Hydrodynamization of shockwave collisons in non-conformal theories

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Holography is a unique and powerful tool capable to explore an unexplored regime of quantum dynamics.

The AdS/CFT correspondence builds a bridge between problems from the Quantum world and methods from General Relativity.

In the fully non-linear regime the use of numerical methods is instrumental.

$\mathcal{N} = 4$ SYM properties:

- $\eta/s \geq 1/4\pi$
- supersymmetric
- conformal
- $\mathcal{N}_c \to \infty$

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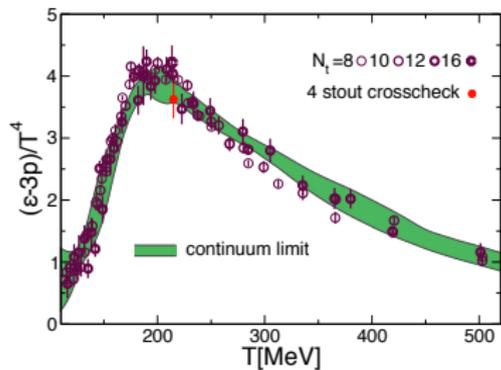- $\mathcal{N}_c \to \infty$

Far from equilibrium dynamics:

- at strong coupling
- non-perturtabitave
- fast hydrodynamization time
- IC hydrodynamics
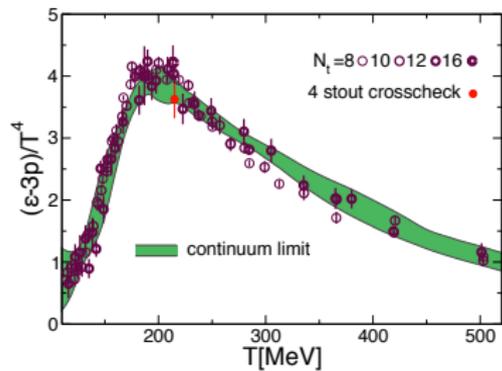
[S. Borsanyi et alii arXiv:1309.5258 [hep-lat]]

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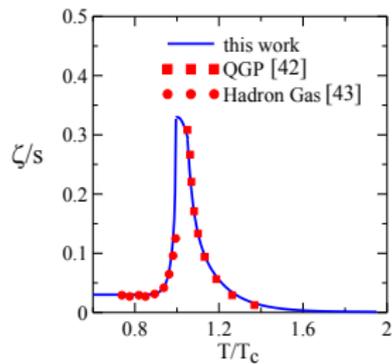
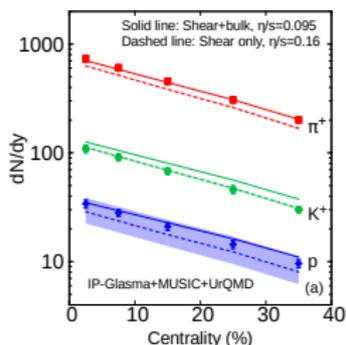

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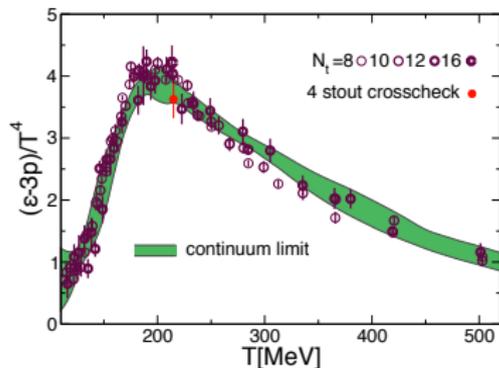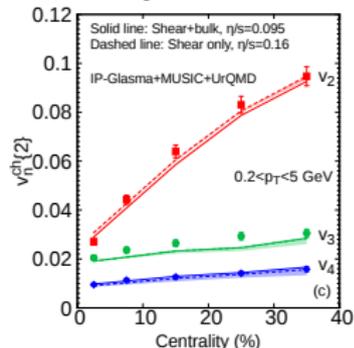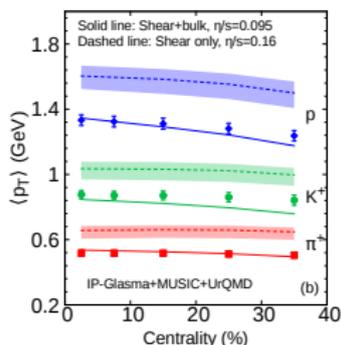
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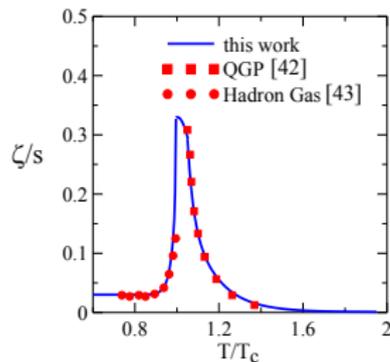
[Denicol *et alii* arXiv:0903.3595 [hep-ph]]

Hydro simulation agreement improves with bulk viscosity:

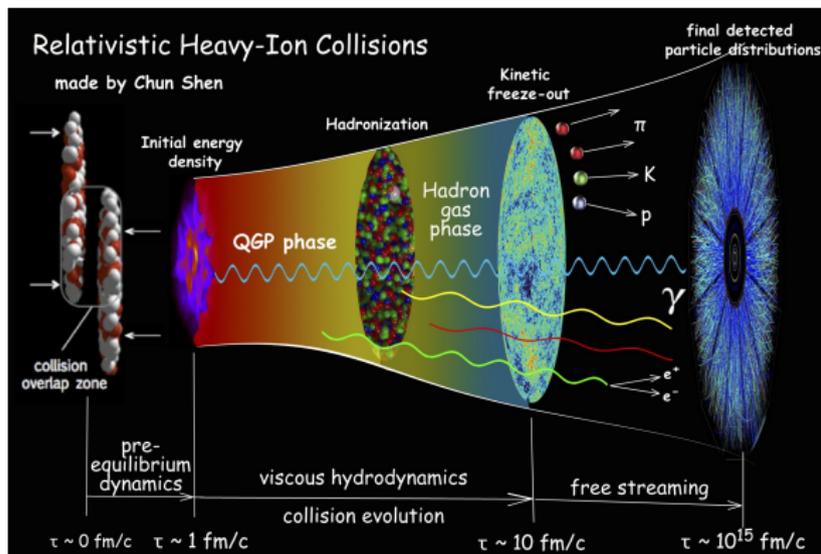[Denicol *et alii* arXiv:1502.01675 [nucl-th]]

Outline

Non-conformal bottom-up thermodynamics

- Heavy-Ion collision
- General Relativity setup
- Scalar potential
- Equilibrium pressure
- Interaction measure
- Speed of sound
- Bulk viscosity
- Buchel bound
- Quasi-Normal-Modes

Dynamics of the scalar potential

Relativistic Heavy-Ion Collisions
made by Chun Shen

Stages:
1) Early out of equilibrium

2) Quark-Gluon Plasma

3) Hot Hadron Gas

Can we describe the first two stages at strong coupling? **Yes!**

AdS/CFT provides setups for strongly coupled phenomena that can be compared to heavy-ion collisions:
Shockwave collisions in N=4 SYM
[Chesler, Yaffe 11; Albacete, Kovchegov, Taliotis 08; Grumiller, Romatschke 08]

Einstein-Hilbert action coupled to a scalar with non-trivial potential in five-dimensional bottom-up model:

$$S = \frac{2}{\kappa_5^2} \int d^5x \sqrt{-g} \left[\frac{1}{4}\mathcal{R} - \frac{1}{2}\left(\nabla\phi\right)^2 - V(\phi) \right].$$

On the gravity side the flow is dual to a domain-wall geometry interpolating between two AdS with L radius of the UV AdS solution:

$$ds^2 = \frac{L_{\text{eff}}(z)^2}{z^2}\left(-dt^2 + d\mathbf{x}^2 + dz^2\right).$$
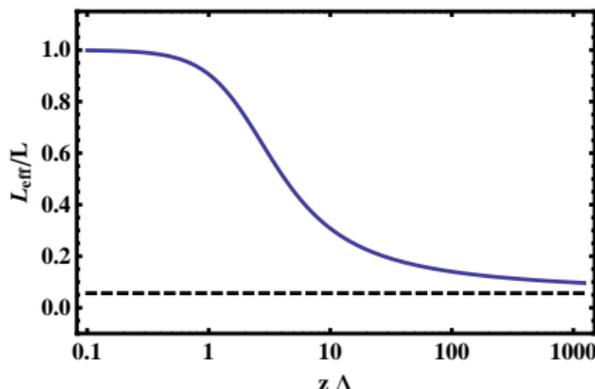
$V(\phi)$ depends on single parameter $\phi_\mathbf{M}$, setting non-conformality for this bottom-up model:
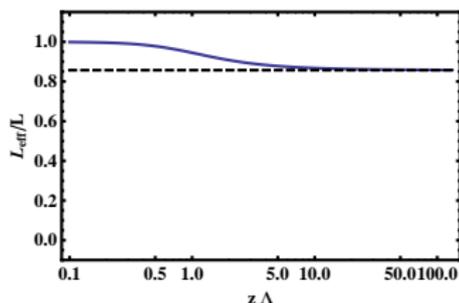
$$L^2 V(\phi) = -\frac{1}{12\phi_M^4}\phi^8 + \left(\frac{1}{2\phi_M^4} + \frac{1}{3\phi_M^2}\right)\phi^6 - \frac{1}{3}\phi^3 - \frac{3}{2}\phi^2 - 3\,.$$

Deforming $\mathcal{N} = 4$ Super Yang-Mills with an dimension 3 operator $\mathcal{O}$ dual to the scalar field ϕ.

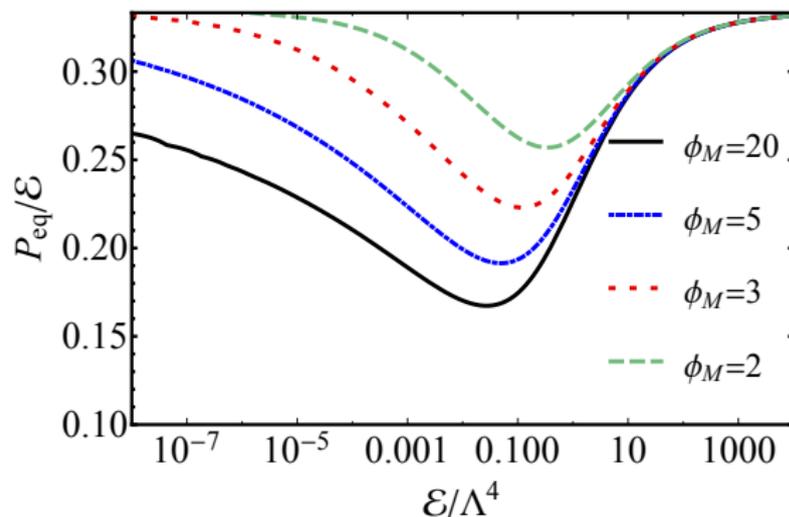$$\left\langle T_\mu^\mu \right\rangle = -\Lambda\mathcal{O}\,.$$

The source Λ triggers RG flow responsible for the breaking of conformal invariance.

Change of variable in the holographic direction shows different approach to IR fixed point for $phi_M = \{1, 3, 10, 100\}$
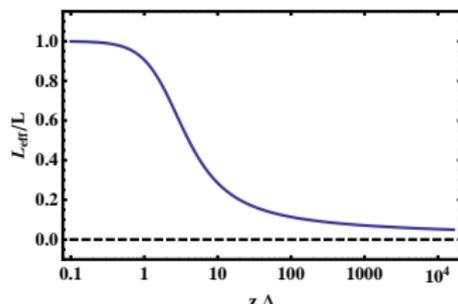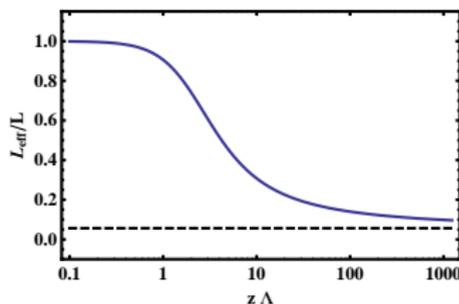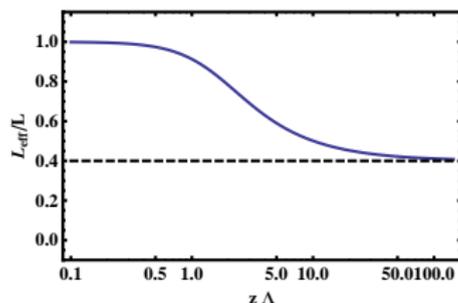
$$ds^2 = \frac{L_{\text{eff}}(z)^2}{z^2} \left(-dt^2 + d\mathbf{x}^2 + dz^2 \right) .$$
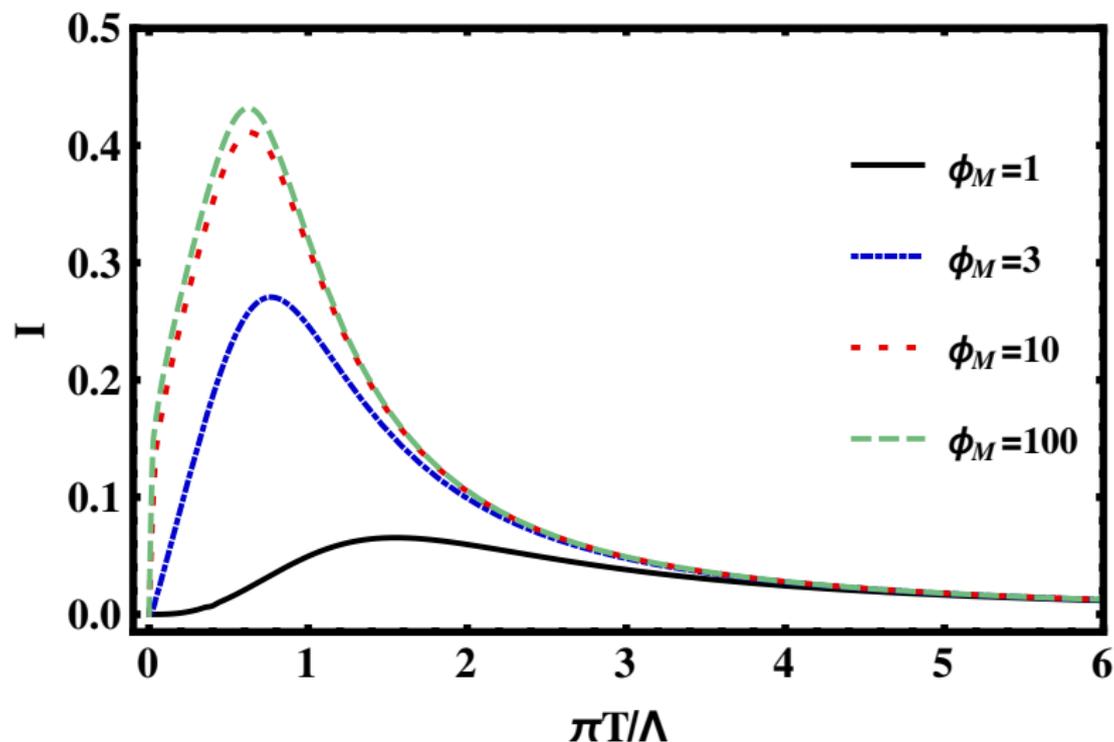
Equilibrium pressure as a function of energy density:

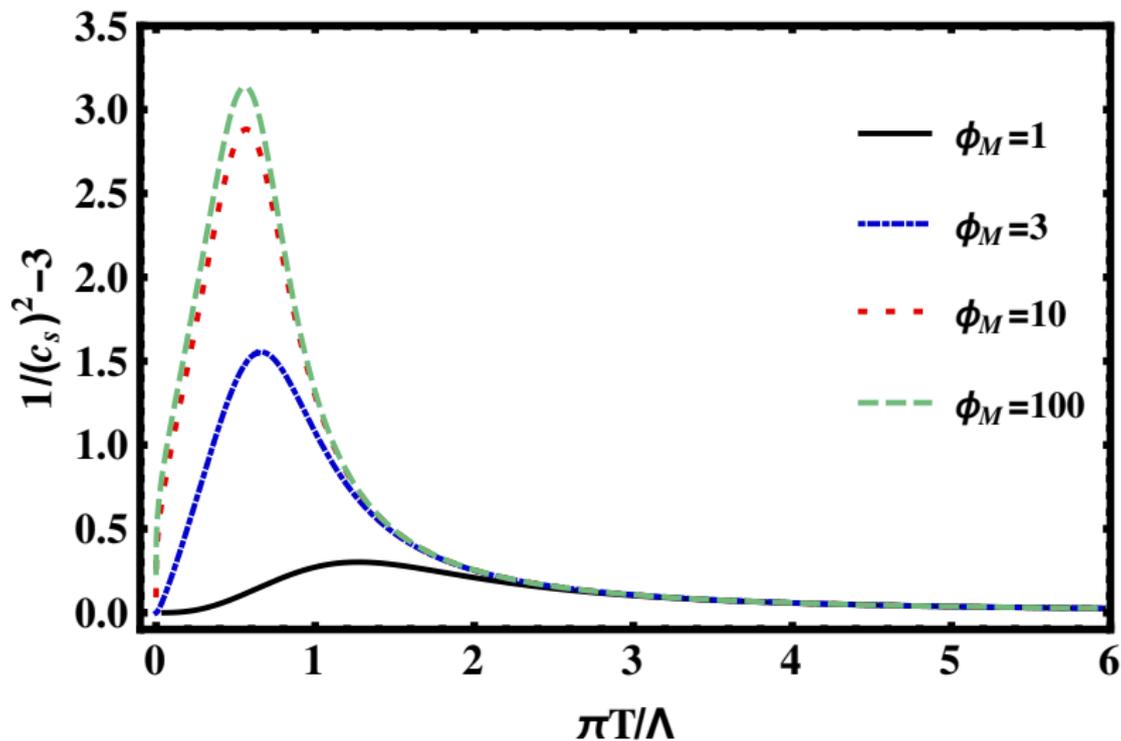$$P_{eq}(\mathcal{E}) = \tfrac{1}{3}\left[\mathcal{E} - \Lambda \mathcal{V}_{eq}(\mathcal{E})\right].$$

Out of equilibrium the average pressure is not determined by the energy density as the scalar expectation value $\mathcal{V}$ fluctuates independently.

Non-conformality measure as function of the temperature:

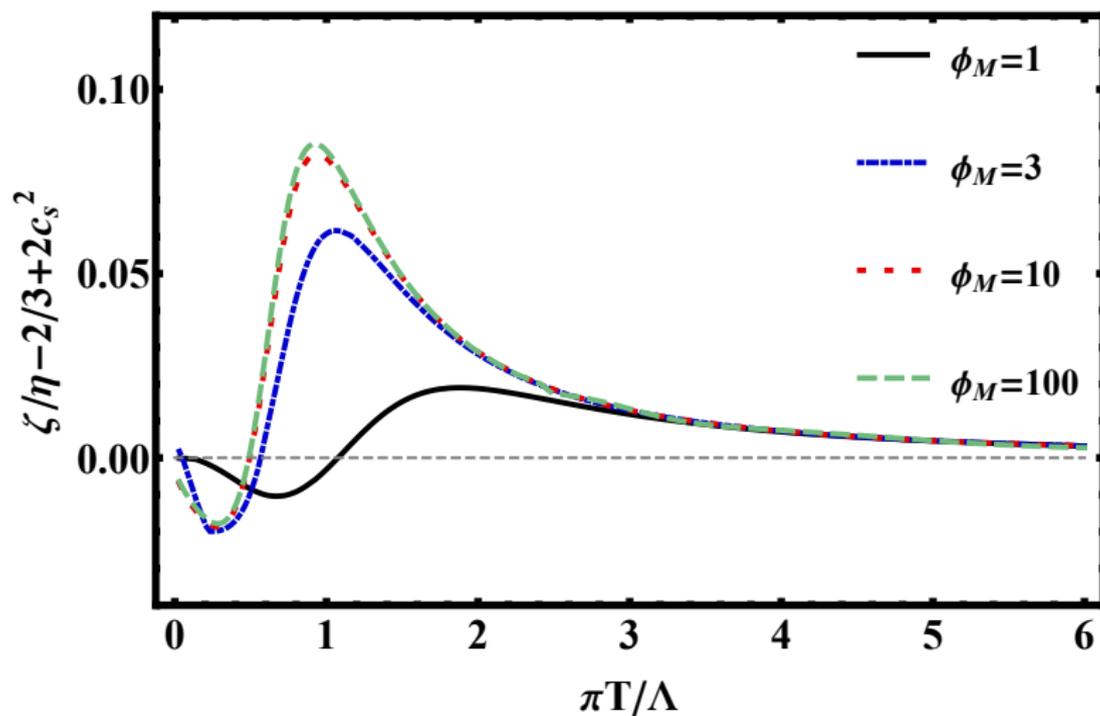conformal at low and high T, **non-conformal** in between

Deviation of c_s from its conformal value:

Temperature behaviour of the inverse speed of sound square

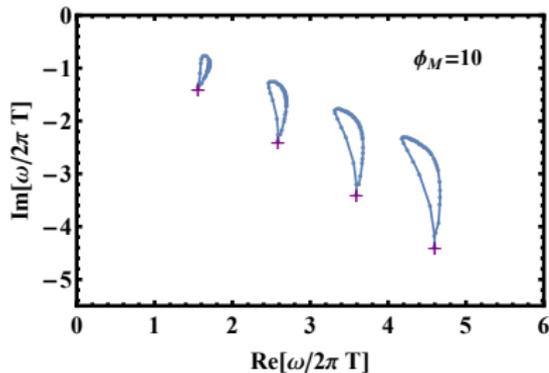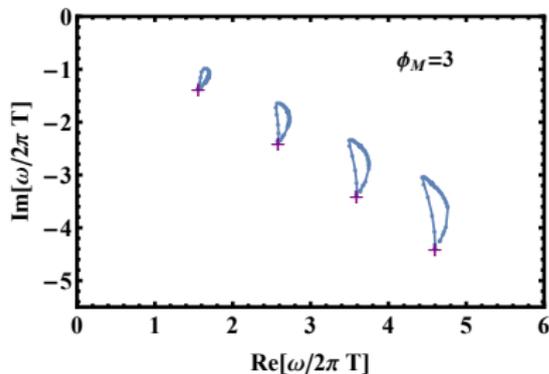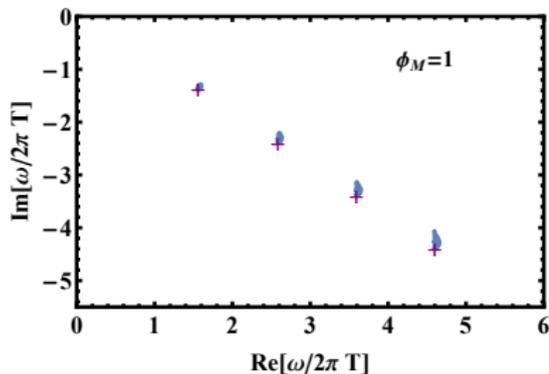Ratio of bulk viscosity over entropy as a function of temperature:
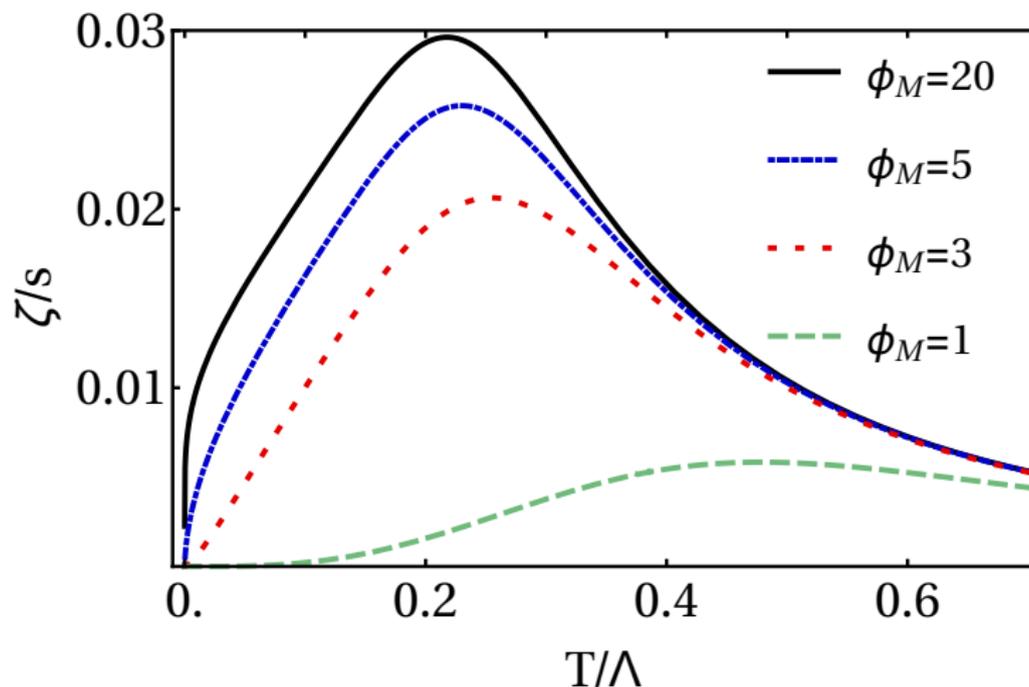
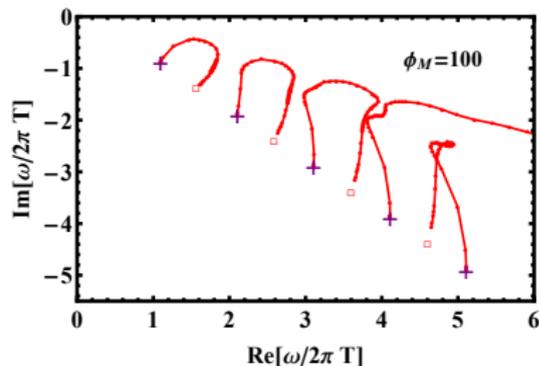non-conformal behaviour reflects in transport coefficients

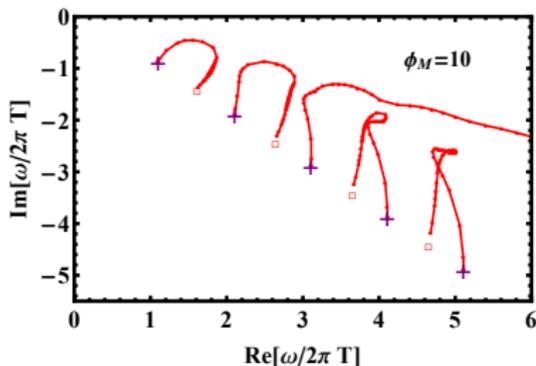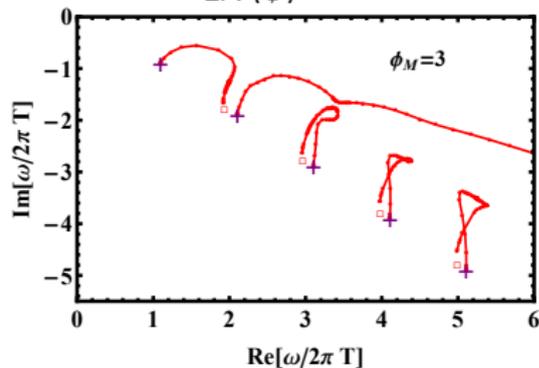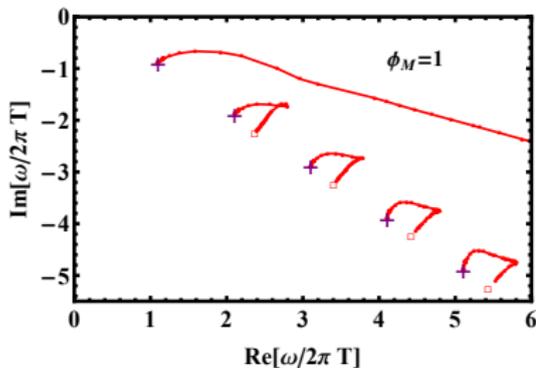Violation of Buchel's bound at low temperatures:

Maxima of speed of sound and bulk to shear viscosity different!

anisotropic perturbation $Z_{aniso} = e^{-2A}(h_{zz} - h_{aa})$
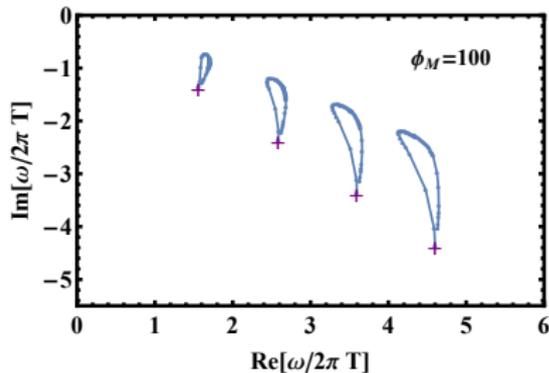
Fluctuations of the stress energy tensor with same IR + UV limit

non-conformal mode $Z_{bulk} = \phi - \frac{e^{-2A(\phi)}}{2A'(\phi)} h_{aa}$

n-th scalar mode decoupling with anti-crossing and different IR/UV limits

Non-conformal bottom-up thermodynamics

Dynamics of the scalar potential
- Shockwaves Initial Conditions
- Hydrodynamization / EoSization time
- Non-conformal temperature scan
- Non-conformal theories
- Rapidity profile
- Scalar relaxation times

5D metric Ansatz in Eddington-Finkelstein:

$$ds^2 = -A dt^2 + \Sigma^2 \left(e^B d\boldsymbol{x}_\perp^2 + e^{-2B} dz^2 \right) + 2dt(dr + Fdz)$$

- Field theory interpretation:
    - Defined by energy density
    - Move in AdS5 space
    - Demand that shockwaves move at speed of light
    - Quantum state/AdS geometry completely fixed for pure gravity
- Homogeneous in transverse plane ('infinite nucleus')
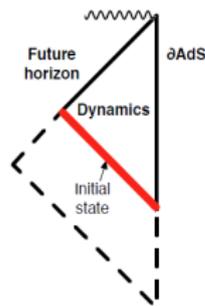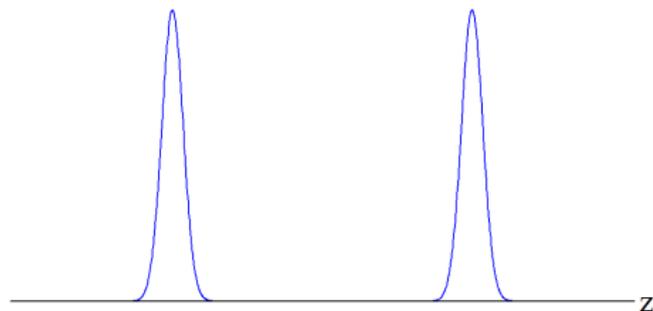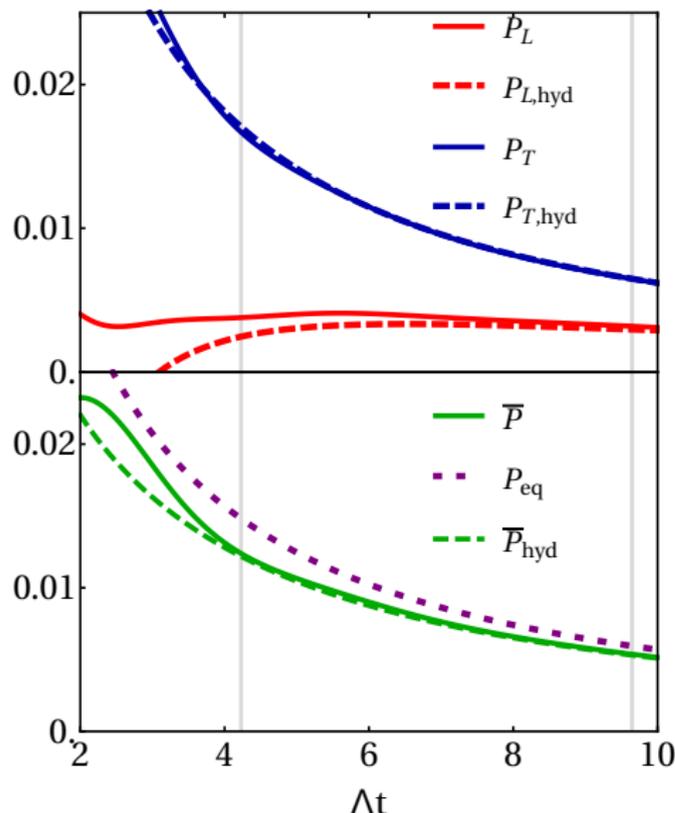
Hydrodynamics assumes mean free path goes to zero:

$$\partial_\mu T^{\mu\nu} = 0$$
$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu + p g^{\mu\nu}$$
$$+ \eta \Pi^{\mu\nu} + \zeta \Pi(g^{\mu\nu} + u^\mu u^\nu)$$
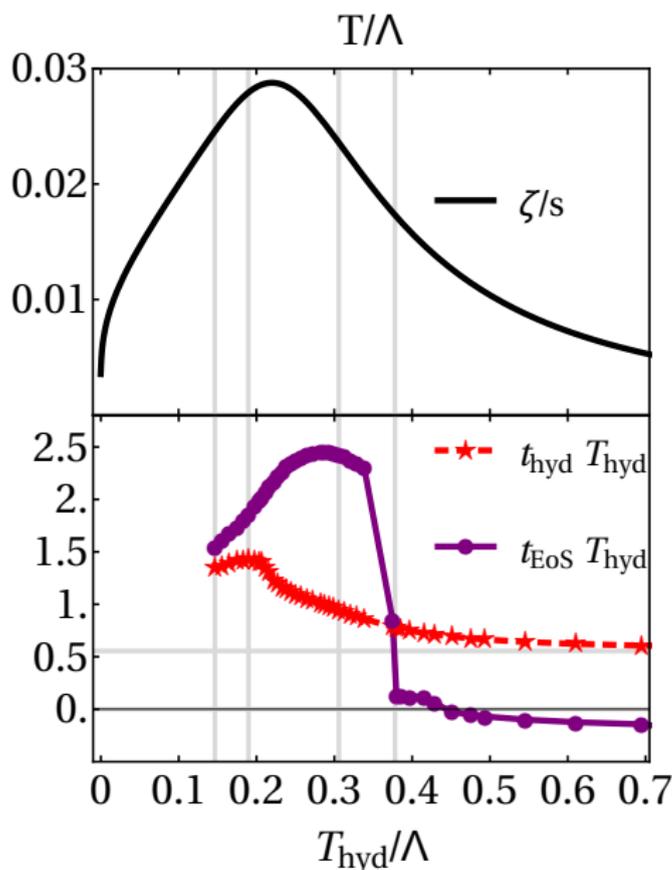
Hydrodynamization:

$$\left| P_{L,T} - P_{L,T}^{\mathrm{hyd}} \right| / \bar{P} < 0.1$$

EoSization:

$$\left| \bar{P} - P_{\mathrm{eq}} \right| / \bar{P} < 0.1$$

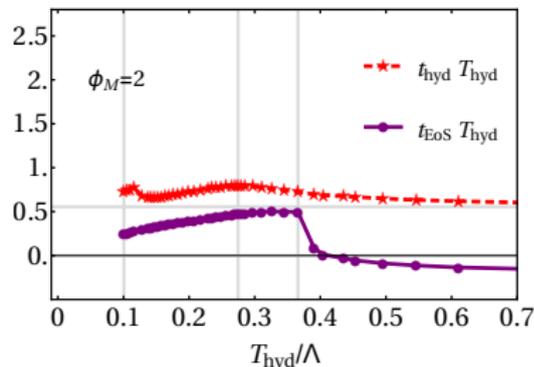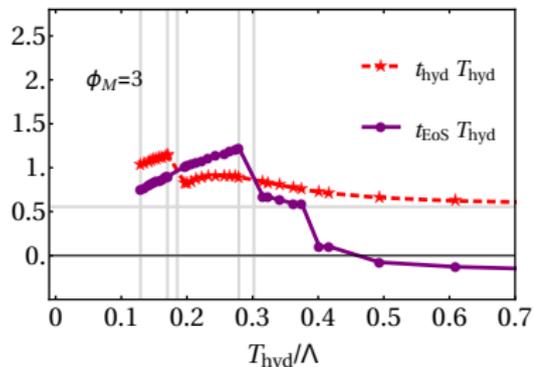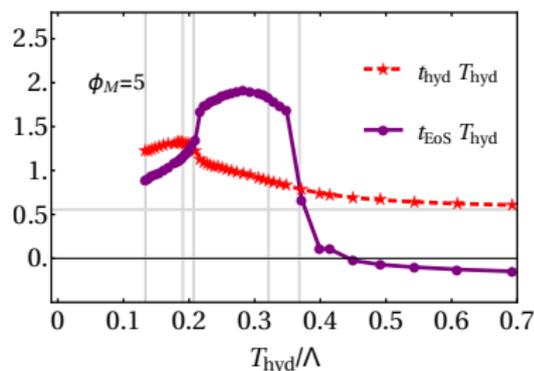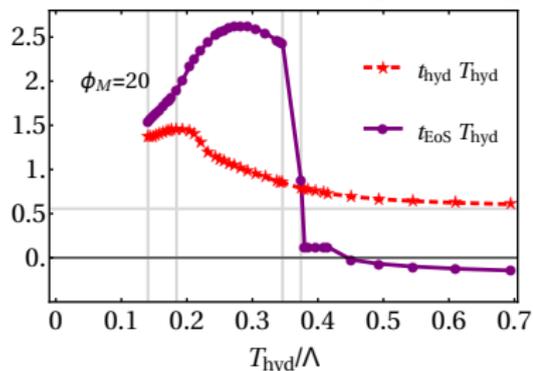hydrodynamization $\neq$ EoSization $\neq$ isotropization
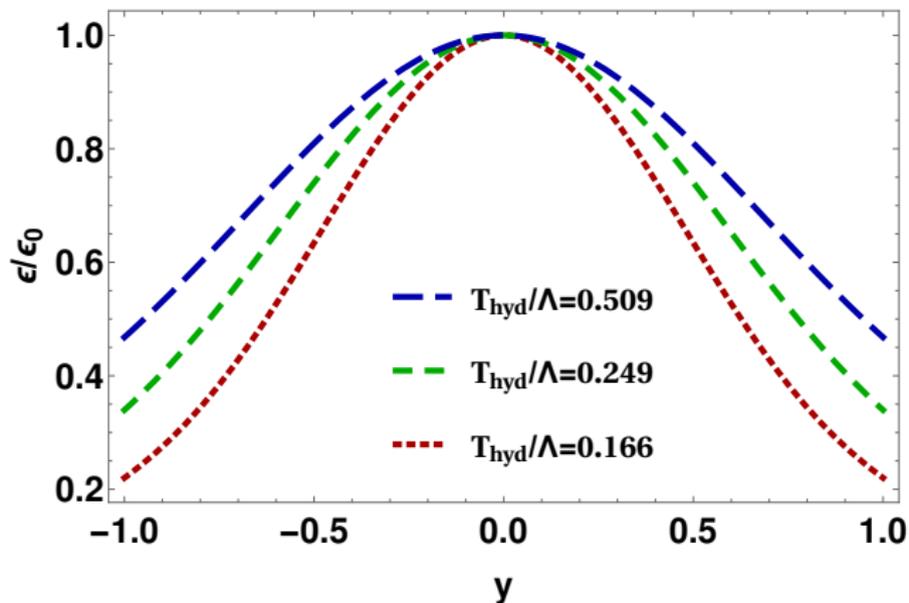
Non-conformal T scan:

- EOS does NOT hold out of equilibrium

- t_{hyd} slow down

- required ζ 1/10 of QCD at T_c

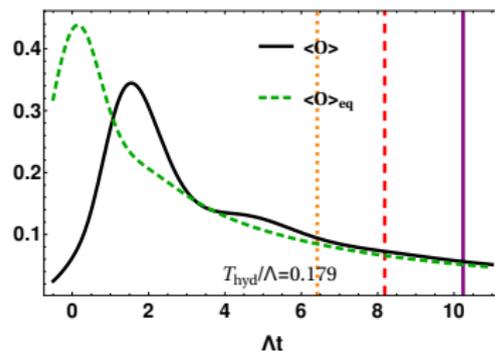- ordering of t_{EoS} and t_{hyd} depends on bulk viscosity

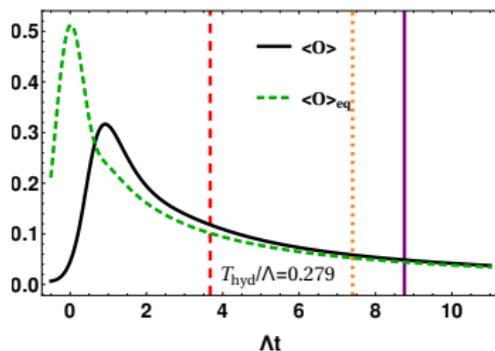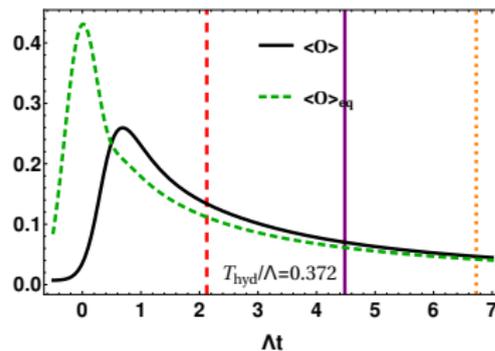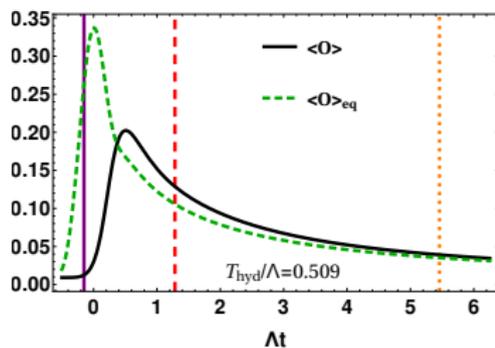Comparing varying non-conformality ϕ_M:

conservative estimate $\zeta/s > 0.025$ needed for $t_{\mathrm{EoS}} > t_{\mathrm{hyd}}$

At Hydrodynamization time allmost gaussian distribution:

Higher energy densities results in broader rapidity profile

$|O - O_{eq}|/O$ relaxation time t_O of the VeV:

4 realized scenarios of t_{hyd}, t_{EoS}, t_O ordering

Conclusions

- First simulation of a holographic non-conformal model for heavy ion collisions

- New relaxation channel from bulk viscosity:
  hydrodynamization without EoSization

- Hydrodynamics works at early time (only max delayed ≈ 3)
  - despite non-trivial equation of state
  - despite sizeable ζ/s bulk viscosity over entropy

- Conservative estimate $\zeta/s \approx 0.025$ for non-conformal effects

- More studies are on the way:
  - Asymmetrical collisions, exploding balls
  - phase transitions
  - Holographic energy scan
 - Different potentials: $\mathcal{N} = 2^*$, Gubser, ..
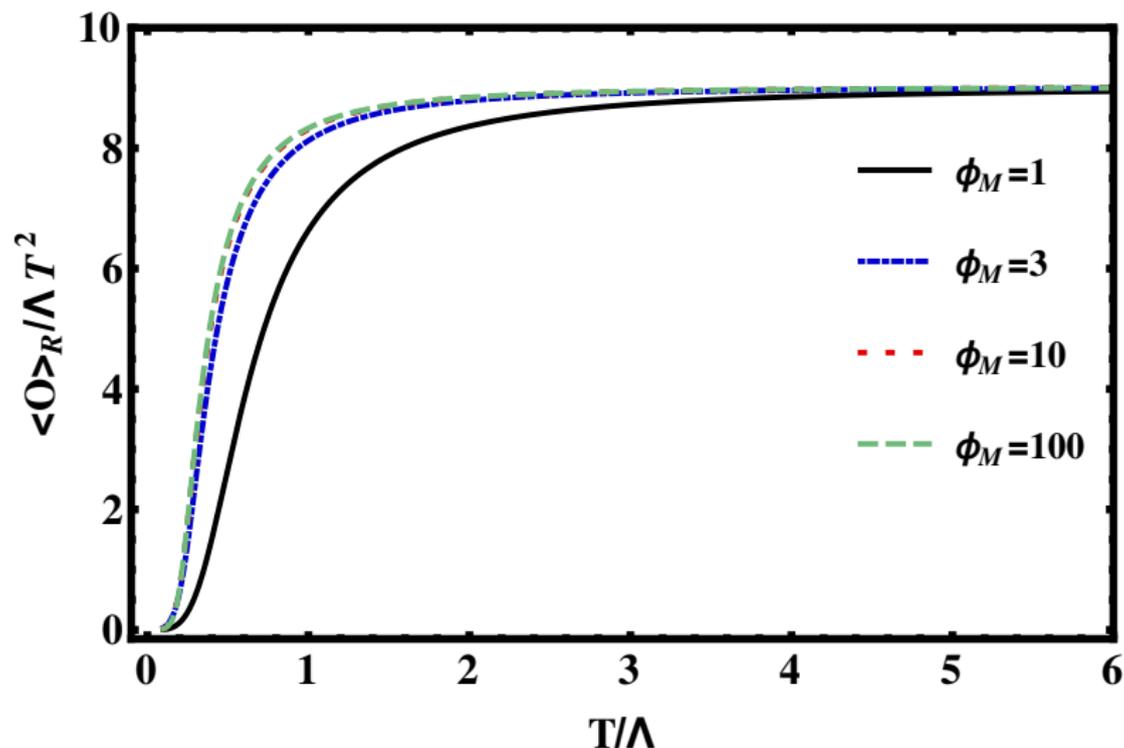
Temperature dependence of the VEV of the scalar operator $\langle \mathcal{O} \rangle_T$ for several values of ϕ_M. $\langle \mathcal{O} \rangle_R = \kappa_5^2 \langle \mathcal{O} \rangle_T / L^3$, $\epsilon - 3p = \Lambda \langle \mathcal{O} \rangle_T$.

ϕ_2 and b_4 as functions of time for a z-independent configuration
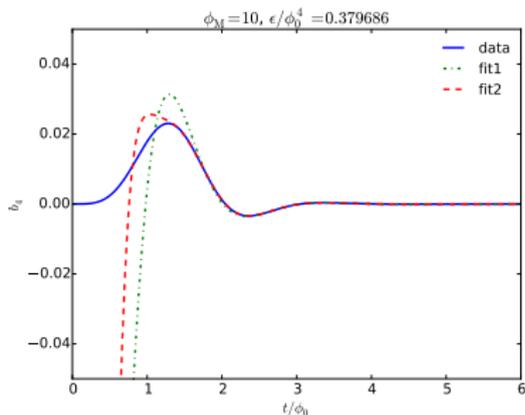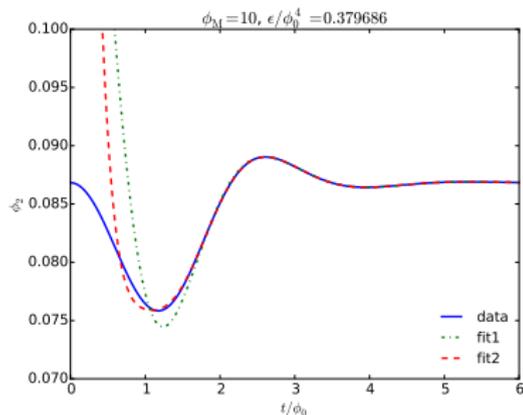
Blue full line corresponds to data from the code, green dash-dotted line correspond to a fit to the data using one QNM, red dashed line corresponds to a fit using two QNMs as explained in the text.

Differences between the coarse and medium (blue solid line) and the medium and fine (red dashed line) resolution run
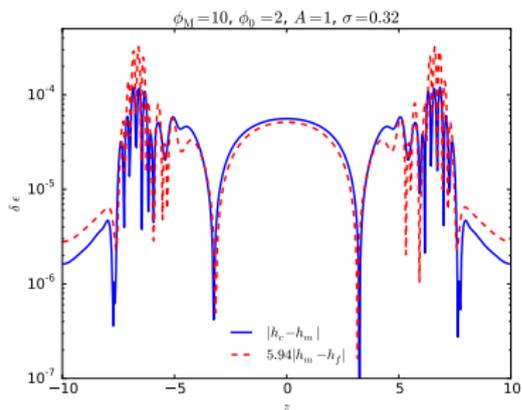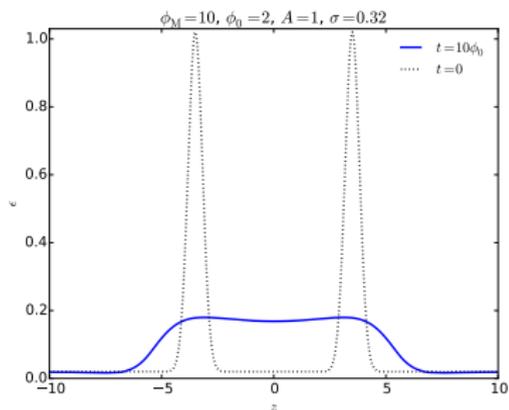
The results show fourth-order convergence.