

Dim.-6 Gluon Interactions in Multijet Events

at a future 100 TeV pp-collider

Silvan Kuttimalai

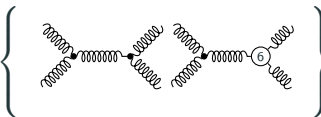
January 17, 2017

1st FCC Physics Workshop, CERN

Dimension-6 Gluon Operator

$$\mathcal{O}_G = \frac{c_G}{\Lambda^2} f_{abc} G_{a,\nu}^\mu G_{b,\kappa}^\nu G_{c,\mu}^\kappa$$

- Only one pure gluonic operator at dimension 6
- Operators involving covariant derivatives can be related to quark operators
- Gives rise to effective 3,4,5, and 6 gluon vertices
- \mathcal{O}_G does NOT contribute to massless $2 \rightarrow 2$ scattering at order Λ^{-2}

$$2\text{Re} \left\{ \begin{array}{c} \text{diagram 1} \\ \text{diagram 2} \\ \text{diagram 3} \end{array} \right\} = 0$$


$$\mathcal{O}_G = \frac{c_G}{\Lambda^2} f_{abc} G_{a,\nu}^\mu G_{b,\kappa}^\nu G_{c,\mu}^\kappa$$

- Some effects in specific kinematic regimes of three-jet events:
experimentally challenging

[Dixon, Shadmi, Nucl. Phys. B 423 (1994)]

- Four-jet production at LEP:
small collider energy

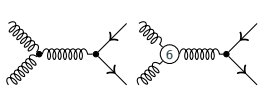
[Duff, Zeppenfeld, Z. Phys. C 53 (1992)] [Dreiner, Duff, Zeppenfeld, Phys. Lett. B 282 (1992).]

- Higgs+jet production:
Little sensitivity

[Ghosh, Wiebusch, Phys. Rev. D 91 (2015)]

Collider Constraints: Top Pair Production

$$\mathcal{O}_G = \frac{c_G}{\Lambda^2} f_{abc} G_{a,\nu}^\mu G_{b,\kappa}^\nu G_{c,\mu}^\kappa$$

$$2\text{Re} \left\{ \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \end{array} \right\} \neq 0 \quad \text{for } m_q \neq 0$$


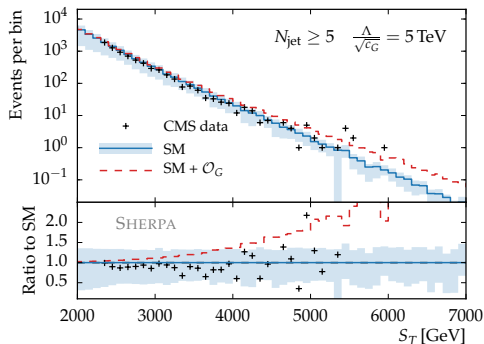
- Very sensitive channel already at Tevatron energies $\Lambda/\sqrt{c_G} > 2 \text{ TeV}$

[Cho, Simmons, Phys. Lett. B 323 (1994)]

- Multitude of other operators affect top sector [Buckley et al., JHEP 04 (2016)]
- Limits in global analyses are weaker: $\Lambda/\sqrt{c_G} > 850 \text{ GeV}$

Collider Constraints: Multijet Production

- Consider higher jet multiplicities $N_{\text{jet}} \geq 4, 5$
- Events produced in abundance at LHC
- 13 TeV-data exists: CMS-PAS-EXO-15-007
- Extended Sherpa to simulate **any** EFT extension of SM
- Simulation of $p p \rightarrow 7$ jets feasible

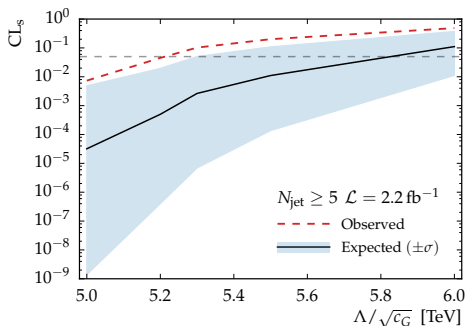


$$S_T = E_T^{\text{miss}} + \sum_{\text{Objects}} E_T$$

- Use S_T distribution
 - Perform CL_s analysis
 - $\text{CL}_s < 5\% \Rightarrow$ exclusion
- $\Rightarrow \Lambda/\sqrt{c_G} > 5.2 \text{ TeV}$

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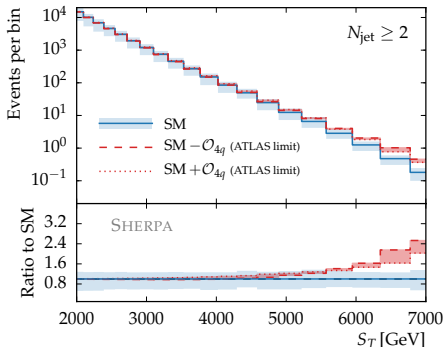
Contamination Due to Four-Quark Operators

Example Operator

$$\mathcal{O}_{q4} = \frac{c_{q4}}{\Lambda^2} \sum_{q,q'} (\bar{q}_L \gamma^\mu q_L) (\bar{q}'_L \gamma^\mu q'_L)$$

ATLAS limits

- $\Lambda/\sqrt{c_{q4}} > +4.8 \text{ TeV}$
- $\Lambda/\sqrt{c_{q4}} < -6.8 \text{ TeV}$



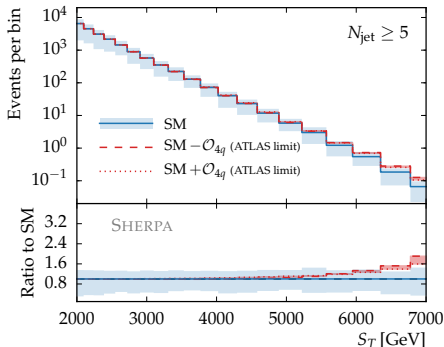
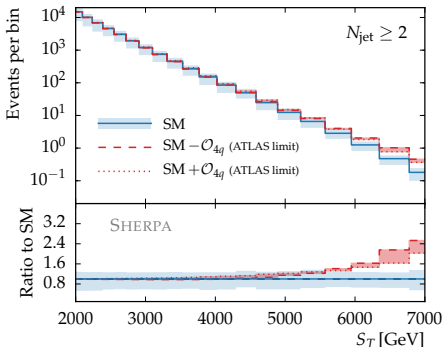
Contamination Due to Four-Quark Operators

Example Operator

$$\mathcal{O}_{q4} = \frac{C_{q4}}{\Lambda^2} \sum_{q,q'} (\bar{q}_L \gamma^\mu q_L) (\bar{q}'_L \gamma^\mu q'_L)$$

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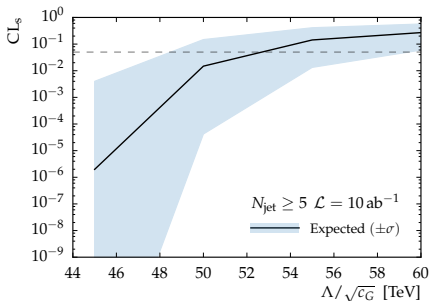
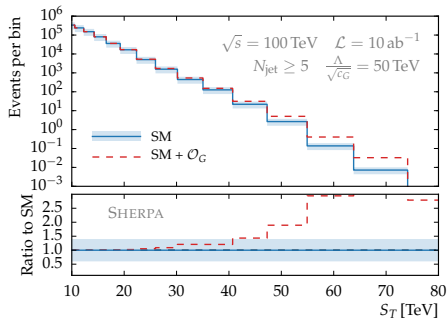
Four-Quark contamination under control in high multiplicity bins

Projected Sensitivity

- Use same observable: S_T
- Anti- k_T jets, $R = 0.4$,
 $p_T^{\min} = 1 \text{ TeV}$
- Parton-level analysis
- Assume 40% / 20% syst. unc.

⇒ Sensitive up to $\Lambda \approx 55 \text{ TeV}$

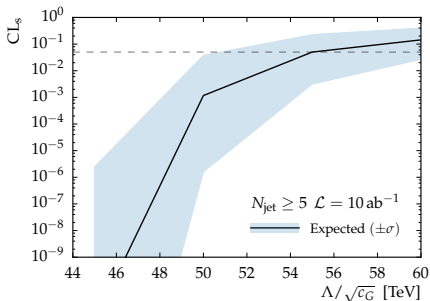
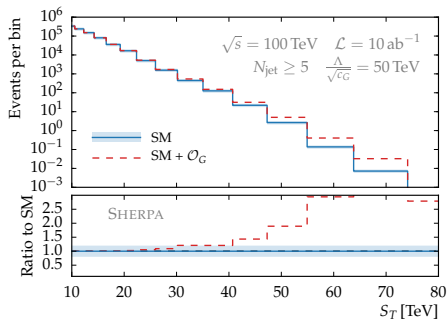
⇒ Insensitive to region $S_T > \Lambda$



Projected Sensitivity

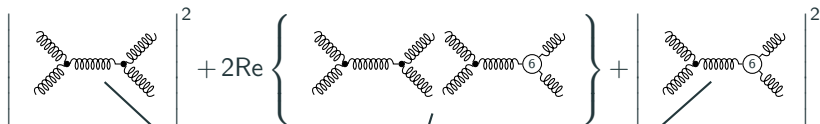
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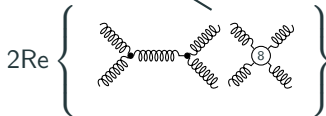
Power Counting

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{c_6}{\Lambda^2} \mathcal{O}_6$$



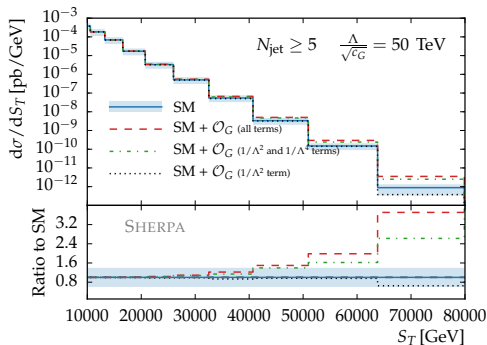
$$d\sigma = d\sigma_{\text{SM}} + \frac{1}{\Lambda^2} d\sigma^{(2)} + \frac{1}{\Lambda^4} d\sigma^{(4)} + \dots$$

- For high final state multiplicities: high powers of Λ^{-1} occur
- Terms of Λ^{-4} and higher technically beyond parametric accuracy



Higher Powers in Λ^{-1}

$$d\sigma = d\sigma_{\text{SM}} + \frac{1}{\Lambda^2} d\sigma^{(2)} + \frac{1}{\Lambda^4} d\sigma^{(4)} + \dots$$



Assess term by term:

Λ^{-2} negligible (but non-zero)

Λ^{-4} numerically leading term
(generic problem for G^3)

Λ^{-6} sizable only for $S_T \geq \Lambda$

Conclusions

- Anomalous dim.-6 interactions in QCD can be probed efficiently in multijet events
- Applies to 4-quark interactions and gluon interactions
- LHC bounds on Λ in $\mathcal{O}(5 \text{ TeV})$ ballpark
- An FCC-pp collider at 100 TeV provides sensitivity up to $\mathcal{O}(50 \text{ TeV})$
- Bounds on \mathcal{O}_G rely on inclusion of terms beyond Λ^{-2}

Backup slides

Background Simulation

Signal and Background Simulation: Multijet-Merging

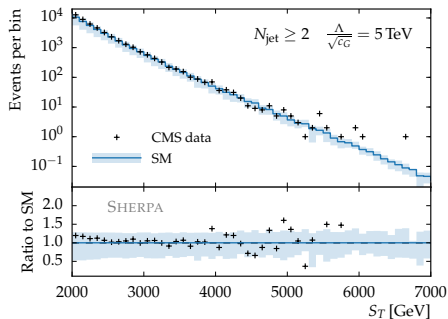
$$d\sigma = d\sigma_{\text{SM}} + \left(\frac{1}{\Lambda^2} d\sigma^{(2)} + \frac{1}{\Lambda^4} d\sigma^{(4)} + \dots \right)$$

$pp \rightarrow jj$	(N)LO+PS		
$pp \rightarrow jjj$	LO+PS	$pp \rightarrow jj$	LO+PS
$pp \rightarrow jjjj$	LO+PS	$pp \rightarrow jjj$	LO+PS
$pp \rightarrow jjjjj$	LO+PS	$pp \rightarrow jjjj$	LO+PS
$pp \rightarrow jjjjjj$	LO+PS	$pp \rightarrow jjjjj$	LO+PS

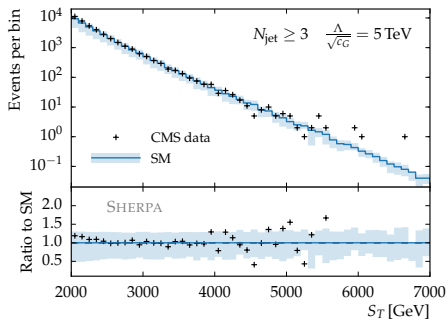
- Combine all parton shower matched samples into one
- Describe jets above hardness threshold Q_{cut} by fixed order ME
- Parton shower fills remaining phase space
- Restore all-order resummation properties of parton shower
- Several schemes on the market, use CKKW (Sherpa) here

Background Simulation: Performance in LHC Environment

Shape and normalization well described, uncertainty $\approx \pm 40\%$



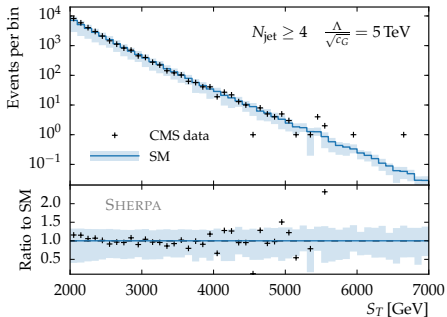
$$N_{\text{jet}} \geq 2$$



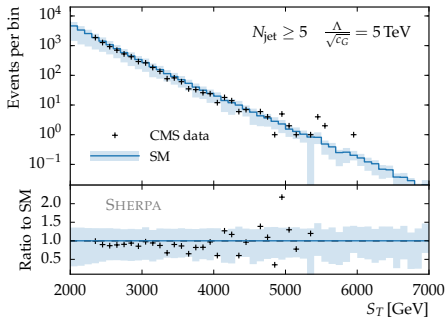
$$N_{\text{jet}} \geq 3$$

Background Simulation: Performance in LHC Environment

Shape and normalization well described, uncertainty $\approx \pm 40\%$



$$N_{\text{jet}} \geq 4$$



$$N_{\text{jet}} \geq 5$$

Uncertainty estimate: μ_f and μ_r variations by factors of 2

CMS Analysis

Existing Run-II Data on Multijets

CMS PAS EXO-15-007

- 2.2 fb^{-1}
- Search for macroscopic black holes
- Main observable: transverse energy

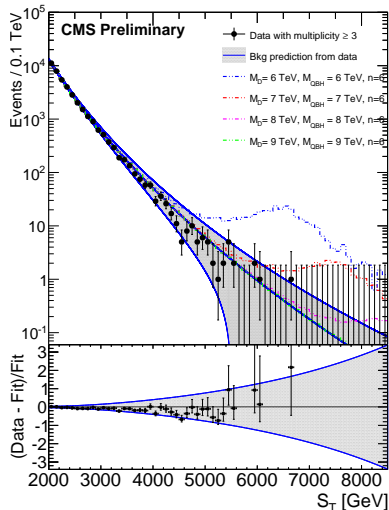
$$S_T = E_T^{\text{miss}} + \sum_{\text{Objects}} E_T$$

- Measured in multi-TeV region
- Binned according to jet multiplicity

$$N_{\text{jet}} \geq 1, \dots, 10$$

- In SM: S_T purely QCD driven

2.2 fb⁻¹ (13 TeV)



Existing Run-II Data on Multijets

CMS PAS EXO-15-007

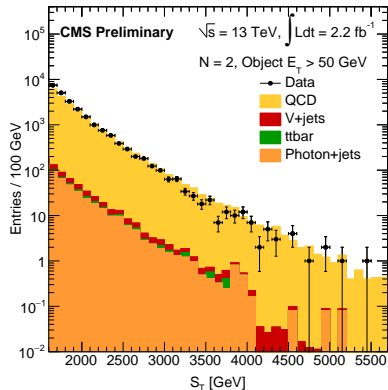
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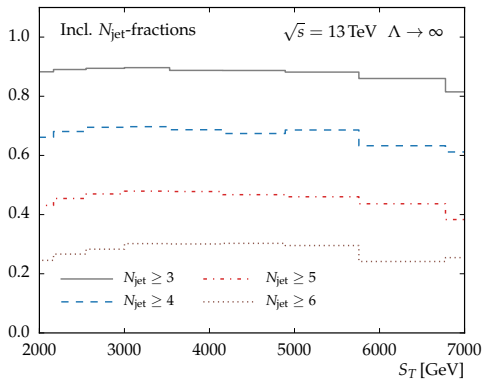
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Jet multiplicites



Statistical analysis

Setting exclusion limits

- Input: individual bins in S_T -distribution
- Treat bins uncorrelated, use likelihood ratio as test statistic:

$$\begin{aligned}\chi(\text{data}) &= 2 \log \frac{P_{s+b}(\text{data})}{P_b(\text{data})} \\ &= 2 \log \prod_{i=1}^{N_{\text{bins}}} \frac{P_{s+b}^i(n_i)}{P_b^i(n_i)}\end{aligned}$$

- Assume Poisson distribution

$$P_b^i(n_i) = b_i^{n_i} \frac{e^{-b_i}}{n_i!} \quad \text{No syst. unc.}$$

$$P_b^i(n_i) = \frac{1}{b_i^1 - b_i^0} \int_{b_i^0}^{b_i^1} b^{n_i} \frac{e^{-b}}{n_i!} db \quad \text{Unc. band } b_i^0 < b_i < b_i^1$$

Setting exclusion limits

- Given some data, define the standard frequentist confidence levels

$$CL_b = P_b(\chi < \chi(\text{data}))$$

$$CL_{s+b} = P_{s+b}(\chi < \chi(\text{data}))$$

- Calculate them by brute force (using Metropolis-Hastings MC)
- Use modified frequentist confidence level to set exclusion limits

$$CL_s = CL_{s+b}/CL_b$$

- Consider “signal+background” hypothesis excluded if

$$CL_s \leq 5\%$$

Signal simulation

Model Implementation via FeynRules and UFO

Lagrangian

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{c_G}{\Lambda^2} G^3 + \dots$$

FeynRules

[Alloul et al., Comput. Phys. Commun. 185 (2014)]

UFO Output

[Degrande et al., Comput. Phys. Commun. 183 (2012)]

Monte Carlo

Sherpa
MadGraph

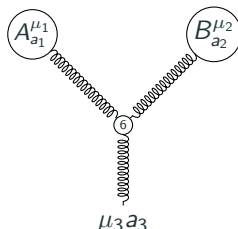
...

Input parameters
Particle spectrum
Vertices
Color structures
Lorentz Structures

Matrix Element Building Blocks

Matrix Elements with Comix/Sherpa

- Non-diagrammatic approach: Berends-Giele type recursions
- Highly efficient for multi-leg amplitudes
- Computes amplitudes in terms of *currents*
- Needs numerical routines for **Lorentz** and **color** structures
- Fully automated for arbitrary (higher dimensional) interactions

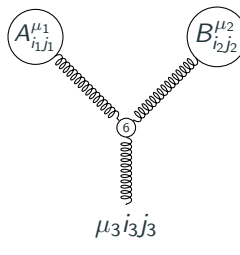

$$= \underbrace{A^{\mu_1}_{a_1} B^{\mu_2}_{a_2} \Gamma^{\mu_1 \mu_2}_{\mu_3}}_{\uparrow} C^{a_1, a_2, a_3}$$

[Höche, SK, Schumann, Siegert, Eur. Phys. J. C 75 (2015)]

Matrix Element Building Blocks

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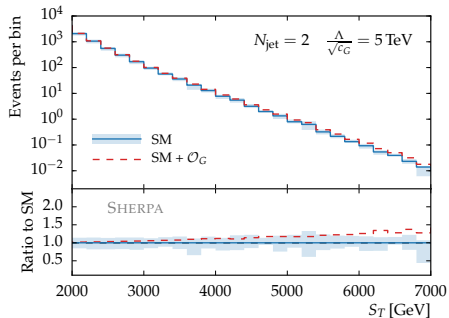
The diagram shows a central vertex labeled with a circled '6'. Three wavy lines (representing gluons) extend from this vertex. The top-left leg is labeled $A_{i_1 j_1}^{\mu_1}$ in a circle. The top-right leg is labeled $B_{i_2 j_2}^{\mu_2}$ in a circle. The bottom leg is labeled $\mu_3 i_3 j_3$.

$$= \underbrace{A^{\mu_1}_{i_1 j_1} B^{\mu_2}_{i_2 j_2} \Gamma^{\mu_1 \mu_2}_{\mu_3}}_{\text{Lorentz and color structures}} \underbrace{C^{a_1, a_2, a_3} [\lambda_{a_1}]_{i_1 j_1} [\lambda_{a_2}]_{i_2 j_2} [\lambda_{a_3}]_{i_3 j_3}}_{\text{New!}}$$

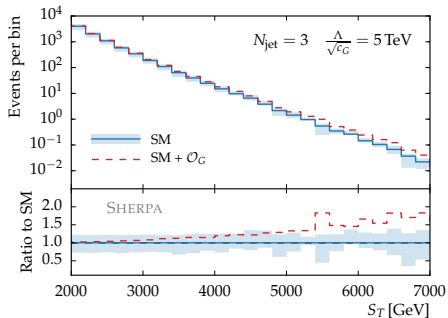
[Höche, SK, Schumann, Siebert, Eur. Phys. J. C 75 (2015)]

Signal Simulation: Results

Signal contributions become sizable for large N_{jet}



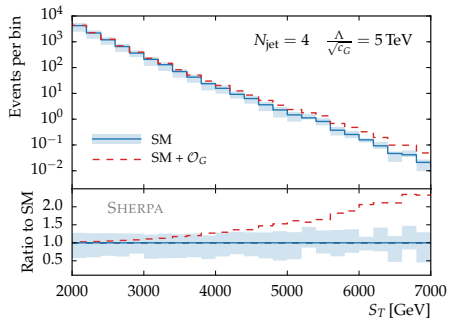
$$N_{\text{jet}} = 2$$



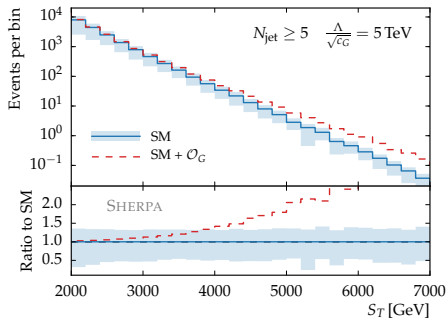
$$N_{\text{jet}} = 3$$

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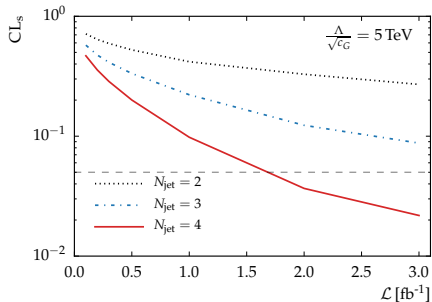
$$N_{\text{jet}} = 4$$



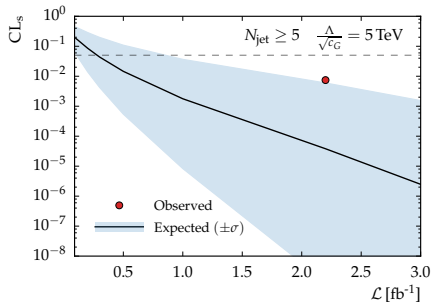
$$N_{\text{jet}} \geq 5$$

Confidence levels for $\Lambda/\sqrt{c_G} = 5\text{TeV}$ hypothesis

Can exclude benchmark point with existing Run-II data



$N_{\text{jet}} = 2, 3, 4$



$N_{\text{jet}} \geq 5$

Higher Powers of Λ^{-1} in $t\bar{t}$

