

From Dark Particle Physics to the Matter Distribution of the Universe

Aspen Particle Physics conference, March 21 2017

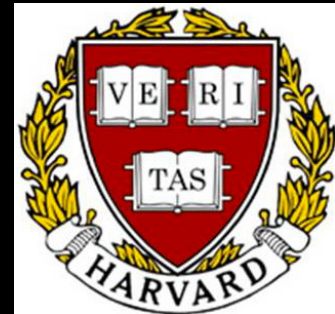
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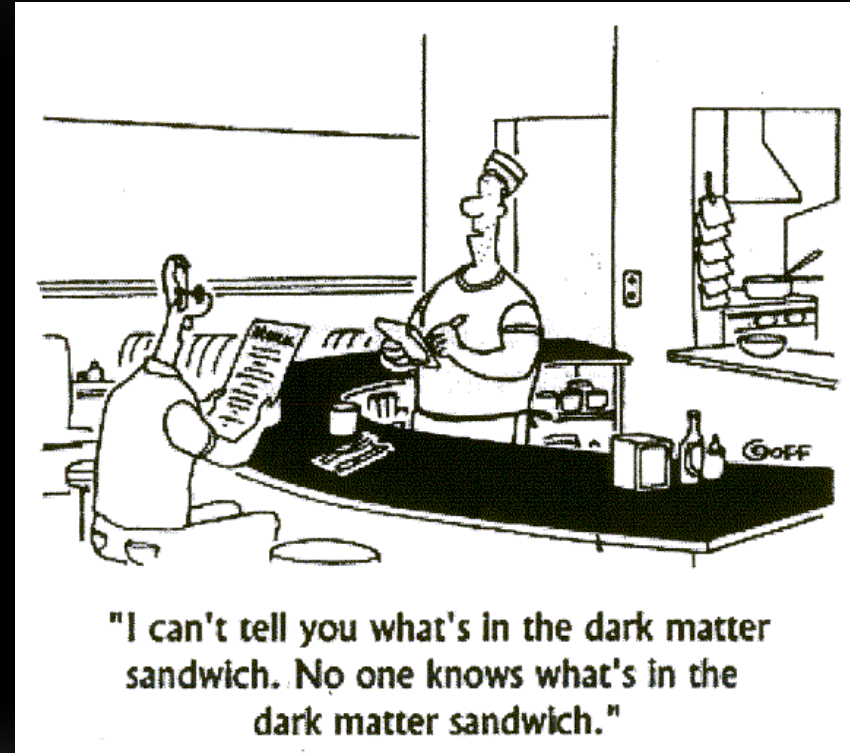
With contributions from:

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Kris Sigurdson, Torsten Bringmann,
Christoph Pfrommer

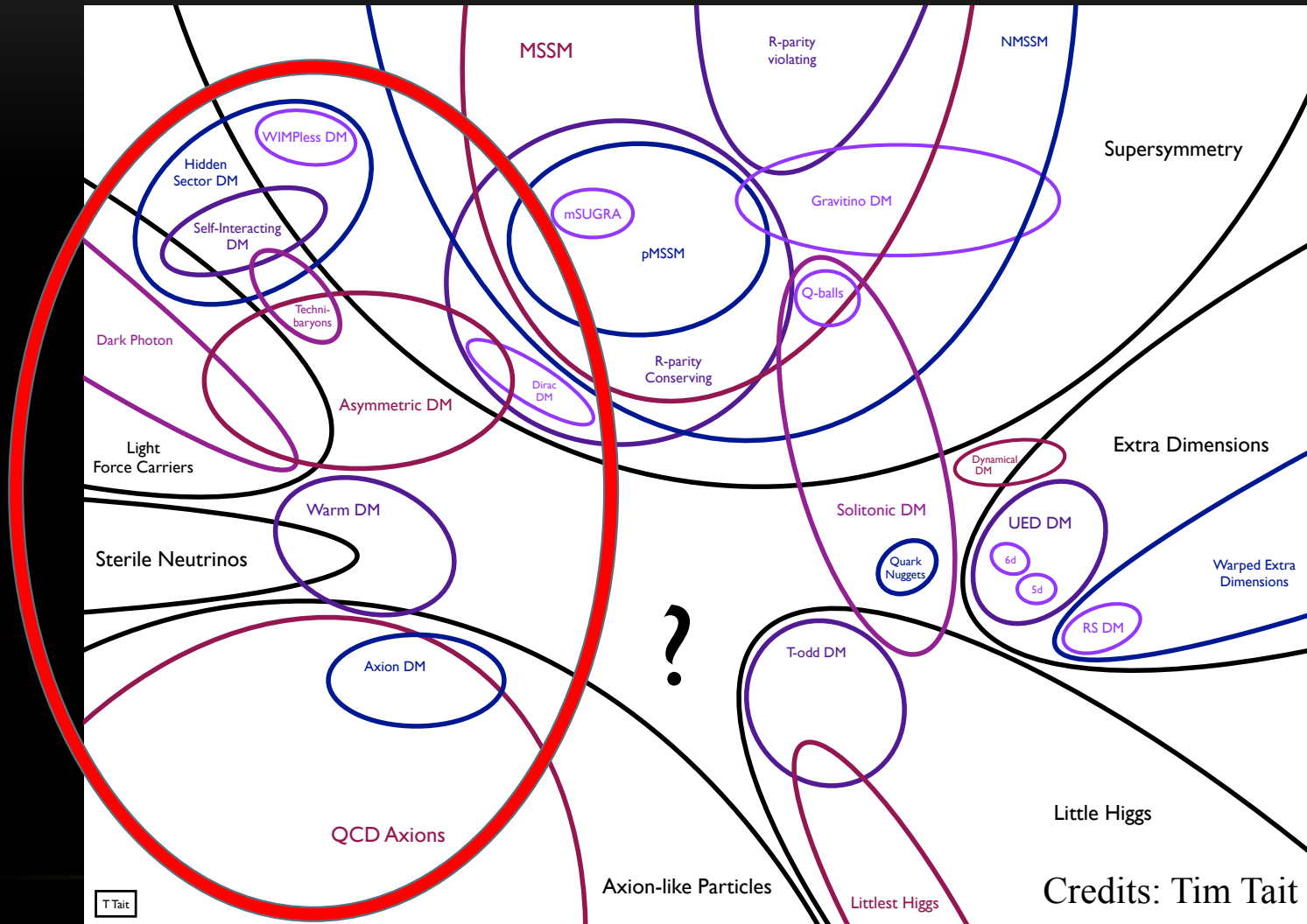


Dark Matter: what is it, and what it isn't

- Dark matter is primarily an **astrophysical and cosmological problem**.
- It is **not** primarily a **particle physics problem**, although it can easily be accommodated within many extensions of the Standard Model.
- Consensus: some kind of new particle(s).

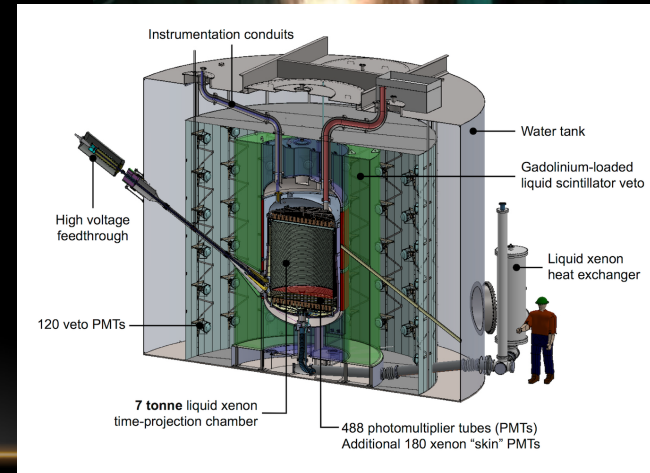
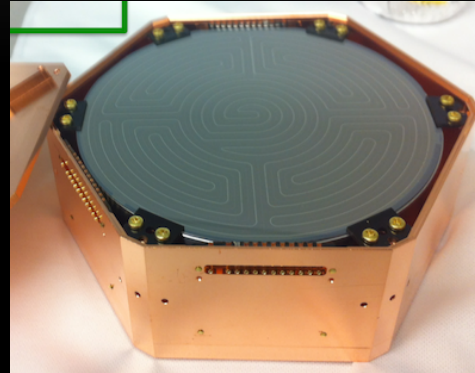
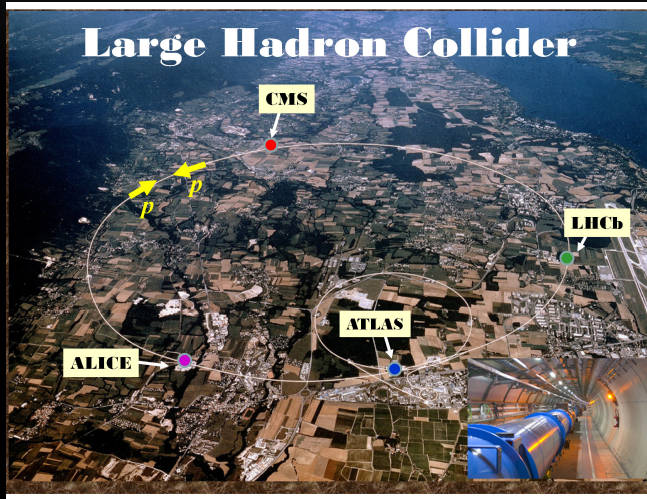


Dark Matter: Possible Ideas



In dark matter science, hope for the best...

- Let's hope we can find dark matter in the lab...



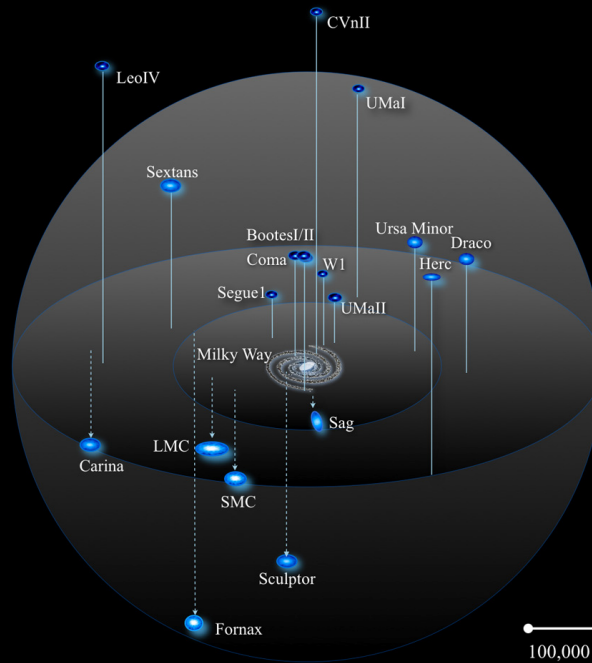
...but prepare for the worst!

- Gravitational signatures might be all we can observe!

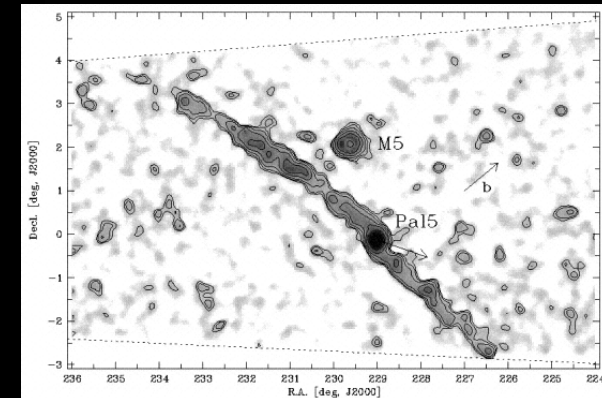
Gravitational Lensing



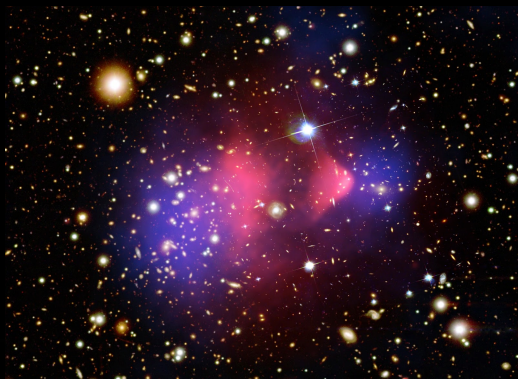
Dwarf galaxies



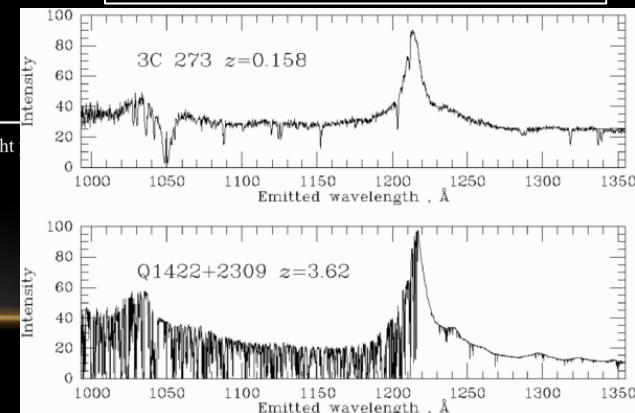
Stellar Streams



Merging Clusters



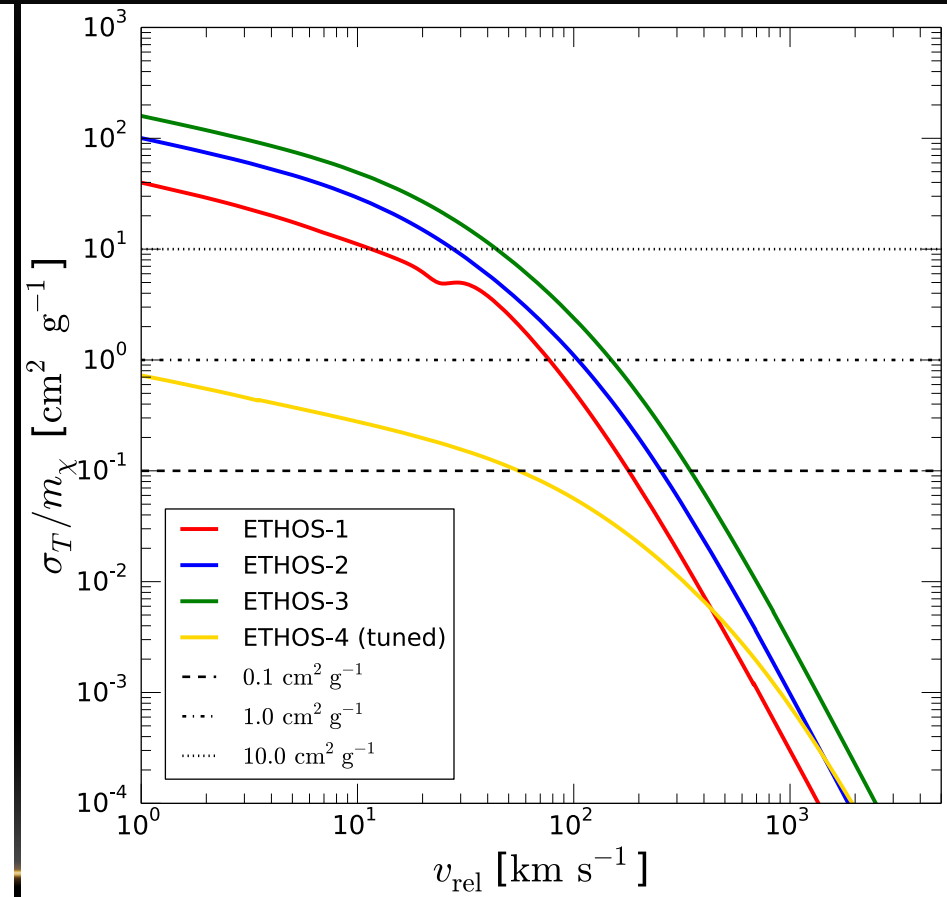
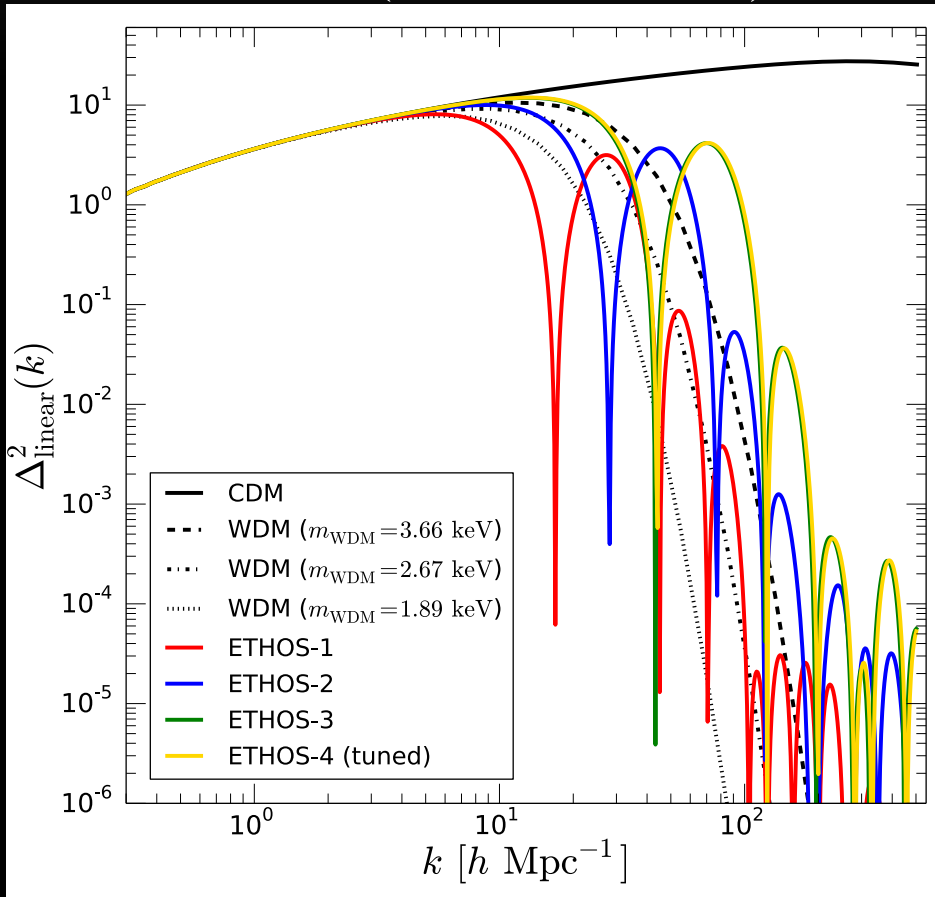
Lyman-alpha forest



Which kind of DM physics can we probe?

1) Physics affecting the DM transfer function (initial conditions)

2) Physics affecting the dynamics of structure formation (self-interaction)

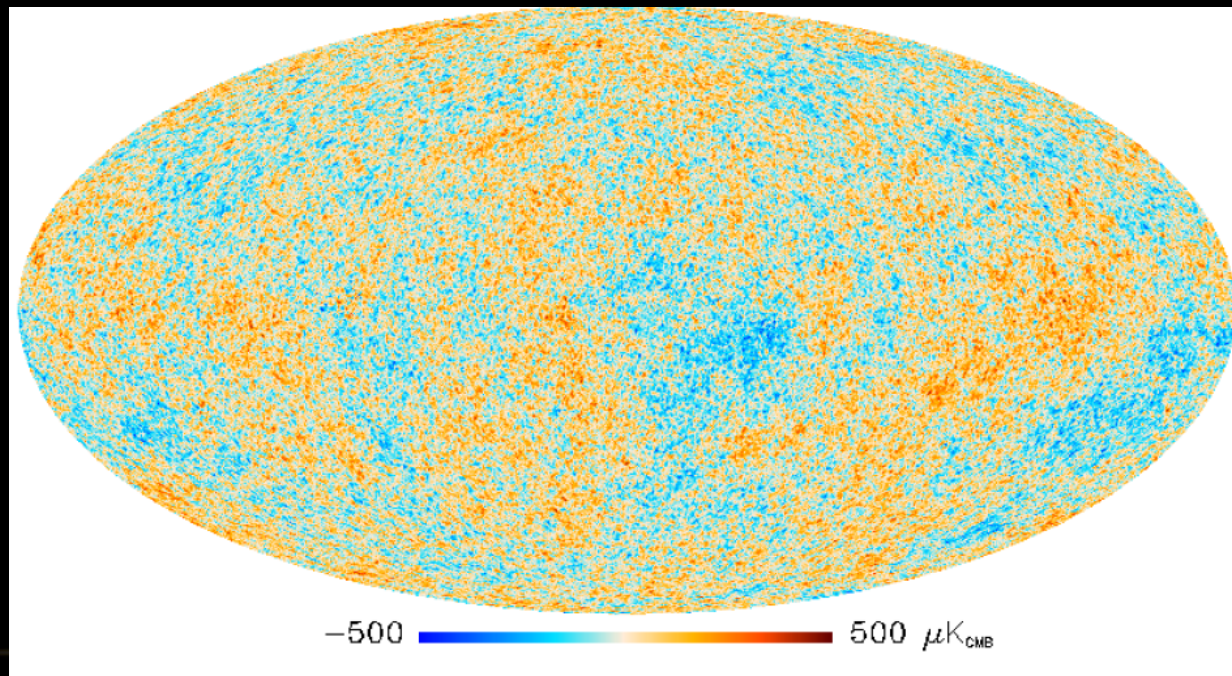


Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349

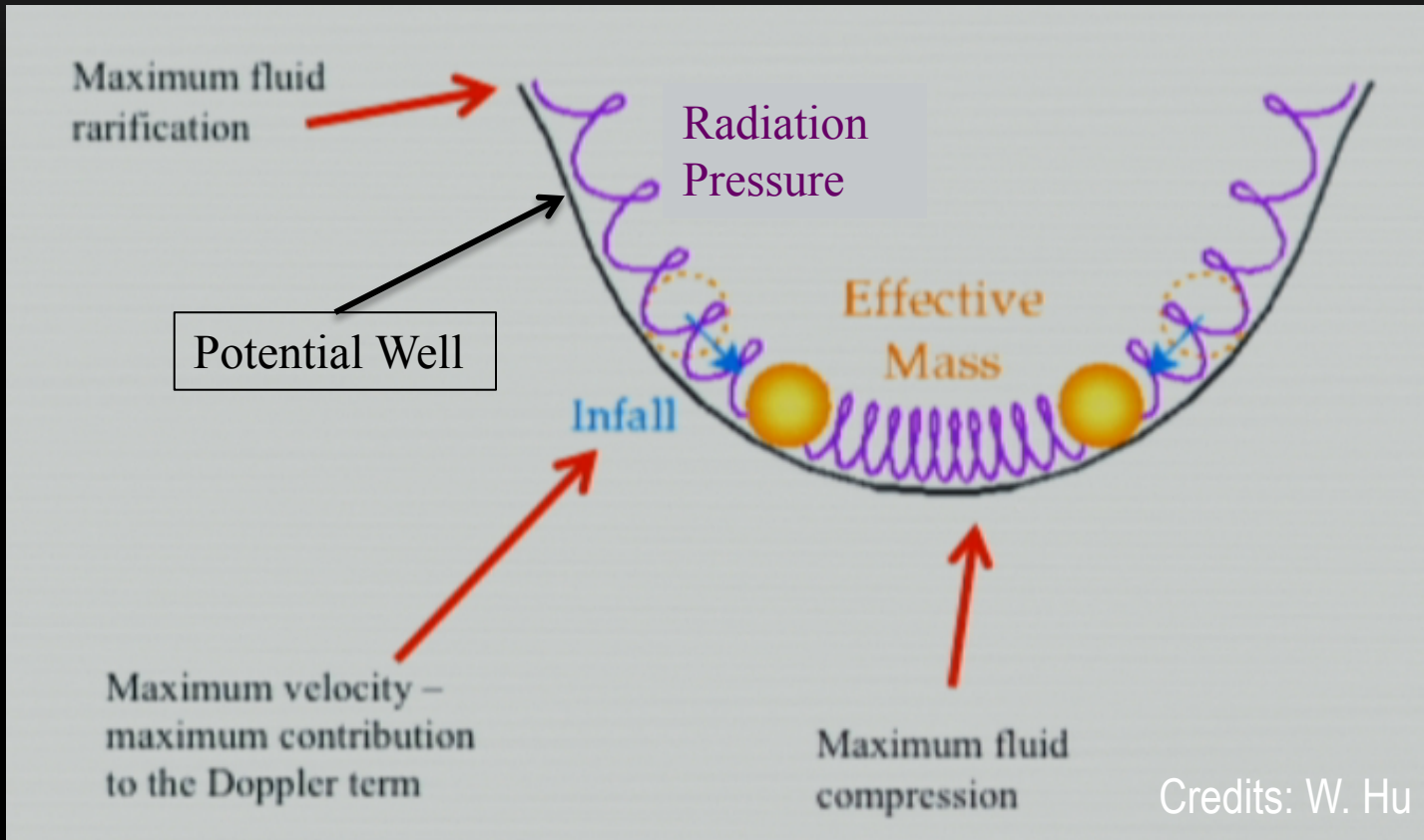
But small scales are messy...

Is there anything we can do with large
scales?

Let's consider the CMB.

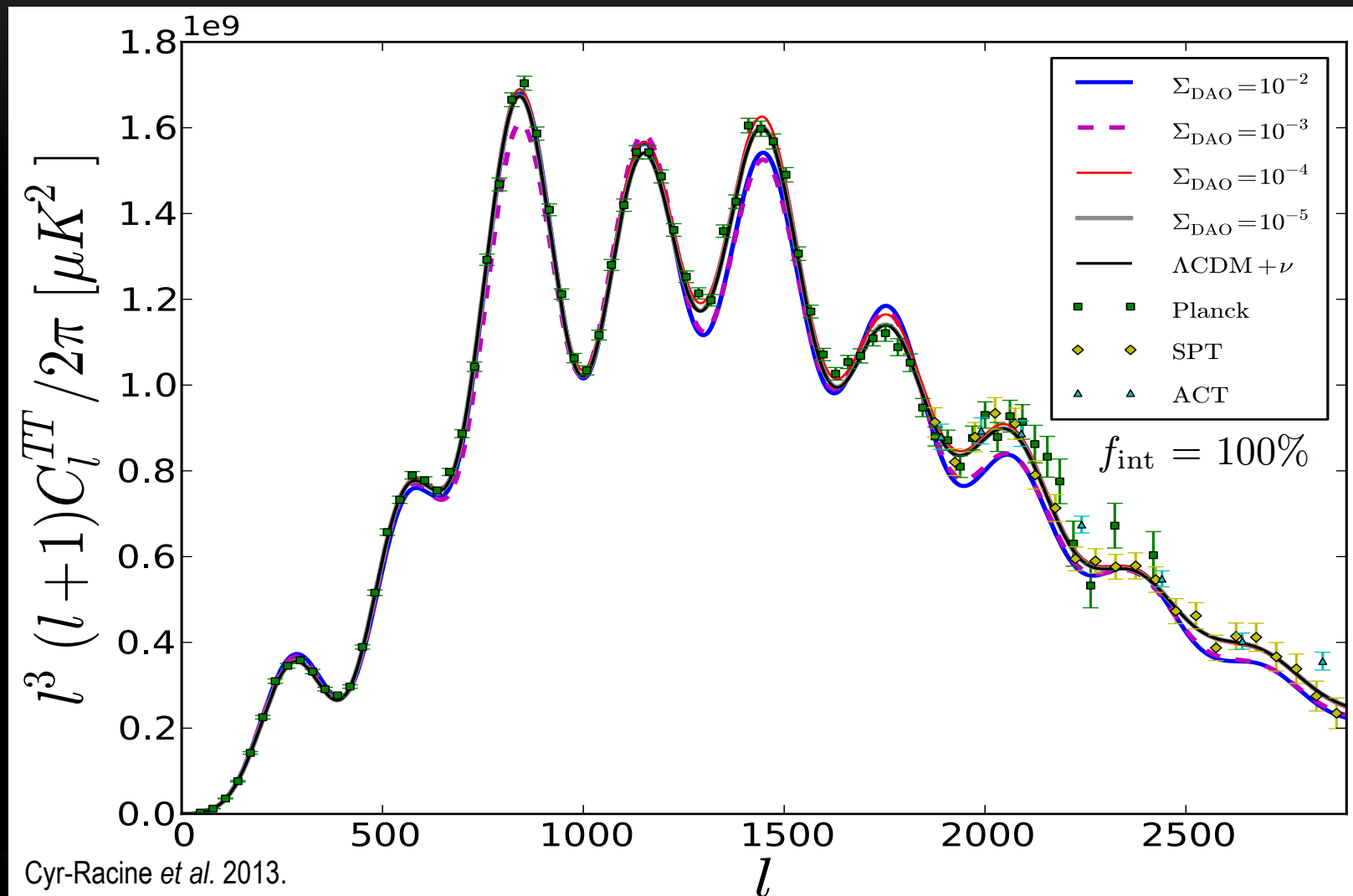


Probing the dark matter transfer function: Cosmic microwave background

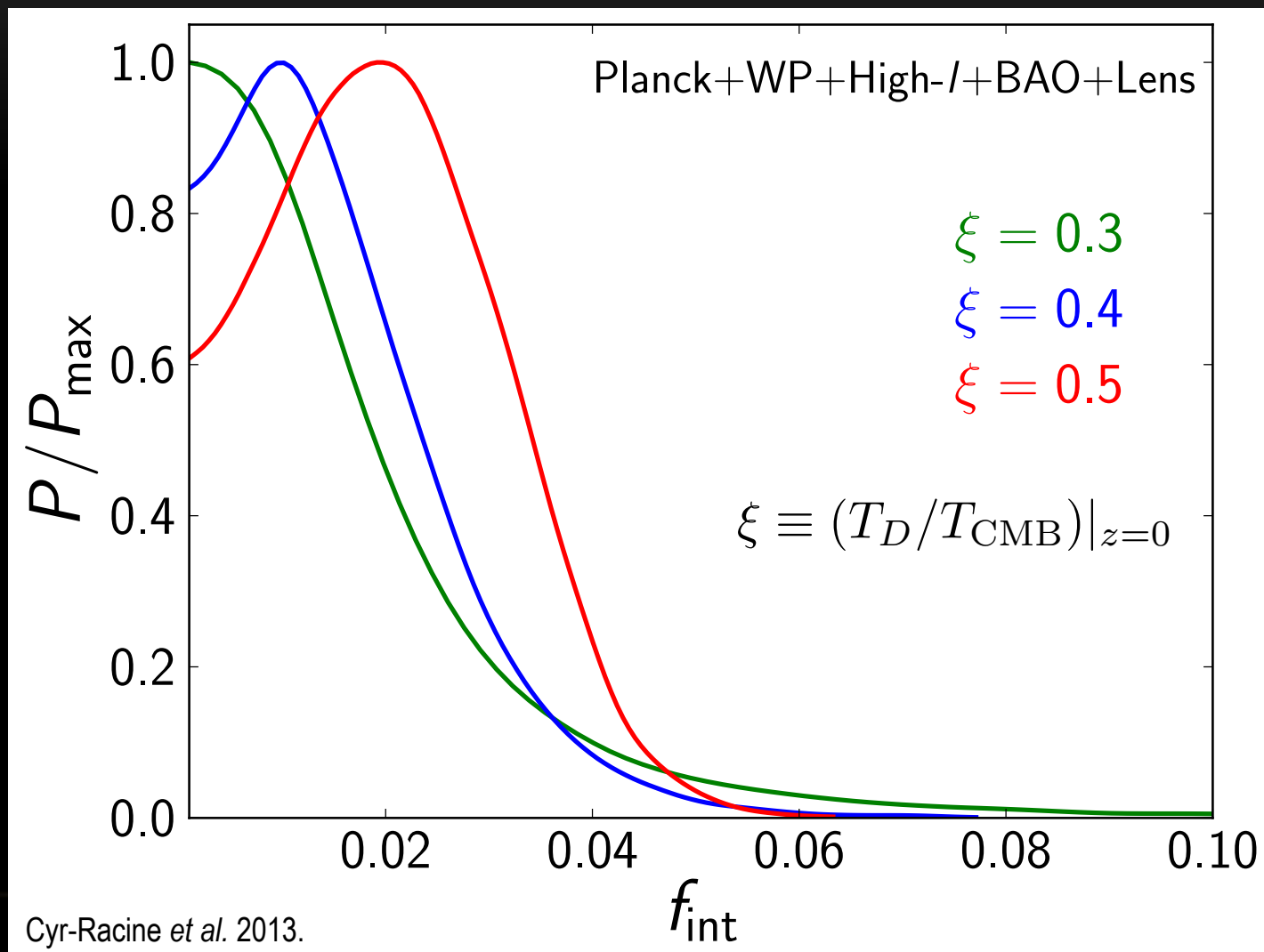


Dark matter is needed to source the gravitational potential

Linear regime: Cosmic microwave background

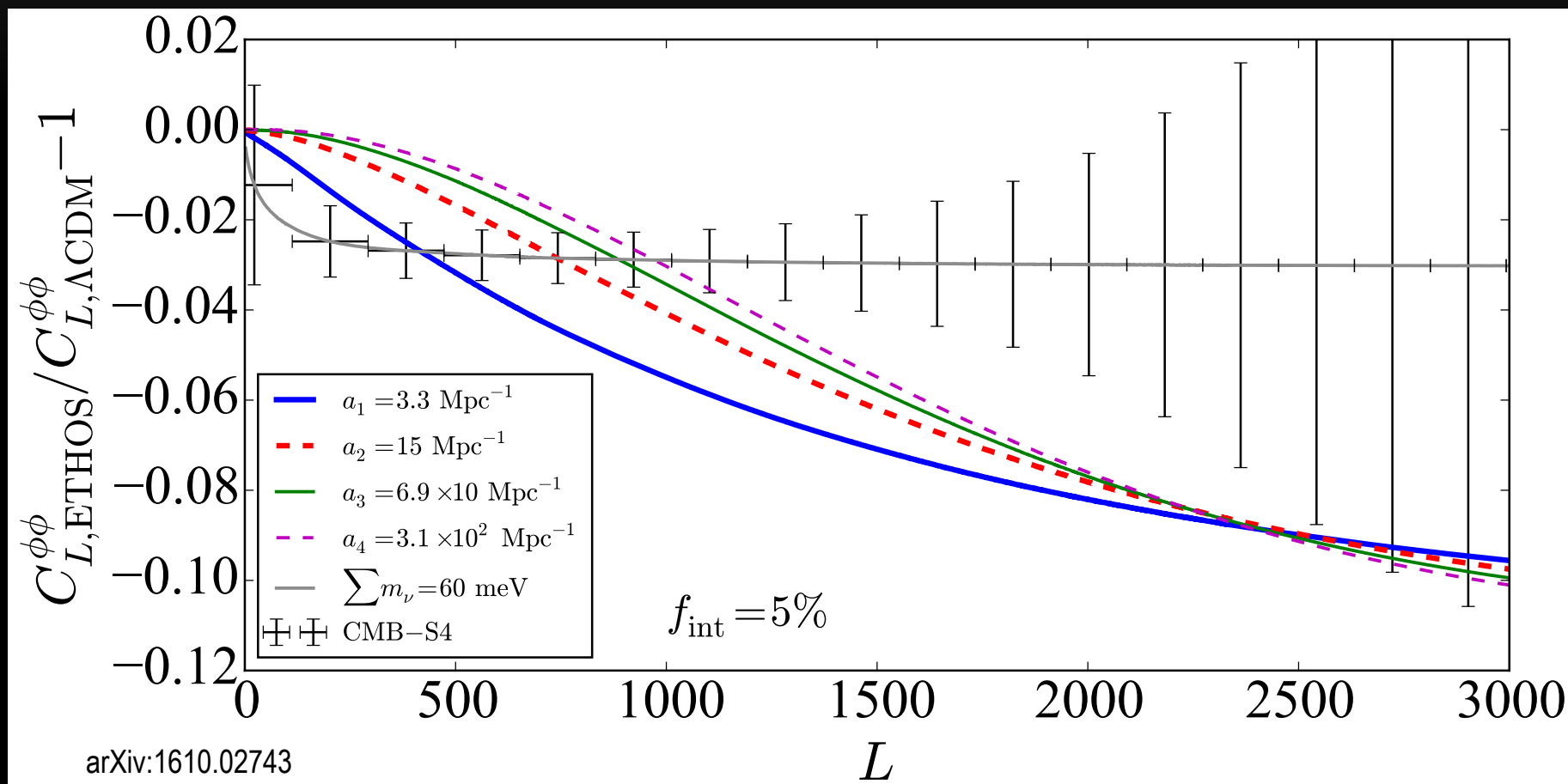


Linear regime: Cosmic microwave background



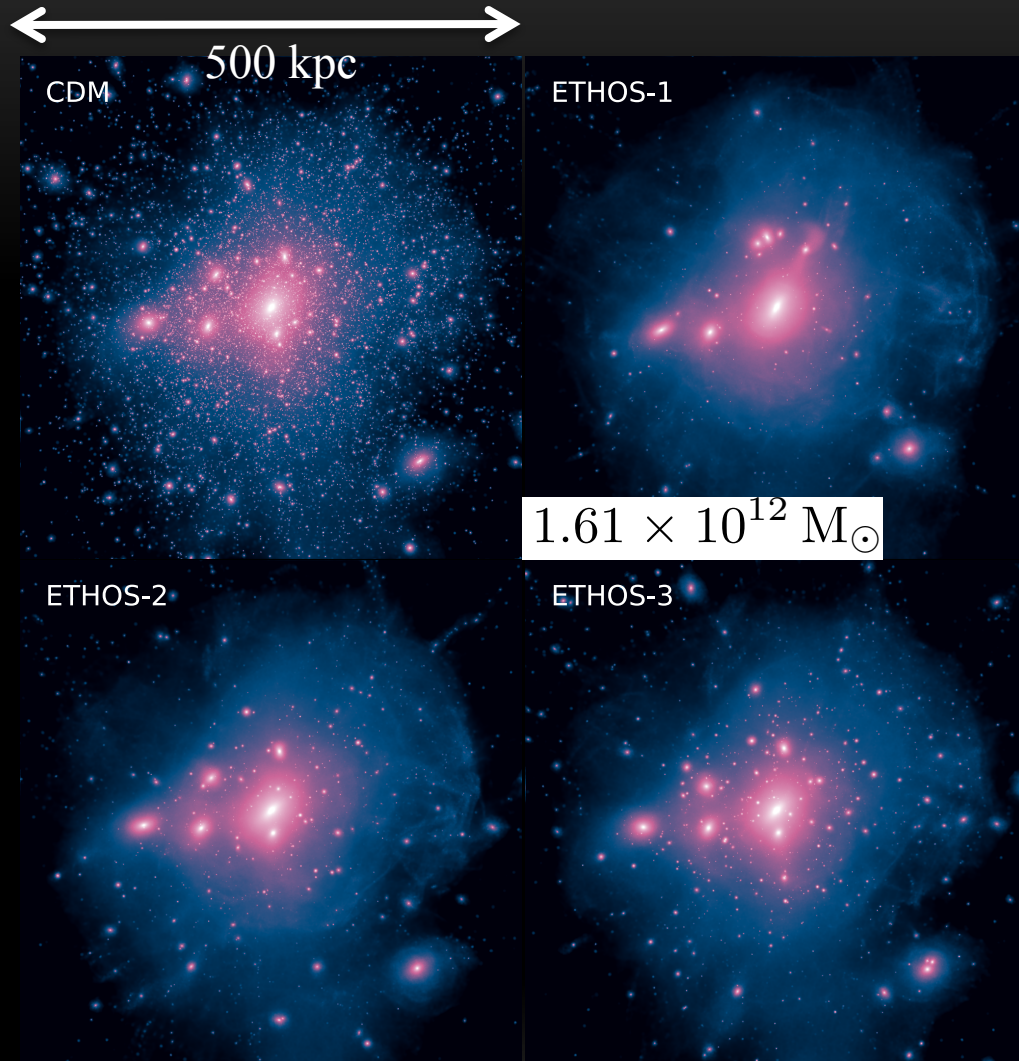
Linear regime: Cosmic microwave background

- The next frontier is to constrain subdominant DM component



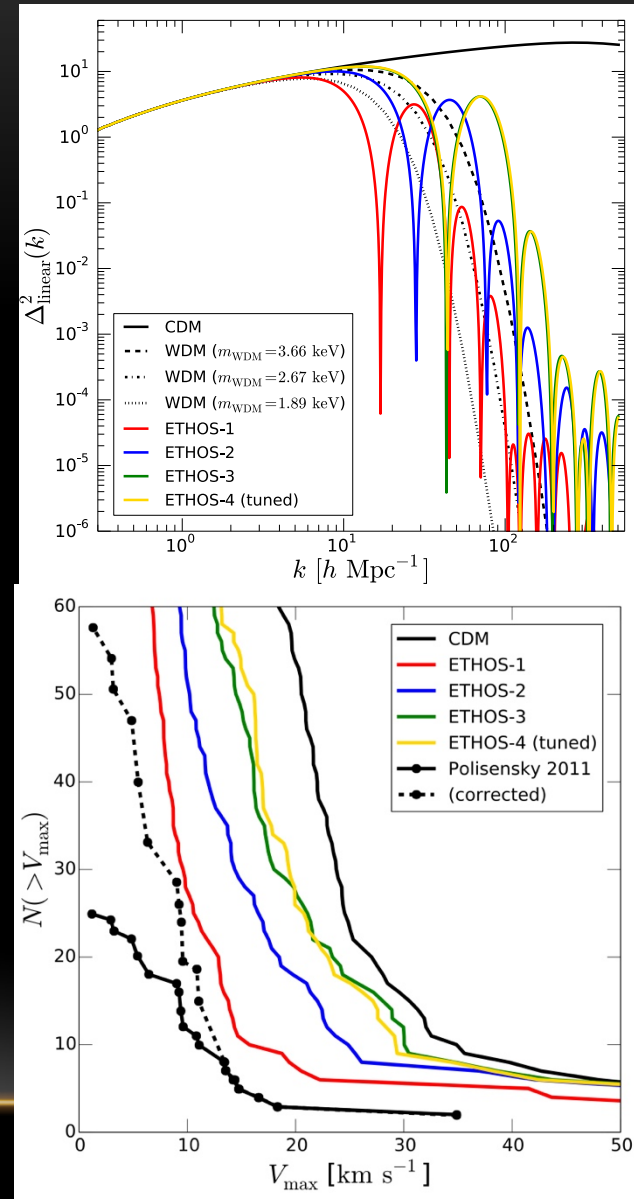
Let's now focus on small, non-linear
scales

Simulations: Subhalo mass function



Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349

Francis-Yan Cyr-Racine, Harvard

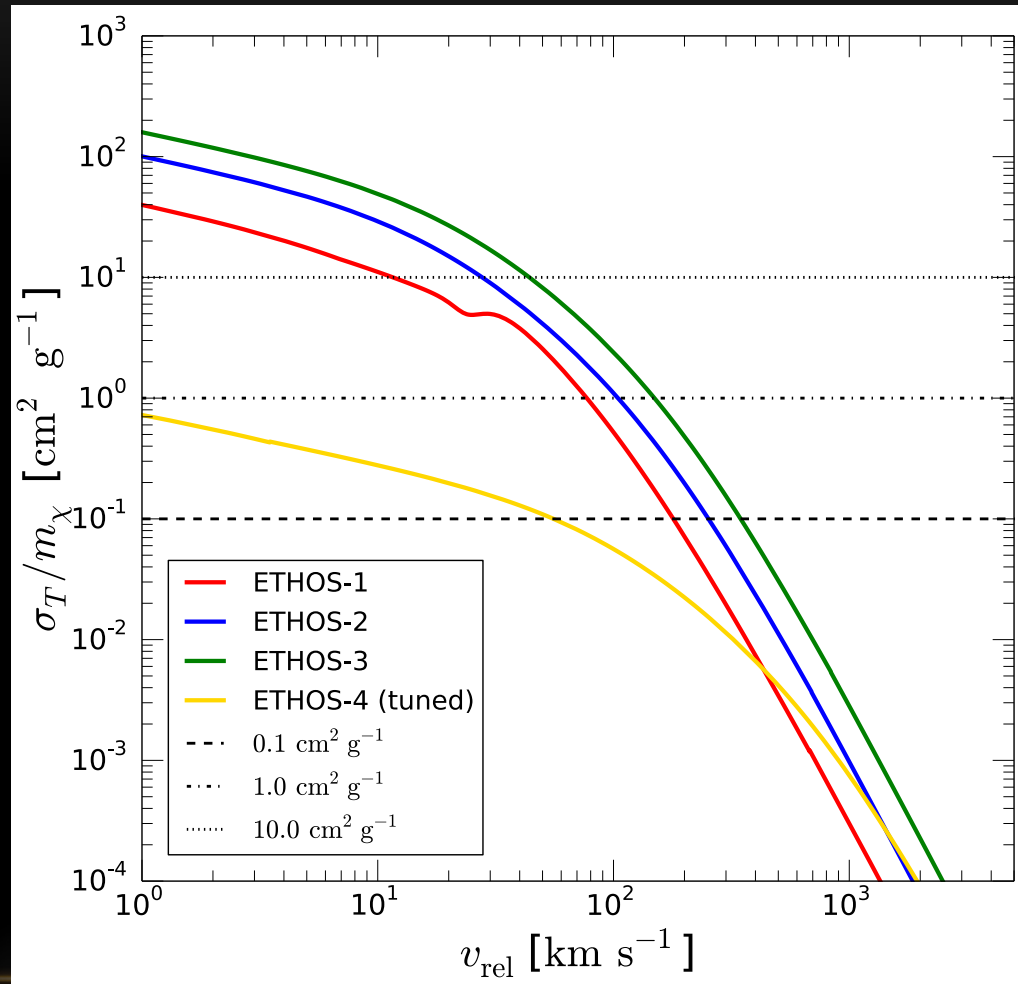


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2) Changing the dynamic of structure formation: self-interaction

Self-
interaction
cross
section over
DM mass

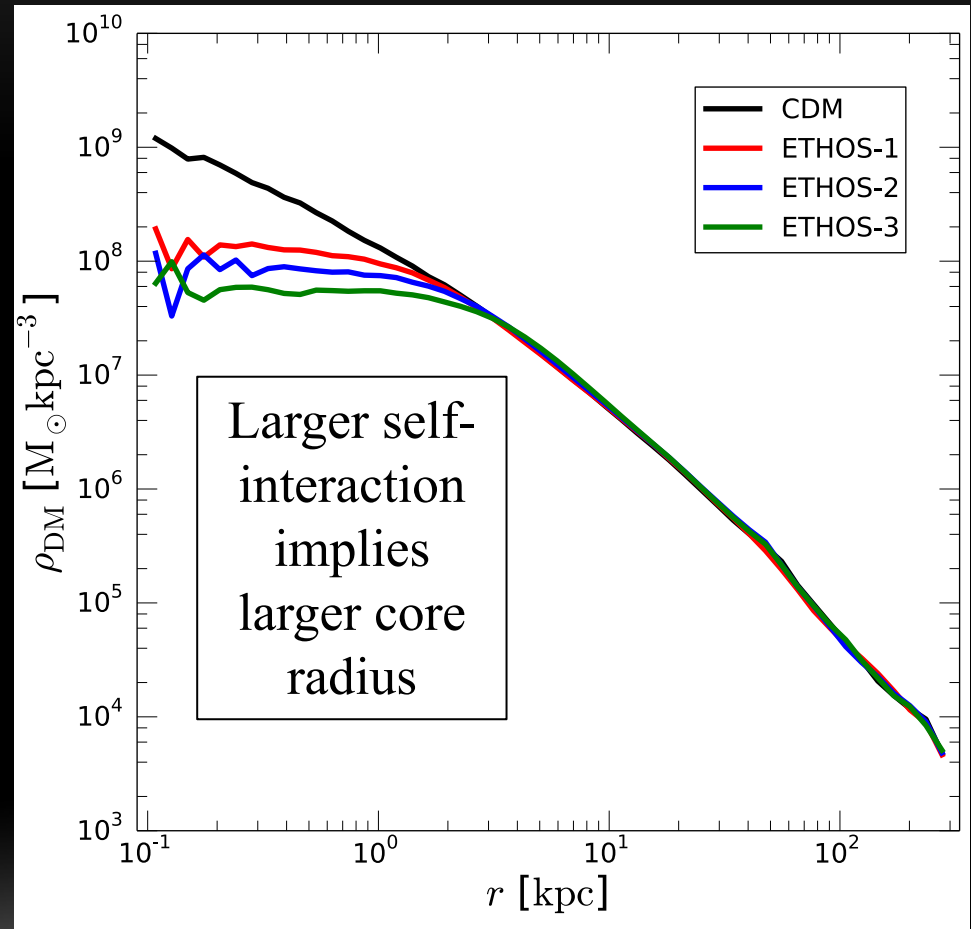


Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349

2) Changing the dynamic of structure formation: self-interaction

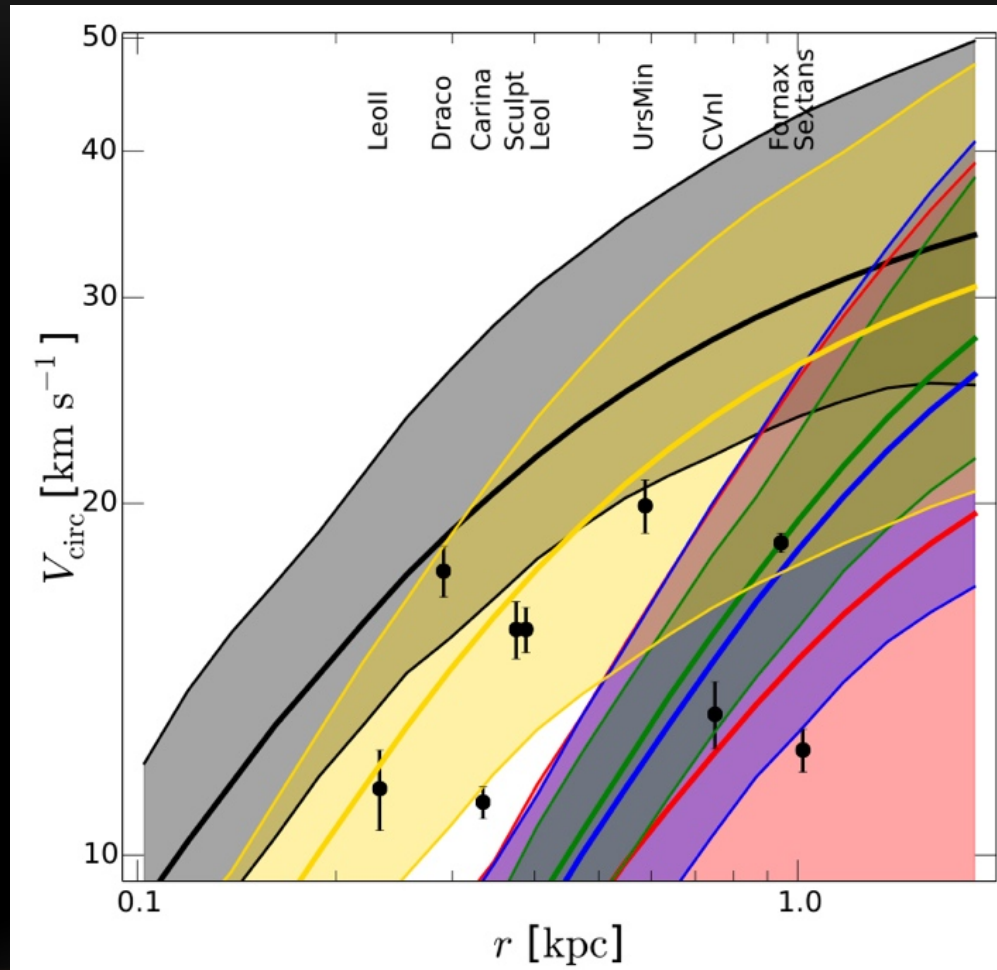
Example:

The density profile of collapsed DM objects develops a constant density core.



Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349

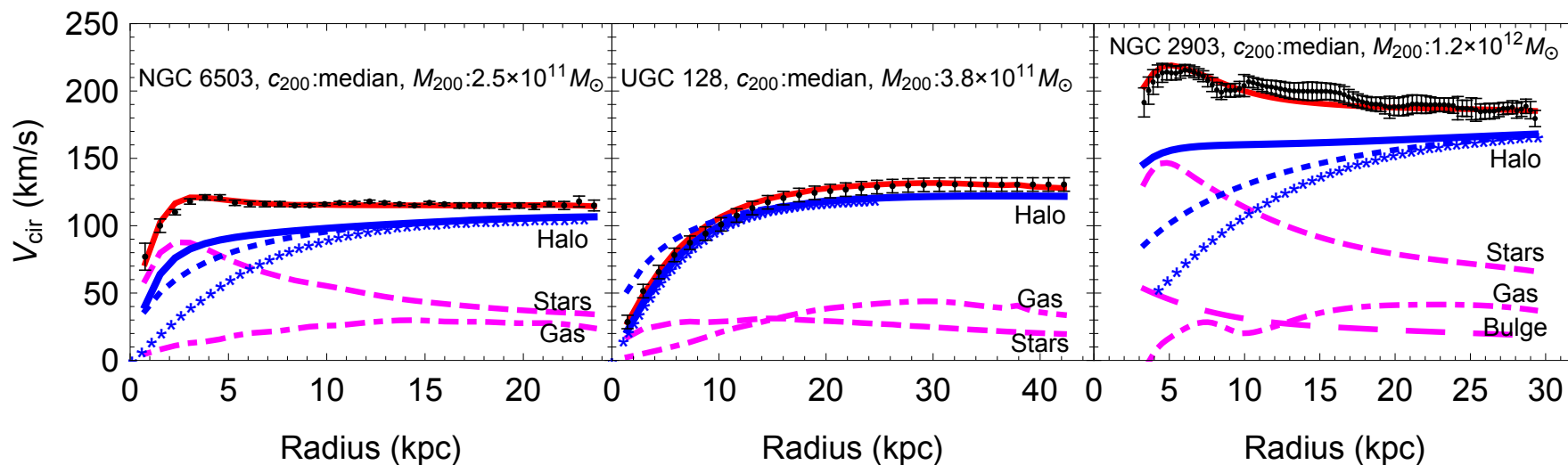
Self-interaction: Velocity profile of local dwarfs



Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349

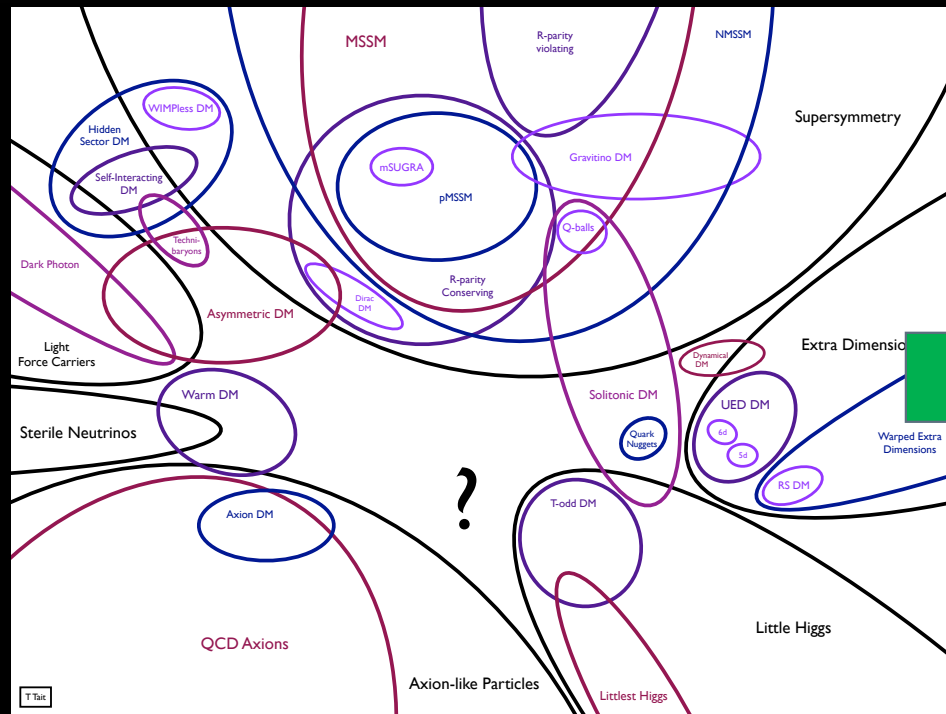
Self-interaction: Interplay with baryons

- Unlike CDM, SIDM can thermalize and equilibrate with the baryons.



From Dark Matter Physics to Predictions

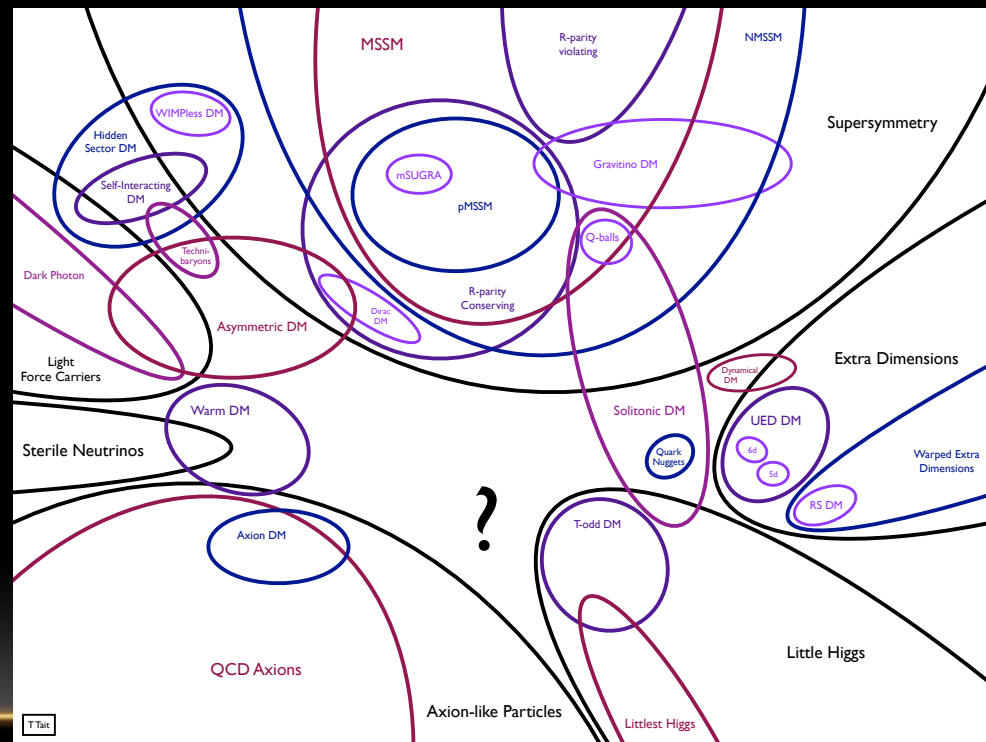
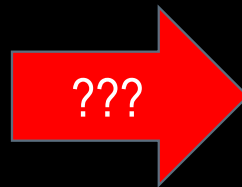
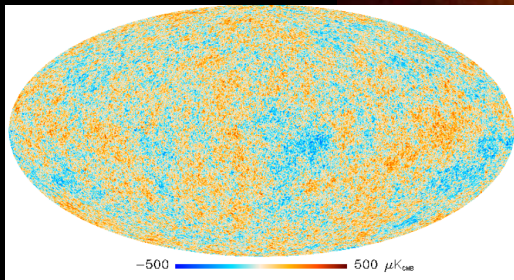
- So far, we have focused on making structure-formation predictions in relevant dark matter models.



From Observations to Dark Matter Physics

- How do we infer the physics of dark matter from observations?

Figure 7 RXJ 1131-1213 (HST)



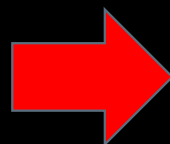
NEEDED:

A classification of dark matter theories according to their structure formation properties.

ETHOS: the Effective Theory Of Structure formation

- ETHOS allows the classification of dark matter theories according to their structure formation properties rather than their intrinsic particle properties.

$$\Xi_{\text{ETHOS}} = \left\{ \omega_{\text{DR}}, \{a_n, \alpha_l\}, \left\{ \frac{\langle \sigma_T \rangle v_{M_i}}{m_\chi} \right\} \right\}$$

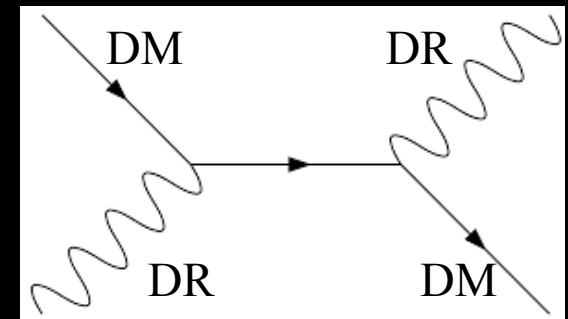
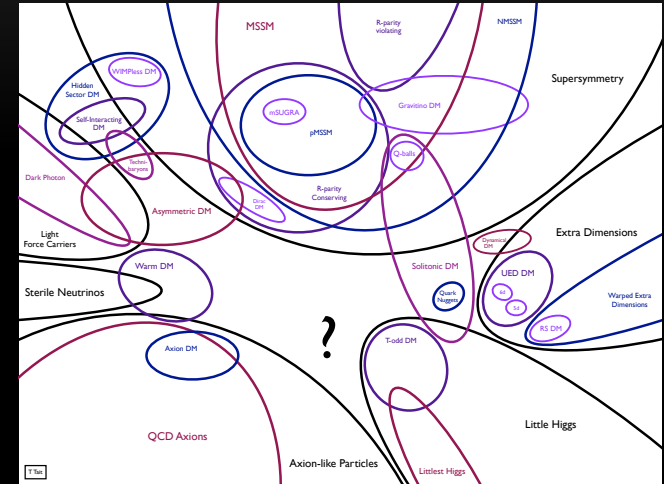


- All dark matter particle models that map to a given effective ETHOS model can be constrained at the same time.

Vogelsberger, Zavala, Cyr-Racine +, arXiv:1512.05349
Cyr-Racine, Sigurdson, Zavala +, arXiv:1512.05344

ETHOS: First Step

- In the first paper, we are primarily concerned with dark matter having **significant interactions with relativistic species**.
- These models are **well-motivated** in the context of **self-interacting dark matter**.
- These models are characterized by a **non-CDM matter power spectrum** and **self-interaction** at late times inside halos.



Cyr-Racine, Sigurdson, Zavala +, arXiv:1512.05344

Classification of dark matter theories according to their structure formation properties

- Where does the particle physics of DM enter the problem?

Dark Matter
Equations

$$\begin{aligned}\dot{\delta}_\chi + \theta_\chi - 3\dot{\phi} &= 0, \\ \dot{\theta}_\chi - c_\chi^2 k^2 \delta_\chi + \mathcal{H}\theta_\chi - k^2\psi &= \dot{\kappa}_\chi [\theta_\chi - \theta_{\text{DR}}]\end{aligned}$$

DM drag opacity

Dark Radiation
Equations

$$\begin{aligned}\dot{\delta}_{\text{DR}} + \frac{4}{3}\theta_{\text{DR}} - 4\dot{\phi} &= 0, \\ \dot{\theta}_{\text{DR}} + k^2(\sigma_{\text{DR}} - \frac{1}{4}\delta_{\text{DR}}) - k^2\psi &= \dot{\kappa}_{\text{DR-DM}}(\theta_{\text{DR}} - \theta_\chi), \\ \dot{\Pi}_{\text{DR},l} + \frac{k}{2l+1}((l+1)\Pi_{\text{DR},l+1} - l\Pi_{\text{DR},l-1}) &= (\alpha_l \dot{\kappa}_{\text{DR-DM}} + \beta_l \dot{\kappa}_{\text{DR-DR}})\Pi_{\text{DR},l},\end{aligned}$$

DR opacity

Angular coefficients

Cyr-Racine, Sigurdson, Zavala +, arXiv:1512.05344

Classification of dark matter theories according to their structure formation properties

- Structure of the opacity terms:

$$\dot{\kappa}_{\text{DR-DM}} = -(\Omega_{\chi} h^2) x_{\chi}(z) \sum_n a_n \left(\frac{1+z}{z_{\text{D}}} \right)^n$$

Envelope Function

Opacity Coefficients

$$\dot{\kappa}_{\chi} = -(\Omega_{\text{DR}} h^2) x_{\chi}(z) \sum_n \left(\frac{2+n}{3} \right) a_n \frac{(1+z)^{n+1}}{z_{\text{D}}^n}$$

Momentum Conservation

Cyr-Racine, Sigurdson, Zavala +, arXiv:1512.05344

What does a_n mean?

- The index n is directly related to the momentum dependence of the scattering:

$$a_n \longrightarrow \sum |\mathcal{M}|^2 \propto \left(\frac{p_{\text{DR}}}{m_\chi} \right)^{n-2}$$

- For example:

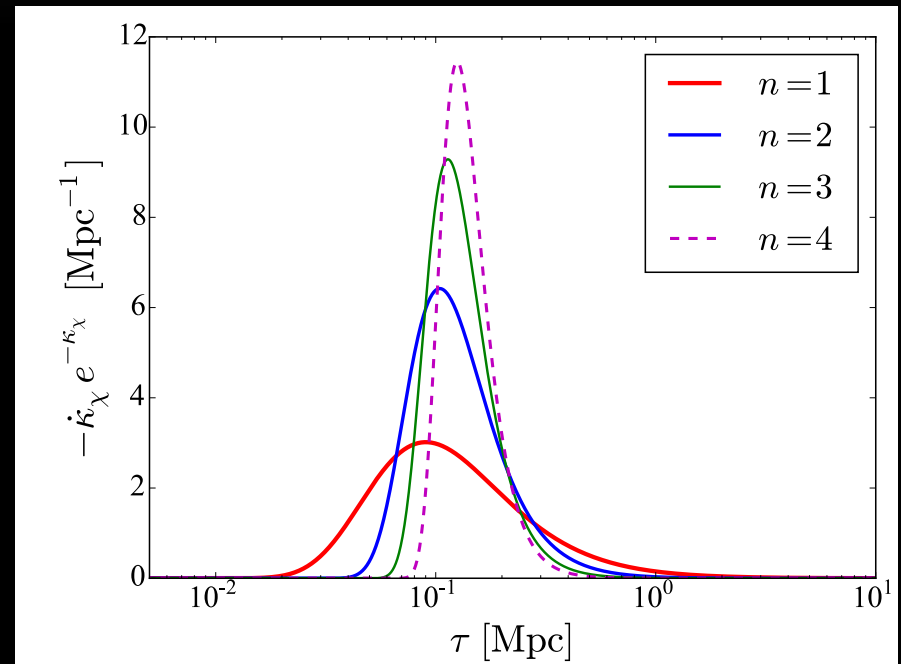
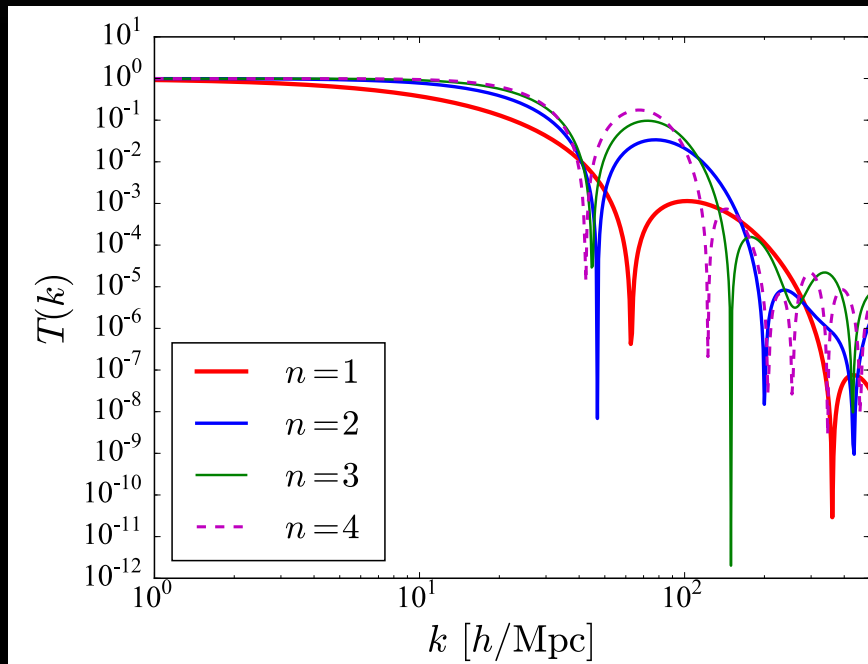
$$a_0 \longrightarrow \sum |\mathcal{M}|^2 \propto \frac{1}{p_{\text{DR}}^2}$$

$$a_2 \longrightarrow \sum |\mathcal{M}|^2 \propto \text{const.}$$

$$a_4 \longrightarrow \sum |\mathcal{M}|^2 \propto p_{\text{DR}}^2$$

Cyr-Racine, Sigurdson, Zavala +, arXiv:1512.05344

Application Shape of linear matter power spectrum for different n

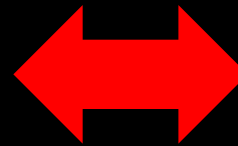


$$-\dot{\kappa}_\chi(z_{\text{drag}}) = \mathcal{H}(z_{\text{drag}})$$

ETHOS: Brief Summary

- Starting from **first principles**, we have identified where the key DM physics enters the evolution of cosmological perturbations.
- We have proposed a **simple parametrization** for the generic case of **DM-DR interaction**.

$$\Xi_{\text{ETHOS}} = \left\{ \omega_{\text{DR}}, \{a_n, \alpha_l\}, \left\{ \frac{\langle \sigma_T \rangle v_{M_i}}{m_\chi} \right\} \right\}$$



ETHOS -- Next Step: Make it an actual “effective” theory

- ETHOS

- Early times: perturbative density fluctuations.
- Near collapse time: k -modes strongly-coupled, need simulations.
- After virialization: new degrees of freedom (halos??) **What is the effective description here?**

- QCD

- High T : perturbative in the quarks.
- Near T_c : strongly-coupled, need Lattice computation.
- $T \ll T_c$: new degrees of freedom (mesons, hadrons), chiral Lagrangian.

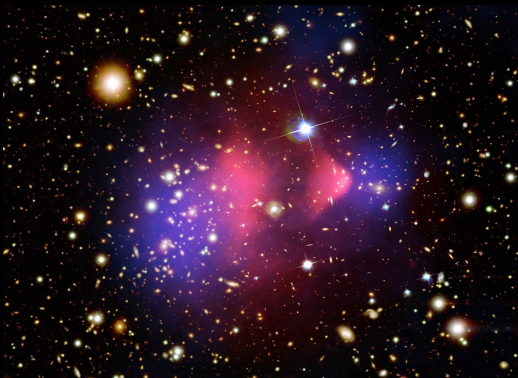
Ultimate goal: Combine constraints

Gravitational Lensing

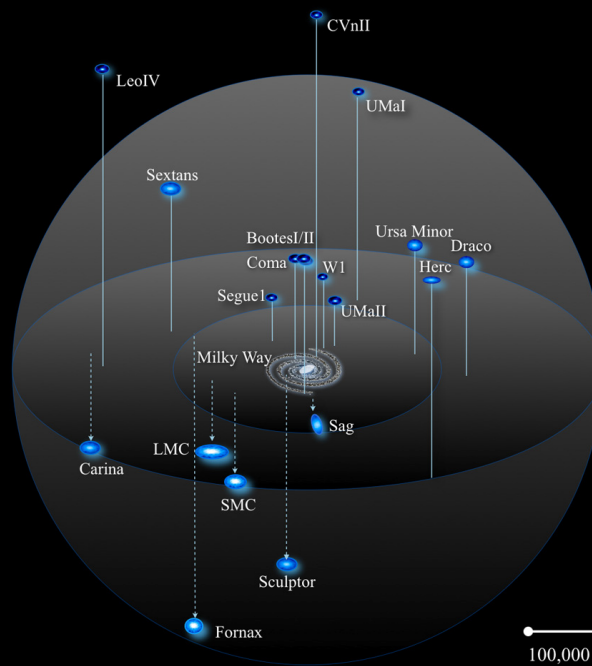
Figure 7



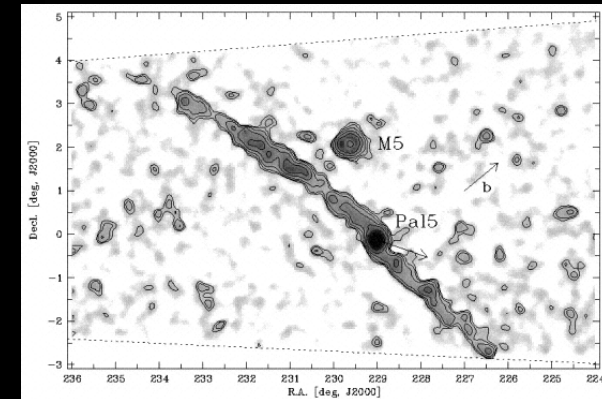
Merging Clusters



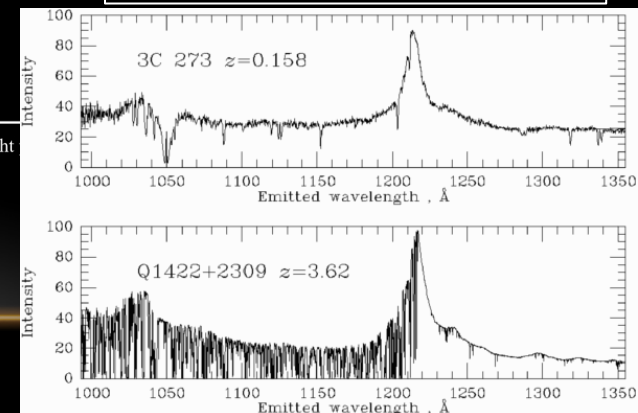
Dwarf galaxies



Stellar Streams



Lyman-alpha forest



Take-Home Message

- In the most pessimistic scenario, **gravitational signatures** of DM might be all we have access to.
- Important clues about DM physics lie on **small-scales**.
- From an **astrophysical perspective**, there is a need to classify and parametrize dark matter models with respect to their **structure formation properties**.
- We have taken a first step in this direction with the **ETHOS framework**.
- Much work remains to be done to make ETHOS an actual **effective theory** and use it as a common language to describe observations.