

Strangeness enhancement from pp to AA

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in collaboration with

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Trying to understand (based on EPOS)

□ **Particle ratios vs $\left\langle \frac{dn_{\text{ch}}}{d\eta}(0) \right\rangle$ for pp, pPb, (PbPb)**

□ **Average transverse momenta vs $\left\langle \frac{dn_{\text{ch}}}{d\eta}(0) \right\rangle$ for pp, pPb, (PbPb)**

□ **Charmed meson production vs $\frac{dn_{\text{ch}}}{d\eta}(0)$ for pp**

$\left\langle \frac{dn_{\text{ch}}}{d\eta}(0) \right\rangle$ for multiplicity classes defined via forward multiplicities

First step to get a global view, to see where EPOS works and where not, in pp, pA, AA, same version

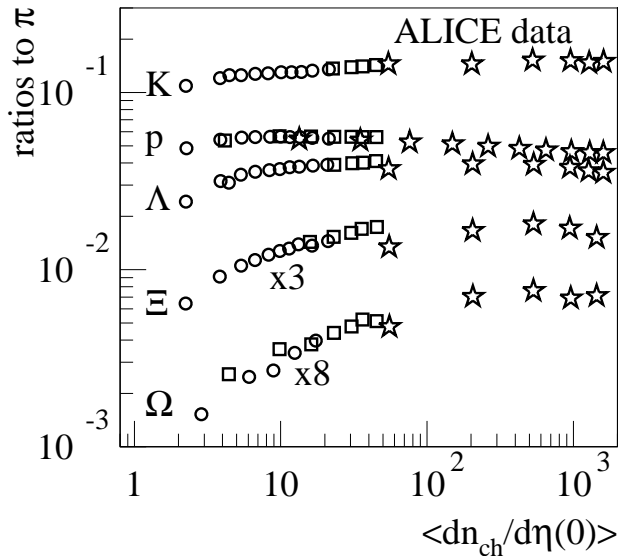
Status 2015: Two parallel developments

EPOS LHC: Gribov Regge approach, parameterized flow as in EPOS1.99, tuned to LHC data (2012), **very much used (and tested) by LHC pp groups, UE, forward physics etc, and used for air shower simulations**

EPOS 3.0xx: Gribov Regge approach, viscous hydro, parton saturation, **mainly used for HI and collectivity in pp**

2015/2016/2017: “Fusion”, to accommodate basic pp and HI features, public version;
Currently: EPOS3.2xx

Particle ratios to pions vs $\left\langle \frac{dn_{ch}}{d\eta}(0) \right\rangle$



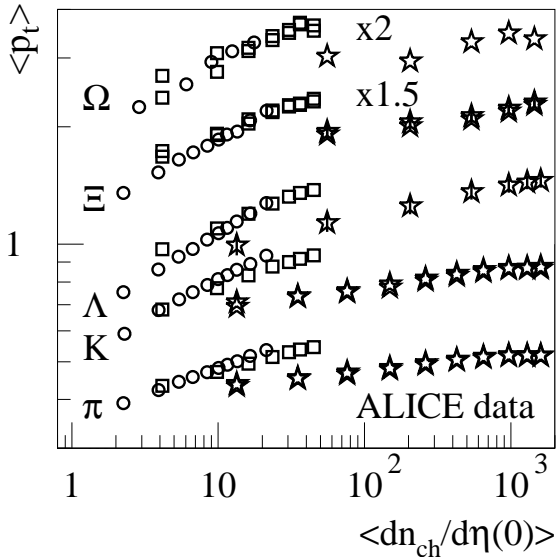
circles = pp (7TeV)

squares = pPb (5TeV)

stars = PbPb (2.76TeV)

Refs: next slide

Mean p_t vs $\left\langle \frac{dn_{ch}}{d\eta}(0) \right\rangle$



circles = pp (7TeV)

squares = pPb (5TeV)

stars = PbPb (2.76TeV)

Data partly collected by A. G. Knospe

Refs:

$\langle dn_{ch}/d\eta \rangle$ in Pb+Pb: Phys. Rev. Lett. 106 032301 (2011)
 π^+ , K^+ , and (anti)protons in Pb+Pb: Phys. Rev. C 88 044910 (2013)

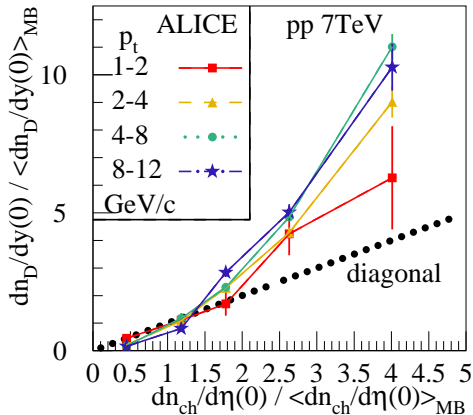
Λ in Pb+Pb: Phys. Rev. Lett. 111 222301 (2013)
 Ξ^- and Ω in p+Pb: Phys. Lett. B 758 389-401 (2016)
 π^+ , K^+ , (anti)protons, and Λ in p+Pb: Phys. Lett. B 728 25-38 (2014)

$\langle dn_{ch}/d\eta \rangle$ in p+Pb: Eur. Phys. J. C 76 245 (2016)
 Ξ^- and Ω in p+Pb: Phys. Lett. B 758 389-401 (2016)
 $\langle dn_{ch}/d\eta \rangle$ in p+p 7 TeV: Eur. Phys. J. C 68 345-354 (2010)

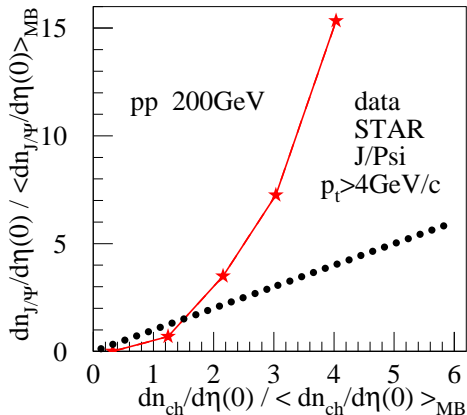
π^+ , K^+ , and (anti)protons in p+p 7 TeV: Eur. Phys. J. C 75 226 (2015)

Ξ^- and Ω in p+p 7 TeV: Phys. Lett. B 712 309 (2012)
 and data points from Rafael Derradi de Souza, SQM2016

D or J/ Ψ multiplicity vs $\frac{dn_{ch}}{d\eta}(0)$ in pp



ALICE JHEP 09 (2015) 148,
arXiv:1505.00664v1

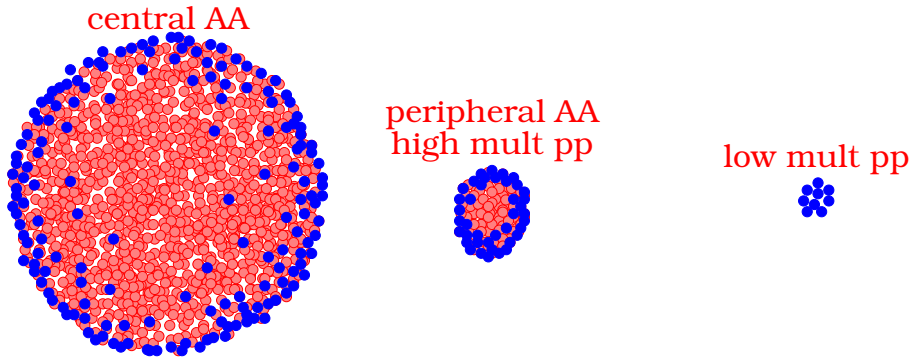


STAR, shown at MPI2016

strongly nonlinear increase

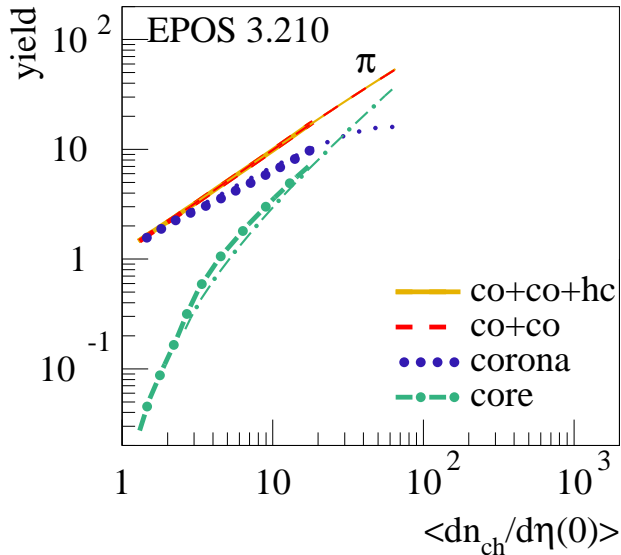
Core-corona picture in EPOS (details later)

Gribov-Regge approach => (Many) kinky strings
=> core/corona separation (based on string segments)



core => hydro => statistical decay ($\mu = 0$)
corona => string decay

Pion yields: core / corona contribution

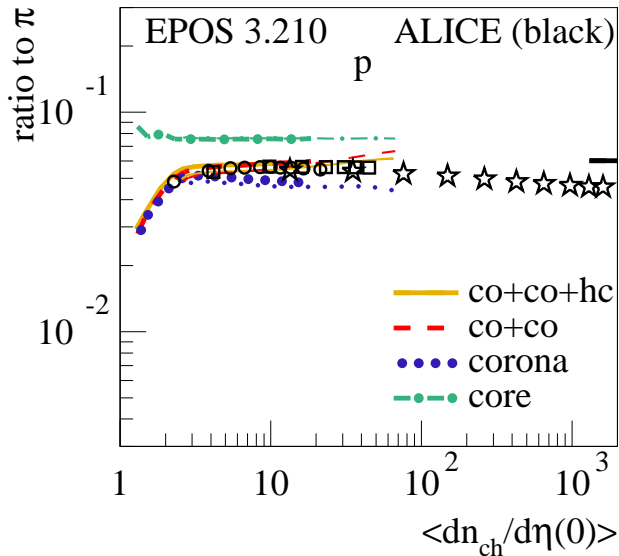


thick lines
= pp (7TeV)

thin lines
= pPb (5TeV)

hc = hadronic cascade
(UrQMD)

Proton to pion ratio



core hadronization:

$T = 164 \text{ MeV}, \mu_B = 0$

statistical model fit

(horizontal black line)

A. Andronic et al.,

arXiv:1611.01347

$T = 156.5 \text{ MeV}, \mu_B = 0.7 \text{ MeV}$

thick lines = pp (7TeV)

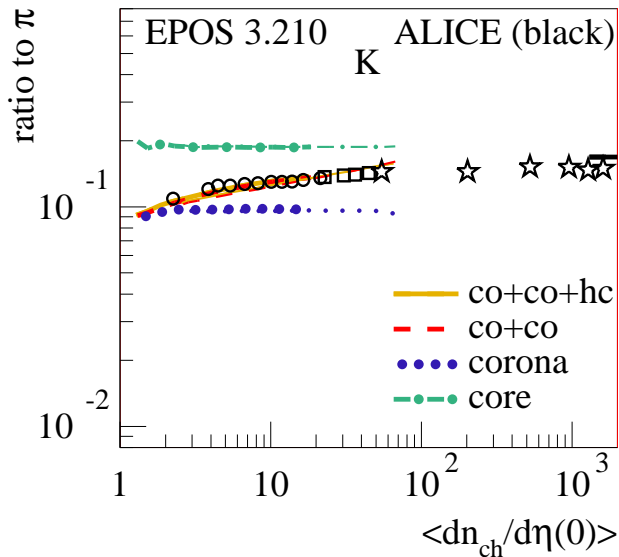
thin lines = pPb (5TeV)

circles = pp (7TeV)

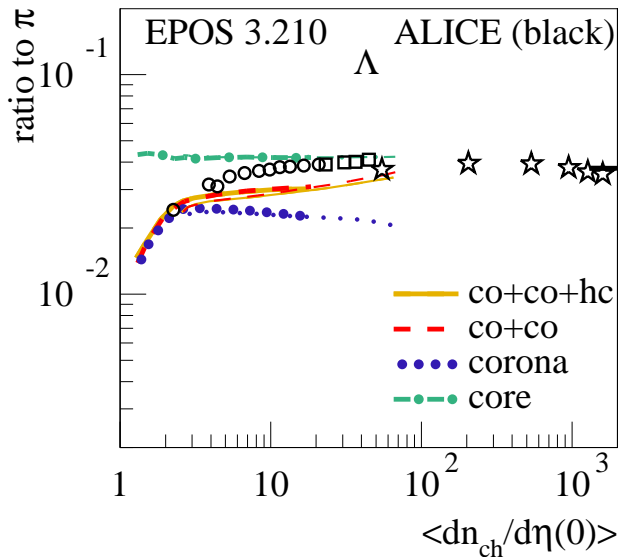
squares = pPb (5TeV)

stars = PbPb (2.76TeV)

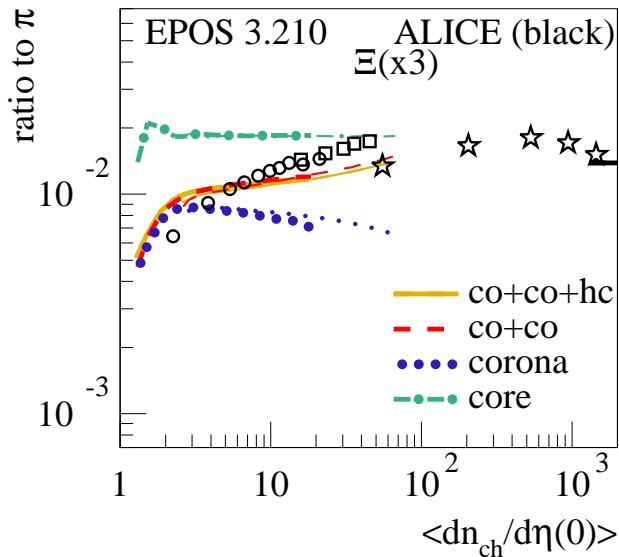
Kaon to pion ratio



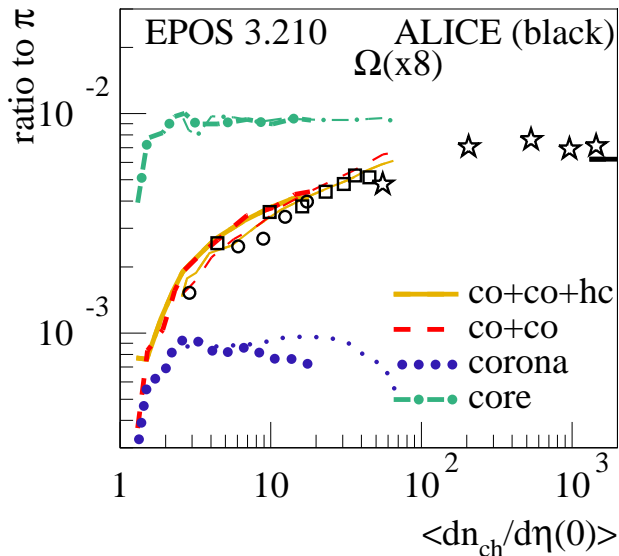
Lambda to pion ratio



Xi to pion ratio



Omega to pion ratio



thick lines = pp (7TeV)
 thin lines = pPb (5TeV)
 circles = pp (7TeV)
 squares = pPb (5TeV)
 stars = PbPb (2.76TeV)

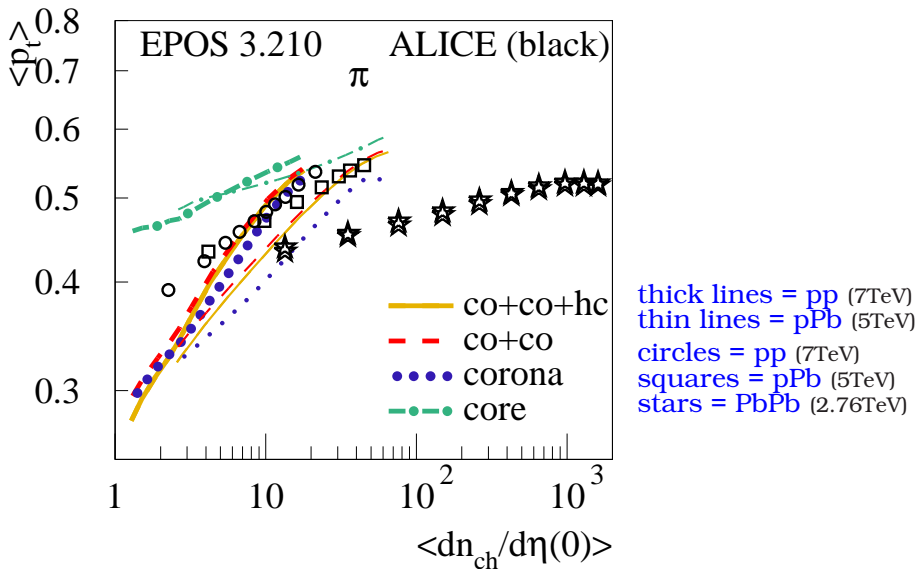
Ratios h/π for $h = p, K, \Lambda, \Xi, \Omega$ vs $\left\langle \frac{dn}{d\eta}(0) \right\rangle$:

Core and corona contributions separately roughly constant

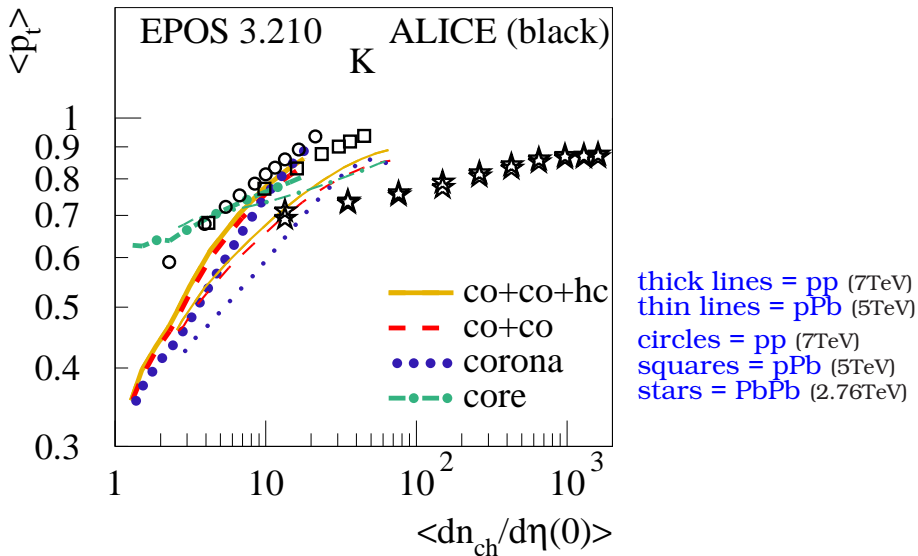
Difference (core - corona) increasing for $p \rightarrow K \rightarrow \Lambda \rightarrow \Xi \rightarrow \Omega$

**=> increasing slope
(not enough for Ξ, Ω)**

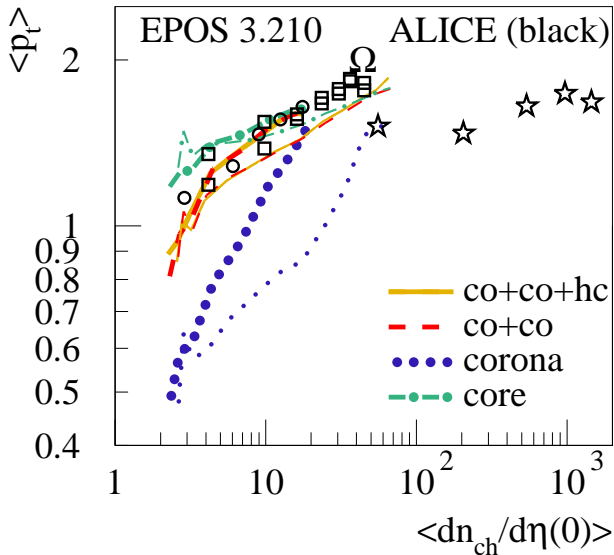
Average p_t of pions



Average p_t of kaons



Average p_t of Omegas



Average p_t of $\pi, K, (p, \Lambda, \Xi), \Omega$ vs $\left\langle \frac{dn}{d\eta} \right\rangle_{(0)}$:

Moderate increase of core contribution
(same for pp and pPb, similar to PbPb)

Strong increase of corona contribution
(stronger for pp compared to pPb)

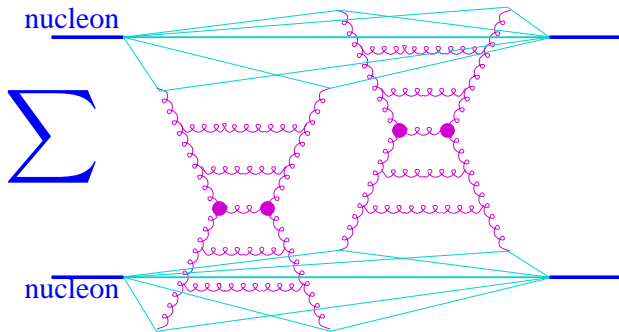
Slope(pp) > slope(pPb) >> slope(PbPb)

The multiplicity dependence of the corona contribution is crucial

Presently: Corona mean p_t too small at small multiplicity

Why such a strong mean pt increase with multiplicity for corona particles?

EPOS: Gribov-Regge approach



S-Matrix based on Pomerons

Pomerons :
Parton ladders (initial and final state radiation, DGLAP)

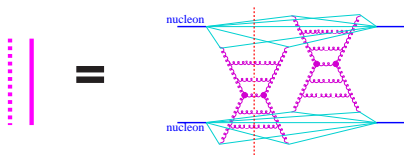
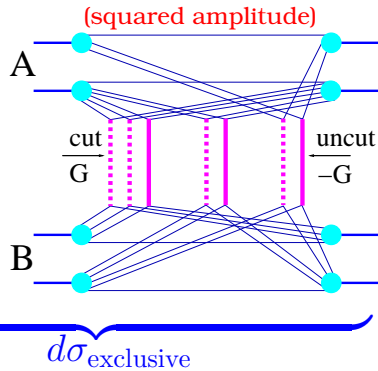
Cutting rules to get inelastic cross sections.

Same principle for AA

Explicite formulas for cross sections

(even partial cross sections)

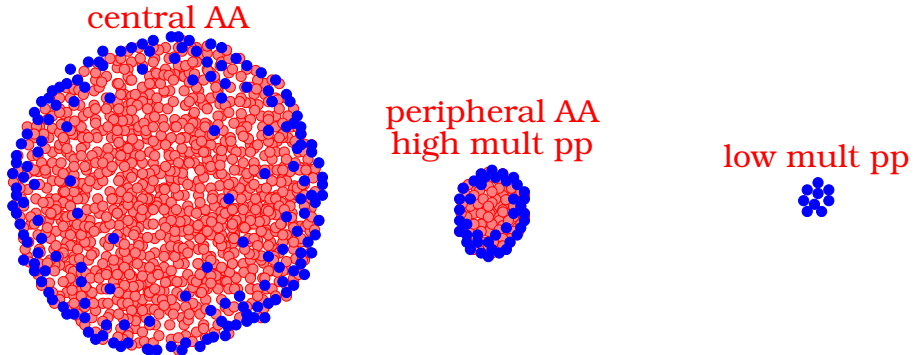
$$\sigma^{\text{tot}} = \sum_{\text{cut } P} \int \sum_{\text{uncut } P} \int$$



=> kinky strings



Based on string segments: core-corona separation



core: string segments “melt” => fluid

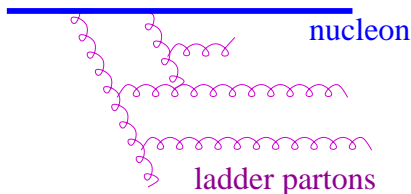
corona: strings survive (ordinary kinky strings from parton ladders)

Parton-ladders⁽¹⁾ are perfectly fitted⁽²⁾ as $G = \alpha (x^+ x^-)^\beta$.

G depends on the virtuality cutoff: $G = G(Q_0)$.

To mimic the effects of gluon fusion, the fits are modified (for pp) as $\alpha (x^+ x^-)^{\beta+\varepsilon}$, referred to as G_{eff} .

The exponent $\varepsilon = \varepsilon(s)$ is chosen to reproduce the energy dependence of cross sections.



Procedure employed in EPOS LHC

(1) Imaginary part G of the corresponding amplitude in b -space

(2) x^+, x^- : light cone momentum fractions of the Pomeron end

But adding an exponent ε

- **must be accompanied by a corresponding modification of the internal structure of the Pomeron**

This can be done by defining a **saturation scale** Q_s via

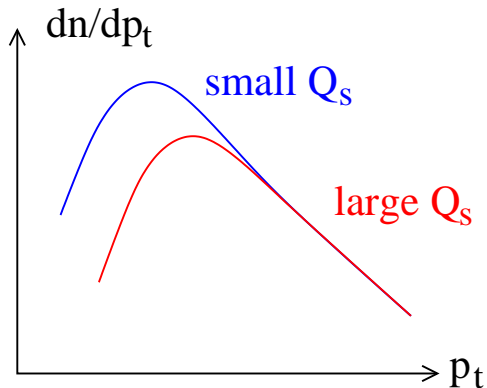
$$G_{\text{eff}} = A (N_{\text{Pom}})^B G(Q_s)$$

and then considering the parton ladder with the cutoff Q_s (thus changing the internal structure! => consistent!)

We find

$$Q_s = Q_s(x^+ x^-) \propto (x^+ x^-)^{0.30}$$

Parton distributions

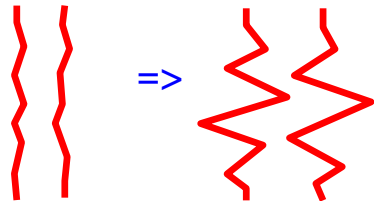


Increasing $\langle dn/d\eta(0) \rangle$
 corresponds to increasing N_{Pom}

=> Increasing Q_s

=> harder Pomerons

=> harder strings



=> more high p_t particles

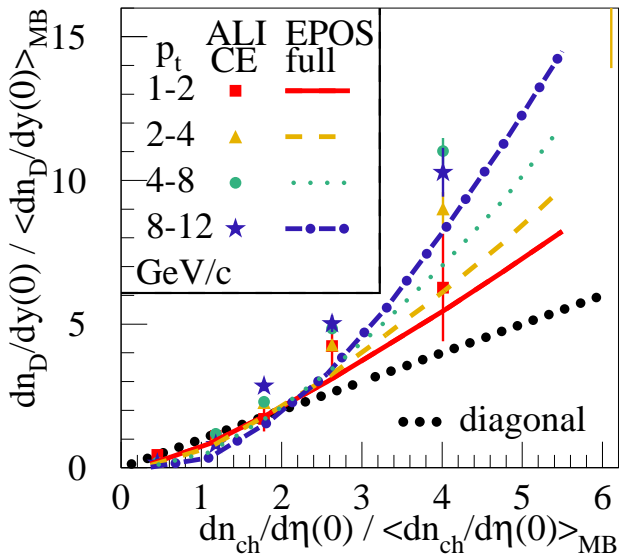
=> Strong increase of $\langle p_t \rangle$ with $\langle dn/d\eta(0) \rangle$

Very closely related to this discussion:

**The multiplicity dependence
of charm production (D , J/Ψ , ...)**

**The “ultimate tool” to test multiple
scattering (and the implementation
of Q_S)**

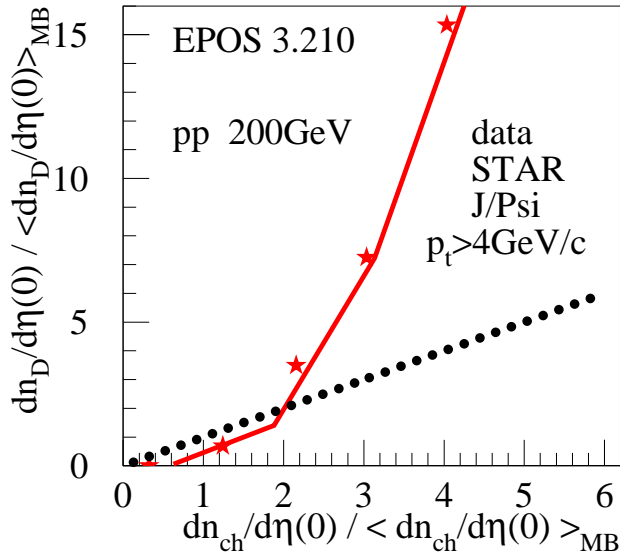
EPOS 3 compared to ALICE data



hadronic cascade
on/off
has no effect

hydro on/off
has small effect

EPOS 3 compared to RHIC data

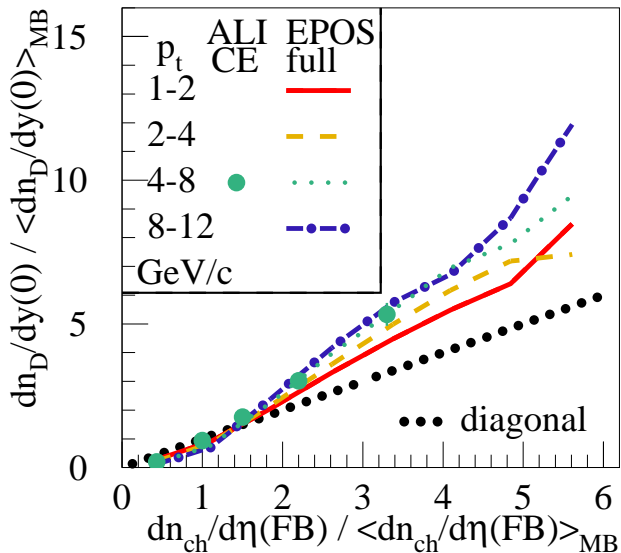


Calculations:
D mesons

Data: J/ Ψ

**Increase
stronger
than at LHC**

Multiplicity at FB rapidity (LHC)



**FB =
forward/backward
rapidity range:**

$$2.8 < \eta < 5.1$$

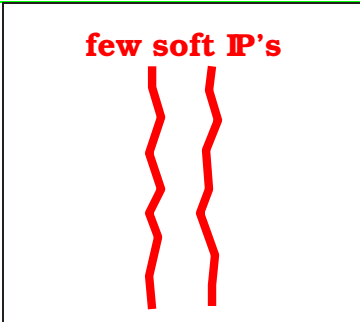
and

$$-3.7 < \eta < -1.7$$

Smaller increase

**Low
multi-
plicity
(LM)**

Small
 N_{Pom}



IP = Pomeron

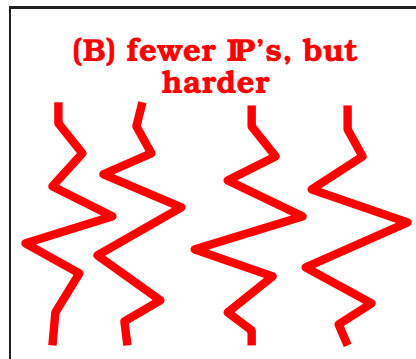
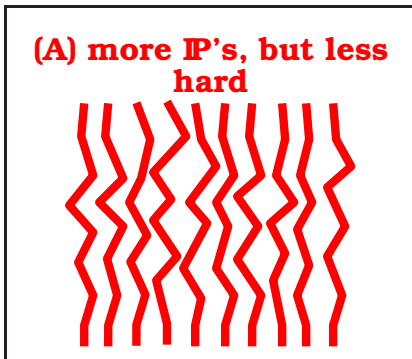
**“Hardness”
increases
with N_{Pom}**

(larger Q_s)

**High
multi-
plicity
(HM)**

**many
hard**

**IP's
on avg**



LM → HM:

Pomerons get harder (larger Q_s)

→ favors high pt or large masse production

**in particular due to case B (fewer IP's, but harder)
for highest pt bins !**

**Bigger effect at RHIC due to much narrower N_{Pom}
distribution (harder IP's are needed)**

Smaller effect for $\frac{dn}{d\eta}(FB)$ as multipl. variable

**(case B is replaced by case C: fewer IP's, but more covering
the FB rapidity range)**

Summary

- **Investigating the multiplicity dependence of particle ratios and mean p_t in pp, pA: EPOS's core-corona picture describes the trend**
- **Strong increase of corona p_t due to the N_{Pom} dependence of the saturation scale ...**
- **... which explains also the strong nonlinearity of the $D(J\backslash\Psi)$ multiplicity vs the charged multiplicity**
- **Multiplicity dep of $D(J/Psi)$ multiplicity:
“Ultimate tool” to test multiple scattering and its implementation of Q_S**

Core => Hydro evolution (Yuri Karpenko)

Israel-Stewart formulation, $\eta - \tau$ coordinates, $\eta/S = 0.08$, $\zeta/S = 0$

$$\partial_{;\nu} T^{\mu\nu} = \partial_\nu T^{\mu\nu} + \Gamma_{\nu\lambda}^\mu T^{\nu\lambda} + \Gamma_{\nu\lambda}^\nu T^{\mu\lambda} = 0$$

$$\gamma (\partial_t + v_i \partial_i) \pi^{\mu\nu} = -\frac{\pi^{\mu\nu} - \pi_{\text{NS}}^{\mu\nu}}{\tau_\pi} + I_\pi^{\mu\nu} \quad \gamma (\partial_t + v_i \partial_i) \Pi = -\frac{\Pi - \Pi_{\text{NS}}}{\tau_\Pi} + I_\Pi$$

$T^{\mu\nu} = \epsilon u^\mu u^\nu - (p + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$,

$\pi_{\text{NS}}^{\mu\nu} = \eta (\Delta^{\mu\lambda} \partial_{;\lambda} u^\nu + \Delta^{\nu\lambda} \partial_{;\lambda} u^\mu) - \frac{2}{3} \eta \Delta^{\mu\nu} \partial_{;\lambda} u^\lambda$

$\partial_{;\nu}$ denotes a covariant derivative,

$\Pi_{\text{NS}} = -\zeta \partial_{;\lambda} u^\lambda$

$\Delta^{\mu\nu} = g^{\mu\nu} - u^\mu u^\nu$ is the projector orthogonal to u^μ ,

$I_\pi^{\mu\nu} = -\frac{4}{3} \pi^{\mu\nu} \partial_{;\gamma} u^\gamma - [u^\nu \pi^{\mu\beta} + u^\mu \pi^{\nu\beta}] u^\lambda \partial_{;\lambda} u_\beta$

$\pi^{\mu\nu}$, Π shear stress tensor, bulk pressure

$I_\Pi = -\frac{4}{3} \Pi \partial_{;\gamma} u^\gamma$

Freeze out: at 164 MeV, Cooper-Frye $E \frac{dn}{d^3p} = \int d\Sigma_\mu p^\mu f(up)$, equilibrium distr

Hadronic afterburner: UrQMD

Marcus Bleicher, Jan Steinheimer