

Forward Physics at the LHC

1. Exclusive/diffractive Higgs signal: $pp \rightarrow p + H + p$
2. Properties of “soft” interactions
(forward/diffractive physics at the LHC)
3. Return to the exclusive processes
(at the Tevatron and the LHC)

SM discoveries with early LHC data
UCL, March 30th - April 1st 2009

Alan Martin, IPPP, Durham

Advantages of $pp \rightarrow p + H + p$ with $H \rightarrow b\bar{b}$

- If outgoing protons are tagged far from IP then $\sigma(M) = 1 \text{ GeV}$ (mass also from H decay products)
- **Unique** chance to study $H \rightarrow b\bar{b}$:
QCD $b\bar{b}$ bkgd suppressed by $J_z=0$ selection rule
S/B~1 for **SM Higgs** $M < 140 \text{ GeV}$
- Very **clean** environment, even with pile-up---10 ps timing
- **SUSY Higgs**: parameter regions with larger signal **S/B~10**, even regions where conv. signal is challenging and diffractive signal enhanced----**h, H both observable**
- Azimuth angular distribution of tagged p's \rightarrow spin-parity **0^{++}**

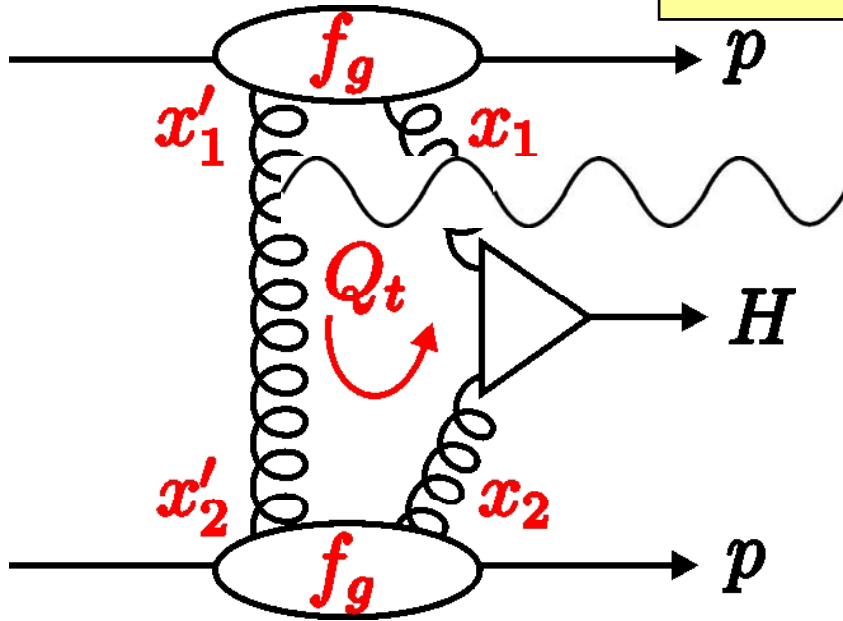
\rightarrow FP420 \rightarrow ATLAS + CMS

Is the cross section large enough ?

How do we calculate $\sigma(pp \rightarrow p + H + p)$?

What price do we pay for an **exclusive** process with large rapidity **gaps** ?

QCD mechanism for $pp \rightarrow p+H+p$



no emission when $(\lambda \sim 1/k_t) > (d \sim 1/Q_t)$
i.e. only emission with $k_t > Q_t$

$$\left(x' \sim \frac{Q_t}{\sqrt{s}}\right) \ll \left(x \sim \frac{M_H}{\sqrt{s}}\right) \ll 1$$

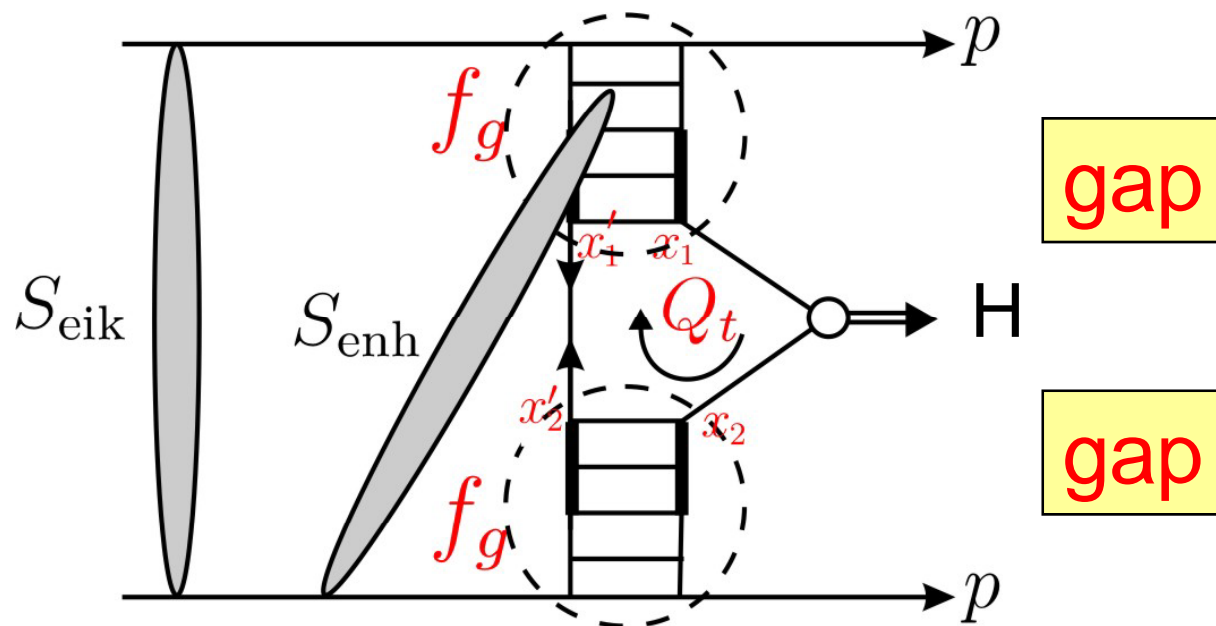
$$\mathcal{M} \sim A(gg \rightarrow H) \int \frac{d^2 Q_t}{Q_t^4} f_g(x_1, x'_1, Q_t^2, \mu^2) f_g(x_2, x'_2, Q_t^2, \mu^2)$$

unintegrated skewed gluons f_g given in terms of $g(x, Q_t^2)$ and Sudakov factor $T = \exp\left(-\int_{Q_t^2}^{\mu^2} dk_t^2 \dots\right)$ which exponentially suppresses infrared region

can use pQCD

$\sigma > 100 \text{ fb} !!$ but....

...but “soft” scatt. can easily destroy the gaps



soft-hard
factorizⁿ

eikonal rescatt: between protons

← conserved

enhanced rescatt: involving intermediate partons

← broken

→ soft physics at high energies

(Dark age?)

Model for “soft” high-energy interactions

needed to ---- understand asymptotics, intrinsic interest
---- describe “underlying” events for LHC jet algor^{ms}
---- calc. rap.gap survival S^2 for exclusive prodⁿ

Model should:

1. be self-consistent theoretically --- satisfy unitarity
 - importance of absorptive corrections
 - importance of multi-Pomeron interactions
2. agree with available soft data
CERN-ISR to Tevatron range $\sigma_{\text{tot}}, \frac{d\sigma_{\text{el}}}{dt}, \frac{d\sigma_{\text{SD}}}{dt dM^2}(pp \rightarrow pX)$
3. include Pomeron comp^{ts} of different size---to study effects of soft-hard factⁿ breaking

Optical theorems

at high energy
use Regge

$$\sigma_{\text{total}} = \sum_X \left| \text{Diagram} \right|^2 = \text{Im} \left[\text{Diagram} \right] = \text{Diagram} \sim g_N^2 \left(\frac{s}{s_0} \right)^{\alpha_P(0)-1}$$

but screening important
so σ_{total} suppressed

High mass diffractive dissociation

$$\left| \text{Diagram} \right|^2 = \text{Diagram} = \text{Diagram} \sim g_N^3 g_{3P} \left(\frac{M^2}{s_0} \right)^{\alpha_P(0)-1} \left(\frac{s}{M^2} \right)^{2\alpha_P(t)-2}$$

triple-Pomeron diag

but screening important

$$(g_{3P})_{\text{effective}} = S^2 (g_{3P})_{\text{bare}} \quad \text{so } (g_{3P})_{\text{bare}} \text{ increased}$$

Must include unitarity

diagonal in $b \sim l/p$

$$S S^\dagger = I \quad \text{with } S = I + iT \quad \rightarrow \quad T - T^\dagger = i T^\dagger T$$

elastic unitarity \rightarrow

$$2 \operatorname{Im} T_{el}(s, b) = |T_{el}(s, b)|^2 + G_{inel}(s, b)$$

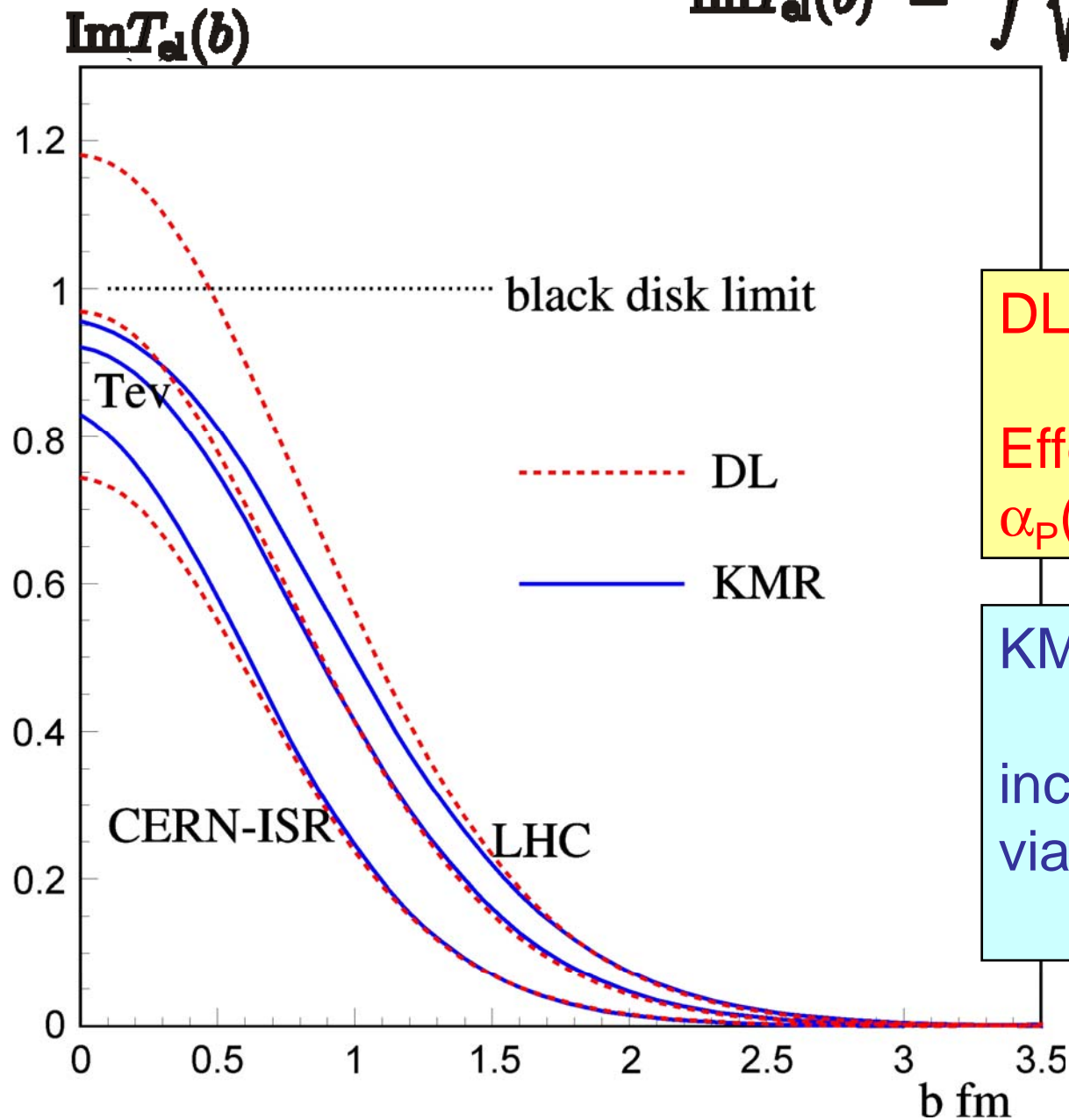
$$\left\{ \begin{array}{l} \frac{d^2 \sigma_{tot}}{d^2 b} = 2 \operatorname{Im} T_{el} = 2(1 - e^{-\Omega/2}) \\ \frac{d \sigma_{el}}{d^2 b} = |T_{el}|^2 = (1 - e^{-\Omega/2})^2 \\ \frac{d \sigma_{inel}}{d^2 b} = 2 \operatorname{Im} T_{el} - |T_{el}|^2 = 1 - e^{-\Omega} \end{array} \right.$$

Opacity / Eikonal $\Omega(s, b) \geq 0$

$$\left. \begin{array}{l} \text{e.g. black disc} \\ \operatorname{Im} T_{el} = 1, \quad b < R \end{array} \right\} \begin{array}{l} \sigma_{tot} = 2\pi R^2 \\ \sigma_{el} = \sigma_{inel} = \pi R^2 \end{array}$$

$e^{-\Omega}$ is the probability of no inelastic interaction

$$\text{Im}T_{el}(b) = \int \sqrt{\frac{d\sigma_{el}}{dt} \frac{16\pi}{1+\rho^2}} J_0(qb) \frac{qdq}{2\pi}$$



DL parametrization:

Effective Pomeron pole
 $\alpha_p(t) = 1.08 + 0.25t$

KMR parametrization

includes absorption
 via multi-Pomeron
 effects

Elastic amp. $T_{el}(s,b)$

bare amp. $\Omega = \overline{\text{I}}$

$$T_{el} = \overline{\text{O}} = 1 - e^{-\Omega/2} = \sum_{n=1}^{\infty} \overline{\text{I} \dots \text{I}} \Omega \quad (-20\%)$$

Low-mass diffractive dissociation

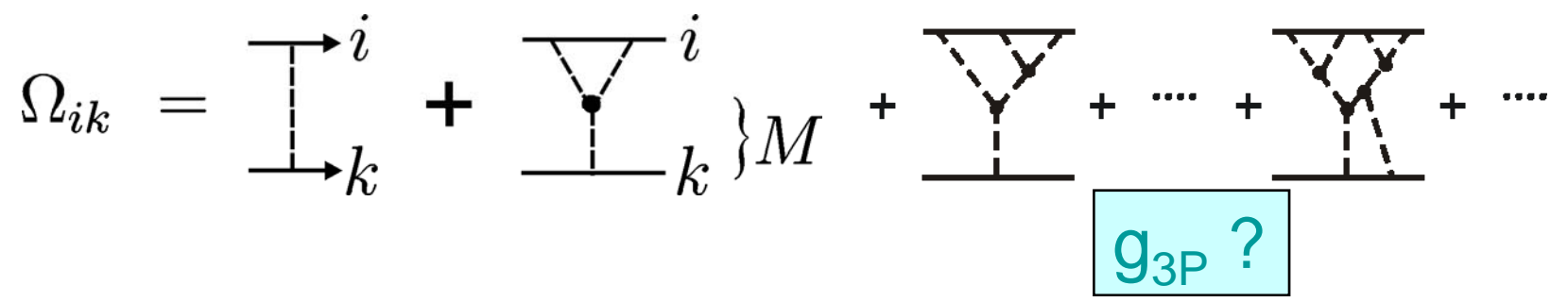


introduce diff^{ve} estates ϕ_i, ϕ_k (comb^{ns} of p, p^*, \dots) which **only** undergo “elastic” scattering (Good-Walker)

$$T_{ik} = \overline{\text{O}}^i_k = 1 - e^{-\Omega_{ik}/2} = \sum \overline{\text{I} \dots \text{I}} \Omega_{ik} \quad (-40\%)$$

include high-mass diffractive dissociation

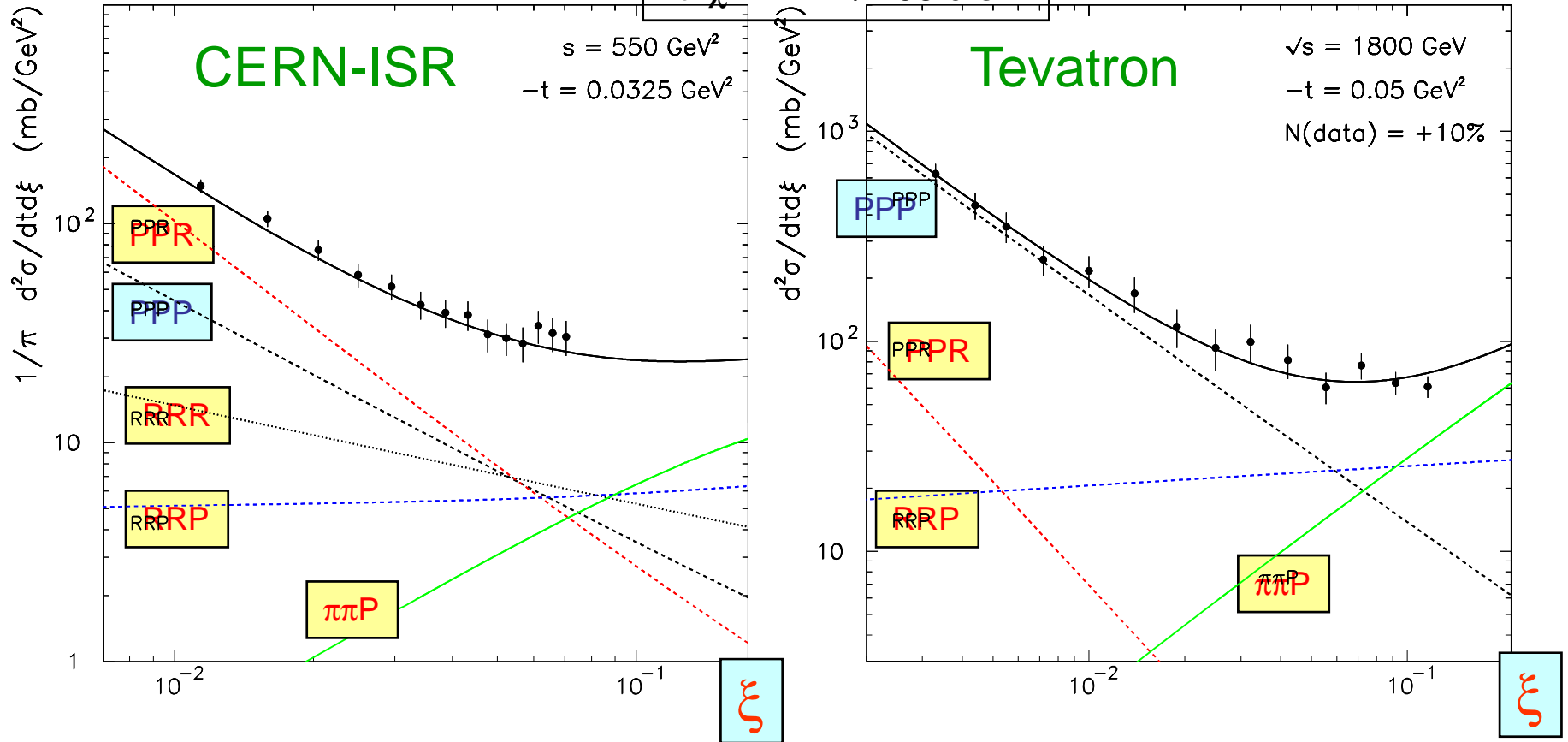
(SD -80%)



triple-Regge analysis of $d\sigma/dtd\xi$, including screening

(includes compilation of SD data by Goulianos and Montanha)

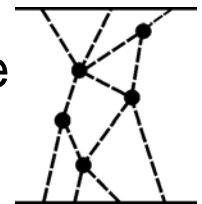
fit: $\chi^2 = 171 / 206$ d.o.f.



LKMR

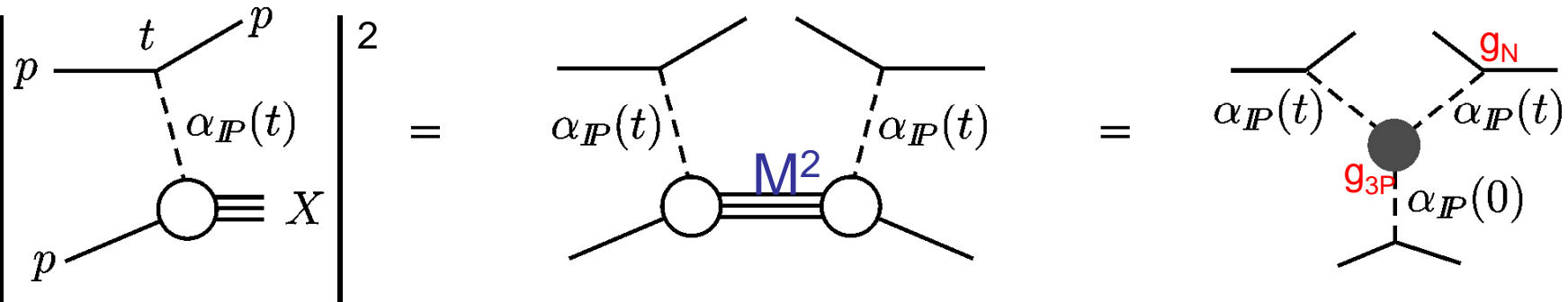
$$g_{3P} = \lambda g_N \quad \lambda \sim 0.2$$

g_{3P} large, need to include multi-Pomeron effects



$$g_{3P} = \lambda g_N \quad \lambda \sim 0.2$$

← large ?



$$M^2 d\sigma_{SD}/dM^2 \sim g_N^3 g_{3P} \sim \lambda \sigma_{el}$$

ln s

$$\sigma_{SD} = \int \frac{M^2 d\sigma_{SD}}{dM^2} \frac{dM^2}{M^2} \sim \underline{\lambda \ln s} \sigma_{el}$$

so at collider energies $\sigma_{SD} \sim \sigma_{el}$

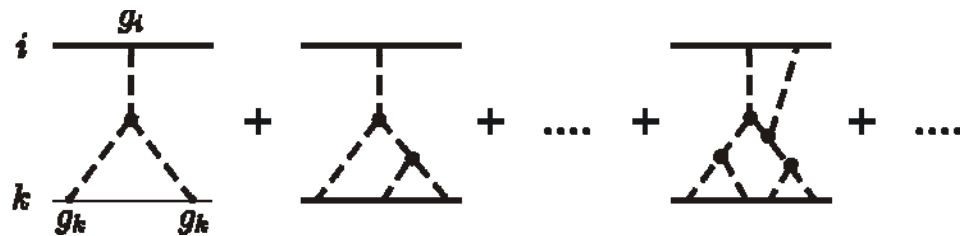
Multi-compt. s- and t-ch analysis of soft data

KMR 2008

$$\sigma_{\text{tot}}, \frac{d\sigma_{\text{el}}}{dt}, \sigma_{\text{SD}}(\text{low } M), \frac{d\sigma_{\text{SD}}}{dt dM^2}$$

model:

- 3-channel eikonal, ϕ_i with $i=1,3$



- include multi-Pomeron diagrams

$$g_m^n = \lambda^{n+m-2} g_N$$

- attempt to mimic BFKL diffusion in $\log q_t$ by including **three components** to approximate q_t distribution – possibility of seeing “soft \rightarrow hard” Pomeron transition

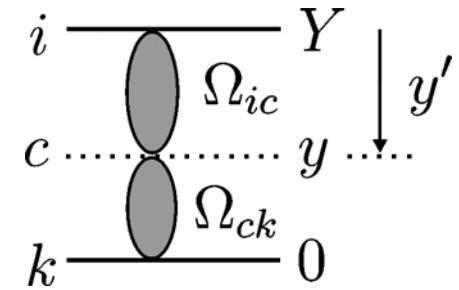
Use four exchanges in the t channel

3 to mimic BFKL diffusion in $\ln q_t$ sec. Reggeon

$$a = P_{\text{large}}, P_{\text{intermediate}}, P_{\text{small}}, R$$

soft \longrightarrow pQCD

average $q_{t1} \sim 0.5, q_{t2} \sim 1.5, q_{t3} \sim 5 \text{ GeV}$



$$V_{RP1} \sim g_{PPR}, g_{RRP}$$

$$V_{PiPj} \sim \text{BFKL}$$

evolve up from $y=0$

bare pole

absorptive effects

$$\left\{ \begin{aligned} \frac{d\Omega_{ck}^a(y)}{dy} &= (\Delta + \alpha' \nabla_b^2) \Omega_{ck}^a(y) e^{-\lambda \Omega_{ck}(y)/2} e^{-\lambda \Omega_{ic}(y')/2} + V_{aa'} \Omega_{ic}^{a'}(y) \\ \frac{d\Omega_{ic}^a(y')}{dy'} &= (\Delta + \alpha' \nabla_b^2) \Omega_{ic}^a(y') e^{-\lambda \Omega_{ic}(y')/2} e^{-\lambda \Omega_{ck}(y)/2} + V_{aa'} \Omega_{ck}^{a'}(y') \end{aligned} \right.$$

evolve down from $y'=Y-y=0$

solve for $\Omega_{ik}^a(y,b)$ by iteration

(arXiv:0812.2407)

Parameters

multi-Pomeron coupling λ from $\xi d\sigma_{SD}/d\xi dt$ data ($\xi \sim 0.01$)

diffractive eigenstates from $\sigma_{SD}(\text{low } M) = 2\text{mb}$ at $\sqrt{s} = 31\text{ GeV}$,
-- equi-spread in R^2 , and t dep. from $d\sigma_{el}/dt$

Results

All soft data well described

$g_{3P} = \lambda g_N$ with $\lambda = 0.25$

$\Delta_{P_i} = 0.3$ (close to the BFKL NLL resummed value)

$\alpha'_{P_1} = 0.05\text{ GeV}^{-2}$

These values of the **bare** Pomeron trajectory yield, after screening, the expected soft Pomeron behaviour ---

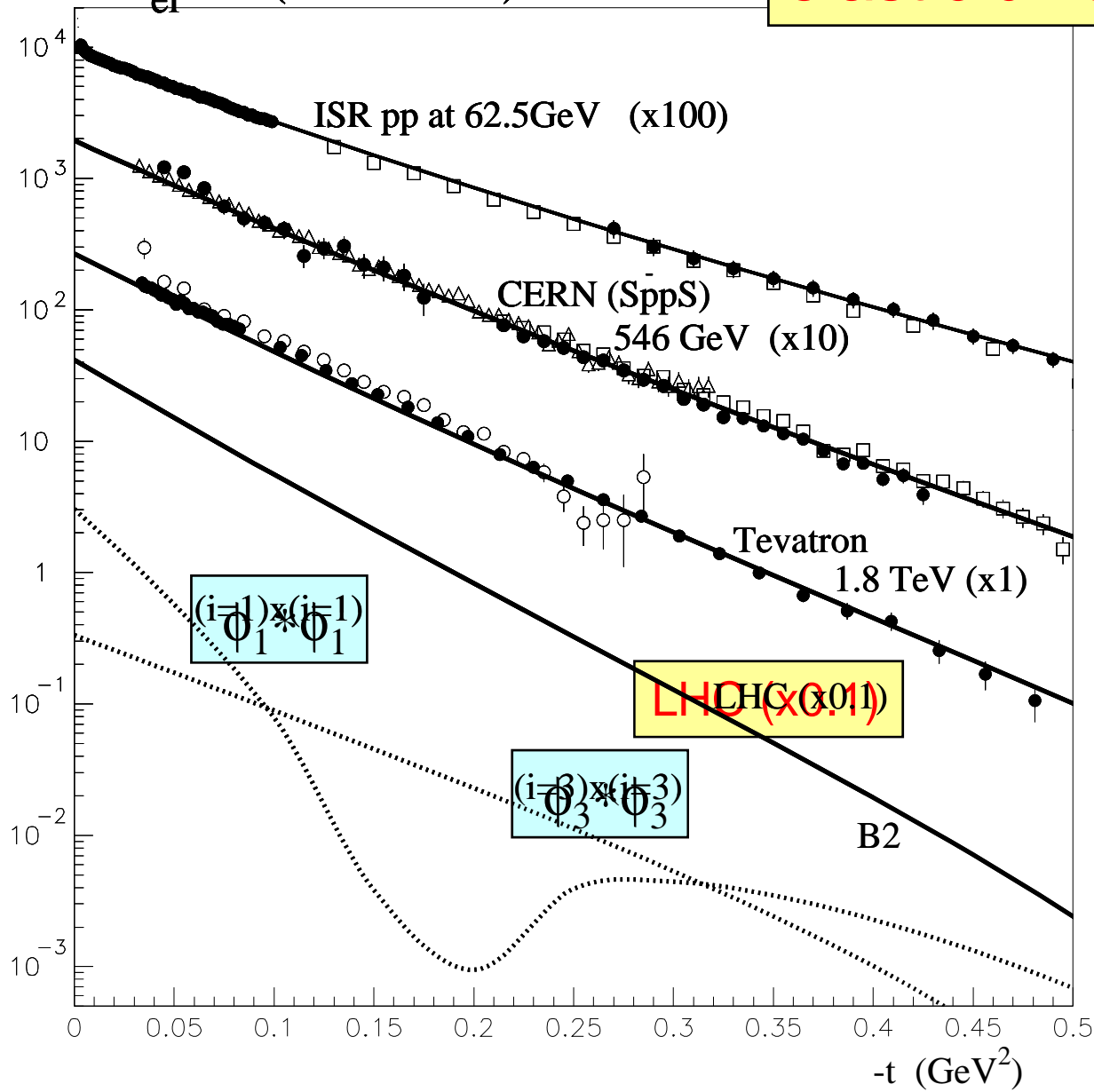
“soft-hard” matching (since P_1 heavily screened, ..., $P_3 \sim$ bare)

$\Delta_R = -0.4$ (as expected for secondary Reggeon)

$$\Delta = \alpha(0) - 1$$

$d\sigma_{el}/dt$ (mb/GeV²)

elastic differential $d\sigma/dt$



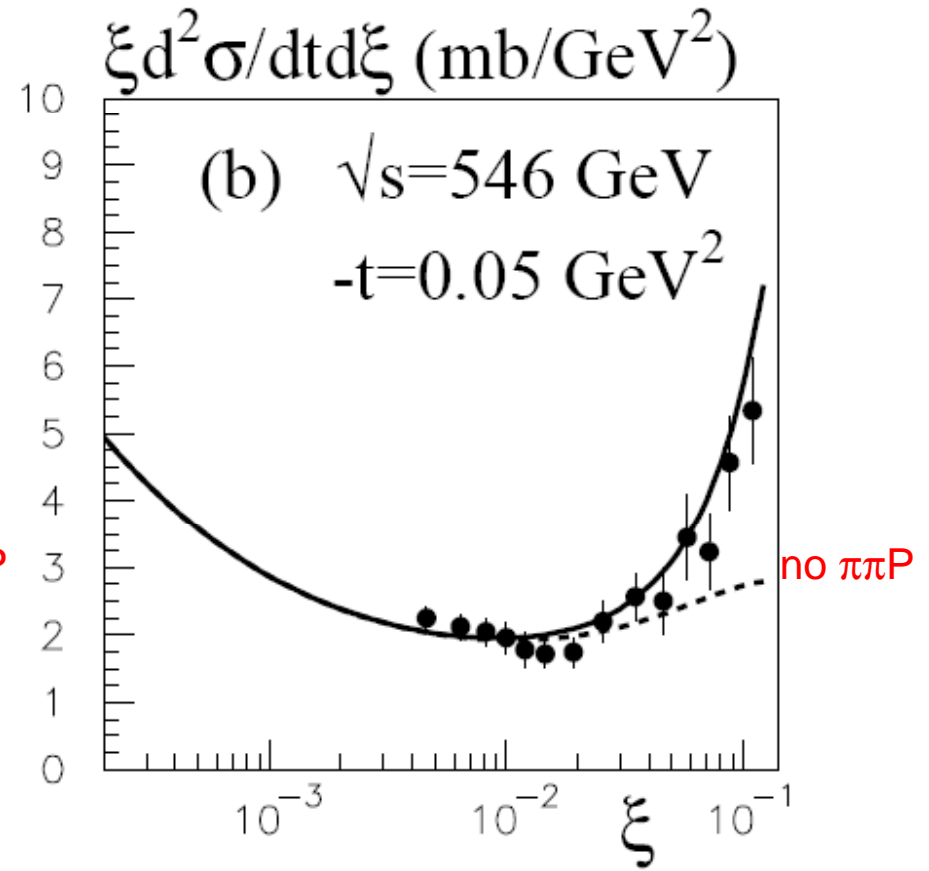
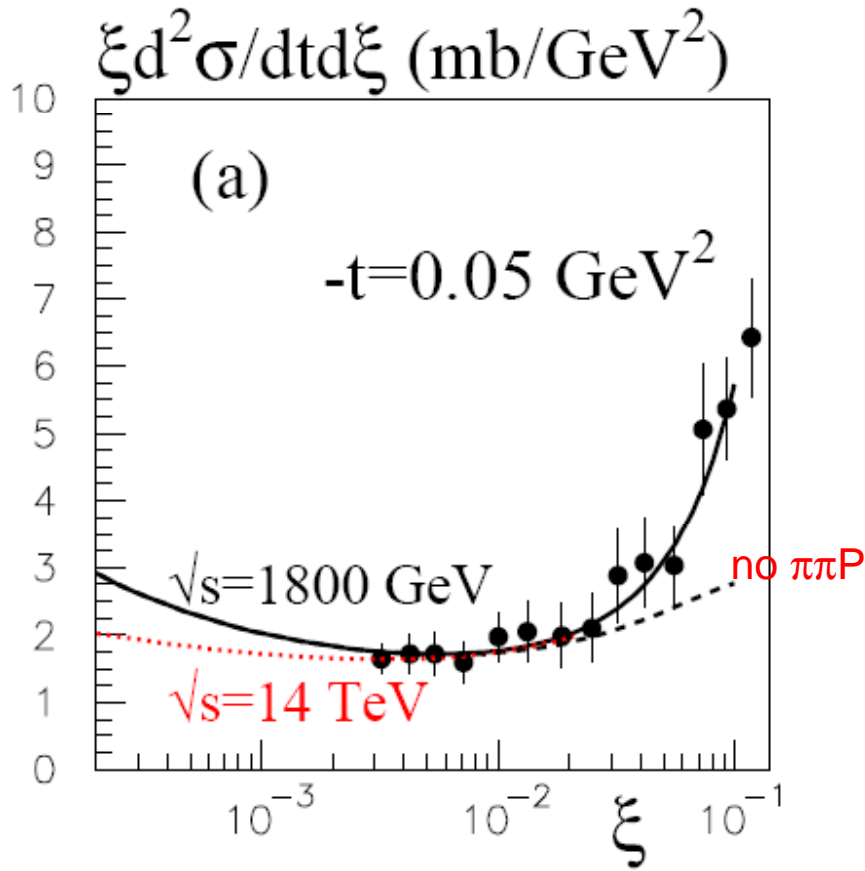
~ g, sea

ϕ_1 : "large"

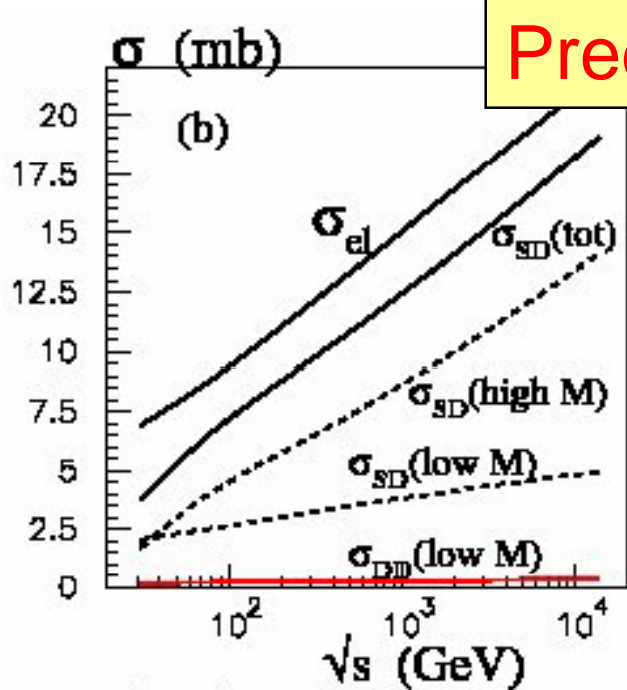
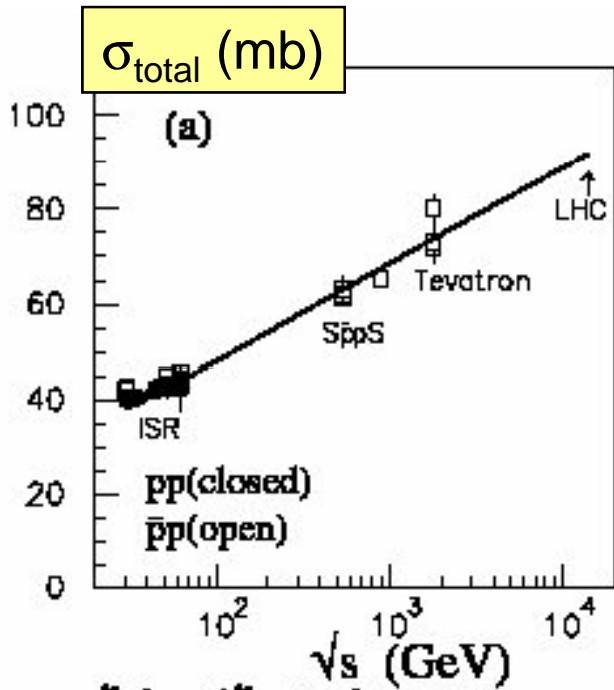
ϕ_3 : "small"

more valence

Description of CDF dissociation data

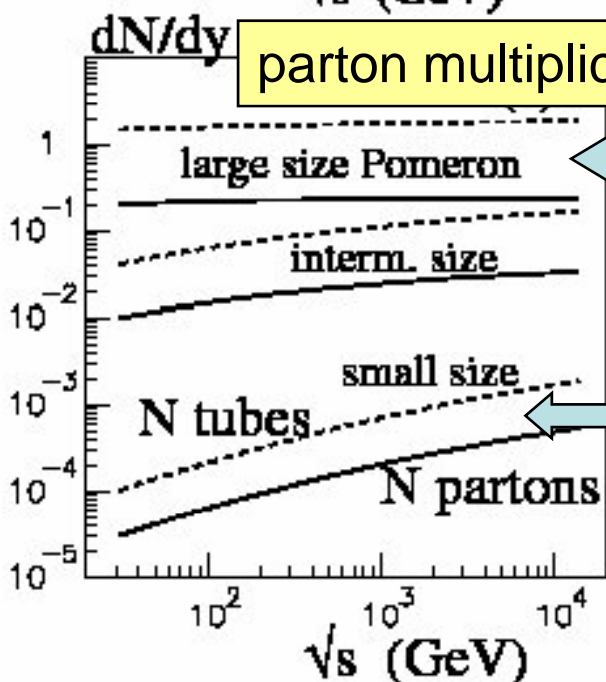
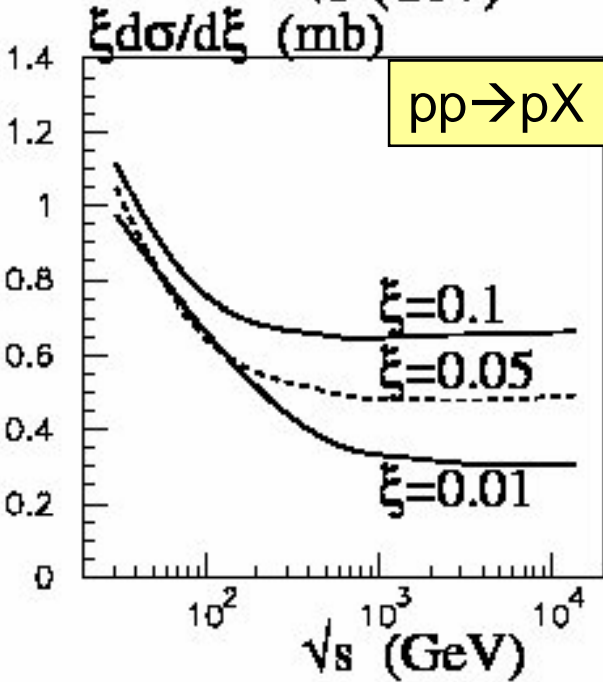


Predictions for LHC



$\sigma_{\text{total}} = 91.7 \text{ mb}^*$
 $\sigma_{\text{el}} = 21.5 \text{ mb}$
 $\sigma_{\text{SD}} = 19.0 \text{ mb}$

*see also Sapeta, Golec-Biernat; Gotsman et al.

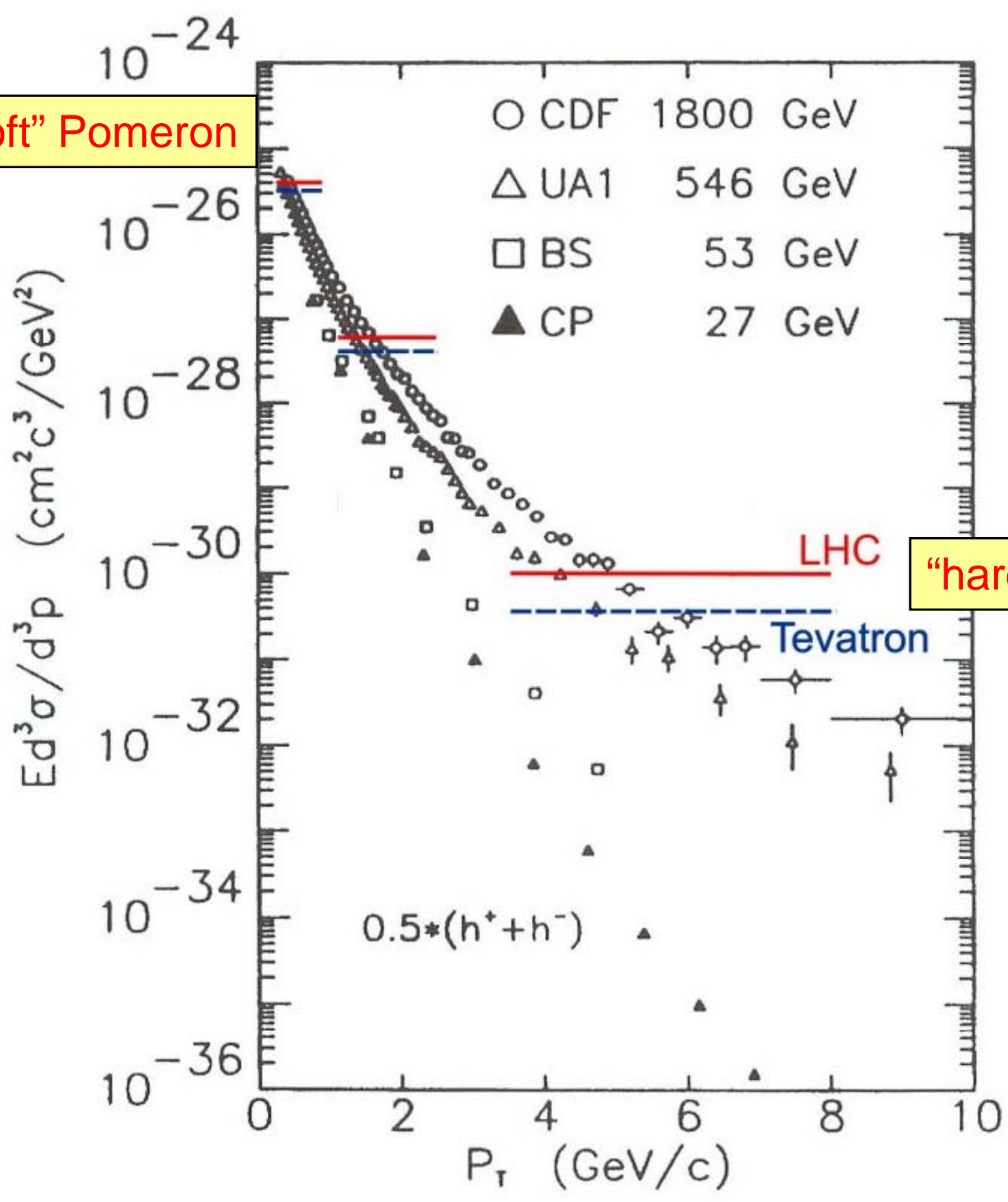


All Pom. compts have $\Delta_{\text{bare}}=0.3$

“soft”, screened, little growth, partons saturated

“hard” ~ no screening much growth, $s^{0.3}$

“soft” Pomeron

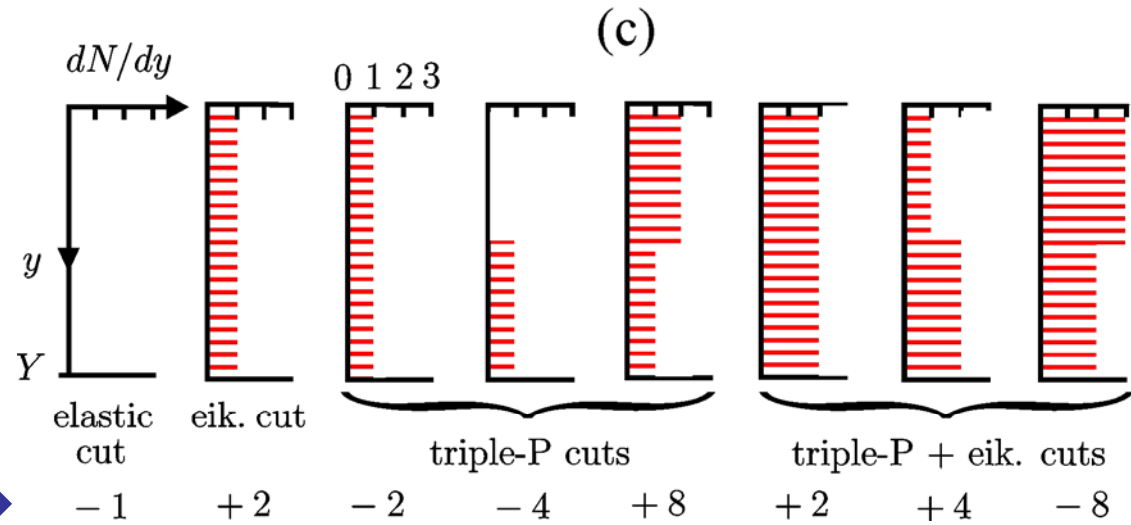
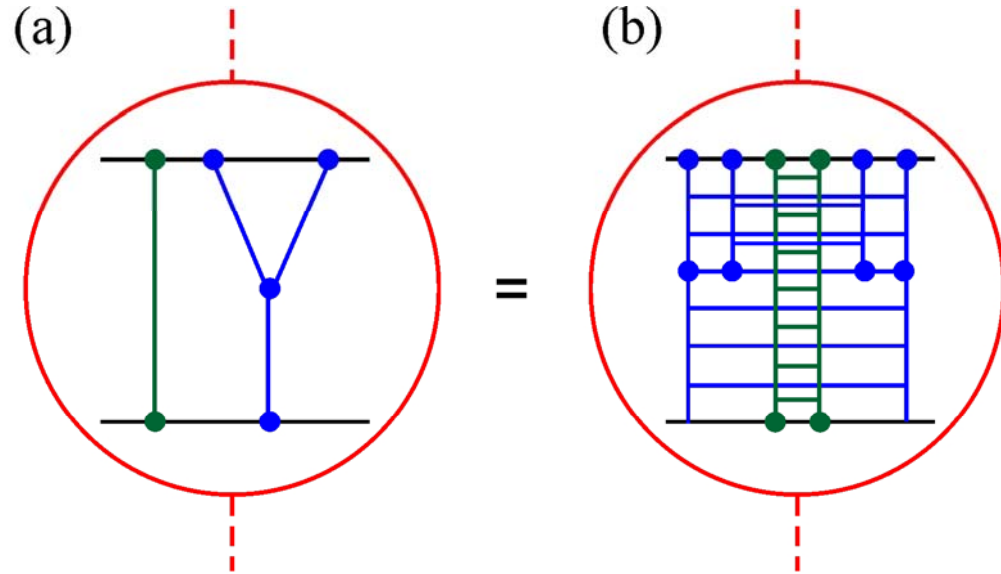


“hard” Pomeron

Multi-Pomeron effects at the LHC

Each multi-Pomeron
diag. simultaneously
describes several
different processes

Example



8 different “cuts”
AGK cutting rules →

Long-range correlations at the LHC

cutting n eikonal Pomerons \rightarrow multiplicity n times that cutting one Pomeron

\rightarrow long range correlation
even for large rapidity differences $|y_a - y_b| \sim Y$

$$R_2 = \frac{\sigma_{\text{inel}} d^2\sigma / dy_a dy_b}{(d\sigma / dy_a)(d\sigma / dy_b)} - 1 = \frac{d^2N / dy_a dy_b}{(dN / dy_a)(dN / dy_b)} - 1$$

$$\rightarrow R_2 > 0$$

without multi-Pomeron exch. $R_2 > 0$ only when two particles are close, e.g. from resonance decays
(short-range correlation)

Calculation of S^2_{eik} for $pp \rightarrow p + H + p$

prob. of proton to be in diffractive estate i

over b

average over diff. estates i,k

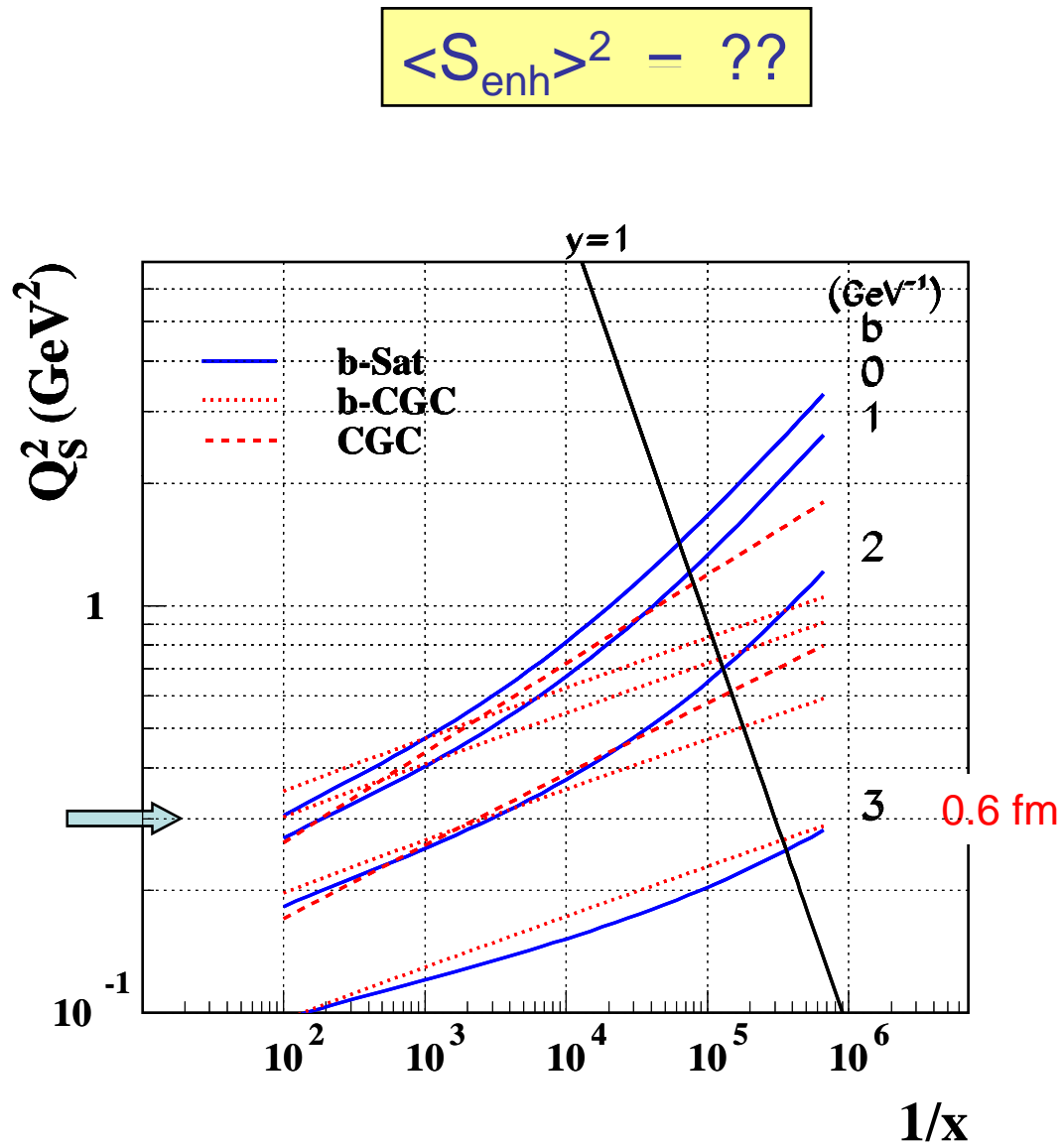
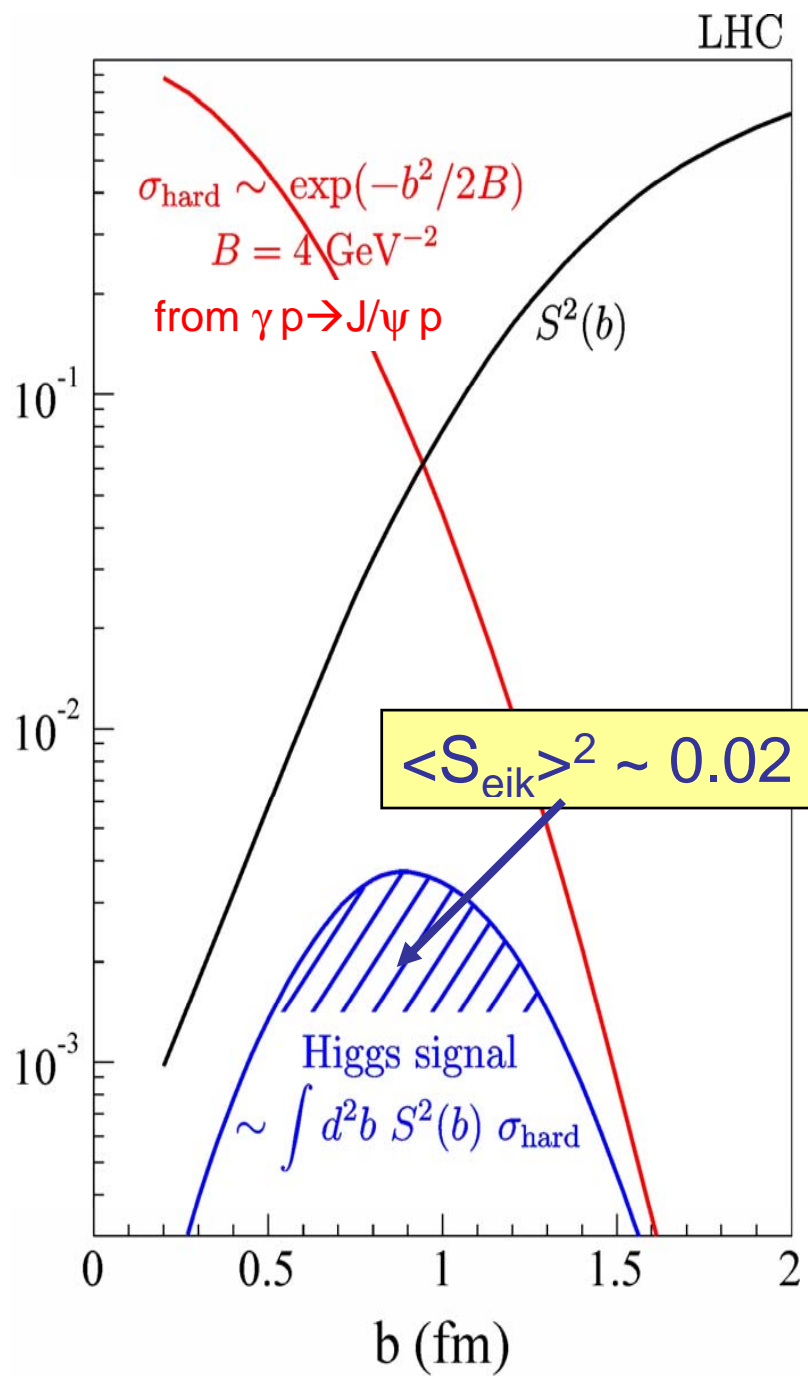
hard m.e. $i k \rightarrow H$

survival factor w.r.t. soft $i-k$ interaction

$$\overline{S^2} = \frac{\sum_{i,k} \int d^2b |a_{pi}|^2 |a_{p'k}|^2 |\mathcal{M}_{ik}|^2 \exp(-\Omega_{ik}(s, b))}{\sum_{i,k} \int d^2b |a_{pi}|^2 |a_{p'k}|^2 |\mathcal{M}_{ik}|^2}$$

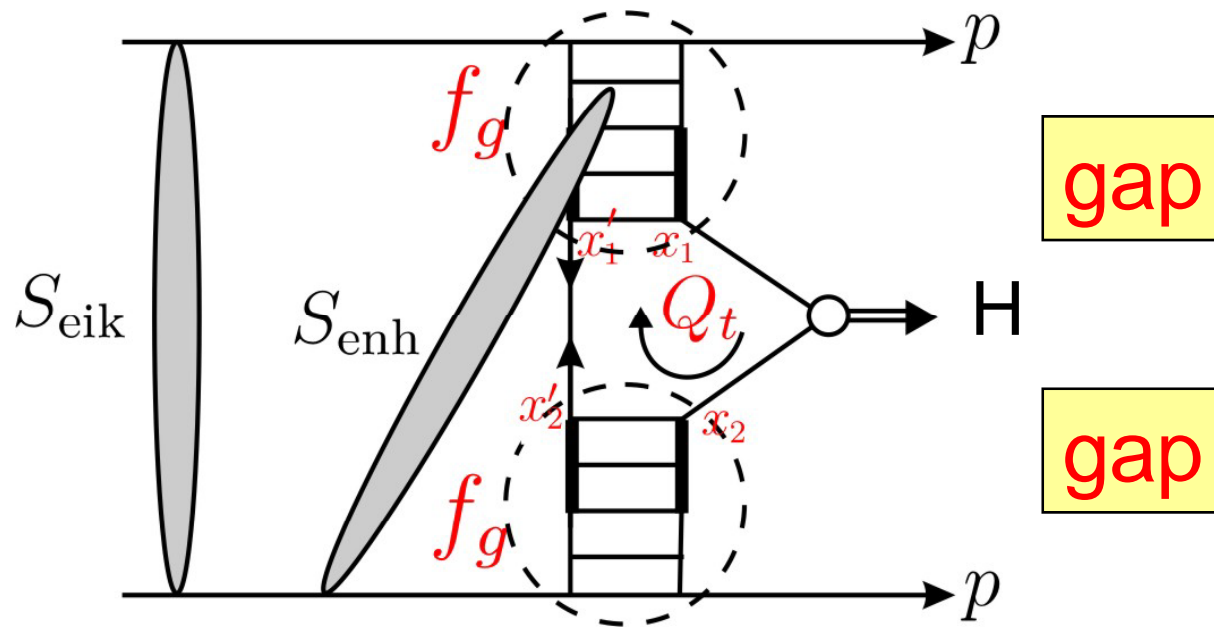
$\overline{S^2}_{eik} \sim 0.02$ for 120 GeV SM Higgs at the LHC

$\rightarrow \sigma \sim 2 - 3 \text{ fb at LHC}$



Watt, Kowalski

Calculation of S^2_{enhanced} for $pp \rightarrow p + H + p$



eikonal rescatt: between protons

enhanced rescatt: involving intermediate partons

soft-hard
factorizⁿ
← conserved
← broken

The new soft analysis, with Pomeron q_t structure,
enables S^2_{enh} to be calculated

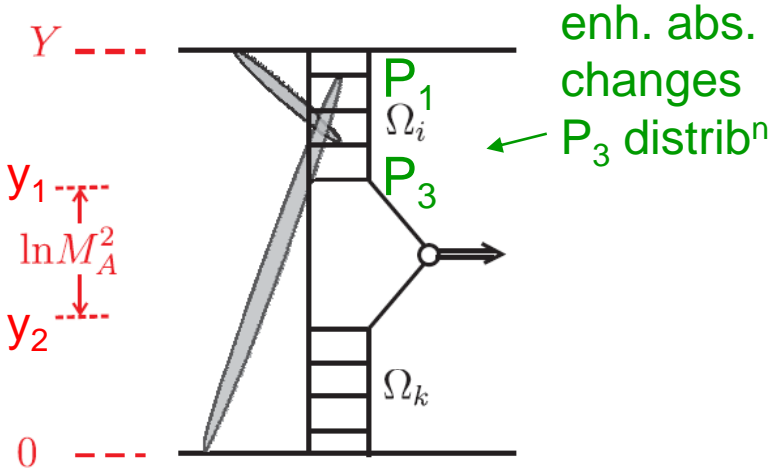
model has 4 t-ch. exchanges

3 to mimic BFKL diffusion in $\ln q_t$

$$a = P_{\text{large}}, P_{\text{intermediate}}, P_{\text{small}}, R$$

soft \longrightarrow pQCD

average $q_{t1} \sim 0.5$, $q_{t2} \sim 1.5$, $q_{t3} \sim 5$ GeV



$$V_{RP1} \sim g_{PPR}, g_{RRP}$$

$$V_{PiPj} \sim \text{BFKL}$$

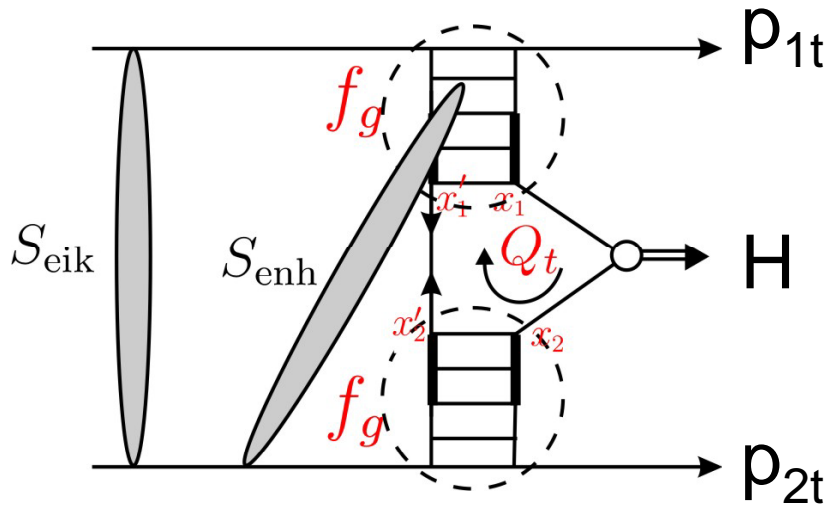
evolve up to y_2

$$\left\{ \begin{aligned} \frac{d\Omega_{ck}^a(y)}{dy} &= \underbrace{(\Delta + \alpha' \nabla_b^2) \Omega_{ck}^a(y)}_{\text{bare pole}} \underbrace{e^{-\lambda \Omega_{ck}(y)/2} e^{-\lambda \Omega_{ic}(y')/2}}_{\text{absorptive effects}} + V_{aa'} \Omega_{ic}^{a'}(y) \\ \frac{d\Omega_{ic}^a(y')}{dy'} &= (\Delta + \alpha' \nabla_b^2) \Omega_{ic}^a(y') \underbrace{e^{-\lambda \Omega_{ic}(y')/2} e^{-\lambda \Omega_{ck}(y)/2}}_{\text{absorptive effects}} + V_{aa'} \Omega_{ck}^{a'}(y') \end{aligned} \right.$$

evolve down to y_1

~ solve with and without abs. effects

Survival prob. for $pp \rightarrow p+H+p$



$$\langle S_{eik}^2 \rangle \sim 0.02 \quad \text{consensus}$$

$$\langle S_{enh}^2 \rangle \sim 0.01 - 1$$

controversy

KMR 2008 \rightarrow

$$\langle S^2 \rangle_{tot} = \langle S_{eik}^2 S_{enh}^2 \rangle \sim 0.015$$

(B=4 GeV⁻²)

However enh. abs. changes p_t behaviour from exp form, so

$$\langle S^2 \rangle_{tot} \langle p_t^2 \rangle^2 = \left\{ \begin{array}{l} 0.0015 \quad \text{LHC} \\ 0.0030 \quad \text{Tevatron} \end{array} \right\} \text{KMR 2000 (no } S_{enh})$$

$$\left\{ \begin{array}{l} 0.0010 \quad \text{LHC} \\ 0.0025 \quad \text{Tevatron} \end{array} \right\} \text{KMR 2008 (with } S_{enh})$$

see [arXiv:0812.2413](https://arxiv.org/abs/0812.2413)

Comments on S^2

1. Enhanced rescattering reduces the signal by ~30%
2. However, the quoted values of S^2 are **conservative** lower limits
3. The very small values of S^2_{enh} in recent literature are **not** valid

The arguments are given in [arXiv:0903.2980](https://arxiv.org/abs/0903.2980)

CDF observation of exclusive processes at the Tevatron offers the first experimental checks of the formalism

Observation of exclusive prodⁿ, $\bar{p}p \rightarrow \bar{p} + A + p$, by CDF

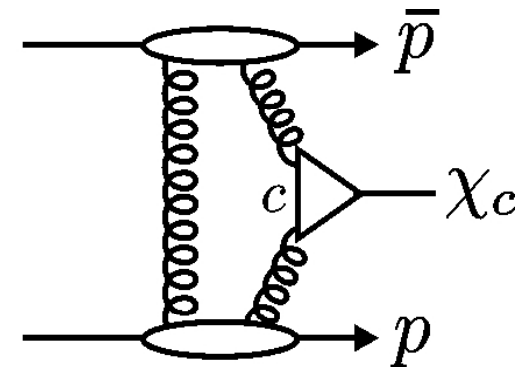
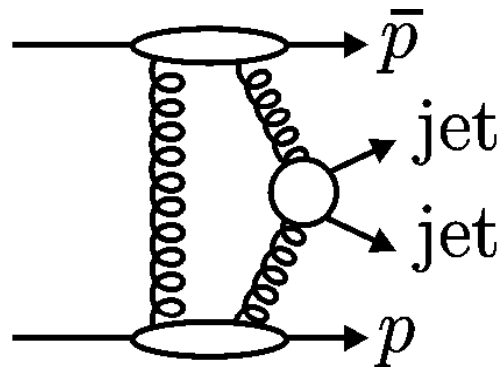
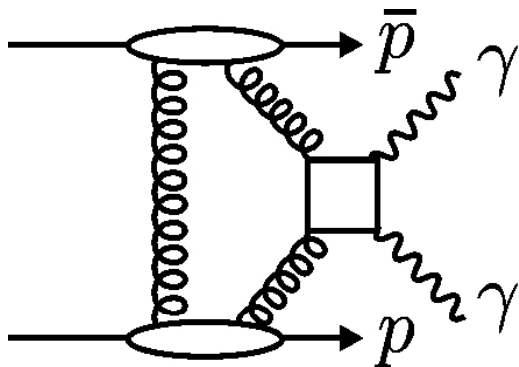
with $A=\gamma\gamma$

or

$A = \text{dijet}$

or

$A = \chi_c \rightarrow J/\psi\gamma \rightarrow \mu+\mu-\gamma$



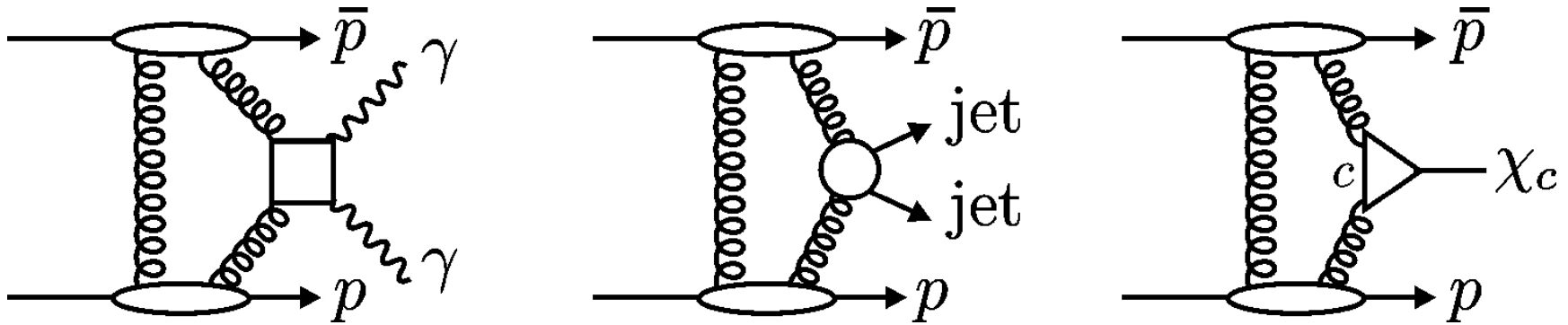
Same mechanism as $pp \rightarrow p+H+p$

$$\mathcal{M}(\bar{p}p \rightarrow \bar{p} + A + p) \sim S^2 \int \frac{d^2 Q_t}{Q_t^4} f_g f_g$$

tho' pred^{ns} become more unreliable as M_A becomes smaller, and infrared Q_t region not so suppressed by Sudakov factor

$$T = \exp \left(- \int_{Q_t^2}^{M_A^2} dk_t^2 \dots \right)$$

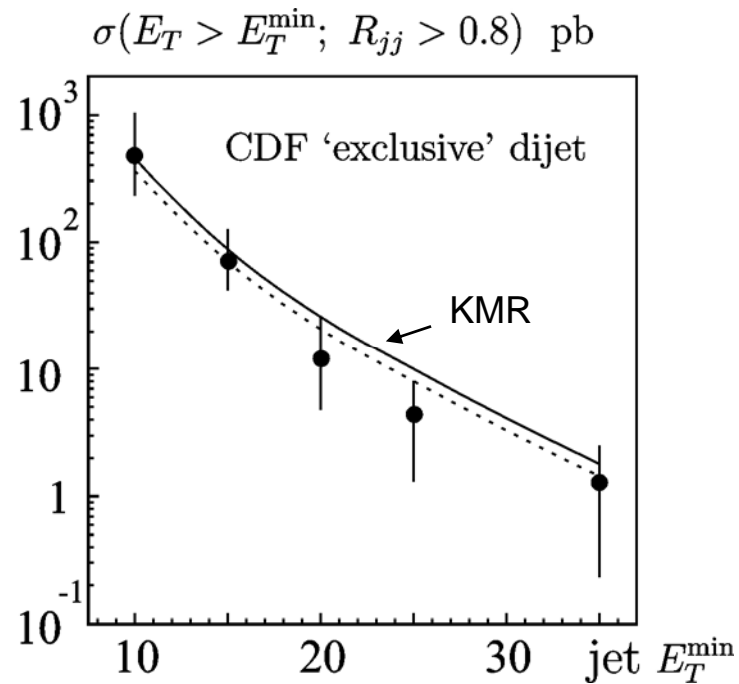
Observation of exclusive prodⁿ, $\bar{p}p \rightarrow \bar{p} + A + p$, at Tevatron



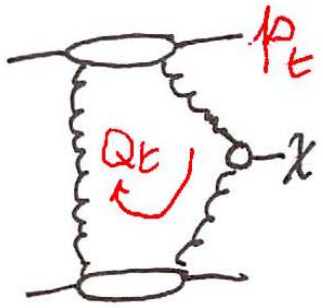
KMR cross section predictions are consistent with CDF data

3 events observed
 (one due to $\pi^0 \rightarrow \gamma\gamma$)
 $\sigma(\text{excl } \gamma\gamma)_{\text{CDF}} \sim 0.09\text{pb}$
 $\sigma(\text{excl } \gamma\gamma)_{\text{KMR}} \sim 0.04\text{pb}$

$\sigma(\gamma\gamma) = 10 \text{ fb}$ for
 $E_T^\gamma > 14 \text{ GeV}$ at LHC



Exclusive $\bar{p}p \rightarrow \bar{p} + \chi_c + p$



CDF: $\chi_c \rightarrow J/\psi \gamma \rightarrow \mu^+ \mu^- \gamma$
 (with 0.06 written below γ)

$$\left. \frac{d\sigma_x}{dy} \right|_{y=0} = 76 \pm 14 \text{ nb}$$

B. fractions: $0^{++} \quad 0.013$
 $1^{++} \quad 0.36$
 $2^{++} \quad 0.20$

$\sim \langle p_E^2 \rangle / M^2$
 $\sim \langle p_E^2 \rangle^3 / Q_E^4$

$\sim 30 \text{ nb}$
 "
 "

even tho'
 $\frac{\chi_0}{\chi_1} \sim \frac{\chi_0}{\chi_2} \sim 10-40$

The KMRS predⁿ is reduced by $S_{\text{enh}}^2 \sim 1/3$ and by 1.45 due to a revised $\Gamma_{\text{tot}}(\chi_c(0))$

KMRS: **only order-mag. pred^{ns}**
 (light M_x , non-p QCD effects)

Better decay channels $\chi_c \rightarrow \pi\pi$ or $K\bar{K}$

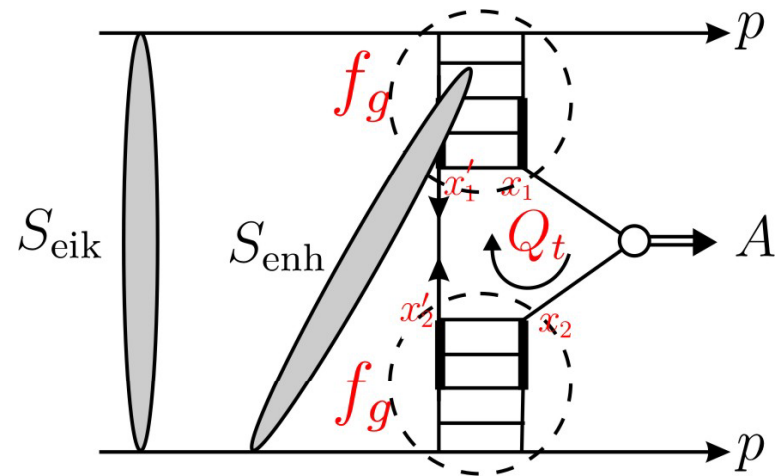
B. fractions $0^{++} \quad \sim 1\%$
 $1^{++} \quad \text{forbidden}$
 $2^{++} \quad \text{suppressed}$

$\chi_b ?$

Early LHC runs can give detailed checks of all of the ingredients of the calculation of $\sigma(pp \rightarrow p + A + p)$, sometimes even **without** proton taggers

Early LHC checks of
theoretical formalism for
 $pp \rightarrow p + A + p$?

$$\sigma \sim S^2 \left| \int \frac{dQ_t^2}{Q_t^4} f_g f_g \right|^2$$



Possible checks of:

(i) survival factor S^2 : W +gaps, Z +gaps

(ii) generalised gluon f_g : $\gamma p \rightarrow Yp$, 3 central jets

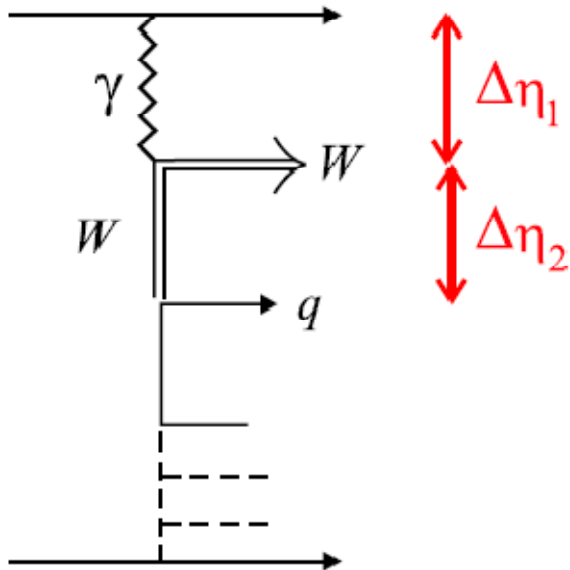
(iii) soft-hard factorisation
(broken by enhanced
absorptive effects)

$\frac{\#(A+\text{gap}) \text{ evts}}{\#(\text{inclusive } A) \text{ evts}}$
 with $A = W, \text{ dijet}, Y \dots$

(arXiv:0802.0177)

W+gaps

$$\int \frac{dk_t^2 k_t^2}{(|t_{\min}| + k_t^2)^2} \quad \text{with} \quad |t_{\min}| \simeq \frac{m_N^2 \xi^2}{1 - \xi}$$



Even without a proton tag
 $x_p = 1 - \xi$ can be measured by

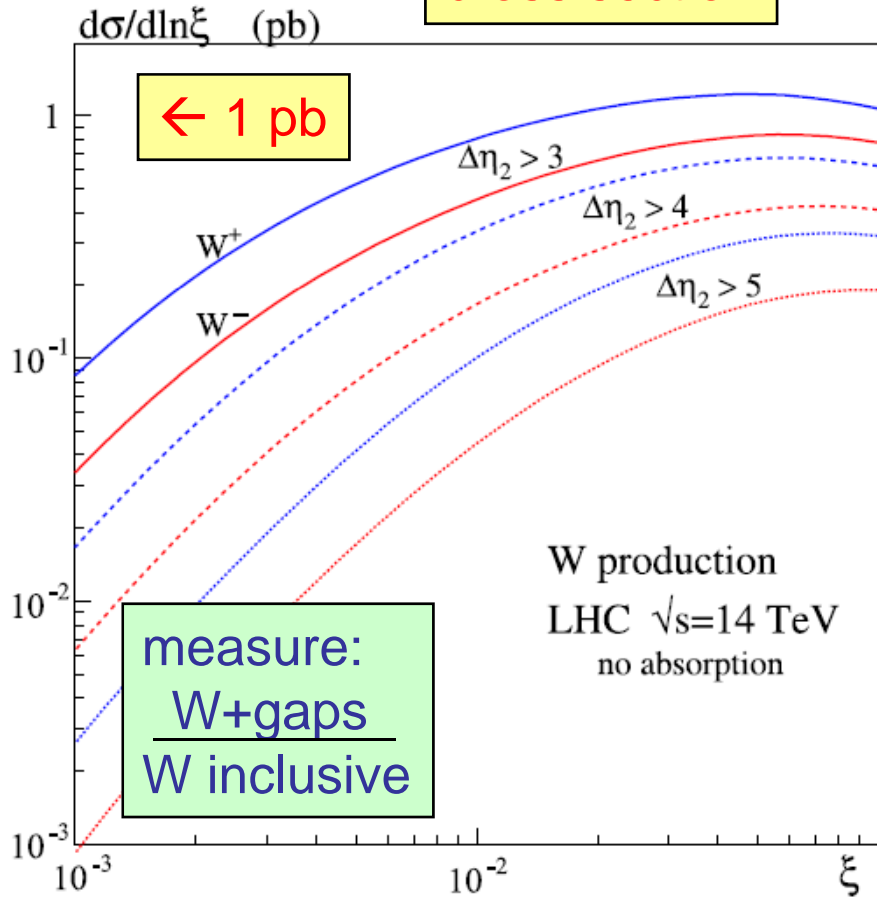
$$\xi = \sum \sqrt{m_i^2 + k_{ti}^2} e^{y_i} / \sqrt{s}$$

successfully used by CDF

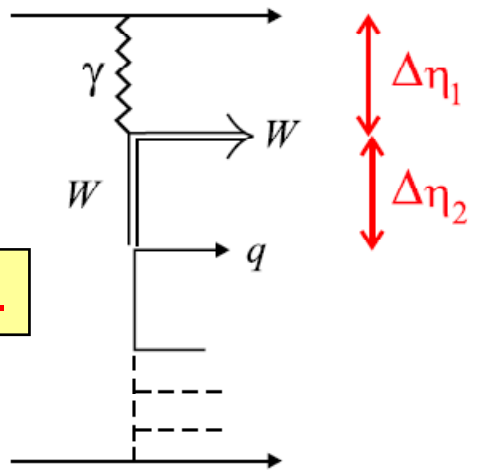
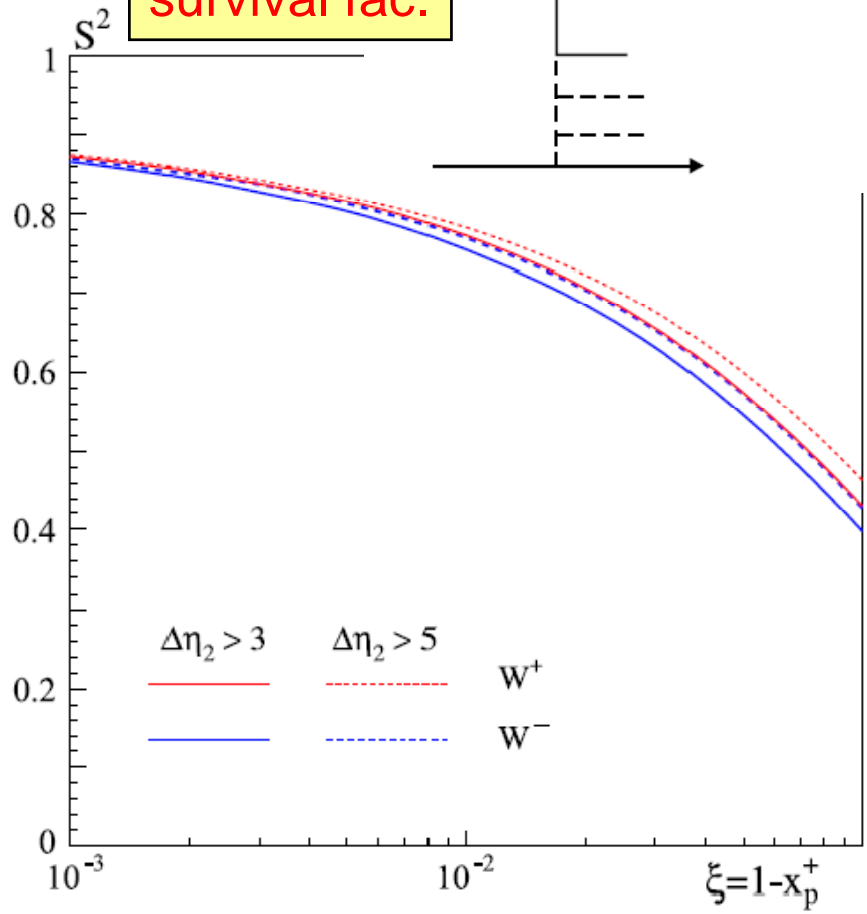
$$\eta_W = 2.3(-2.3) \longrightarrow \xi \sim 0.1(0.001)$$

W+gaps

cross section



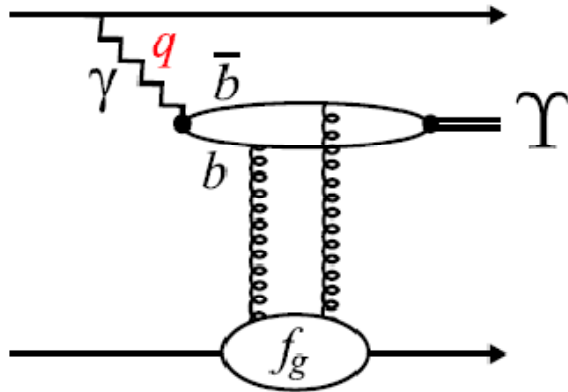
survival fac.



S^2 large, as large b_t (small opacity)

Exclusive Y production as probe of f_g

(a) γ exch

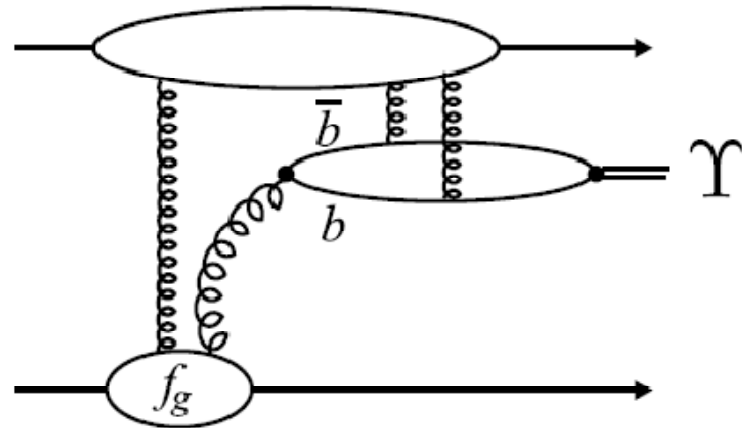


$$\left. \frac{d\sigma^{(a)}}{dy} \right|_{y=0} \simeq 50 \text{ pb}$$

$\times 0.025$ (br for $Y \rightarrow \mu\mu$)

Bzdak, Motyka, Szymanowski, Cudell

(b) odderon exch



comparable ?

can separate by p_t if a tag
of upper proton is done
(odderon has larger p_t)

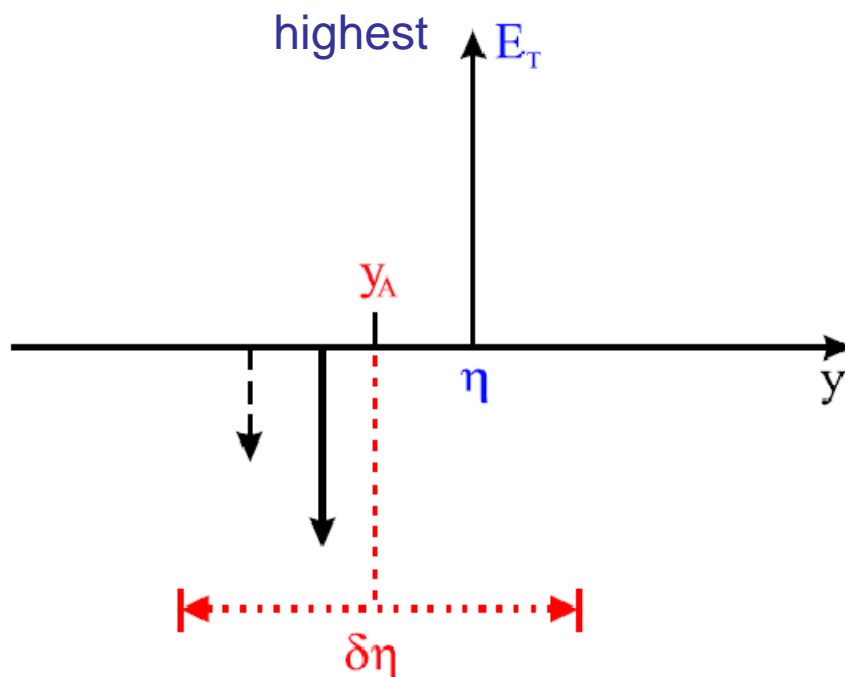
If $|y_Y| < 2.5$, then sample
 $f_g(x_1, x_2)$ with x_i in $(10^{-4}, 10^{-2})$

3-jet events as probe of Sudakov factor T

T is prob. **not** to emit additional gluons in gaps: $pp \rightarrow p + A + p$
 $T = \exp(-n)$, where **n** is the mean # gluons emitted in gap

3 central jets → allow check of additional gluon emission

System A must be **colourless** – so optimum choice is emission of **third** jet in high E_T dijet production

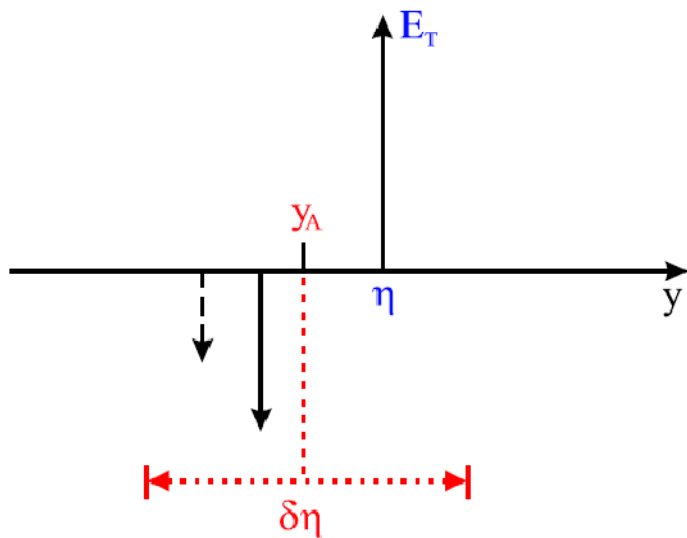


$$R_j = 2E_T (\cosh \eta^*) / M_A$$

$$M_A = \sqrt{s} \xi^+ \xi^-$$

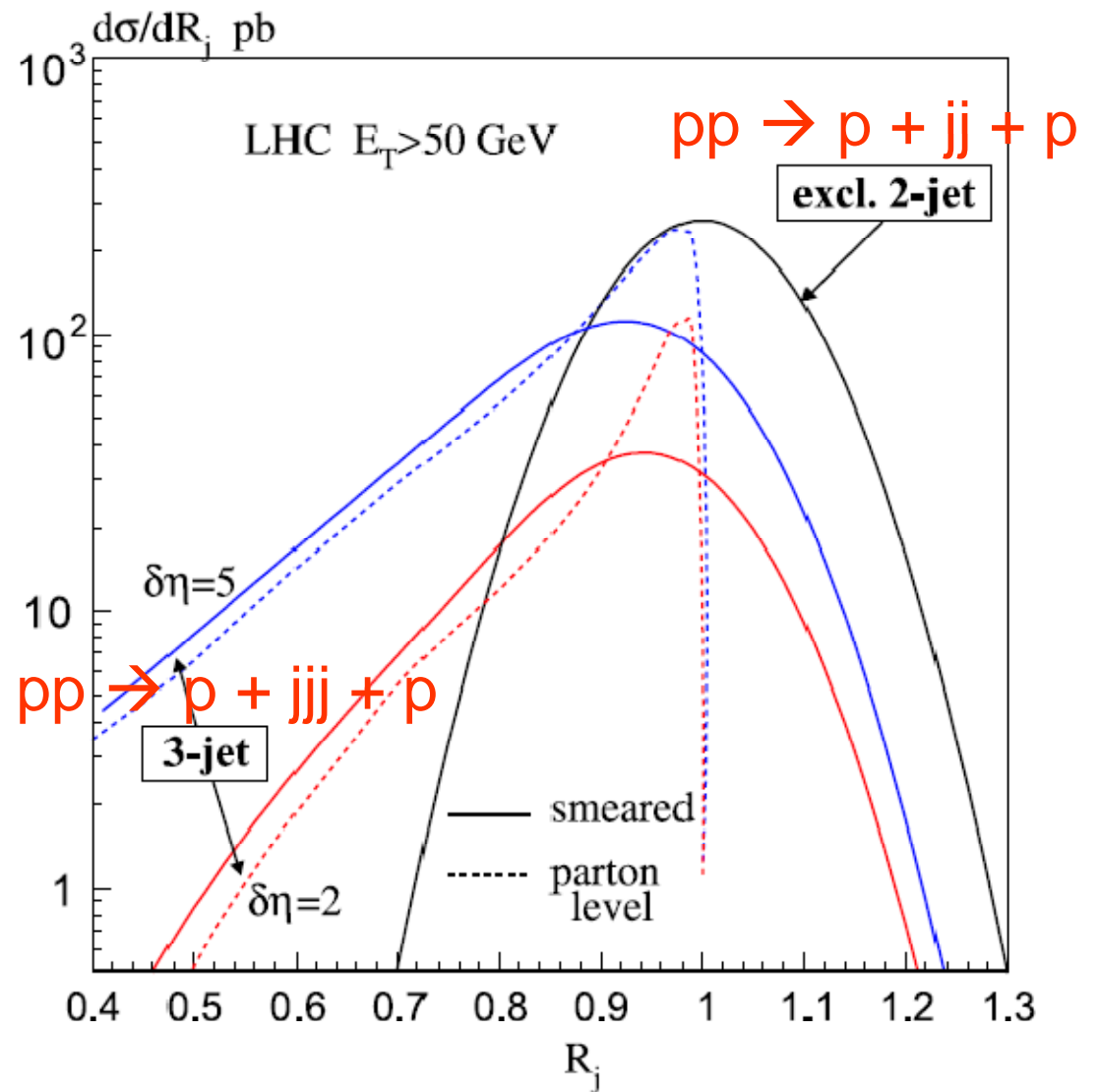
$$\eta^* = \eta - y_A$$

only highest E_T jet used –
 stable to hadronization,
 final parton radiation...



study both $\delta\eta$ and E_T dependence of central 3-jet production

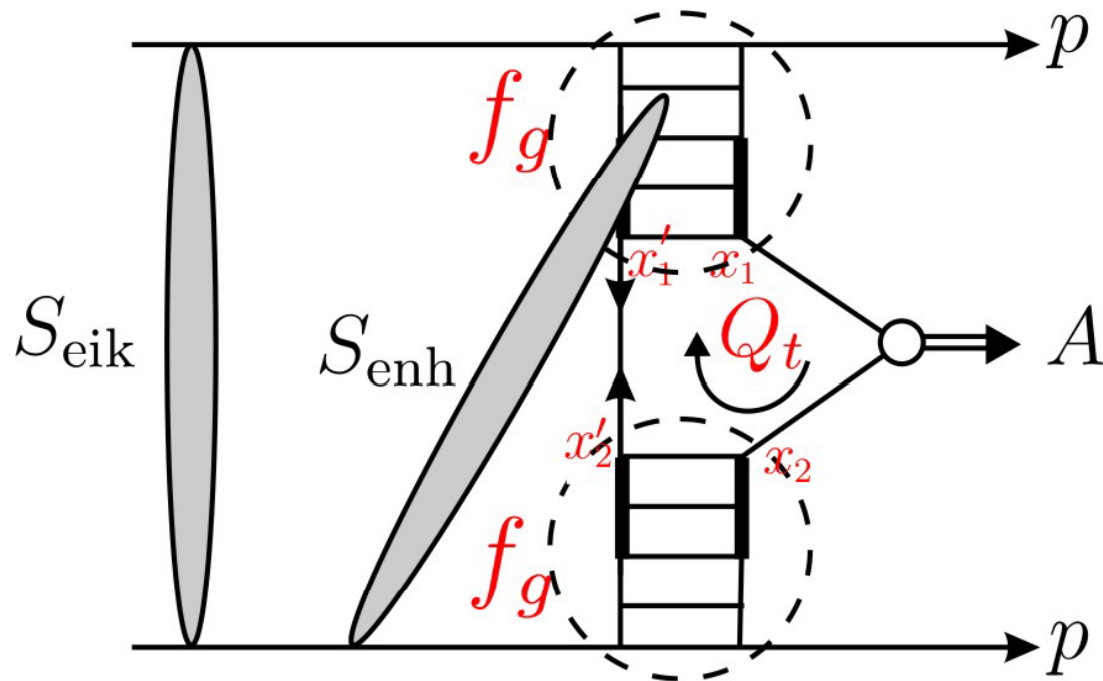
(negligible DPE background)



“Enhanced” absorptive effects

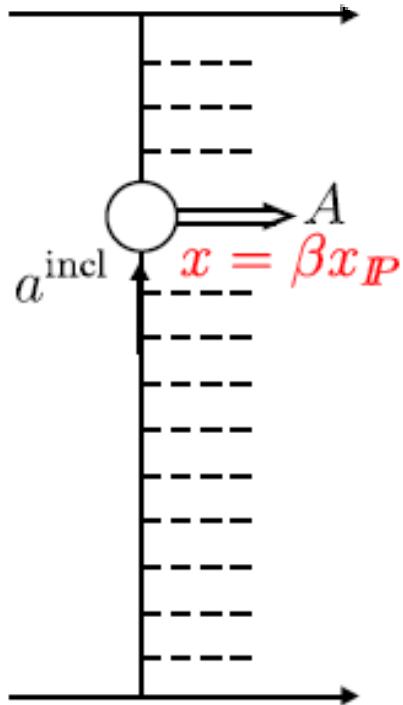
(break soft-hard factorization)

rescattering on an intermediate parton:

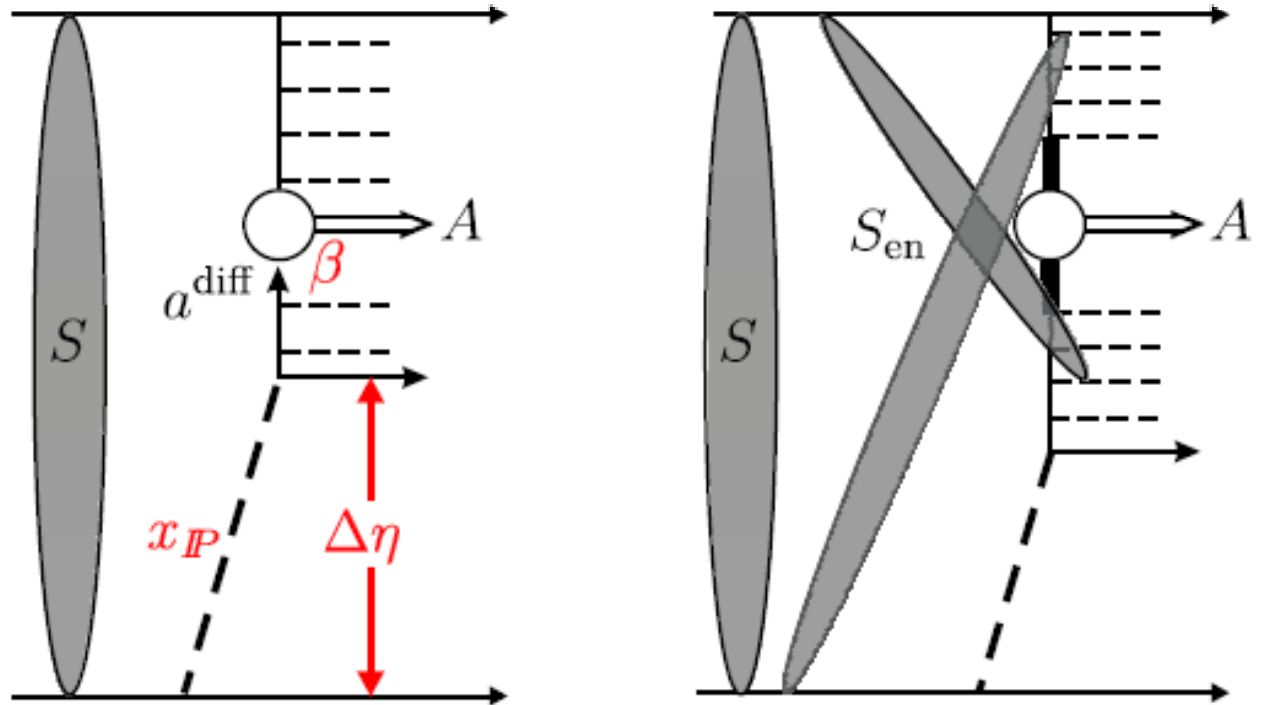


can LHC probe this effect ?

inclusive



diffractive



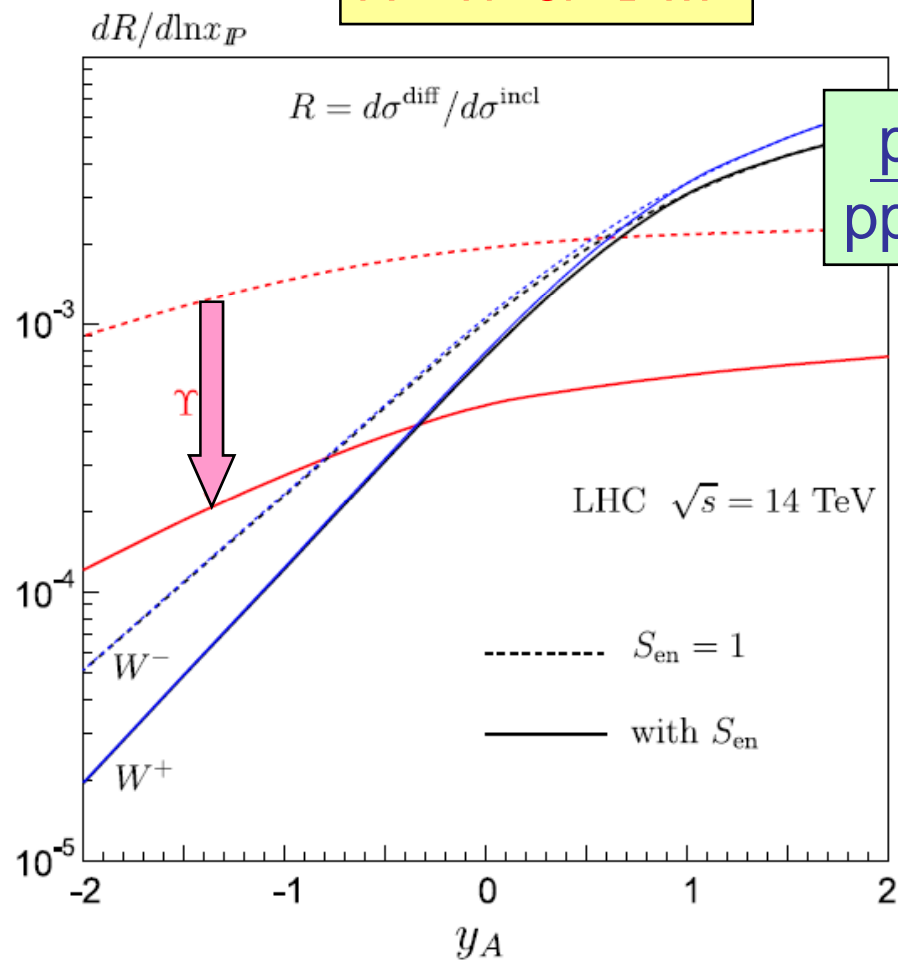
A – W or dijet or Y

$$R = \frac{\text{no. of } (A + \text{gap}) \text{ events}}{\text{no. of (inclusive } A) \text{ events}} = \frac{a^{\text{diff}}(x_{\mathbb{P}}, \beta, \mu^2)}{a^{\text{incl}}(x = \beta x_{\mathbb{P}}, \mu^2)} \langle S^2 S_{\text{en}}^2 \rangle_{\text{over } b_t}$$

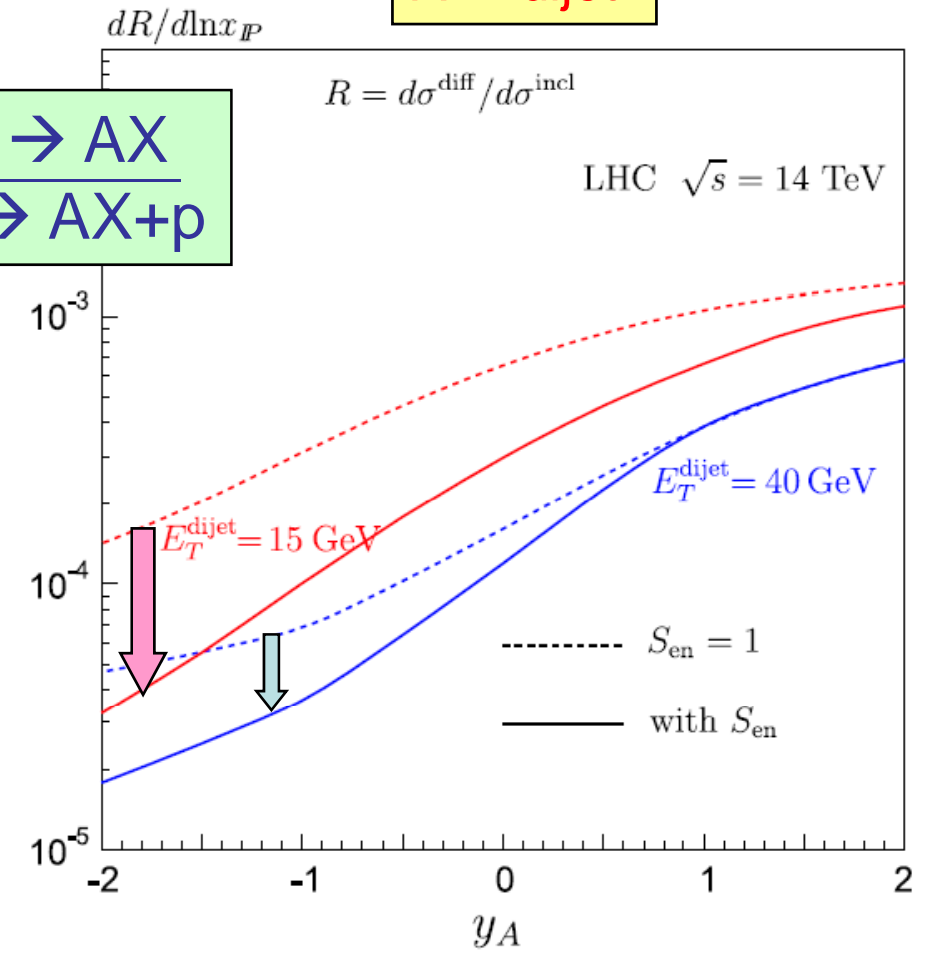
known from HERA

A = W or Y

A = dijet



pp → AX
pp → AX+p



rough estimates of enhanced absorption S_{en}^2

Conclusions – soft processes at the LHC

- screening/unitarity/absorptive corrections are **vital**
- Triple-Regge analysis with screening $\rightarrow g_{3P}$ increased by ~ 3
 \rightarrow **importance of multi-Pomeron diagrams**
- Latest analysis of all available “soft” data:
multi-ch eikonal + multi-Regge + compts of Pom. to mimic BFKL
(showed some LHC predictions $\sigma_{\text{total}} \sim 90 \text{ mb}$)
soft-hard Pomeron transition emerges
 - “soft” compt. --- heavily screened --- little growth with s
 - “intermediate” compt. --- some screening
 - “hard” compt. --- little screening --- large growth (\sim pQCD)
- **LHC can probe “soft” int^{ns}** \rightarrow i.e. probe multi-Pomeron struct.
via long-range rapidity correlations
or via properties of multi-gap events etc.

Conclusions – exclusive processes at the LHC

soft analysis allows rapidity gap survival factors to be calculated for any hard diffractive process

Exclusive central diffractive production, $pp \rightarrow p+H+p$, at LHC has great advantages, $S/B \sim O(1)$, but $\sigma \sim \text{few fb}$ for SM Higgs. However, some SUSY-Higgs have signal enhanced by 10 or more. **Very exciting possibility, if proton taggers installed at 420 m**

Formalism consistent with CDF data for $pp(\text{bar}) \rightarrow p + A + p(\text{bar})$
with $A = \text{dijet}$ and $A = \gamma\gamma$ and $A = \chi_c$
More checks with higher M_A valuable.

Processes which can probe all features of the formalism used to calculate $\sigma(pp \rightarrow p+A+p)$, may be observed in the **early LHC runs**, even without proton taggers