

Two scales in Bose-Einstein correlations

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BEC is the way to probe the size/structure of the domain which radiates secondaries.

Huge ATLAS statistics

($\sim 10^7$ events corresponding to $\sim 10^9$ same sign pairs)
allows the *detailed* study of the space-time structure of multiparticle production.

Consider the amplitude $M(r_1, r_2)$ in coordinate representation in which two identical pions, π_1 and π_2 , are emitted at the points r_1 and r_2 correspondingly. The coordinates r_1, r_2 can not be measured. So we have to take the Fourier transform

$$M_a(p_1, p_2) = \int \frac{d^4 r_1}{(2\pi)^4} e^{ip_1 r_1} \frac{d^4 r_2}{(2\pi)^4} e^{ip_2 r_2} M(r_1, r_2).$$

Besides this we have to consider the permutation of two identical pions. That is we have to add to M_a the amplitude

$$M_b(p_1, p_2) = \int \frac{d^4 r_1}{(2\pi)^4} e^{ip_2 r_1} \frac{d^4 r_2}{(2\pi)^4} e^{ip_1 r_2} M(r_1, r_2),$$

where the pion with momentum p_2 was emitted from the point r_1 and wise versa. This can be written as

$$M(p_1, p_2) = M_a + M_b = M_a \cdot (1 + e^{irQ}),$$

where the 4-vectors $r = r_1 - r_2$ and $Q = p_2 - p_1$ ^a.

Finally the cross section takes the form

$$\frac{E_1 E_2 d^2 \sigma}{d^3 p_1 d^3 p_2} = \frac{1}{2!} |M_a|^2 \langle 2 + 2e^{irQ} \rangle = |M_a|^2 (1 + \langle e^{irQ} \rangle).$$

Inclusive cross section for the two identical particles takes the form

$$\frac{E_1 E_2 d^2\sigma}{d^3p_1 d^3p_2} = \frac{1}{2!} |M|^2 \langle 2 + 2e^{irQ} \rangle = |M|^2 \langle 1 + e^{irQ} \rangle,$$

where M is the production amplitude, and $Q = p_2 - p_1$ and $r = r_1 - r_2$.

The proximity in phase space between final state particles with 4-momenta p_1 and p_2 can be quantified by

$$Q = \sqrt{-(p_1 - p_2)^2}.$$

The BEC effect is observed as an enhancement at low $Q \sim 300MeV$. To extract the effect one can compare measured Q-spectra with similar one but without BEC, with reference spectra. Then the ratio

$$R(Q) = \frac{\frac{dN}{dQ} - \frac{dN_{ref}}{dQ}}{\frac{dN_{ref}}{dQ}} \quad (1)$$

can be fitted with an appropriate formulae

$$R(Q) = \lambda F(rQ) + a + bQ$$

$F(rQ)$ is the Fourier transformation of the spatial distribution of the emission region with an effective size $\langle r \rangle$. The function simplest parametrization can be a linear exponent

$$SF1 = \lambda e^{-r1Q} + a + bQ$$

Parameter λ can be called as a strength of BEC, and a and b describe a simplest background to BEC.

Do we expect the homogeneous distribution? – **NO**
we expect few/many small size sources of secondaries
distributed over a larger area.

$$B_{el} = B_0 + 2\alpha' \ln(s/s_0)$$

$$B_0 \sim 10 \text{GeV}^{-2}, \quad 2\alpha' \sim 0.5 \text{GeV}^{-2}$$

Two different scales!

$\alpha' \implies$ transverse size of the Pomeron

. (BFKL/multiperipheral ladder)

$B_0 \implies$ size of beam/target hadrons.

Multipomeron exchange

(Multiple Interaction option in Monte Carlo)

is caused by the unitarity $2ImA(b) = |A(b)|^2 + G_{in}(b)$

The solution is

$$A(b) = i(1 - \exp(2i\delta_l)) = i(1 - \exp(-\Omega(b)/2))$$

• $(l = b\sqrt{s}/2)$

mean number of Pomerons, $N_P = \Omega(b)$; $\Omega_{LHC}(b=0) = 9$

N_P depends on b

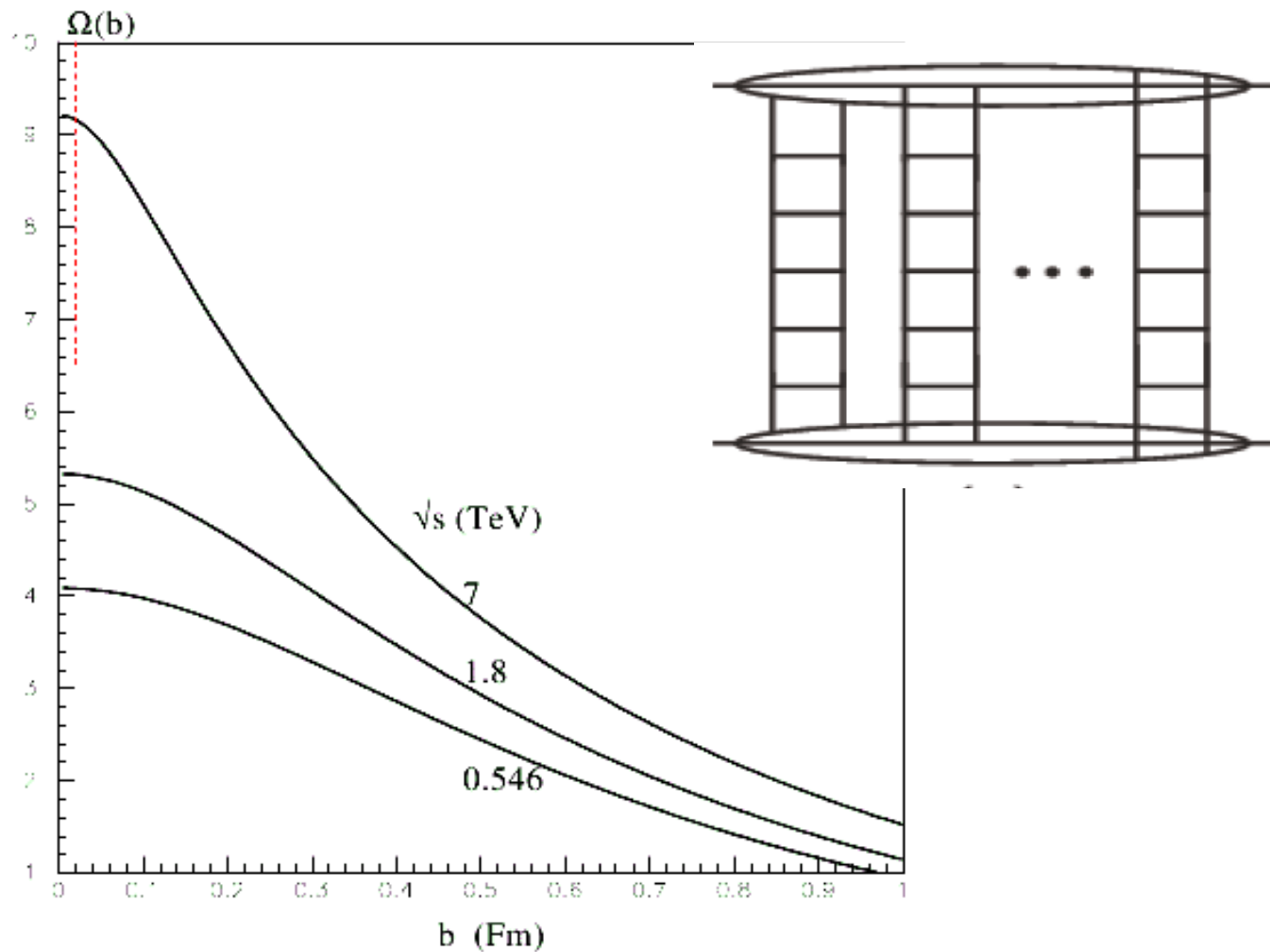


Figure 1: The proton opacity $\Omega(b)$ determined directly from the pp $d\sigma_{el}/dt$ data at 546 GeV [6], 1.8 TeV [7] and 7 TeV [2] data. The uncertainty on the LHC value at $b = 0$ is indicated by a dashed line.

Expect few small size sources
distributed over a much larger area
(overlap of beam and target hadrons)

***** Try to fit BEC with two radii *****

$$R(Q) = \lambda e^{-r_1 Q} + (1 - \lambda) e^{-r_2 Q}$$

Expectations: $r_1 \sim 2 - 3$ fm (corresponds to the whole B_{el})

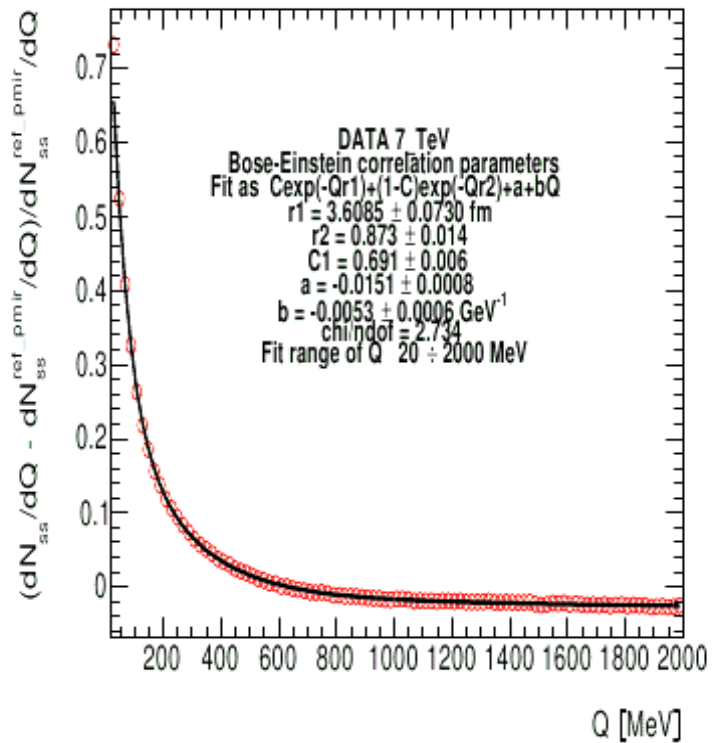
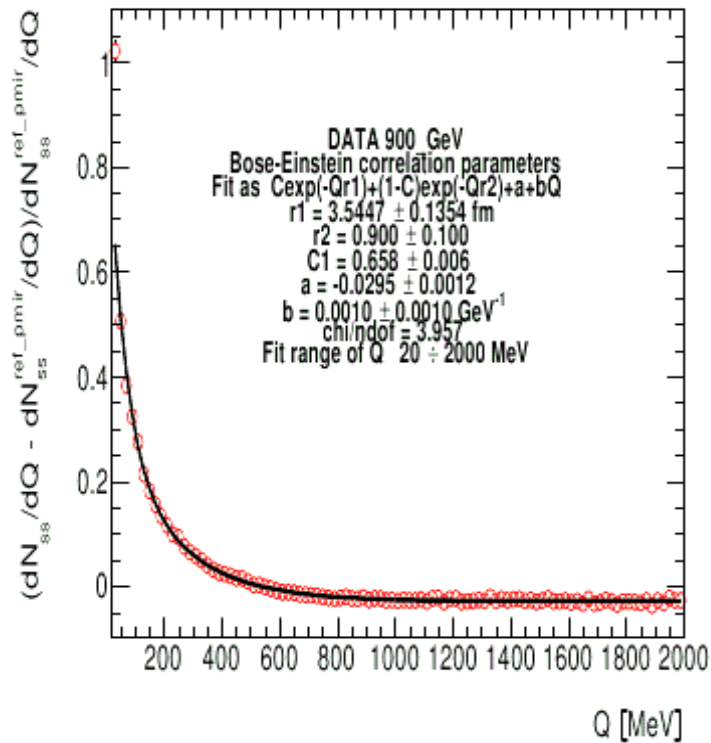
• $r_2 \sim 0.5 - 1$ fm (size of the Pomeron/pion)

r_1 and λ_1 increase with N_{ch}

$r_2 \sim const(N_{ch})$; do not depend on incoming particle

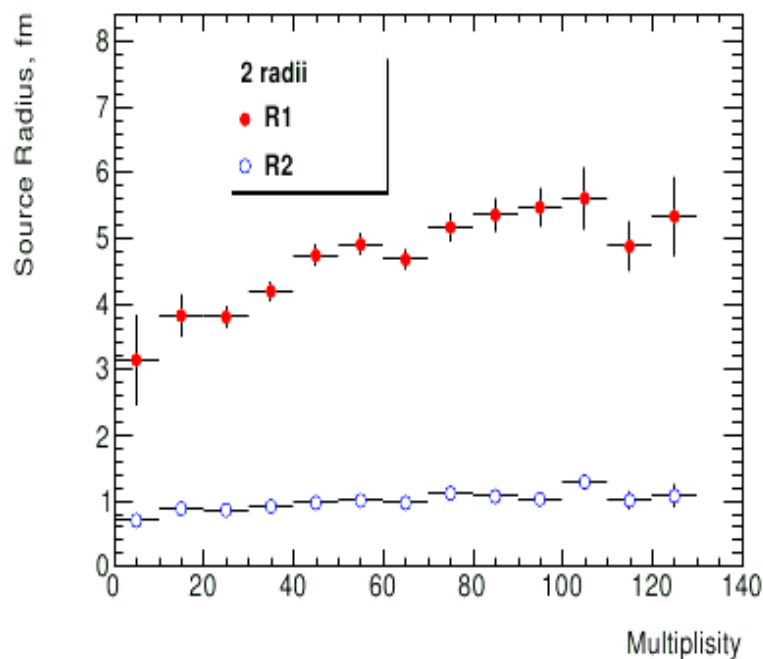
λ_2 close to 1 at low N_{ch} (one Pomeron); decreases with N_{ch}

Fit with two radii. An acceptable quality of a statistical description indicates that radii of a radiation zone is a weak (if any) function of hidden parameters.

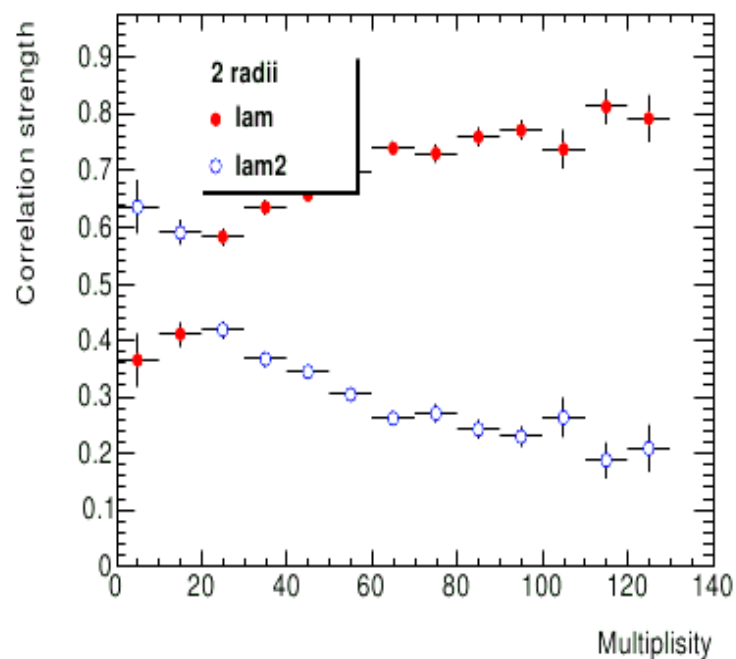


(a) 900 GeV data Q-distribution with two radii structure functions (b) 7TeV data Q-distribution with two radii structure function

In the following, only 7TeV pictures will be shown. 900 GeV statistics is too low. One can see that r_1 dependence on the multiplicity is rather smooth. Value r_2 is independent on the multiplicity. Correlation strength spectacular dependence on multiplicity "regulates" error bars in radii estimations.

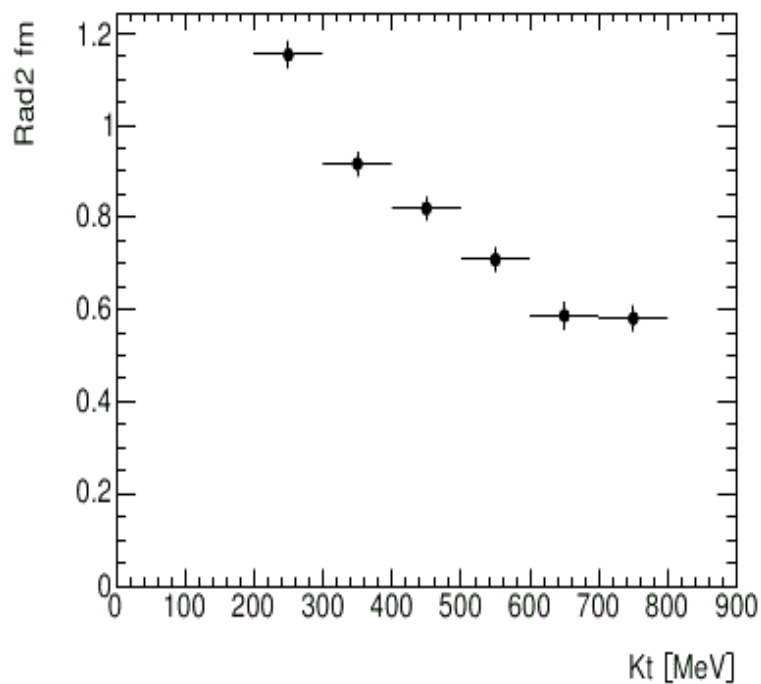


(a) Radiation source size as multiplicity

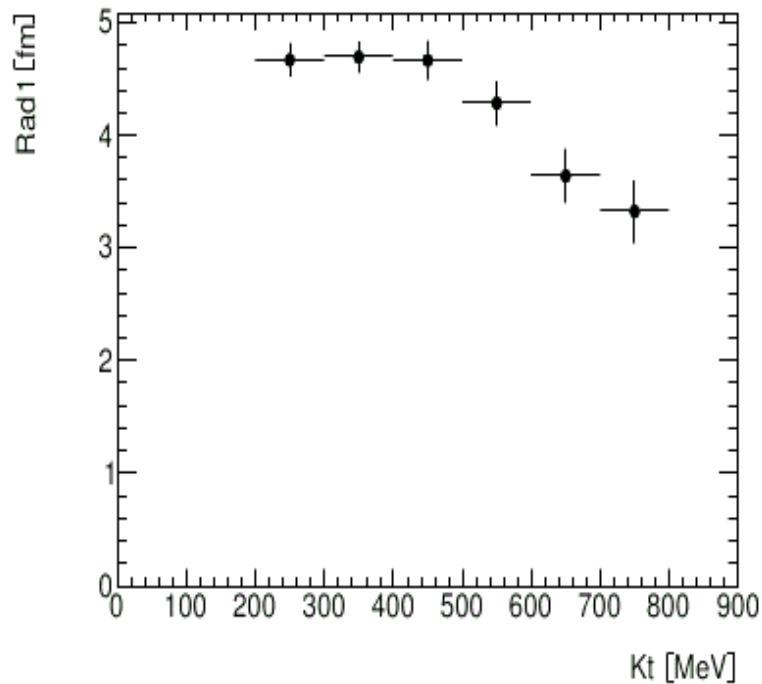


(b) Correlation strength

One may have a look radii dependencies on Kt - half of BEC pair transverse momentum - under suggestion that radii are independent of multiplicity



(a) r_2 Kt dependence



(b) r_1 Kt dependence

Conclusion:

It is First time observed that the secondaries are produced by a number of **SMALL size sources (hot spots)**

($r_2 \sim R_\pi$, $r_1 \sim (1 - 2)R_N$ is the distance between 'hot spots')

For a low particle density $dN_{ch}/d\eta < 10$ the radiation area is less than the proton size - (one hot spot)

At high $dN_{ch}/d\eta$ the size of radiation zone may be larger than that given by B_{el} due to the final state rescattering.

Radiative parton/Pomeron size decreases with
 $k_T = (p_{1T} + p_{2T})/2$ (Femtoscope resolution)
approaching to the pion radius ~ 0.6 fm

Due to small size (r_2) of radiative sources
the probability of this "hot spots" to overlap is
rather small even in nuclear-nuclear collisions
while the energy density inside each "hot spot" is large.

Outlook

It would be interesting:

to have particle identification

Are BEC between kaons and/or protons the same as that in pion case ?

to study BEC in events with Large Rapidity Gap

to study BEC in events with high E_T jet and/or W/Z bosons

Backup

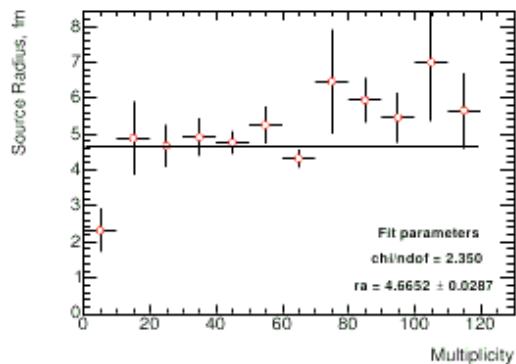
high $k_T = (p_{1T} + p_{2T})/2 \implies$ both pions from one minijet

r_2 component dominates

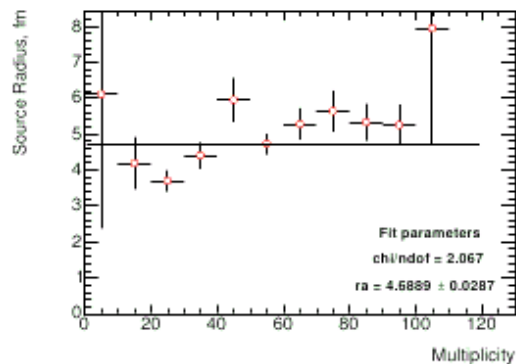
The structure may be changed by the final state interactions

IF so - r_2 = size of bubbles in QGLiquid

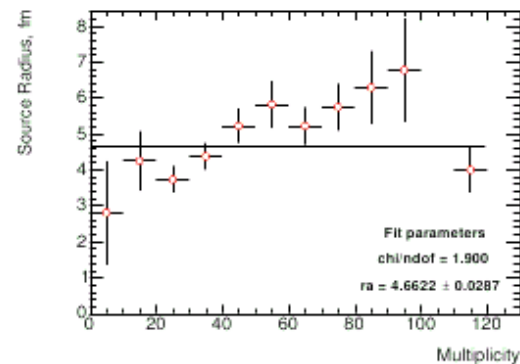
$$r_1 \propto N_{ch}^{1/3}$$



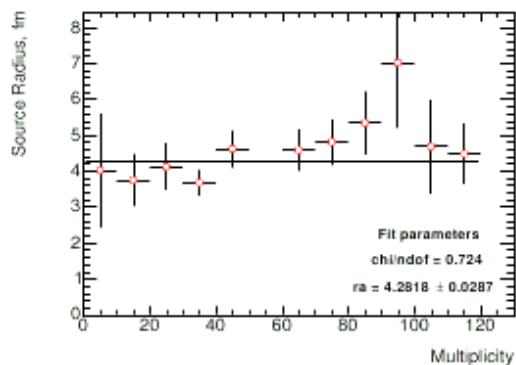
(a) r_1 as multiplicity with K_t in the interval (200,300) MeV



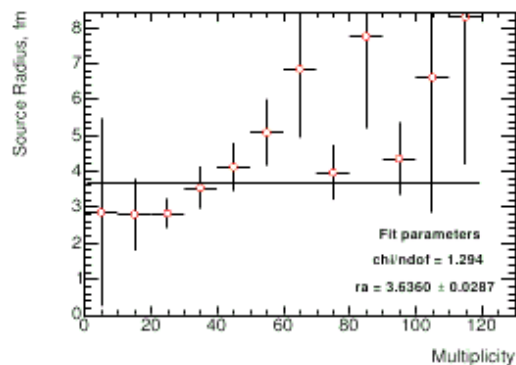
(b) r_1 as multiplicity with K_t in the interval (300,400) MeV



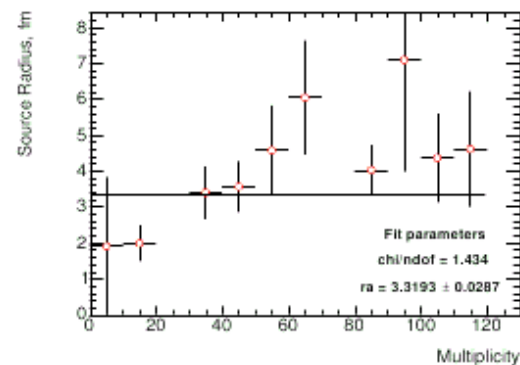
(c) r_1 as multiplicity with K_t in the interval (400,500) MeV



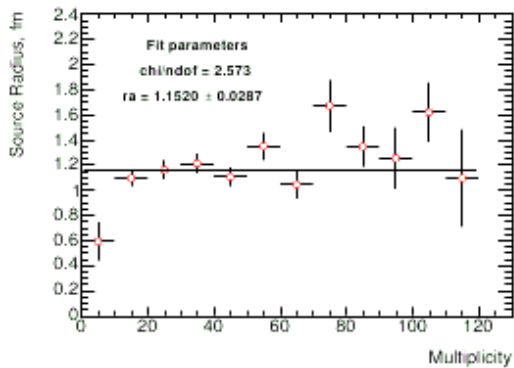
(d) r_1 as multiplicity with K_t in the interval (500,600) MeV



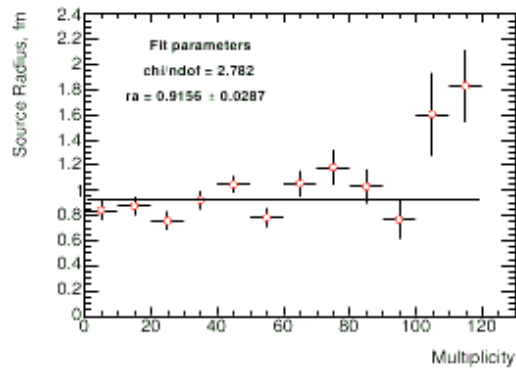
(e) r_1 as multiplicity with K_t in the interval (600,700) MeV



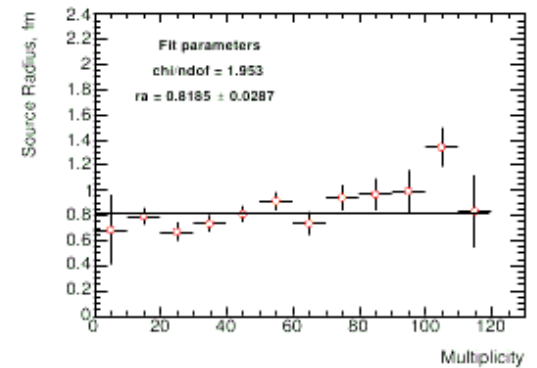
(f) r_1 as multiplicity with K_t in the interval (700,800) MeV



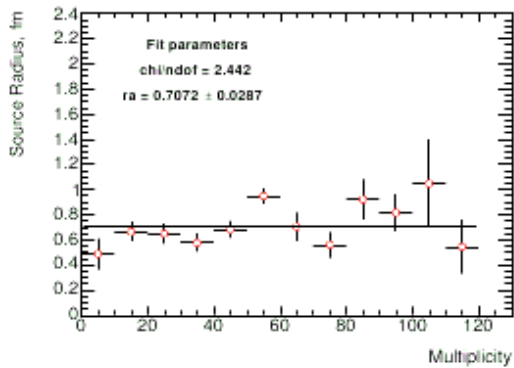
(a) r_2 as multiplicity with K_t in the interval (200,300) MeV



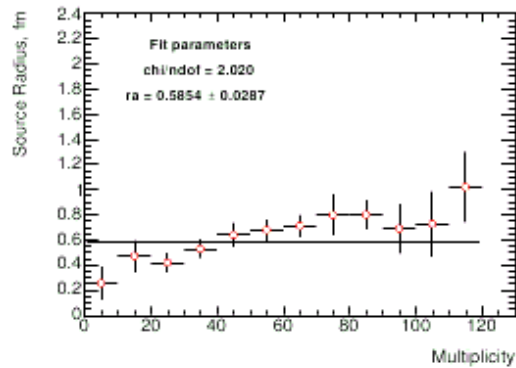
(b) r_2 as multiplicity with K_t in the interval (300,400) MeV



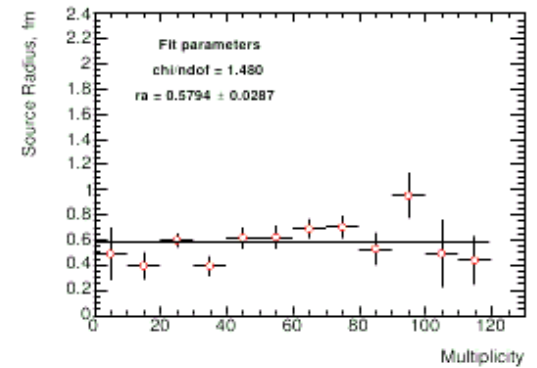
(c) r_2 as multiplicity with K_t in the interval (400,500) MeV



(d) r_2 as multiplicity with K_t in the interval (500,600) MeV



(e) r_2 as multiplicity with K_t in the interval (600,700) MeV



(f) r_2 as multiplicity with K_t in the interval (700,800) MeV

