Heavy Flavour Spectroscopy and Exotic States From Lattice QCD

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Implications of LHCb Measurements and Future Prospects CERN, 13 October 2016

Motivation

'A modern day November revolution'





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Lattice QCD in a Nutshell

$$L = \bar{\psi}(i\gamma^{\mu}D_{\mu} - m)\psi - \frac{1}{4}G^{a}_{\mu\nu}G^{\mu\nu}_{a}$$

- Gluons \longrightarrow SU(3) matrices ('links'): $U_{\mu}(x) = e^{-iagA^b_{\mu}(x)T^b}$
- Quarks live on 'sites' with colour, flavour and spinor indices
- **Derivatives** finite differences: $\nabla_{\mu}\psi(x) = \frac{1}{a}[U_{\mu}(x)\psi(x+a\hat{\mu})-\psi(x)]$
- Monte Carlo estimation of the path integral in a finite Euclidean space-time
- Measure desired observables . . .





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Spectroscopic observables continued . . .

So what can we do?









- Extensive calculations for mesons and baryons
- Access to different structures and states

'Limiting cases' 'Static' heavy-quarks, (p)NRQCD

Recent calculations of 4 and 5 quark states

'Límítíng cases' 'Static' heavy-quarks, (p)NRQCD

$$\langle W(r,t) \rangle = \left\langle 0 \left| \mathcal{Q}_r \ \mathcal{T}^{t/a} \ \mathcal{Q}_r^{\dagger} \right| 0 \right\rangle$$
Wilson loop
Wilson loop
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 $\frac{\text{Static potential}}{V_0(r) = -\lim_{t \to \infty} \frac{d}{dt} \langle W(r, t) \rangle}$

Does the static potential become more attractive in the presence of light hadrons?

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$$\begin{split} \Delta V_H(r) &= V_H - V_0 \\ &= -\lim_{t \to \infty} \frac{d}{dt} \ln \left[\frac{\langle H | \mathcal{Q}_r \ \mathcal{T}^{t/a} \ \mathcal{Q}_r^{\dagger} | H \rangle}{\langle 0 | \mathcal{Q}_r \ \mathcal{T}^{t/a} \ \mathcal{Q}_r^{\dagger} | 0 \rangle} \right] \\ &= -\lim_{t \to \infty} \frac{d}{dt} \ln \left[\frac{\langle W(r,t) \ C_H^{2pt}(t+2\delta t) \rangle}{\langle W(r,t) \rangle \ \langle C_H^{2pt}(t+2\delta t) \rangle} \right] \end{split}$$

High statistics:

Ensemble:

•

- 1552 configurations
- 12 time-sources

- $M_{\pi} \approx 220 \text{ MeV}$
 - $M_K \approx 480 \text{ MeV}$

-800

0.2

0.4

0.6

 $r \, [\mathrm{fm}]$

0.8

1

1.2

• 12 time-sources

 $M_K \approx 480 \text{ MeV}$

•

Similar effects for the π , K, ρ , K^* mesons

fit

0.1

0.2

0.3

0.4

 $r \, [\mathrm{fm}]$

0.5

0.6

0.7

-10^L0

Similar effects for all octet and decouplet baryons

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Modified potentials — Schrödinger equation:

Charmonium 1S, 1P and 2S states reduce in mass

by < 10 **MeV**

Tetra-quarks from NRQCD

$$C_{ij}(t) = \left\langle 0 \left| \mathcal{O}_i(t) \; \mathcal{O}_j^{\dagger}(0) \right| 0 \right\rangle = \sum_n A e^{-E_n t}$$

• A single 'meson-meson' type operator

• A single 'diquark-antidiquark' operator

'Single hadron spectroscopy'

Excitation spectrum of finite-volume energy eigenstates

Light quark mass dependence:

- M_{π} : 400 MeV \implies 240 MeV
- Small quantitative changes
- No qualitative changes!

The Charm Sector - Baryons

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'Lüscher formalism' $\det[t_{ij}^{-1}(E) + M_{ij}(E,L)] = 0$

Infinite-volume t-matrix —

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Lattice QCD spectrum - infinite-volume t-matrix

The bad news:

- N channels \longrightarrow N(N+1)/2 unknowns per energy!
- Under-constrained for N > 1

A work-around:

- Parametrise the t-matrix with a 'few' free parameters
- Use >> 'few' parameters to constrain the t-matrix as a function of energy

Infinite-volume t-matrix —

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- Preserve Unitarity
- Examine pole content of t-matrix

0.98

0.96

-4.6

-4.8

2523

2525

2527

2529

2525

2527

2529

-4.6

-4.8

2523

Poles on all sheets with $\text{Im}[k_{D\pi}] < 0$

DK Scattering

S-wave:

- Bound-state pole $\approx 2380 \text{MeV}$; $\approx 55 \text{ MeV}$ below DK threshold (at $M_{\pi} = 391 \text{ MeV}$)
- Expt: $D_{s0}^{*}(2317) = 2317.7 \pm 0.6 \text{ MeV}$; $\approx 45 \text{ MeV}$ below *DK* threshold

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- c.f: S-wave pole in the $D\pi$ channel $\approx 1 \text{MeV}$ below threshold

Summary and Outlook

Lattice QCD calculations now probing exotic states and structures with heavy quarks

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Thank you for your attention!

