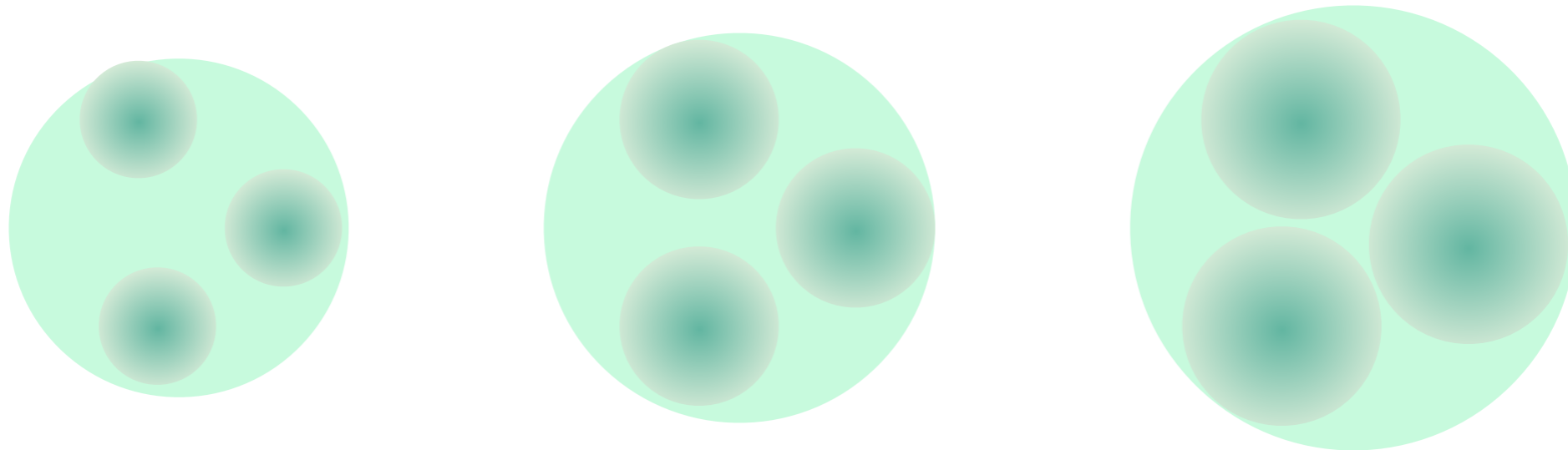


HOT SPOTS AND THE HOLLOWNESS OF PROTON-PROTON INTERACTIONS



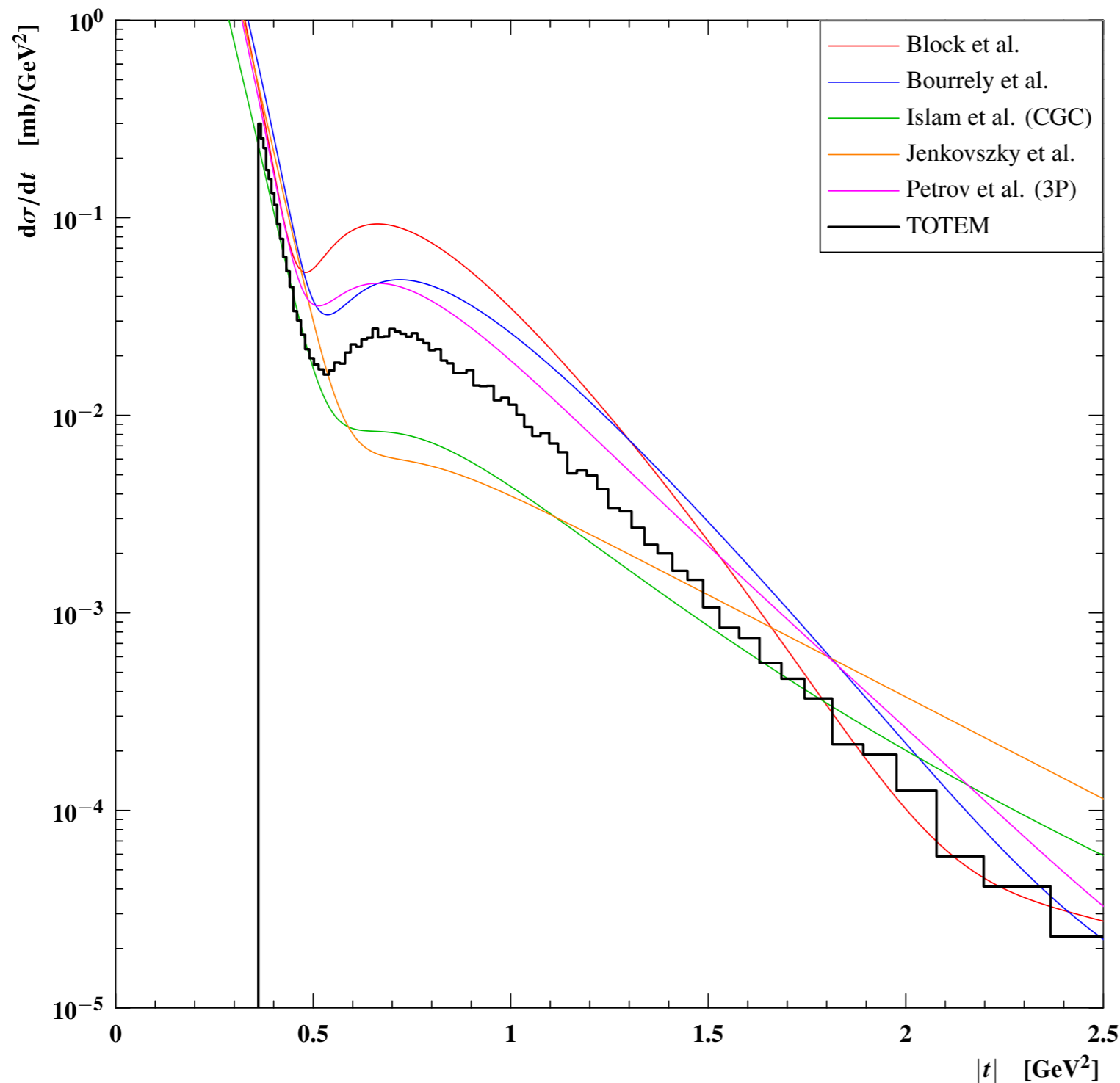
Javier L. Albacete

+ Alba Soto-Ontoso

based on arXiv:1605.09176



- TOTEM data on elastic differential cross section in *pp* collisions at 7 TeV



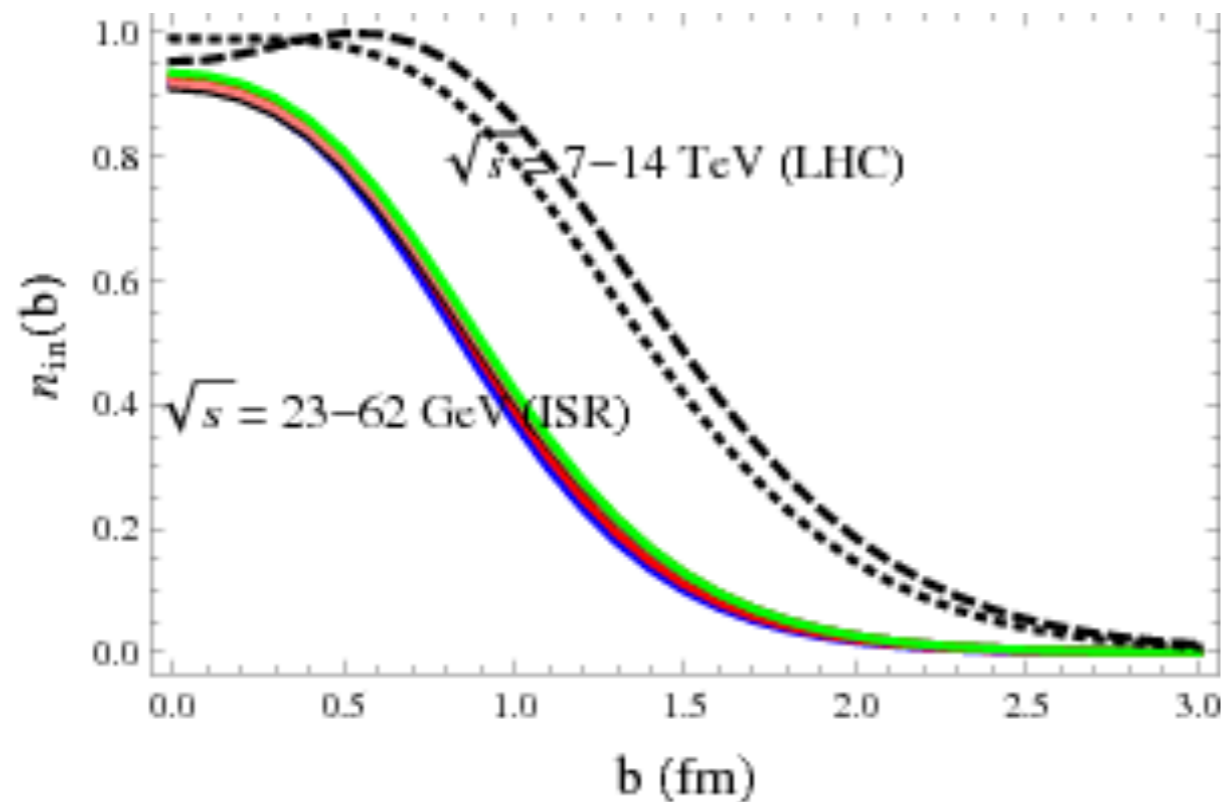
$$\frac{d\sigma_{\text{el}}}{dt} = \frac{1}{4\pi} |T_{\text{el}}(s, t)|^2$$

Hollowness effect

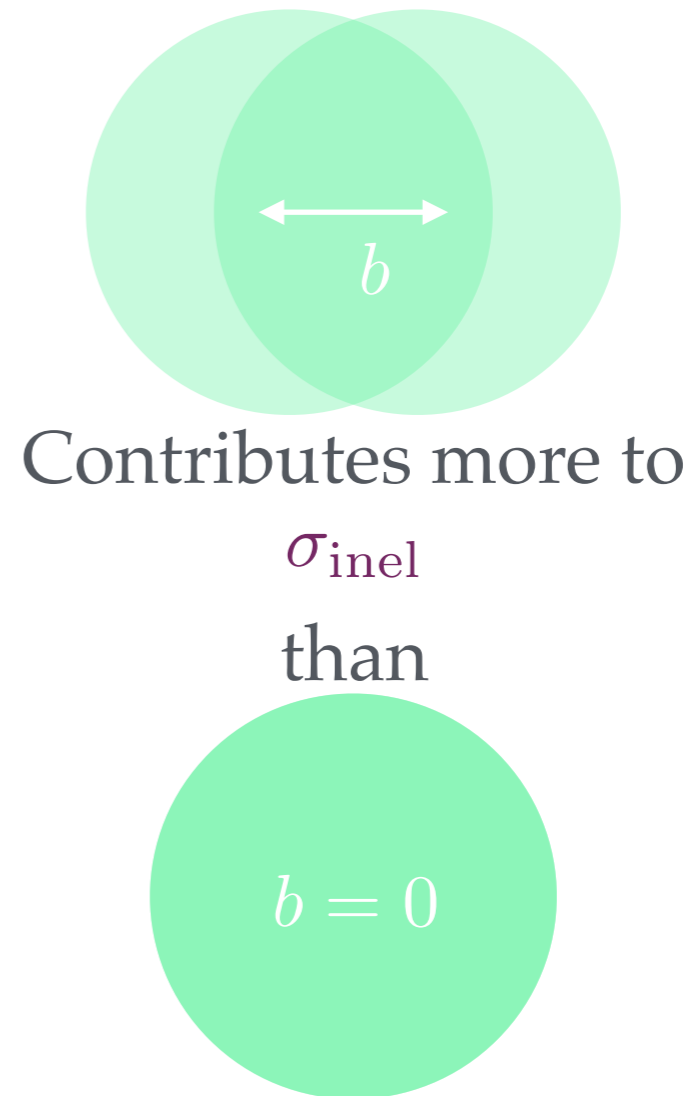
- The *hollowness/grayness* effect in pp interactions @LHC

$$G_{\text{in}}(s, \vec{b}) = 2\text{Im}\tilde{T}_{\text{el}}(s, \vec{b}) - |\tilde{T}_{\text{el}}(s, \vec{b})|^2$$

$$G_{\text{in}} = d^2\sigma_{\text{inel}}/d^2b$$



[Ruiz-Arriola et Al. '15]



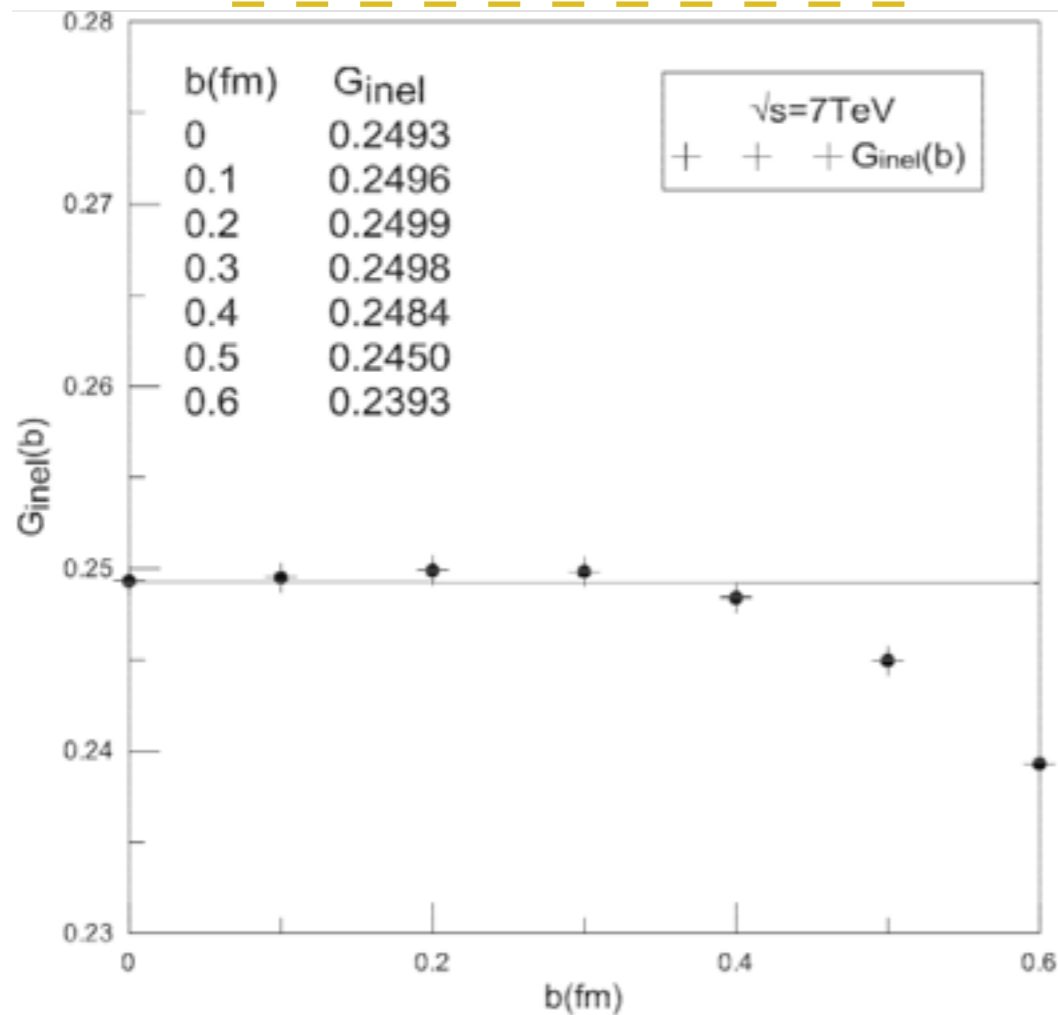
- Not observed @ISR and no dynamical explanation @market

Hollowness effect

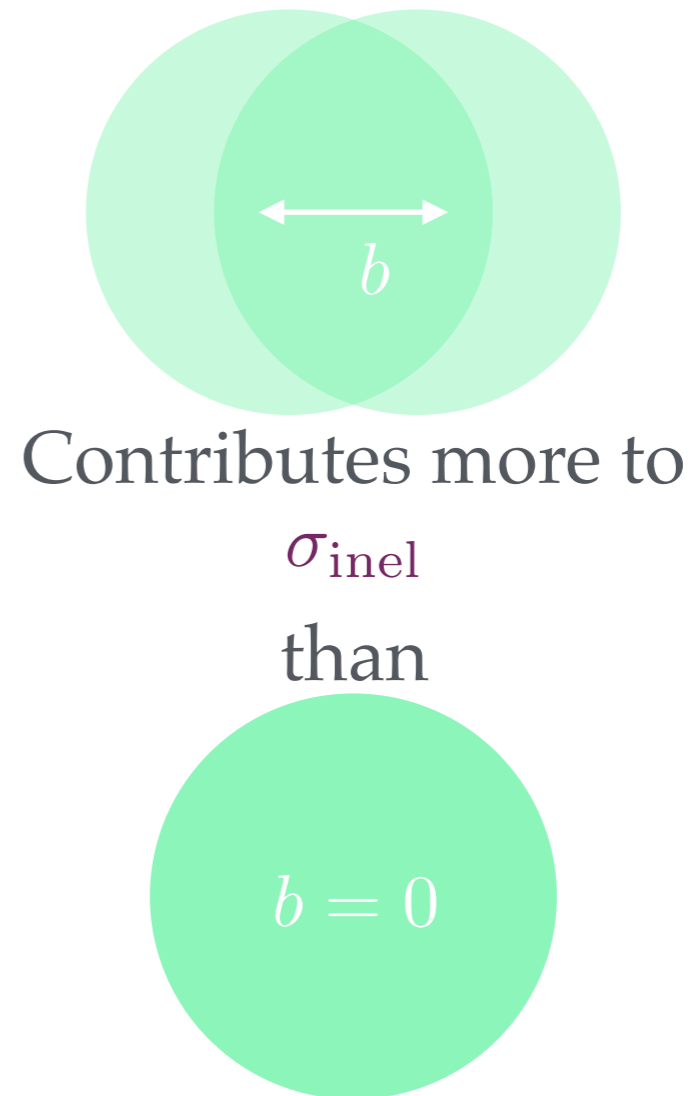
- The *hollowness/grayness* effect in pp interactions @LHC

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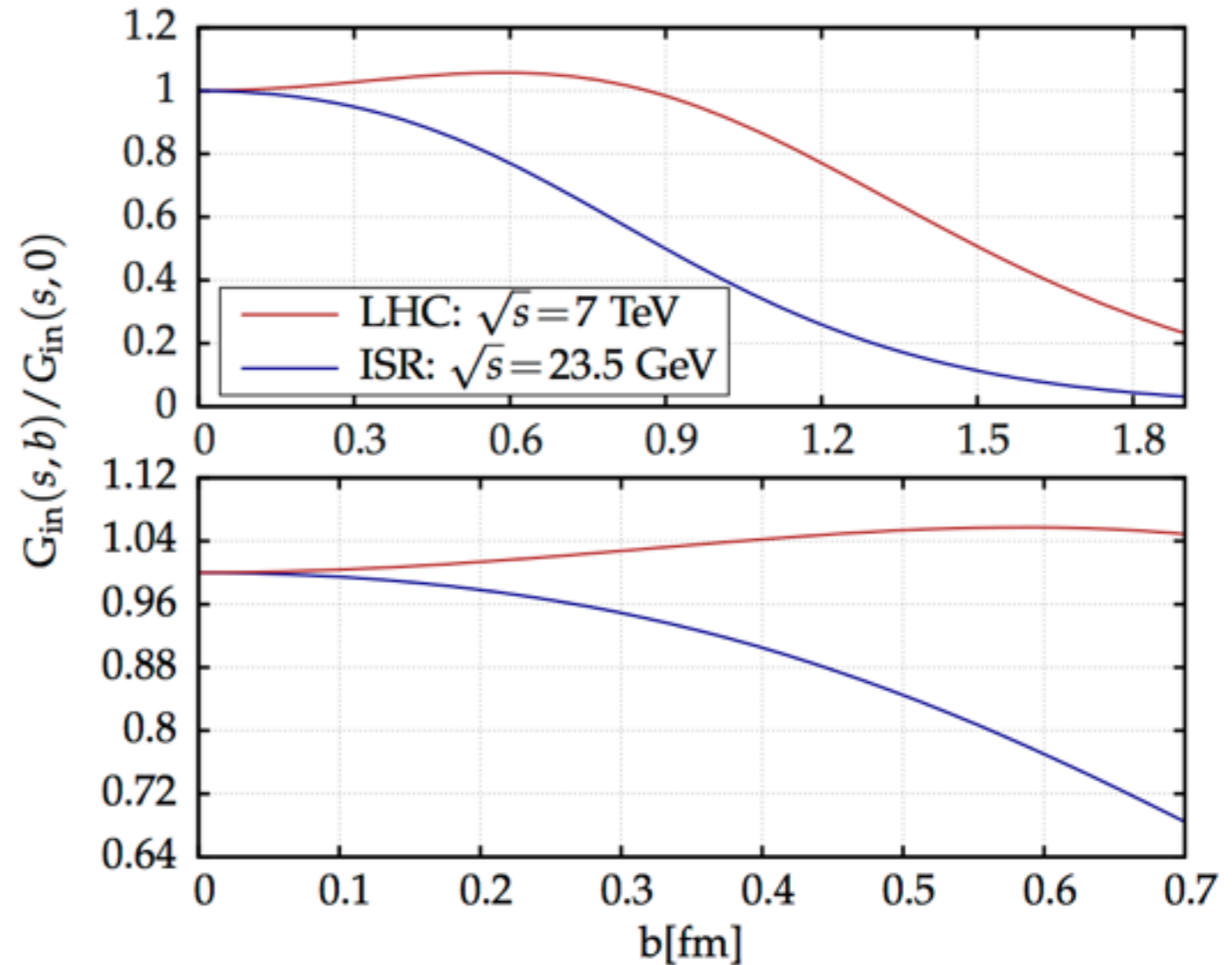
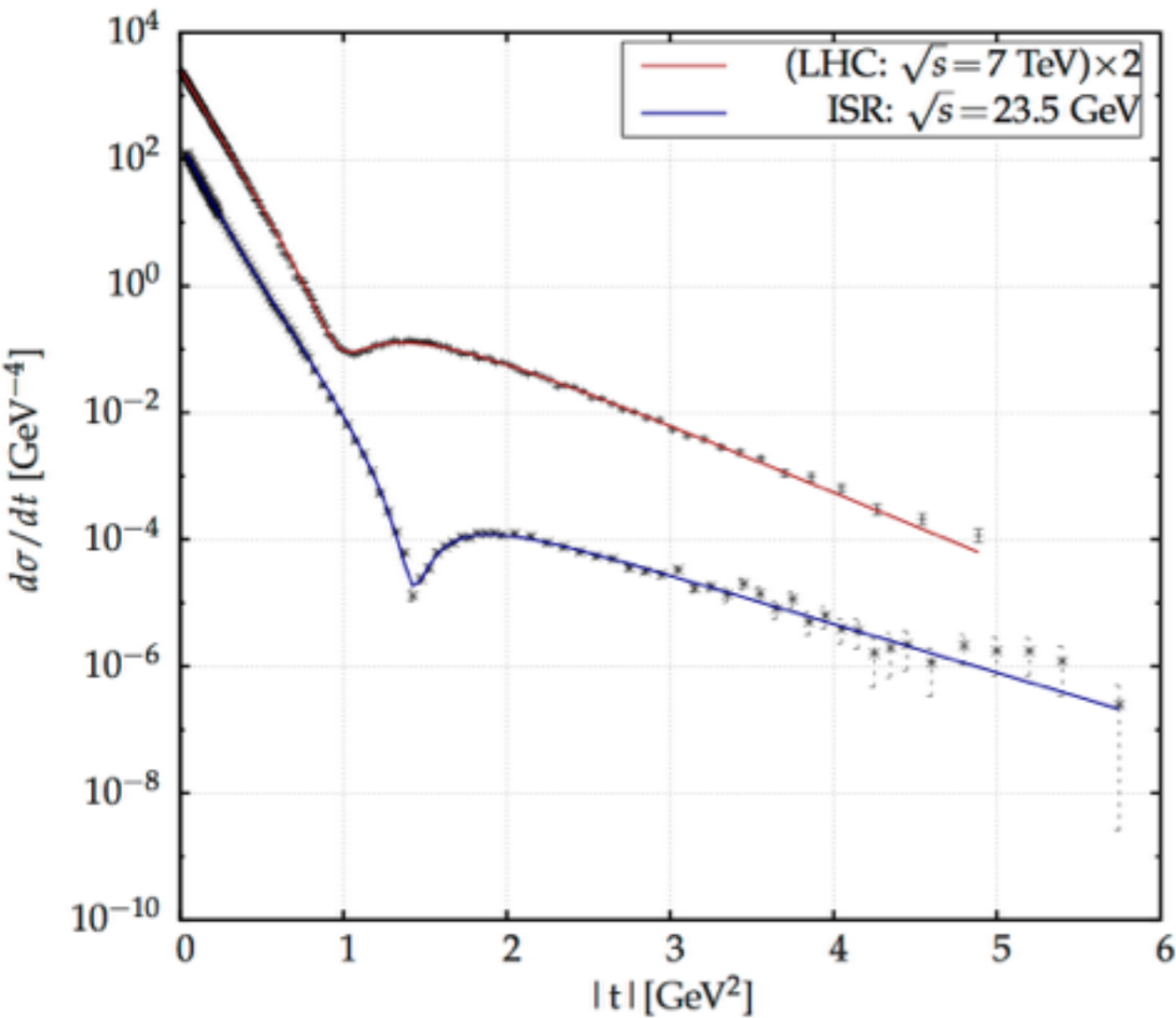
[A.Alkin et Al. '14]



- Not observed @ISR and no dynamical explanation @market

Hollowness effect

- We have performed an independent analysis



- \ddagger The inelasticity density of the collision does not reach a maximum at $b=0$!!

Problem to solve

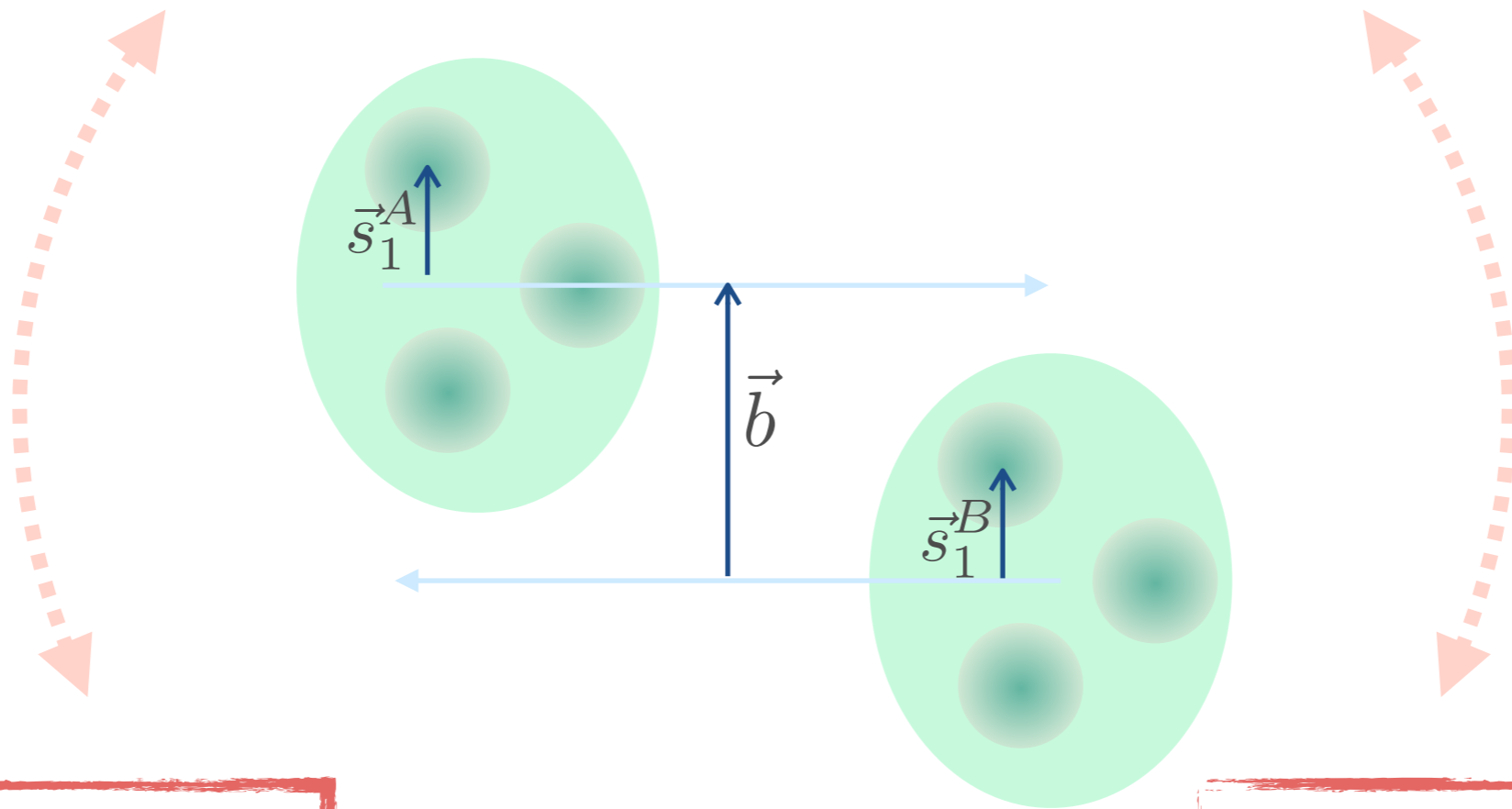
- The inelasticity density exhibits a maximum at $b > 0$: *hollowness* effect
 - Peripheral collisions are more *destructive*.
 - Pure convolution models are precluded.
 - It disappears at ISR energies.
- Constrain the *transverse structure* of the proton
 - Implications in harmonic flow coefficients.

2. Ingredients

The model

- To construct the elastic scattering amplitude in pp collisions

Gluonic hot-spots as effective d.o.f



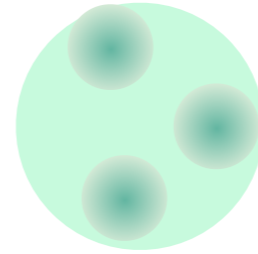
Glauber multiple scattering expansion

Spatial correlations in transverse space

Hot spots

- **Assumption:** the gluon content of the proton concentrated in small domains

$$R_{hs} \ll R_p$$



- **Open debate:** they may be radiatively generated from valence quarks in DGLAP or BFKL-like cascades (growth with energy)

Hot spot  Fock space of the valence partons

/ instantons / combination of perturbative and non perturbative physics

[Kopeliovich et Al. '99, Braun et Al. '93, Schafer et Al. '98, Kovner '02, Shuryak'04, Schenke et Al.'15...]

- ✓ Smallness of the correlation length of the gluon field in lattice QCD.

[DiGiacomo et Al. '92]

- ✓ Phenomenological tool [Kopeliovich et Al. '07]

Glauber model

- pp interactions as a collision of two systems A and B, each one composed of 3 hot spots. [Similar to A. Bialas et al '70s]. Model works for $N_{hs} \cong 3$

$$\tilde{T}_{el}(\vec{b}) = \int \prod_{k,l} d^2 s_k^A d^2 s_l^B D_A(\{\vec{s}_k^A\}) D_B(\{\vec{s}_l^B\}) \left(1 - \prod_i \prod_j \left[1 - \Theta_{ij}(\vec{b} + \vec{s}_i^A - \vec{s}_j^B) \right] \right)$$

- \vec{b} : impact parameter of the collision.
- \vec{s}_i : transverse positions the hot spots.
- $D(\vec{s}_1, \vec{s}_2, \vec{s}_3)$: density distribution of hot spots.
- $\Theta_{ij}(\vec{b} + \vec{s}_i^A - \vec{s}_j^B)$: elastic amplitude of the i -th and j -th hot spot interaction.

$$\Theta(s_{ij}) = \exp(-s_{ij}^2/2R_{hs}^2) (1 - i\rho_{hs})$$

Spatial correlations

- The general structure that we consider for $D(\vec{s}_1, \vec{s}_2, \vec{s}_3)$

$$D(\{\vec{s}_i\}) = C \left(\prod_{i=1}^3 d(\vec{s}_i; R) \right) \times f(\vec{s}_1, \vec{s}_2, \vec{s}_3)$$

- C : normalization constant.

- $d(\vec{s}_i; R)$: uncorrelated probability distribution for a single hot spot.

$$d(\vec{s}_i; R) = \exp(-s_i^2/R^2)$$

- $f(\vec{s}_1, \vec{s}_2, \vec{s}_3)$: correlation structure.

$$f(\vec{s}_1, \vec{s}_2, \vec{s}_3) = \delta^{(2)}(\vec{s}_1 + \vec{s}_2 + \vec{s}_3) \prod_{\substack{i < j \\ i, j = 1}}^3 \left(1 - e^{-\mu |\vec{s}_i - \vec{s}_j|^2 / R^2} \right)$$

Spatial correlations

$$f(\vec{s}_1, \vec{s}_2, \vec{s}_3) = \delta^{(2)}(\vec{s}_1 + \vec{s}_2 + \vec{s}_3) \prod_{\substack{i < j \\ i, j = 1}}^3 \left(1 - e^{-\mu |\vec{s}_i - \vec{s}_j|^2 / R^2} \right)$$

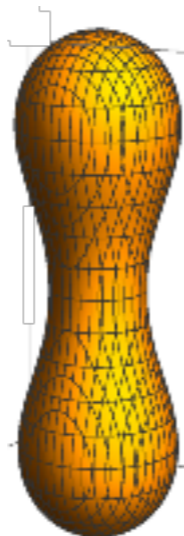
- $\delta^{(2)}(\vec{s}_1 + \vec{s}_2 + \vec{s}_3)$: fixes the center of mass of the hot spots system.

- $\prod_{\substack{i < j \\ i, j = 1}}^3 \left(1 - e^{-\mu |\vec{s}_i - \vec{s}_j|^2 / R^2} \right)$: repulsive short-range correlations controlled by

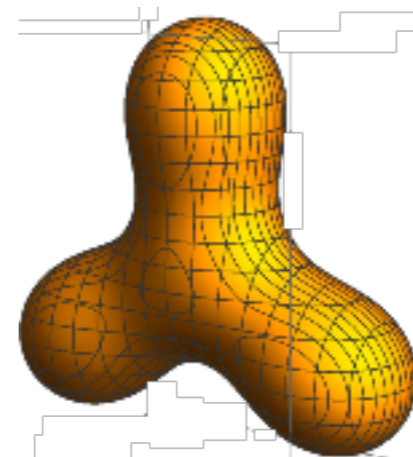
$$r_c^2 \propto R^2 / \mu$$

- Similar correlation structure than 3D models (when projected)

Quark-Diquark:



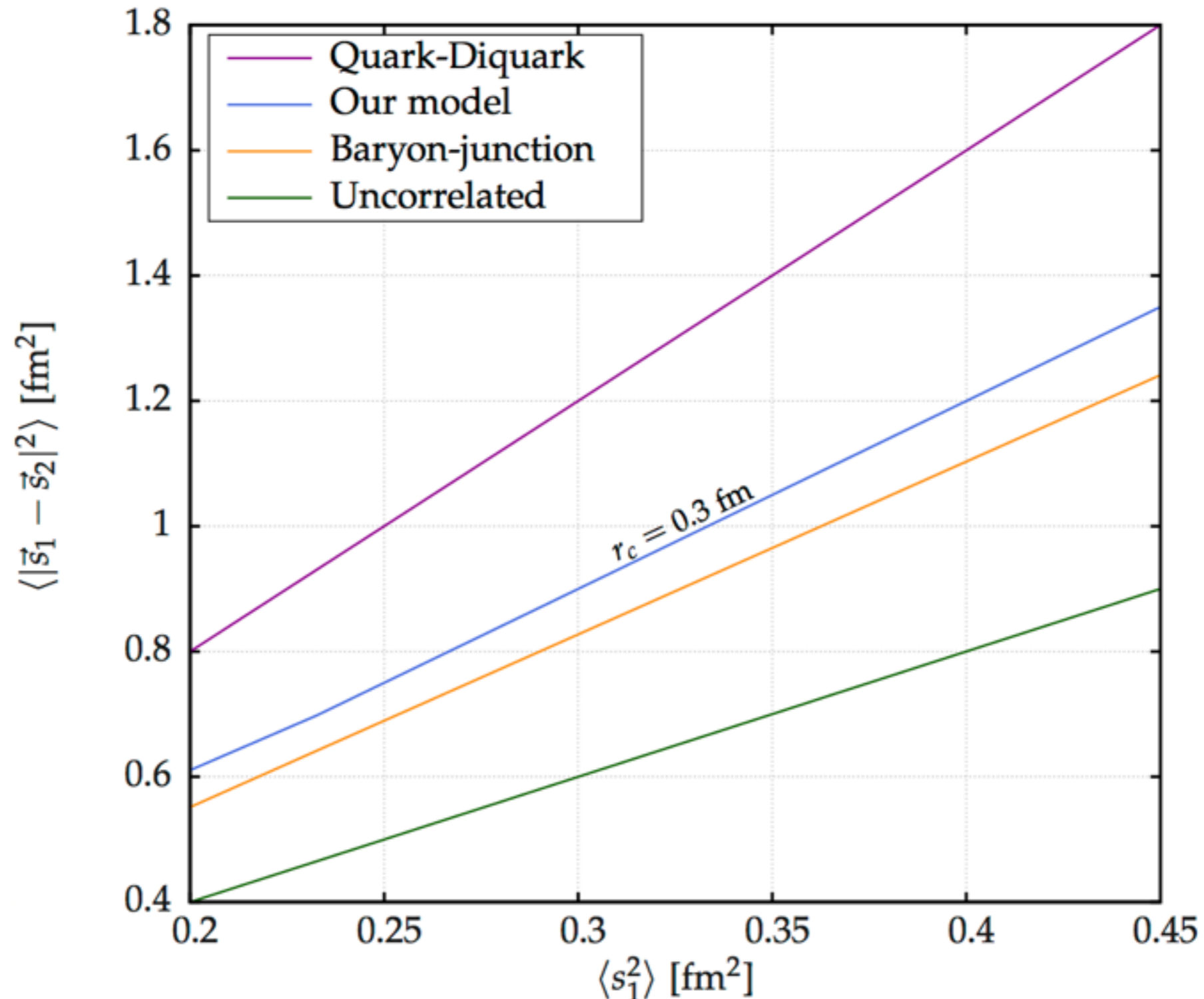
Baryon junction:



[Kubiczek et Al. '15]

Spatial correlations

- Averaged hot spot-hot spot transverse distance for different $D(\vec{s}_1, \vec{s}_2, \vec{s}_3)$



Conventions

- $\frac{d\sigma_{\text{el}}}{dt} = \frac{1}{4\pi} |T_{\text{el}}(s, t)|^2$
- $T_{\text{el}}(s, t) = \int d^2b \tilde{T}_{\text{el}}(s, \vec{b}) e^{-i\vec{q}\cdot\vec{b}}$
- $\sigma_{\text{el}} = \int d^2b |\tilde{T}_{\text{el}}(s, \vec{b})|^2$
- $\sigma_{\text{tot}} = 2\text{Im}T_{\text{el}}(s, 0) = 2 \int d^2b \text{Im}\tilde{T}_{\text{el}}(s, \vec{b})$
- $\sigma_{\text{in}} = \sigma_{\text{tot}} - \sigma_{\text{el}} = \int d^2b 2\text{Im}\tilde{T}_{\text{el}}(s, \vec{b}) - |\tilde{T}_{\text{el}}(s, \vec{b})|^2$
- $\rho = \frac{\text{Re}T_{\text{el}}(s, 0)}{\text{Im}T_{\text{el}}(s, 0)}$
- $G_{\text{in}}(s, \vec{b}) = 2\text{Im}\tilde{T}_{\text{el}}(s, \vec{b}) - |\tilde{T}_{\text{el}}(s, \vec{b})|^2$

3. Results

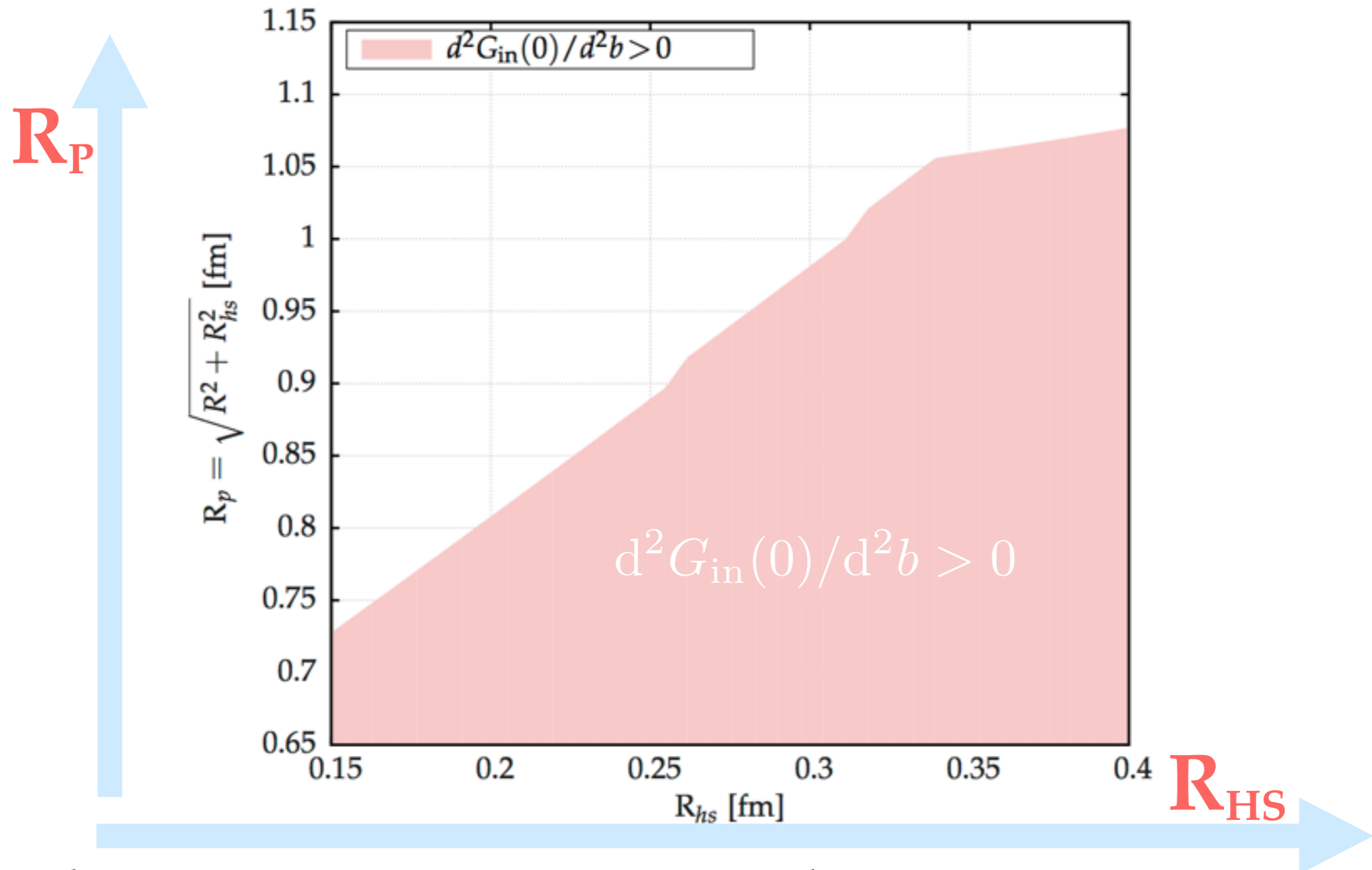
- We scan the **parameter space** with the conditions

▶ Maximum of the elastic amplitude: $\frac{d^2 \tilde{T}(s, 0)}{d^2 b} < 0$

▶ Maximum of the inelastic density: $\frac{d^2 G_{\text{in}}(s, 0)}{d^2 b} > 0$

R_p vs R_{hs}

- For $r_c = 0.5$ fm and $\rho_{hs} = 0.1$,



- Up to this point, purely geometric approach. No energy dependence.

- To be compatible with the phenomenology

▶ Maximum of the elastic amplitude: $\frac{d^2\tilde{T}(s, 0)}{d^2b} < 0$

▶ Maximum of the inelastic density: $\frac{d^2G_{\text{in}}(s, 0)}{d^2b} > 0 \Big|_{\text{LHC}}$

$$\frac{d^2G_{\text{in}}(s, 0)}{d^2b} < 0 \Big|_{\text{ISR}}$$

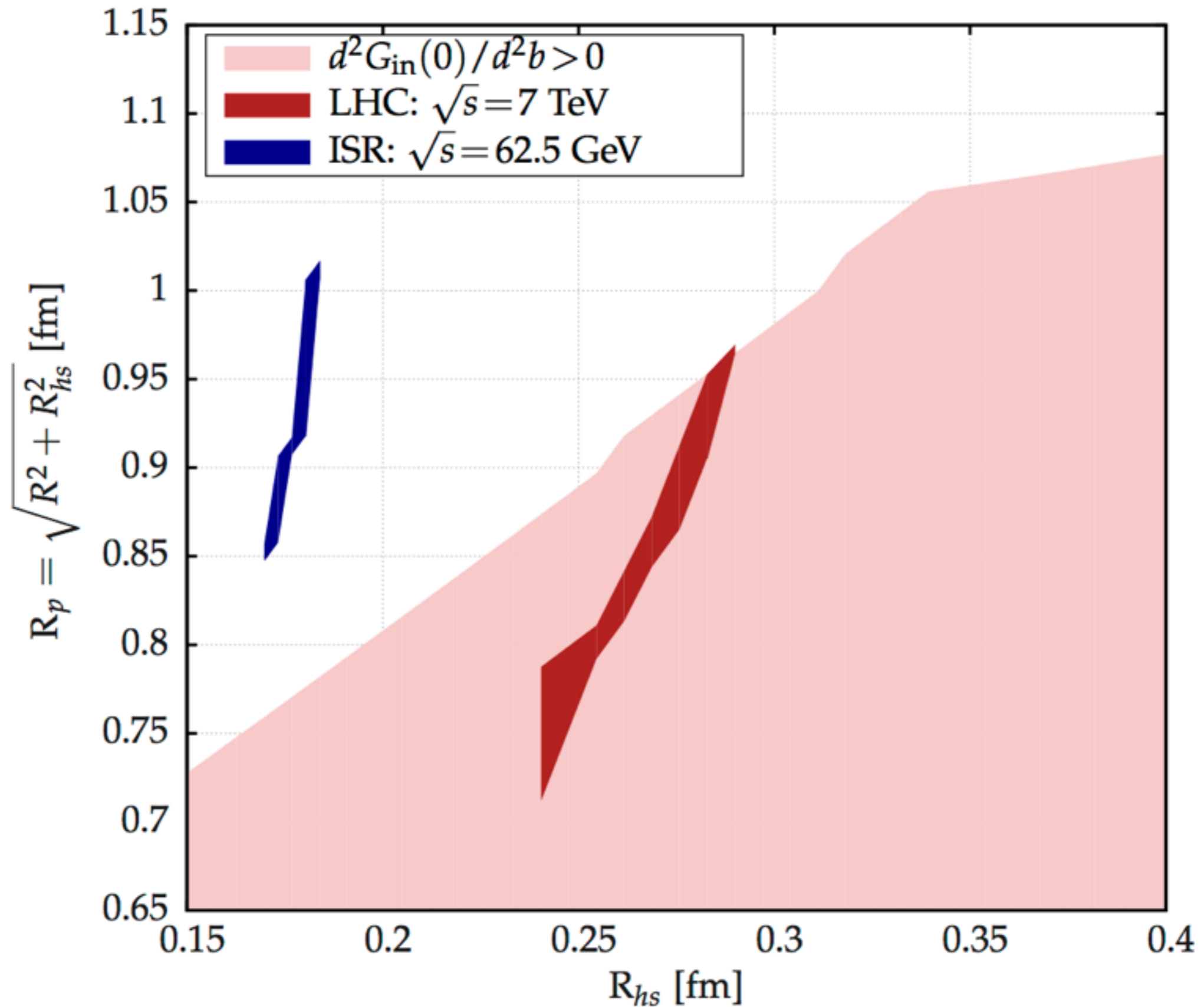
▶ Phenomenological constraints:

- **LHC, 7 TeV:** $\sigma_{\text{tot}} = 9,83 \pm 0.28 \text{ [fm}^2\text{]} \quad \rho = 0.14^{+0.01}_{-0.08}$

- **ISR, 62.5 GeV:** $\sigma_{\text{tot}} = 4,332 \pm 0.023 \text{ [fm}^2\text{]} \quad \rho = 0.095 \pm 0.018$

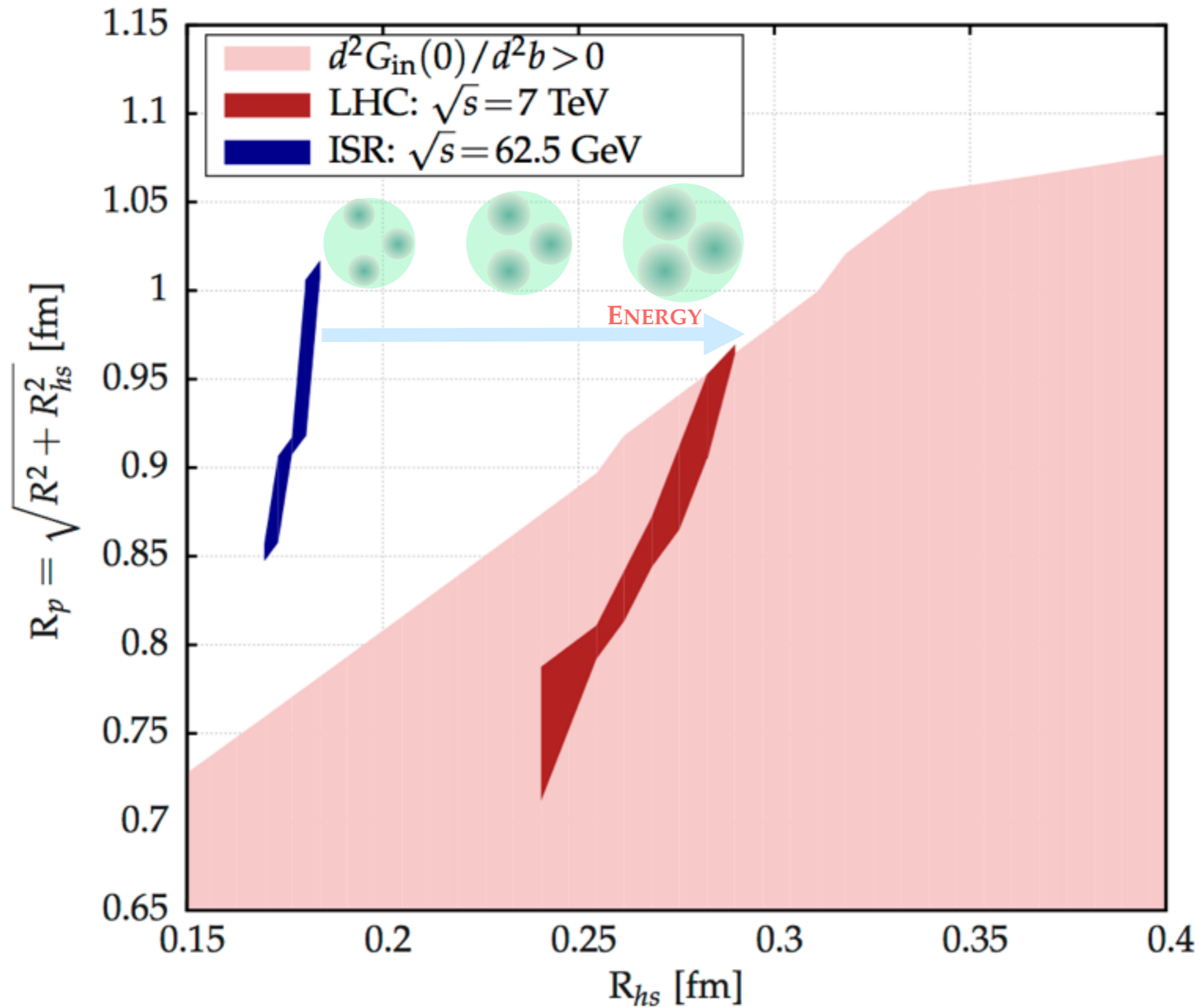
R_p vs R_{hs}

$r_c=0.5$ fm



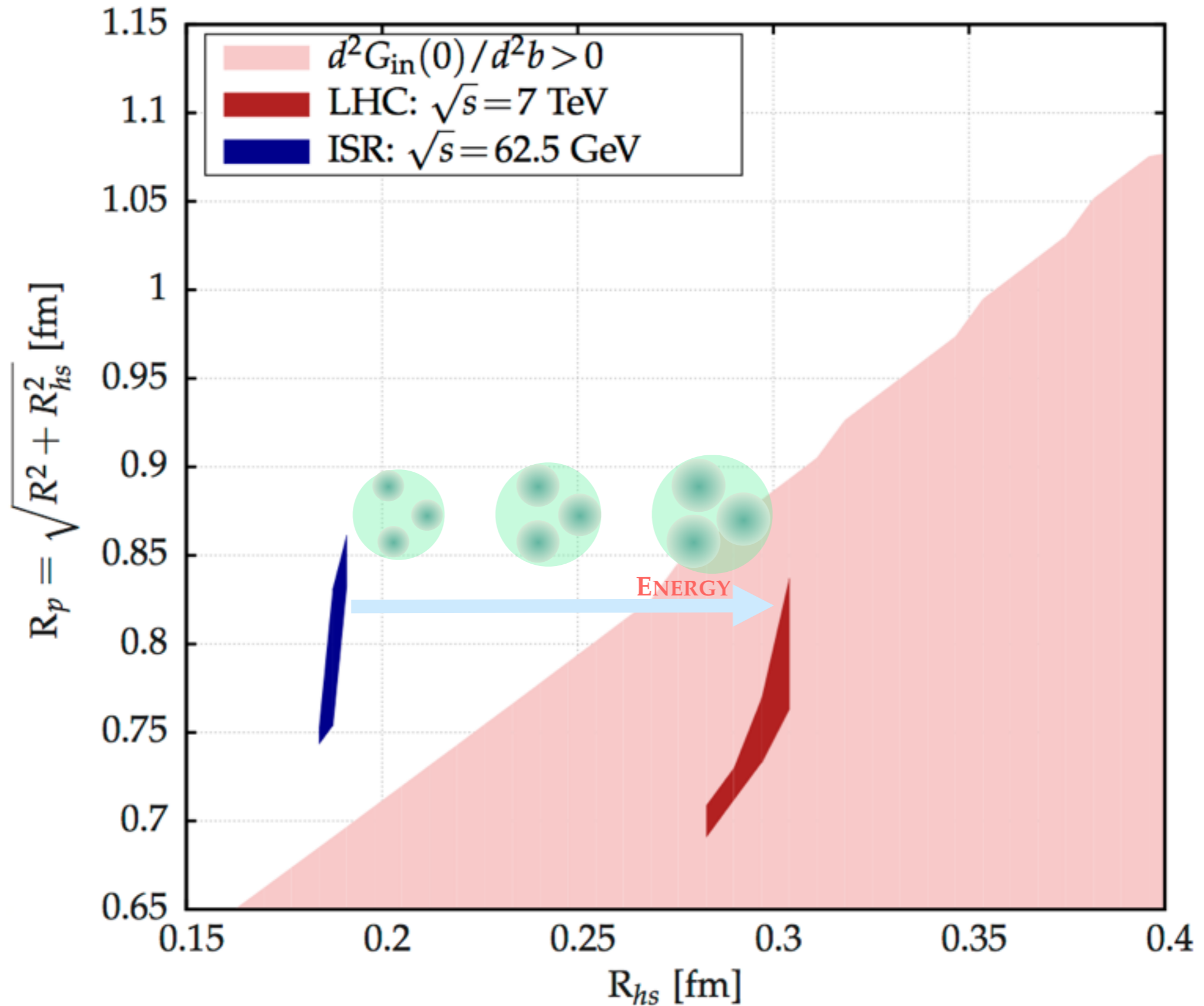
R_p vs R_{hs}

$r_c=0.5$ fm



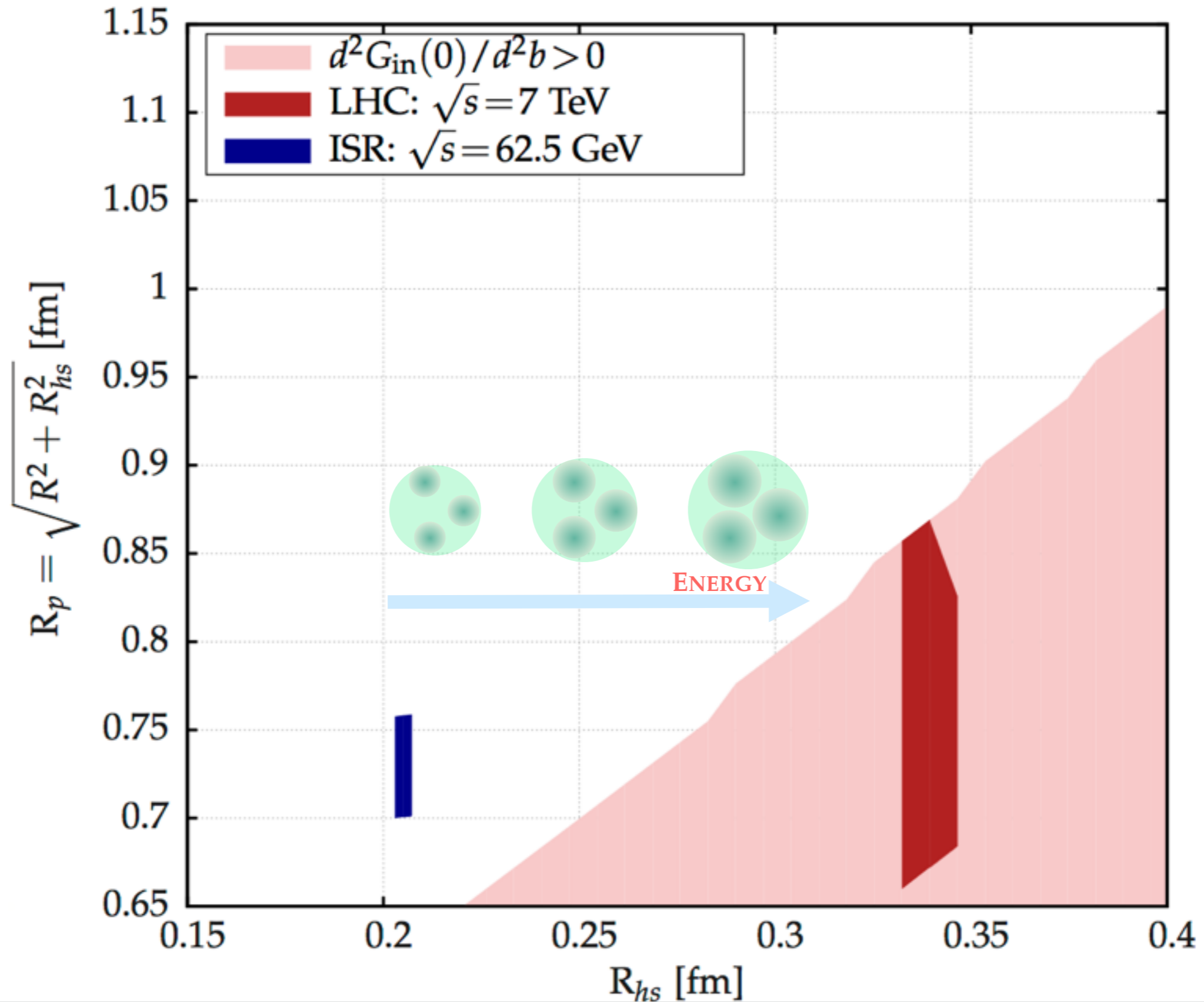
R_p vs R_{hs}

$r_c=0.4$ fm



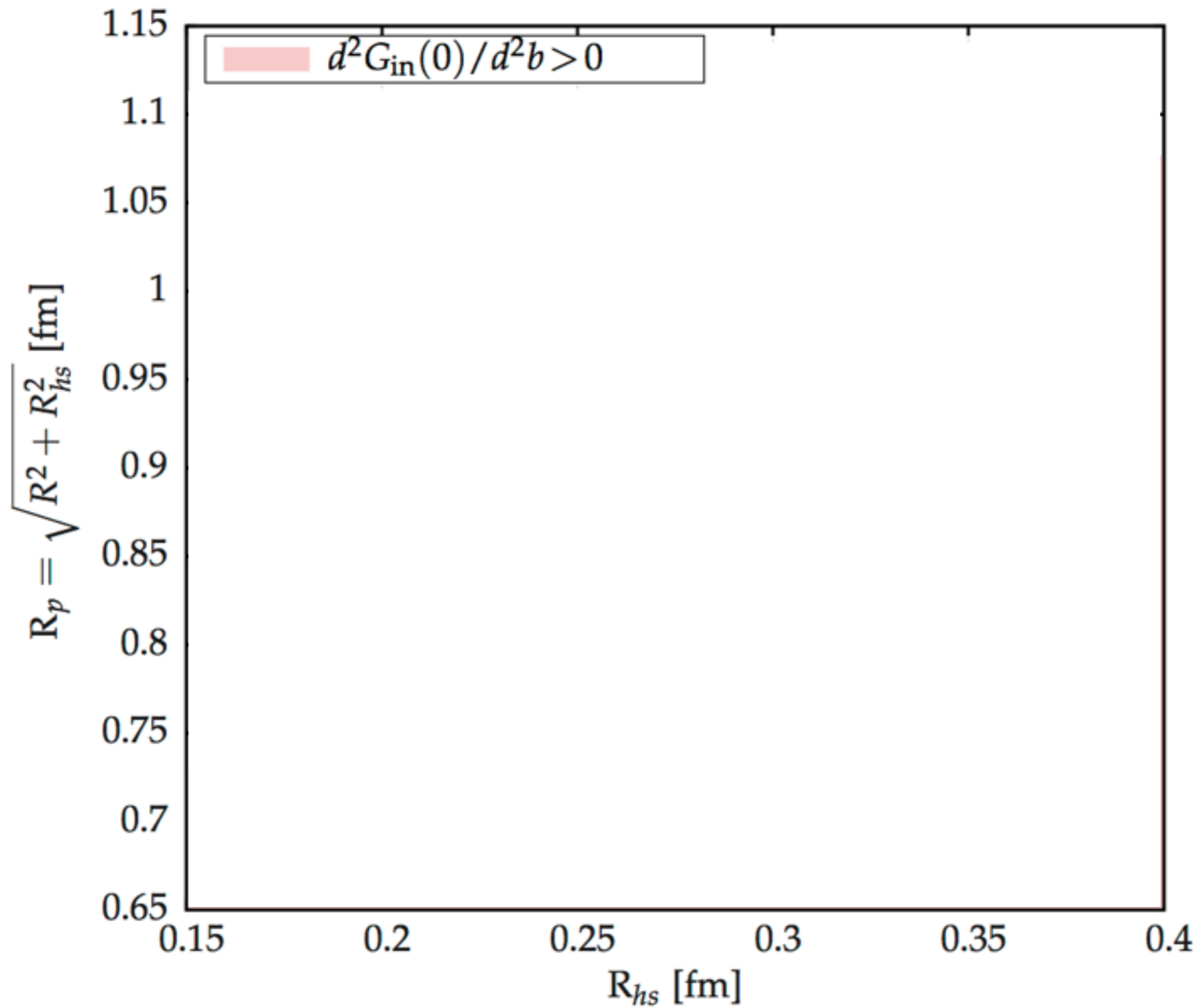
R_p vs R_{hs}

$r_c=0.3$ fm



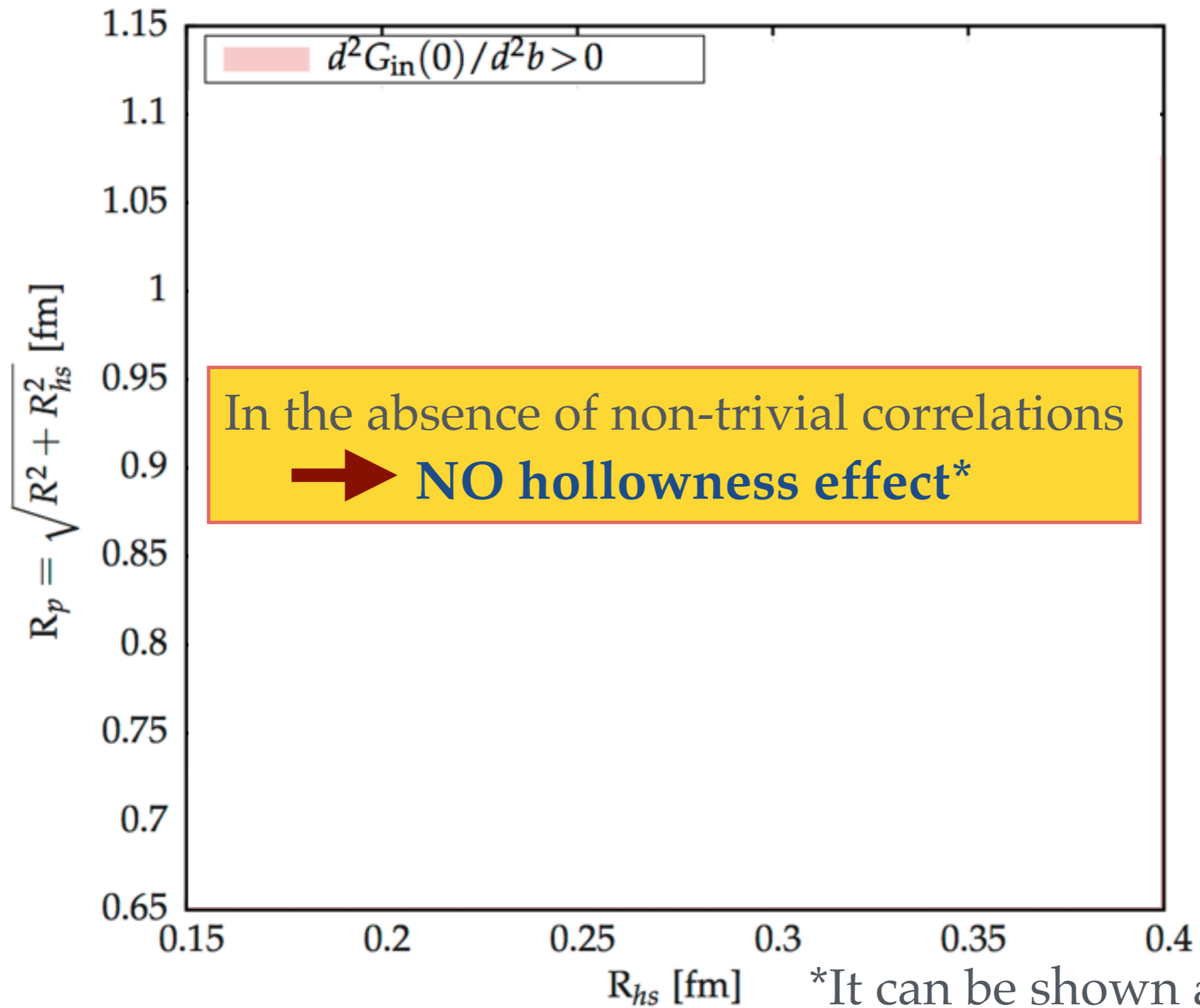
R_p vs R_{hs}

$r_c=0$ fm



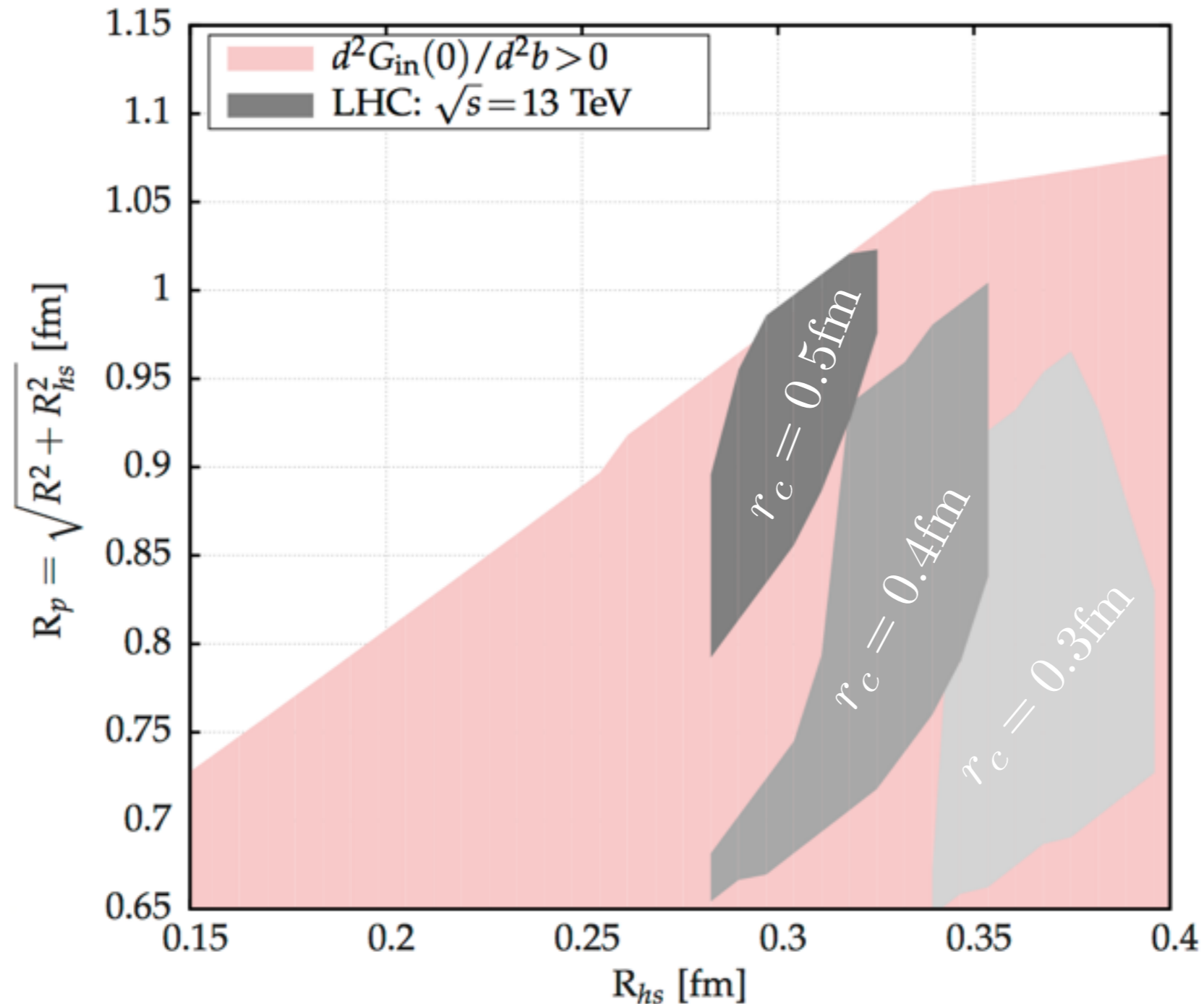
R_p vs R_{hs}

$r_c=0$ fm



*It can be shown analytically

- LHC, 13 TeV: $\sigma_{\text{tot}} = 11,15 \pm 1 [\text{fm}^2]$ $\rho = 0.14_{-0.08}^{+0.01}$
[COMPETE Collab. '02]



Take home message

- New and intriguing feature of hadronic interactions: *hollowness* effect.
- We propose a dynamical explanation based on:
 - Hot spots as the effective degrees of freedom.
 - Non-trivial correlations between the transverse positions of the hot spots.
 - Scattering amplitude from a Glauber-like multiple scattering series.
- Diffusion/ growth of the hot spots in the transverse plane with increasing collision energy is the key mechanism to explain the *hollowness* effect.
- Future work: impact of this new effect in other observables in pp and heavy ion collisions: flow harmonics, multiplicities...

However...

- Present data compatible with **NO hollowness** effect within error bars.
Full saturation of the unitarity bound at $b = 0$ ($G_{in}(b=0)=1$)
- Data at 13 TeV will (?) clarify the situation. Also data at $t > 2-3 \text{ GeV}^2$ would help to reduce the uncertainties associated to the Fourier transform at $b=0$

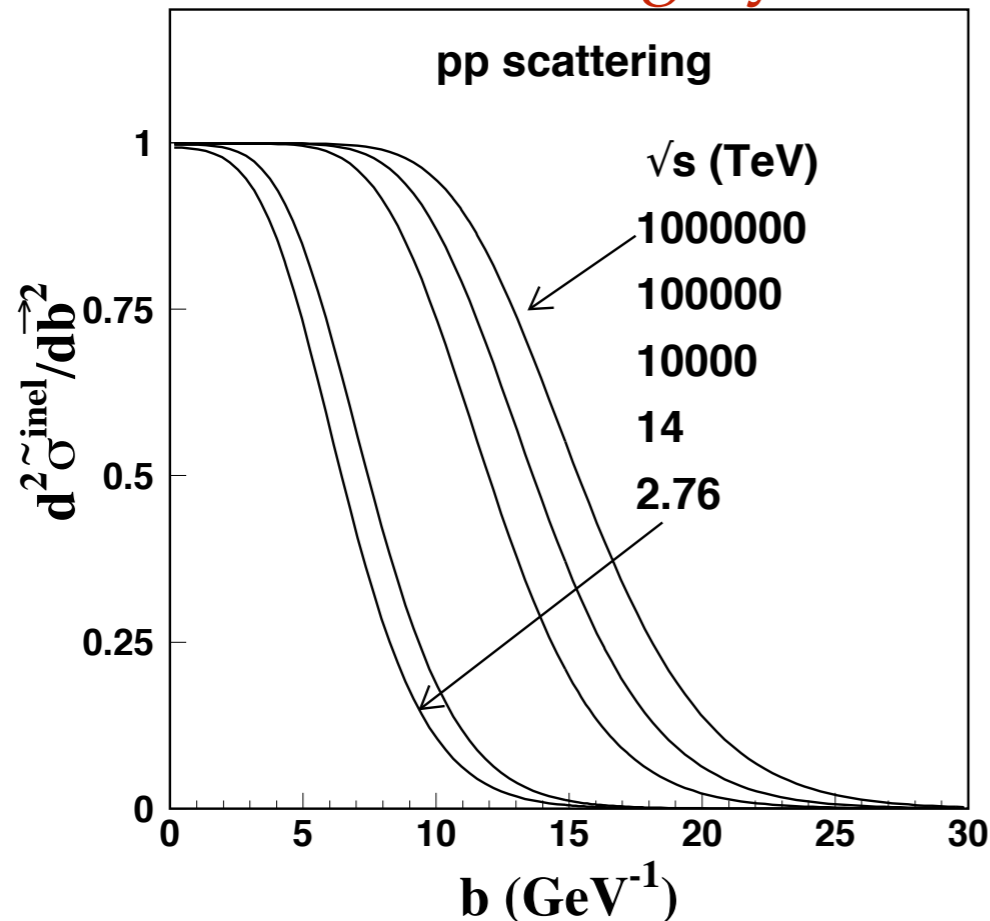
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$$\sigma_{el}/\sigma_{tot}(s \rightarrow \infty) =$$

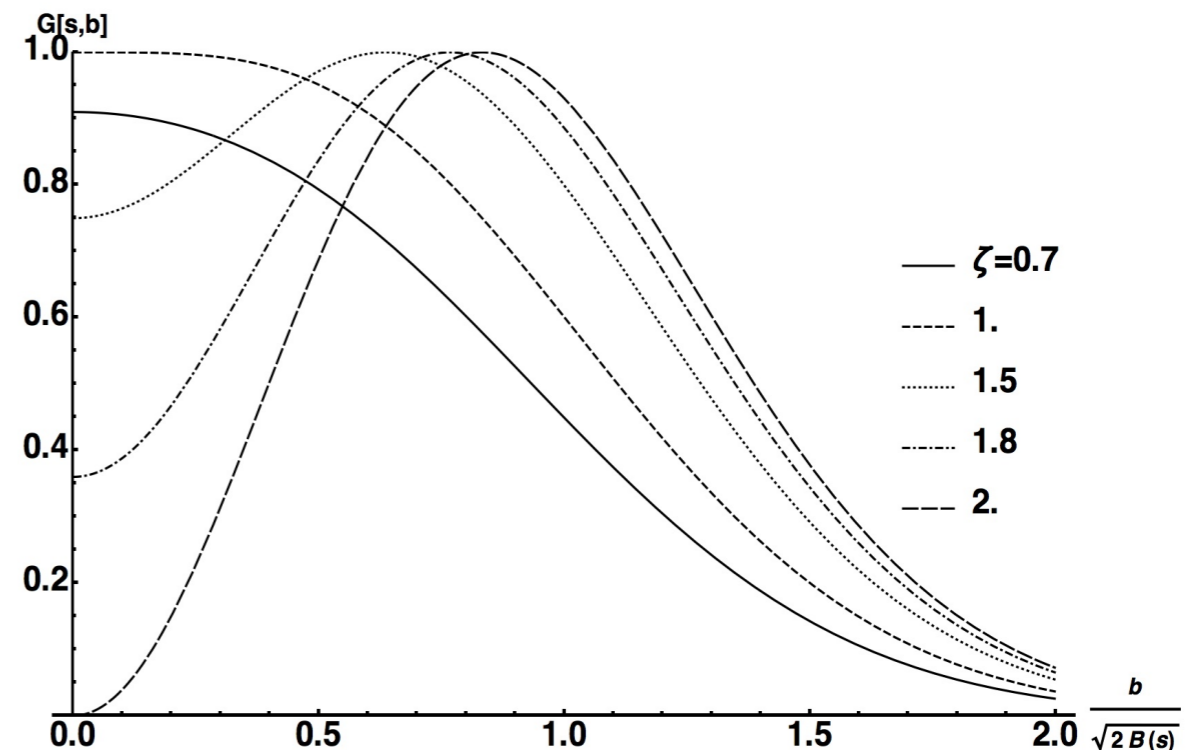
Fagundes et al '15]

0.35 (grey disk)



1 ("torus")

I M Dremin, '16]



However...

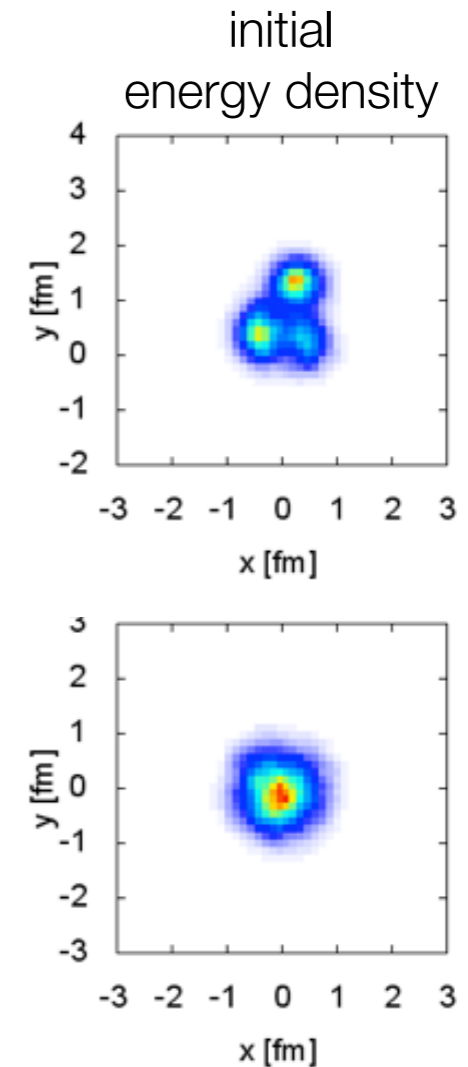
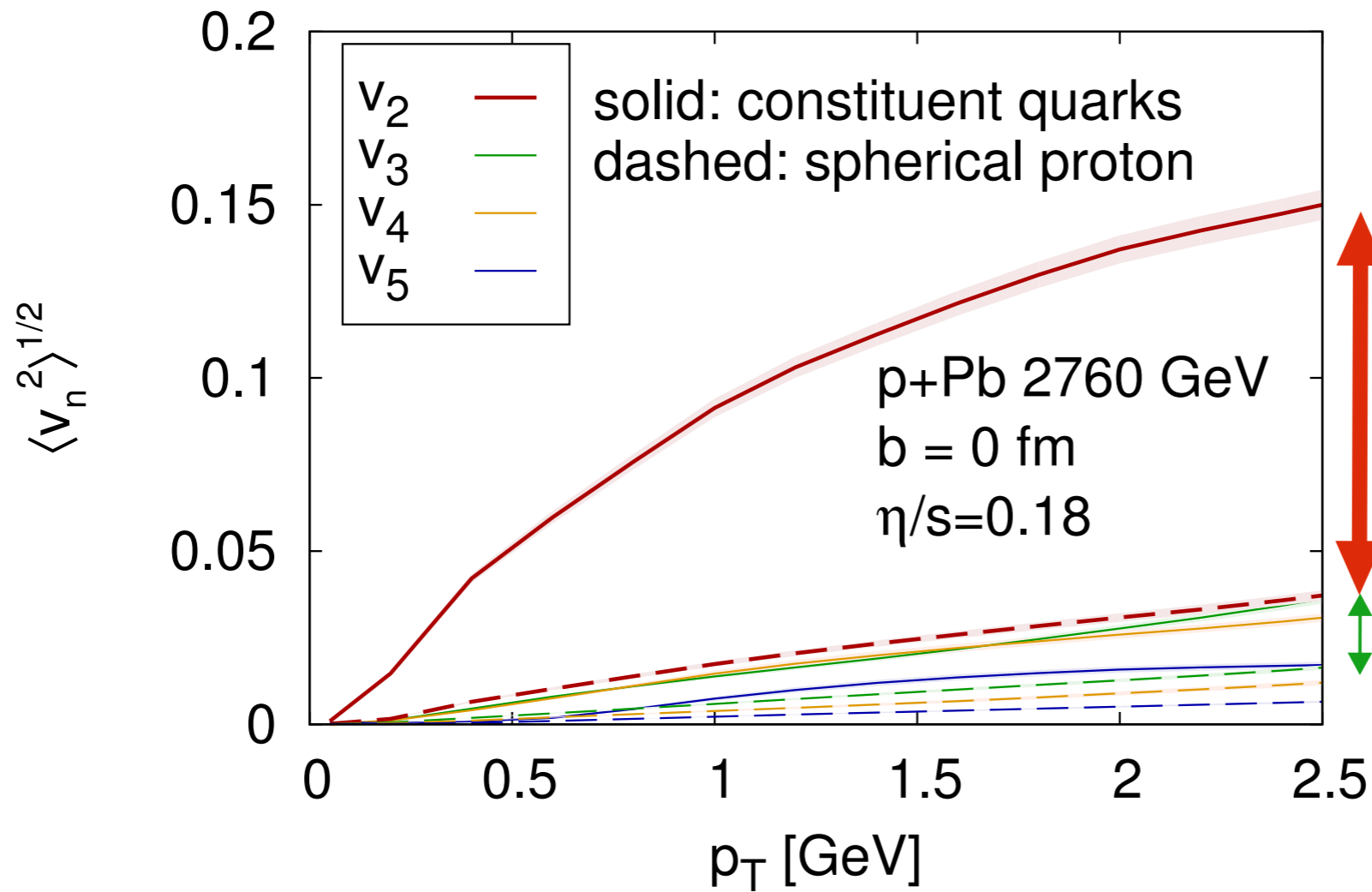
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- Data at 13 TeV will (?) clarify the situation. Also data at $t > 2-3 \text{ GeV}^2$ would help to reduce the uncertainties associated to the Fourier transform at $b=0$
- **Future work:** Better understanding of the role of correlations in a many-body scattering problem. Artefact of the eikonal unitarization?
- **Future work:** impact of this new effect in other observables in pp and **heavy ion collisions:** flow harmonics, multiplicities, soft-hard correlations etc...

Take home message

[Schenke'15]

Round vs. structured proton: IP-Glasma + MUSIC

It makes a huge difference!



22

Back up

pp elastic scattering

- Our approach starts from a generic parametrization

$$\text{Im}T_{el}(s, t) = a_1 e^{b_1 t} + a_2 e^{b_2 t} + a_3 e^{b_3 t}$$

$$\text{Re}T_{el}(s, t) = c_1 e^{d_1 t}$$

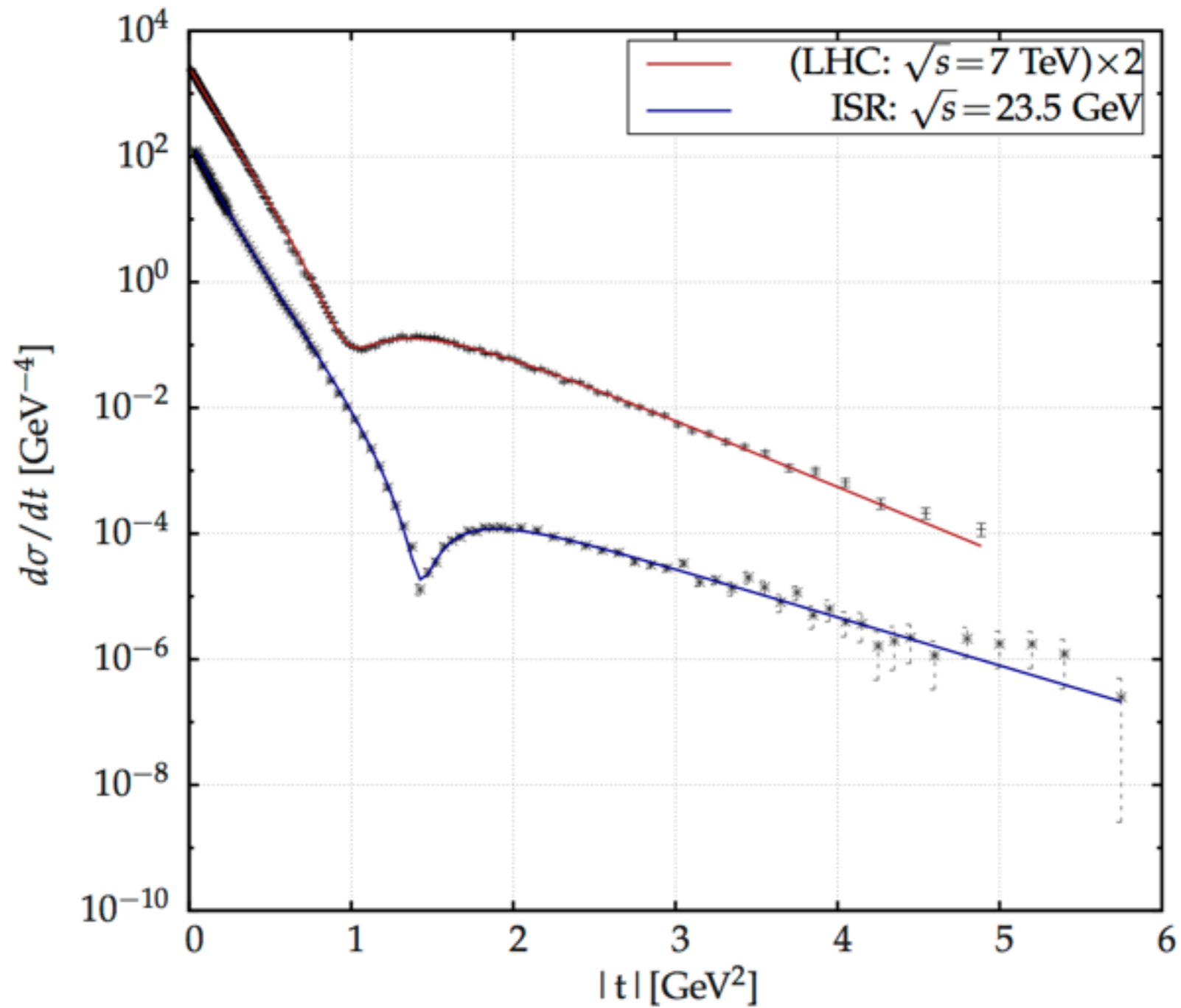
- Fit parameters are subject to two phenomenological constraints

$$\sigma_{\text{tot}} = 2 \sum_i a_i$$

$$\rho = \sum_i c_i / a_i$$

- Minimal number of parameters to reduce correlations

pp elastic scattering



$$\chi^2/\text{d.o.f} \sim 1.1 - 2$$