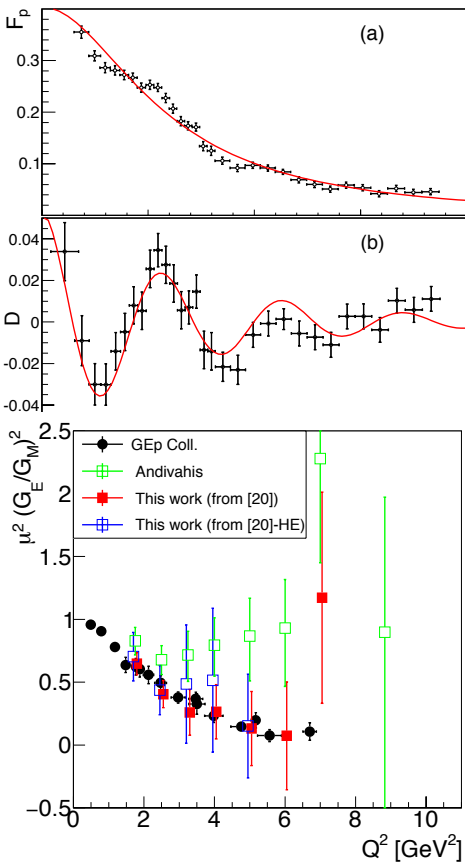


HSQCD
2016

Nucleon Form Factors Recent findings

Egle Tomasi-Gustafsson, *IRFU, SPhN-Saclay*
S. Pacetti, *Università di Perugia (Italy)*
A. Bianconi, *Università di Brescia (Italy)*

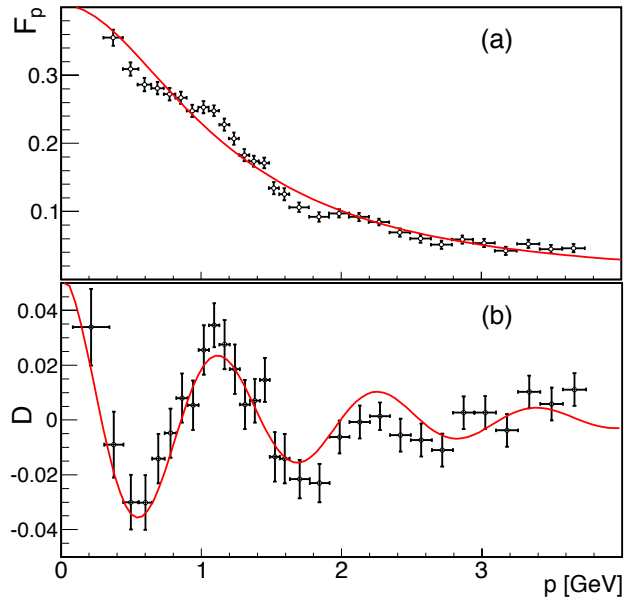


- Periodic structures in TL region
...accessing the hadron creation?
- Deviation from dipole in SL region?



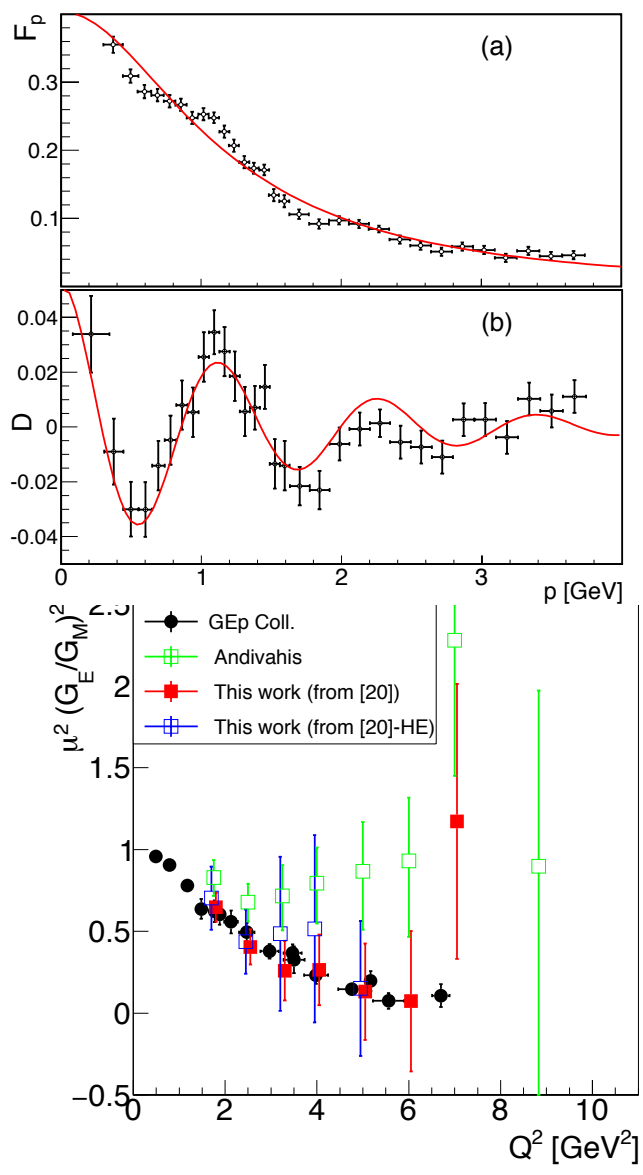
[S. Pacetti and R. Baldini-Ferrolì, ETG, Phys. Rept. 514 \(2014\) 1](#)
[A. Bianconi, ETG, PRL 114, 232301 \(2015\), PRC 93, 035201 \(2016\)](#)

Two questions



Periodic structures recently discovered in TL region
- Hadron creation from vacuum (Resonances?)

Proton EM Form Factors: two issues



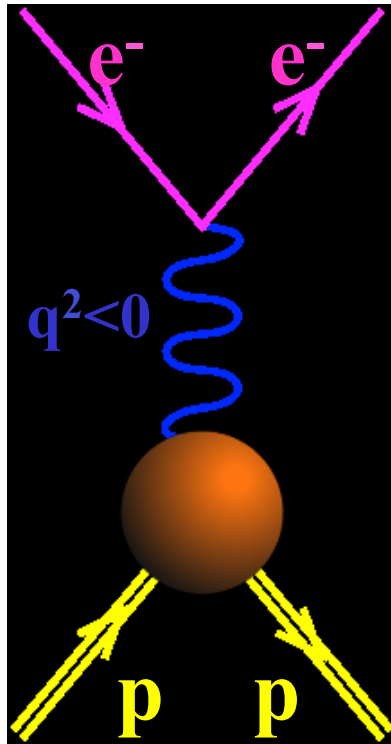
Periodic structures recently discovered in TL region

- *Hadron creation from vacuum (Resonances?)*

Discrepancy between polarized and unpolarized measurements of elastic EMFFs:

- *Is it real?*
- *Two photon exchange?!*

Electromagnetic Interaction



The electron vertex is known, γ_μ

The interaction is carried by a virtual photon of mass q^2

The proton vertex is parametrized in terms of FFs: Pauli and Dirac F_1, F_2

$$\Gamma_\mu = \gamma_\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}q^\nu}{2M} F_2(q^2)$$

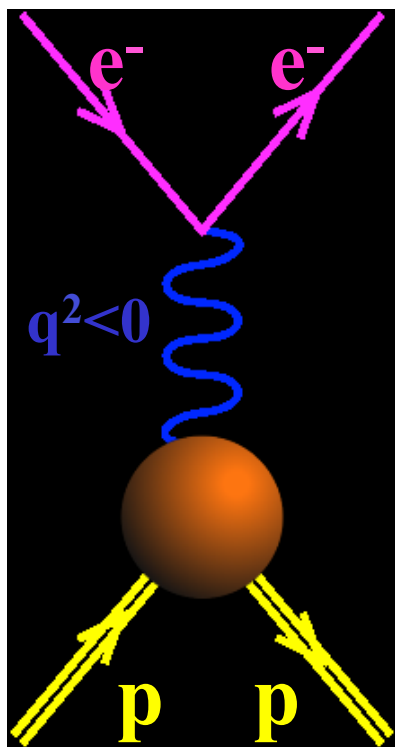
or in terms of Sachs FFs:

$$GE = F_1 - \tau F_2, \quad GM = F_1 + F_2, \quad \tau = -q^2/4M^2$$

What about high order radiative corrections?

Hadron Electromagnetic Form factors

$$\Gamma_\mu = \gamma_\mu F_1(q^2) + \frac{i\sigma_{\mu\nu}q^\nu}{2M} F_2(q^2)$$

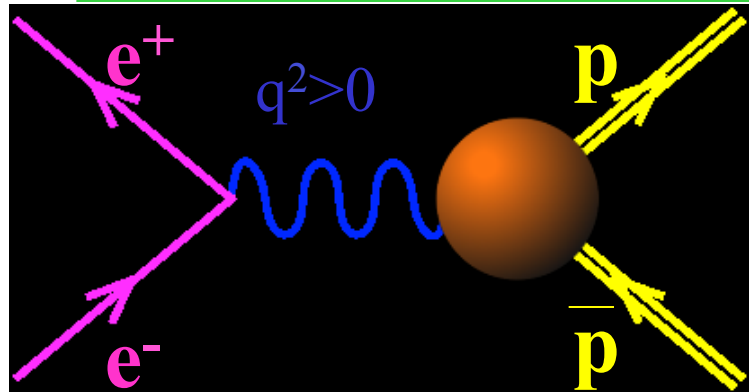


$GE(0)=1$
 $GM(0)=\mu_p$

*Space-like
 FFs are real*

*Unphysical region
 $p+\bar{p} \leftrightarrow e^+ + e^- + \pi^0$*

*Asymptotics
 - QCD
 - analyticity*



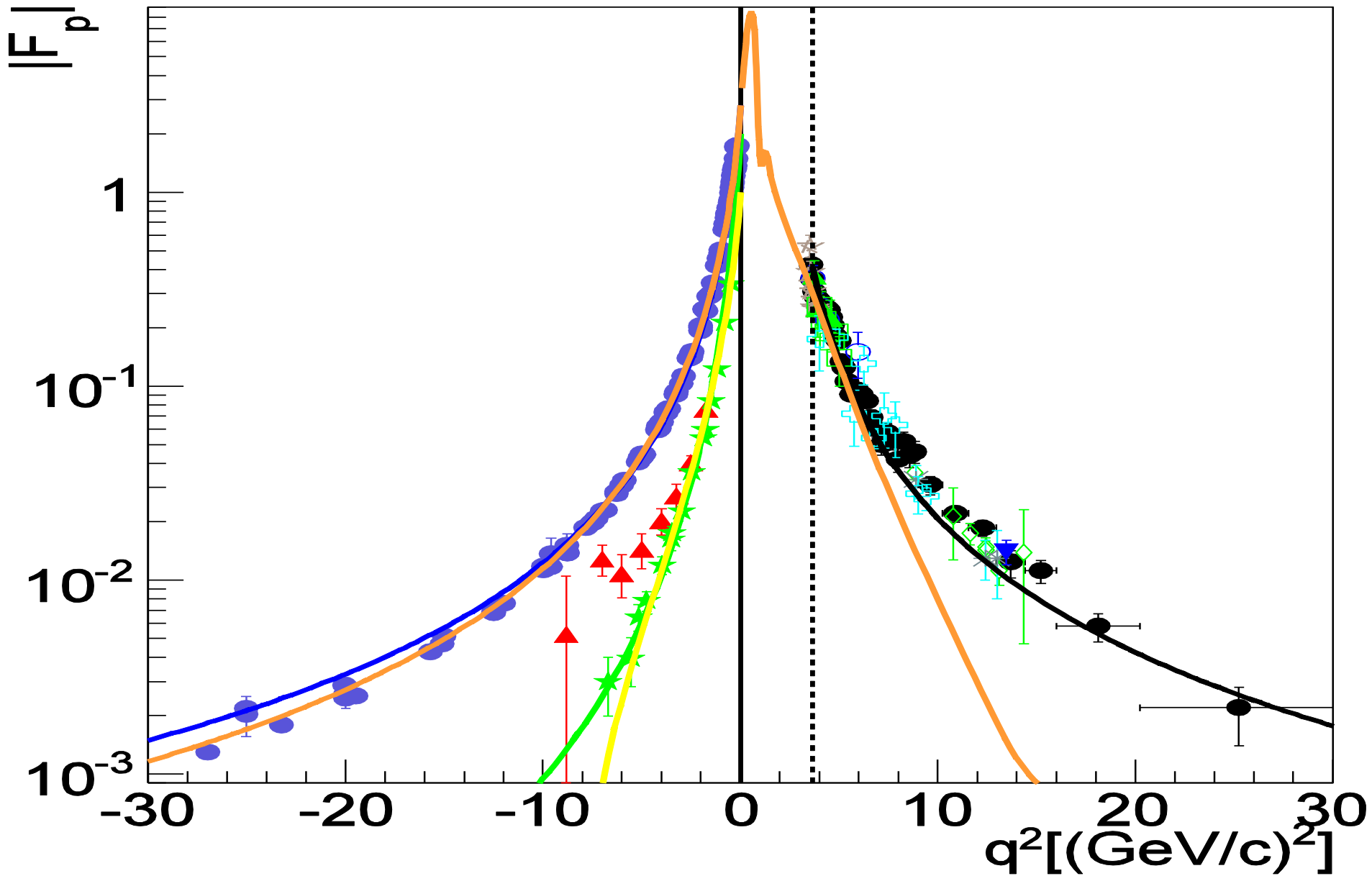
*Time-Like
 FFs are complex*

$e+p \rightarrow e+p$

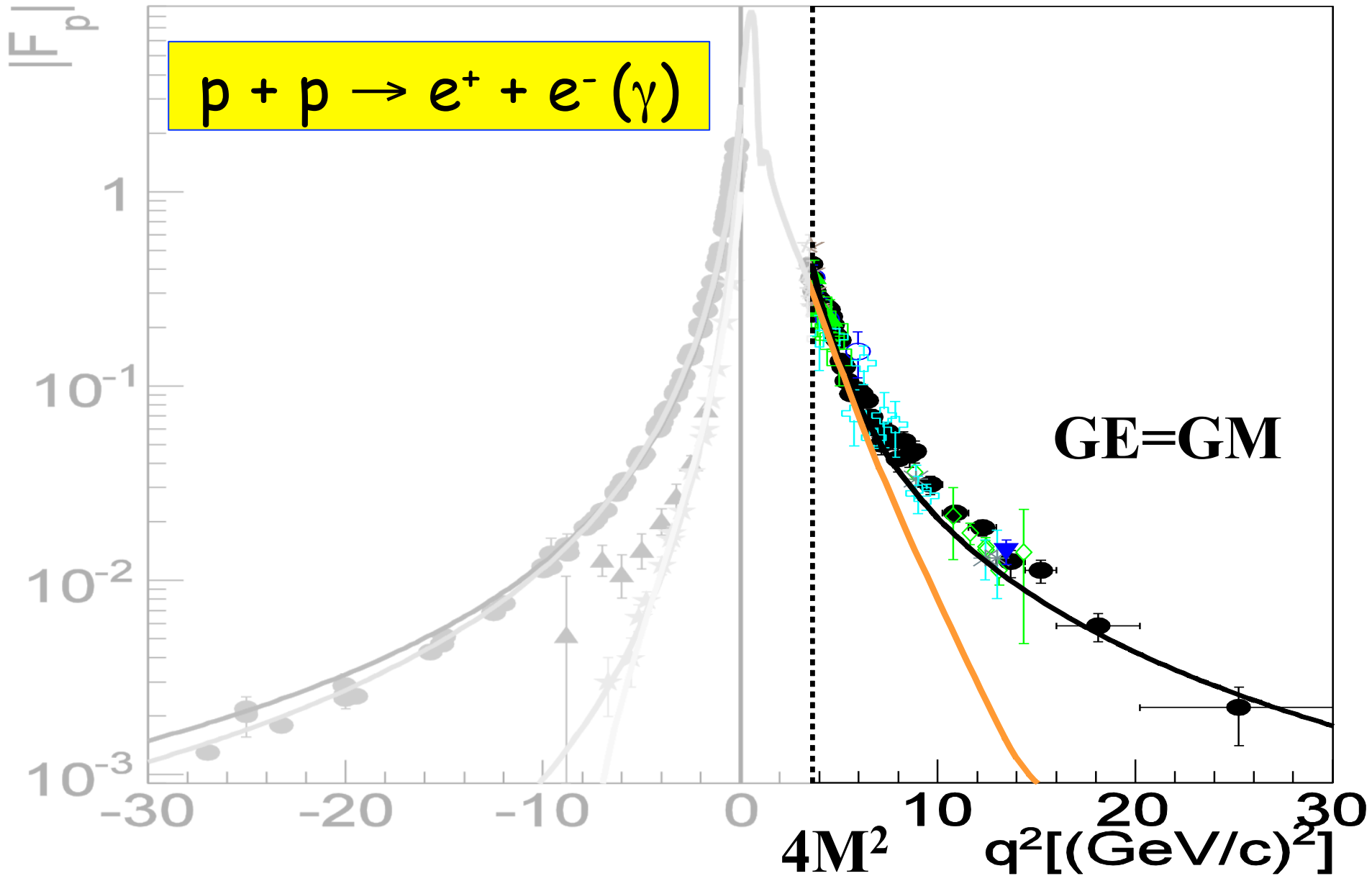
$q^2=4m_p^2$
 $GE=GM$

$p+\bar{p} \leftrightarrow e^+ + e^-$ q^2

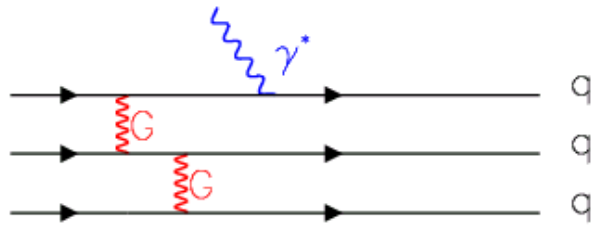
Proton Electromagnetic Form factors



The Time-Like region

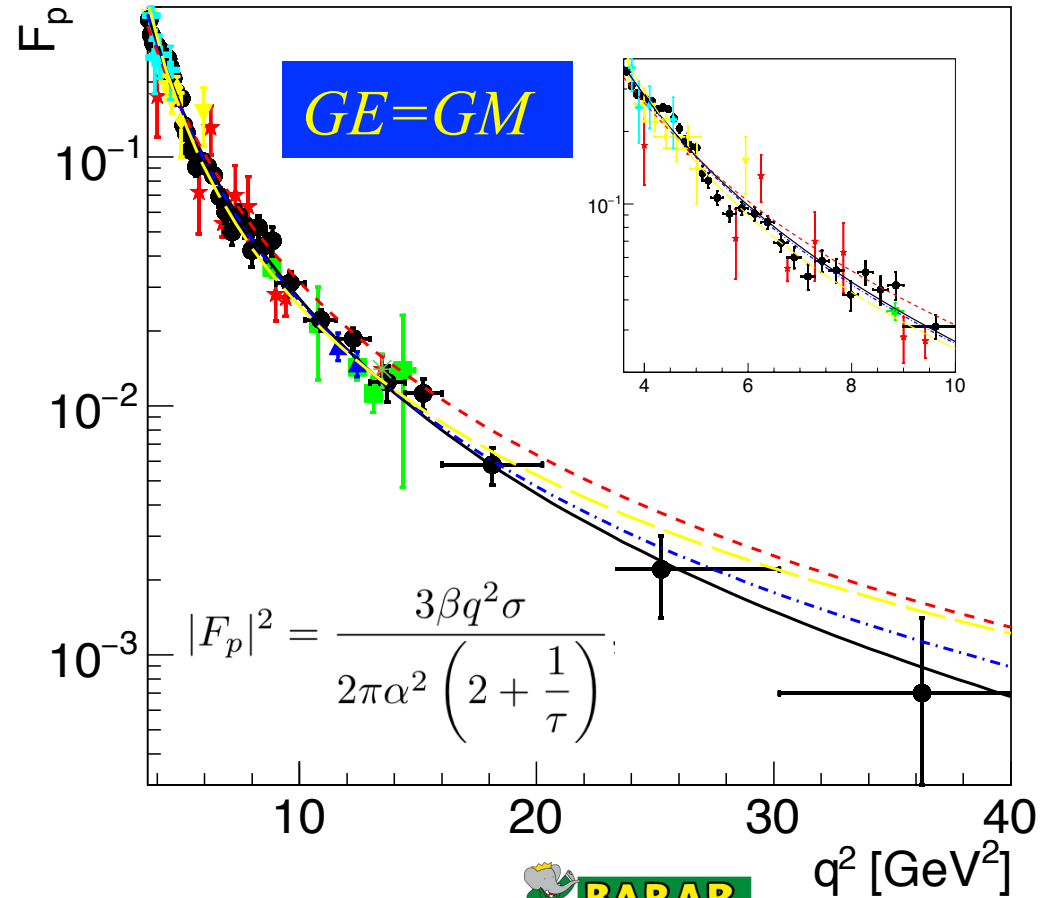
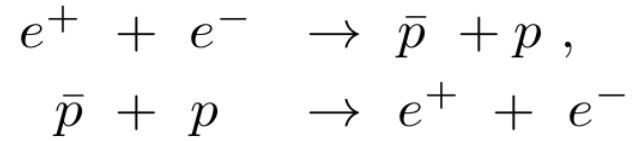


The Time-like Region

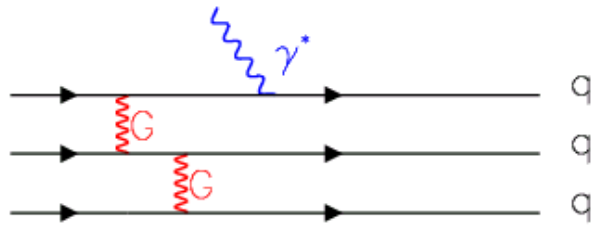


Expected QCD scaling $(q^2)^2$

$$|F_{scaling}(q^2)| = \frac{\mathcal{A}}{(q^2)^2 \log^2(q^2/\Lambda^2)}$$



The Time-like Region

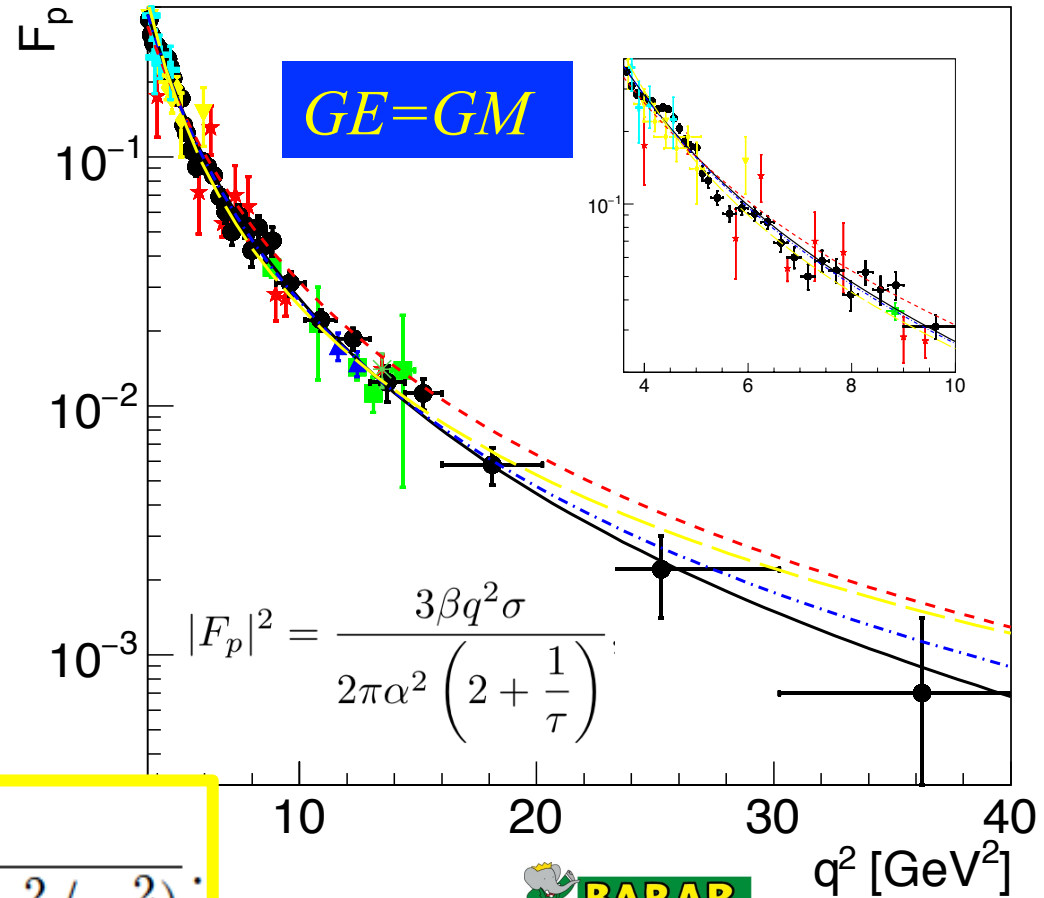
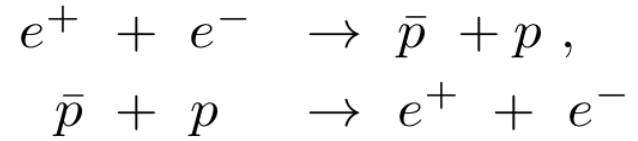


Expected QCD scaling $(q^2)^2$

$$\frac{A}{(q^2)^2 [\log^2(q^2/\Lambda^2) + \pi^2]}$$

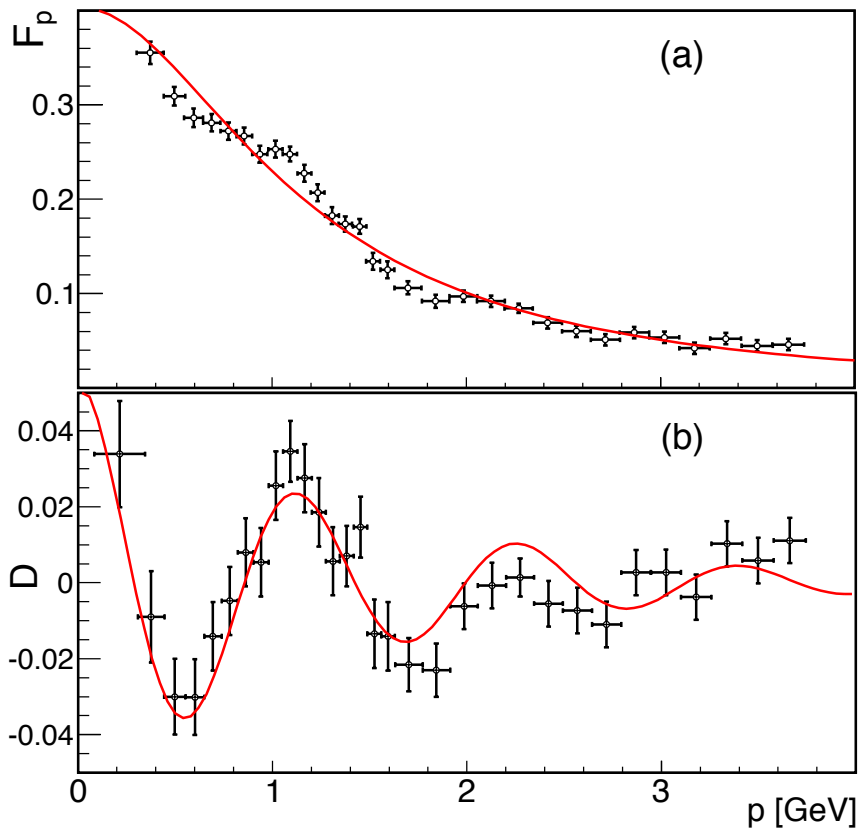
$$\frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2}$$

$$|F_{T3}(q^2)| = \frac{A}{(1 - q^2/m_1^2)(2 - q^2/m_2^2)}$$



Oscillations : regular pattern in P_{Lab}

The relevant variable is p_{Lab} associated to the relative motion of the final hadrons.



$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$

$A \pm \Delta A$	$B \pm \Delta B$	$C \pm \Delta C$	$D \pm \Delta D$	$\chi^2/n.d.f$
	$[GeV]^{-1}$	$[GeV]^{-1}$		
0.05 ± 0.01	0.7 ± 0.2	5.5 ± 0.2	0.03 ± 0.3	1.2

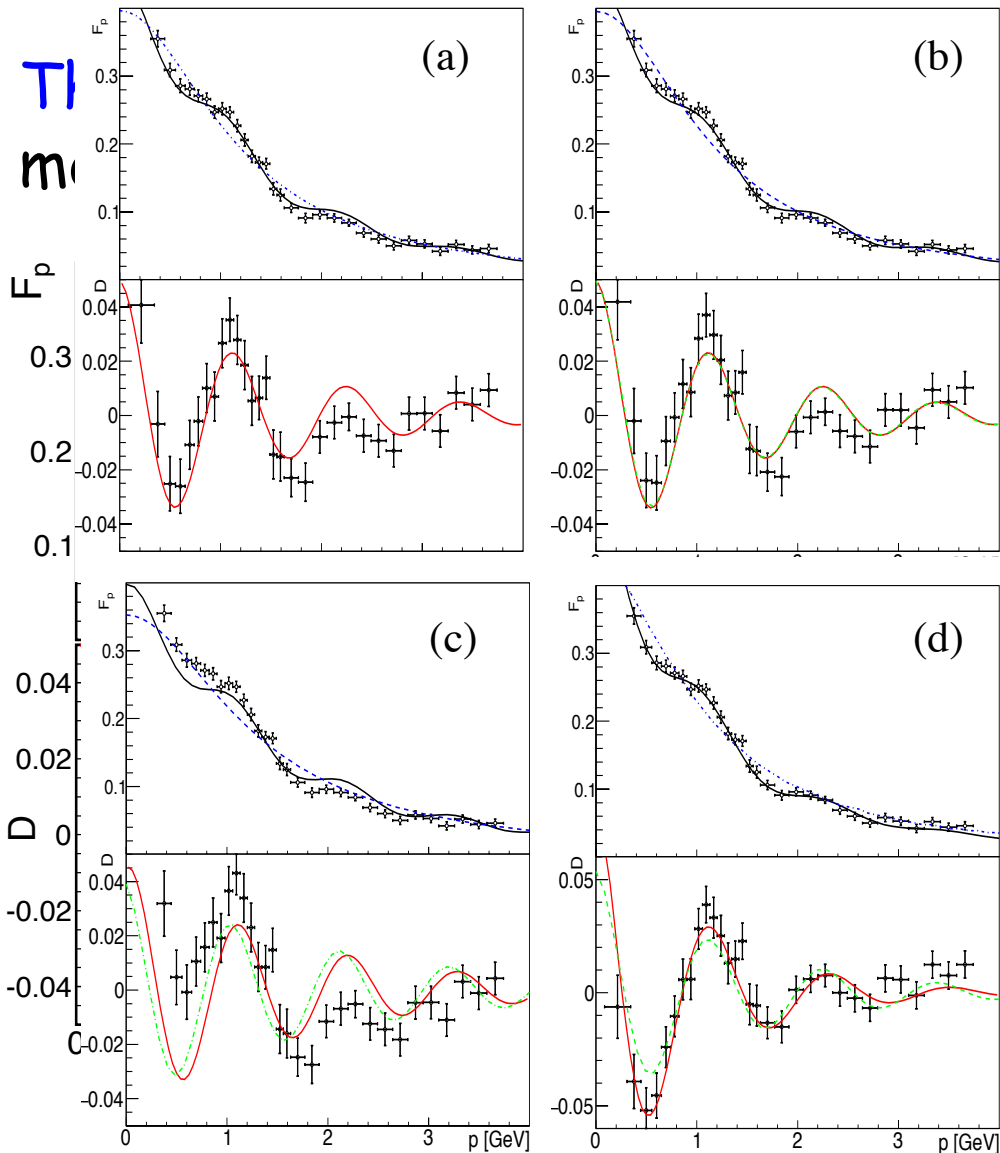
A: Small perturbation B: damping
C: $r < 1$ fm D=0: maximum at $p=0$

Simple oscillatory behaviour
Small number of coherent sources

A. Bianconi, E. T-G. Phys. Rev. Lett. 114,232301 (2015)

Oscillations : regular pattern in P_{Lab}

TI
m_i



related to the relative

$$p) \equiv A \exp(-Bp) \cos(Cp + D).$$

	$B \pm \Delta B$	$C \pm \Delta C$	$D \pm \Delta D$	$\chi^2/n.d.f$
	$[GeV]^{-1}$	$[GeV]^{-1}$		
1	0.7 ± 0.2	5.5 ± 0.2	0.03 ± 0.3	1.2

all perturbation B: damping
fm D=0: maximum at p=0

oscillatory behaviour
number of coherent sources

Bianconi, E. T-G. Phys. Rev. Lett. 114,232301 (2015), PRC 93, 035201 (2016)

Fourier Transform

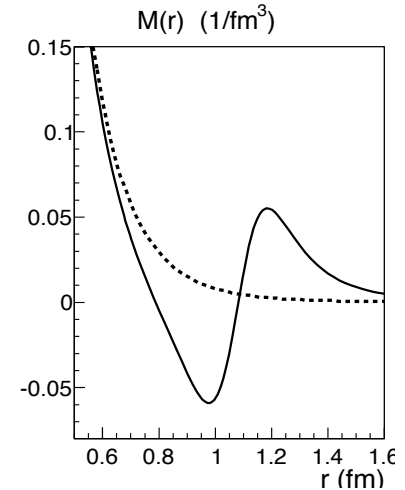
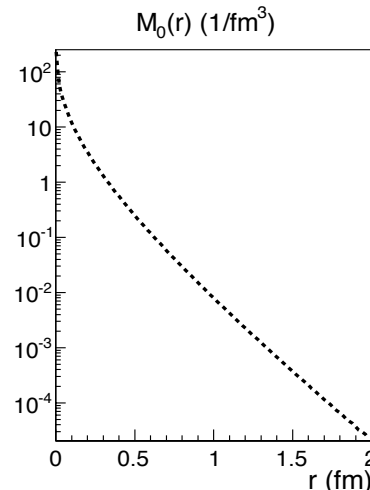


$$F_0(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M_0(r)$$

$$F(p) = F_0(p) + F_{osc}(p) \equiv \int d^3\vec{r} \exp(i\vec{p} \cdot \vec{r}) M(r).$$

$$F_0 = \frac{A}{(1 + q^2/m_a^2) [1 - q^2/0.71]^2},$$

$$F_{osc}(p) \equiv A \exp(-Bp) \cos(Cp + D).$$



- *Rescattering processes*
- *Large imaginary part*
- *Related to the time evolution of the charge density?*
(E.A. Kuraev, E. T.-G., A. Dbeyssi, PLB712 (2012) 240)
- *Consequences for the SL region?*
- *Data from BESIII confirm the structure*
- *Expected from PANDA*

Optical model analysis

The excited vacuum created by e^+e^- annihilation decays in multi-quark states: $p\bar{p}$ is one of them

- feeding at small r by decay of higher mass states in $p\bar{p}$
- depletion at large r from $p\bar{p}$ annihilation into mesons

From the $p\bar{p}$ point of view, the coupling with the other channels transforms into an imaginary potential that

- destroys flux (absorption - negative potential)
- generates flux (creation - positive potential)

Optical model analysis:

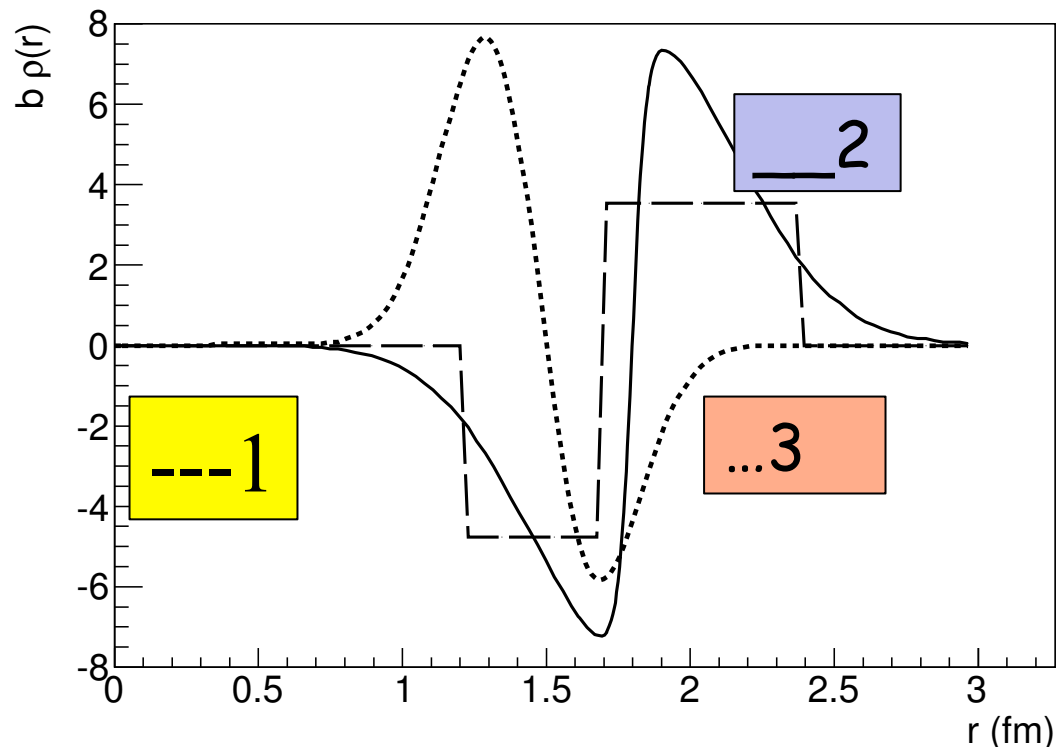
two component imaginary potential:

absorbing outside, regenerating inside

with steep change of sign.

Double layer potentials

Double layer rescattering densities : combination of two hollow potentials: one absorbing and one generating (imaginary potentials).



1) Multiple step function

2) Soft multistep

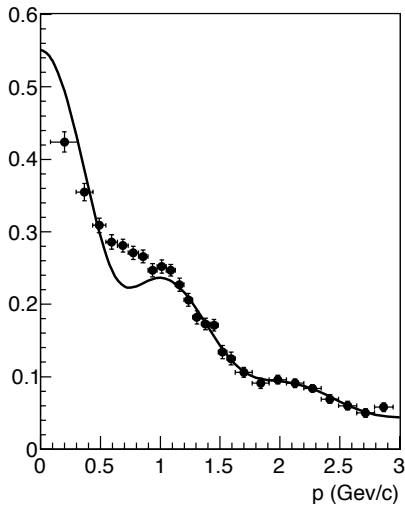
3) Two-gaussian opposite sign potential

A. Bianconi, E. T-G., PRC 93, 035201 (2016)

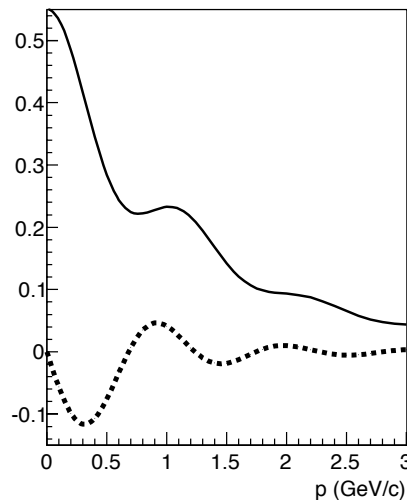
Optical model analysis

1) Multiple step function

Model fit

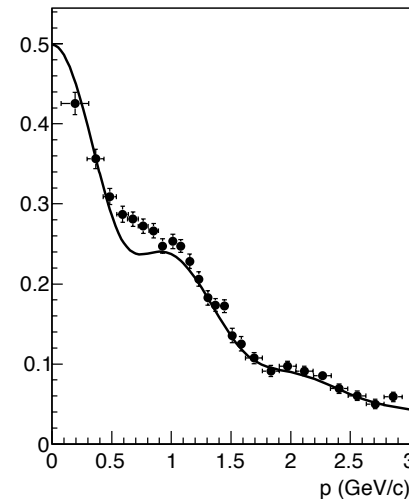


Re(F) and Im(F)

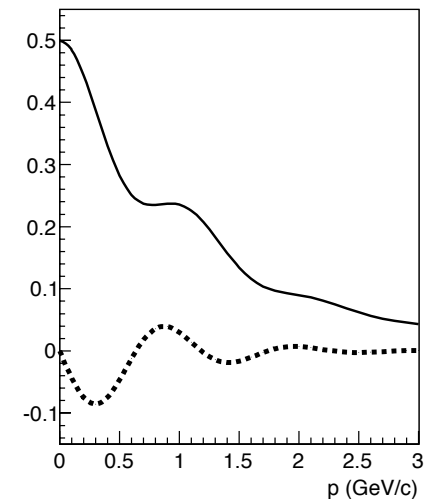


2) Soft multistep

Model fit

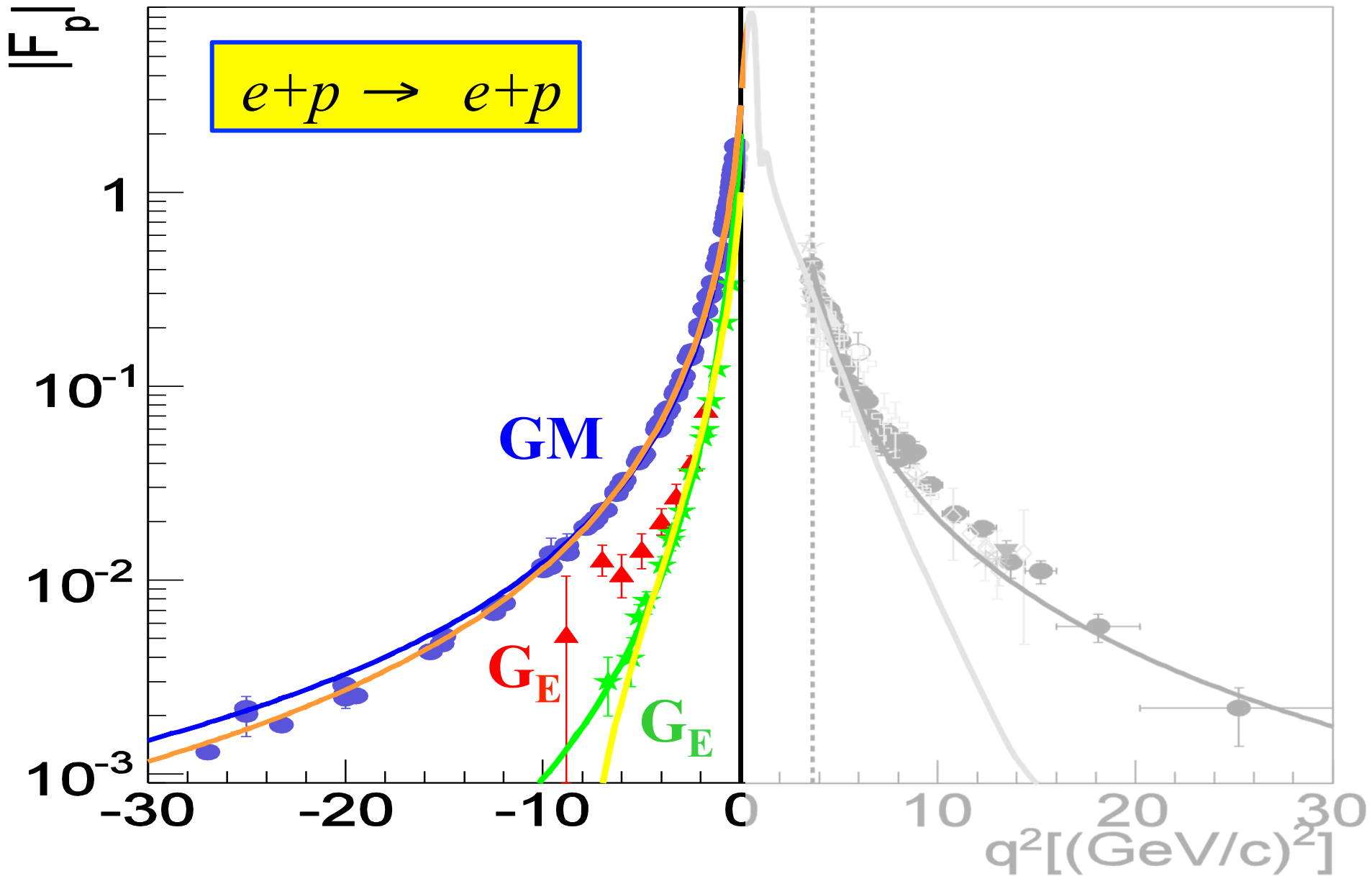


Re(F) and Im(F)



- At large r : purely absorptive
- At small r : the product $D(r)M(r)$ "resonates" with the FT factor
- Importance of the steep behavior (oscillation period)
- Related to threshold enhancement

The Space-Like region

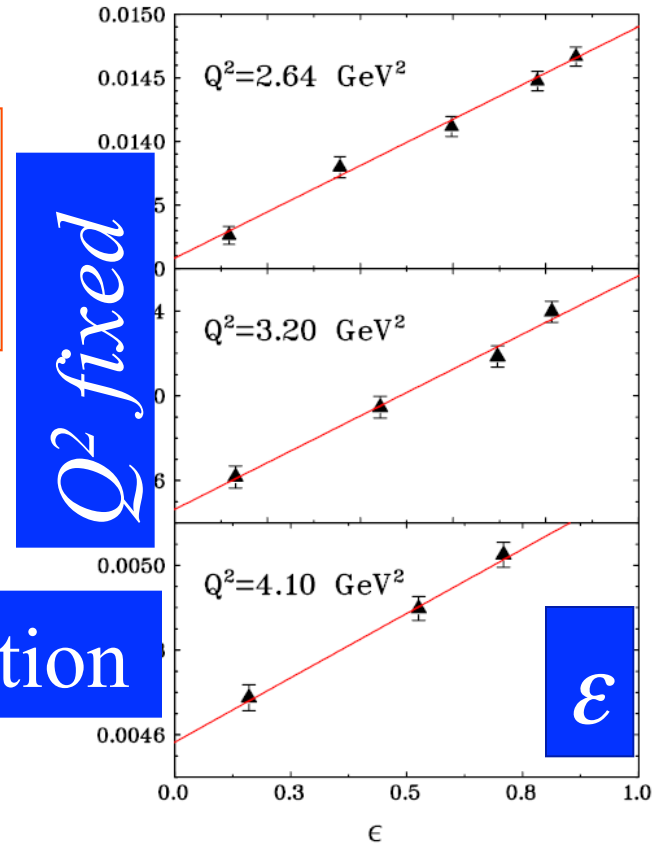


The Rosenbluth separation

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \frac{1}{(1+\tau)} \left(G_E^2(Q^2) + \frac{\tau}{\varepsilon} G_M^2(Q^2) \right)$$

$$\varepsilon = \left(1 + 2(1+\tau) \tan^2 \left(\frac{\theta_e}{2} \right) \right)^{-1}, \quad \tau = \frac{Q^2}{4M^2}$$

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2$$



Linearity of the reduced cross section

→ $\tan^2 \theta_e$ dependence

→ Holds for 1γ exchange only

PRL 94, 142301 (2005)

The polarization method (theory:1967)

SOVIET PHYSICS - DOKLADY

VOL. 13, NO. 6

DECEMBER, 1968

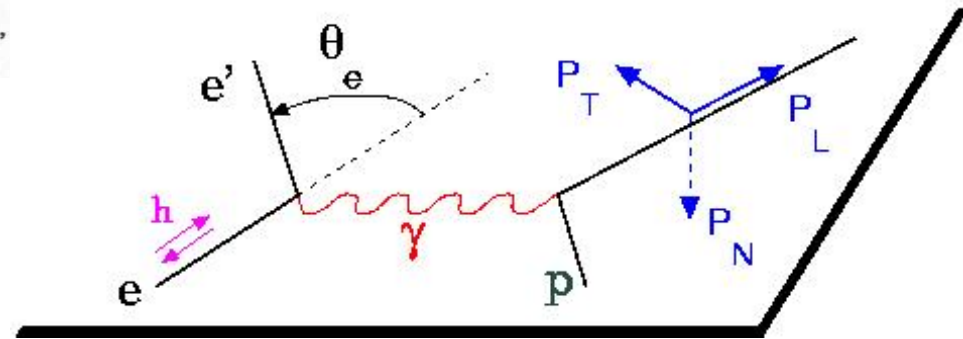
PHYSICS

POLARIZATION PHENOMENA IN ELECTRON SCATTERING BY PROTONS IN THE HIGH-ENERGY REGION

Academician A. I. Akhiezer* and M. P. Rekalo

Physicotechnical Institute, Academy of Sciences of the Ukrainian SSR
Translated from Doklady Akademii Nauk SSSR, Vol. 180, No. 5,
pp. 1081-1083, June, 1968
Original article submitted February 26,

$$s_2 \frac{d\sigma}{d\Omega_R} = 4p_2 \frac{(\mathbf{s} \cdot \mathbf{q})}{1 + \tau} \Gamma(\theta, \varepsilon_1) \left[\tau G_M (G_M + G_E) - \frac{1}{4\varepsilon_1} G_M (G_E - \tau G_M) \right],$$



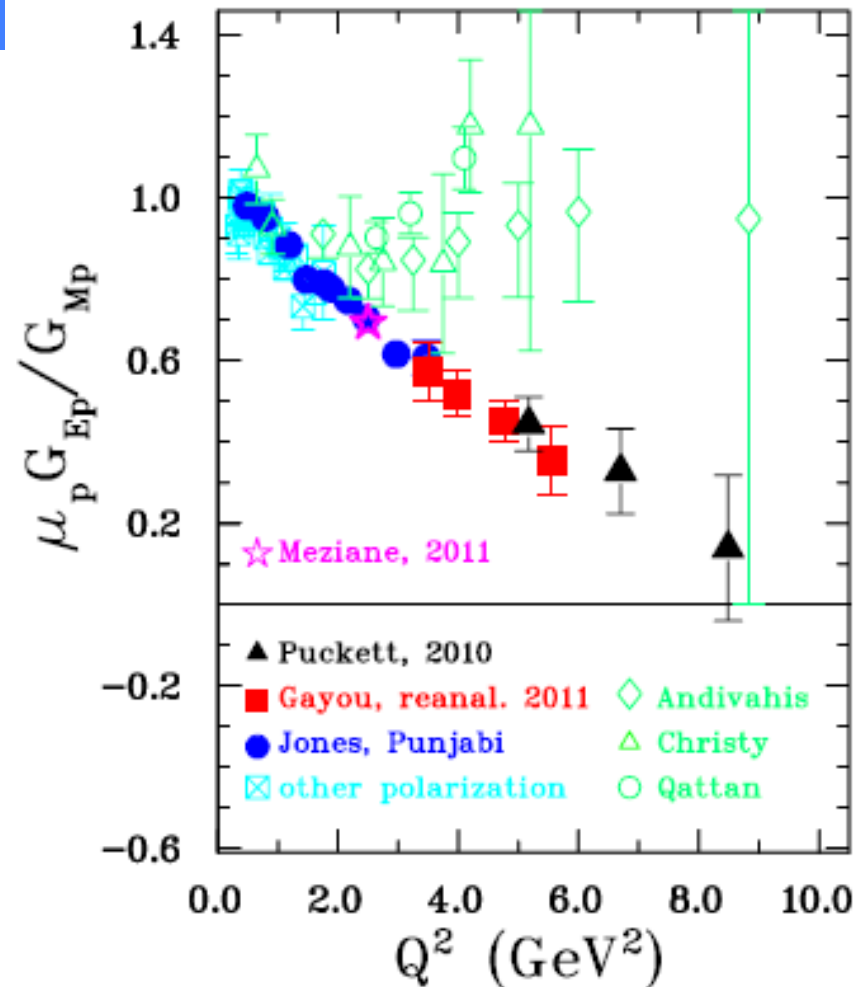
The polarization induces a term in the cross section proportional to $G_E G_M$
Polarized beam and target or
polarized beam and recoil proton polarization

Polarization Experiments

A.I. Akhiezer and M.P. Rekalo, 1967

Jlab-GEp collaboration

- 1) "standard" **dipole function** for the nucleon magnetic FFs **G_{Mp}** and **G_{Mn}**
- 2) **linear deviation** from the dipole function for the electric proton FF **G_{Ep}**
- 3) **QCD scaling** not reached
- 3) **Zero crossing** of G_{Ep} ?
- 4) **contradiction between polarized and unpolarized measurements**



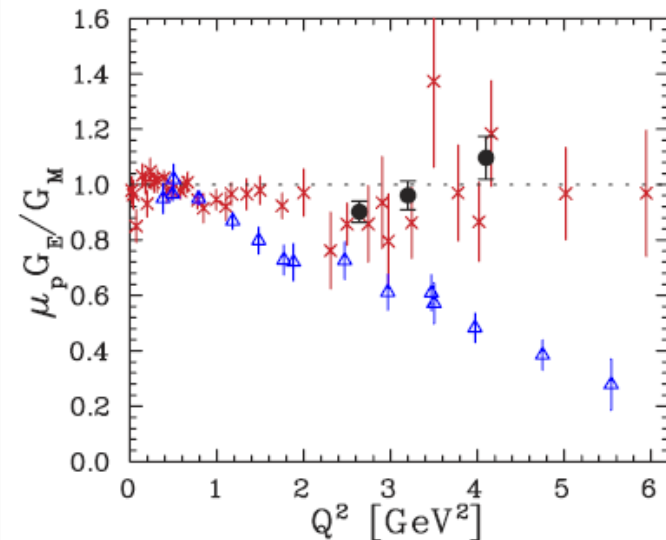
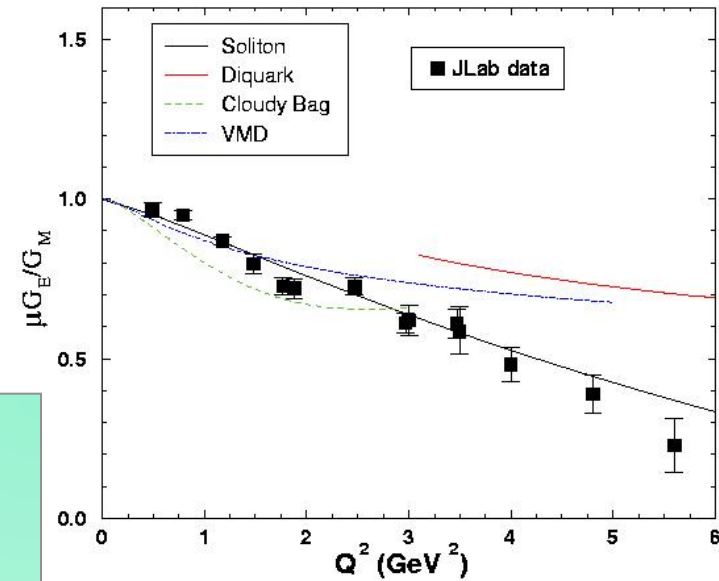
A.J.R. Puckett et al, PRL (2010), PRC (2012)

Issues

- Some models (IJL 73, Diquark, soliton..) predicted such behavior before the data appeared

BUT

- Simultaneous description of the four nucleon form factors...
- ...in the space-like and in the time-like regions
- Consequences for the light ions description
- When pQCD starts to apply?
- Source of the discrepancy



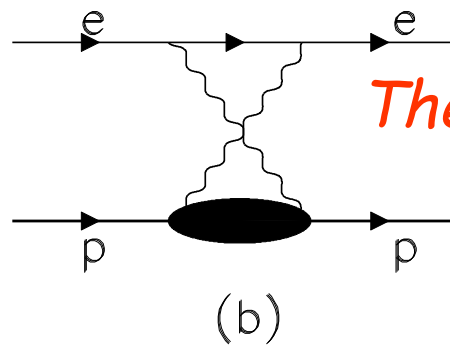
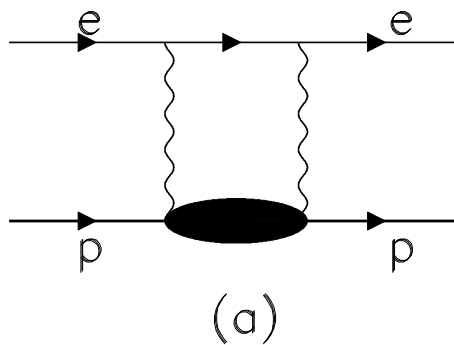
***Reaction mechanism:
1 γ -2 γ interference ?***

Radiative corrections?

Two photon exchange

- 1γ - 2γ interference is of the order of $\alpha=e^2/4p=1/137$

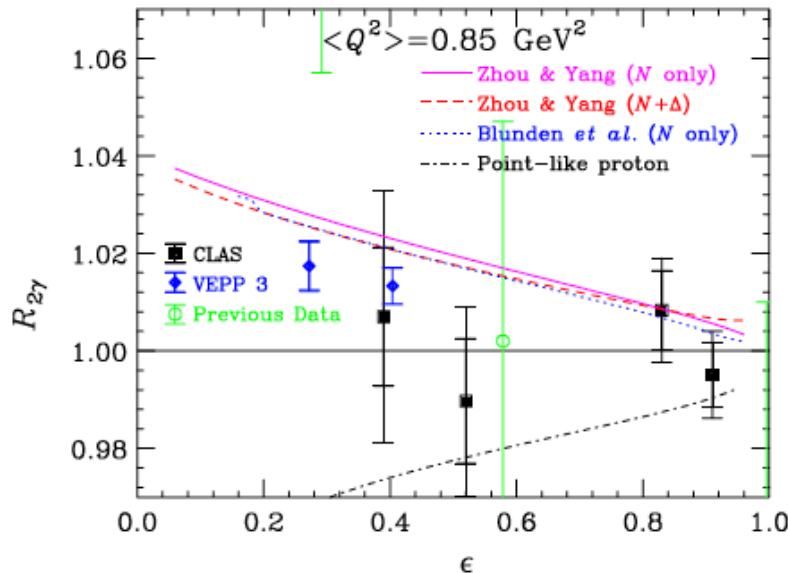
- In the 70's it was shown [J. Gunion and L. Stodolsky, V. Franco, F.M. Lev, V.N. Boitsov, L. Kondratyuk and V.B. Kopeliovich, R. Blankenbecker...] that, at large momentum transfer, the sharp decrease of the FFs, if the momentum is shared between the two photons, may compensate α
- The calculation of the box amplitude requires the description of intermediate nucleon excitation and of their FFs at any Q^2 ..
- Different calculations give quantitatively different results ·



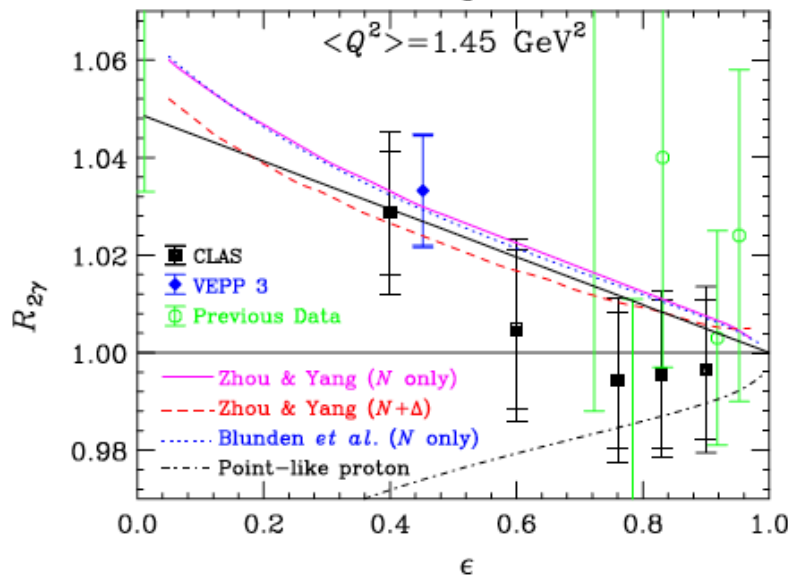
Theory not enough constrained!

CLAS, VEPP, OLYMPUS...

V. Rimal, ArXiv 1603.003151



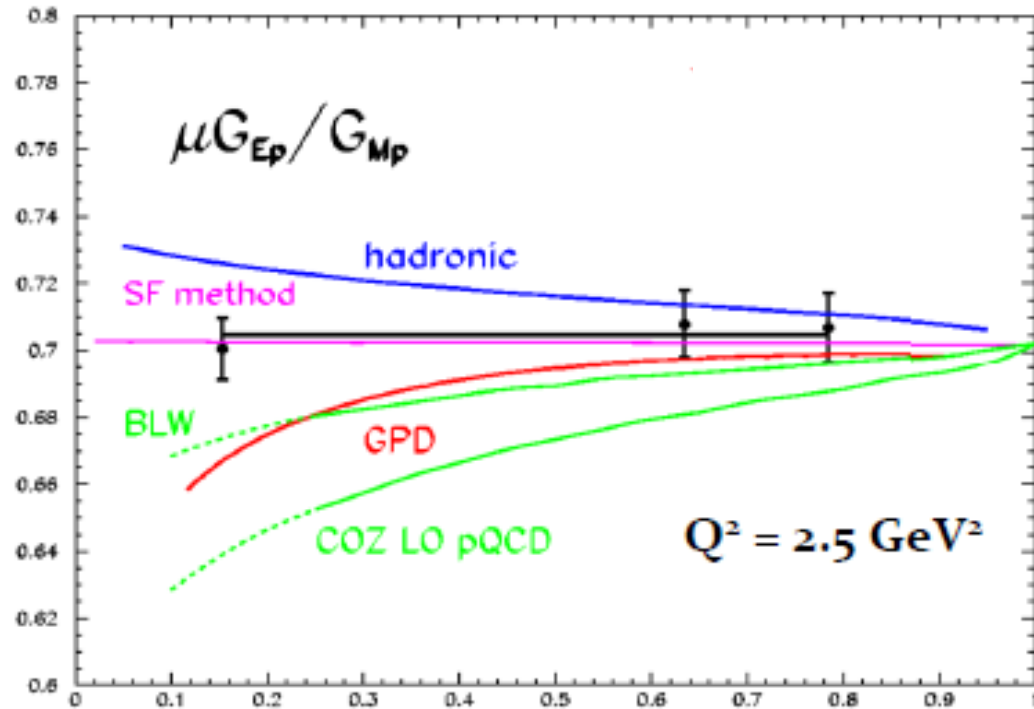
- $Q^2 < 2 \text{ GeV}^2$
- Effect $< 2\%$
- No evident increase with Q^2



Polarization ratio (ε -dependence)

- **DATA:** No evidence of ε -dependence at 1% level

- **MODELS:** large correction (opposite sign) at small ε



- **SF method:** ε -(almost)independent corrections

- **Theory:** corrections to the Born approximation at $Q^2 = 2.5 \text{ GeV}^2$

Y. Bystritskiy, E.A. Kuraev and E.T.-G, Phys.Rev.C75: 015207 (2007)

P. Blunden et al., Phys. Rev. C72:034612 (2005) (mainly GM)

A. Afanasev et al., Phys. Rev. D72:013008 (2005) (mainly GE)

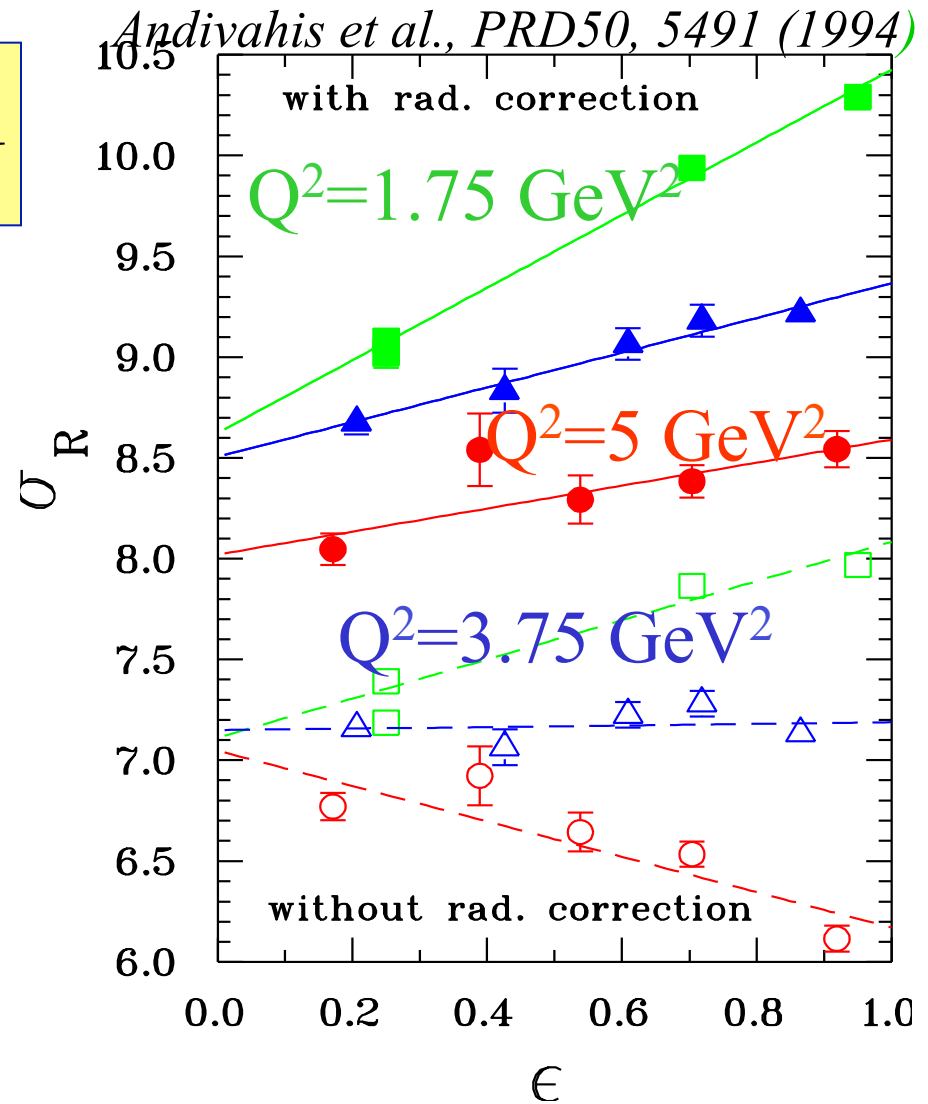
N.Kivel and M.Vanderhaeghen, Phys. Rev. Lett.103:092004 (2009). (high Q_2)

Radiative Corrections (ep)

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2$$

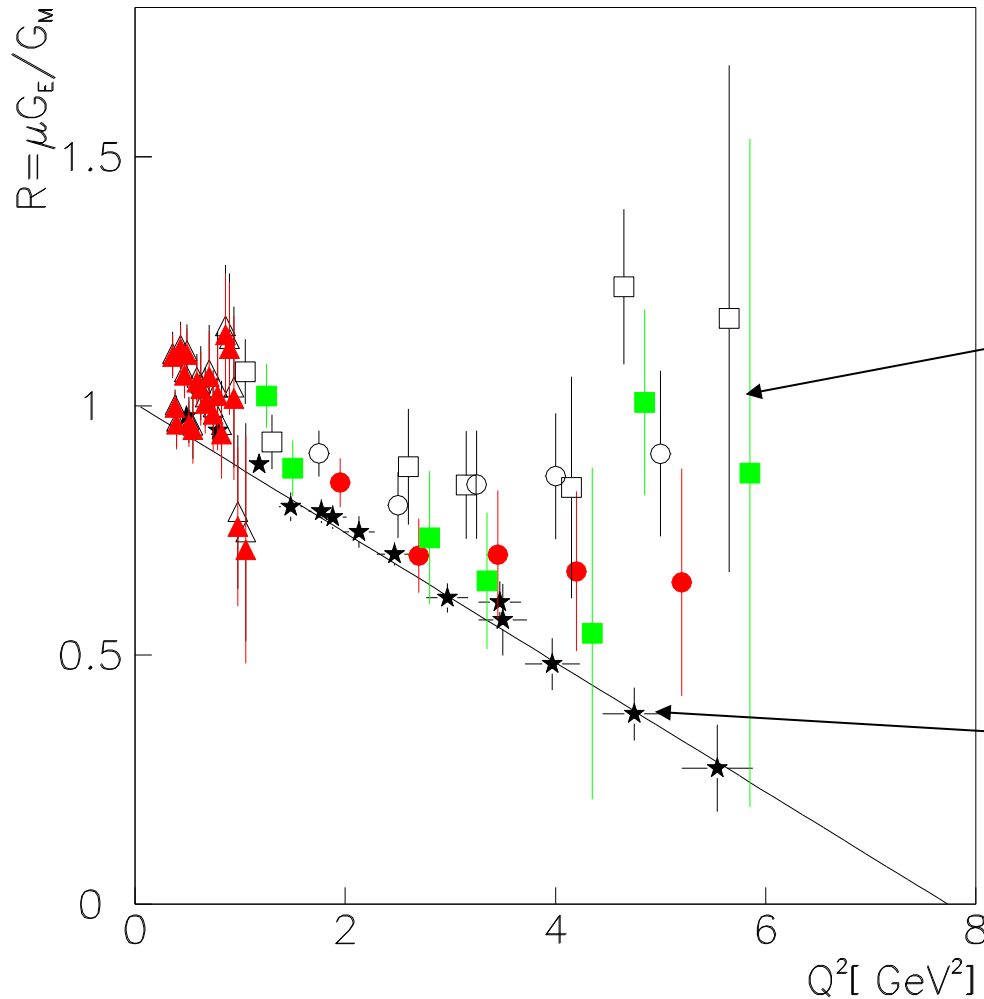
*May change
the slope of σ_R
(and even the sign !!!)*

*RC to the cross section:
- large (may reach 40%)
- ε and Q^2 dependent
- calculated at first order*



E. T.-G., G. Gakh, PRC 72, 015209 (2005)

Radiative Corrections (SF method)



Andivahis et al., PRD50, 5491 (1994)

SLAC data

SLAC data
corrected by SF

Jlab Polarization
data

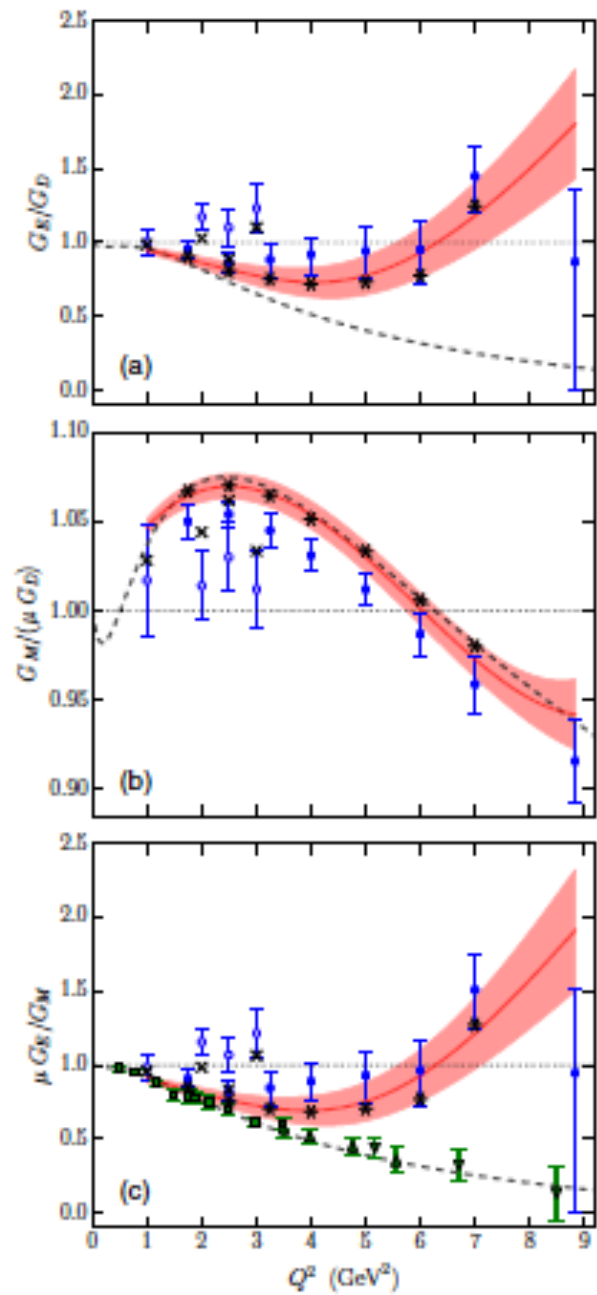
Yu. Bystricky, E.A.Kuraev, E. T.-G., Phys. Rev. C 75, 015207 (2007)

Reanalysis of Rosenbluth measurements of the proton form factors

A. V. Gramolin* and D. M. Nikolenko

Budker Institute of Nuclear Physics, 630090 Novosibirsk, Russia

(Received 28 March 2016; published 10 May 2016)



V. Fadin, R.E. Gerasimov

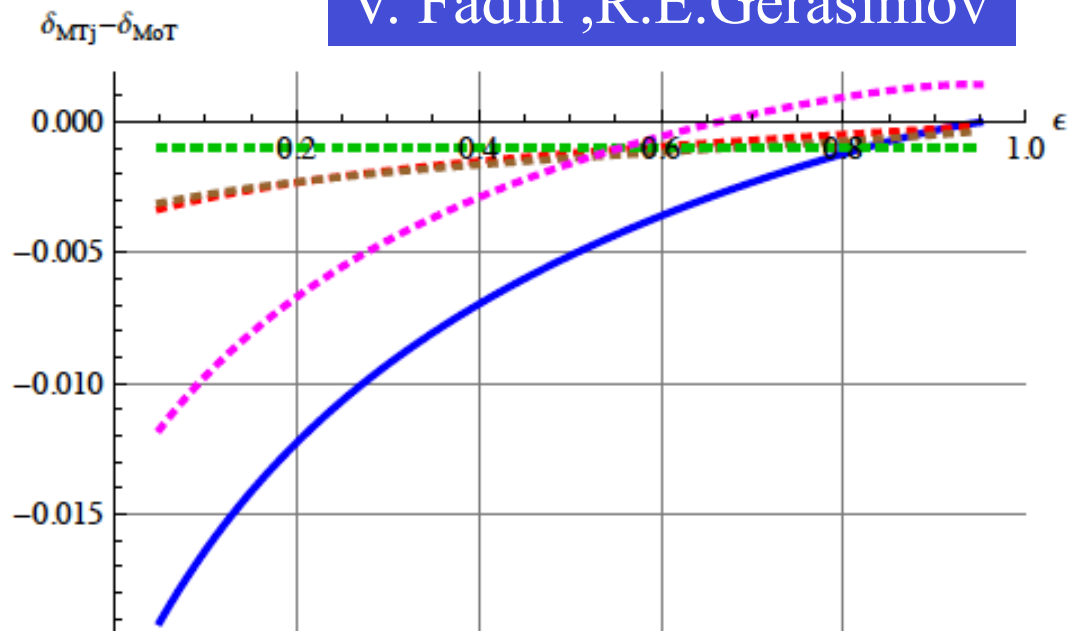
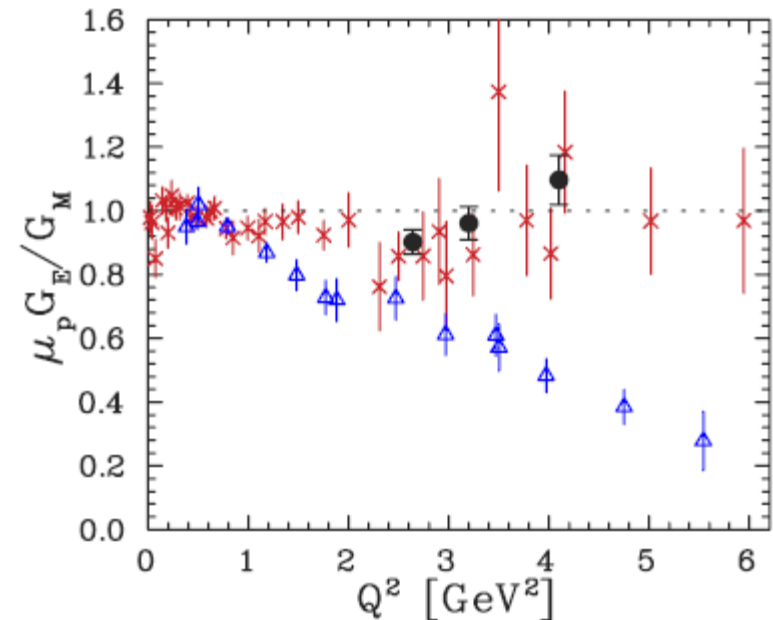
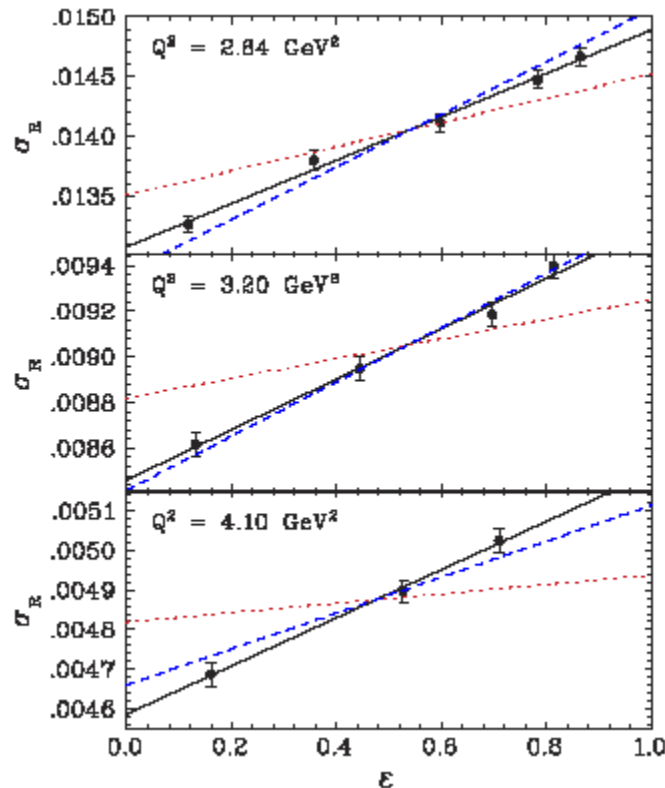


Figure 3: Difference at $Q^2 = 5 \text{ GeV}^2$.

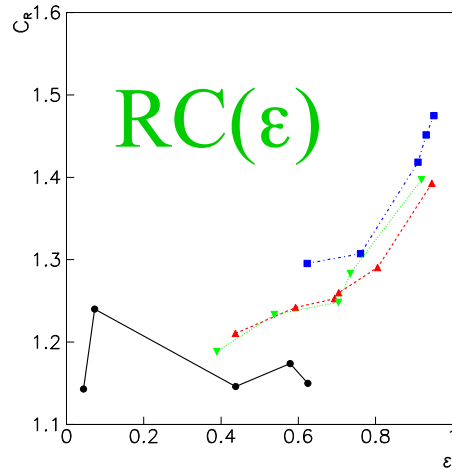
Precision Rosenbluth Measurement of the Proton Elastic Form Factors

I. A. Qattan,^{1,2} J. Arrington,² R. E. Segel,¹ X. Zheng,² K. Aniol,³ O. K. Baker,⁴ R. Beams,² E. J. Brash,⁵ J. Calarco,⁶ A. Camsonne,⁷ J.-P. Chen,⁸ M. E. Christy,⁴ D. Dutta,⁹ R. Ent,⁸ S. Frullani,¹⁰ D. Gaskell,¹¹ O. Gayou,¹² R. Gilman,^{13,8} C. Glashauser,¹³ K. Hafidi,² J.-O. Hansen,⁸ D. W. Higinbotham,⁸ W. Hinton,¹⁴ R. J. Holt,² G. M. Huber,⁵ H. Ibrahim,¹⁴ L. Jisonna,¹ M. K. Jones,⁸ C. E. Keppel,⁴ E. Kinney,¹¹ G. J. Kumbartzki,¹³ A. Lung,⁸ D. J. Margaziotis,³ K. McCormick,¹³ D. Meekins,⁸ R. Michaels,⁸ P. Monaghan,⁹ P. Moussiegt,¹⁵ L. Pentchev,¹² C. Perdrisat,¹² V. Punjabi,¹⁶ R. Ransome,¹³ J. Reinhold,¹⁷ B. Reitz,⁸ A. Saha,⁸ A. Sarty,¹⁸ E. C. Schulte,² K. Slifer,¹⁹ P. Solvignon,¹⁹ V. Sulkosky,¹² K. Wijesooriya,² and B. Zeidman²

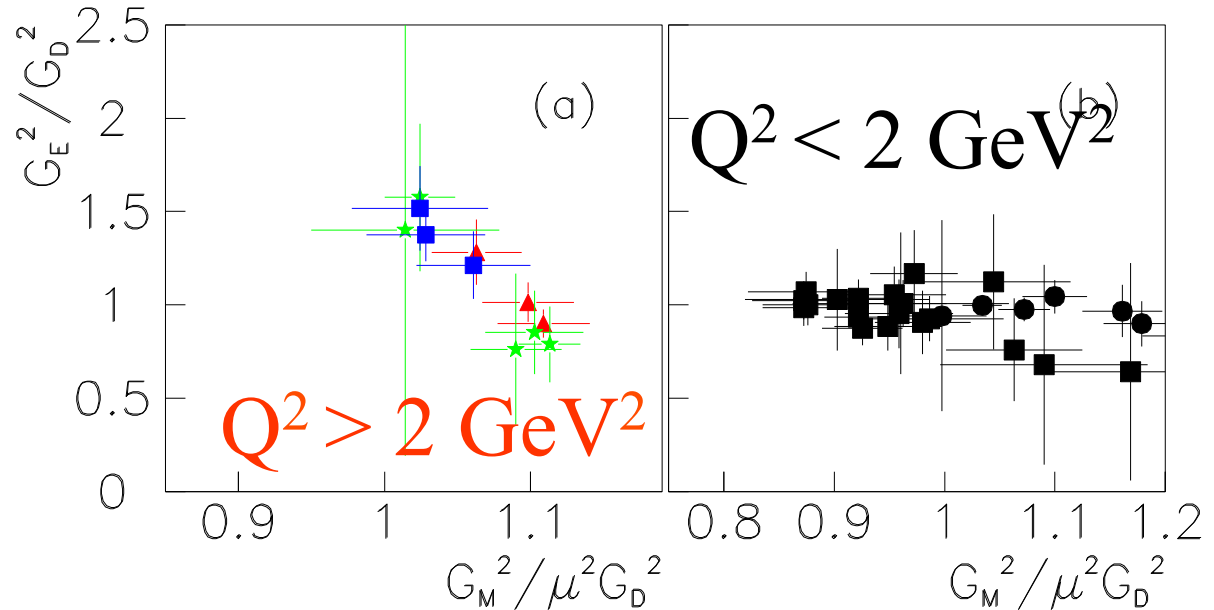


Other issues in data

- Correlations



G_E^2 versus G_M^2



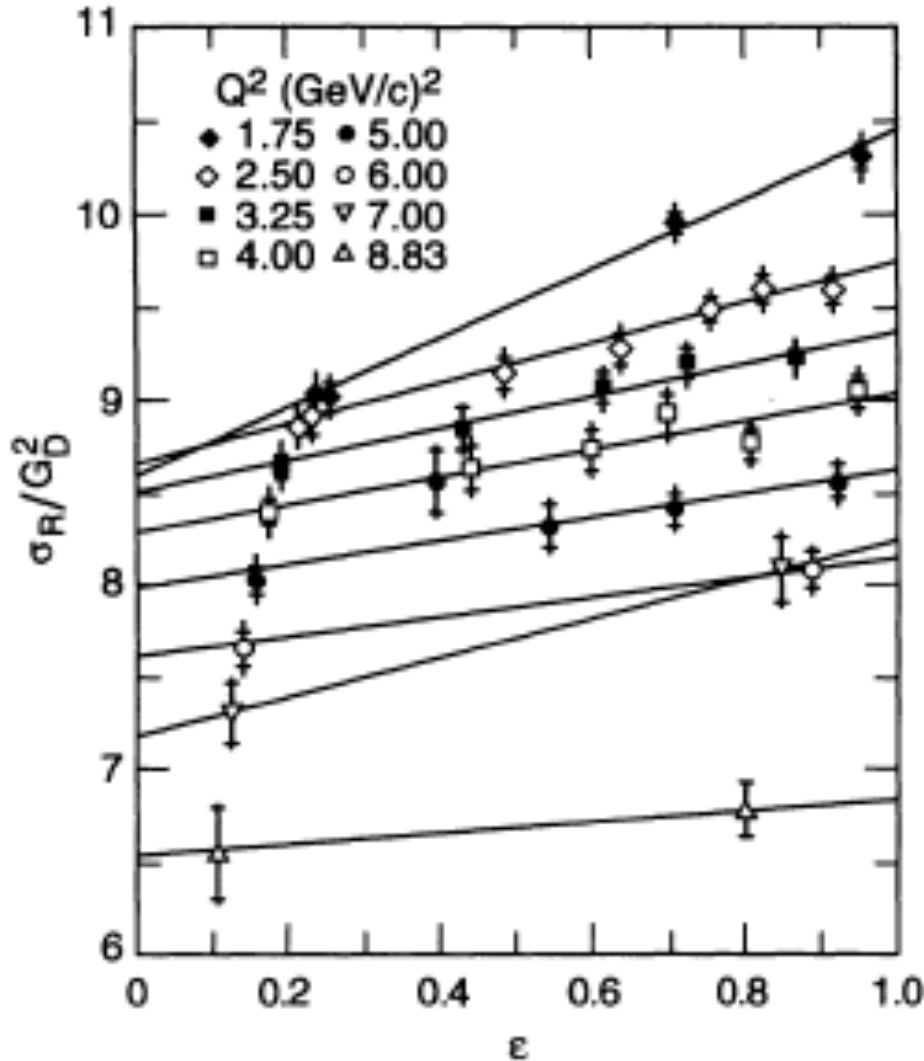
E.T-G, Phys. Part. Nucl. Lett. 4, 281 (2007)

- Normalizations

- of different sets of data
- within a set of data

Normalization

Andivahis et al., PRD50, 5491 (1994)



Two spectrometers
(8 and 1.6 GeV)

2 points at low ϵ

Fixed renormalization
for the lowest ϵ point
 $c=0.956$

(acceptance correction)

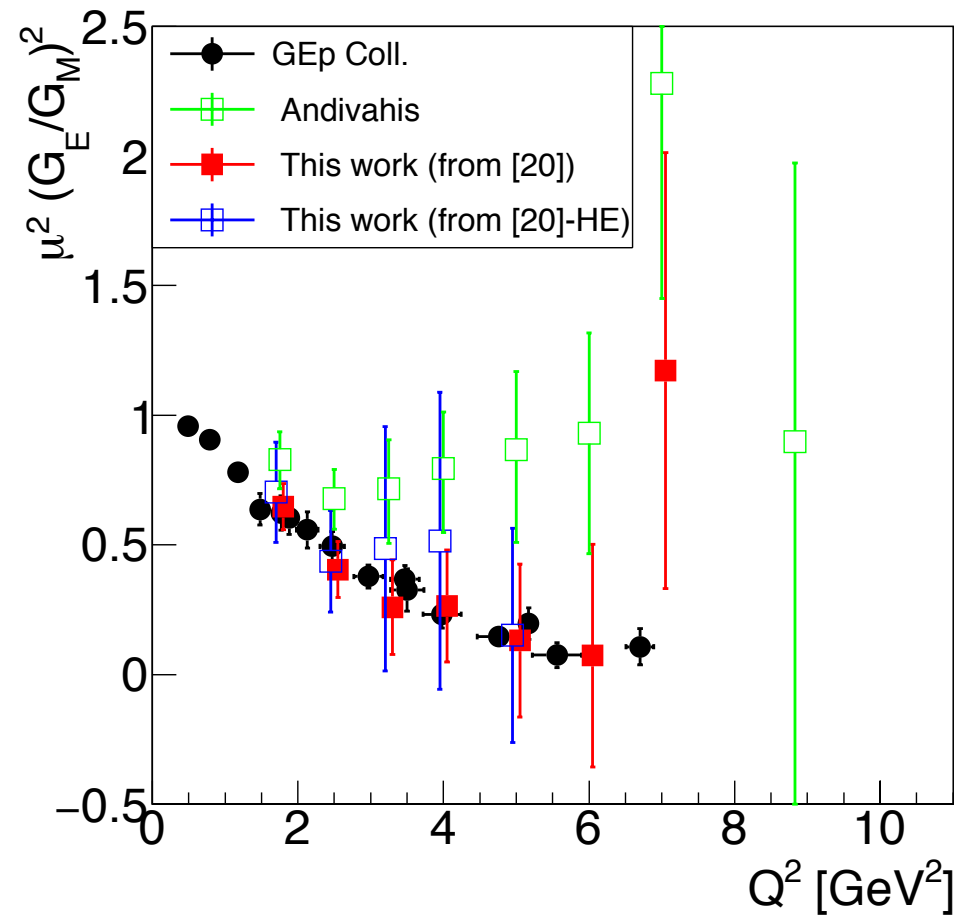
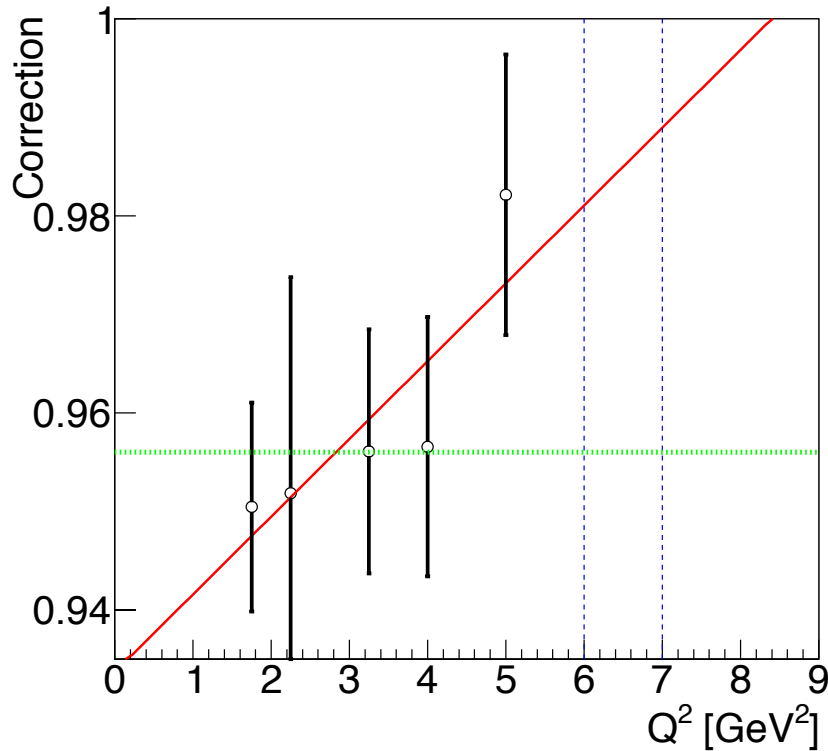
Increases the slope!

$$G_E \approx G_D$$

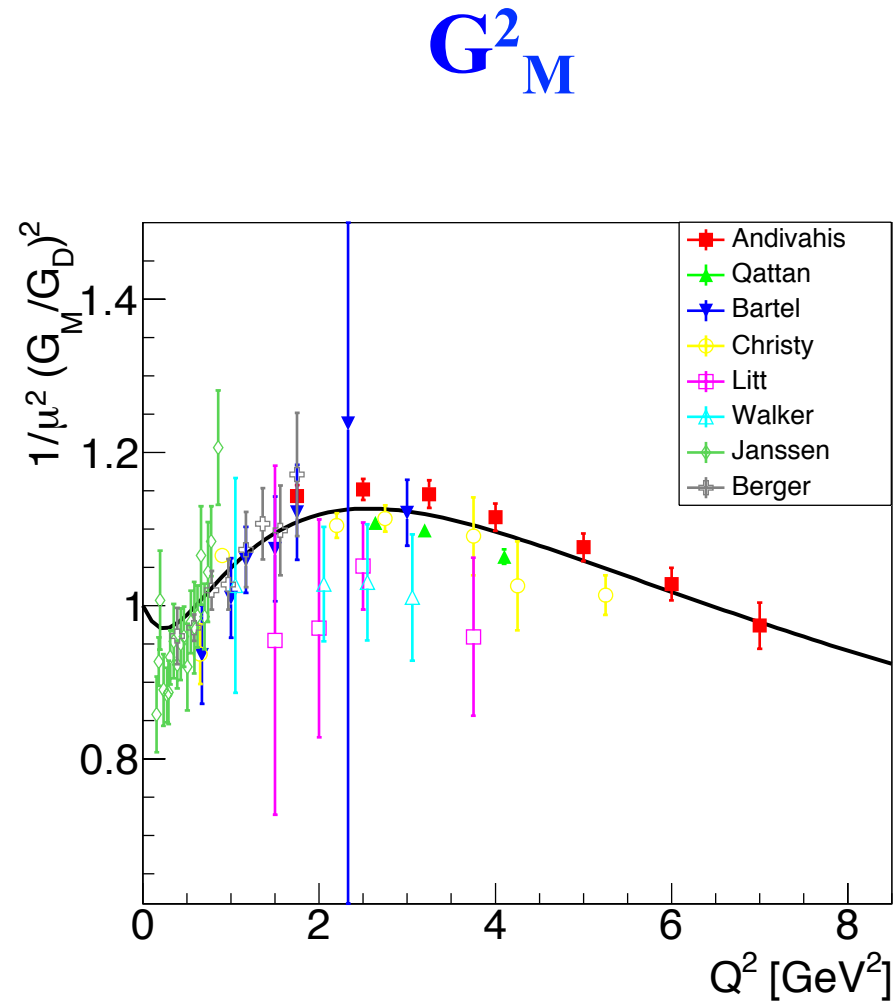
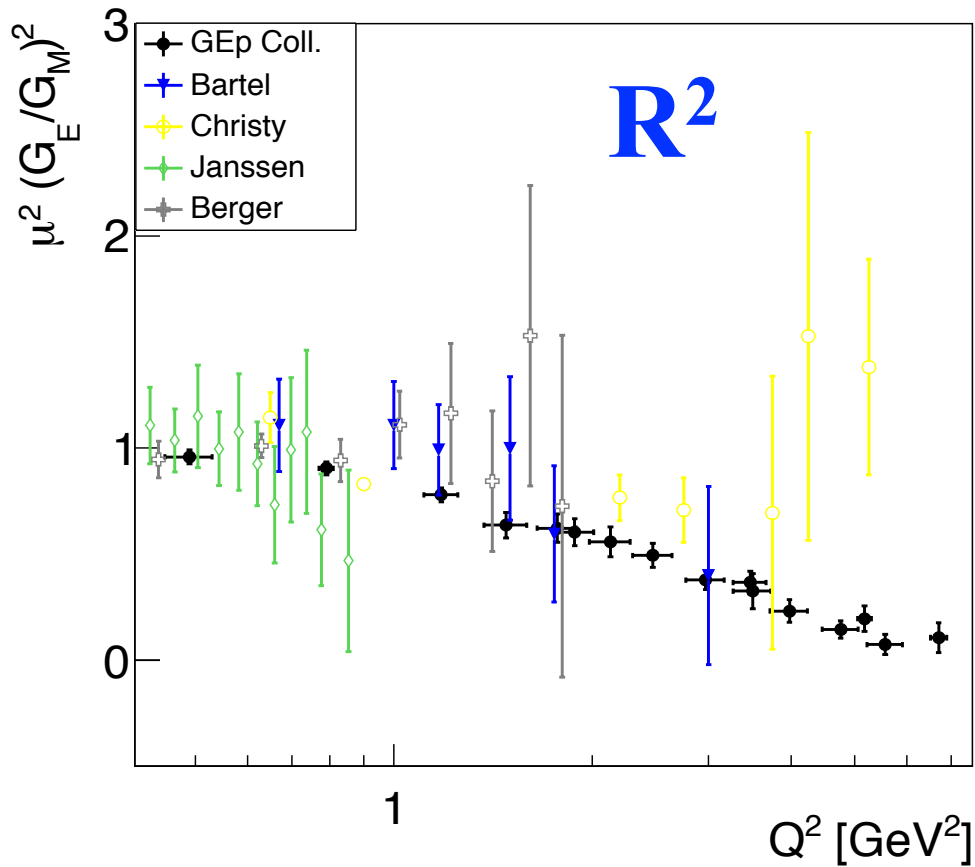
Direct extraction of the Ratio

Andivahis et al., PRD50, 5491 (1994)

$$\sigma_{\text{red}} = G_M^2 (R^2 \epsilon + \tau),$$



Different Data Sets



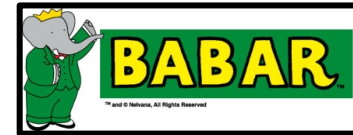
Conclusion - Discussion

- Large activity in Space and Time-like regions increase precision or extend q^2 range

Jefferson Lab

- Unified models in SL and TL

VEPP-3
Novosibirsk



To explore:

- Neutron/proton EM structure: FFs contain essential information in SL and TL
- Effect of deviation of both GE and GM from dipole
- If problems were not in observables... but in derivatives?