



”About Neutrinos and their oscillations”

Ewa Rondio

National Centre for Nuclear Research
Warsaw, Poland



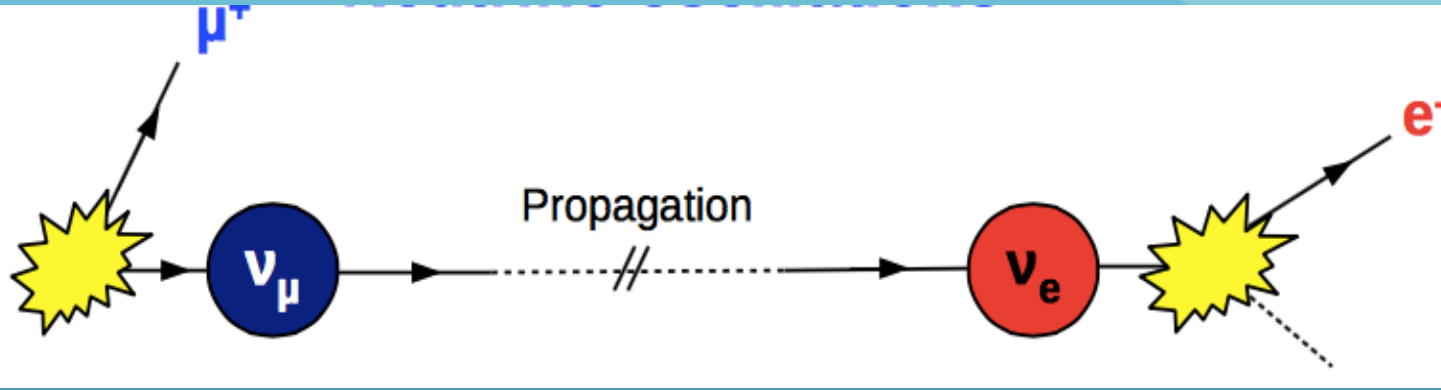
12-th Polish Workshop on Relativistic Heavy-Ion Collisions

Kielce, November 5-th , 2016

What is “oscillation”??

Quarks	u up	c charm	t top
	d down	s strange	b bottom
Leptons	ν_e e- Neutrino	ν_μ μ - Neutrino	ν_τ τ - Neutrino
	e electron	μ muon	τ tau
	I	II	III

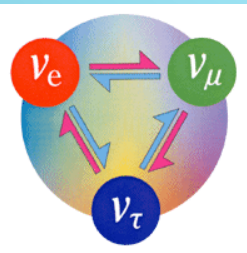
The Generations of Matter



The flavour eigenstate is not the same as the mass eigenstate

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \times \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Mixing matrix (Pontecorvo-Maki-Nakagawa-Sakata=PMNS)



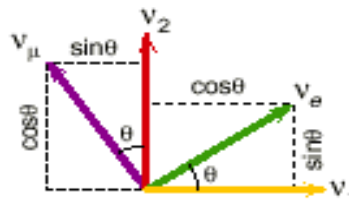
First look at two neutrino case

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 L}{E} \right)$$

L – dist. to the detector
E - neutrino energy

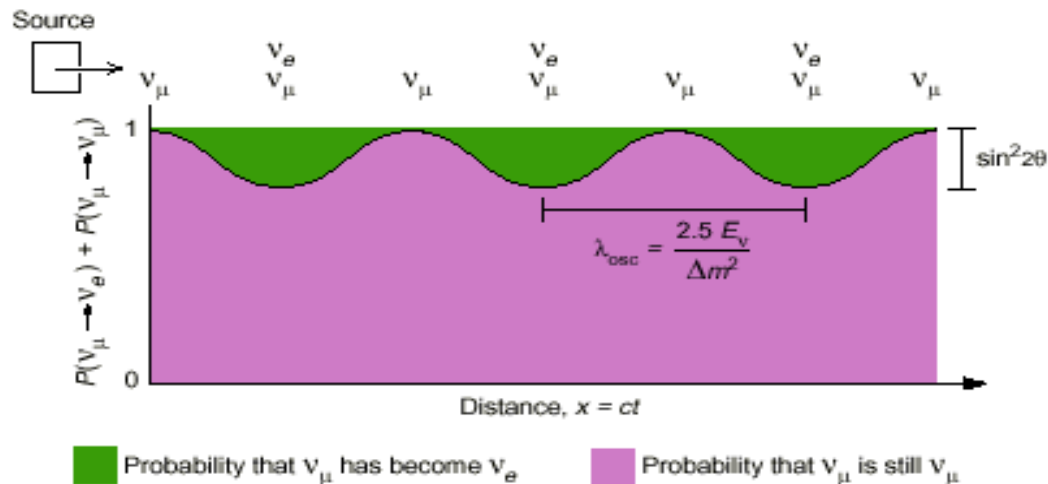
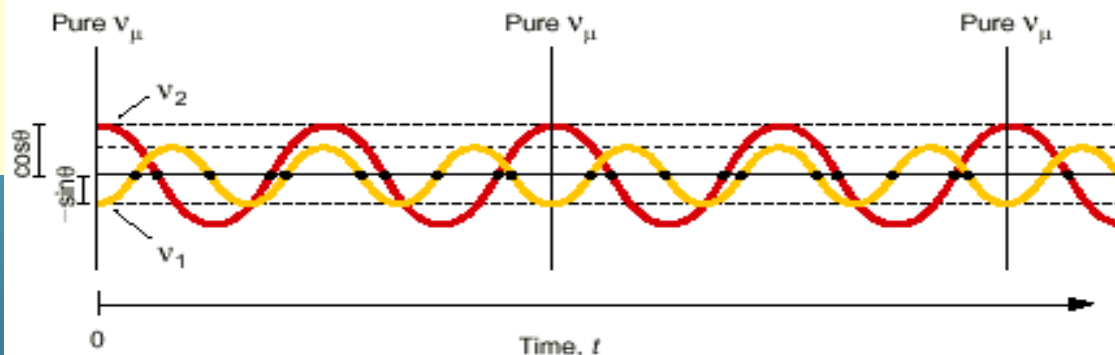
Maximal effect when

$$\sin^2 \left\{ \frac{1.27 \Delta m^2 L}{E} \right\} = 1$$



$$\nu_e = \cos \vartheta \nu_1 + \sin \vartheta \nu_2$$

$$\nu_\mu = -\sin \vartheta \nu_1 + \cos \vartheta \nu_2$$



Sensitivity to oscillations

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2 2\theta \sin^2 \left(\frac{1.27 \Delta m^2 L}{E_\nu} \right)$$

ν energy - E and distance L define range of sensitivity

	E_ν (MeV)	L (m)	Range of Δm^2
Supernovae	<100	> 10^{19}	10^{-19} - 10^{-20}
Solar	<14	10^{11}	10^{-10} ???
Atmospheric	>100	10^4 - 10^7	10^{-3} - 10^{-4}
Reactor	<10	< 10^6	10^{-5}
Accelerator - SB	>100	10^3	10^{-1}
Accelerator - LB	>100	< 10^6	10^{-3}

Two mass differences and three neutrino types oscillating

→ full description in 3x3 oscillation matrix,

→ studies in many experiments to get full picture....

Neutrino oscillations

– picture as of today

FLAVOR

PMNS mixing matrix

MASS

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \begin{pmatrix} \cos\theta_{12} & 0 & \sin\theta_{12}e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin\theta_{13}e^{-i\delta} & \sin\theta_{13} & \cos\theta_{13} \end{pmatrix} \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ \sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

2-3 **Sub-leading** **1-3** **1-2**

„atmospheric”
SK, K2K, T2K, MINOS
Nova

$$\Delta m_{31}^2 = \begin{cases} 2.53^{+0.08}_{-0.10} \\ -(2.40^{+0.10}_{-0.07}) \end{cases} \times 10^{-3} \text{ eV}^2$$

CHOOZ,
DayaBay,
Reno,
DbIChooz,
T2K

$$\theta_{12} = 34^\circ \pm 1^\circ$$

$$\theta_{23} = 40^\circ + 5^\circ / - 2^\circ$$

$$\theta_{13} = \boxed{9.1^\circ \pm 0.6^\circ!}$$

Based on PDG 2012

„solar”
SNO, KamLand,
SK, Borexino

$$\Delta m_{21}^2 = (7.62 \pm 0.19) \times 10^{-5} \text{ eV}^2$$

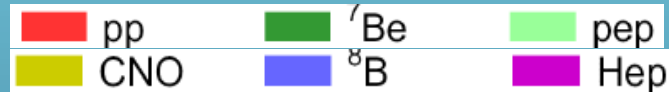
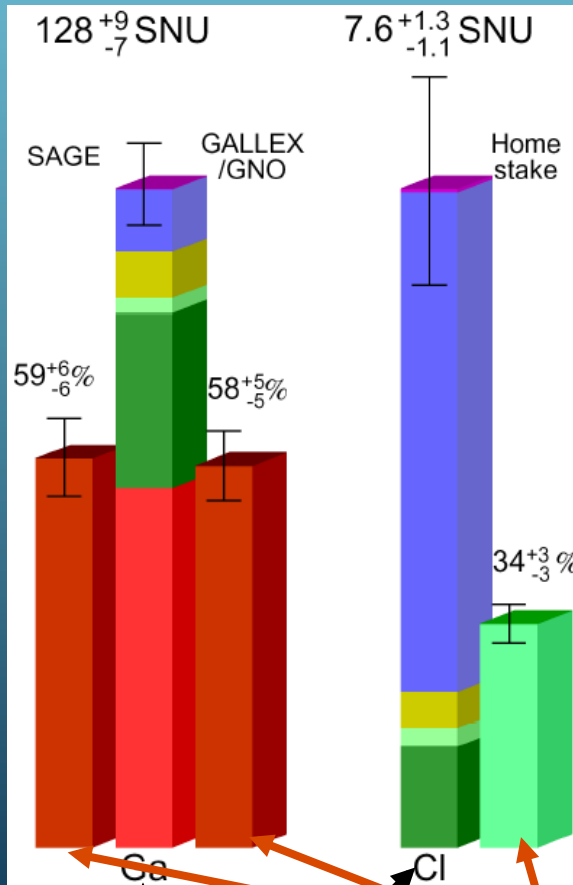
parameter θ_{13}
found to be non zero !!!

mixing angles, squared mass differences, CP violation phase - fundamental parameters of nature

* $\Delta m_{ji}^2 = m_j^2 - m_i^2$
 Two free parameters for the three Δm^2 's.
 ($\Delta m_{31}^2 = \Delta m_{21}^2 + \Delta m_{32}^2$)

What we know about each sector?

1-2 – solar neutrinos



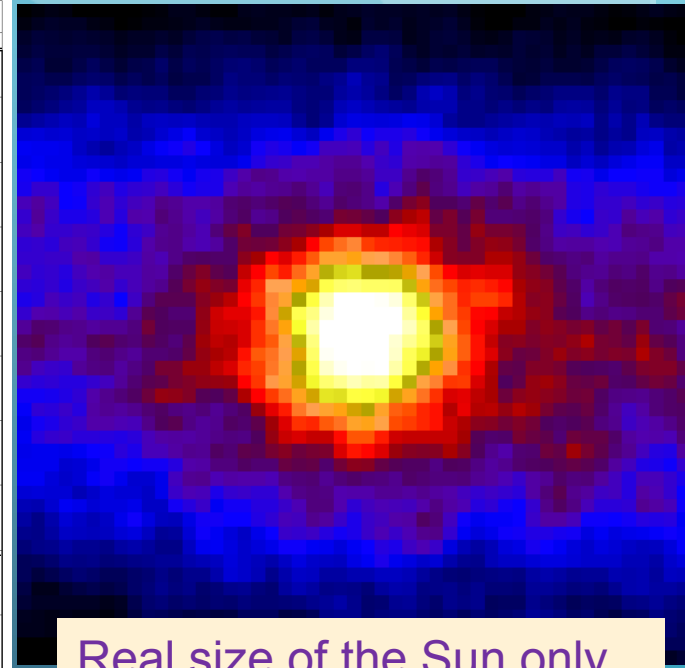
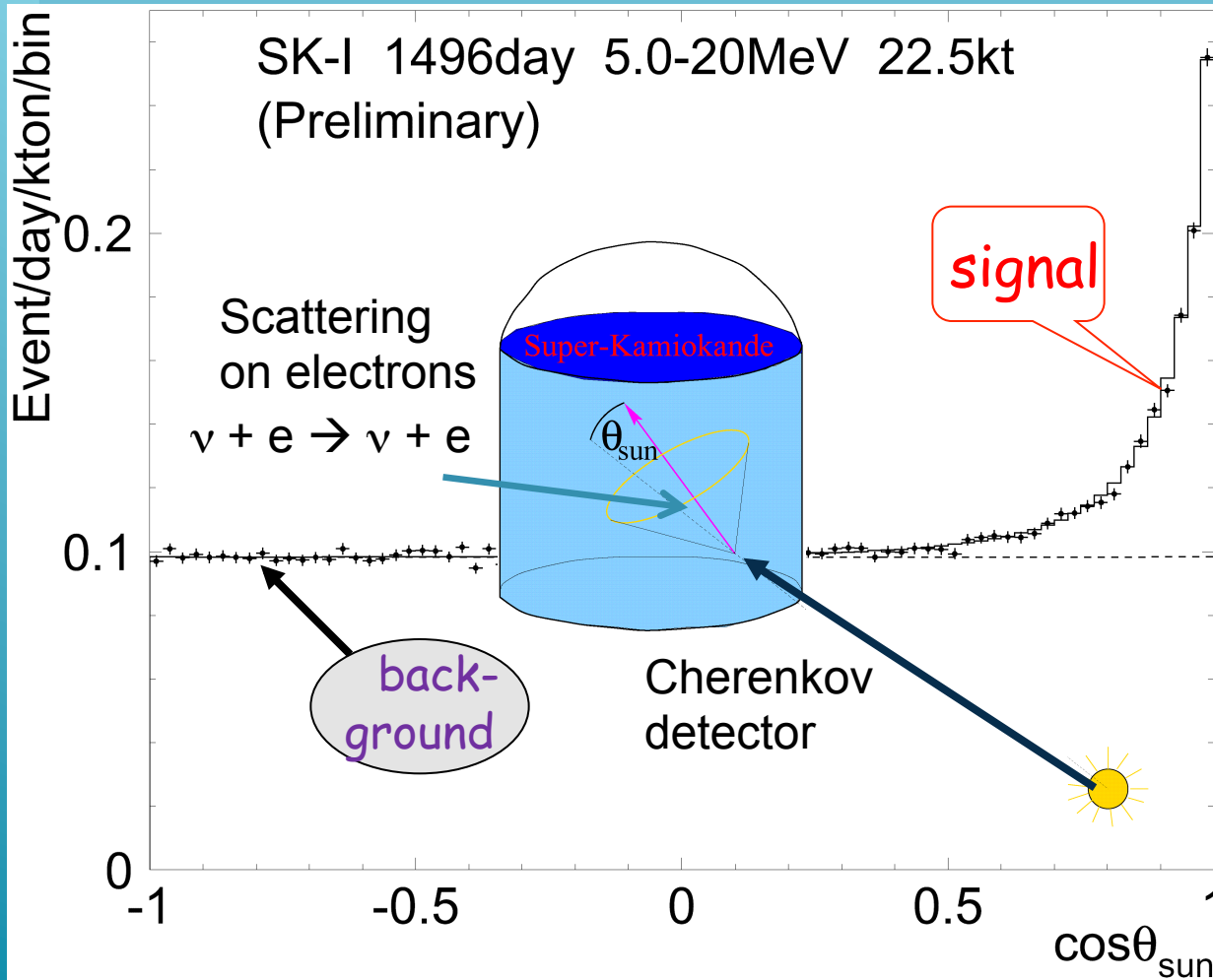
Contribution calculated from „SSM” - Standard Solar Model (SSM):

- observations in radio-chemical experiments (counting interactions, no time and direction information)
- First data – Davis experiment

Neutrino deficit

measurements

Super-K: neutrino's arriving from the Sun ⁷



Real size of the Sun only
– ½ pixel. The smearing is
due to electron scattering

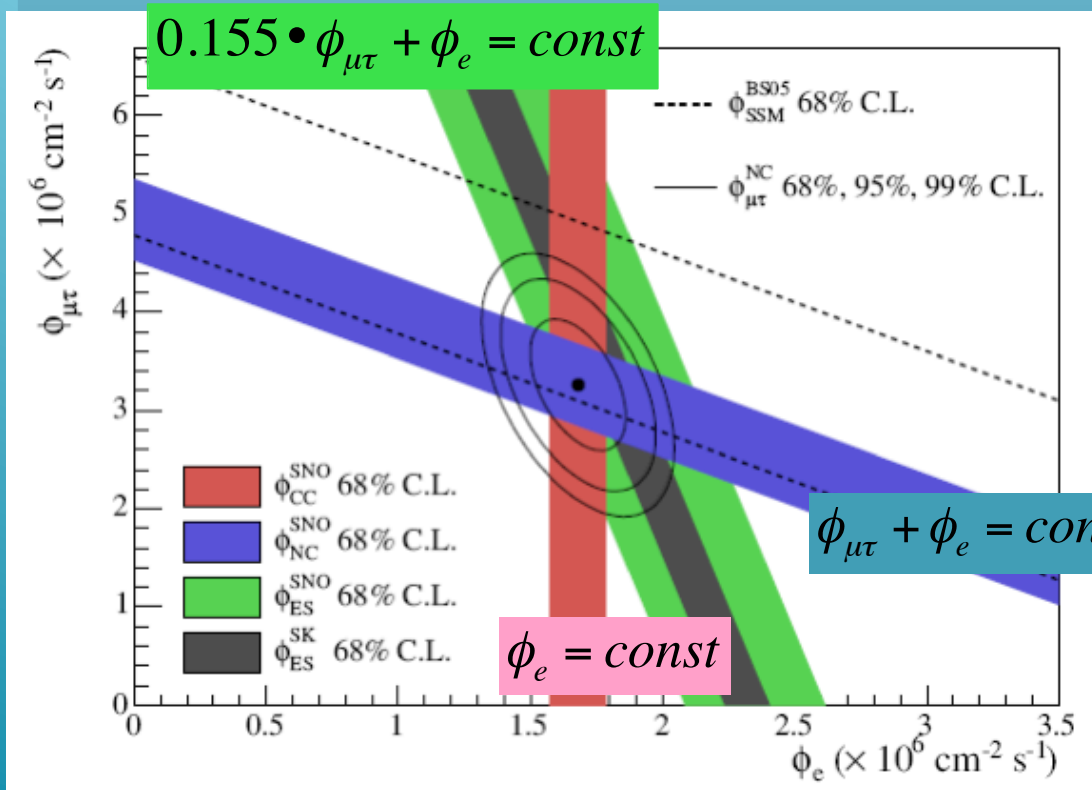
**neutrino deficit
....again...**

Here also observed number of neutrinos is too small
Observed about ½ of what was expected

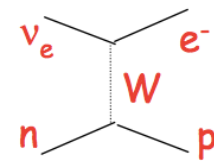
Solution – SNO experiment - heavy water

(½ Nobel Prize 2015)

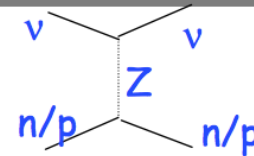
SNO experiment : measurement sensitive to 3 reactions
Including sensitivity to Neutral Currents → it “sees all neutrinos”



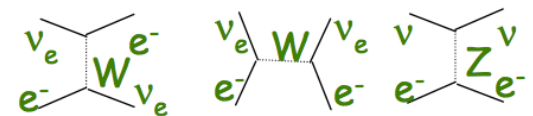
CC



NC

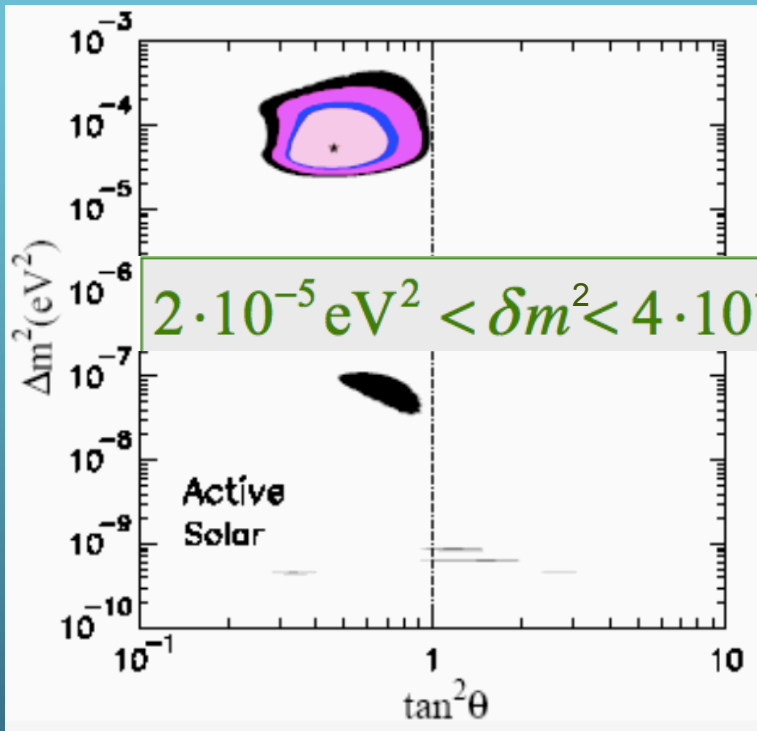


ES



Here it is seen directly that ν_e flux is smaller than expected from Solar Model, but sum over 3 ν flavours gives expected flux

Parameters for 1-2 sector:



- From sensitivity for Sun-Earth distance – expected 10^{-10}eV^2
- Why it is different?

we needed to make correction for

MATTER EFFECTS

→ Modification of potential due to electrons density in matter

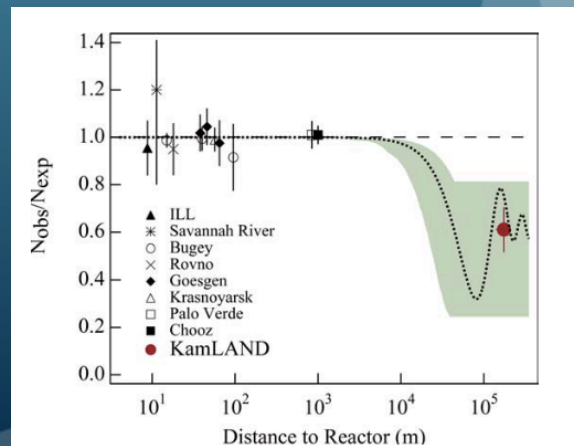
$$\Delta m^2_{matter} = \sqrt{(\Delta m^2 \cos 2\theta - A)^2 + (\Delta m^2 \sin 2\theta)^2}$$

KAMLAND – sector 1-2 for anti-neutrinos from reactors

Obs/exp = 0.631 +/- 0.014 (stat)
 +/- 0.027 (syst)

Corresponding to
 exclusion of non-oscillation at
 10.2 σ CL

no matter effects !!!

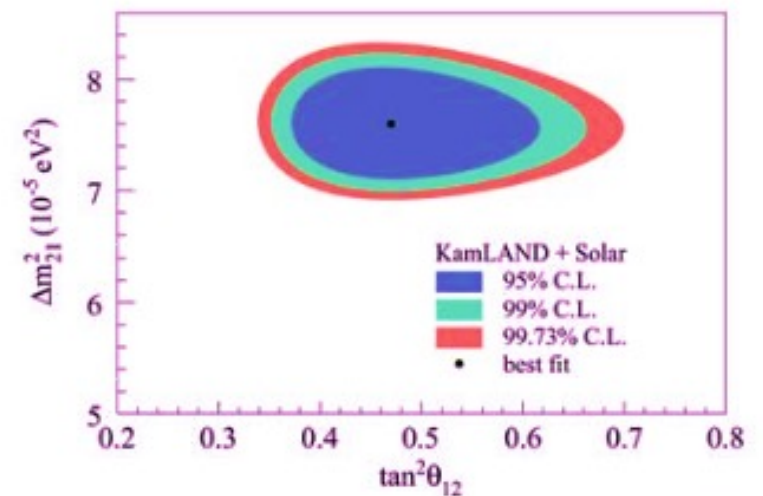
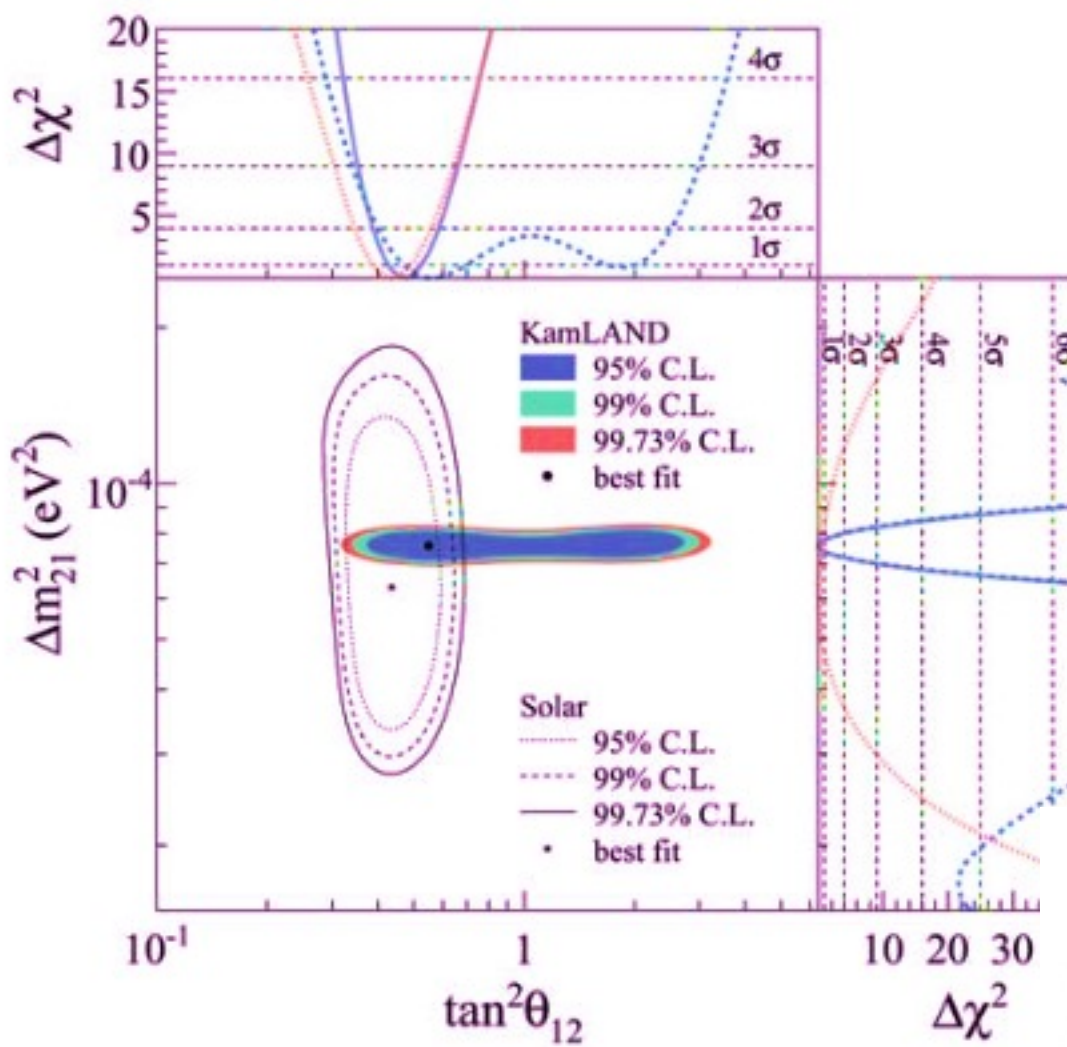


Determination of 1-2 mixing

Assuming same values for Neutrinos and anti-neutrinos

We obtain:

Δm_{21}^2 (10^{-5} eV^2)	$\tan^2 \theta_{12}$
$7.53^{+0.18}_{-0.18}$	$0.436^{+0.029}_{-0.025}$

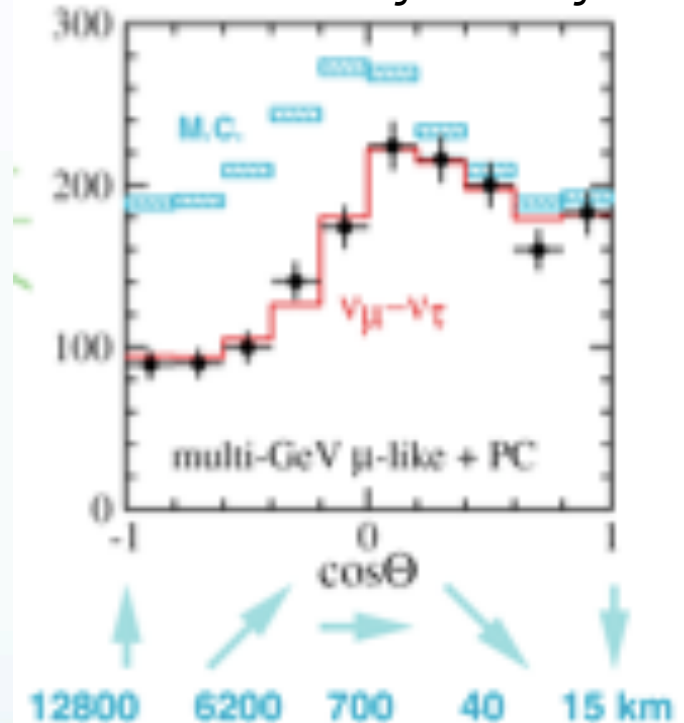
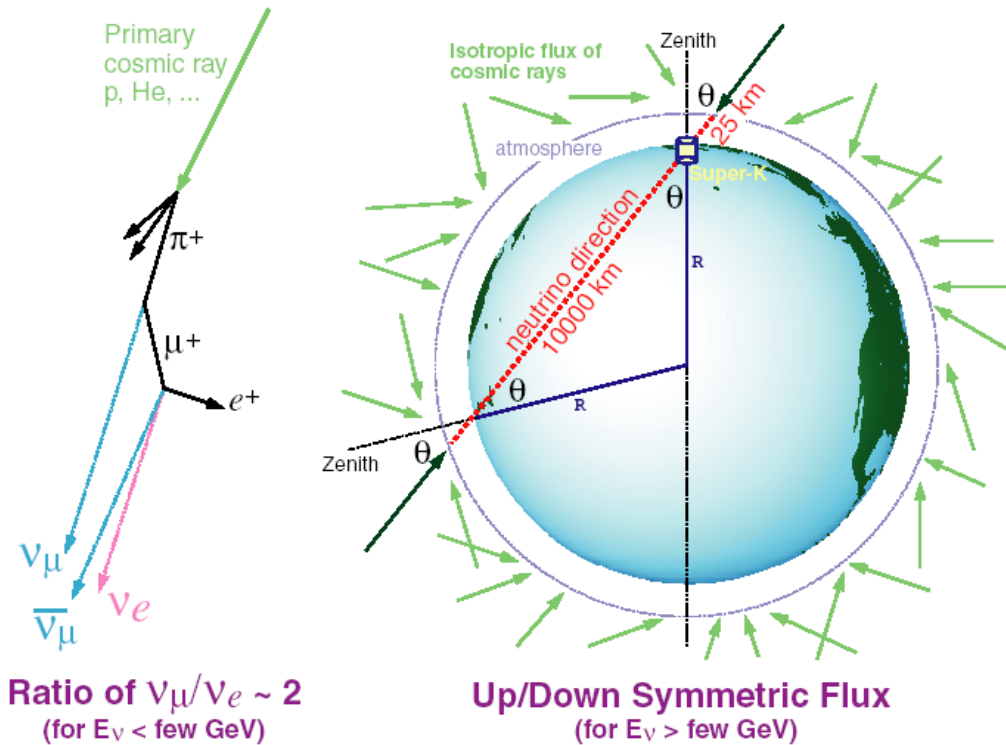


- Kamland - for anti-neutrinos

Sector 2-3 → first signal – atmospheric neutrinos

studies of background for proton decay

Observation of strong UP-DOWN asymmetry



Compare ν_μ to ν_e - take ratios to cancel out errors on absolute neutrino fluxes:

$$R = \frac{(\mu/e)_{data}}{(\mu/e)_{MC}} = 0.638 \pm 0.016 \pm 0.050$$

$$R_{highE} = \frac{(\mu/e)_{data}}{(\mu/e)_{MC}} = 0.658^{+0.030}_{-0.028} \pm 0.078$$

Too few muon neutrinos observed!

Evidence for Oscillation of Atmospheric Neutrinos

interpretation of the deficit of ν_μ after passing the Earth

$\nu_\mu \rightarrow \nu_x$ What is x?
Not e as we do not observe excess of ν_e

So we observe

$$\nu_\mu \rightarrow \nu_\tau$$

The neutrino interaction in SK is identified by observation of a charged lepton

$$\nu_\mu + N \rightarrow \mu^- + X$$

$$\nu_\tau + N \rightarrow \tau^- + X$$

But $m_\mu \ll m_\tau$, so if energy is too small to produce τ, ν_τ is **not observed**

This is why ν_μ are "missing"

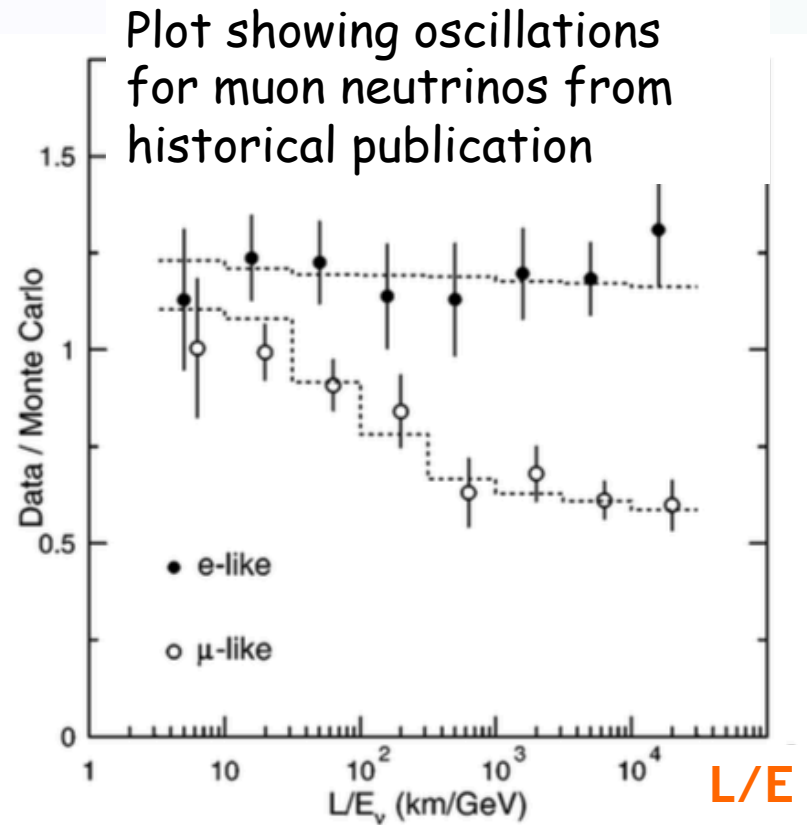
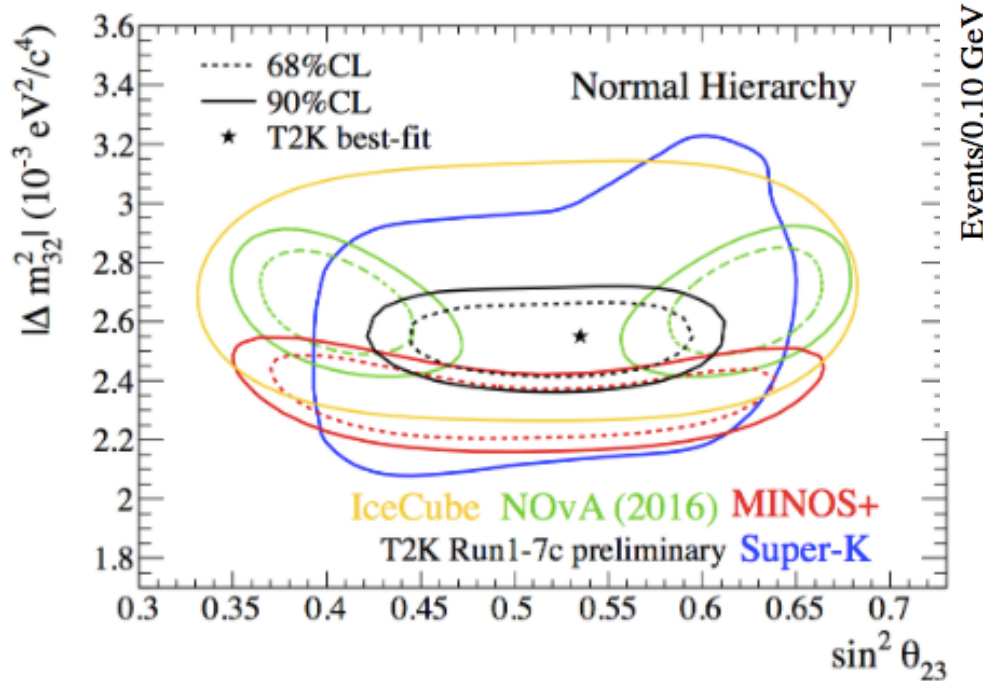


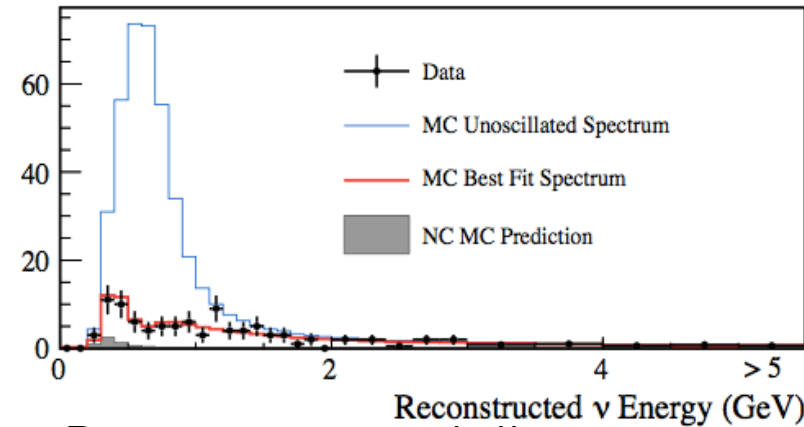
FIG. 4. The ratio of the number of FC data events to FC Monte Carlo events versus reconstructed L/E_ν . The points show the ratio of observed data to MC expectation in the absence of oscillations. The dashed lines show the expected shape for $\nu_\mu \leftrightarrow \nu_\tau$ at $\Delta m^2 = 2.2 \times 10^{-3} \text{ eV}^2$ and $\sin^2 2\theta = 1$. The slight L/E_ν dependence for e -like events is due to contamination (2–7%) of ν_μ CC interactions.

Present knowledge about 2-3 sector parameters:

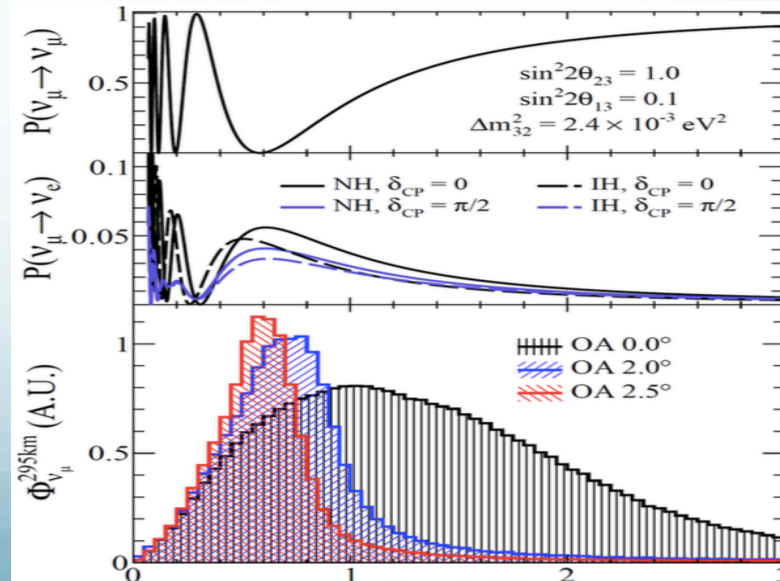
Summary of present knowledge:



Most precise measurement



Beam energy and distance tuned
To get maximal oscillation effect



	NH	IH
$\sin^2 \theta_{23}$	$0.532^{+0.044}_{-0.060}$	$0.534^{+0.041}_{-0.059}$
$ \Delta m^2_{32} $ ($/10^{-3} \text{eV}^2$)	$2.545^{+0.084}_{-0.082}$	$2.510^{+0.082}_{-0.083}$

1-3 ... and ways of measuring θ_{13}

- disappearance -> reactor experiments

$$\bar{\nu}_e \rightarrow \bar{\nu}_e$$

$$P_{\text{sur}} \approx 1 - \sin^2 2\theta_{13} \sin^2(1.267 \Delta m_{31}^2 L/E),$$

Energy ~ a few MeV
Distance ~ a few km

- appearance -> long-baseline experiments with ν_μ beam

$$\nu_\mu \rightarrow \nu_e$$

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2(1.27 \Delta m_{23}^2 L/E)$$

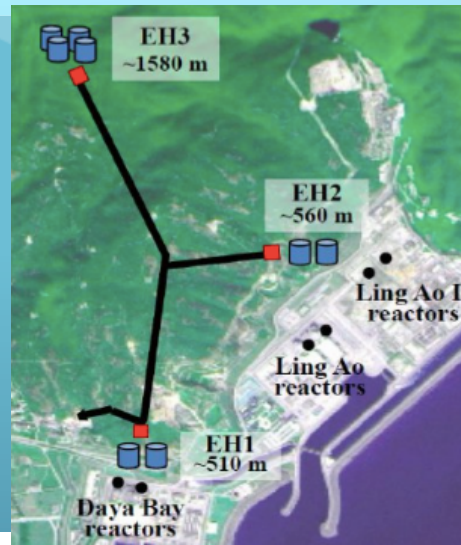
Second order terms depend on δ and mass hierarchy

Energy ~ a few GeV
Distance ~ a few hundred km

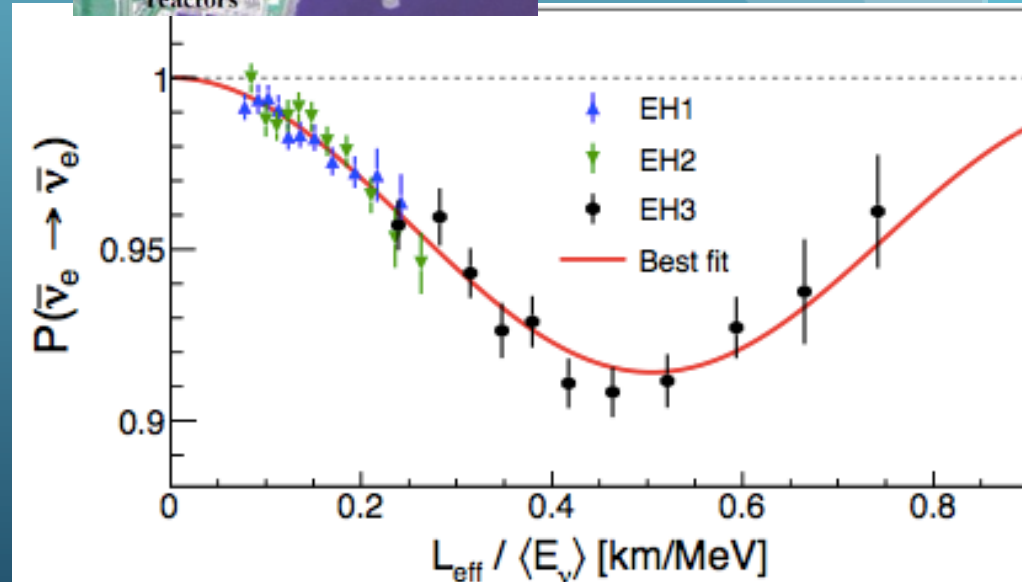
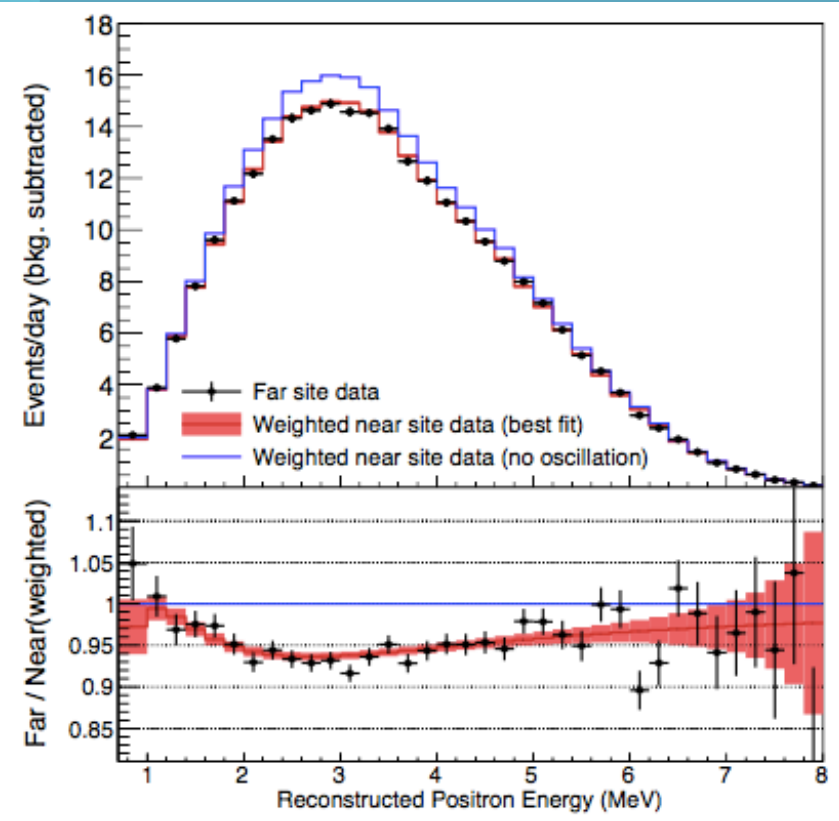
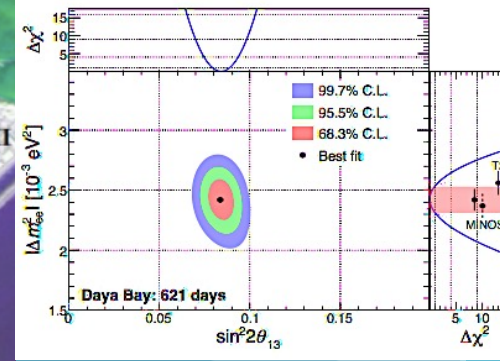
Leading terms!

Sector 1-3 reactor data

Daya Bay, RENO, Double CHOOZ
most precise measurements of θ_{13}



far and near
detectors



$$\sin^2 2\theta_{13} = 0.084 \pm 0.005$$

$$|\Delta m_{ee}^2| = (2.42 \pm 0.11) \times 10^{-3} \text{ eV}^2$$

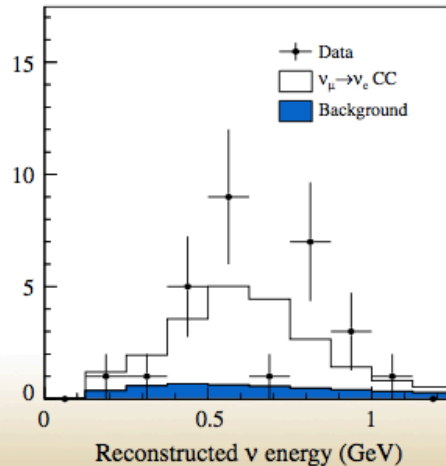
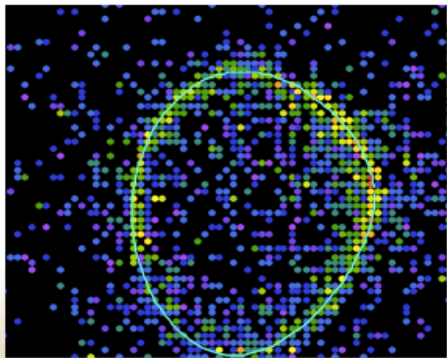
Observation in muon-neutrino beam

Selecting candidates for
electrons

e-like

- Electron-like ring
- $E_{\text{vis}} > 100 \text{ MeV}$
- $N_{\text{Michels}} == 0$
- $E_{\text{v, rec}} < 1250 \text{ MeV}$
- Not π^0 -like

28 ν_e
events



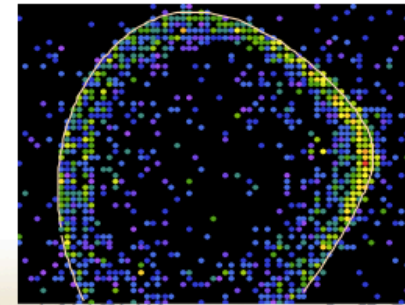
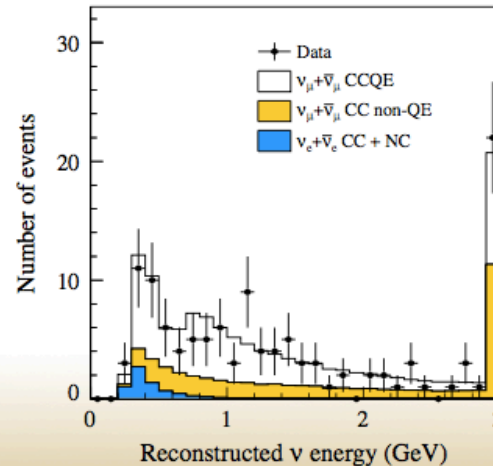
appearance

selecting candidates for
muons

μ -like

- Muon-like ring
- $P_\mu > 200 \text{ MeV}/c$
- $N_{\text{Michels}} < 2$

120 ν_μ
events



disappearance

What's next?

CPV

$$\delta \neq 0, \pi?$$

MH

$$m_3 \gtrless m_2?$$

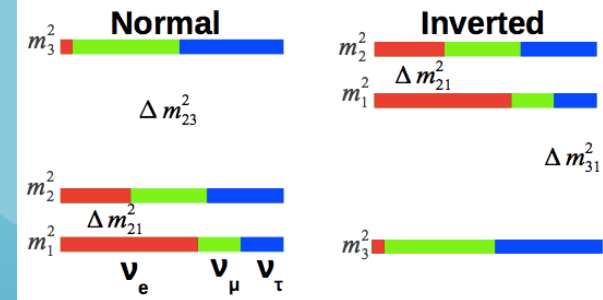
$$\theta_{23} \gtrless 45^\circ?$$

Differences in neutrino vs antineutrino oscillation probabilities

Changes the contribution from matter effects

(important for neutrinos travelling through dense matter e.g through Earth)

Additional source of degeneracies



Measurement strategies (for LBL):

- Looking for appearance

$$P(\nu_\mu \rightarrow \nu_e) \text{ vs. } P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)$$

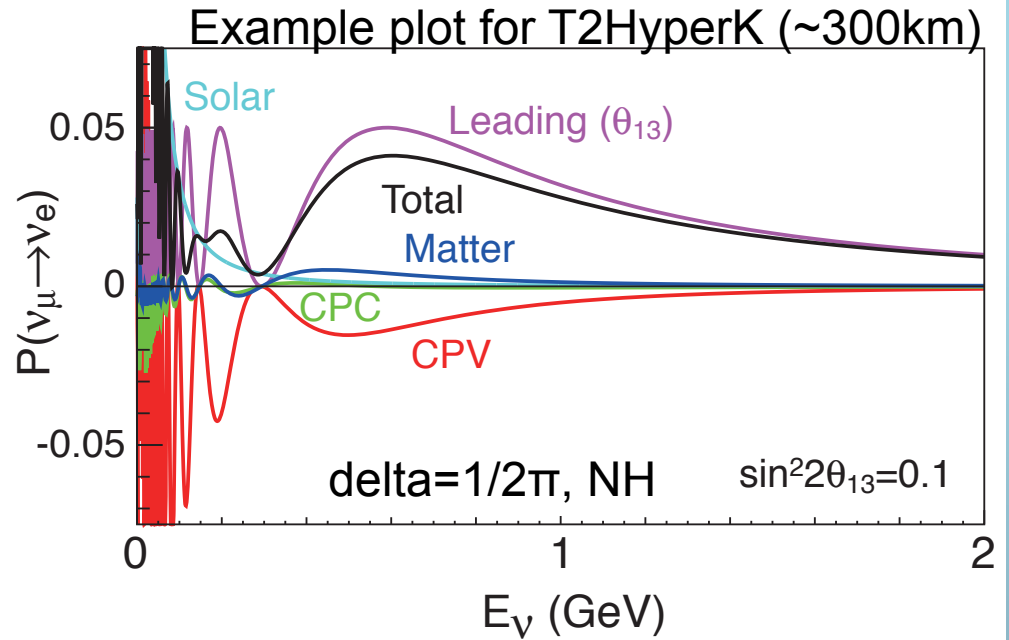
- The longer the baseline the better (matter effects!)
- Study more than one oscillation maximum to disentangle the effects

An unknown hierarchy usually leads to a reduced ability to observe CP violation

CPV and MH

In long baseline neutrino experiments

→ Many contributions, for precisions all need to be considered



$$\begin{aligned}
 P(\nu_\mu \rightarrow \nu_e) = & 4C_{13}^2 S_{13}^2 S_{23}^2 \cdot \sin^2 \Delta_{31} \quad \text{leading term} \quad \text{CP conserving} \\
 & + 8C_{13}^2 S_{12} S_{13} S_{23} (C_{12} C_{23} \cos \delta - S_{12} S_{13} S_{23}) \cdot \cos \Delta_{32} \cdot \sin \Delta_{31} \cdot \sin \Delta_{21} \\
 & - 8C_{13}^2 C_{12} C_{23} S_{12} S_{13} S_{23} \sin \delta \cdot \sin \Delta_{32} \cdot \sin \Delta_{31} \cdot \sin \Delta_{21} \quad \text{CP violating} \\
 & + 4S_{12}^2 C_{13}^2 (C_{12}^2 C_{23}^2 + S_{12}^2 S_{23}^2 S_{13}^2 - 2C_{12} C_{23} S_{12} S_{23} S_{13} \cos \delta) \cdot \sin^2 \Delta_{21} \\
 & - 8C_{13}^2 S_{13}^2 S_{23}^2 \cdot \frac{aL}{4E_\nu} (1 - 2S_{13}^2) \cdot \cos \Delta_{32} \cdot \sin \Delta_{31} \quad \text{solar term} \\
 & + 8C_{13}^2 S_{13}^2 S_{23}^2 \frac{a}{\Delta m_{31}^2} (1 - 2S_{13}^2) \cdot \sin^2 \Delta_{31}, \quad \text{matter effects}
 \end{aligned}$$

for $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$

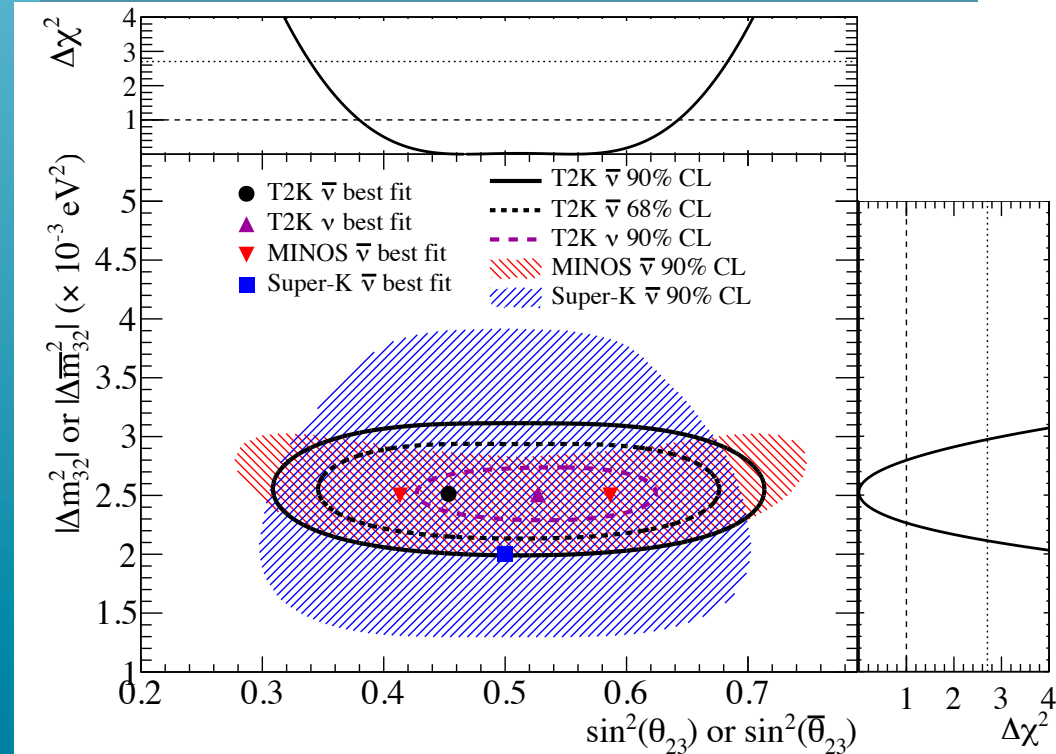
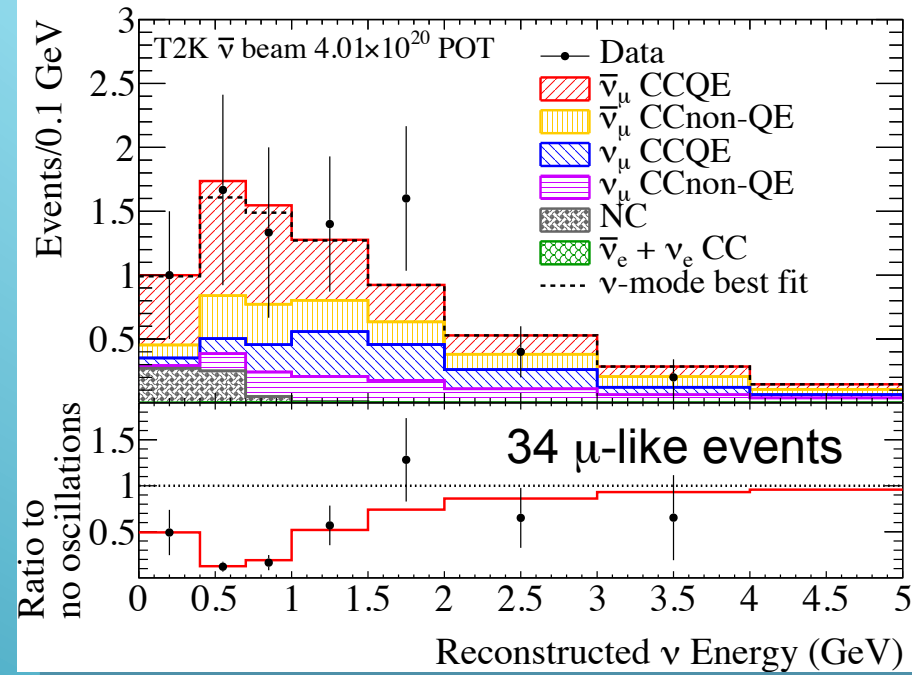
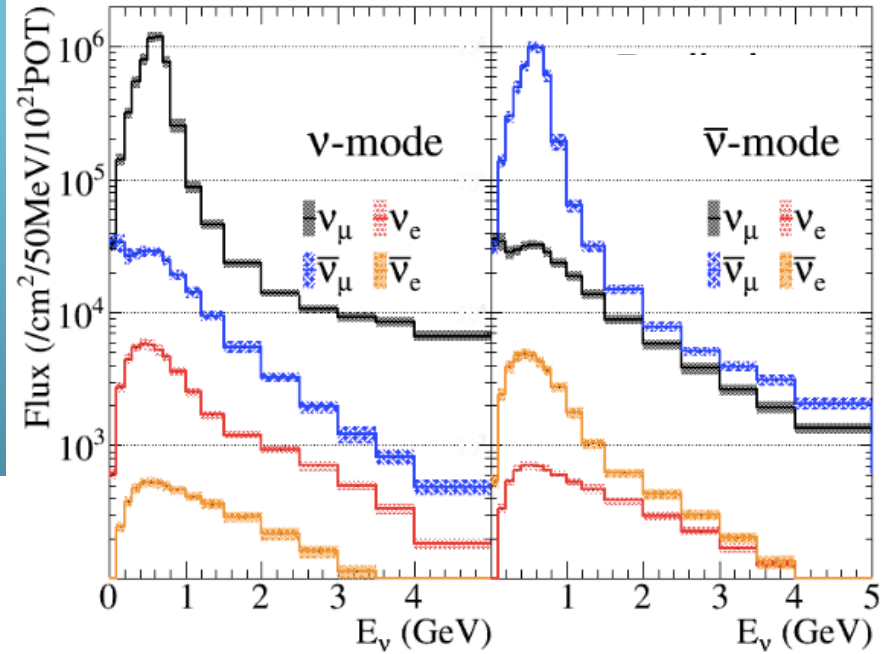
$\delta \rightarrow -\delta$

$a \rightarrow -a$

$C_{ij}, S_{ij}, \Delta_{ij}$
 $\cos \theta_{ij}, \sin \theta_{ij}, \Delta m_{ij}^2 L/4E_\nu$

$a \sim \rho * E_\nu$

SK flux prediction



Analysis

with anti- ν beam mode

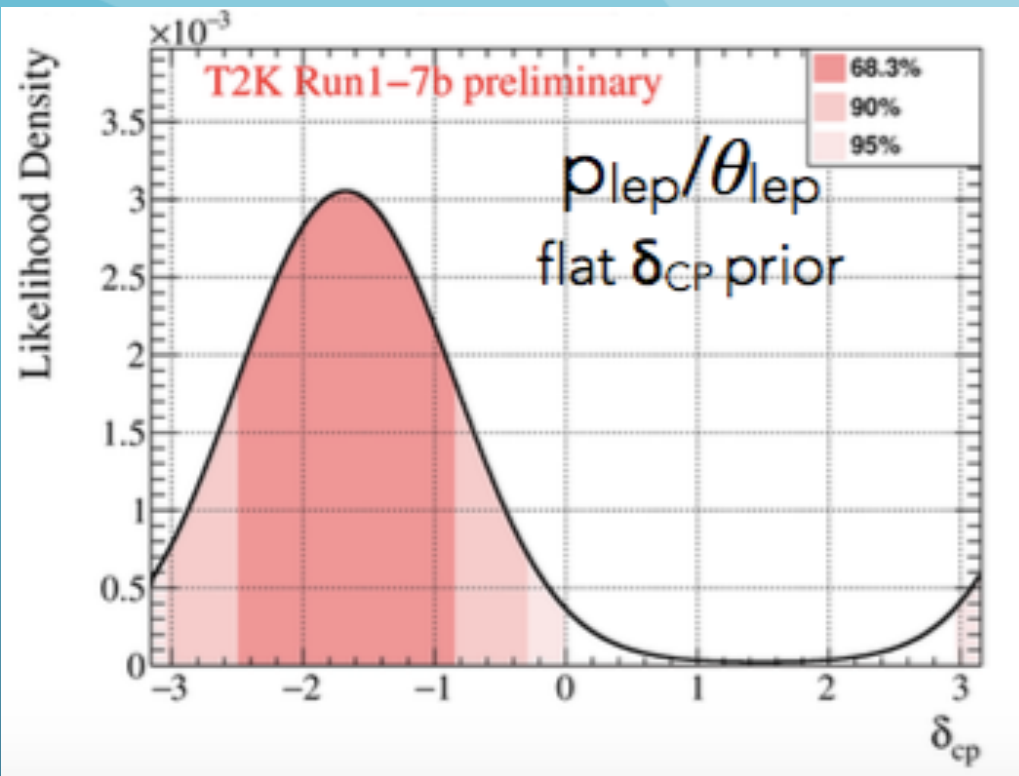
Disappearance:

parameters consistent with

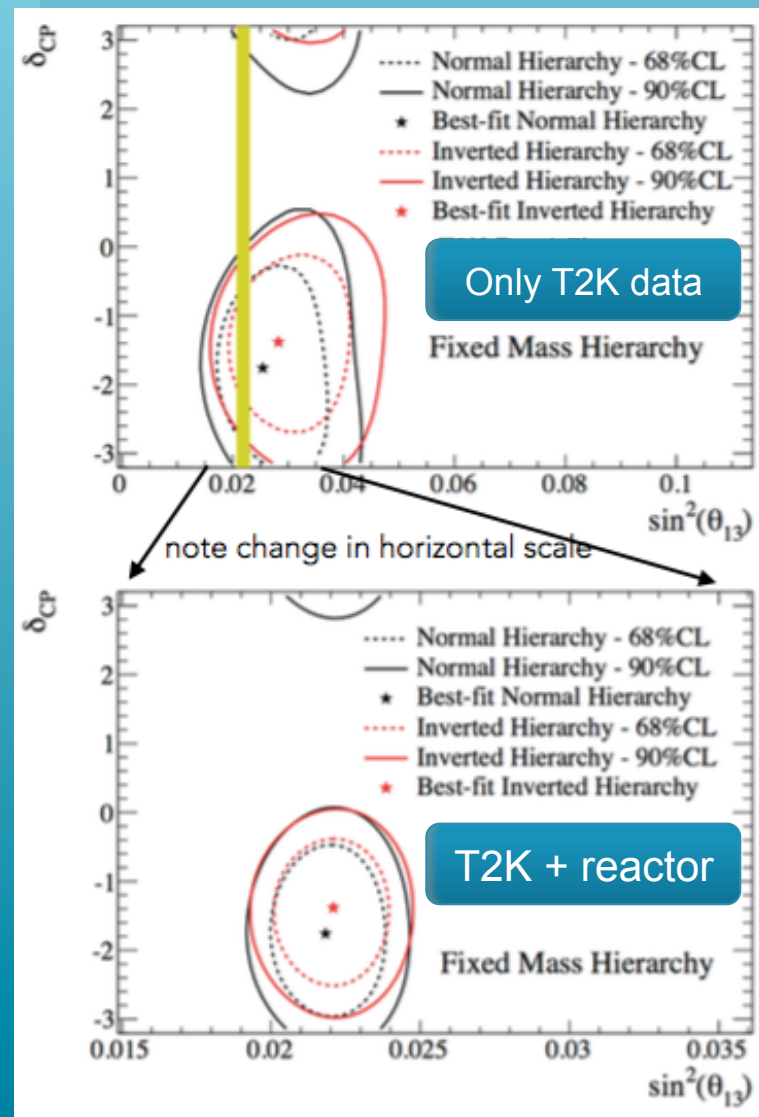
- maximal mixing
- parameters same as neutrino
- other experiments

**Appearance – 3 e-like events seen
more data needed**

New analysis ν and $\bar{\nu}$



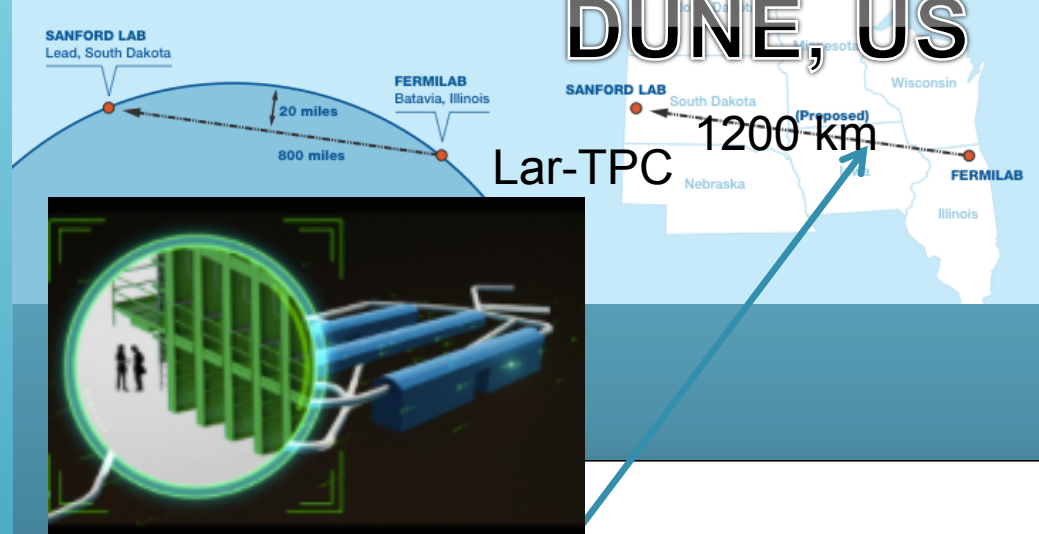
- First time from single experiment 90% CL range for CP violation phase
- Pointing to value corresponding to CP violation (0 and π outside 90% limit)
- More data needed \rightarrow T2K II
- Combined fits for all experiments will come soon (teoretica groups)



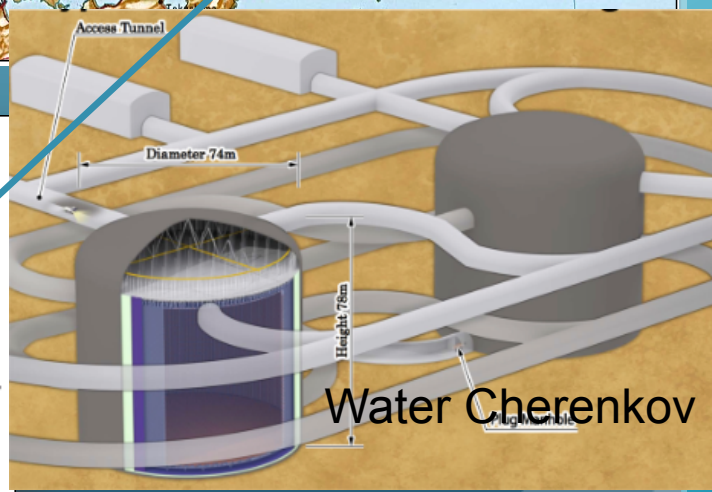
Long Baseline Future

Long term
ie. after/around 2025

Long-Baseline Neutrino Experiment



DUNE, US



And how about connection of neutrinos with work of Stanisław??



First try.... **Small oscillations**
of a collisionless quark plasma

Paper in Physics Letters B, 188 (1987) 129-132

Abstract:

The oscillations of a collisionless quark plasma are studied on the basis of the gauge covariant kinetic equations. The small oscillation approach provides the dispersion relations which coincide with those predicted by the finite-temperature QCD in one-loop approximation.

→ This connection is weak and not natural (only by playing words)

Next try.... Started with popularization:

“Tachyons: Particles faster than light”, Postępy Fizyki **32** (1981) 351

and more in this field.....



сообщения
 объединенного
 института
 ядерных
 исследований
 дубна

E2-83-299

St. Mrówczyński *

THE PHASE-SPACE OF TACHYONS

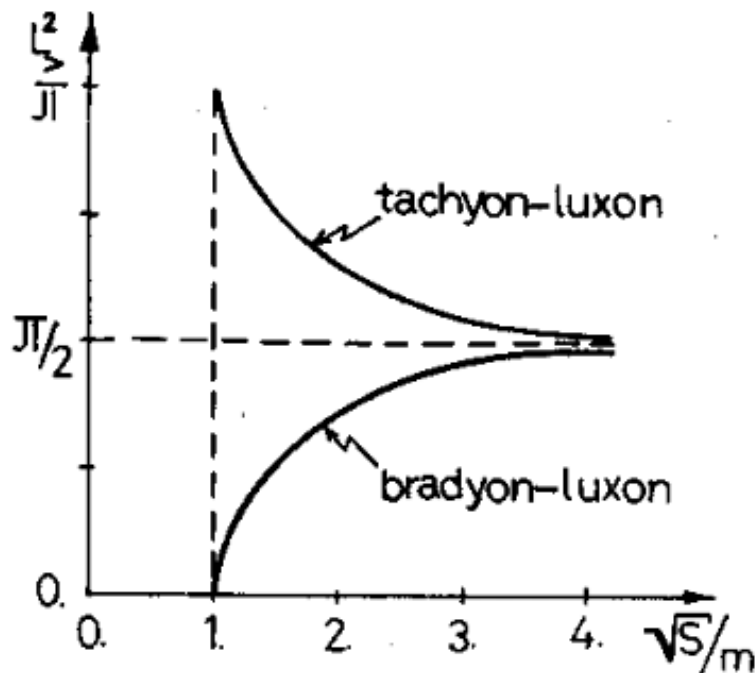
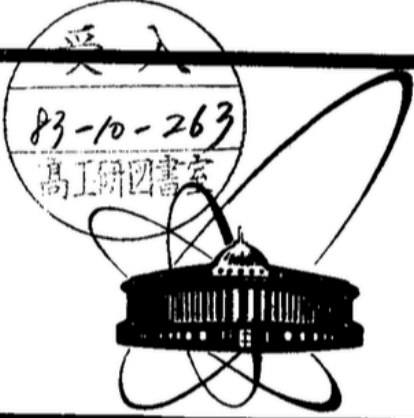


Fig.3. The volume of PS $L^2_{>}$ in CM for tachyon-luxon and bradyon-luxon systems

We conclude that it is easy to construct PS that is free of difficulties discussed in the previous section, however at the expense of Lorentz invariancy. We accept the view of the authors of /7/ that "the first major problem to be overcome in developing a quantum theory of tachyons is in reconciling the apparent conflict between Lorentz invariancy and need to have only positive energies capable of being observed to the theory".



сообщения
 объединенного
 института
 ядерных
 исследований
 дубна

E2-83-476

On the ideal gas of tachyons

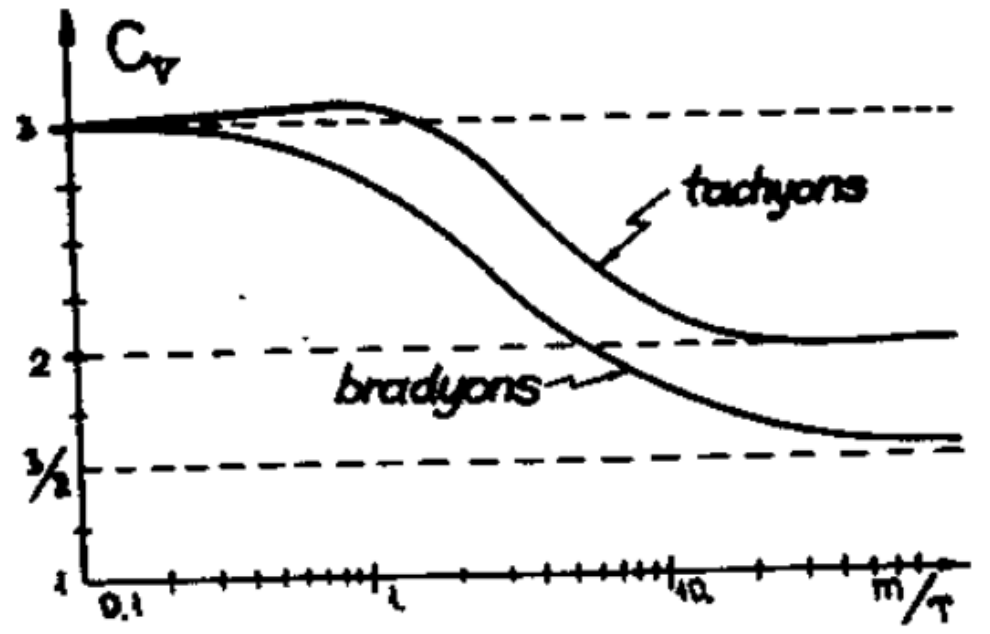
$$c_V = \begin{cases} = 3 + \mathcal{O}\left(\frac{m^2}{T^2} \ln(m/T)\right); & \text{for } T \gg m \\ = 2 + \mathcal{O}(T/m); & \text{for } T \ll m. \end{cases}$$

St.Mrówczyński*

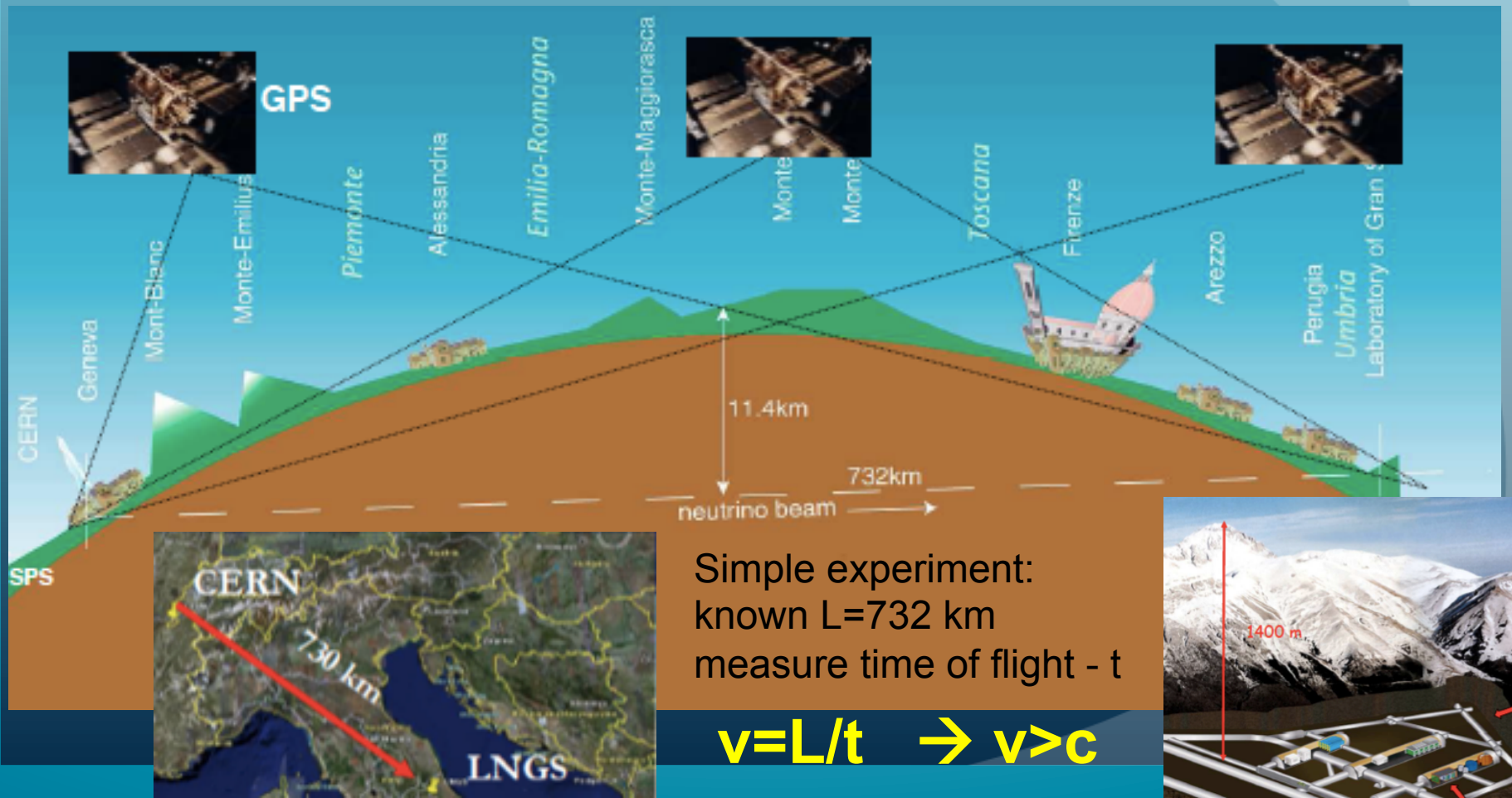
ON THE IDEAL GAS OF TACHYONS

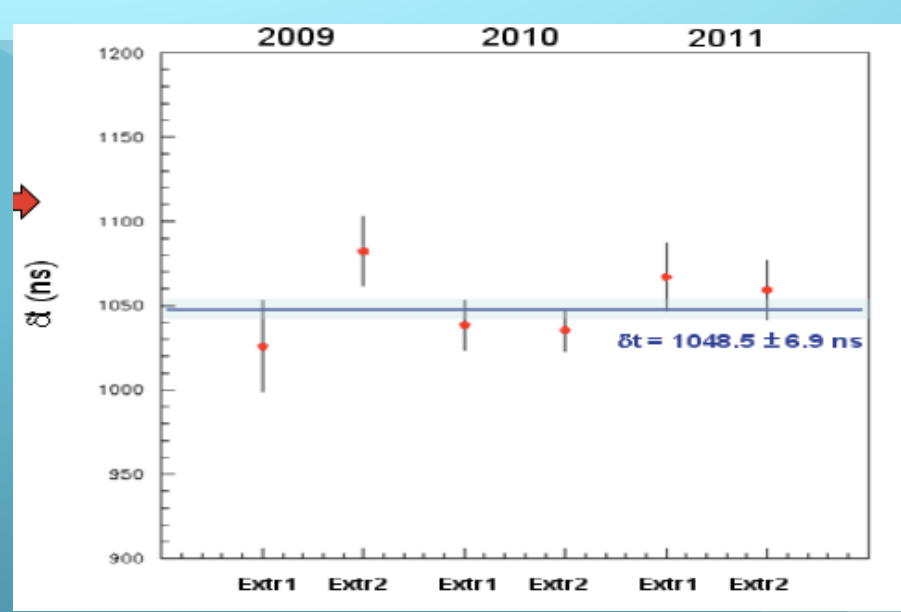
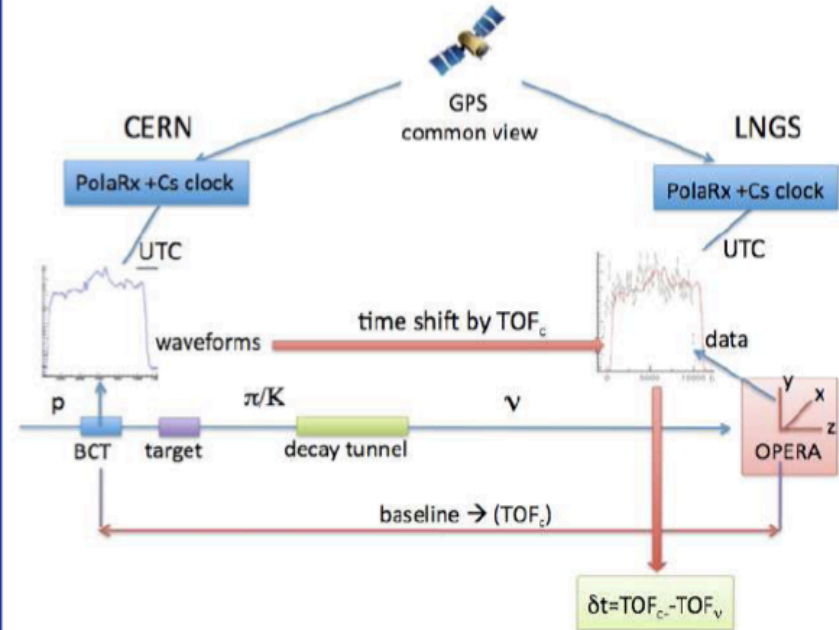
The specific heat of the idea gas of tachyons and the gas of bradyons as a function of m/T .

- 1982 - 1986 - asyistent, Joint Institut for Nuclear Research, Dubna, ZSRR



...and here we have connection:
in 2011 for some time v looked like
tachyons





$$\delta t = TOF_c - TOF_v = (1048.5 \pm 6.9(stat.))ns - 987.8ns$$

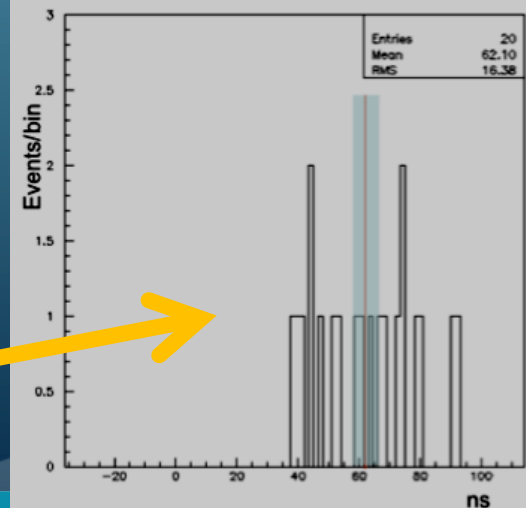
$$= (60.7 \pm 6.9(stat.))ns$$

$$(v - c) / c = \delta t / (TOF_c - \delta t) =$$

$$= (2.49 \pm 0.28(stat.) \pm 0.30(sys.)) \cdot 10^{-5} s$$

for test – run with narrow beam bunches,

measure $\delta = ToF_v - ToF_c$



Team admits to possible errors in faster-than-light finding.

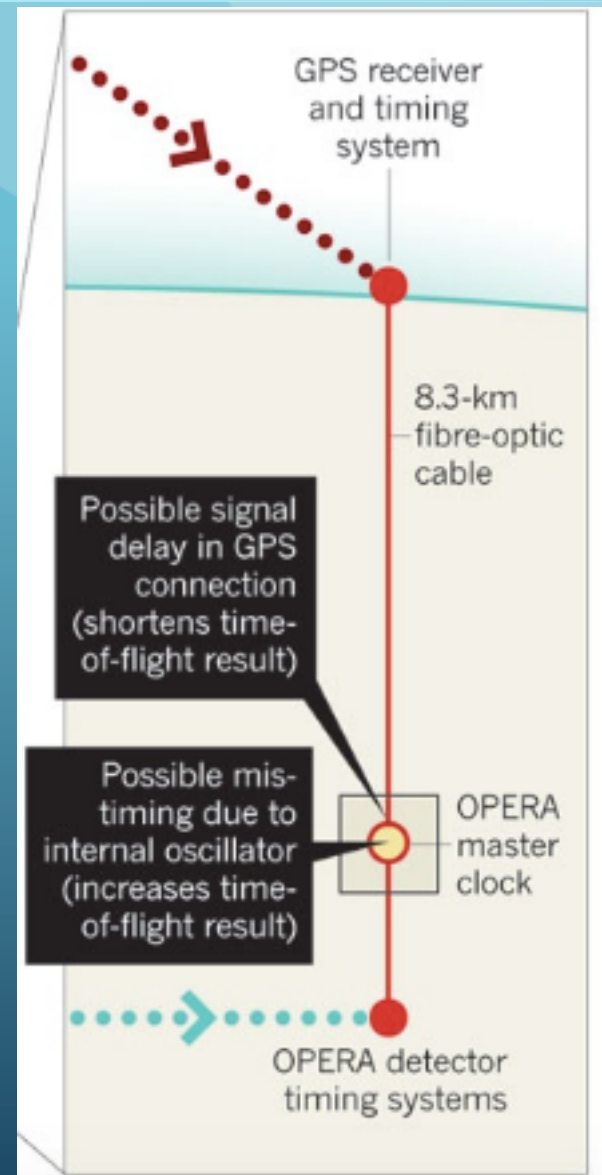
NATURE, 27 February 2012

Fault was found in the delay in GPS connection,
after that v speed $< c$ again

so present status:

no tachyons found experimentally

(again..)



With best wishes for Stanisław !!!!!