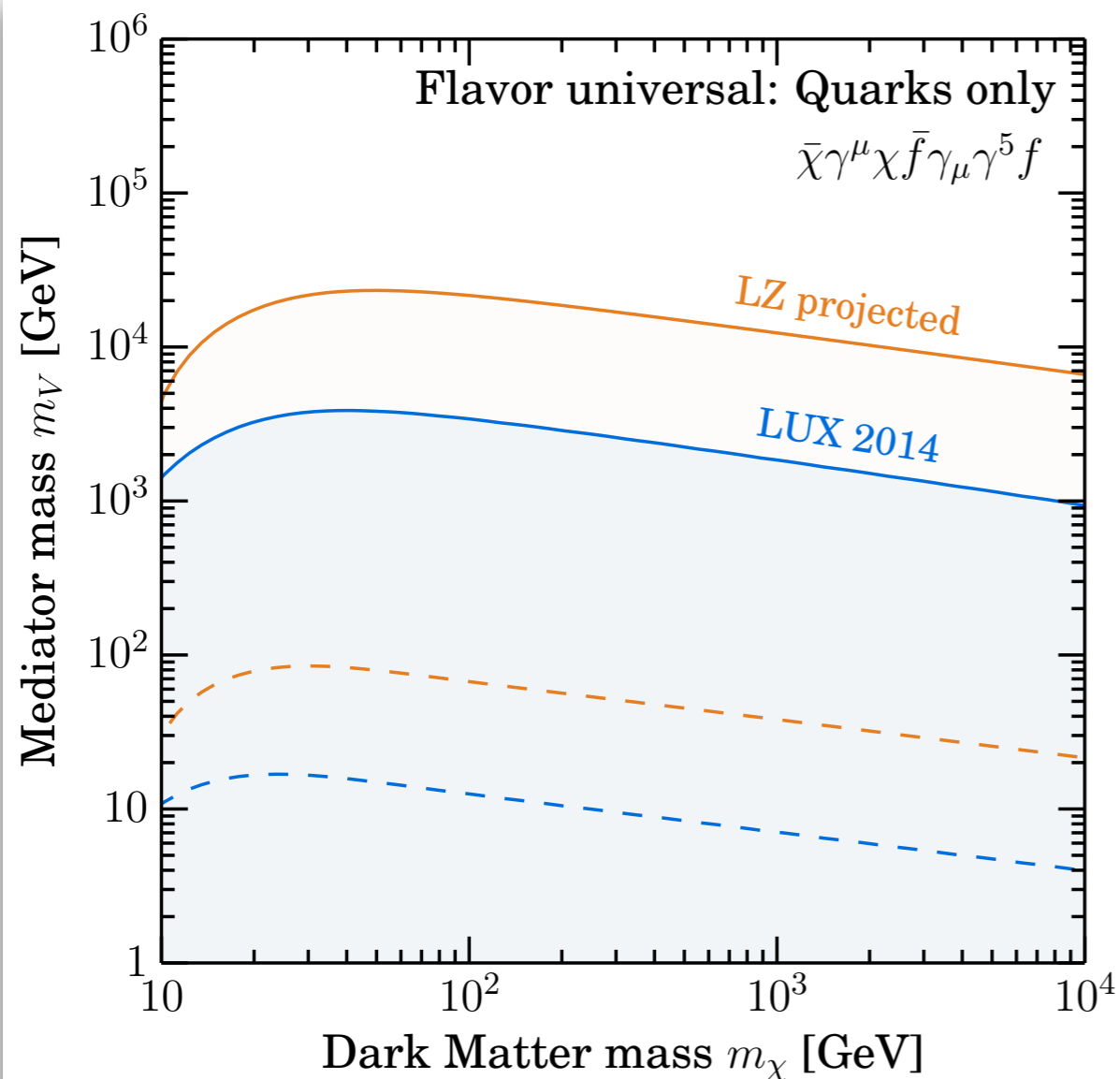
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Flavor universal couplings  
to quark axial currents

$$\mathcal{L}_{\text{EFT}} = -\frac{1}{m_V^2} J_{\text{DM} \mu} \sum_{i=1}^3 \left[ \bar{u}^i \gamma^\mu \gamma^5 u^i + \bar{d}^i \gamma^\mu \gamma^5 d^i \right]$$

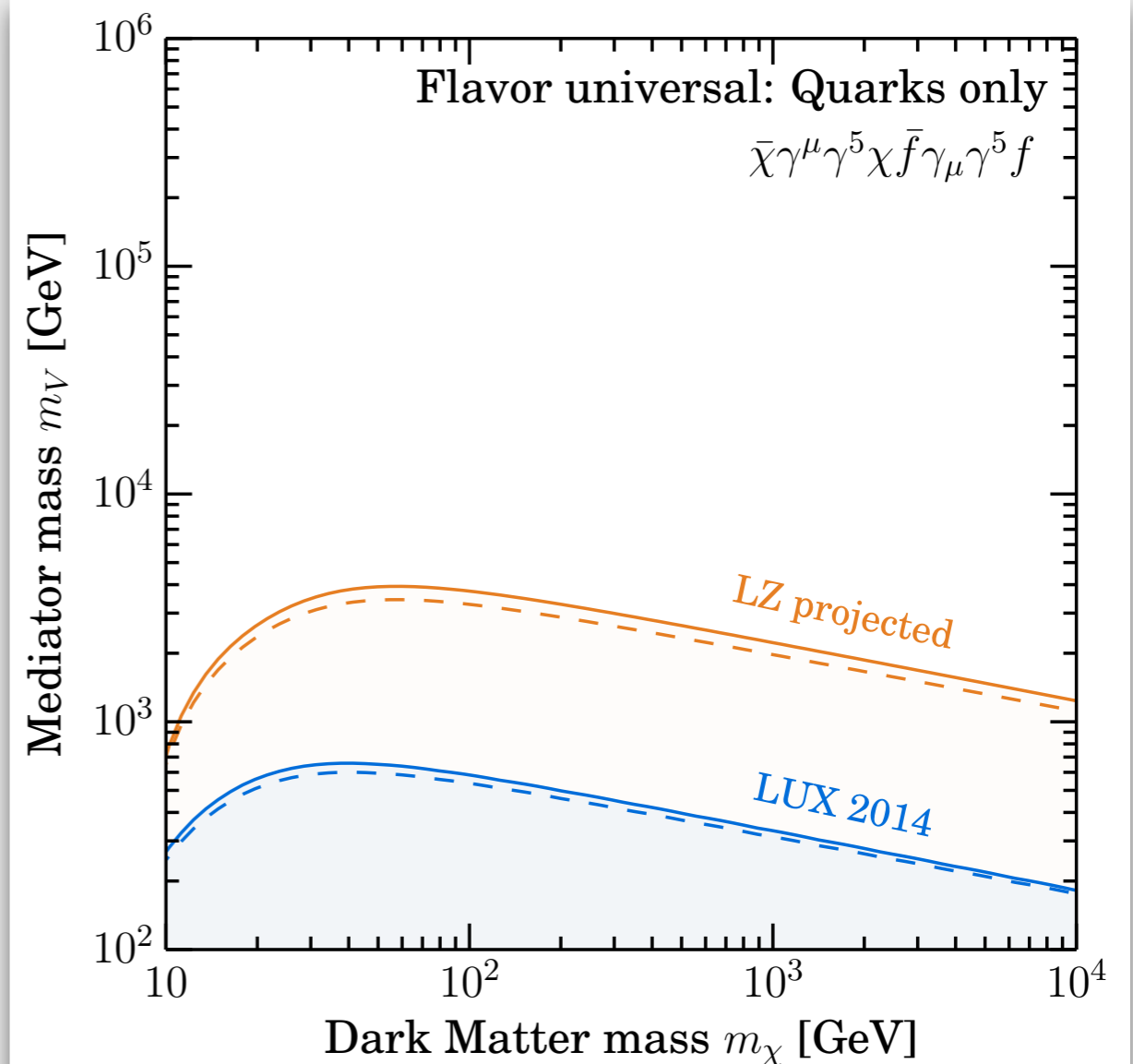
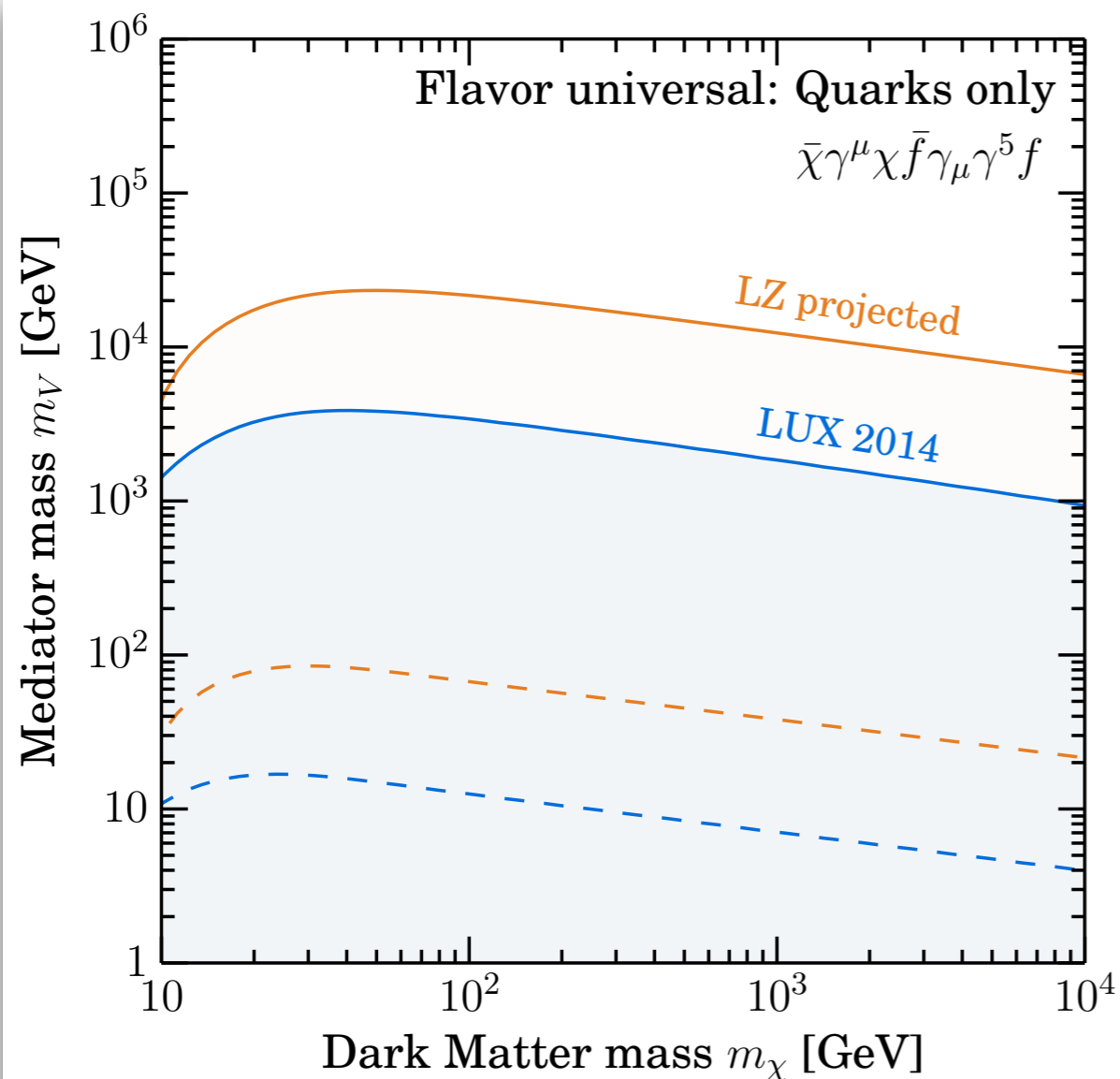




# Results: quarks axial

Flavor universal couplings  
to quark axial currents

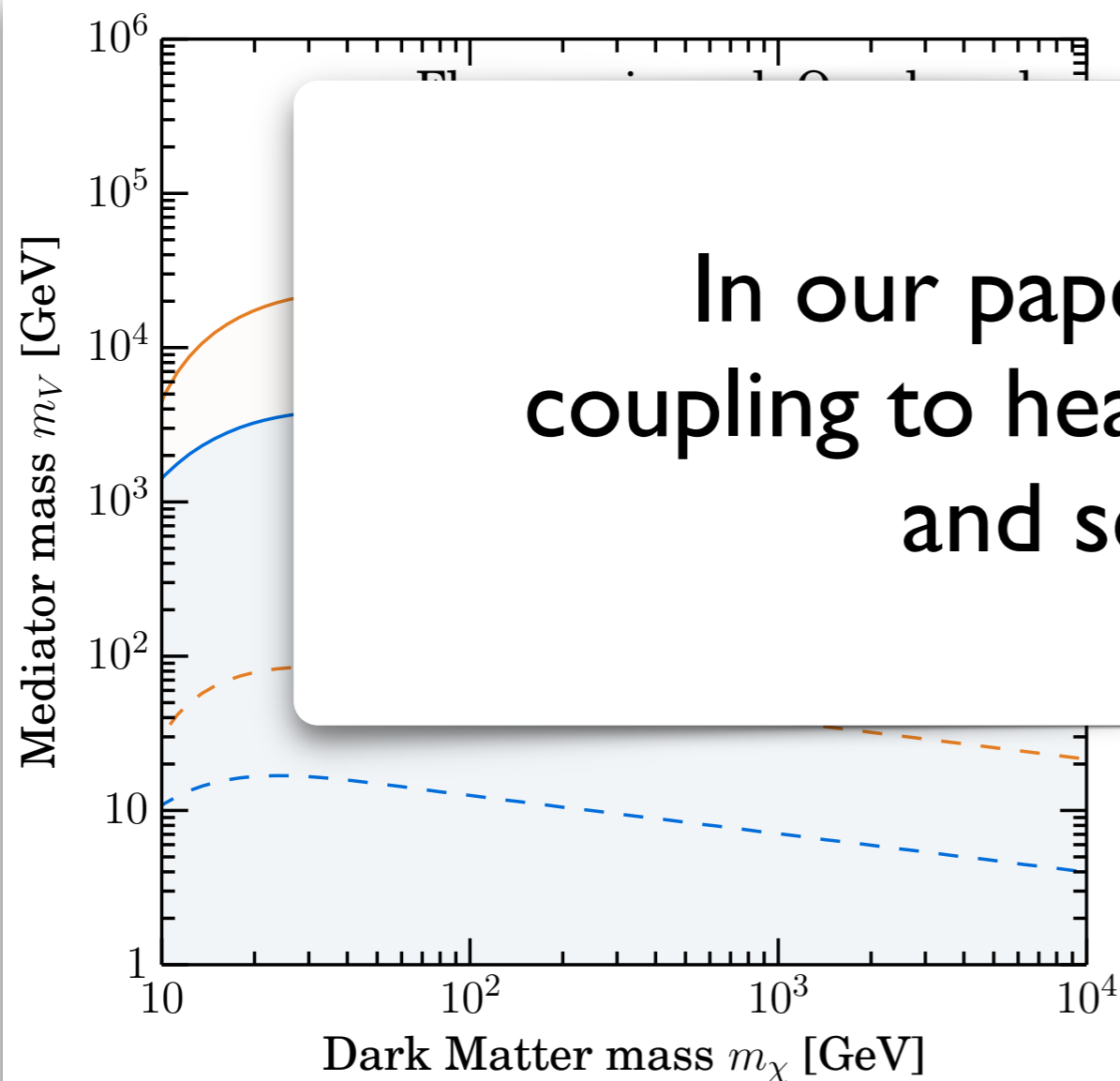
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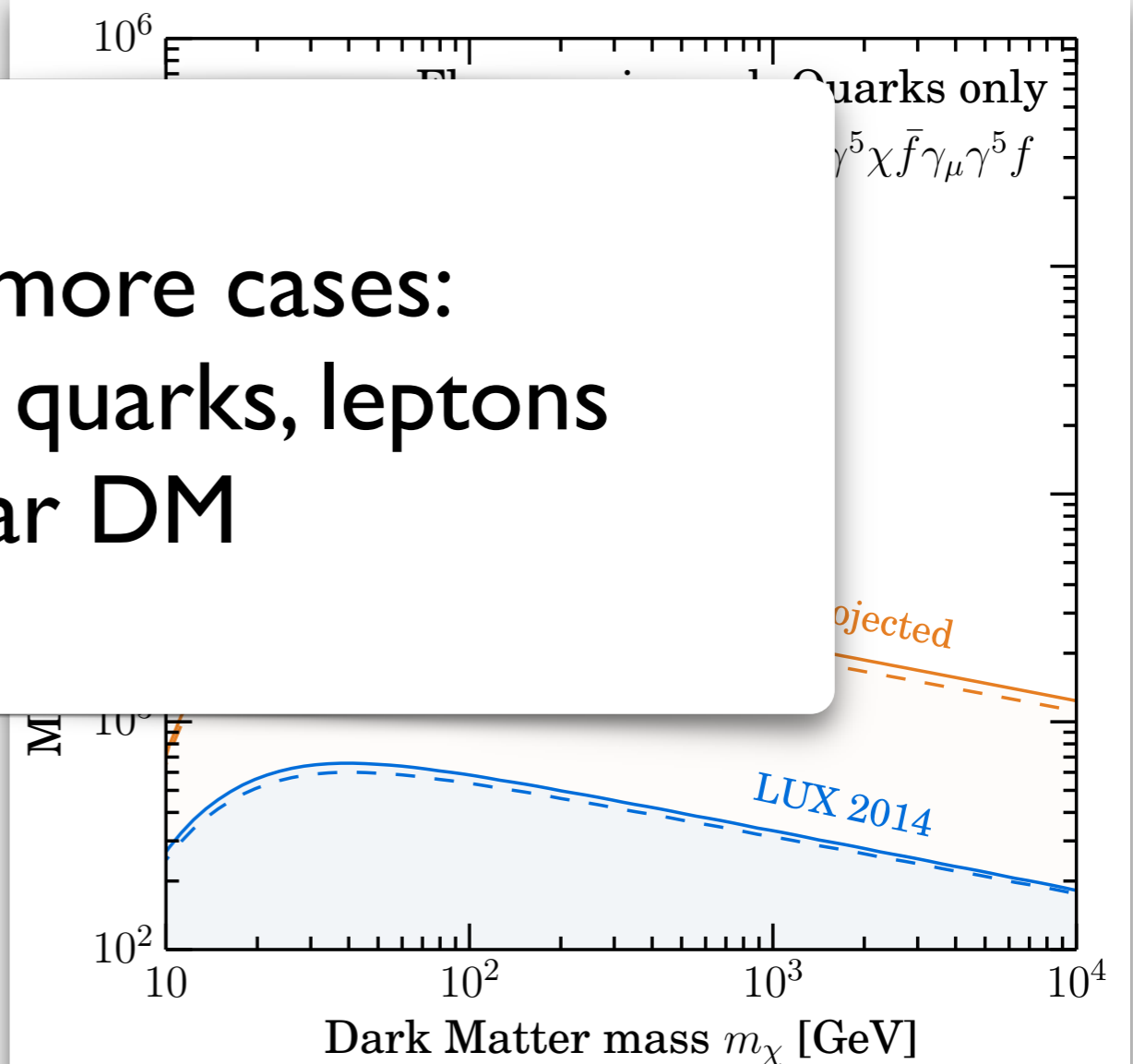
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Flavor universal couplings  
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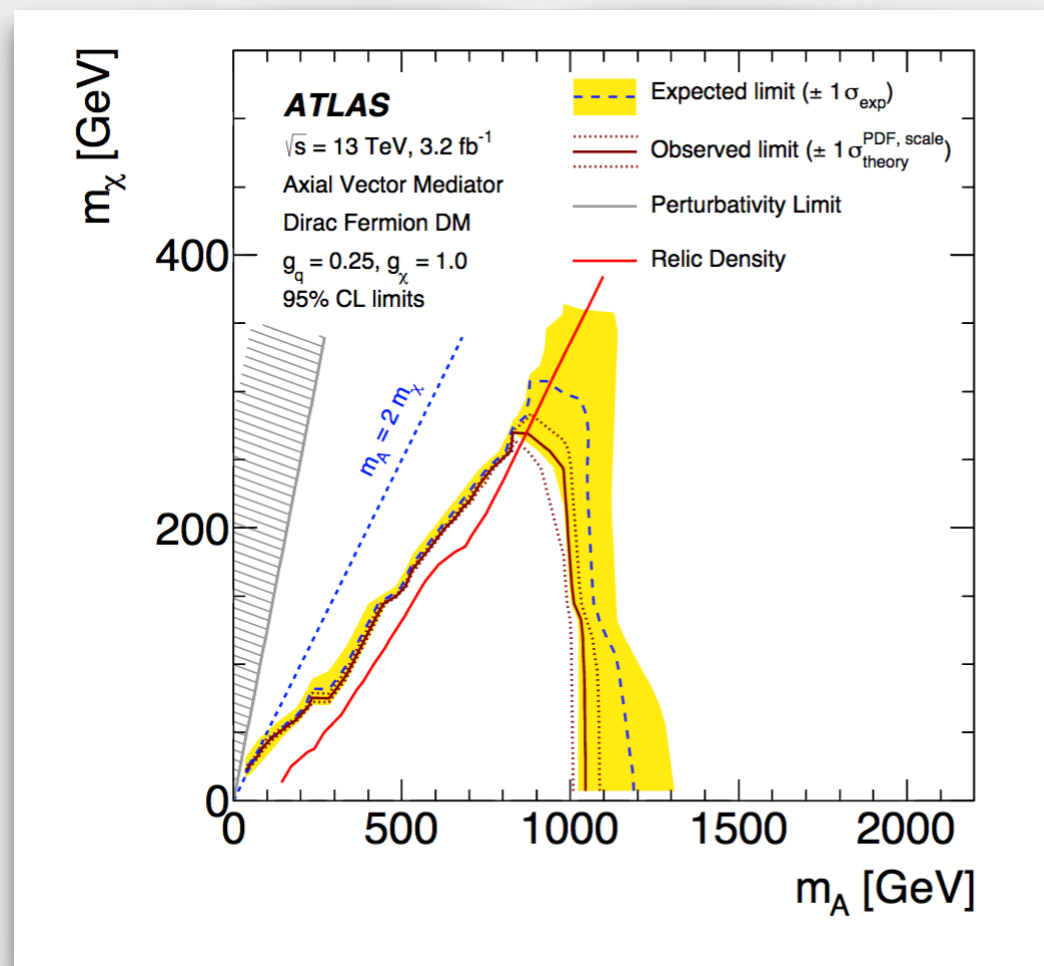


In our paper more cases:  
coupling to heavy quarks, leptons  
and scalar DM

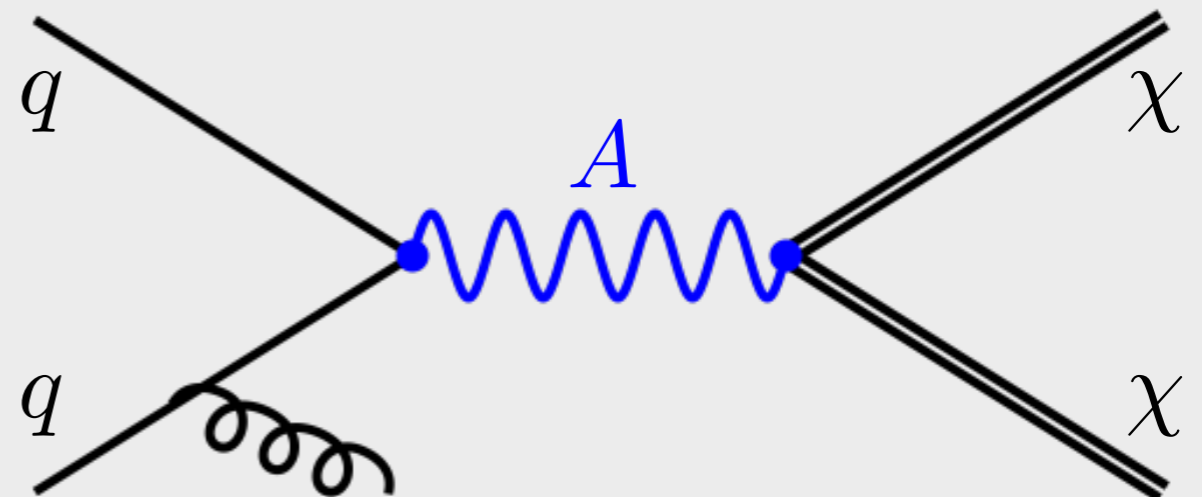


# Comparison with LHC

LHC limits on mediator mass translated into limits on direct detection cross section



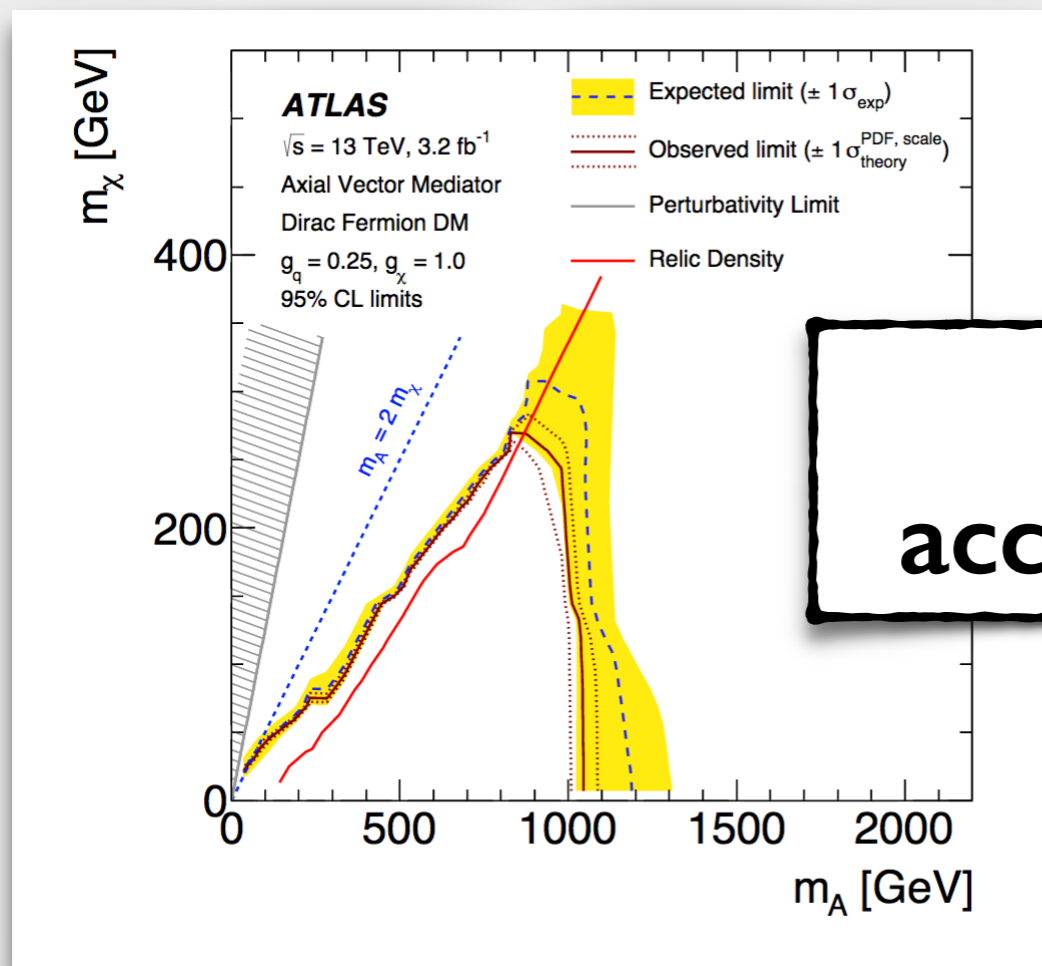
$$\mathcal{L}_{\text{Axial-Vector}} = g_\chi A_\mu \bar{\chi} \gamma^\mu \gamma^5 \chi + g_q A_\mu \bar{q} \gamma^\mu \gamma^5 q$$



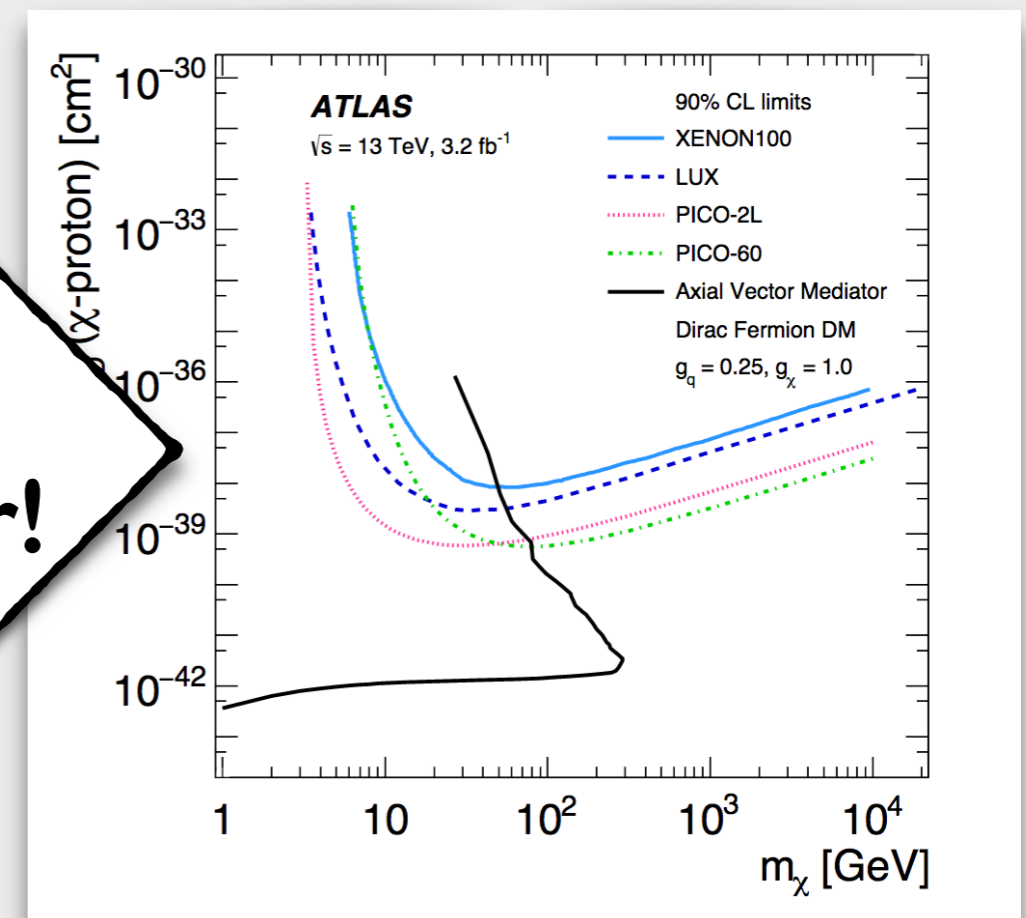
LUX, arXiv:1602.03489 and PICO-2L, arXiv:1601.03729  
ATLAS, arXiv:1604.07773 (mono-jet) and arXiv:1604.01306 (mono-photon)

# Comparison with LHC

LHC limits on mediator mass translated into limits on direct detection cross section

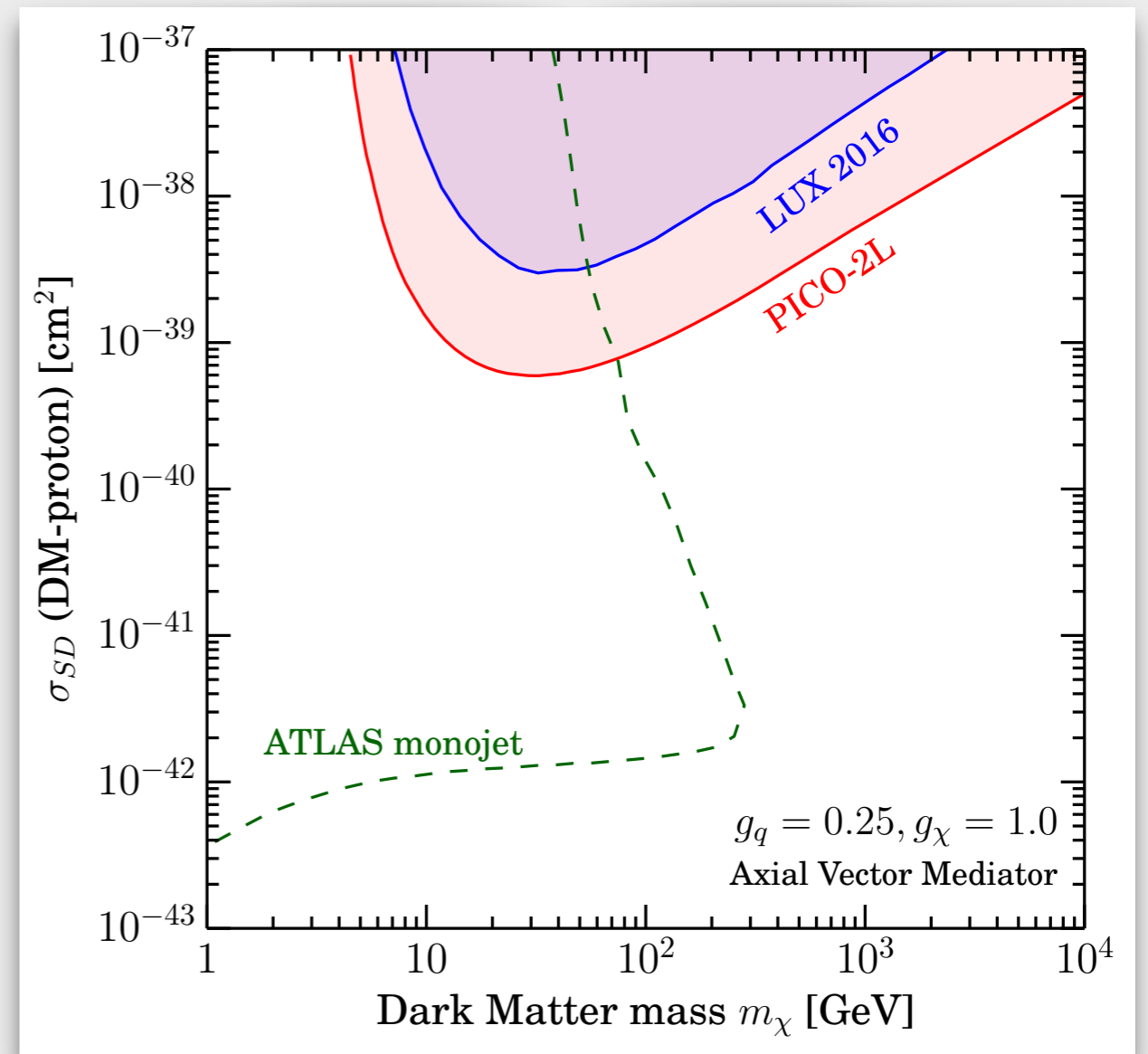
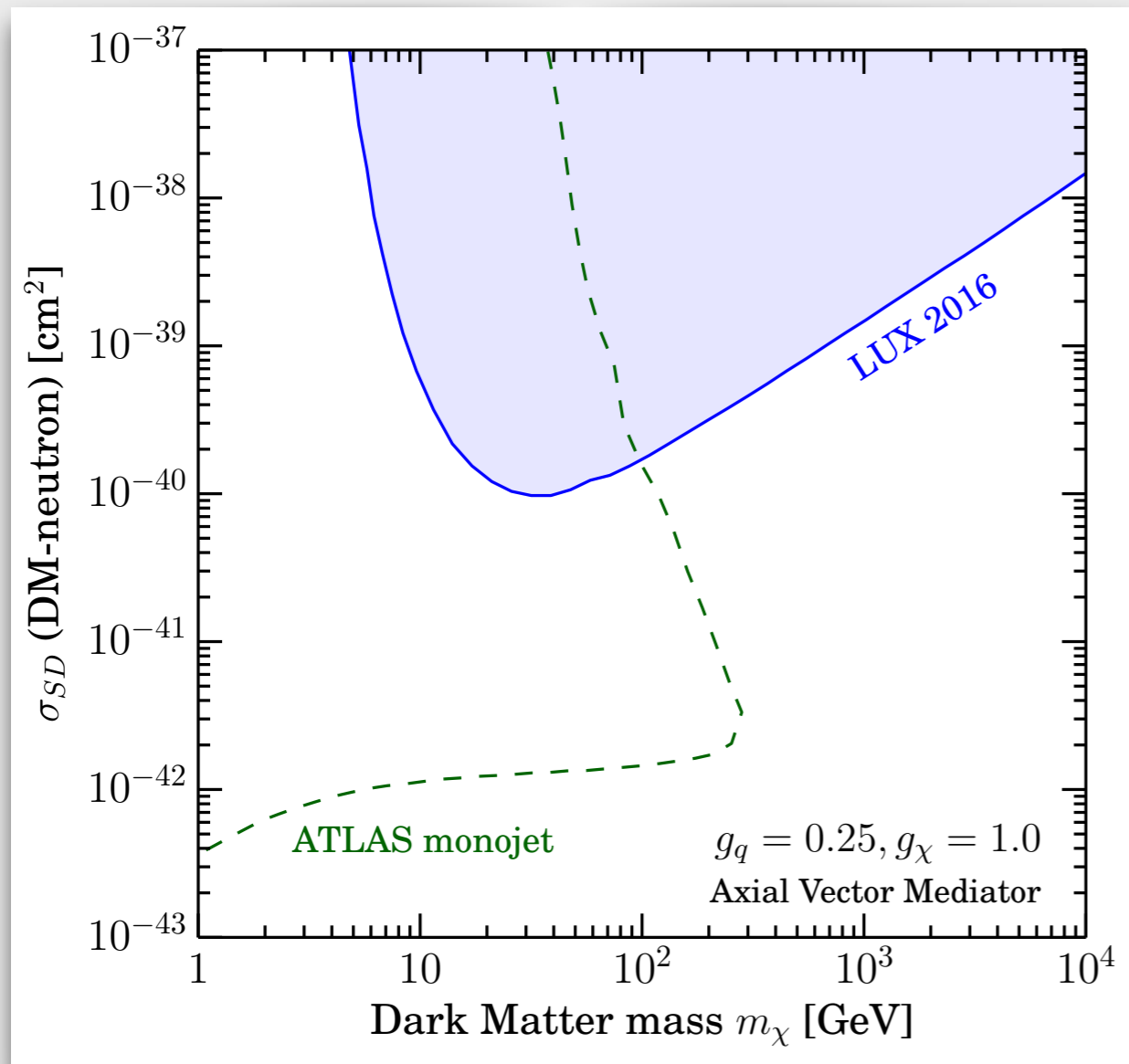


RGE not accounted for!



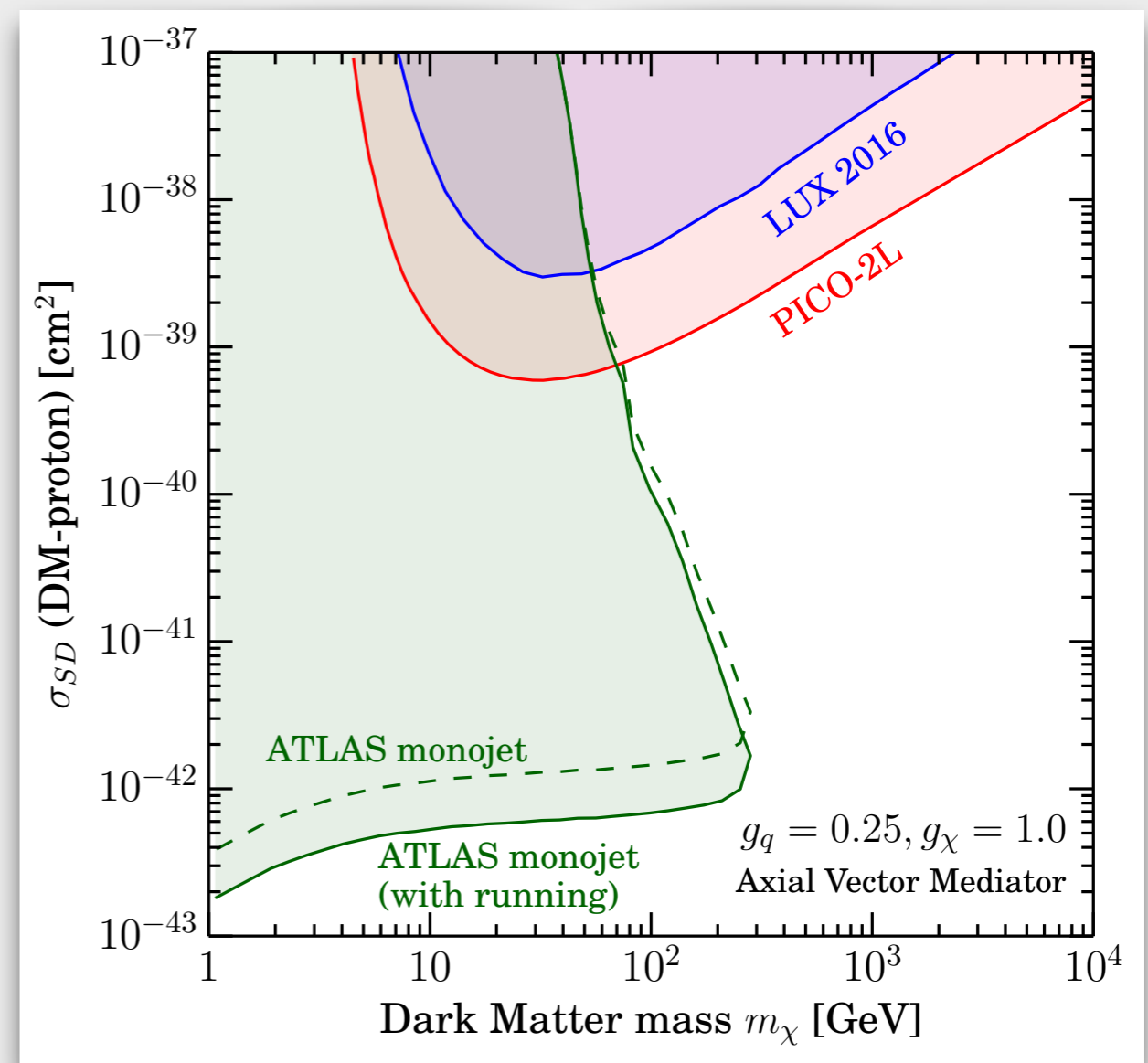
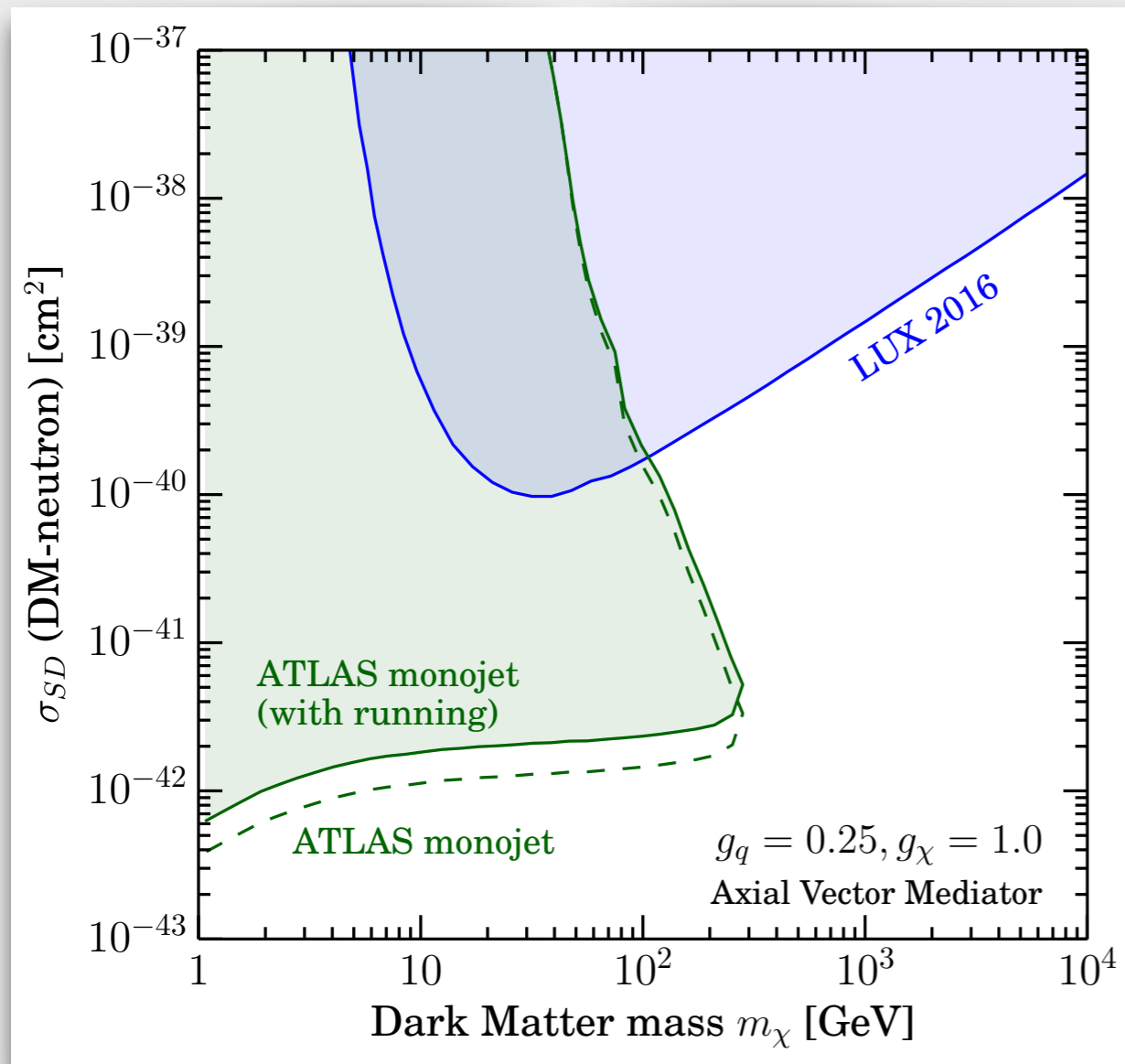
LUX, arXiv:1602.03489 and PICO-2L, arXiv:1601.03729  
ATLAS, arXiv:1604.07773 (mono-jet) and arXiv:1604.01306 (mono-photon)

# Mono-jet Searches



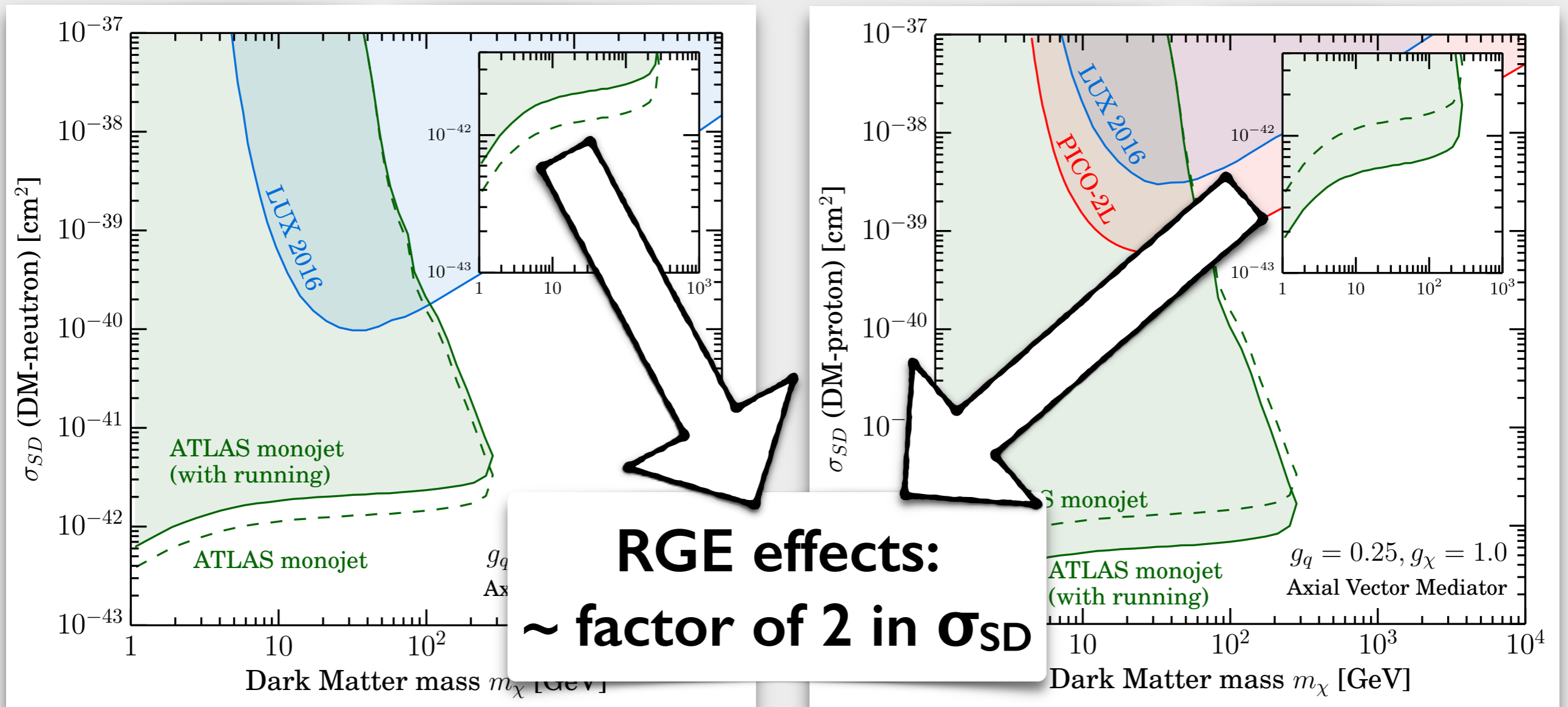
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# Mono-jet Searches



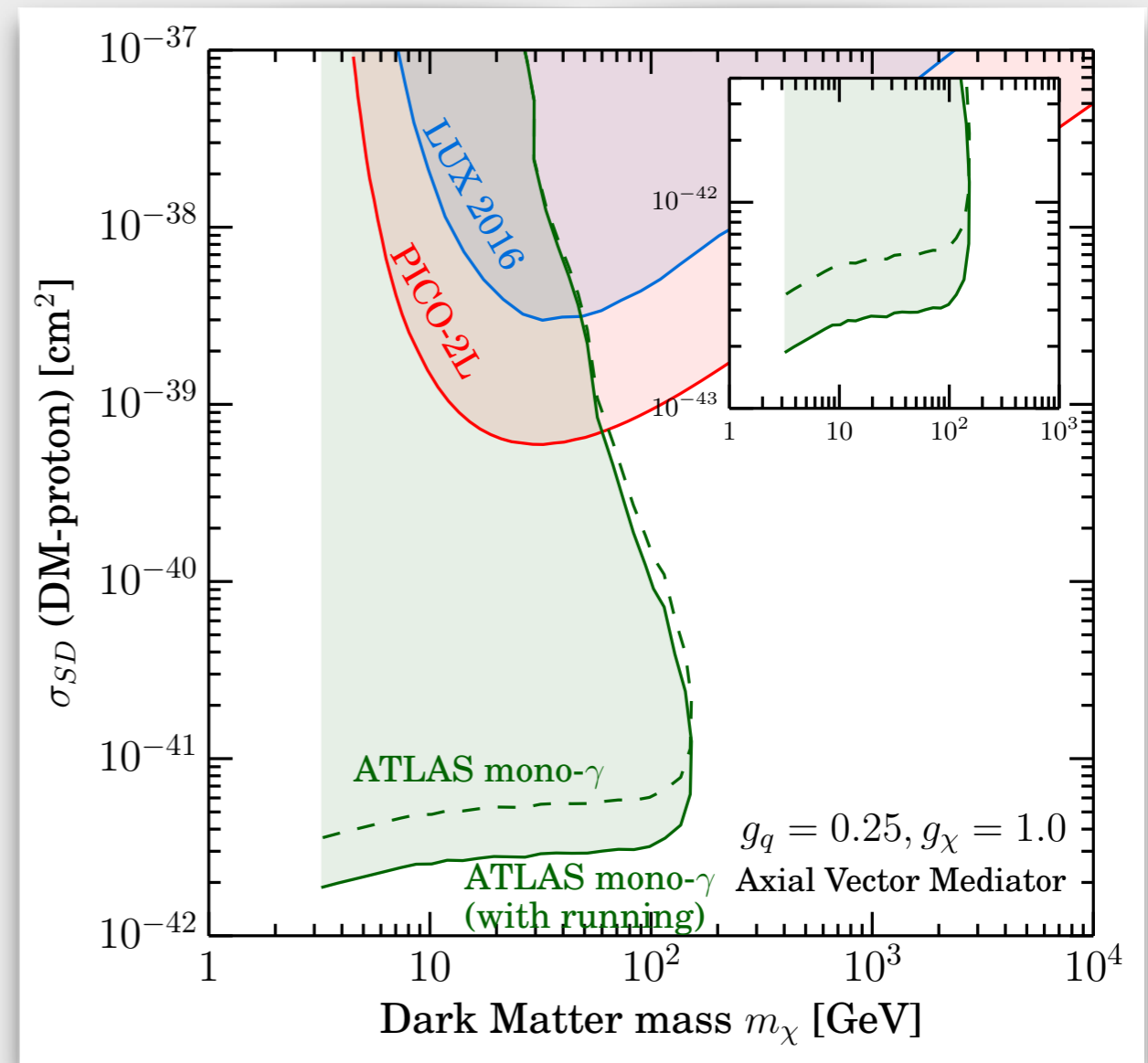
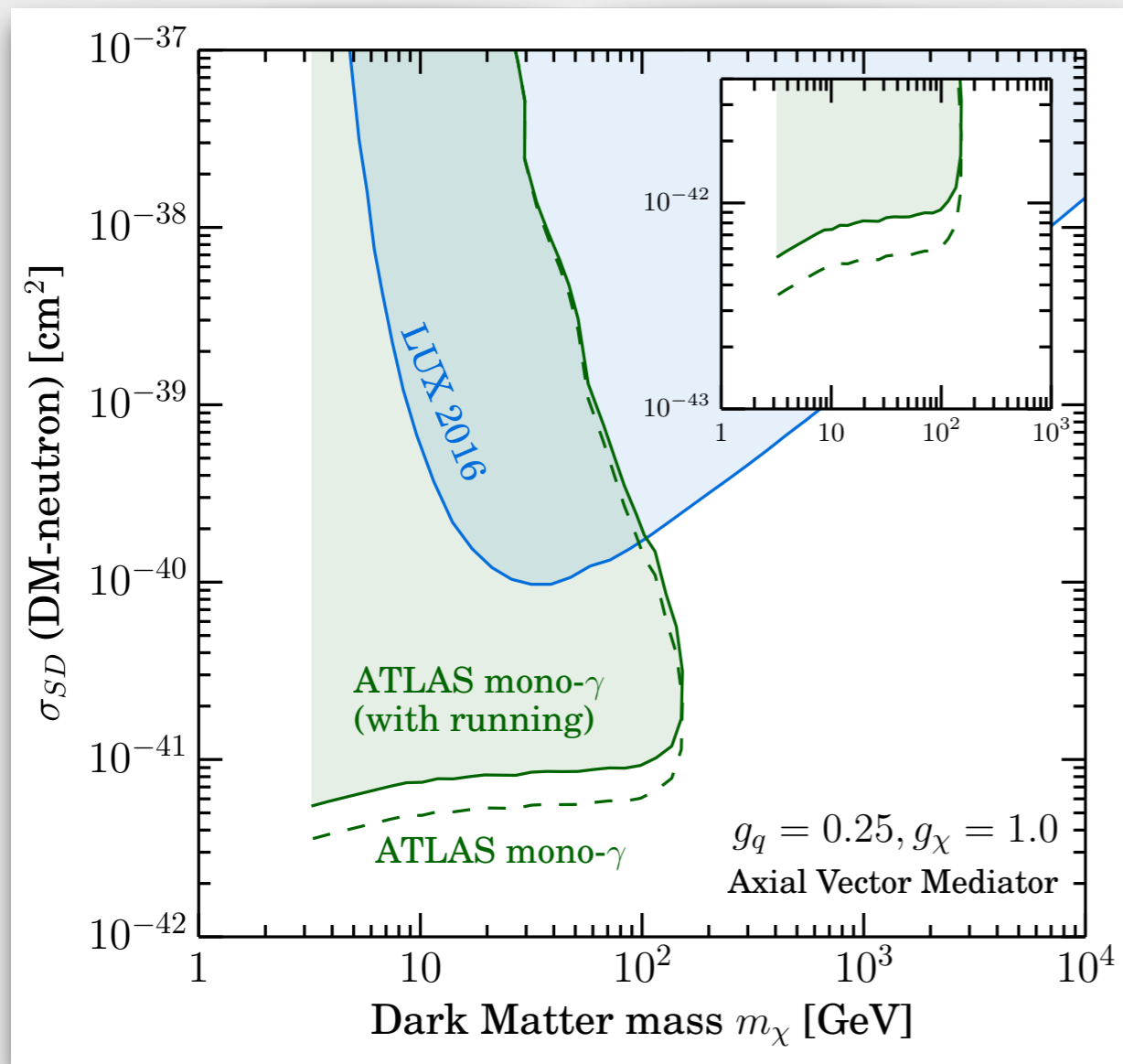
LUX, arXiv:1602.03489 and PICO-2L, arXiv:1601.03729  
ATLAS, arXiv:1604.07773 (mono-jet) and arXiv:1604.01306 (mono-photon)

# Mono-jet Searches



LUX, arXiv:1602.03489 and PICO-2L, arXiv:1601.03729  
 ATLAS, arXiv:1604.07773 (mono-jet) and arXiv:1604.01306 (mono-photon)

# Mono-photon Searches



LUX, arXiv:1602.03489 and PICO-2L, arXiv:1601.03729  
ATLAS, arXiv:1604.07773 (mono-jet) and arXiv:1604.01306 (mono-photon)



# More Complementarity

How well can be constrained by  
**collider** and **direct detection**?

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$$

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$$

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$$

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$$

# More Complementarity

How well can be constrained by  
**collider** and **direct detection**?

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$$

~ same as the others

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$$

~ same as the others

# More Complementarity

How well can be constrained by  
**collider** and **direct detection**?

**Standard  
Lore**

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$$

~ same as the others

Spin-Independent (SI), no suppression

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

Spin-Dependent (SD), no suppression

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

SD with  $v$

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu q$$

~ same as the others

SI with  $v$

# More Complementarity

How well can be constrained by  
**collider** and **direct detection**?

**Standard  
Lore**

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu q$$

~ same as the others

Spin-Independent (SI), no suppression

$$\bar{\chi}\gamma^\mu\gamma^5\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

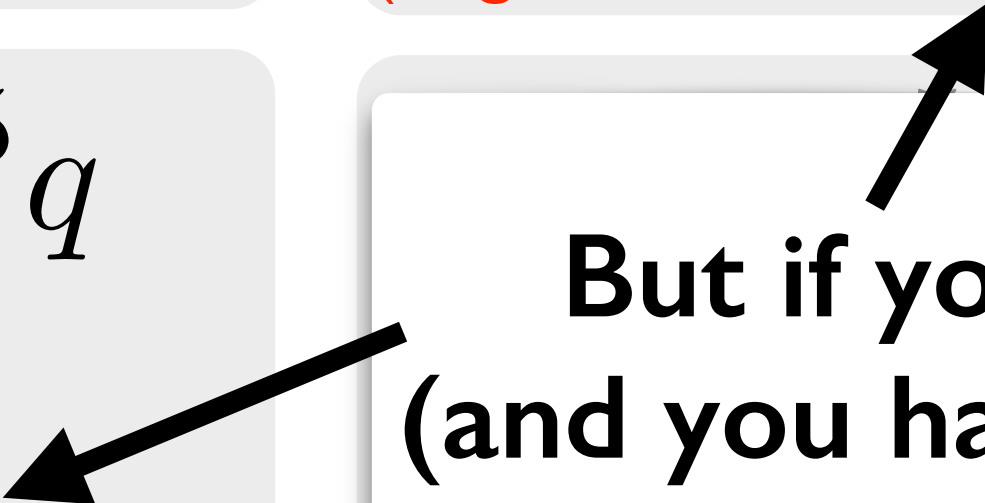
Spin-Dependent (SD), no suppression  
(large RGE corrections)

$$\bar{\chi}\gamma^\mu\chi\bar{q}\gamma_\mu\gamma^5q$$

~ same as the others

Spin-Independent (SI) via RGE

**But if you run...  
(and you have to run!)**

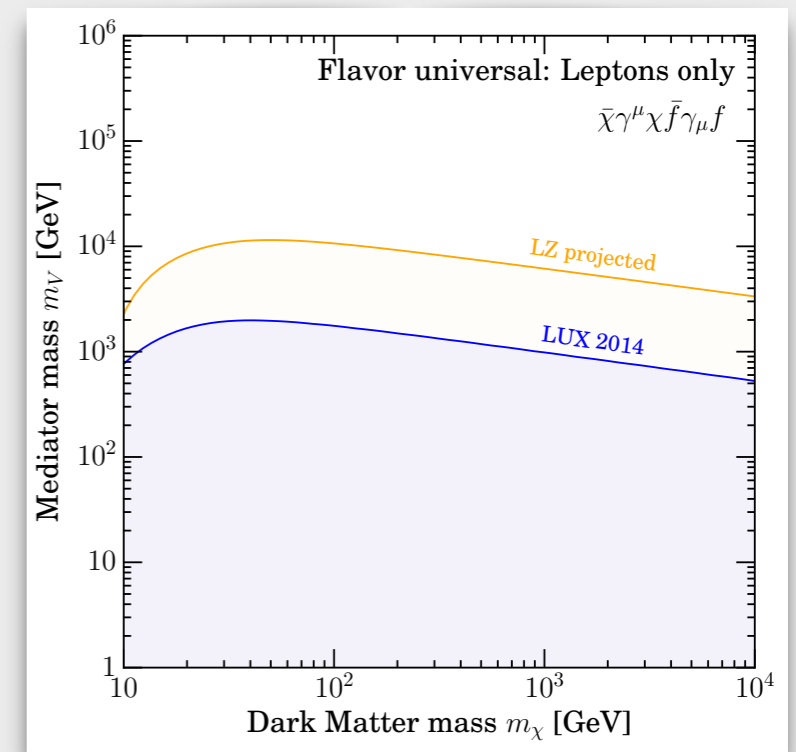
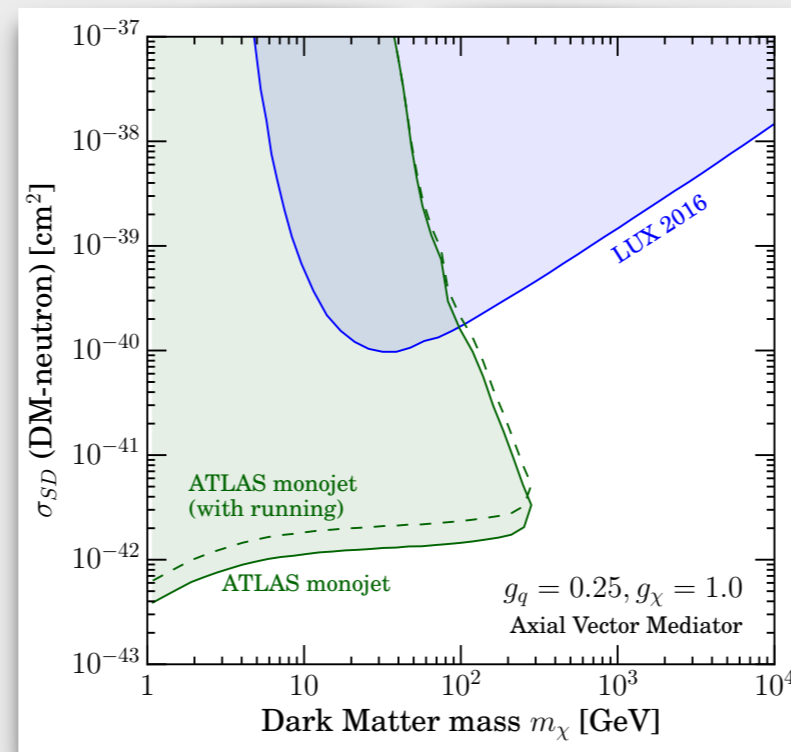
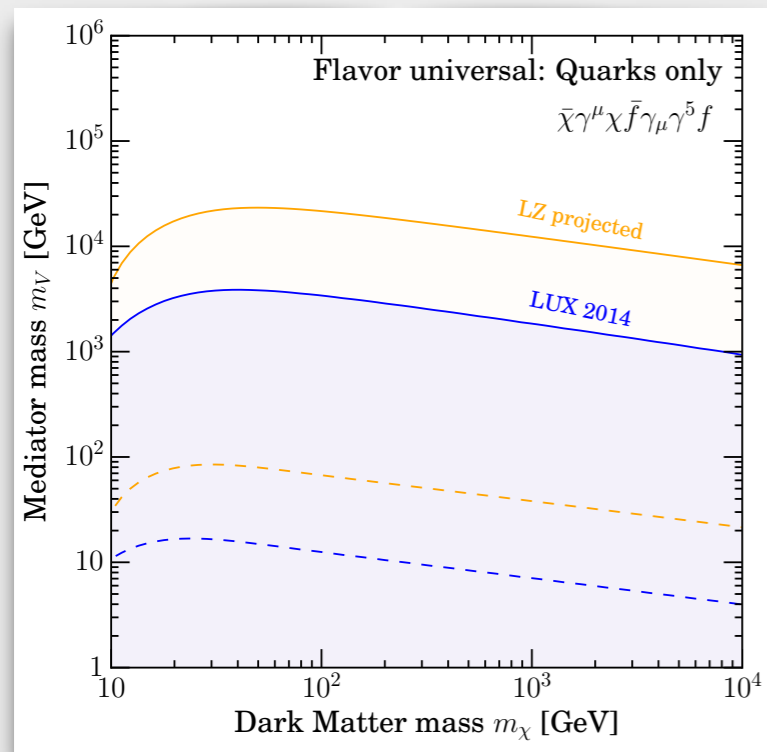


# Outlook

**Direct detection rates and comparison with LHC:  
not always straightforward**

# Outlook

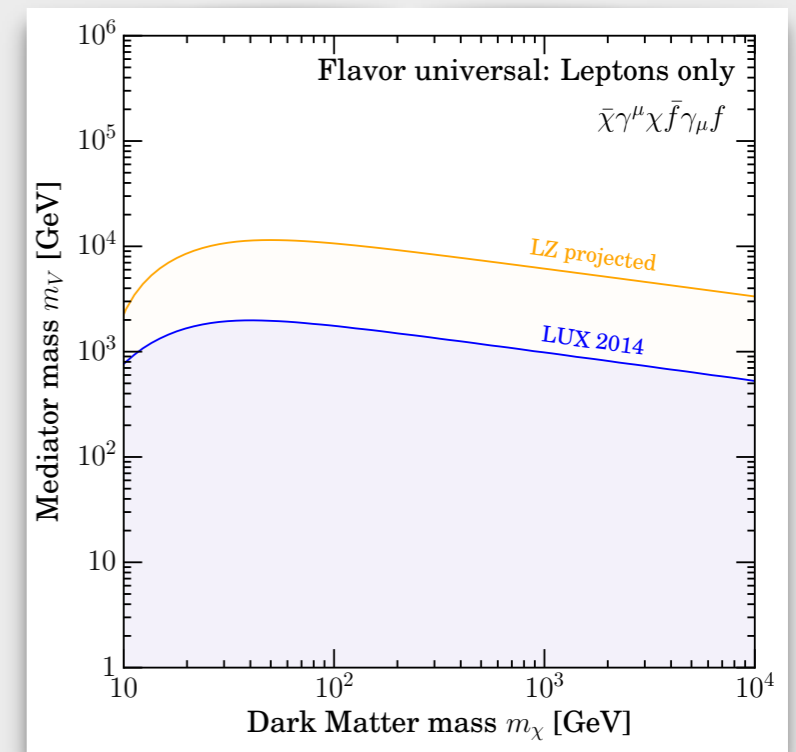
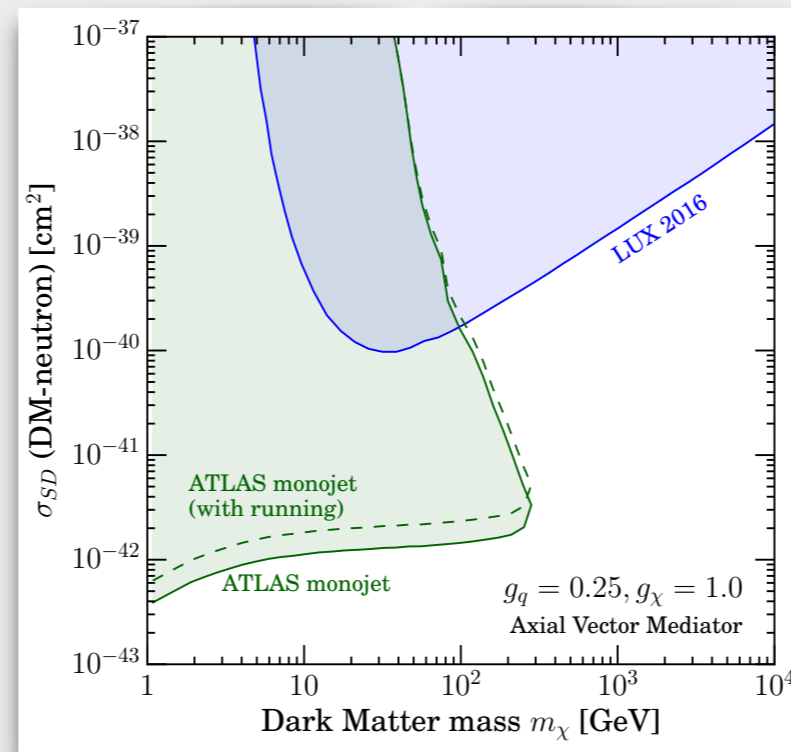
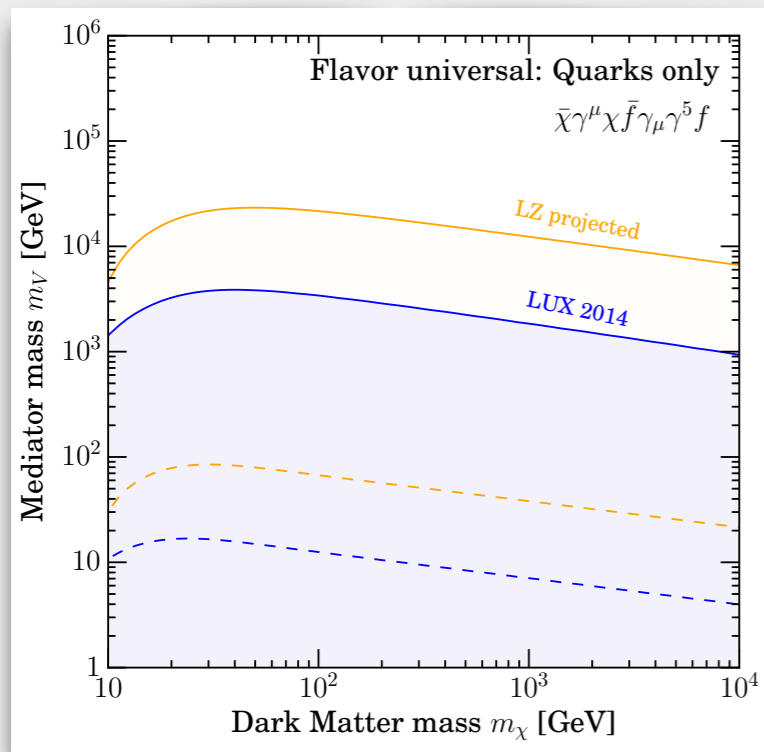
Direct detection rates and comparison with LHC:  
not always straightforward



FD, Kavanagh, Panci, arXiv:1605.04917

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Thank You!