

Search for Sterile Neutrinos in Tripleton Final States at Hadron Colliders

Kechen Wang

DESY, Hamburg, Germany
CFHEP, IHEP, CAS, Beijing, China

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Based on [Phys. Rev. D **94** (2016) 013005, hep-ph/1701.XXXXX]
with C.S. Kim, C. Dib and J. Zhang

Outline

- ★ Introduction of Sterile Neutrinos
 - ◆ Simplified Model

- ★ Discovering/Excluding N
 - ◆ Dirac
 - ◆ Majorana

- ★ Distinguishing Dirac/Majorana N

- ★ Summary

→ Results @ HL-LHC today
→ Can be extended to 100 TeV, like FCC-hh

Theory Model

Discovery of neutrino oscillations => neutrinos have mass

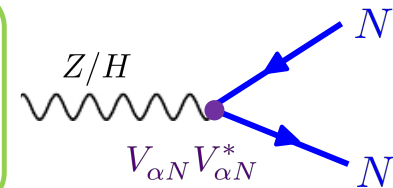
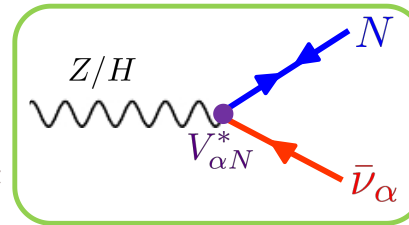
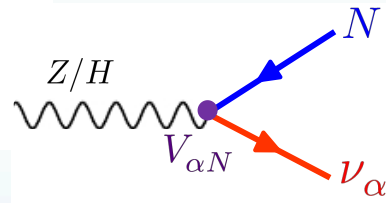
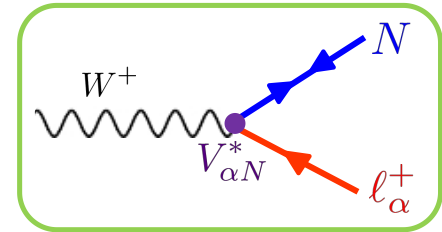
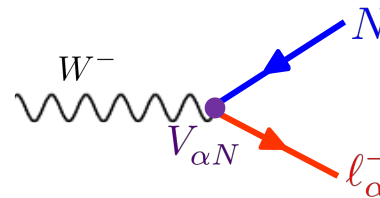
→ In SM, neutrinos are massless

→ A window to BSM physics

Type-I see-saw: Singlet (Sterile) Fermions Interactions: [0901.3589]

$$-\mathcal{L} = h_{\ell\alpha} \bar{L}_\ell \tilde{\Phi} N_\alpha + \frac{1}{2} M_{N\alpha\beta} \bar{N}_\alpha^C N_\beta + \text{H.c.}$$

$$M_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix}$$



Simplified model with assumptions:

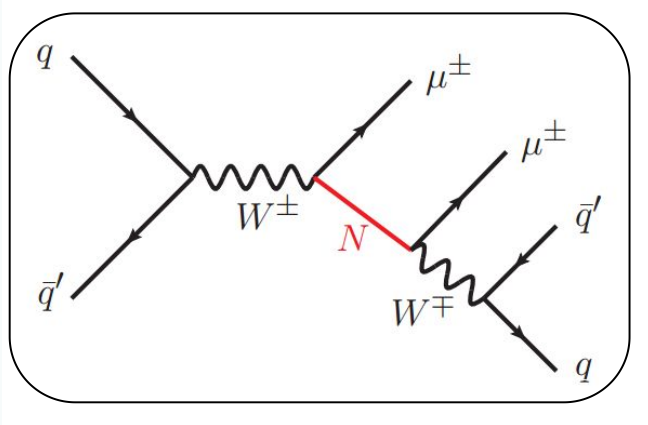
Only 1 generation of sterile neutrinos is light & within experimental reach;

$$U_{NT} = 0;$$

3 free parameters: m_N , U_{Ne} , $U_{N\mu}$, Dirac/Majorana.

Studies @ LHC

Main Search Channels:



2l + 2j

need well isolated energetic 2 jets;
need SS di-lepton to suppress BG;

→ better for Majorana N with $m_N > m_W$.

[CMS: 1207.6079, 1501.05566]
[ATLAS-CONF-2012-139]

Majorana:

$$pp \rightarrow W^\pm \rightarrow l^\pm N \rightarrow l^\pm l^\pm jj \quad (l = e, \mu)$$

Dirac:

$$pp \rightarrow W^\pm \rightarrow l_1^\pm N \rightarrow l_1^\pm l_2^\mp jj \quad (l_{1,2} = e, \mu)$$

3l + MET

→ better for Majorana or Dirac N with $m_N < m_W$

$m_N < m_W$: [1504.02470]
 $m_N > m_W$: [0809.2096,
0910.2720, 1112.6419 ...]

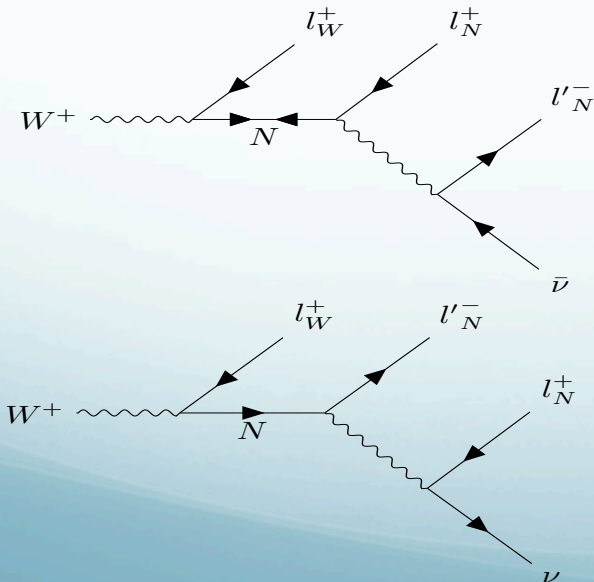
$$\text{LNV: } W^+ \rightarrow e^+ e^+ \mu^- \bar{\nu}_\mu$$

non-trivial ←
flavor of ν undetectable

distinguishable?

Today !

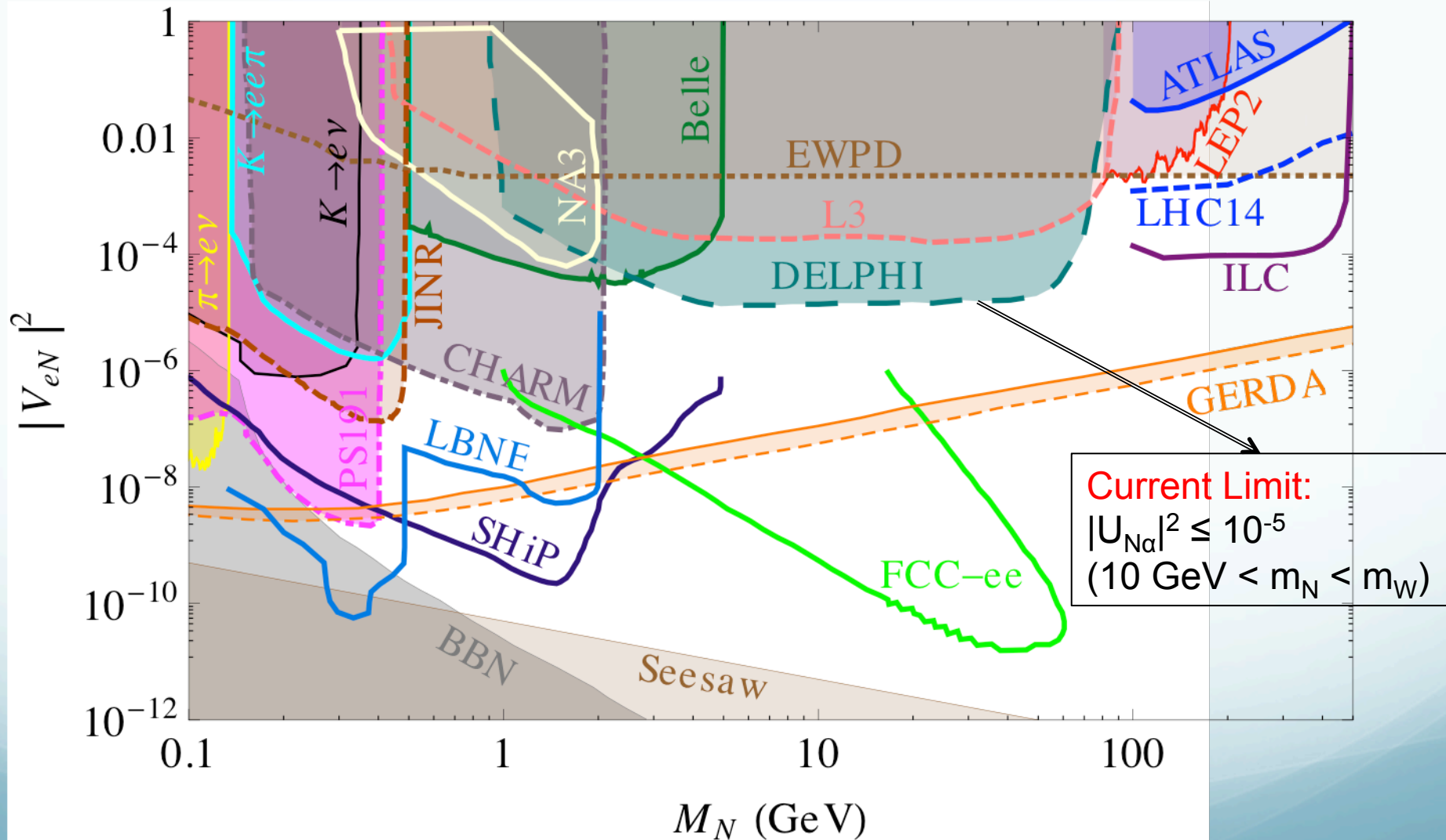
$$\text{LNC: } W^+ \rightarrow e^+ e^+ \mu^- \nu_e \quad [0809.2096, 1509.05981]$$



Global Constraints

from [Deppisch, Dev and Pilaftsis, New J. Phys. 17 (2015) 085019]

m_N : 0.1 ~ 500 GeV



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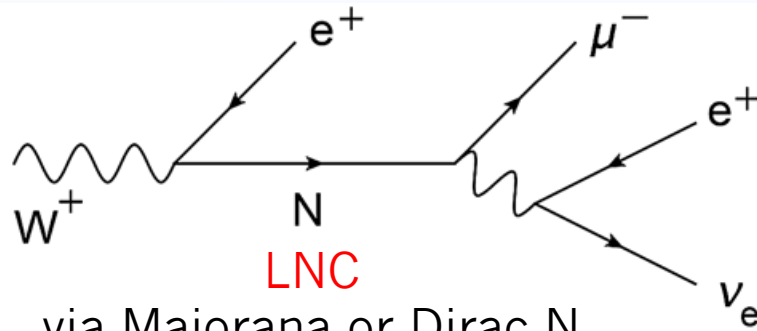
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 - ◆ Dirac
 - ◆ Majorana

- ★ Distinguishing Dirac/Majorana N

- ★ Summary

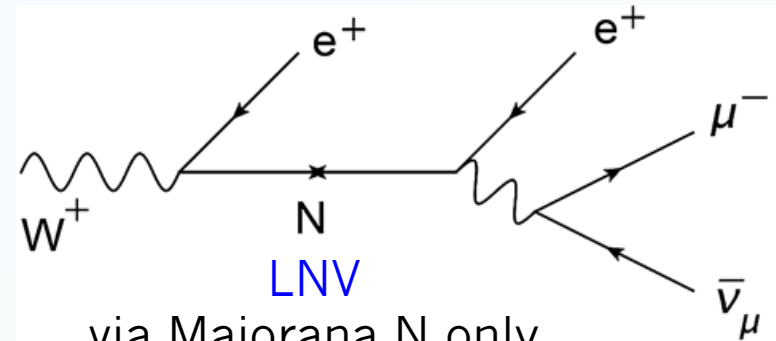
Production Rates

tri-lepton + MET with no-OSSF lepton pairs



via Majorana or Dirac N

$$\text{Br}(W^+ \rightarrow e^+ e^+ \mu^- \nu_e) \propto \frac{|U_{Ne} U_{N\mu}|^2}{|U_{Ne}|^2 + |U_{N\mu}|^2}$$



via Majorana N only

$$\text{Br}(W^+ \rightarrow e^+ e^+ \mu^- \bar{\nu}_\mu) \propto \frac{|U_{Ne}|^4}{|U_{Ne}|^2 + |U_{N\mu}|^2}$$

Scale factors for different tri-lepton states

	Dirac (LNC)	Majorana (LNC+LNV)
$e^+ e^+ \mu^- \nu$	s	s (1 + r)
$\mu^+ \mu^+ e^- \nu$	s	s (1 + $\frac{1}{\mathbf{r}}$)

normalization factor

$$s \equiv 2 \times 10^6 \times \frac{|U_{Ne} U_{N\mu}|^2}{|U_{Ne}|^2 + |U_{N\mu}|^2}$$

disparity factor $r \equiv \frac{|U_{Ne}|^2}{|U_{N\mu}|^2}$

For benchmark point

$$|U_{Ne}|^2 = |U_{N\mu}|^2 = 10^{-6} \rightarrow r = s = 1$$

Collider Simulation

Simulation

MadGraph (jet matching up to 2 extra partons) + **PYTHIA** + **Delphes**

Signal:

tri-lepton + MET with no OSSF lepton pairs
 $e^+ e^+ \mu^- / \mu^+ \mu^+ e^- / e^- e^- \mu^+ / \mu^- \mu^- e^+ + \text{MET}$.

SM backgrounds:

→ **Leptonic τ decay:**

WZ $\rightarrow (l \nu) (\tau \tau) \rightarrow 3 l + \text{MET}$

→ **Fake leptons** from jets containing heavy-flavor mesons:

γ^*/Z +jets: $\gamma^*/Z \rightarrow \tau \tau$ + a 3rd faked lepton

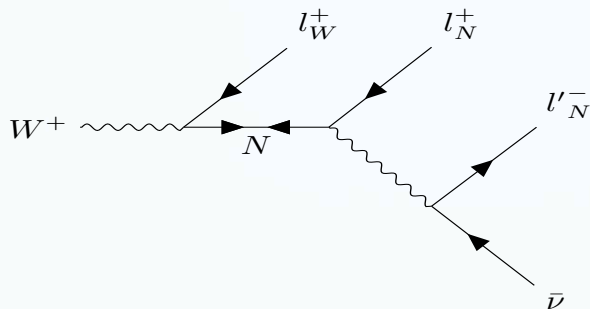
t tBar+jets: prompt decay of t tBar + a 3rd fake lepton

Strategy

Apply various cuts to reduce BG

→ 3 leptons $l^\pm l^\pm l'^\mp$, veto b-jets.

→ Cut-and-Count or MVA



correct lep from N & $p_{z,v}$

← by minimizing the

$$\chi^2 = \left(\frac{M_W - m_W}{\sigma_W} \right)^2 + \left(\frac{M(l^\pm l'^\mp \nu)}{\sigma_N} \right)^2$$

The equation shows the chi-squared function used for lepton identification. The first term is the W mass constraint, and the second term is the neutrino mass constraint. Blue arrows point from the terms to their respective mass variables: $M(l^\pm l'^\mp \nu)$ and $M(l'_N)$.

MVA input **observables**:

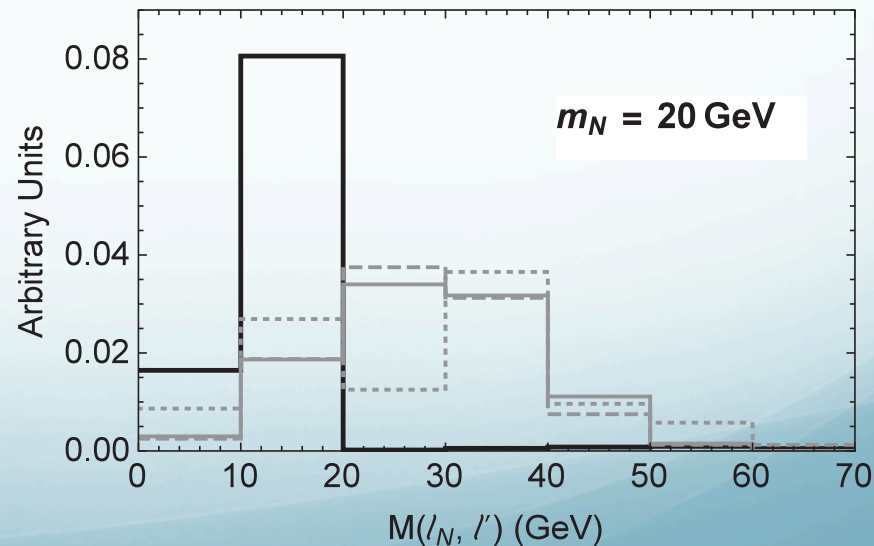
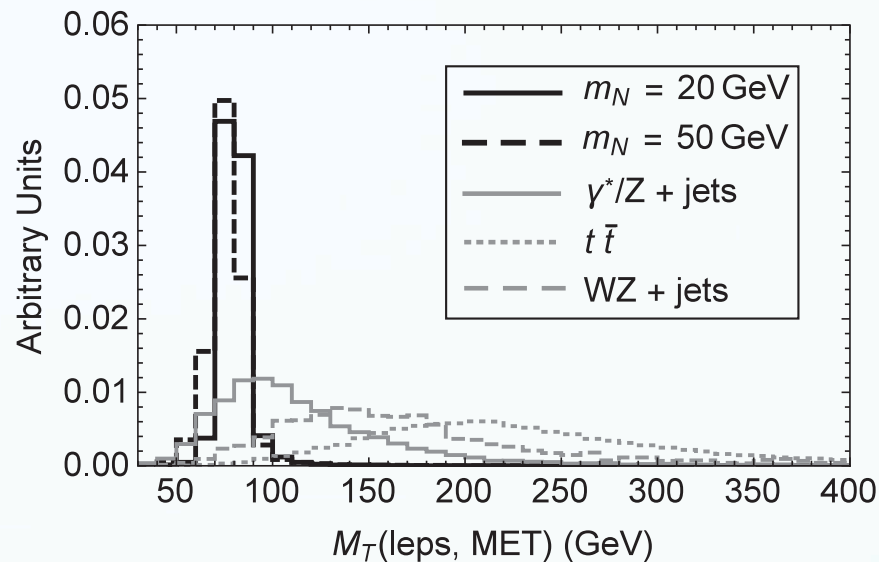
→ met, H_T ;

→ $M_T(\text{met}, \text{lep}(s))$;

→ $\Delta\phi(\text{met}, \text{lep}(s))$;

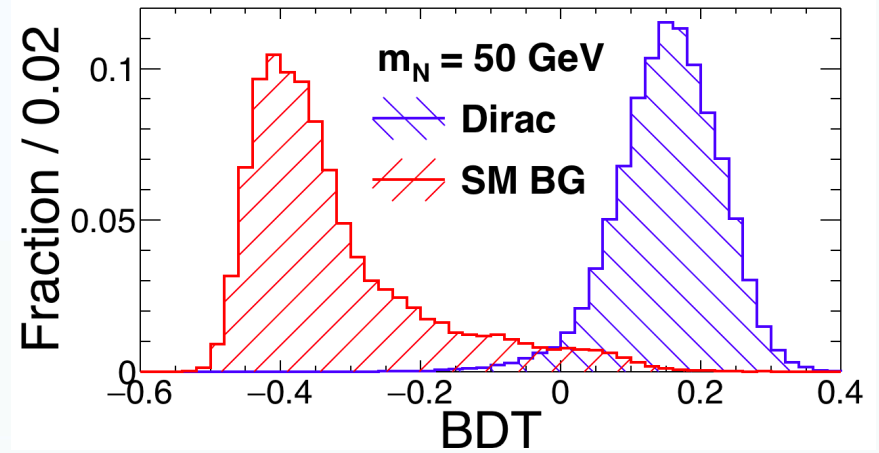
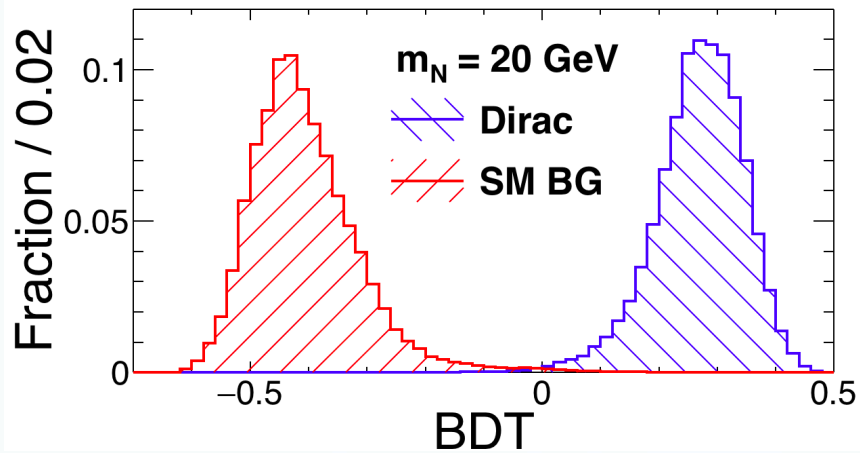
→ $M(\text{leps})$;

→ $\Delta\phi(\text{lep}, \text{lep})$.



Cut Flow Tables

BDT from TMVA package

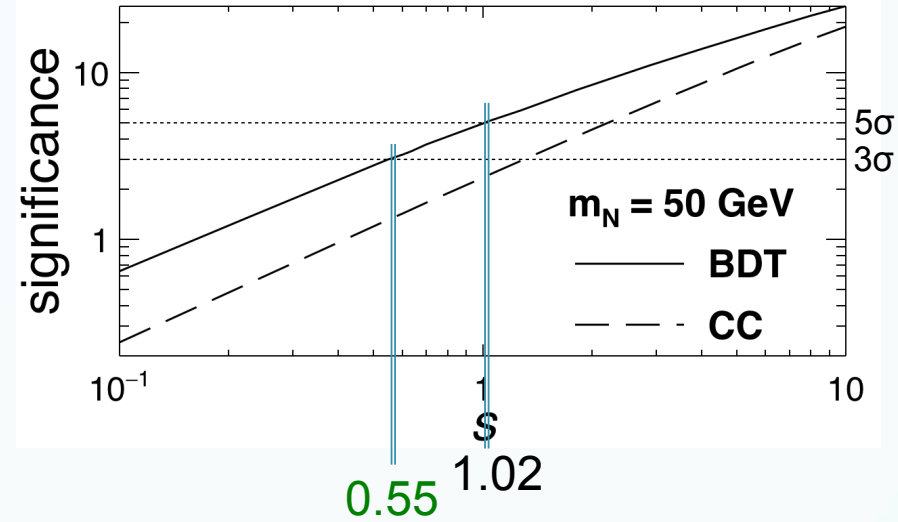
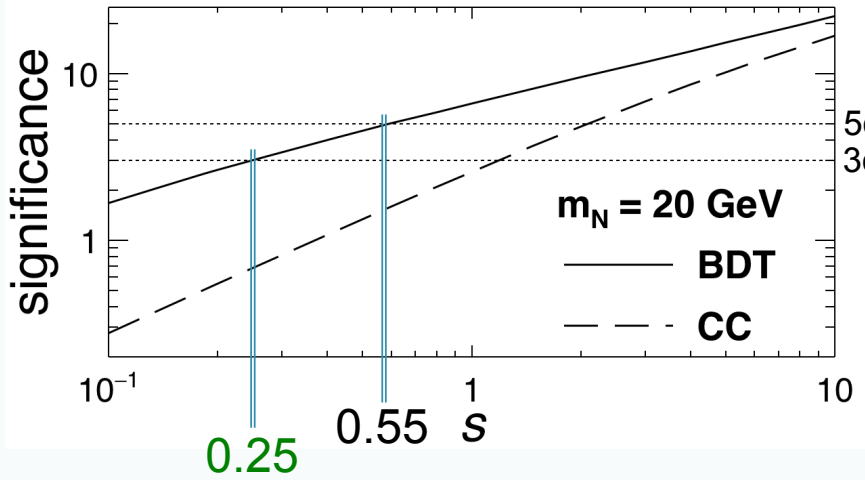


of events, 14 TeV, 3000 fb⁻¹ $m_N = 20$ GeV

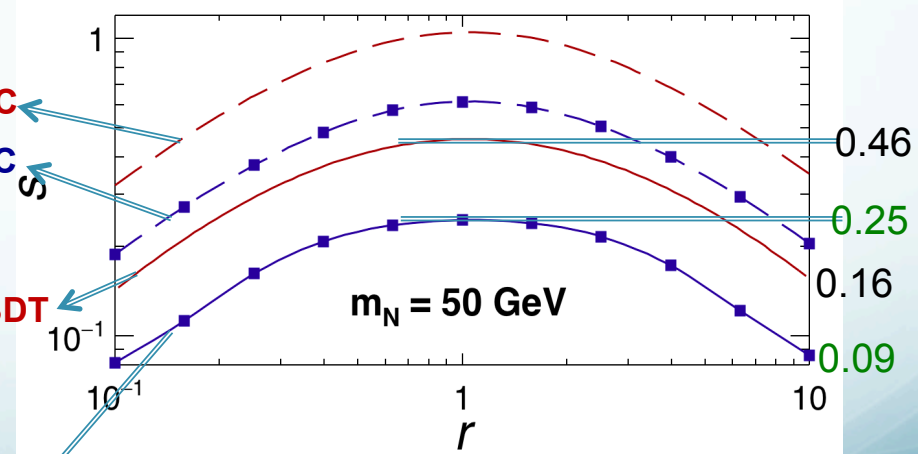
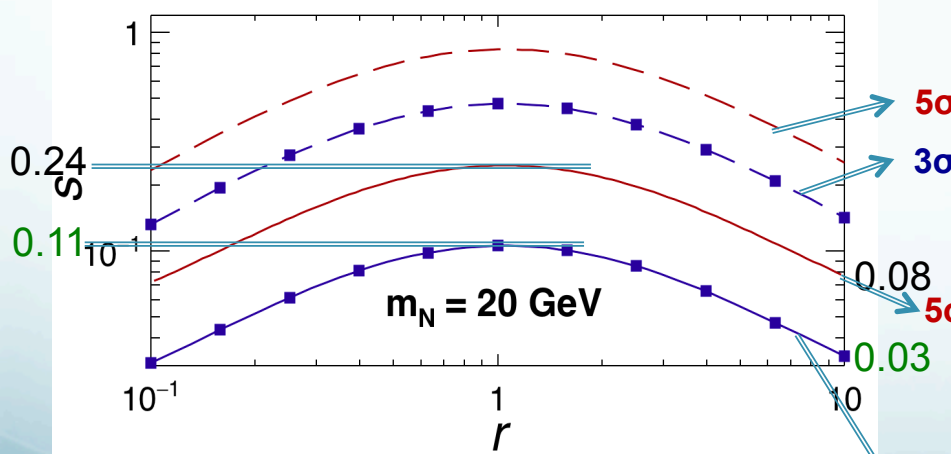
Cuts	Dirac	Majorana	γ^*/Z	WZ	$t\bar{t}$	SS
Basic cuts	54.0	133.2	4220	2658	68588	
N(b-jets)=0	53.1	131.1	4063.0	2497.1	31953.5	
CC	44.2	110.9	209.8	25.3	16.9	2.6 (5.8)
BDT > 0.1825	46.7	-	1.9	1.3	0.0	6.6
BDT > 0.1705	-	120.7	5.1	1.7	0.8	10.7

Limits

Dirac N



Majorana N



$$|U_{Ne}|^2 = \frac{1}{2 \times 10^6} \times (1 + r) \times s$$

$$|U_{N\mu}|^2 = \frac{1}{2 \times 10^6} \times \left(1 + \frac{1}{r}\right) \times s$$

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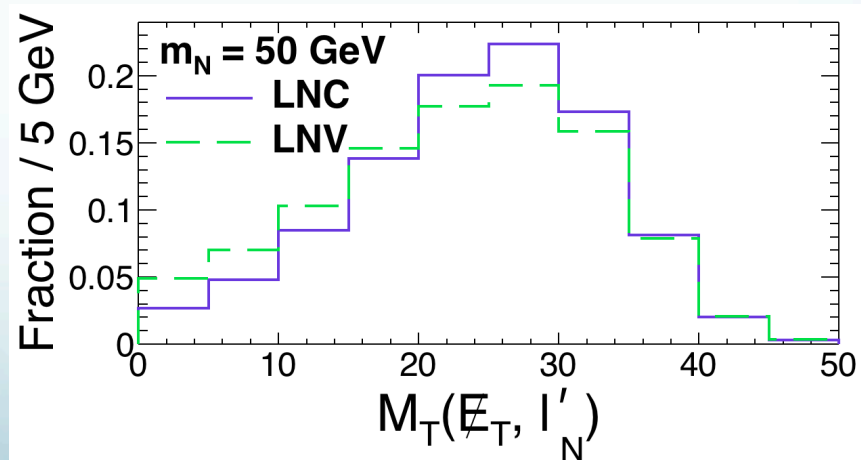
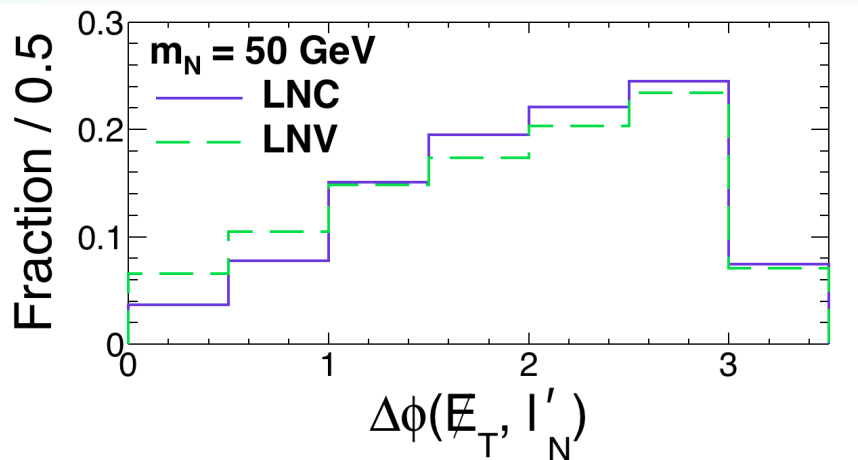
Basic Idea

(1) **1st BDT** -> reduce SM BG

of events, 14 TeV, 3000 fb⁻¹ $m_N = 20$ GeV

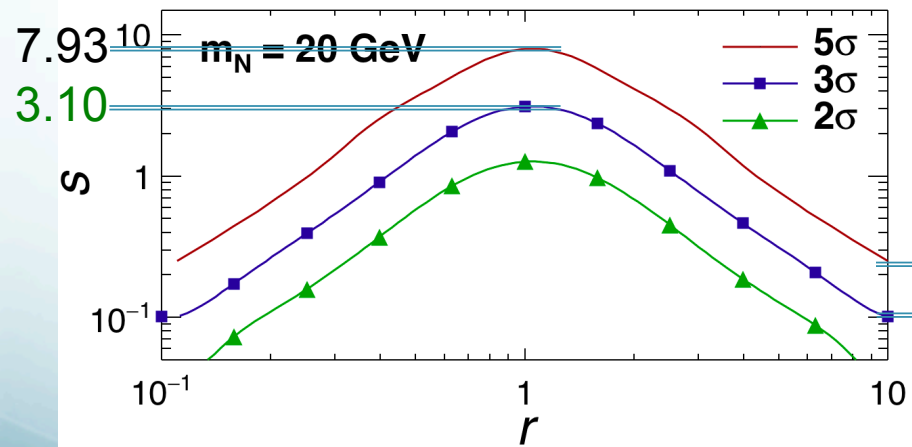
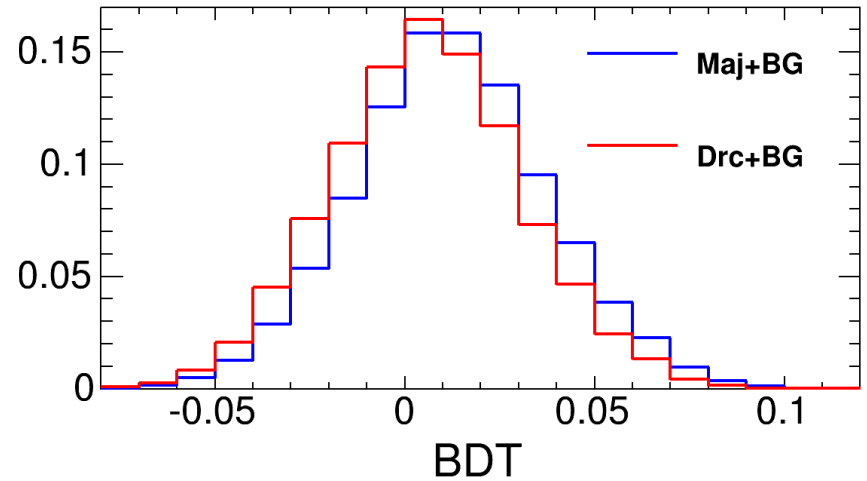
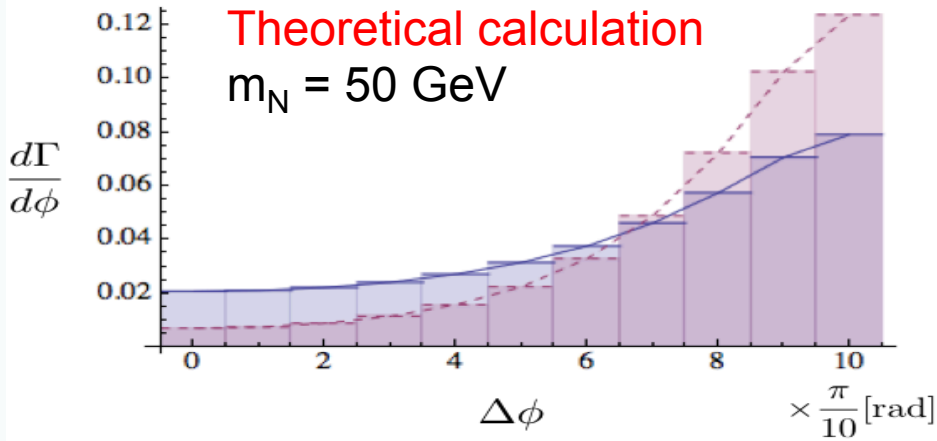
Cuts	$e^+e^+\mu^-$		$\mu^+\mu^+e^-$		$e^-e^-\mu^+$		$\mu^-\mu^-e^+$		$l^\pm l^\pm l'^\mp$	$l^+l^+l'^-$	$l^-l^-l'^+$	$l^\pm l^\pm l'^\mp$
	LNC	LNV	LNC	LNV	LNC	LNV	LNC	LNV	γ^*/Z	W^+Z	W^-Z	$t\bar{t}$
Basic cuts	13.6	19.5	15.0	22.0	12.1	18.2	13.3	19.5	1055.0	779.0	550.0	17147.0
N(b-jets)=0	13.4	19.2	14.7	21.7	11.9	17.9	13.1	19.2	1015.8	731.8	516.7	7988.4
BDT1 > 0.1709	12.2	17.7	13.5	20.0	10.9	16.5	12.0	17.7	1.2	0.5	0.4	0.2

(2) **2nd BDT**: kinematical distributions differing between LNC & LNV, $M_T(\text{met}, \text{lep}(s))$ & $\Delta\phi(\text{met}, \text{lep}(s))$

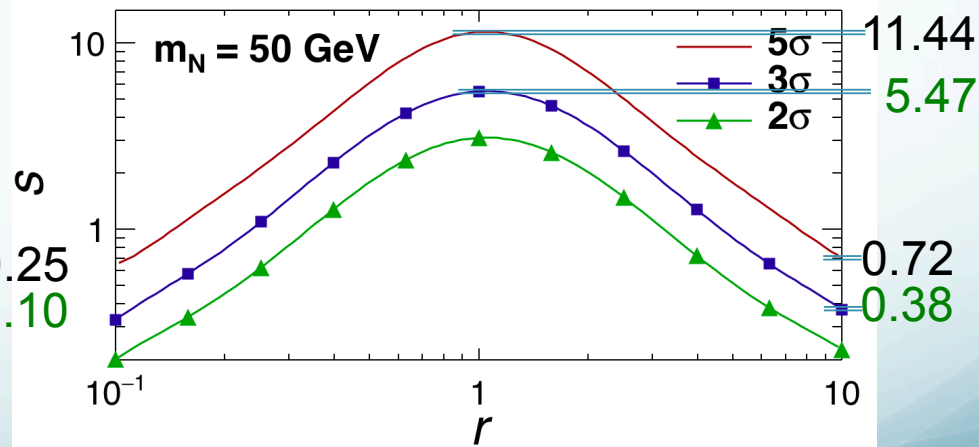


Limits

$m_N = 20 \text{ GeV}$



$$|U_{Ne\ell}|^2 = \frac{1}{2 \times 10^6} \times (1 + r) \times s$$



$$|U_{N\mu}|^2 = \frac{1}{2 \times 10^6} \times \left(1 + \frac{1}{r}\right) \times s$$

Summary

- ★ A complete search strategy for N from $W^\pm \rightarrow e^\pm e^\pm \mu^\mp \nu / \mu^\pm \mu^\pm e^\mp \nu$
 - $m_N = 20, 50$ GeV @ 14 TeV LHC, 3000 fb⁻¹;
 - CC & MVA;
 - MVA greatly enhance the limits;
 - limits for 20 GeV better.

- ★ Discovering / Excluding N
 - Dirac: $s \geq 0.25$ @ 3 σ ;
 - Majorana $r \sim 1$: $s \geq 0.11$ @ 3 σ
 $r \sim 10$, limits on $|U_{N\mu}|^2$ stronger.

$$|U_{Ne}|^2 = \frac{1}{2 \times 10^6} \times (1 + r) \times s$$

$$|U_{N\mu}|^2 = \frac{1}{2 \times 10^6} \times \left(1 + \frac{1}{r}\right) \times s$$

- ★ Discriminating Dirac / Majorana
 - using kinematical distri. & MVA
 - non-trivial, flavor of ν undetectable;
 - $r \sim 1$: $s \geq 3.10$ @ 3 σ ;
 - $r \sim 10$, limits on $|U_{N\mu}|^2$ stronger.
- ★ Analysis can be extended to 100 TeV pp.

The background features several faint, overlapping postmarks and stamps in a light brown or tan color. These include circular postmarks with dates and times, and rectangular stamps with text like 'POSTAGE' and 'AIR MAIL'. The overall aesthetic is that of an old, weathered document or envelope.

Thank you for your attention !

Any Questions ?

Backup Slide

Fake-lepton Simulation

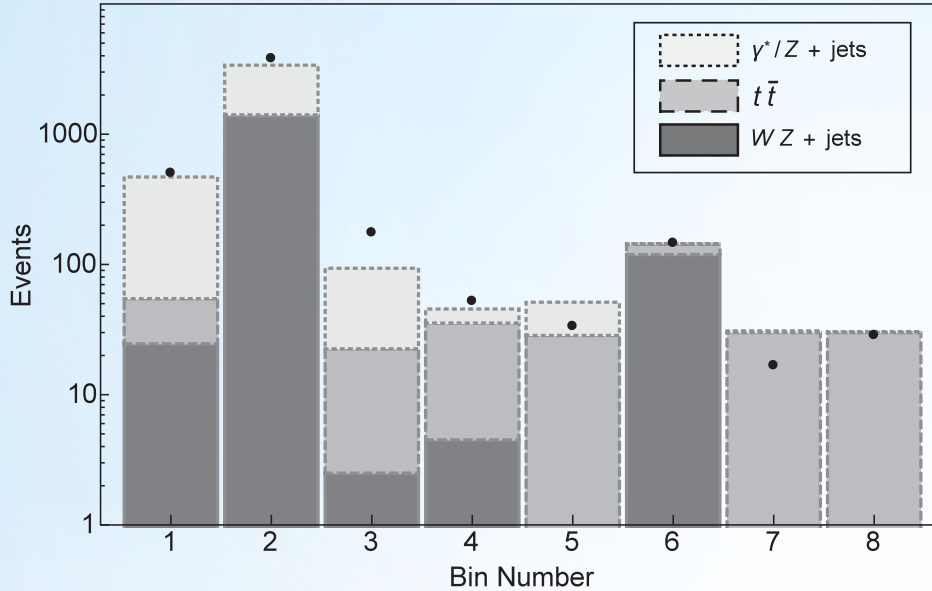


FIG. 6. Validation results for fake lepton simulation. Black dots indicate experimental results in Ref. [31]. Our simulated results for $\gamma^*/Z + \text{jets}$, $t\bar{t}$, and $WZ + \text{jets}$ are given by upper light gray bars, middle brown bars, and bottom pink bars, respectively. Eight bin categories are (1) 0-bjet, 1-OSSF, $M_{\ell^+, \ell^-} < 75$ GeV, (2) 0-bjet, 1-OSSF, $|M_{\ell^+, \ell^-} - M_Z| < 15$ GeV, (3) 0-bjet, 1-OSSF, $M_{\ell^+, \ell^-} > 105$ GeV, (4) 0-bjet, 0-OSSF, (5–8) are the same as the first four bins, but with at least one b-jet.

A pheno. FL simulation method

[E. Izaguirre, B. Shuve, hep-ph/1504.0247]

→ data-driven methods to estimate the fake lepton contributions

→ modeling parameters, pinned down by validating simulated results against actual experimental ones.

1. Mistag rate

(probability of converting a jet to a lepton)

$$\epsilon_{j \rightarrow \ell}(p_{Tj}) = \epsilon_{200} \left[1 - (1 - r_{10}) \frac{200 - p_{Tj}/\text{GeV}}{200 - 10} \right]$$

2. Transfer function

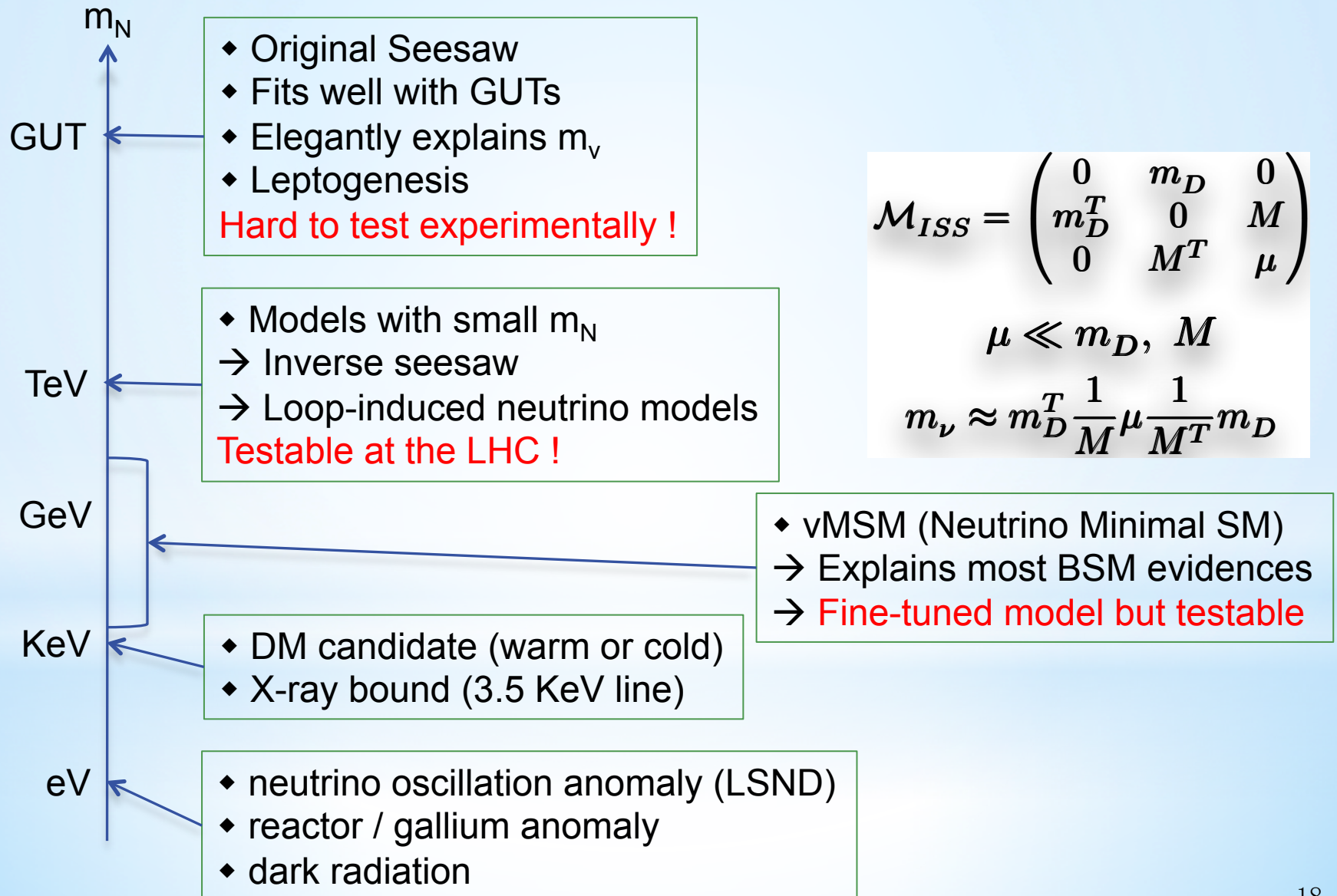
(how much p_T is transferred into the lepton)

$$p_{T\ell} \equiv (1 - \alpha)p_{Tj}$$

$$\mathcal{T}_{j \rightarrow \ell}(\alpha) = \frac{1}{\mathcal{N}} \exp \left[-\frac{(\alpha - \mu)^2}{2\sigma^2} \right]$$

Backup Slide

Interesting Mass Scales of m_N



$$\mathcal{M}_{ISS} = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & 0 & M \\ 0 & M^T & \mu \end{pmatrix}$$

$$\mu \ll m_D, M$$

$$m_\nu \approx m_D^T \frac{1}{M} \mu \frac{1}{M^T} m_D$$

Backup Slide

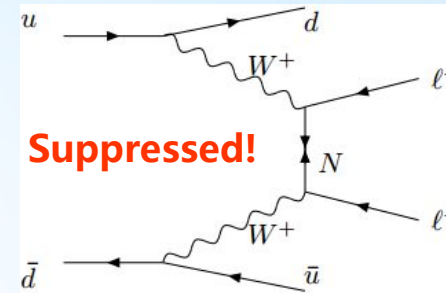
Productions @ pp Colliders

$$q\bar{q} \rightarrow Z^{(*)} \rightarrow \nu N$$

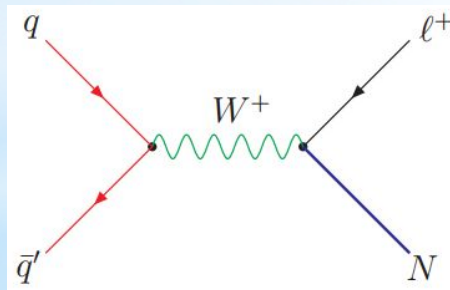
$$gg \rightarrow H^{(*)} \rightarrow \nu N$$

almost unobserved

(final states l^+l^- , l^\pm suffer from huge background)

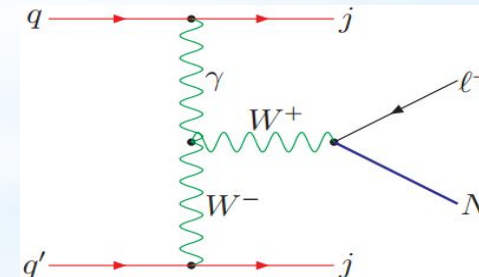


(no resonance enhancement)



Mostly studied

(important for $m_N < 1$ TeV)



More important for $m_N > 1$ TeV

Backup Slide

from [S. Antusch, E. Cazzato, O. Fischer, hep-ph/1612.02728]

