

Aspen 2017 Winter Conference "From the LHC to Dark Matter and Beyond"

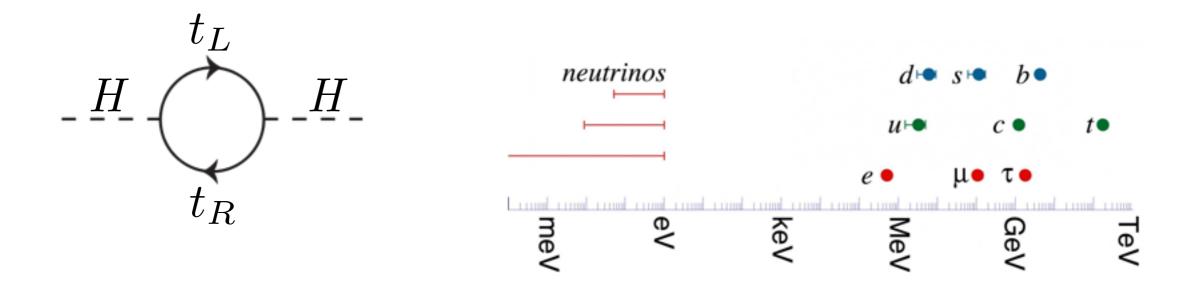
Disambiguating New Physics in the Top Sector

Jernej F. Kamenik



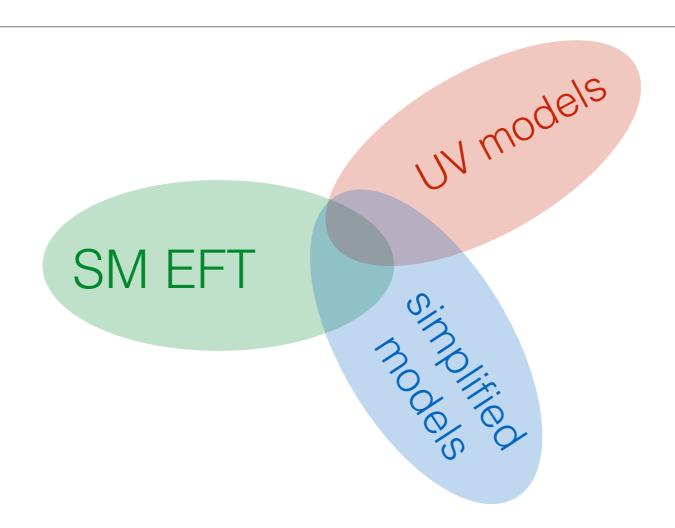
Top quark sector

Plays crucial role in EW hierarchy problem, SM flavor puzzle



- One of most promising experimental windows to NP
- LHC the only running experiment able to produce tops

Interpretation of NP searches at LHC



- Challenging use of EFT (energy gap, degeneracy)
- Limited scope/generality of simplified models
- Dependence on high scale parameters/assumptions in UV complete models

Two top examples

- 'top-philic' dark matter
- is tth really probing y_t?

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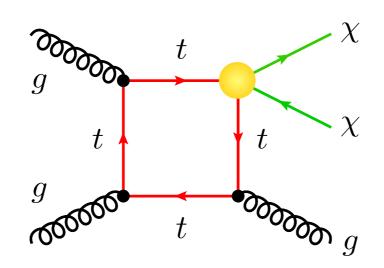
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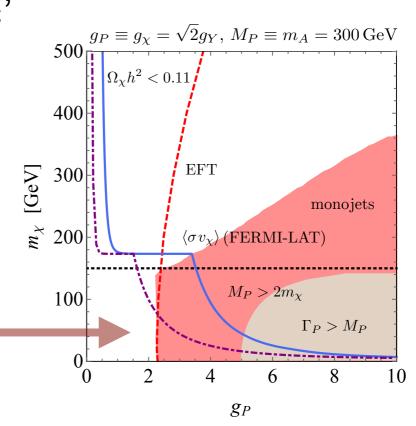
Simplified DM model with (pseudo)scalar mediator

$$\mathcal{L}_{DM} = i g_{\chi} A \bar{\chi} \gamma^{5} \chi + \sum_{f=q,\ell,\nu} i g_{f} A \bar{f} \gamma^{5} f$$

- Direct DM detection (spin, momentum) suppressed $\sigma_{\mathrm{SD}}^{N} \propto \frac{q^4}{m_N^4}$
- Naturally SM Yukawa-like couplings: $g_f = \sqrt{2} g_Y m_f/v$
 - DM becomes effectively 'top-philic'
 - → Reduced LHC mono-X sensitivity

 Haisch & Re, 1503.00691





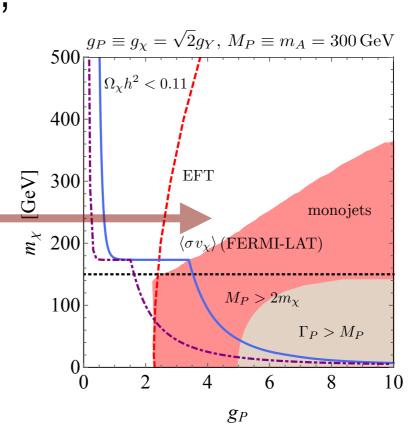
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- Suppressed missing E_T signals for

$$m_{\chi} > m_A/2$$



Bounding simplified NP with less assumptions?

LHC signatures based on resonant A production need to make several assumptions:

Non-resonant (off-shell) effects can be complementary:

- Well known (standard) in flavor physics see e.g. Bauer et al. 1701.07427
- At LHC used to bound Higgs width from $\sigma(pp \to 4\ell)$

Non-resonant NP in LHC 4t production

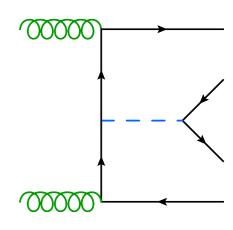
Alvarez, Faroughy, J.F.K., Morales & Szynkman, 1611.05032

In SM, LHC 4t production rare $\mathcal{O}(\alpha_S^4)$ QCD process:

see e.g. Bevilacqua & Worek, 1206.3064

$$\sigma^{NLO}(t \bar{t} t \bar{t}) = 12.32 \; \mathrm{fb}$$
 (LHC@13TeV using aMC@NLO)

- Off-shell $\mathcal{O}(\alpha_S^2 y_t^4)$ Higgs, $\mathcal{O}(\alpha_S^2 \alpha^2)$ Z contributions ~10% each!
- → Potential sensitivity to NP with O(1) couplings to tops
 - for masses below 2m_t effects off-shell
 - only depend on mass and coupling to tops
 - for masses above 2m_t resonant production
 - → reintroduce decay width dependence

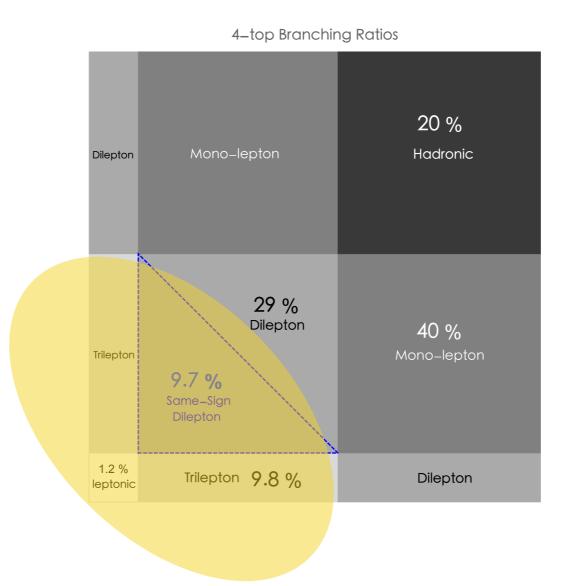


Gerbush et al., 0710.3133, Acharya et al., 0901.3367, Gregoire et al., 1101.1294, Liu & Mahbubani, 1511.09452 Gori et al., 1602.02782 Kim et al., 1604.07421

Non-resonant LHC 4t search strategy

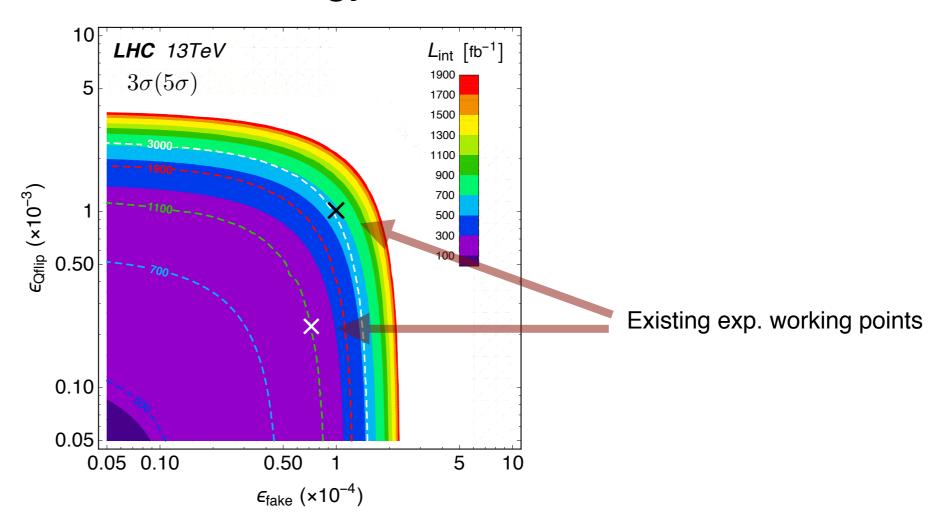
Challenging signatures (many jets, affecting lepton isolation)

clean s.s. dilepton & trilepton channels cover ~20% of total



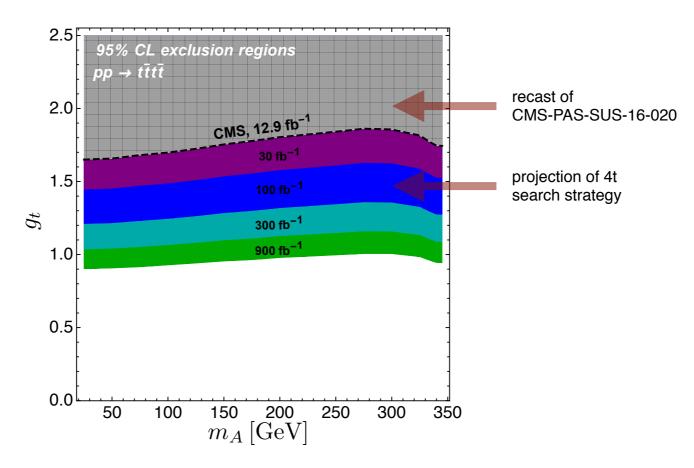
Separate strategies for s.s. dilepton & trilepton channels

- dominant backgrounds due to charge flip & jet-lepton mis-id
- cut-based, data-driven strategy



→ observation of SM rate possible in LHC Run II

- Existing (SUSY optimized) search combining both channels by CMS using early 13TeV data CMS-PAS-SUS-16-020
 - → already non-trivial bounds on (light) top-philic NP



• More general compared to resonant constraints ($\gamma\gamma$).

Two top examples

- 'top-philic' dark matter
- → is tth really probing y_t?

Mass generation in SM through Higgs mechanism

⇒ Higgs has hierarchical couplings to fermions

$$y_f^{\rm SM} = \sqrt{2}m_f/v$$

How well have we tested this?

A. Dery et al., 1302.3229

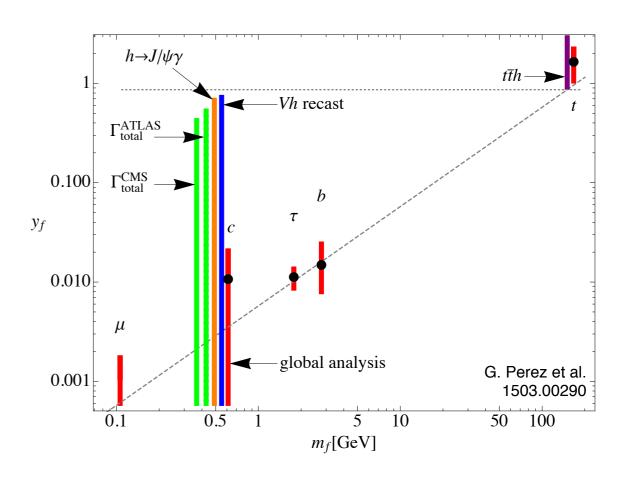
- proportionality $y_{ii} \propto m_i$
- factor of proportionality

$$y_{ii}/m_i = \sqrt{2}/v$$

diagonality

$$y_{ij} = 0, \quad i \neq j$$

G. Blankenburg et al., 1202.5704, R. Harnik et al., 1209.1397, ...



Many recent proposals...

1306.5770, 1406.1722, 1503.04830, 1505.03870, 1606.09621

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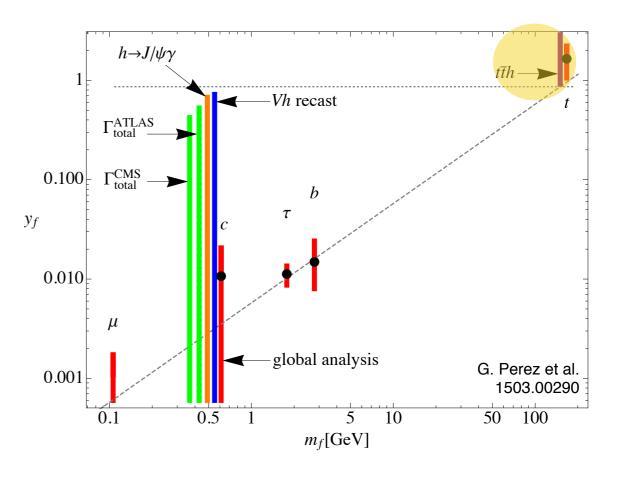
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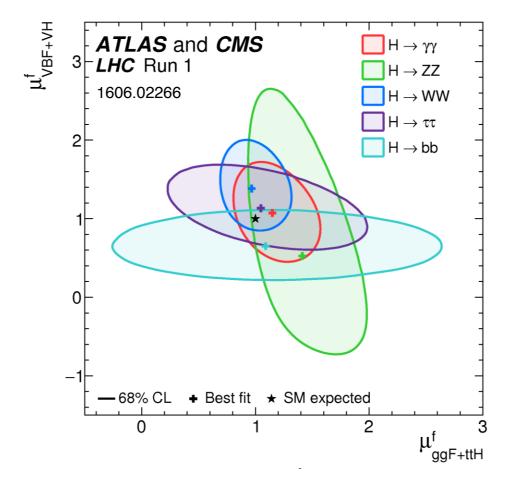
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LHC bounds on top Yukawa

Current global Higgs fits extract bound on y_t (mostly) from its parametric dependence of $\sigma(pp \to h)_{\rm incl.}$, $\sigma(pp \to t\bar{t}h)$

$$\mu \equiv \frac{\sigma_{\rm obs.}}{\sigma_{\rm SM}}$$



Anomalous top Yukawa in EFT

Non-standard value of yt effectively due to higher-dim operator

$$\mathcal{L}_{EFT} = \mathcal{L}_{SM} + \sum_{i} \frac{C_i}{\Lambda^2} O_i + \mathcal{O}(\Lambda^{-4}) + h.c.,$$

Several operators can contribute to relevant observables

$$O_{t\phi} = y_t^3 \left(\phi^{\dagger} \phi \right) \left(\bar{Q} t \right) \tilde{\phi} ,$$

$$O_{\phi G} = y_t^2 \left(\phi^{\dagger} \phi \right) G_{\mu\nu}^A G^{A\mu\nu} \,,$$

$$O_{tG} = y_t g_s(\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{\phi}G^A_{\mu\nu}.$$

$$\delta y_t \equiv \frac{C_{t\phi} v^2}{2\sqrt{2}\Lambda^2}$$

 $O_{t\phi} = y_t^3 \left(\phi^\dagger \phi\right) \left(\bar{Q}t\right) \tilde{\phi}$, modified y_t: $\delta y_t \equiv \frac{C_{t\phi} v^2}{2\sqrt{2}\Lambda^2}$ $O_{\phi G} = y_t^2 \left(\phi^\dagger \phi\right) G_{\mu\nu}^A G^{A\mu\nu}$, point-like contribution to hGG $O_{tG} = y_t g_s (\bar{Q}\sigma^{\mu\nu}T^At) \tilde{\phi} G_{\mu\nu}^A$. (can compensate y_t effects in $\sigma_{\rm GGF}$)

Chromo MDM of top, also affects tt production

Anomalous top Yukawa in EFT

Non-standard value of y_t effectively due to higher-dim operator

 $\mathcal{L}_{\mathrm{EFT}} = \mathcal{L}_{\mathrm{SM}} + \sum_{i} \frac{C_i}{\Lambda^2} O_i + \mathcal{O}(\Lambda^{-4}) + h.c.,$

see e.g. Zhang & Willenbrock, 1008.3869

Several operators can contribute to relevant observables

$$O_{t\phi} = y_{t}^{3} \left(\phi^{\dagger}\phi\right) (\bar{Q}t) \tilde{\phi}, \qquad O_{qq}^{(8,1)} = \frac{1}{4} (\bar{q}^{i} \gamma_{\mu} \lambda^{A} q^{j}) (\bar{q} \gamma^{\mu} \lambda^{A} q) \qquad O_{qq}^{(8,3)} = \frac{1}{4} (\bar{q}^{i} \gamma_{\mu} \tau^{I} \lambda^{A} q^{j}) (\bar{q} \gamma^{\mu} \tau^{I} \lambda^{A} q) \\ O_{\phi G} = y_{t}^{2} \left(\phi^{\dagger}\phi\right) G_{\mu\nu}^{A} G^{A\mu\nu}, \qquad O_{ut}^{(8)} = \frac{1}{4} (\bar{u}^{i} \gamma_{\mu} \lambda^{A} u^{j}) (\bar{t} \gamma^{\mu} \lambda^{A} t) \qquad O_{dt}^{(8)} = \frac{1}{4} (\bar{d}^{i} \gamma_{\mu} \lambda^{A} d^{j}) (\bar{t} \gamma^{\mu} \lambda^{A} t) \\ O_{tG} = y_{t} g_{s} (\bar{Q} \sigma^{\mu\nu} T^{A} t) \tilde{\phi} G_{\mu\nu}^{A}. \qquad O_{qt}^{(1)} = (\bar{q}^{i} t) (\bar{t} q^{j}) \qquad \qquad O_{qd}^{(1)} = (\bar{q} d^{i}) (\bar{d}^{j} q) \\ O_{tG} = y_{t} g_{s} (\bar{Q} \sigma^{\mu\nu} T^{A} t) \tilde{\phi} G_{\mu\nu}^{A}. \qquad O_{qt}^{(1)} = (\bar{q}^{i} t) (\bar{t} q^{j}) \qquad \qquad \text{(assume CP, ~U(2)F)}$$

4fermion operators also enter tt and tth production in fixed

 $C_d^2 = C_{ad}^{(1)} + C_{at}^{(1)}$

Combinations
$$C_u^1 = C_{qq}^{(8,1)} + C_{qq}^{(8,3)} + C_{ut}^{(8)}$$
 $C_u^2 = C_{qu}^{(1)} + C_{qt}^{(1)}$ $C_d^1 = C_{qq}^{(8,1)} - C_{qq}^{(8,3)} + C_{dt}^{(8)}$

known at NLO in QCD

Maltoni et al., 1607.05330

linearized (EFT truncated) effects

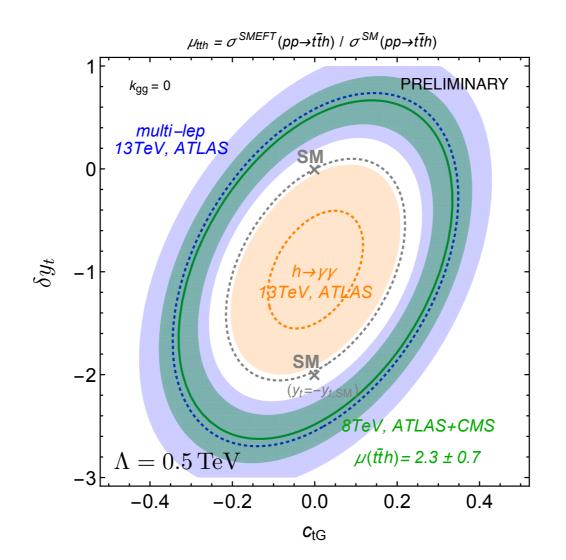
4fermion & CMDM effects in inclusive and boosted tt

see also Aguilar-Saavedra et al., 1412.6654 13 TeV ATLAS already systematics **PRELIMINARY** Inclusive tt 0.6 & theory limited 0.4 1606.02699 0.2 boosted regime $C_4 \equiv 4(C_u^{(2)} + C_d^{(2)})$ 3 Boosted tt more sensitive to 0.0 4f compared to -0.2**CMDM** ATLAS-CONF-2016-100 -0.4 $\Lambda = 0.5 \, \mathrm{TeV}$ -0.60.0 0.2 0.4 0.6 8.0 1.0 see e.g. Contino et al., 1604.06444

Disentangling NP effects in tt & tth

Competitive effects in tth due to yt and CMDM

- 4fermion operator effects already largely negligible
- current measurements allow for sizable degeneracy

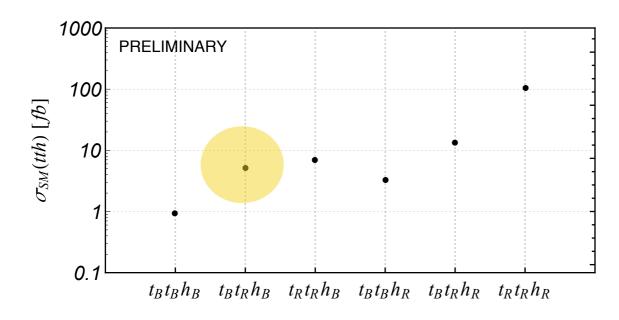


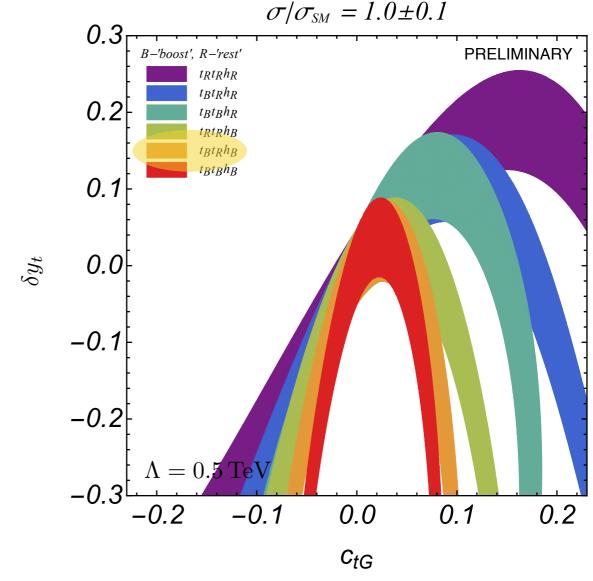
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Resolve degeneracy using boosted regime

t_B - top-tagged R=1.5 jet h_B - Higgs-tagged R=1.2 jet

t_R,h_R - resolved decays

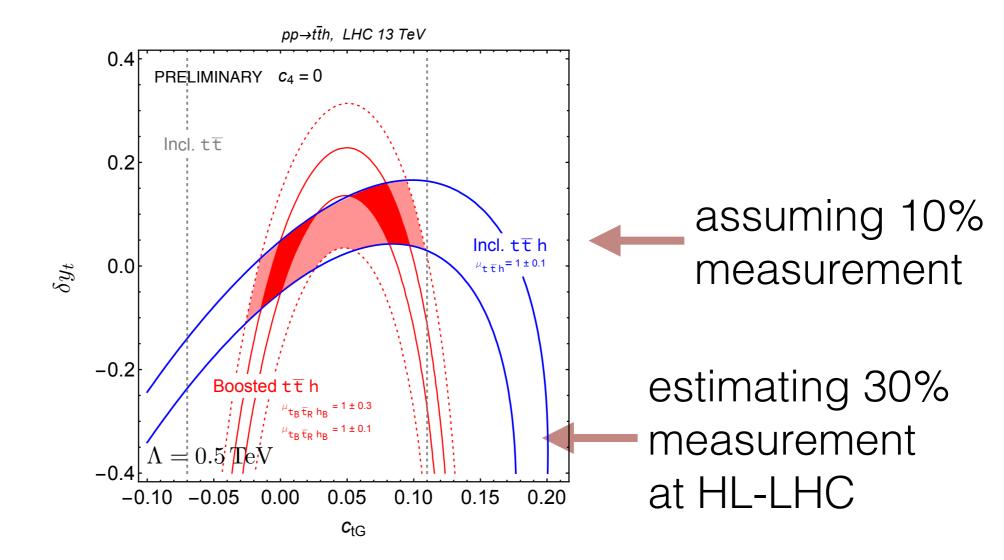




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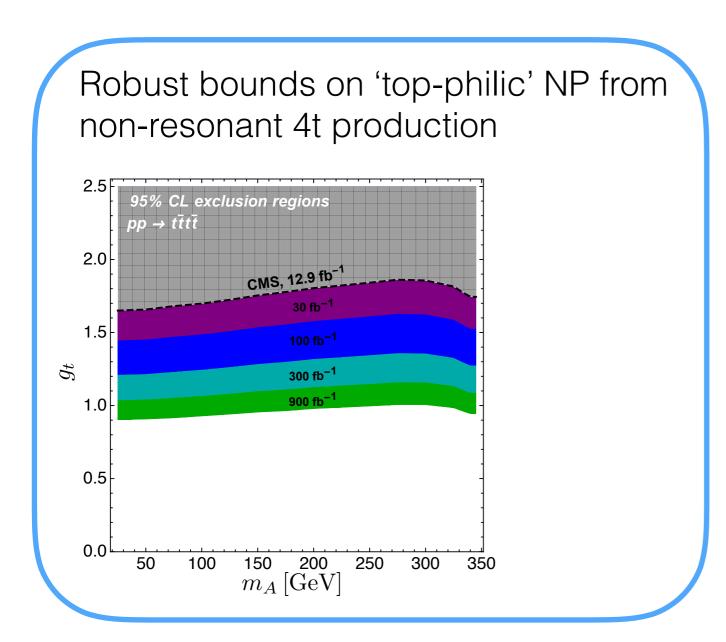
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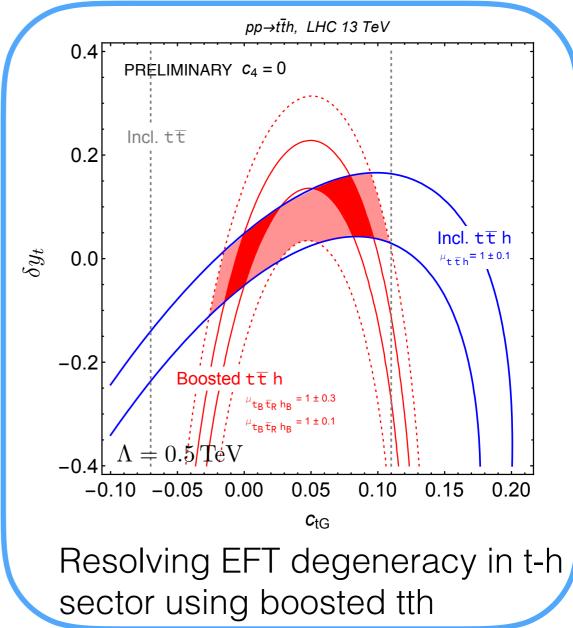
Resolve degeneracy using boosted regime: projections



Conclusions

Examples extending NP reach in the top sector in light of large (HL-)LHC projected statistics



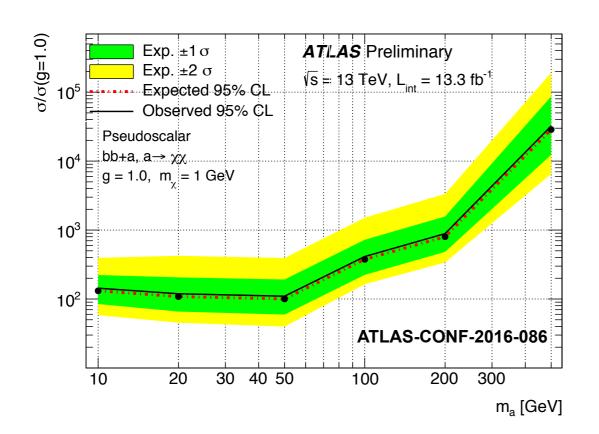


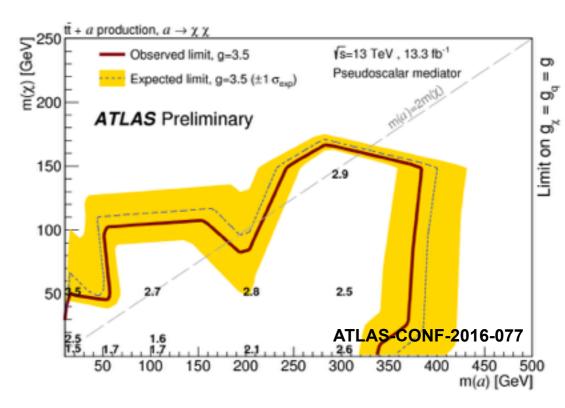
Additional material

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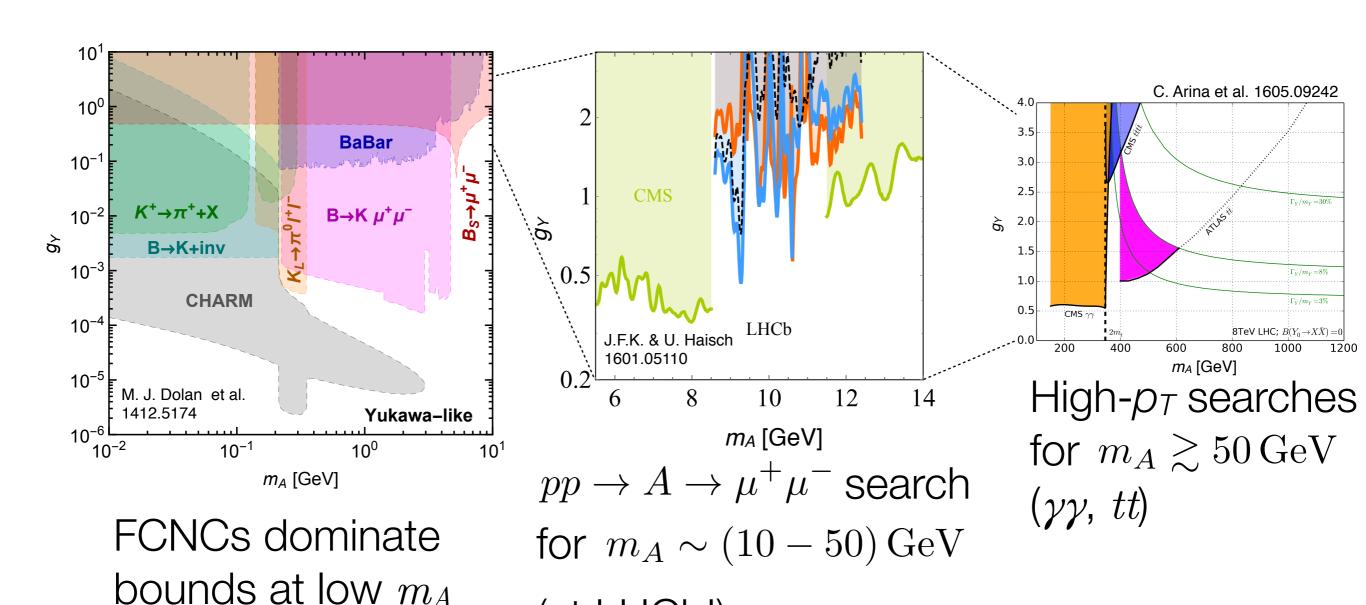
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 - Reduced LHC mono-X sensitivity (better than bbE_T, ttE_T)





Complementary flavor & high-p_T probes

 $A \to f \bar{f}, gg, \gamma \gamma$ dominate, can be probed in different regimes



see also P. Ilten et al. 1603.08926

(at LHCb!)

Disentangling NP effects in tt & tth

Competitive effects in tth due to y_t and CMDM Resolve degeneracy using boosted regime

