

# S-MATRIX APPROACH TO HADRON GAS

POK MAN LO

**University of Wrocław**

XII POLISH WORKSHOP ON RHIC  
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# IN COLLABORATION WITH

Bengt Friman (GSI)

Chihiro Sasaki (U. of Wroclaw)

Pasi Huovinen (U. of Wroclaw)

Krzysztof Redlich (U. of Wroclaw)

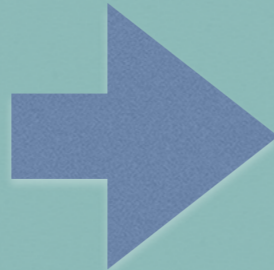
Michal Marczenko (U. of Wroclaw)

# S-MATRIX APPROACH

# HADRON RESONANCE GAS MODEL

- Confinement

physical  
quantities



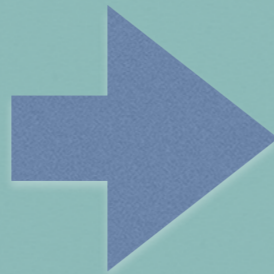
hadronic states  
representation

$$Z = \sum_{\alpha=B,M} \langle \alpha | e^{-\beta H} | \alpha \rangle$$

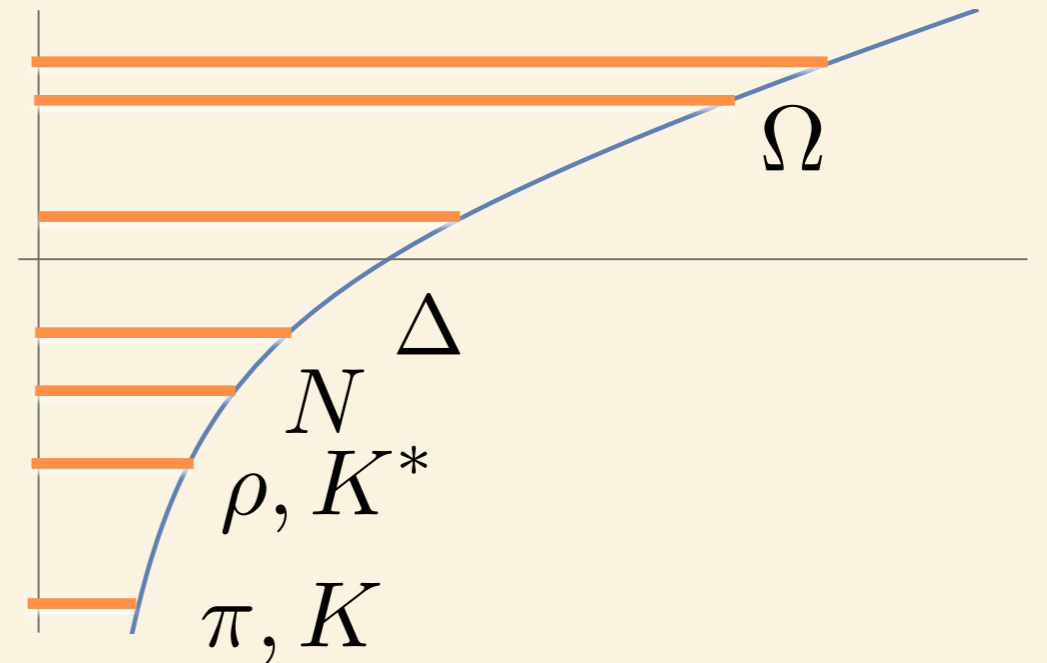
# HADRON RESONANCE MODEL

- Confinement

physical  
quantities

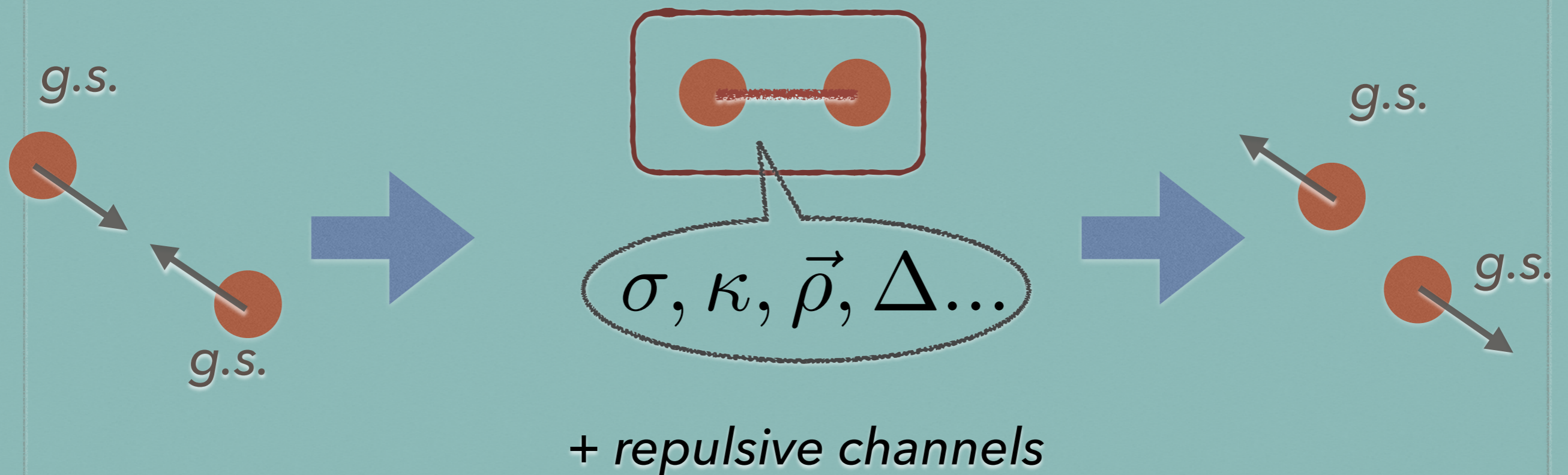


QCD spectrum



$$Z = \sum_{\alpha=B,M} \langle \alpha | e^{-\beta H} | \alpha \rangle$$

# S-MATRIX APPROACH



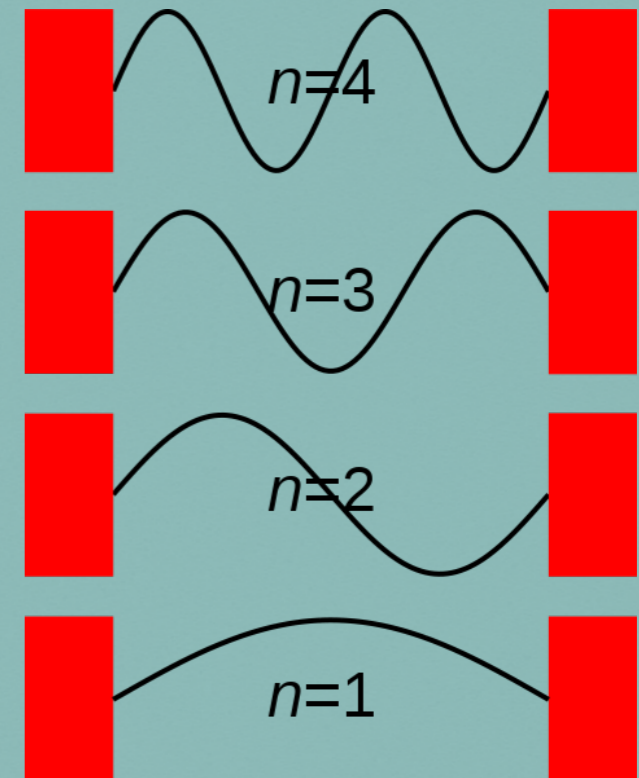
interaction:  
attractive + repulsive

# PHASE SHIFT AND DENSITY OF STATES

*particle in a box*

$$\psi \sim \sin(k^{(0)} x)$$

$$k^{(0)} = \frac{n\pi}{L}$$



# PHASE SHIFT AND DENSITY OF STATES

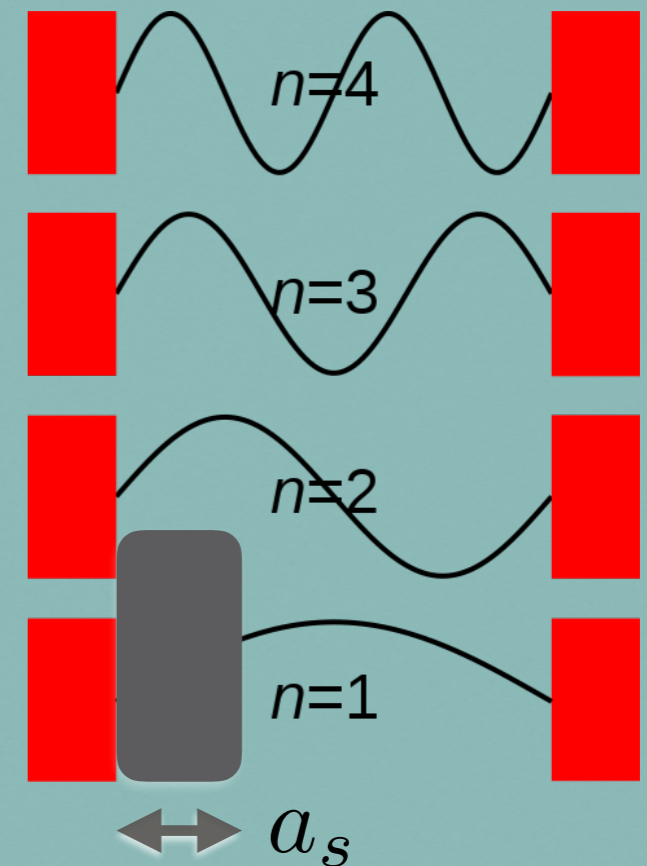
*particle in a box*

$$\psi \sim \sin(k^{(0)}x) \quad k^{(0)} = \frac{n\pi}{L}$$

*in the presence of a scattering potential*

$$\psi \sim \sin(kx + \delta(k))$$

density of states



$$kL + \delta(k) = n\pi \quad \longrightarrow \quad \frac{dn(k)}{dk} = \frac{L}{\pi} + \frac{1}{\pi} \frac{d\delta}{dk}$$

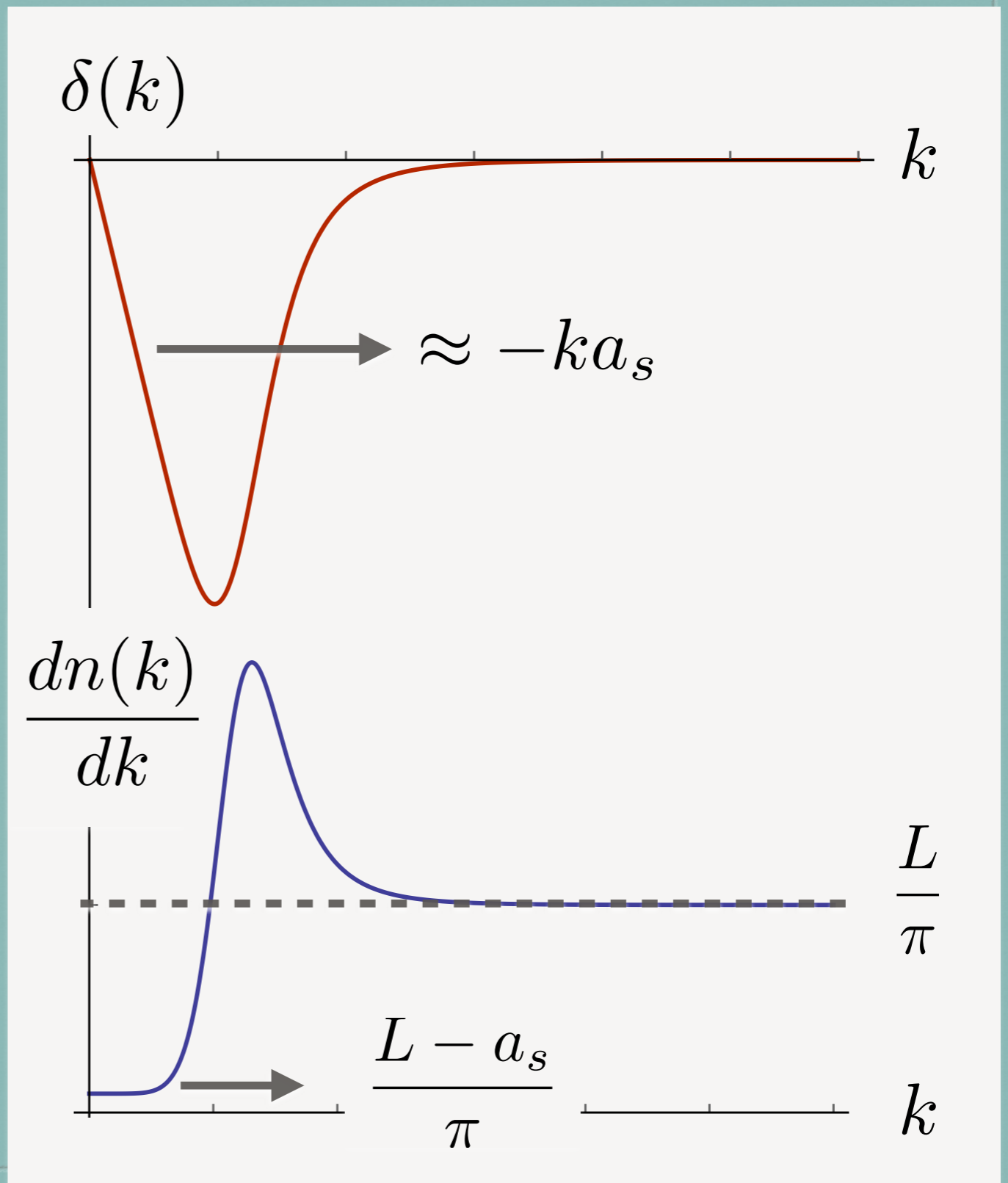


# PHASE SHIFT AND DENSITY OF STATES

$$\frac{dn(k)}{dk} = \frac{L}{\pi} + \frac{1}{\pi} \frac{d\delta}{dk}$$

change in d.o.s.  
due to int.

Effect of repulsive  
interaction:  
pushing states from low  $k$   
to high  $k$



phase shift and d.o.s. (schematics)

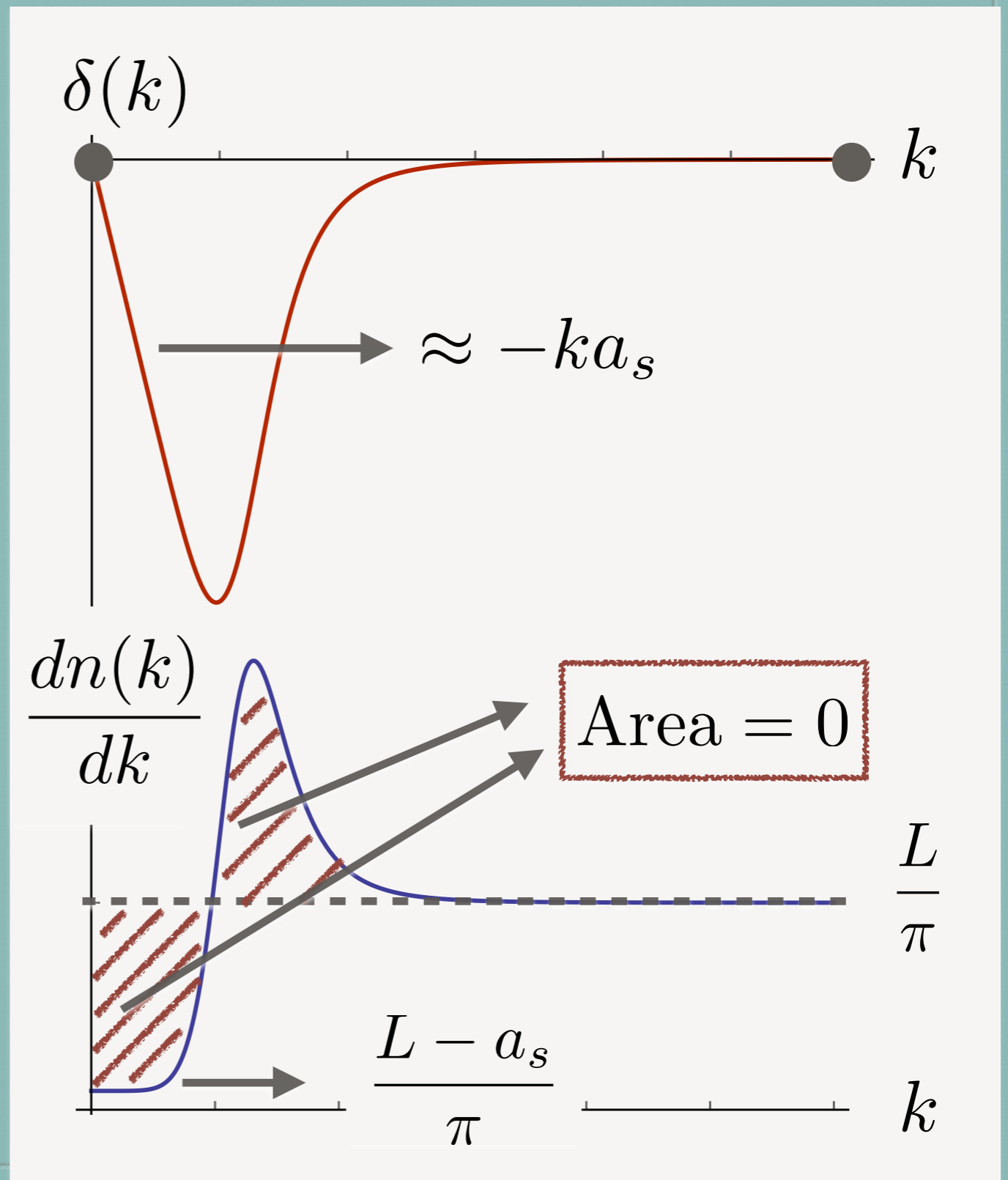
# PHASE SHIFT AND DENSITY OF STATES

$$\frac{dn(k)}{dk} = \frac{L}{\pi} + \frac{1}{\pi} \frac{d\delta}{dk}$$

sum rule  
(Levinson's theorem)

$$\int_0^\infty dk \frac{1}{\pi} \delta' = \frac{\delta(\infty) - \delta(0)}{\pi}$$

$n_{\text{int}}$



phase shift and d.o.s. (schematics)

# FORMULATION

given the exact phase shift  $\delta_l$

from theory  
or  
from experiment



thermodynamics

$$B_l = 2 \frac{d}{dq} \delta_l$$

eff. spectral function

$$P = P^{(0)} + \Delta P^{B.U.}$$

free gas + interaction

# FORMULATION

**dynamical**

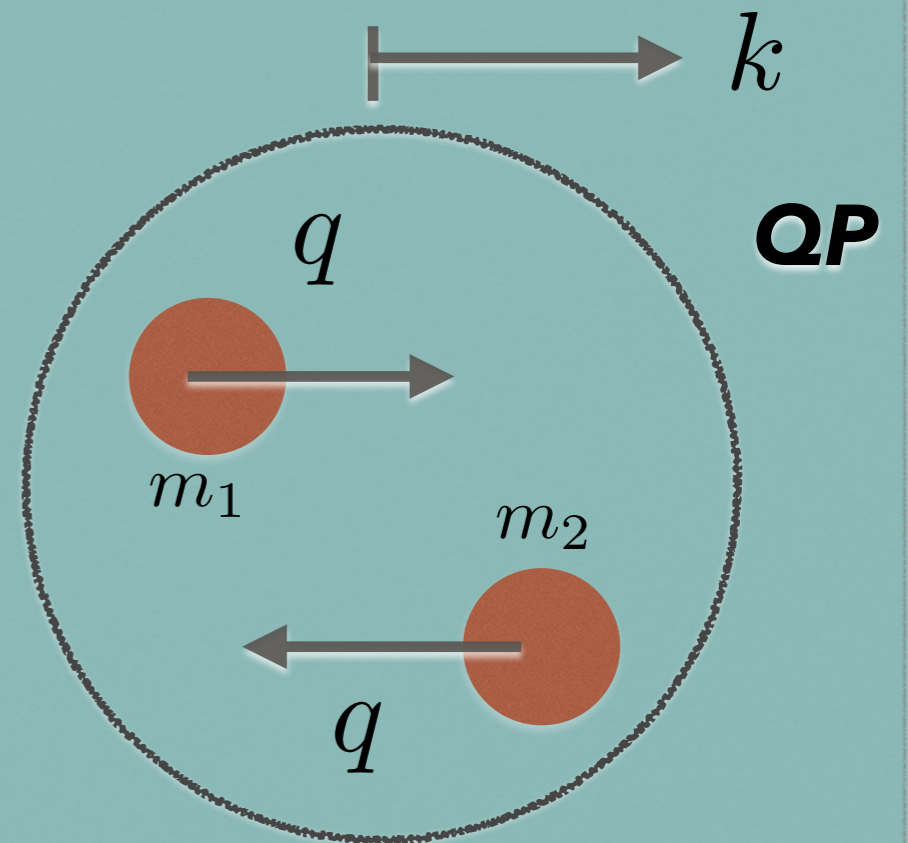
**statistical (thermal weight)**

$$\Delta P^{\text{B.U.}} = (2l + 1) \int \frac{dq}{2\pi} B_l(q) \int \frac{d^3 k}{(2\pi)^3} T \ln(1 + e^{-\beta E(k, q, m_i)})$$

$$E = \sqrt{k^2 + M(q)^2}$$

$$M(q) = \sqrt{q^2 + m_1^2} + \sqrt{q^2 + m_2^2}$$

$$B_l = 2 \frac{d}{dq} \delta_l$$



$$M(q) = \sqrt{q^2 + m_1^2} + \sqrt{q^2 + m_2^2}$$

**APPLICATION**  
**PION + NUCLEON + DELTA SYSTEM**

# WHAT'S IN A NAME? THAT WHICH WE CALL A RESONANCES?

- A resonance is MORE than a **MASS** and a **WIDTH**

$\Delta(1232) 3/2^+$

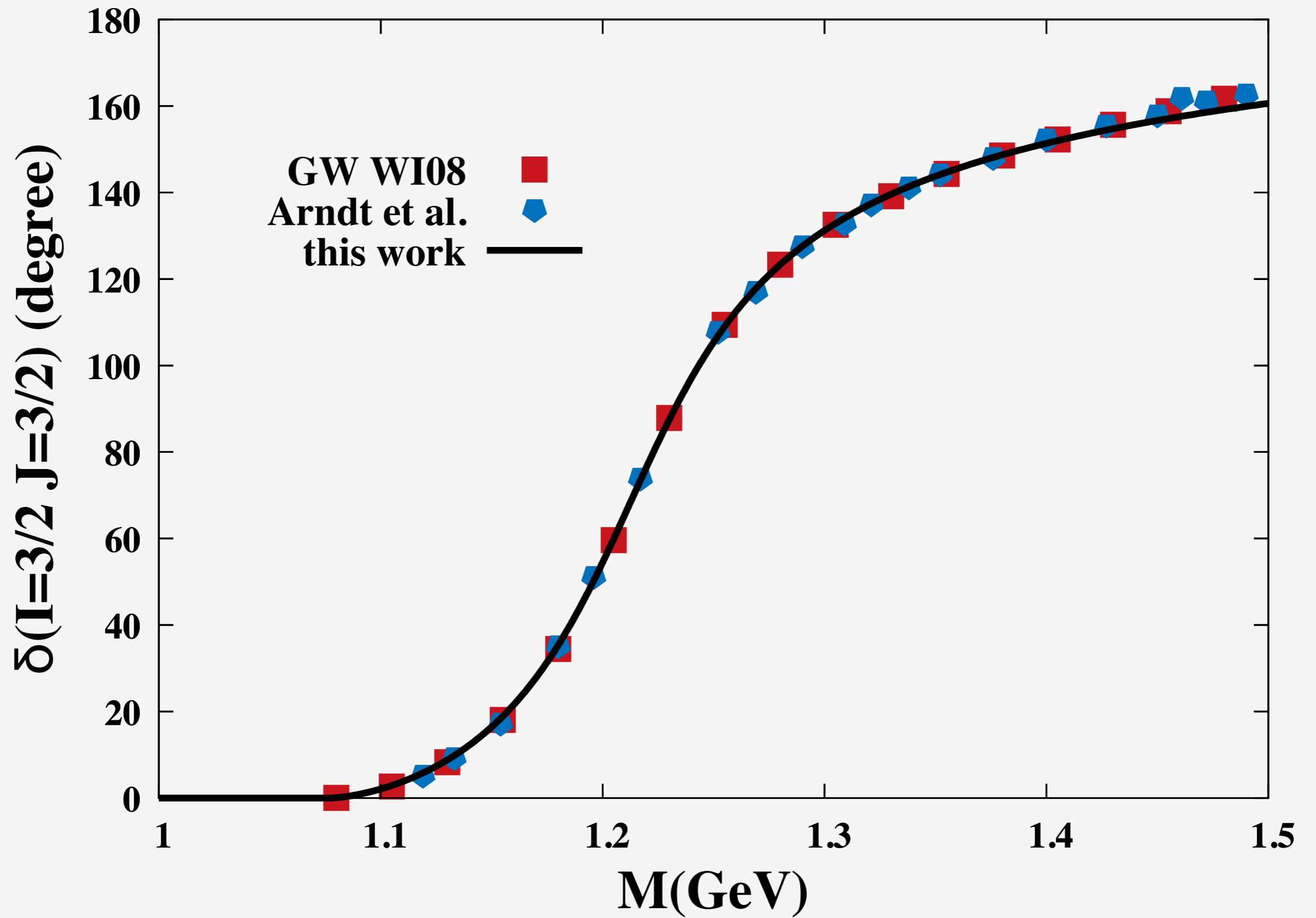
$$I(J^P) = \frac{3}{2}(\frac{3}{2}^+)$$

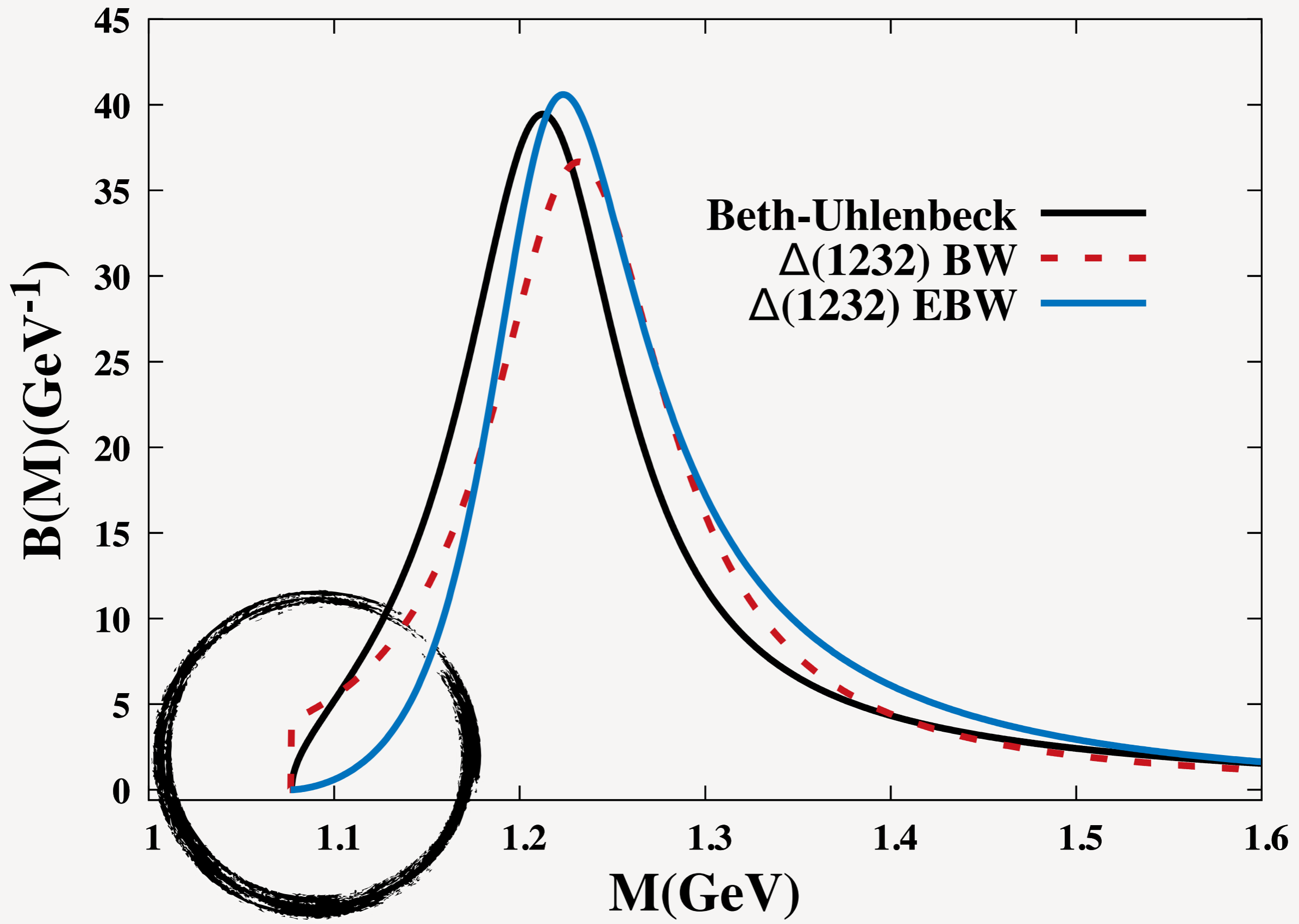
Breit-Wigner mass (mixed charges) = 1230 to 1234 ( $\approx 1232$ ) MeV

Breit-Wigner full width (mixed charges) = 114 to 120 ( $\approx 117$ ) MeV

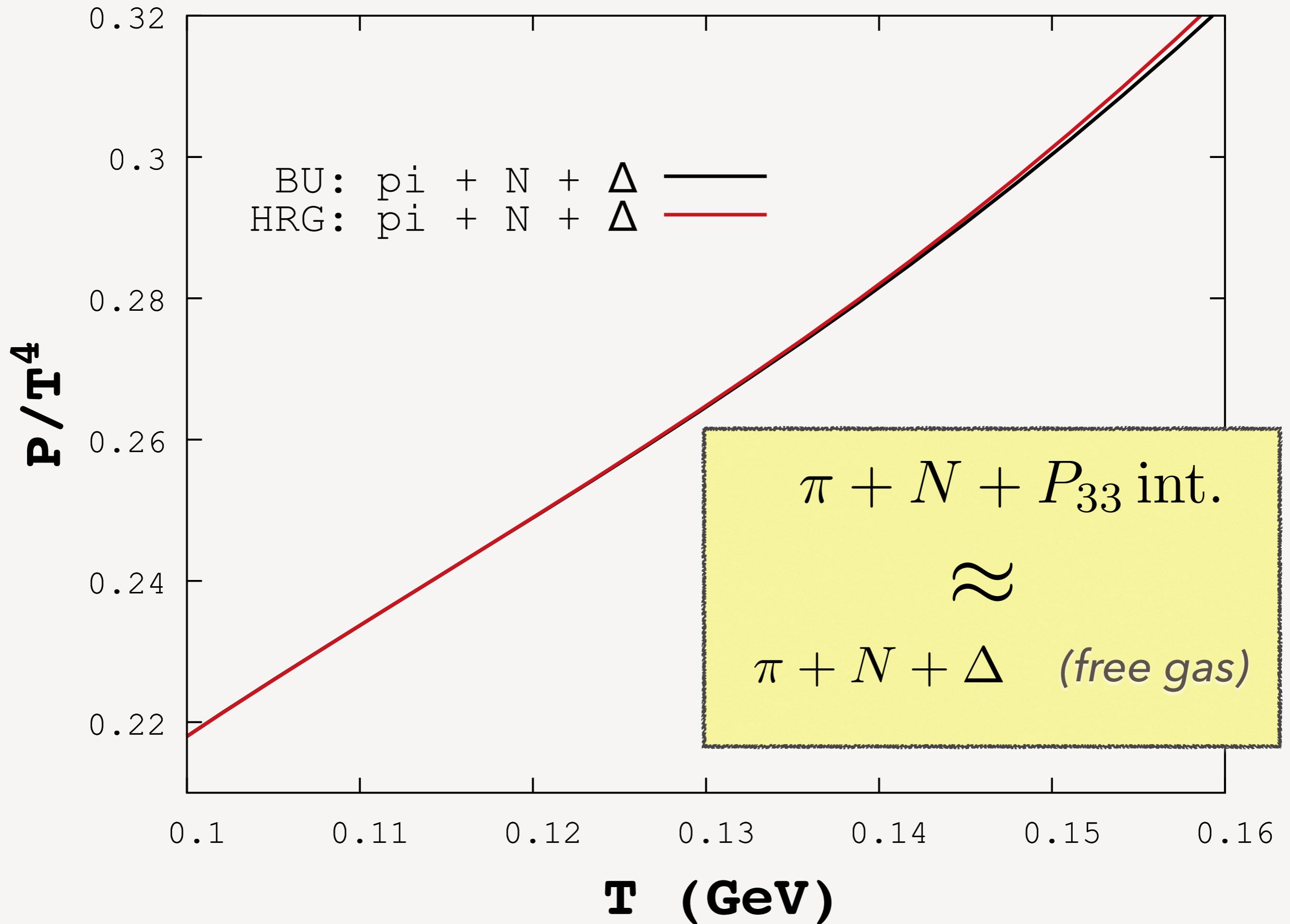
Re(pole position) = 1209 to 1211 ( $\approx 1210$ ) MeV

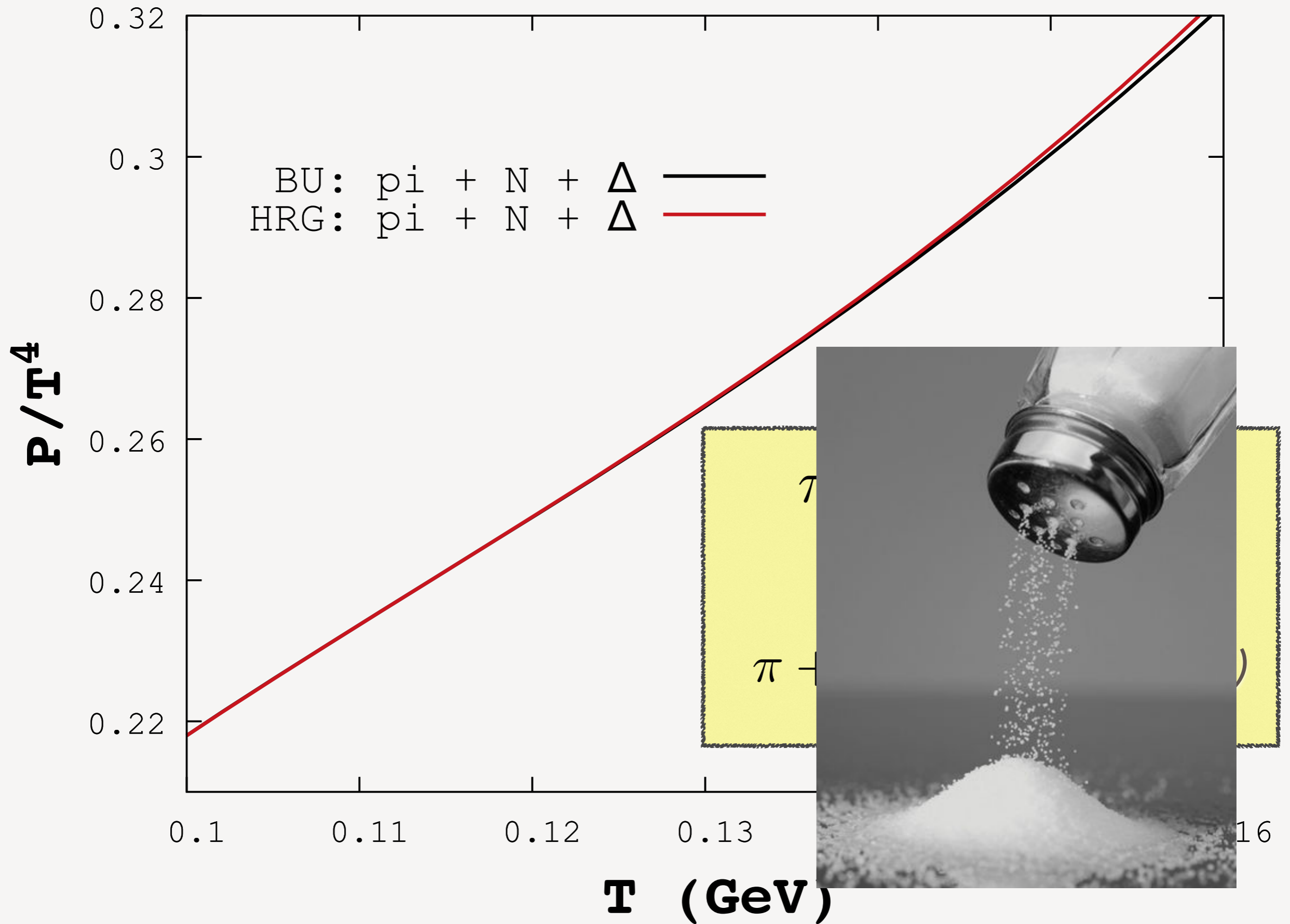
$-2\text{Im}(\text{pole position}) = 98$  to  $102$  ( $\approx 100$ ) MeV









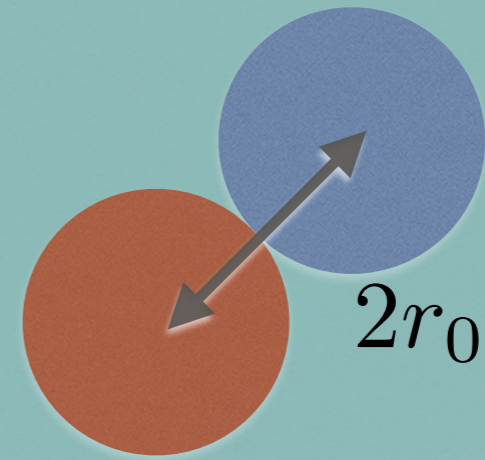


# REPULSIONS AND EXCLUDED VOLUMES

# MODELING SHORT-RANGE REPULSION BETWEEN HADRONS

- Excluded volume approach

eigenvolume  $v_0$



- Beth-Uhlenbeck approach

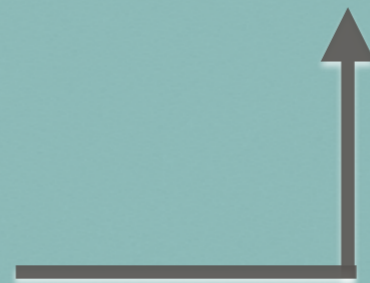
- 1) QM problem of a hard-core potential  $\rightarrow \delta_l(q a_S)$
- 2) Thermodynamics by S-matrix

# FORMULATION #1

$$P(T, \mu) = P^{id}(T, \tilde{\mu} = \mu - v_0 P(T, \mu))$$



To be solved self-consistently.



non-linear effects

# FORMULATION #1

$$P(T, \mu) = P^{id}(T, \tilde{\mu} = \mu - v_0 P(T, \mu))$$

$$\mu = \mu_B B + \mu_S S + \mu_Q Q$$



conserved  
charges

# FORMULATION #1

$$P(T, \mu) = P^{id}(T, \tilde{\mu} = \mu - v_0 P(T, \mu))$$

$$\mu = \mu_B B + \mu_S S + \mu_Q Q$$

conserved  
charges

$\mu_{\text{sym}}$

multiplicity

# FORMULATION

## #2

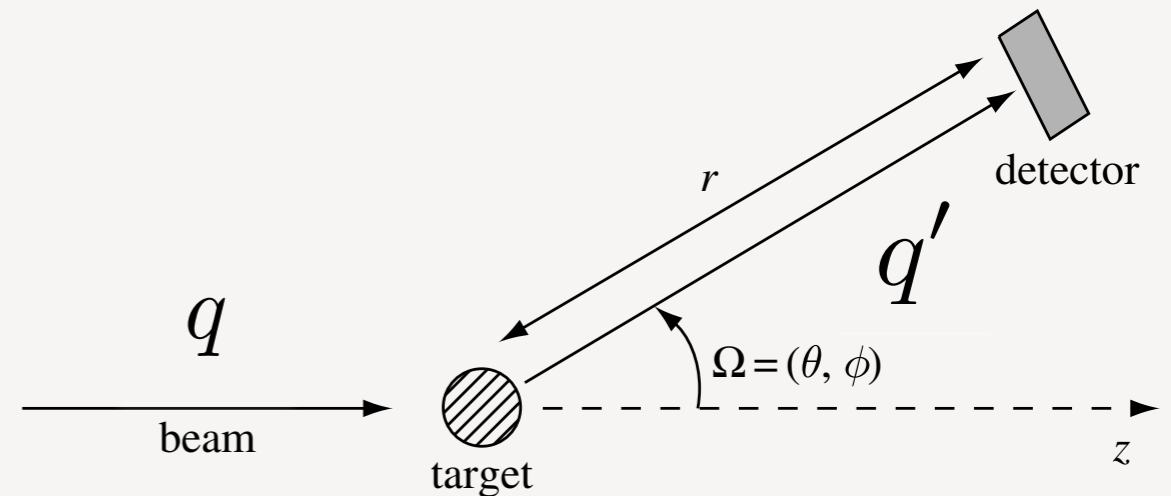
- Starting point:  
hard-core potential in QM

$$V = \infty \quad r < a$$

$$= 0 \quad r > a$$

$$\tan(\delta_l) = \frac{j_l(qa)}{n_l(qa)}$$

$$\psi^q(r \rightarrow \infty) \longrightarrow e^{iqr \cos(\theta)} + \frac{e^{iqr}}{r} \sum_l (2l + 1) P_l \frac{e^{i\delta_l}}{q} \sin(\delta_l)$$



Momentum  $q$  enters through the scattering Schroedinger equation with a centrifugal term ( $l$ -dependence)



# FORMULATION #2

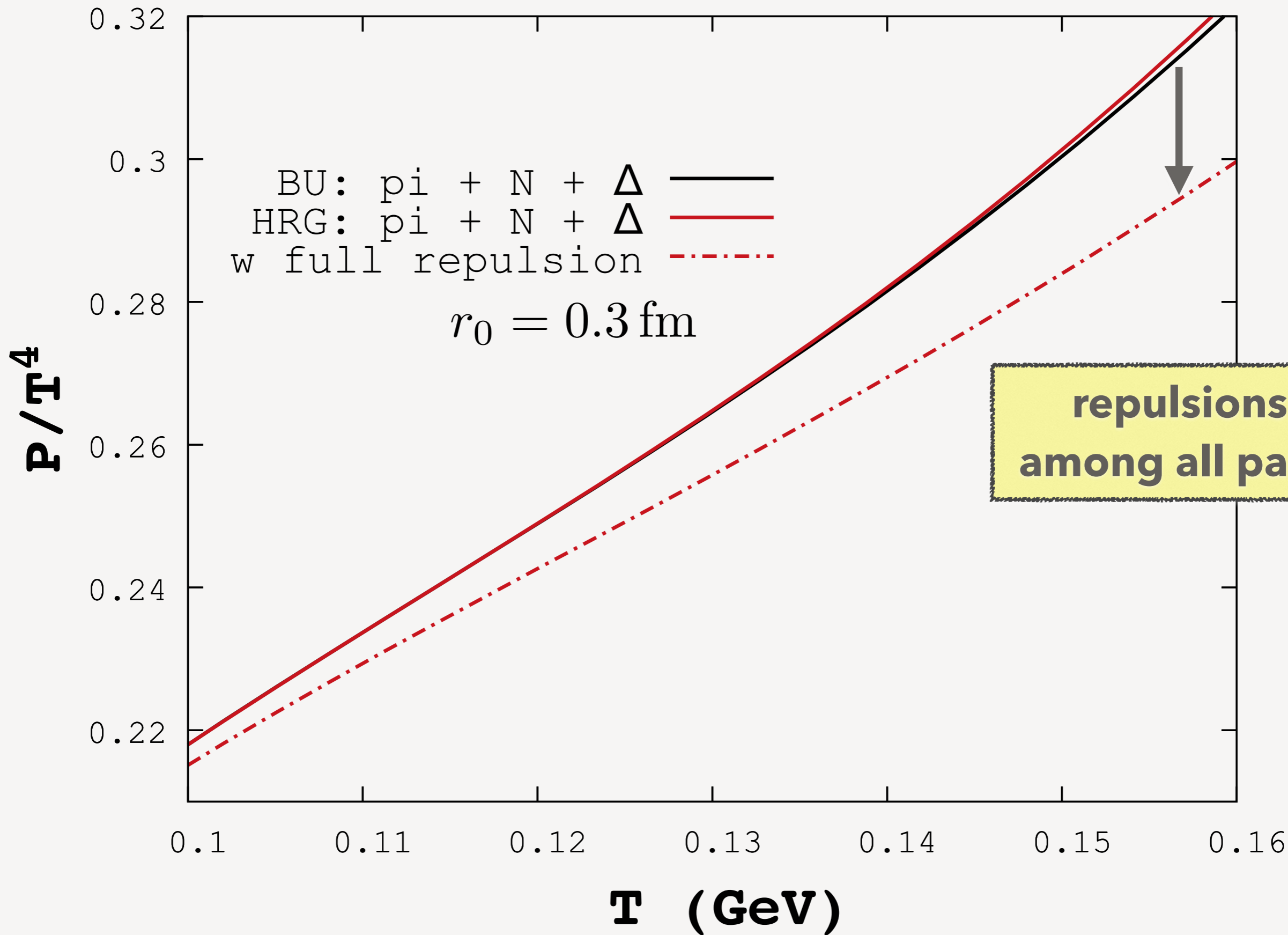
$$\tan(\delta_l) = \frac{j_l(qa)}{n_l(qa)}$$

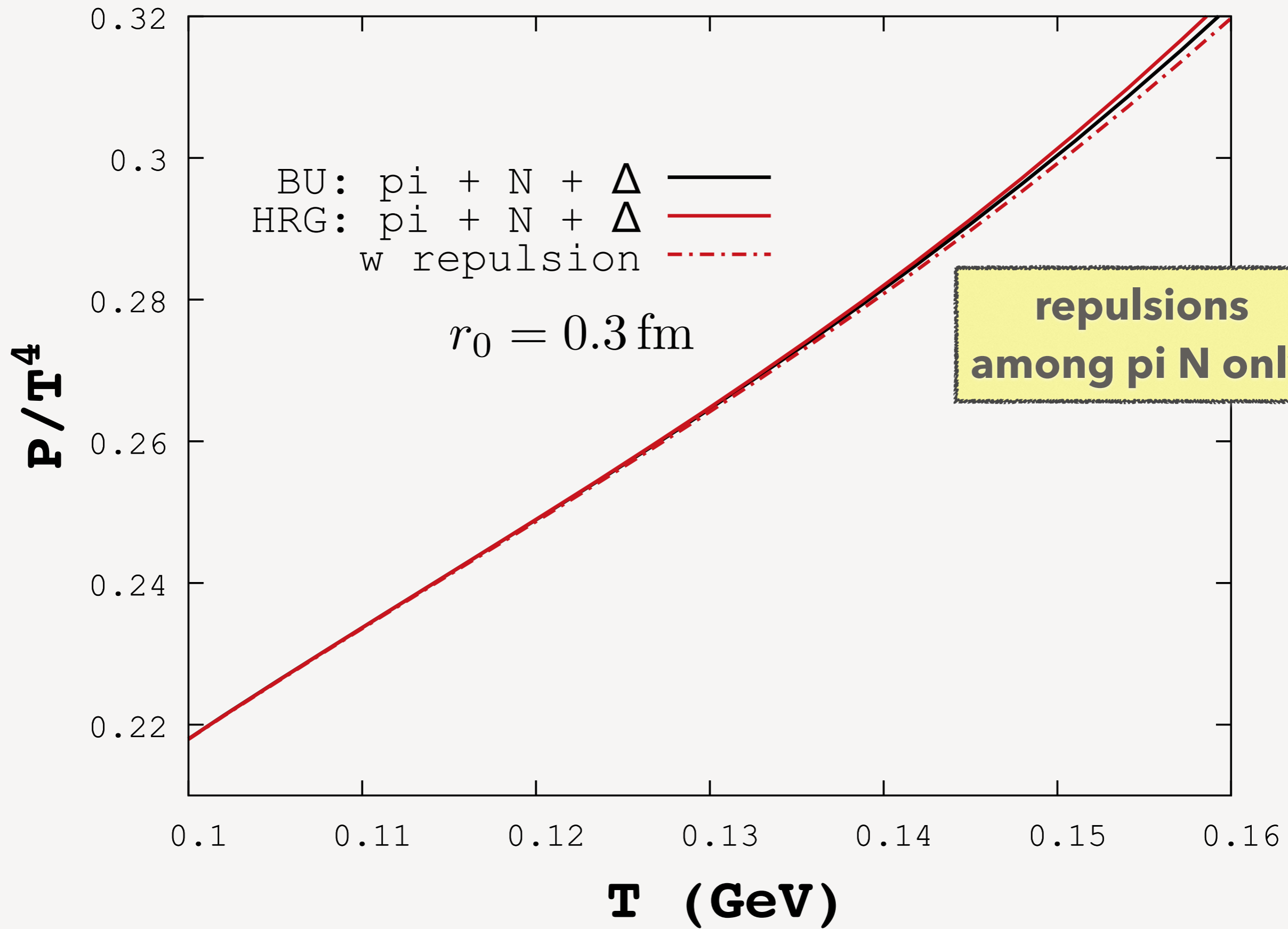
*for small  $x = qa$  (near threshold)*

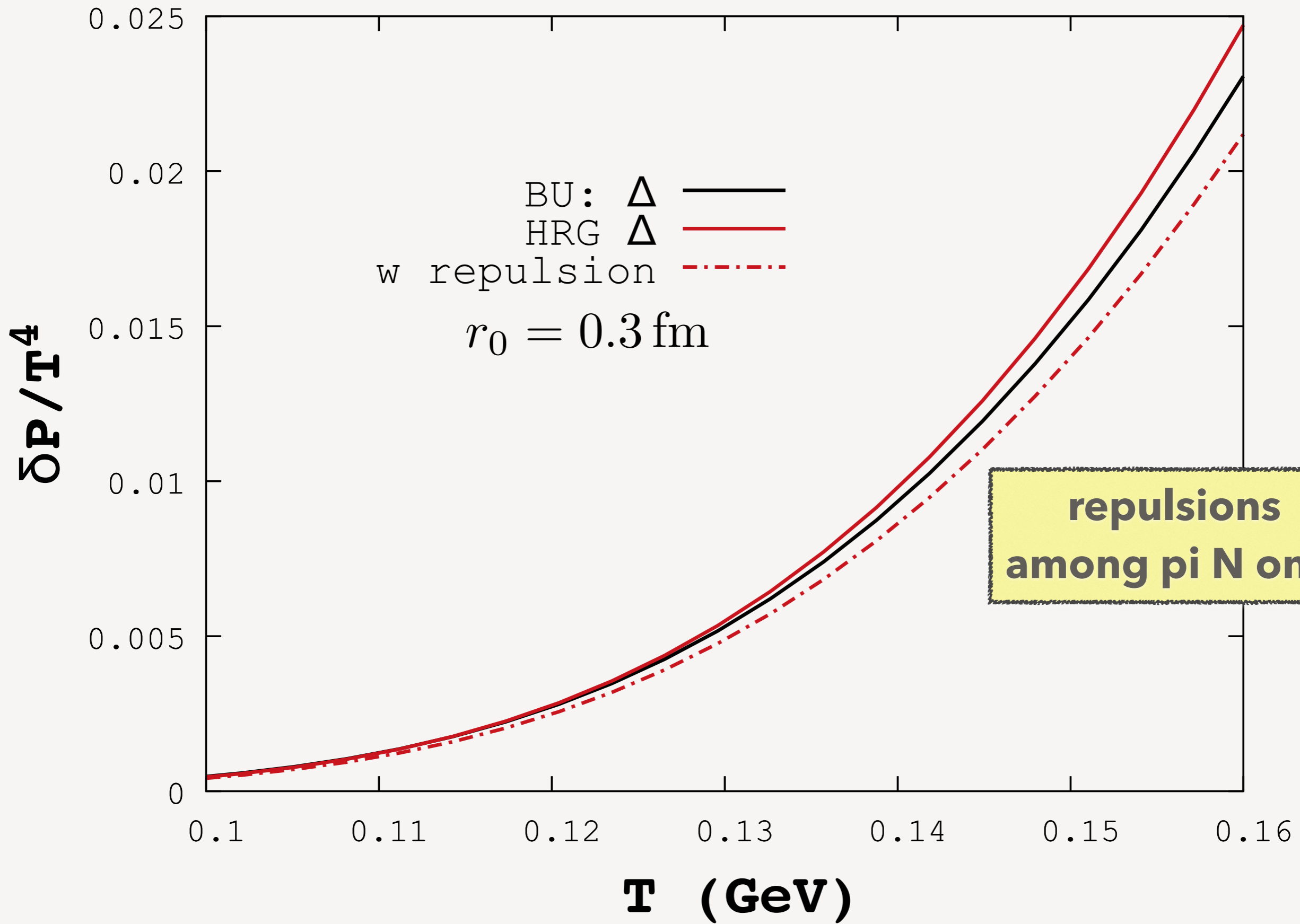
$$\tan(\delta_l) \rightarrow \frac{-x^{2l+1}}{(2l+1)((2l-1)!!)^2}$$

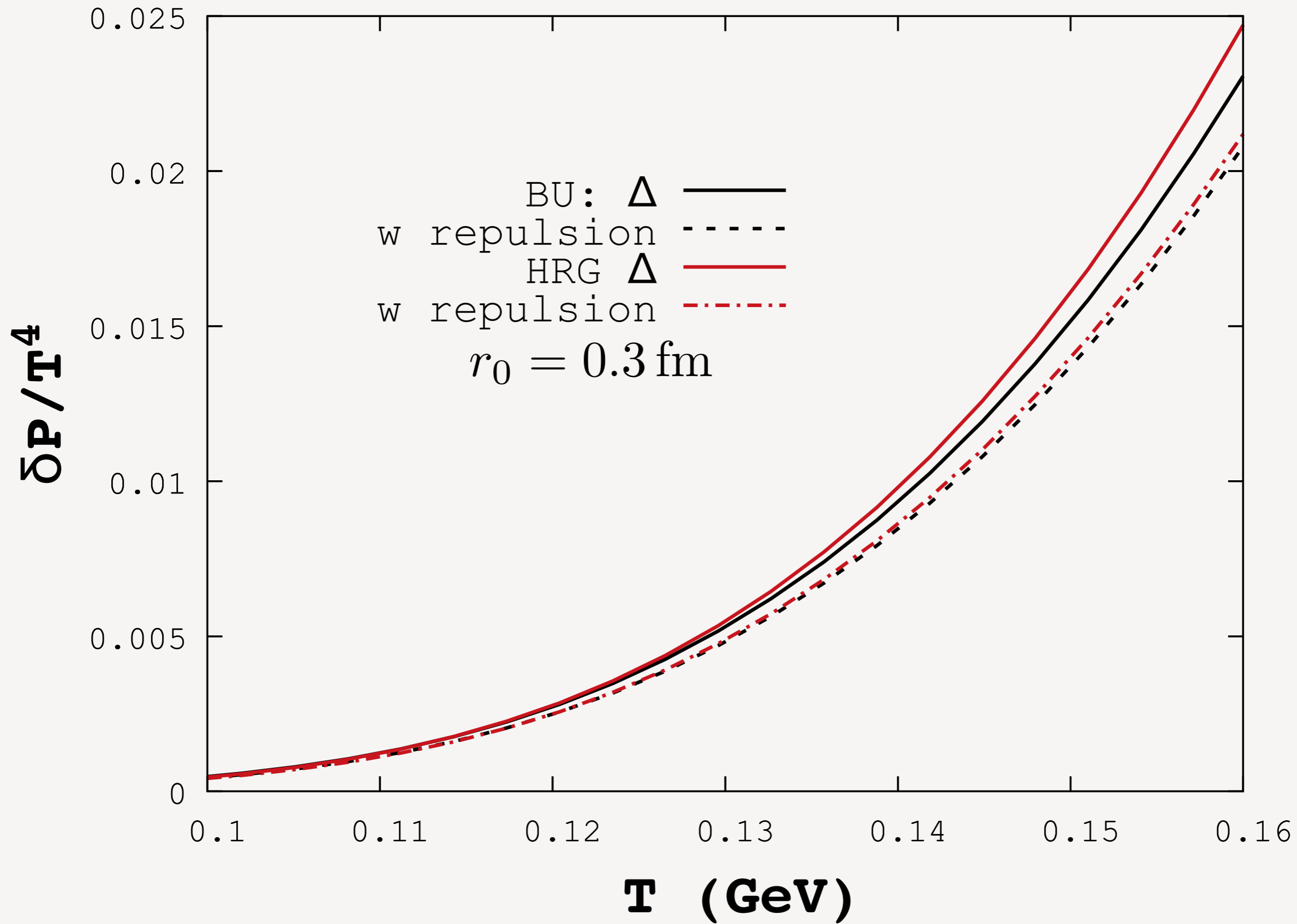
$$\delta_l \propto (qa)^{2l+1}$$

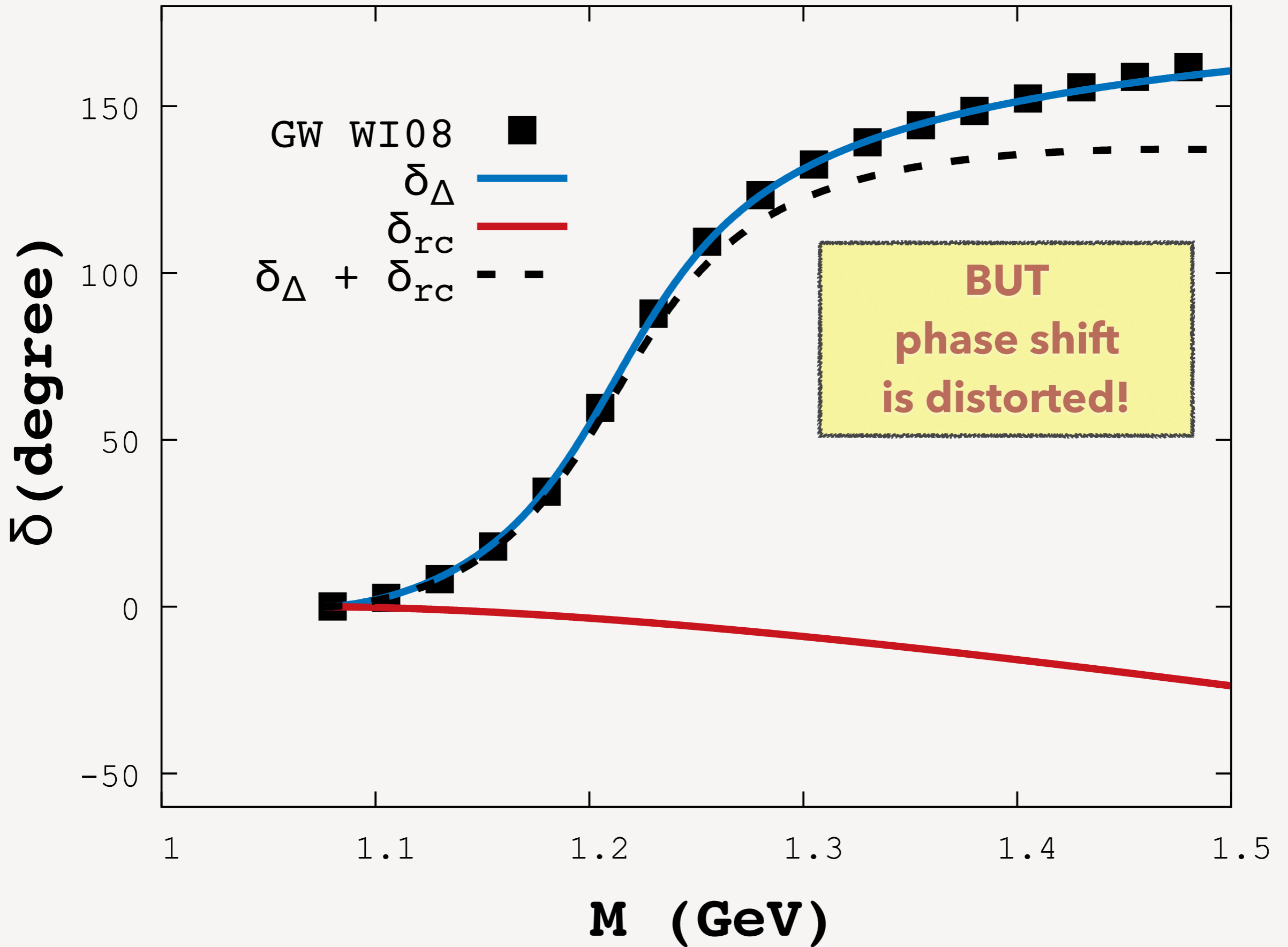
*(near threshold)*

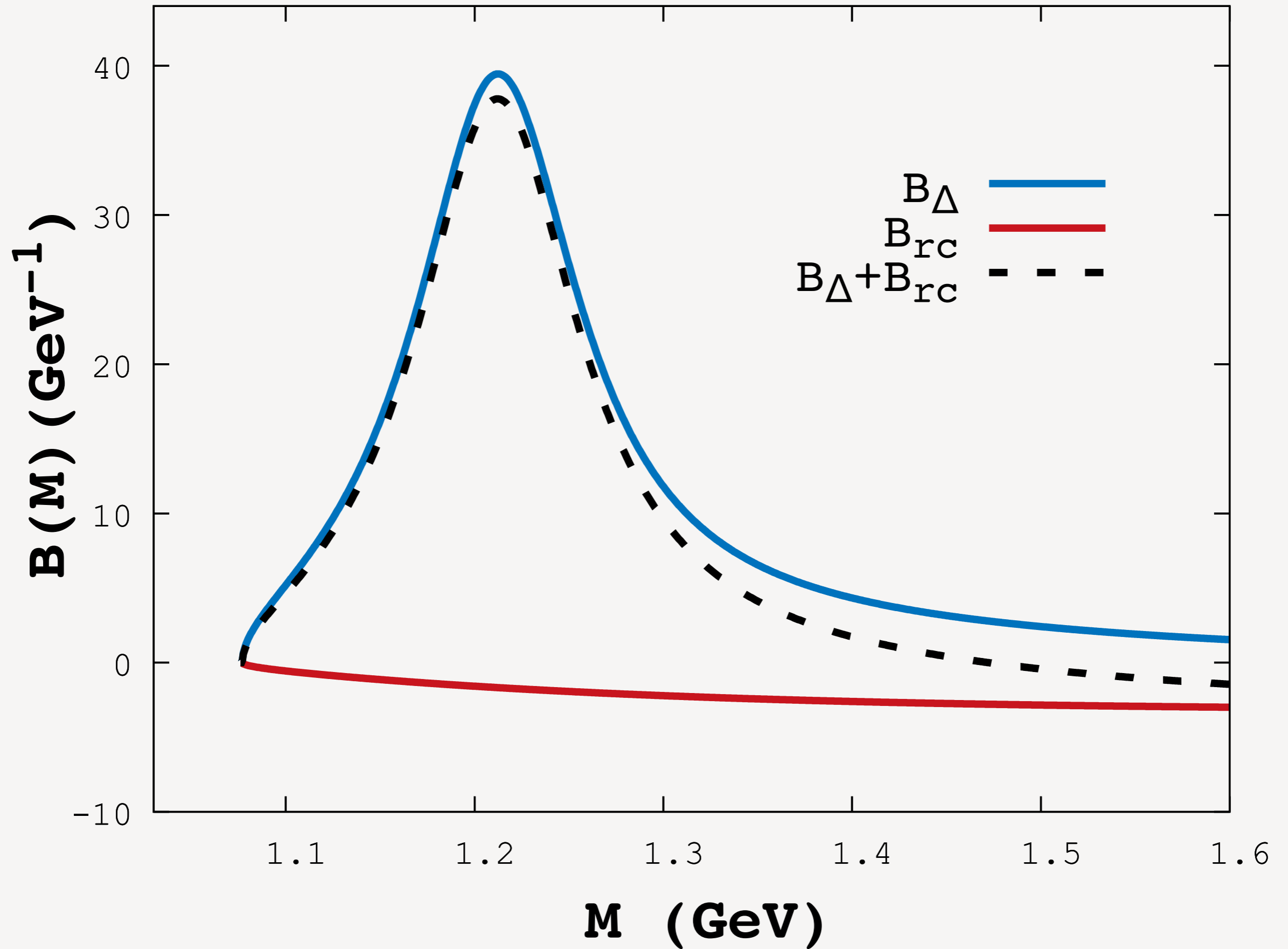






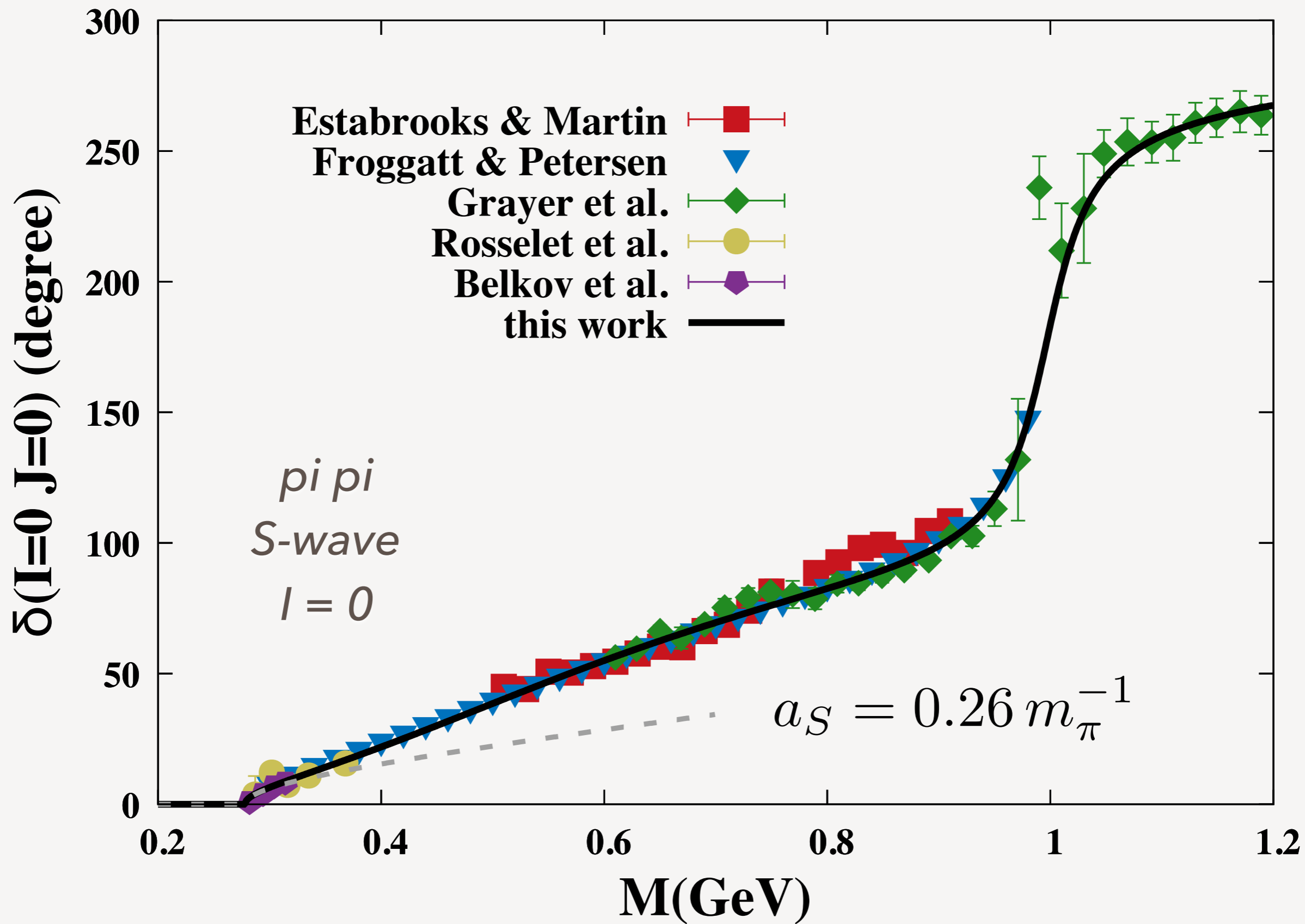


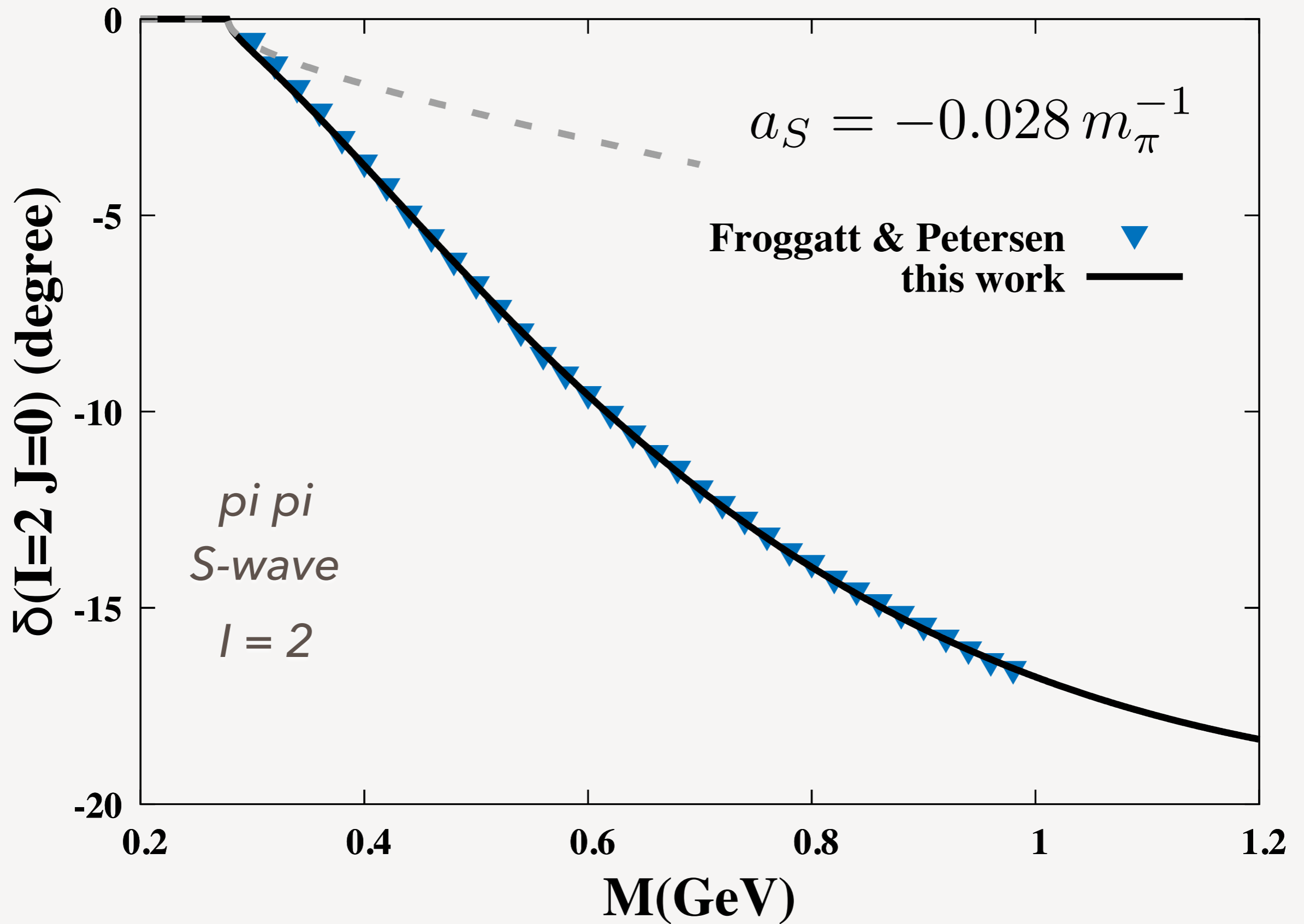


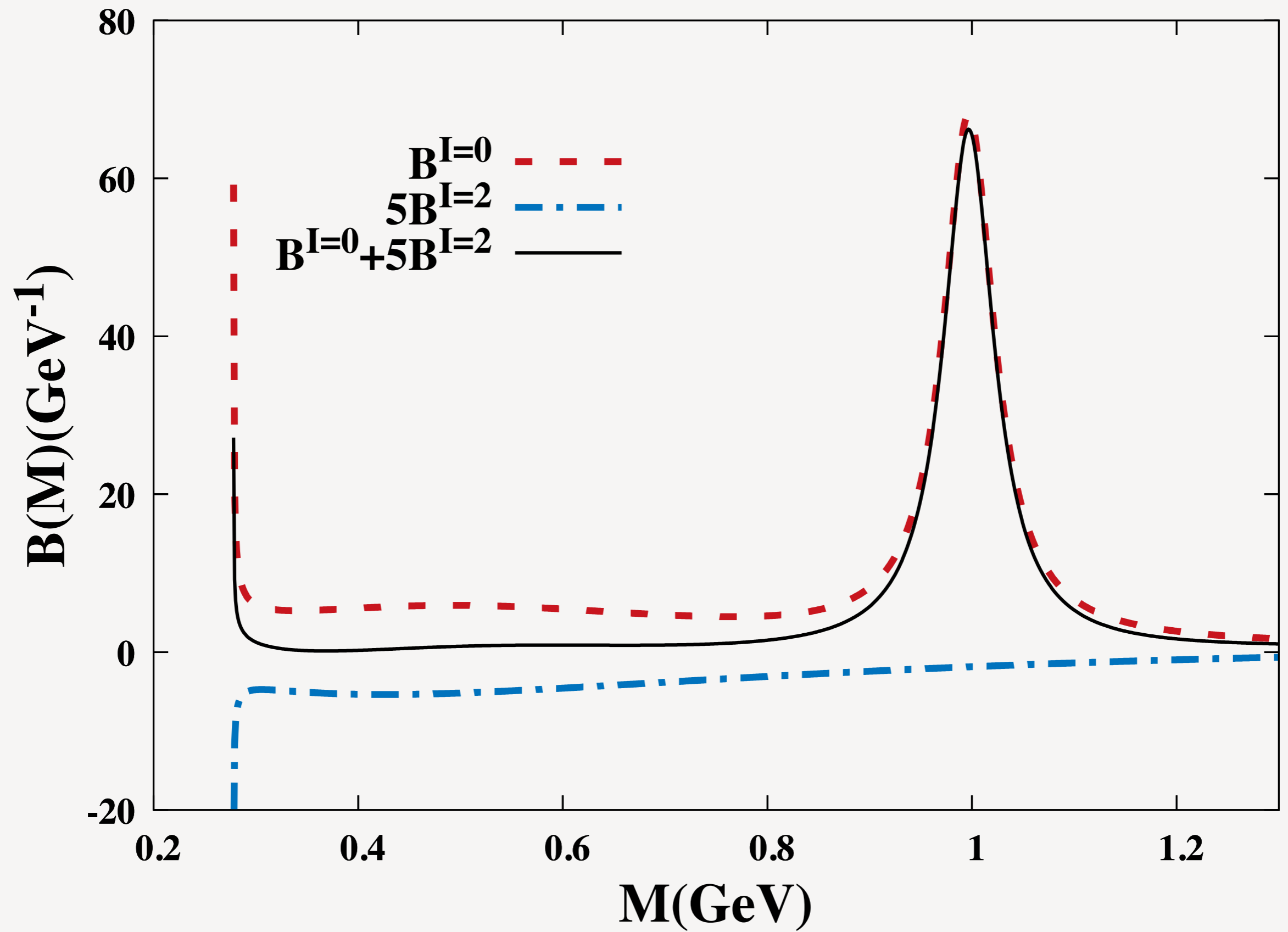


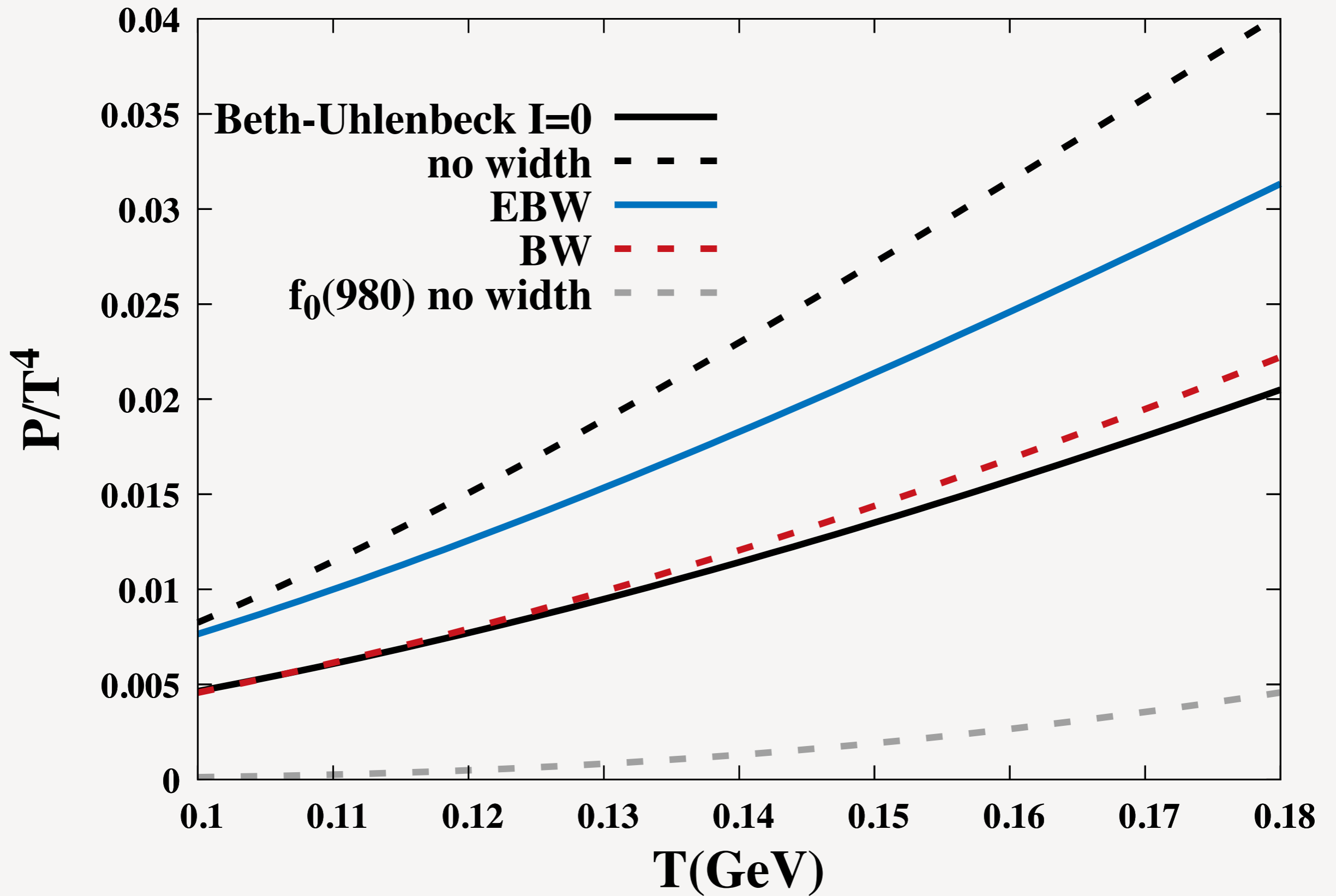
# PI PI SCATTERING

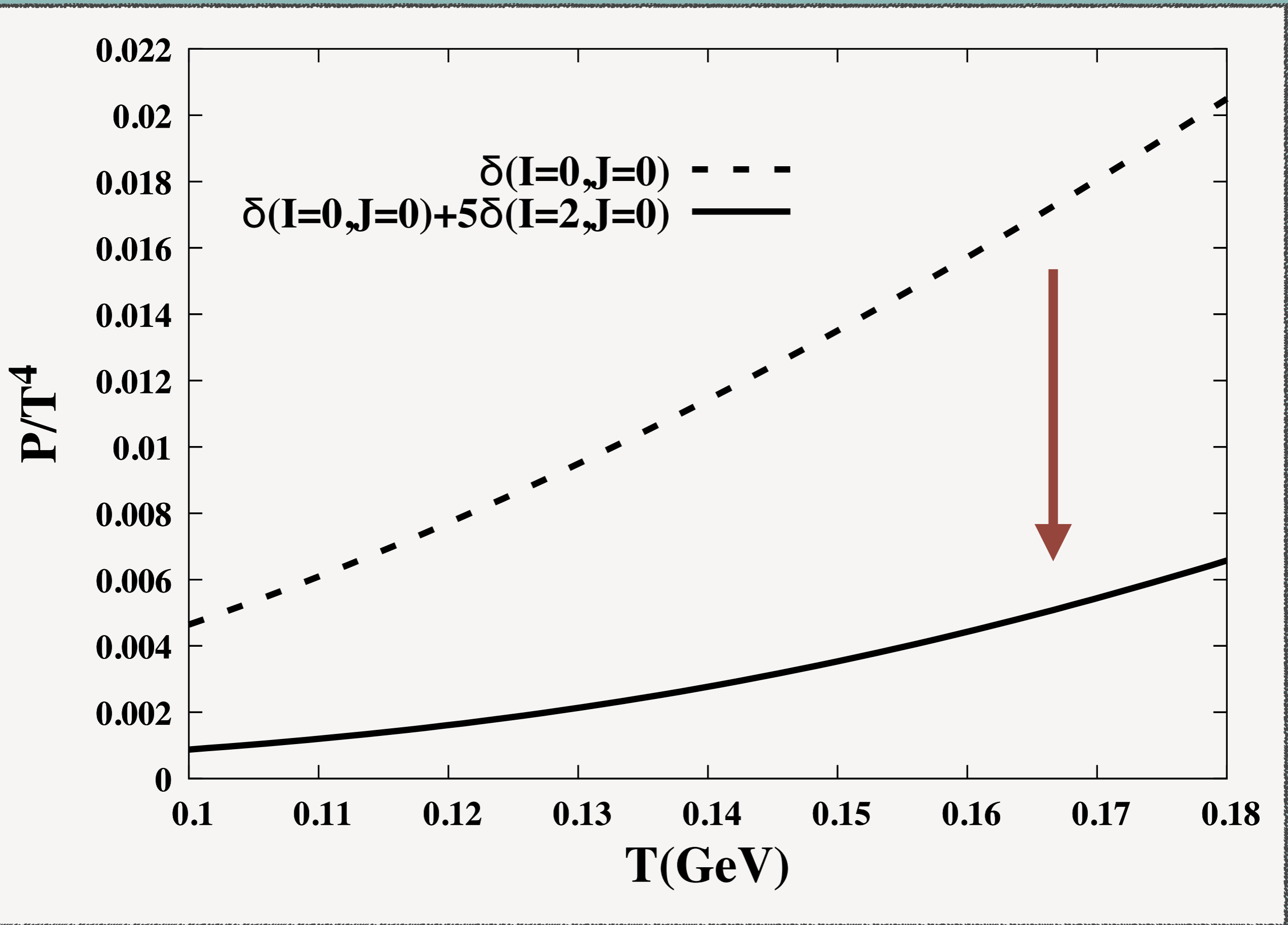


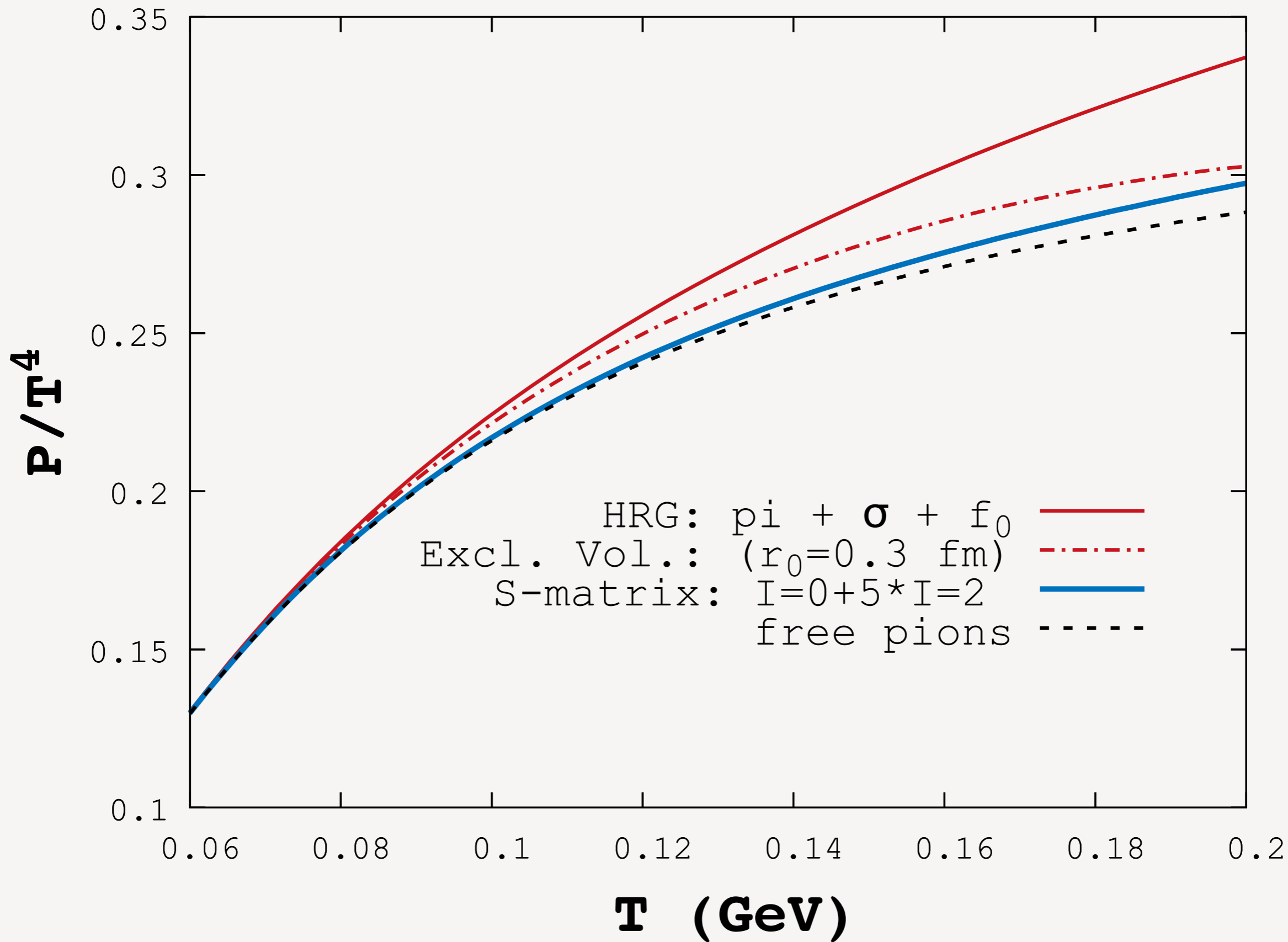












# SUMMARY

- Repulsion strengths are channel-dependent.
- They are automatically incorporated in phase shifts, no need for additional excluded volume
- Flexibility in modeling  
e.g. scattering length is isospin dependent =>  
channel dependent!

$$a_{I=0}^{\pi\pi} > a_{I=2}^{\pi\pi}$$