Helicity Selection Rules and Non-Interference for BSM Amplitudes
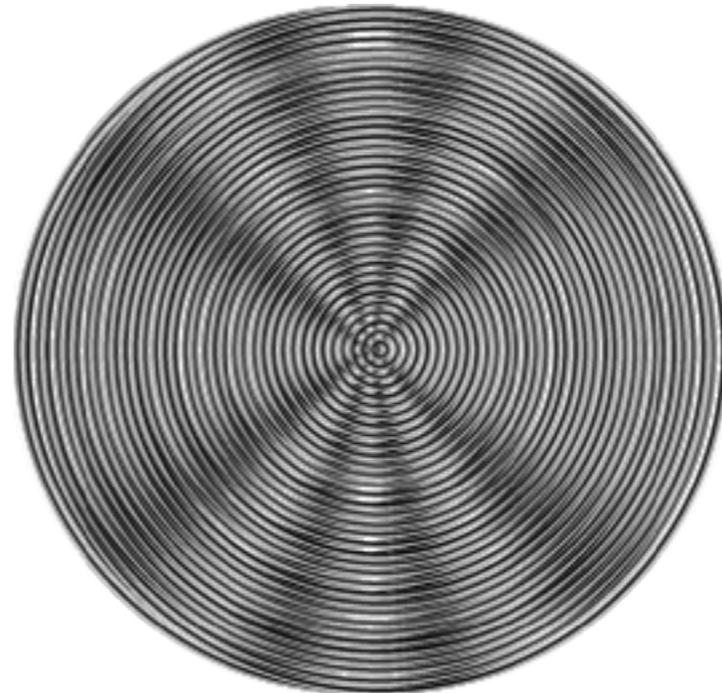
Francesco Riva
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!

The following arguments hold for any parametrization for SM precision tests that involves a Derivative Expansion:

EFT, Non-linear EFT, Higgs-characterization, Pseudo Observables,...

(a consequence of Lorentz invariance and the equivalence theorem)

Interference

When SM and BSM contribute to the **same** amplitude:

$$Amp = SM + BSM = SM(1 + \delta_{BSM})$$

▶ $\quad \sigma \propto |Amp|^2 \simeq SM^2(1 + \delta_{BSM} + \delta_{BSM}^2)$

For **small** BSM effects $1 \gg \delta_{BSM}$, **interference dominates** $\delta_{BSM} \gg \delta_{BSM}^2$

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e.g. EW fits to LEP-1 observables:

Experiment very precise ➡ $1 \gg \delta_{BSM}$

▶ **(EFT) expansion in** $\delta_{BSM} = \dfrac{v^2}{\Lambda^2}$ **well defined, often approximated at linear order (interference only)**

Non-Interference

Derivative (energy) expansion:

dimension-6

$$\delta_{BSM} = c_i \frac{E^2}{\Lambda^2} + \cdots$$

If SM and BSM contribute to different amplitudes:

interference
vanishes

$$\sigma \propto \sum |Amp|^2 \simeq SM^2(1 + c_i \frac{E^2}{\Lambda^2} + c_i^2 \frac{E^4}{\Lambda^4})$$

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The leading effects BSM are $O\left(\frac{1}{\Lambda^4}\right)$:

the same order as dimension-8 that do interfere

▶ **Question of interference relevant for expansion consistency**

LHC

High-energy $(2 \to 2, 2 \to 3, ...)$ processes important part of LHC search program:

e.g. Drell-Yann, TGC, QGC, VH associated production, VBF, tth, h+j,...

This Talk: **Interference in High-energy processes and its implications**

Non-Interference for BSM$_6$ amplitudes

Exploit:

For $E \gg m_W$ states have well defined helicity

Amplitudes for 2→2 with different total h don't interfere

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Theorem:

Any BSM dim-6 operator

A_4	$|h(A_4^{\mathrm{SM}})|$	$|h(A_4^{\mathrm{BSM}})|$
$VVVV$	0	4,2
$VV\phi\phi$	0	2
$VV\psi\psi$	0	2
$V\psi\psi\phi$	0	2
$\psi\psi\psi\psi$	2,0	2,0
$\psi\psi\phi\phi$	0	0
$\phi\phi\phi\phi$	0	0

Massless limit + tree level + at least one transverse vector
- SM and BSM$_6$ contribute to different helicity amplitudes
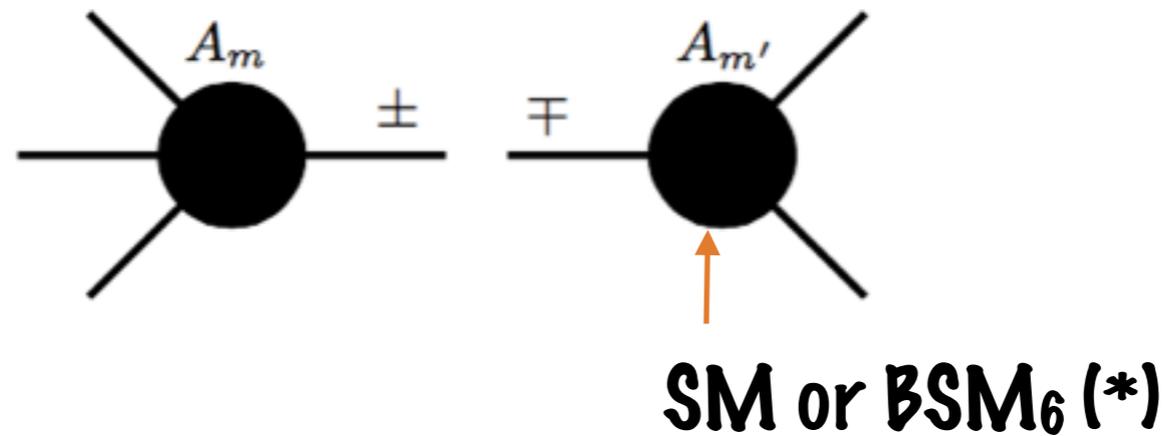- No interference

Non-Interference for BSM$_6$ amplitudes

How? **i) Helicity sums:** $h(A_n) = h(A_m) + h(A_{m'})$

n=m+m'-2 legs

SM or BSM$_6$ (*)

Non-Interference for BSM₆ amplitudes

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A_m $\pm$ $\mp$ $A_{m'}$

SM or BSM₆ (*)

$p \in \mathbb{C}$ (outgoing)so that on-shell condition $p^2 = 0$ satisfied also for A₃

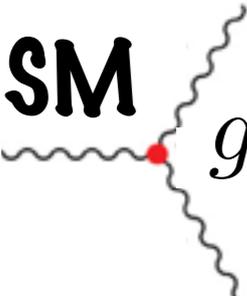(*)=In a basis where all effects proportional to the Equations of Motion have been eliminated

Non-Interference for BSM$_6$ amplitudes

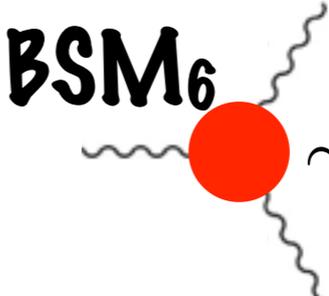How? **ii) Helicity of 3-point$\leftrightarrow$coupling dimension:**

$$|h(A_3)| = 1 - [g]$$

SM $g \Rightarrow |h^{SM}| = 1$

BSM$_6$ $\sim \dfrac{1}{\Lambda^2} \Rightarrow |h^{BSM}| = 3$

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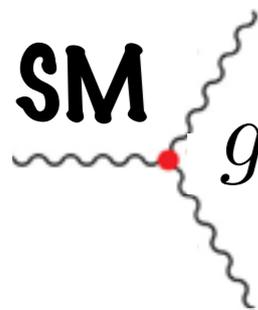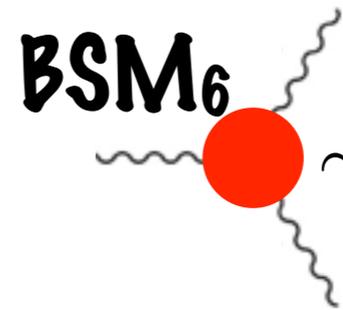**BSM$_6$** $\sim \dfrac{1}{\Lambda^2} \Rightarrow |h^{BSM}| = 3$

Spinor Helicity Formalism:
$$v_+(p) = (|p]_a, 0) \quad \epsilon_+^\mu(p;q) = \frac{\langle q|\gamma^\mu|p]}{\sqrt{2}\langle qp\rangle}$$
$$v_-(p) = (0, |p\rangle^{\dot{a}})$$
$$-\not{p} = |p\rangle[p| + |p]\langle p|$$

irreps of $SU(2) \times SU(2)$
$$\simeq SO(3,1)$$
(familiar in SUSY)

Little Group scaling (Poincarré) $|p_i\rangle \to t_i|p_i\rangle \quad |p_i] \to t_i^{-1}|p_i] \quad \to \quad p_i^\mu$ invariant $\to \quad A \sim t_i^{-2h_i}$

Dimensional Analysis $[|p\rangle] \sim [|p]] \sim \mathrm{GeV}^{1/2} \qquad [A_{3point}] \sim \mathrm{GeV}$

3-point kinematics $\langle p_1 p_2\rangle[p_1 p_2] = 2p_1 \cdot p_2 = (p_1 + p_2)^2 = (-p_3)^2 = 0 \to \begin{cases} A_3 \sim \langle p_1 p_2\rangle \cdots \\ A_3 \sim [p_1 p_2] \cdots \end{cases}$ or

▶ $A_3(1^{h_1} 2^{h_2} 3^{h_3}) = g\langle 12\rangle^{h_3 - h_1 - h_2}\langle 13\rangle^{h_2 - h_1 - h_3}\langle 23\rangle^{h_1 - h_2 - h_3}$ **uniquely fixed**

Non-Interference for BSM$_6$ amplitudes

Azatov,Contino,Machado,FR'16

How? **iii) SUSY* Ward Identities:** $|h(A_4^{SM})| < 2$ (except$\psi^+\psi^+\psi^+\psi^+$)

SM : $A(V^+V^+V^+V^-) = A(V^+V^+\psi^+\psi^-) = A(V^+V^+\phi\phi) = A(V^+\psi^+\psi^+\phi) = 0$

BSM : **Operators with transverse V not supersymmetrizable**

Elias-Miro,Espinosa,Pomarol'14

▶ **no restriction on total helicity for BSM**

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$BSM:$ **Operators with transverse V not supersymmetrizable**

▶ **no restriction on total helicity for BSM**

*in the limit of either y$_u$=0 or y$_d$=y$_l$=0:

SM upliftable to SUSY+R-parity (with 1 Higgs doublett)

→ **e.g.** $0 = \langle 0|[Q, \Psi^+V^+V^+V^+]|0\rangle = \sum_i \langle 0|\Psi^+...[Q,V^+]...V^+|0\rangle \propto \langle 0|V^+V^+V^+V^+|0\rangle$

$[Q,\psi] \sim V \quad [Q,V] \sim \psi \qquad \psi^+\psi^+\cdots = 0$

Non-Interference for BSM$_6$ amplitudes
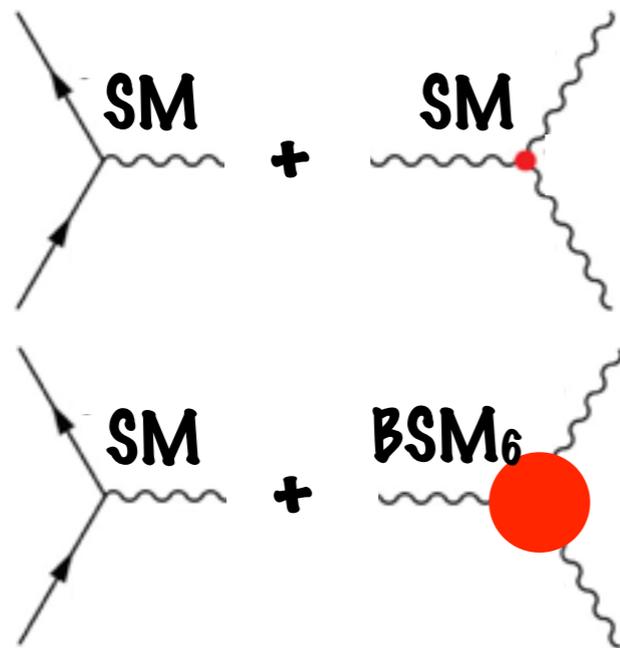
Azatov,Contino,Machado,FR'16

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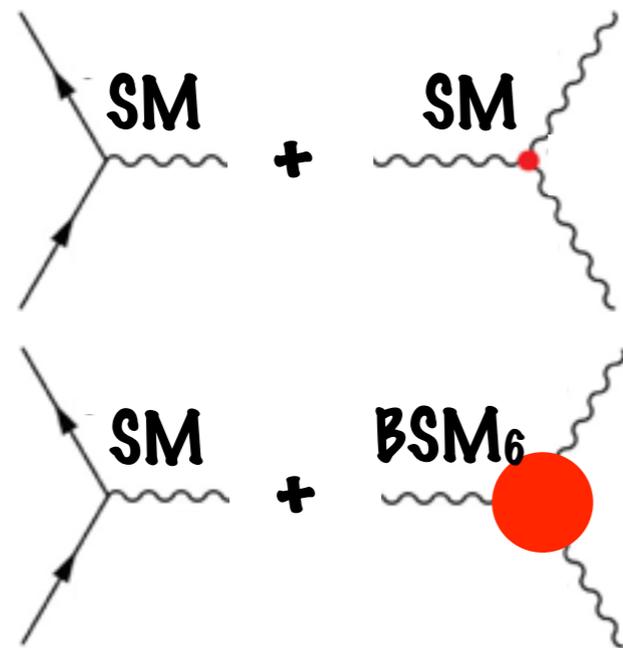# Non-Interference for BSM$_6$ amplitudes

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SM $+$ SM $\qquad h^{tot} = 1 - 1 = 0$

SM $+$ BSM$_6$ $\qquad h^{tot} = 1 \pm 3 = 2, 4$

No Interference
(dim-6, 4-point)

Implications

At $E \gg m_W$ dimension-6 with transverse vectors don't interfere in 2 → 2

▶ Naive derivative expansion truncated at dim-6 inconsistent

Processes Concerned:

- di-boson **WW,WZ,W**γ pair production*
- **VH**, tth associated production
- gg->tt
- **VBF**
- **VV** scattering
- ...

see Eleni's Talk

What to do?

(*= Notice that this concerns also LEPII TGC analyses)

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What to do?
1) Ignore analyses that already test these dim-6
2) Fit also dimension-8 that interfere
3) Study 2→3 processes with extra, e.g., jets that interfere
4) Find BSM context in which these searches actually make sense
5) Forget consistency and take dimension-6 as a complete benchmark

see Dixon,Shadmi'93;
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BSM context in which dim-6 searches make sense

Implications?

$$\bar{q}q \to VV'$$

$$V^{(')} = W^{\pm}, Z, \gamma$$

dim-6

$$\frac{g_*}{M^2} \epsilon_{abc} W_\mu^{a\,\nu} W_{\nu\rho}^b W^{c\,\rho\mu}$$

doesn't interfere for $m_W \to 0$

dim-8

$$\frac{g_*}{M^4} \epsilon_{abc} \Box W_\mu^{a\,\nu} W_{\nu\rho}^b W^{c\,\rho\mu}$$

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(3-point vertex has dimension of coupling → weight with g* for illustration)

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$$\sigma_T \sim \frac{g_{SM}^4}{E^2}\left[1 + \overbrace{\frac{g_*}{g_{SM}}\frac{m_W^2}{\Lambda^2}}^{BSM_6 \times SM} + \overbrace{\frac{g_*^2}{g_{SM}^2}\frac{E^4}{\Lambda^4}}^{BSM_6{}^2} + \underbrace{\frac{g_*}{g_{SM}}\frac{E^4}{\Lambda^4}}_{BSM_8 \times SM} + \dots\right]$$
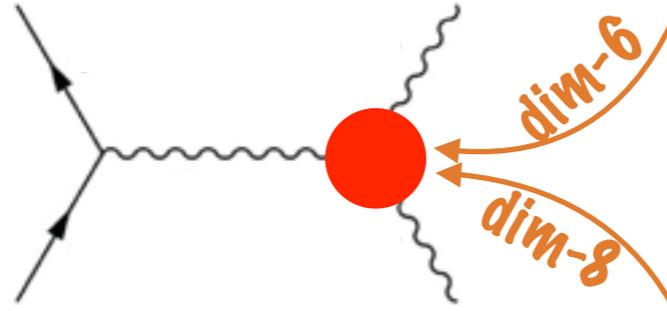

Non-interference:
$O(BSM_6{}^2) \approx O(BSM_8{}^2)$

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Liu,Pomarol,Rattazzi,FR'16

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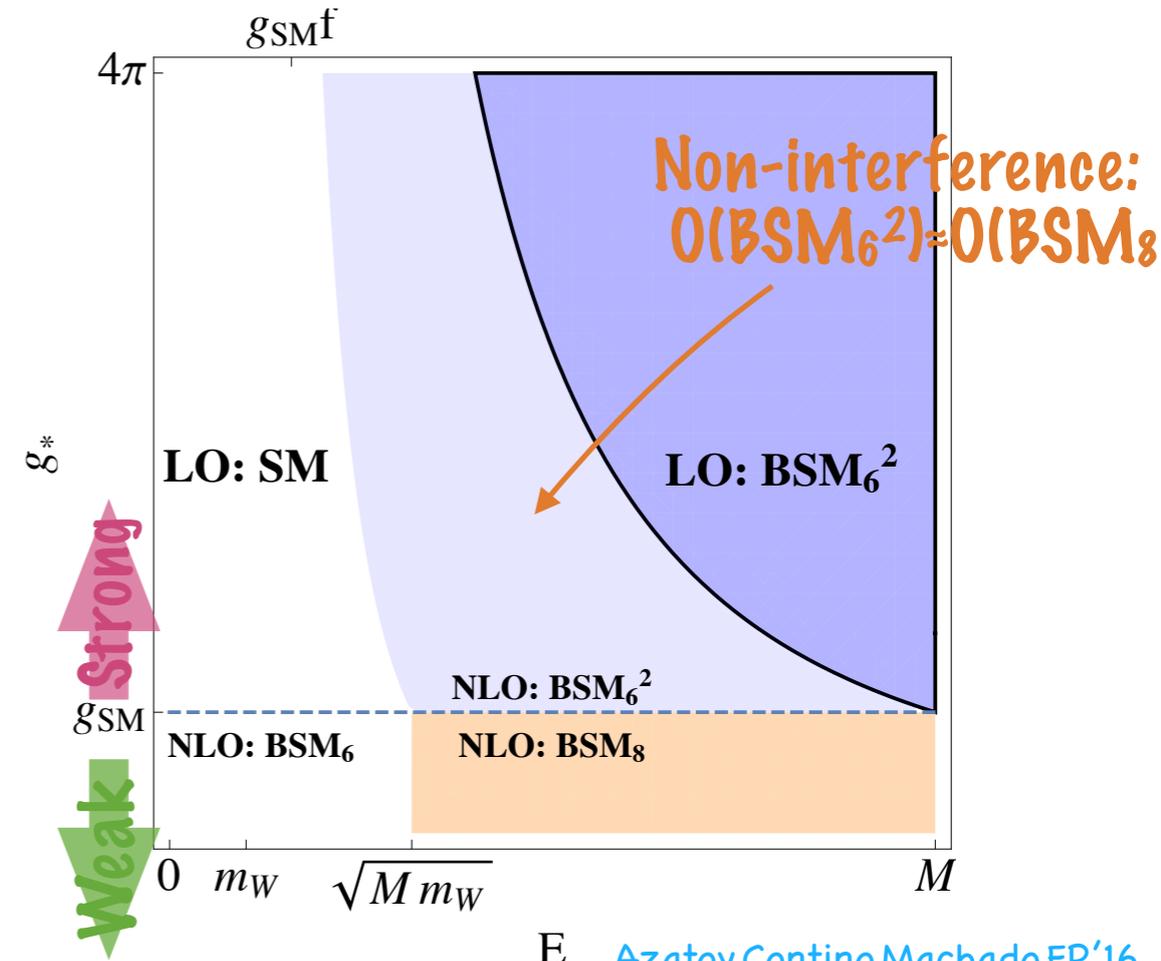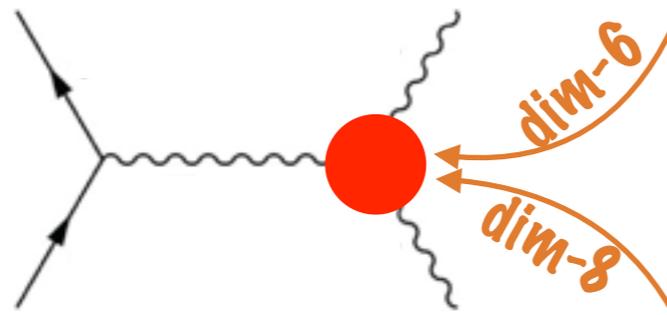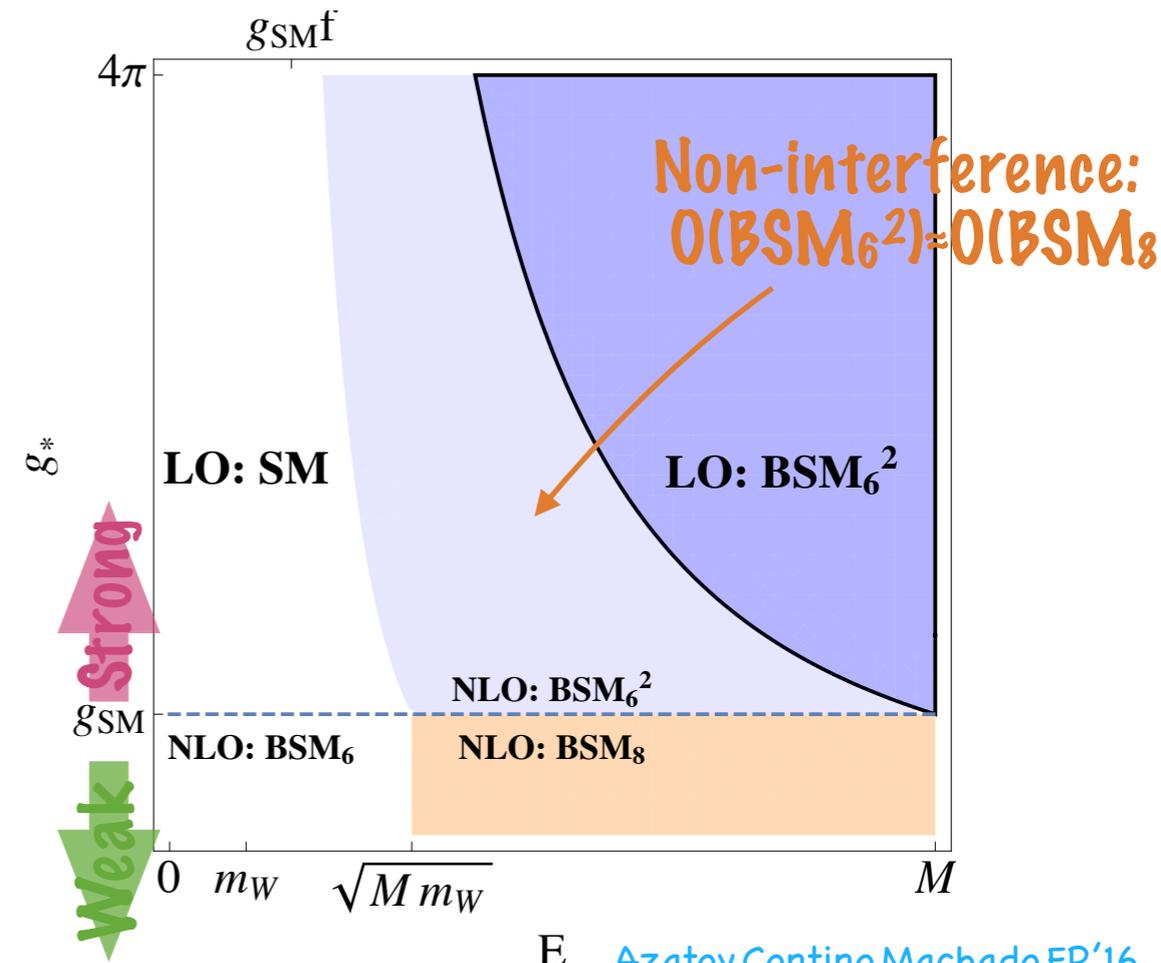doesn't interfere for $m_W \to 0$

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$$\left. + \underbrace{\frac{g_*}{g_{\rm SM}} \frac{E^4}{\Lambda^4}}_{\rm BSM_8 \times SM} + \cdots \right]$$

Non-interference: $O(\rm BSM_6{}^2) \approx O(\rm BSM_8{}$

$g_{\rm SM}$

Strong

Weak

LO: SM

LO: $\rm BSM_6{}^2$

NLO: $\rm BSM_6{}^2$

NLO: $\rm BSM_6$ NLO: $\rm BSM_8$

0 m_W $\sqrt{M\,m_W}$ M

E

Azatov,Contino,Machado,FR'16

In Strongly coupled theories dim-6² might be bigger than dim-8 and analysis with dim-6 only is consistent

Conclusions

In some high-energy (off-resonance) processes, SM and the "leading" BSM (dimension-6 in derivative expansion) do not interfere.

- ▶ (dim-6)^2 same order $O(1/\Lambda^4)$ as dim-8

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In strongly coupled theories dim-6 is coupling-enhanced: $\quad \dfrac{g_*^2}{g^2}\dfrac{E^4}{\Lambda^4} \gg \dfrac{g_*}{g}\dfrac{E^4}{\Lambda^4}$

▶ **In this context present dim-6 analyses make sense**

Liu,Pomarol,Rattazzi,FR'16

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Liu,Pomarol,Rattazzi,FR'16

It would be interesting to explore strategies to make these analyses consistent also for weakly coupled physics, e.g:

1) Study 2 → 3 processes that interfere   see Dixon,Shadmi'93; Azatov et al' in progress
2) Tag polarization to reduce non-interfering piece
3) Improve sensitivity at small energy where mass-effects relevant
4) ...?

...in progress