Introduction to Higgs & Beyond

The Standard Model The Higgs boson What may lie above and beyond it?





The Brout-Englert- Higgs Mechanism

• Postulated effective potential:

$$V[\phi] = -\mu^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2$$

• Minimum energy at non-zero value:

$$\phi_0 = <0|\phi|0> = \frac{1}{\sqrt{2}} \begin{pmatrix} 0\\ +v \end{pmatrix} v = \sqrt{\frac{-\mu^2}{\lambda}}$$

- Components of Higgs field: $\phi(x) = \frac{1}{\sqrt{2}}(v + \sigma(x))e^{i\pi(x)}$
- π m = 0, σ massive: $m_H^2 = 2\mu^2 = 2\lambda v$ Higgs boson
- After gauging: $M_W = \frac{g v}{2}$ Massive gauge boson
- Couple to fermions: non-zero masses: $M_f = y_f \frac{v}{\sqrt{2}}$

Summary of the Standard Model

• Particles and SU(3) × SU(2) × U(1) quantum numbers:

+ $|D_{\mu}\phi|^2 - V(\phi)$

L_L E_R	$ \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L, \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}_L, \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}_L \\ e_R^-, \mu_R^-, \tau_R^- \end{pmatrix} $	(1,2, -1) (1,1, -2)
Q_L U_R D_R	$ \begin{pmatrix} u \\ d \end{pmatrix}_{L}, \begin{pmatrix} c \\ s \end{pmatrix}_{L}, \begin{pmatrix} t \\ b \end{pmatrix}_{L} $ $ u_{R}, c_{R}, t_{R} $ $ d_{R}, s_{R}, b_{R} $	$(\mathbf{3,2,+1/3})$ $(\mathbf{3,1,+4/3})$ $(\mathbf{3,1,-2/3})$

• Lagrangian: $\mathcal{L} = -\frac{1}{4} F^{a}_{\mu\nu} F^{a\ \mu\nu} + i\bar{\psi} D\psi + h.c. + \psi_{i}y_{ij}\psi_{j}\phi + h.c.$

matter fermions Yukawa interactions Higgs potential

gauge interactions

Untested before 2012

1975 A Phenomenological Profile of the Higgs Boson

• First attempt at systematic survey

A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS ** CERN, Geneva

Received 7 November 1975

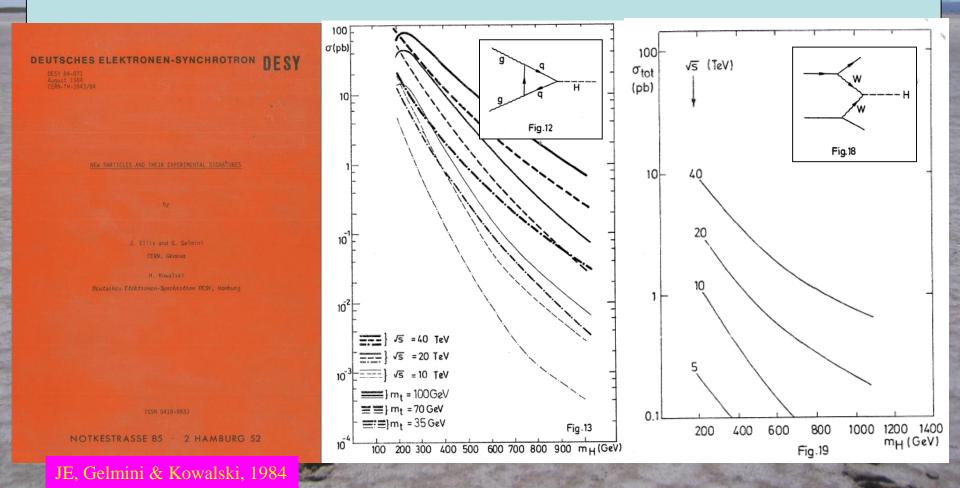
A discussion is given of the production, decay and observability of the scalar Higgs boson H expected in gauge theories of the weak and electromagnetic interactions such as the Weinberg-Salam model. After reviewing previous experimental limits on the mass of

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

A Preview of the Higgs Boson @ LHC

• Prepared for LHC Lausanne workshop 1984

1984



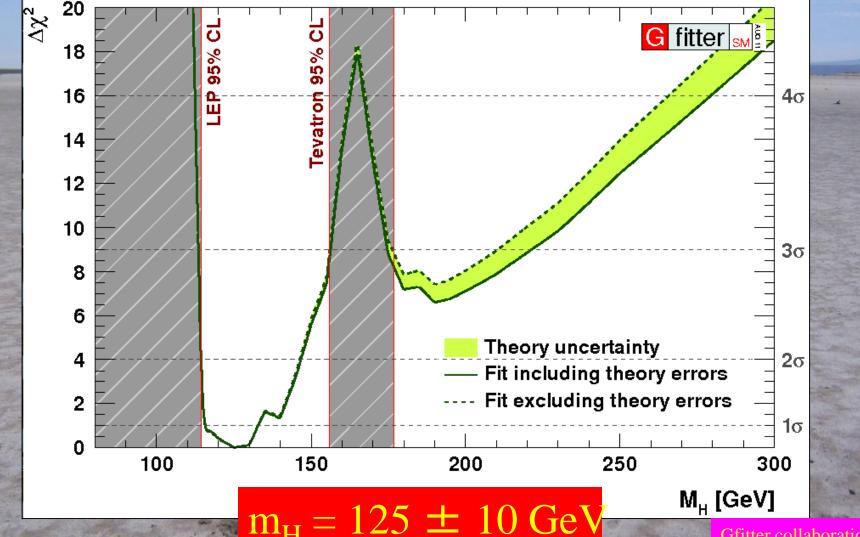
Constraints on Higgs Mass

- Electroweak observables sensitive via quantum loop corrections: $m_W^2 \sin^2 \theta_W = m_Z^2 \cos^2 \theta_W \sin^2 \theta_W = \frac{\pi \alpha}{\sqrt{2} G_F} (1 + \Delta r)$
- Sensitivity to top, Higgs masses:

$$\frac{3\mathbf{G}_F}{8\pi^2\sqrt{2}}m_t^2 \qquad \frac{\sqrt{2}\mathbf{G}_F}{16\pi^2}m_W^2(\frac{11}{3}\ln\frac{M_H^2}{m_Z^2}+\ldots), \, M_H >> m_W$$

 Preferred Higgs mass: m_H ~ 100 ± 30 GeV
 Compare with lower limit from direct search at LEP: m_H > 114 GeV and exclusion around (160, 170 GeV) at TeVatron

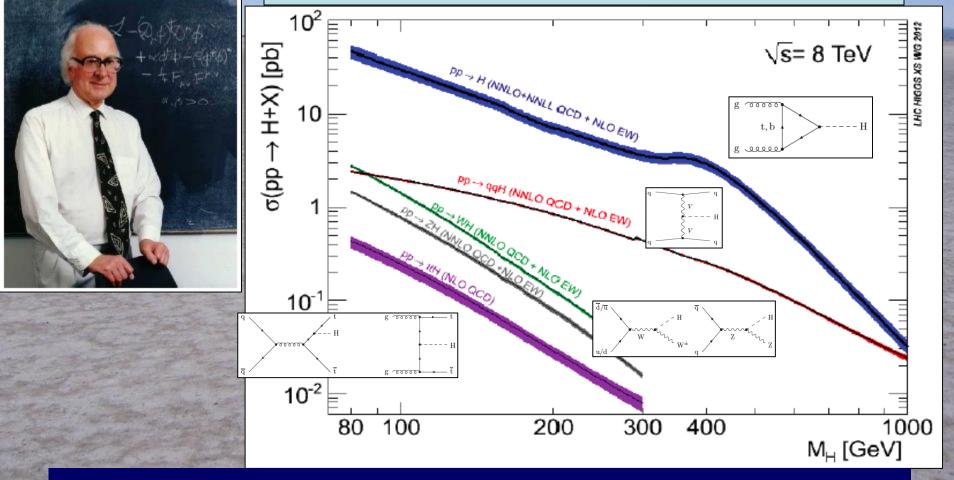
2011: Combining Information from Previous Direct Searches and Indirect Data



Gfitter collaboration

A la recherche du Higgs perdu ...

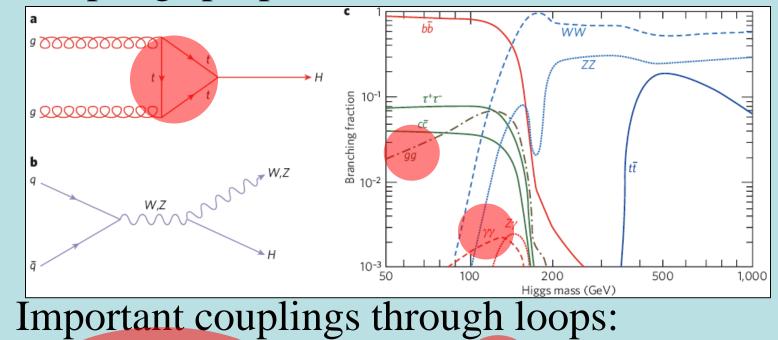
Higgs Production at the LHC



Many production modes measurable if $M_h \sim 125 \text{ GeV}$

Higgs Decay Branching Ratios

• Couplings proportional to masses (?)

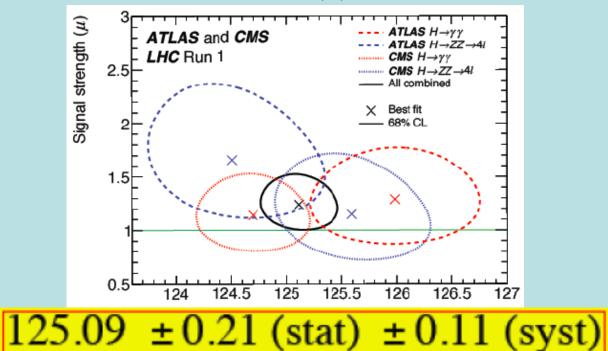


 $-gluon + gluon \rightarrow Higgs \rightarrow \gamma\gamma$

Many decay modes measurable if $M_h \sim 125 \text{ GeV}$

Higgs Mass Measurements

• ATLAS + CMS ZZ^* and $\gamma\gamma$ final states



• Statistical uncertainties dominate

• Allows precision tests

Crucial for stability of electroweak vacuum

The Particle Higgsaw Puzzle

Did LHC find the missing piece? Is it the right shape? Is it the right size?

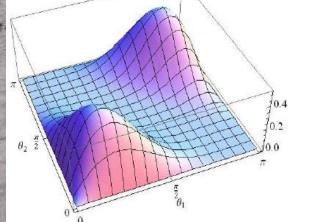
What is it?

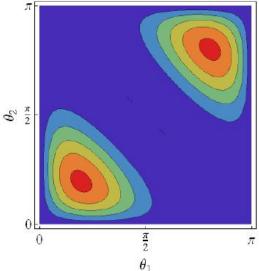
- Does it have spin 0 or 2?
- Is it scalar or pseudoscalar?
- Is it elementary or composite?
- Does it couple to particle masses?
- Quantum (loop) corrections?
- What are its self-couplings?

Does the 'Higgs' have Spin Zero?

- Polar angle distribution for $X_2 \rightarrow W^+W^-$
- Polar angle distribution for $X_0 \rightarrow W^+W^-$ (for $\phi = \pi$)

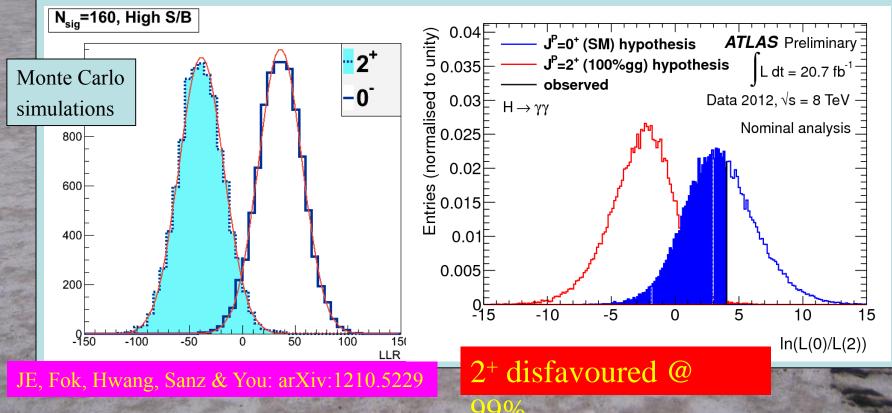
JE, Hwang: arXiv:1202.6660



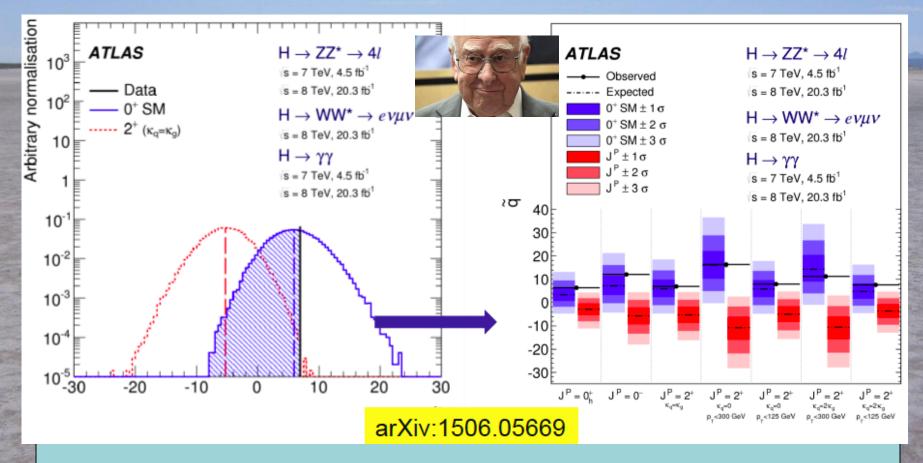


Does the 'Higgs' have Spin Two?

• Discriminate spin 2 vs spin 0 via angular distribution of decays into $\gamma\gamma$ JE & Hwang: arXiv:1202.6660



H Spin-Parity Tests: 0⁺ AOK



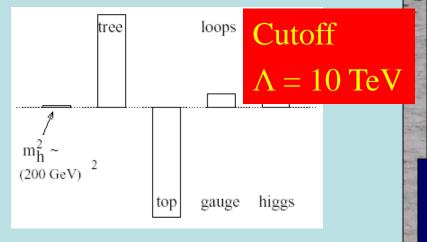
• Alternative spin-parities disfavoured > 99.9%

What is it?

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Elementary Higgs or Composite?

- Higgs field: $<0|H|0> \neq 0$
- Quantum loop problems



Cut-off $\Lambda \sim 1$ TeV with Supersymmetry?

- Fermion-antifermion condensate
- Just like QCD, BCS superconductivity
- Top-antitop condensate? needed m_t > 200 GeV
- New technicolour force?
- Heavy scalar resonance?
- Inconsistent with

precision electroweak data?

Higgs as a Pseudo-Goldstone Boson

UV completion ? sigma model cut-off

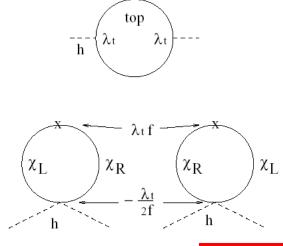
colored fermion related to top quark
 new gauge bosons related to SU(2)
 new scalars related to Higgs

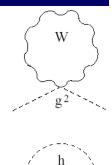
'Little Higgs' models(breakdown of larger symmetry)

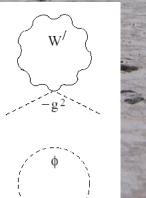
200 GeV- 1 or 2 Higgs doublets, possibly more scalars

Loop cancellation mechanism

Little Higgs

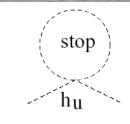


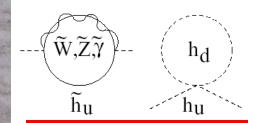




0 TeV

1 TeV





Supersymmetry

Phenomenological Framework

• Assume custodial symmetry:

 $SU(2) \times SU(2) \rightarrow SU(2)_V$ $(\rho \equiv M_W/M_Z \cos \theta_w \sim 1)$

• Parameterize gauge bosons by 2×2 matrix Σ :

$$\begin{split} \mathcal{L} &= \frac{v^2}{4} \text{Tr} D_{\mu} \Sigma^{\dagger} D^{\mu} \Sigma \left(1 + 2 \frac{a}{v} \frac{h}{v} + \frac{b}{v^2} \frac{h^2}{v^2} + ... \right) - m_i \bar{\psi}_L^i \Sigma \left(1 + \frac{c}{v} \frac{h}{v} + ... \right) \psi_R^i + \text{h.c.} \\ &+ \frac{1}{2} (\partial_{\mu} h)^2 + \frac{1}{2} m_h^2 h^2 + \frac{d_3}{6} \left(\frac{3m_h^2}{v} \right) h^3 + \frac{d_4}{24} \frac{1}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 + ... \quad , \end{split}$$

$$\Sigma = \exp\left(i\frac{\sigma^a\pi^a}{v}\right) \quad \mathcal{L}_{\Delta} = -\left[\frac{\alpha_s}{8\pi}b_sG_{a\mu\nu}G_a^{\mu\nu} + \frac{\alpha_{em}}{8\pi}b_{em}F_{\mu\nu}F^{\mu\nu}\right]\left(\frac{h}{V}\right)$$

• Coefficients a = c = 1 in Standard Model

Examples of Higgs as Pseudo-Goldstone Boson

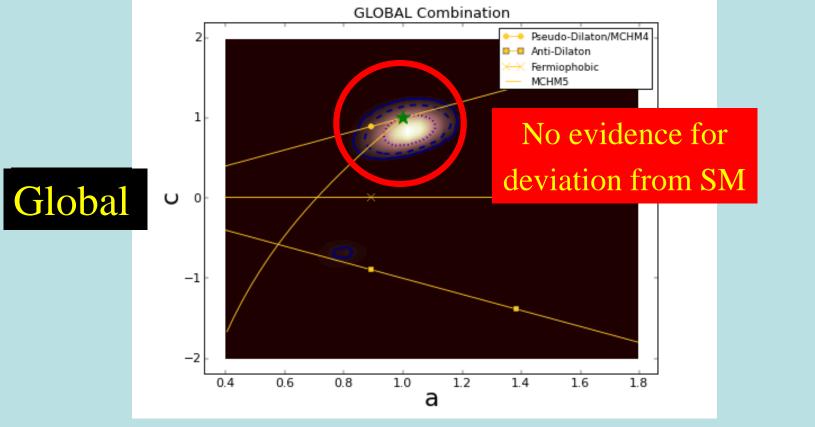
- Sample models:
- Dependences of couplings on model parameters:
- To be measured!
- Translation to experimental parameters:

$$a = \kappa_V, c = \kappa_F$$

Mode1	Symmetry Pattern		Goldstones		
SM	SO(4)/SO(3)		W_L, Z_L		
	SU(3)/SU(2)×U	(1)	1	W_L, Z_L, h	
MCHM	SO(5)/SO(4)×U(1)		W_L, Z_L, h		
NMCHM	SO(6)/SO(5)×U(1)		W_L, Z_L, h, a		
MCTHM	HM $SO(6)/SO(4) \times SO(2) \times U(1)$		$W_L, Z_L, h, H, H^{\pm}, a$		
Parameter	s SILH	MCHI	M4	MCHM5	
а	$1 - c_H \xi/2$	√1 -	ξ	$\sqrt{1-\xi}$	
Ь	$1 - 2c_H \xi$	1 – 2	ξ	1 — 2Ę	
<i>b</i> ₃	$-\frac{4}{3}\xi$	$-\frac{4}{3}\xi \sqrt{1}$	-ξ	$-\tfrac{4}{3}\xi\sqrt{1-\xi}$	
С	$1-(c_H/2+c_y)\xi$	$\sqrt{1}$	ξ	$\frac{1-2\xi}{\sqrt{1-\xi}}$	
<i>c</i> ₂	$-(c_H+3c_y)\xi/2$	-ξ/2	2	-2 <u></u>	
d_3	$1+(c_6-3c_H/2)\xi$	$\sqrt{1}$ –	ξ	$\frac{1-2\xi}{\sqrt{1-\xi}}$	
d_4	$1 + (6c_6 - 25c_H/3)\xi$	1 – 7 <u></u>	/3	$\frac{1-28\xi(1-\xi)/3}{1-\xi}$	

Global Analysis of Higgs-like Models

• Rescale couplings: to bosons by a, to fermions by c

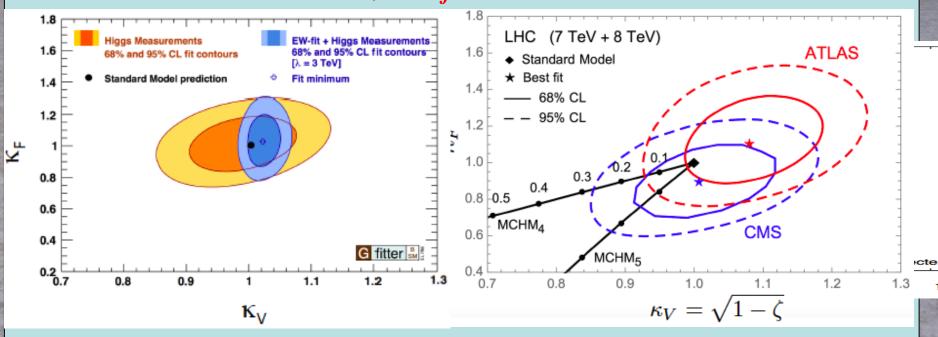


• Standard Model: a = c = 1

JE & Tevong You, arXiv:1303.3879

Global Analysis of Higgs-like Models

- Rescale couplings: to bosons by κ_V , to fermions by κ_f
- Standard Model: $\kappa_V = \kappa_f = 1$



Consistency between Higgs and EW measurements

• Must tune composite models to look like SM

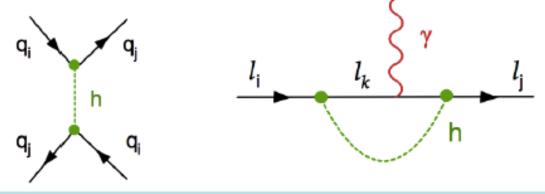
What is it?

- Does it have spin 0 or 2?
 - Spin 2 strongly disfavoured
- Is it scalar or pseudoscalar?
 - Pseudoscalar strongly disfavoured
- Is it elementary or composite?
 - No significant deviations from Standard Model
- Does it couple to particle masses?
- Quantum (loop) corrections?
- What are its self-couplings?

It Walks and Quacks like a Higgs • Do couplings scale ~ mass? With scale = v? ATLAS and CMS LHC Run 1 $\lambda_f = \sqrt{2} \left(\frac{m_f}{M}\right)^{1+\epsilon}, \ g_V = 2 \left(\frac{m_V^{2(1+\epsilon)}}{M^{1+2\epsilon}}\right)^{1+\epsilon}$ Ĕ |> 10⁻¹ 10⁻² ATLAS+CMS Global SM Higgs boson [M, ε] fit fit 68% CL 95% CL 10^{-4} **10**⁻¹ 10² 10 Par • Blue dashed line = Standard Model

Flavour-Changing Couplings?

• Upper limits from FCNC, EDMs, ...

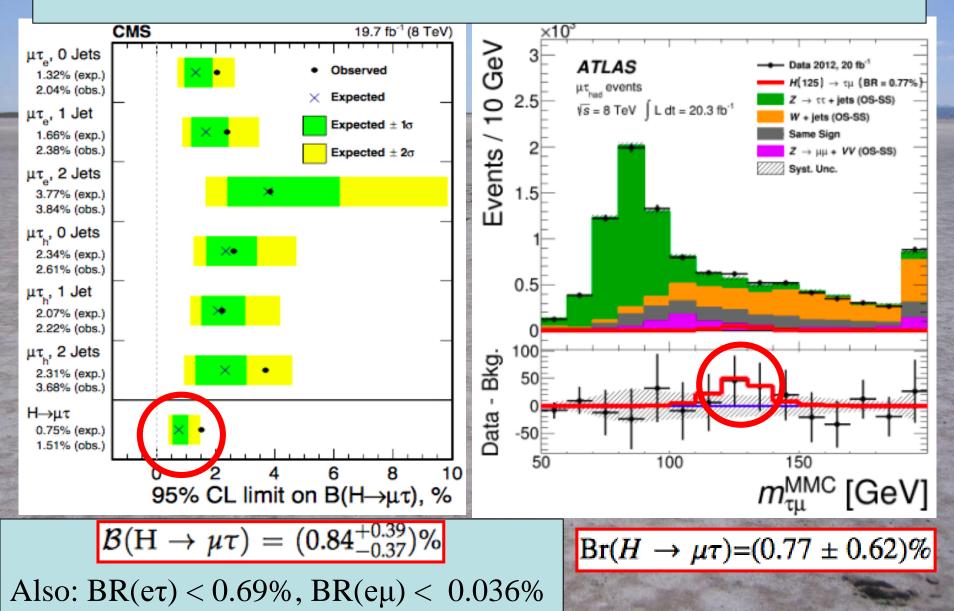


- Quark FCNC bounds exclude observability of quark-flavour-violating *h* decays
- Lepton-flavour-violating *h* decays could be large:
 BR(τμ) or BR(τe) could be O(10)%

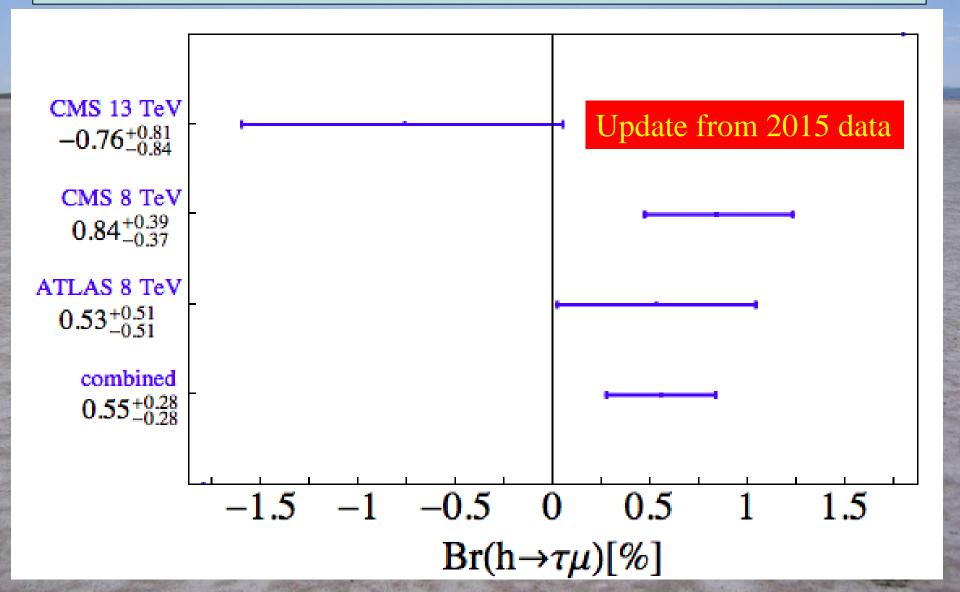
Blankenburg, JE, Isidori: arXiv:1202.5704

BR(μe) must be $< 2 \times 10^{-5}$

Flavour-Changing Higgs Coupling?



Flavour-Changing Higgs Coupling?



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- Is it elementary or composite?
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- Does it couple to particle masses?
 - Prima facie evidence that it does
- Quantum (loop) corrections?
- What are its self-couplings?

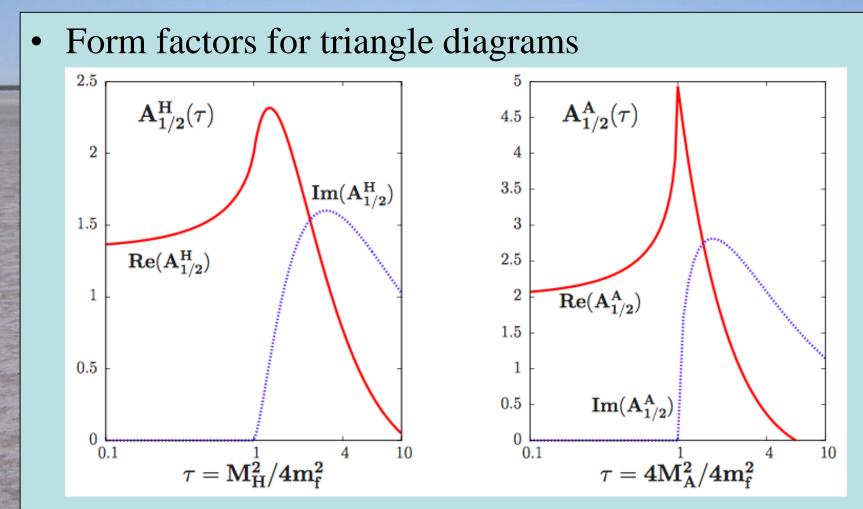
Triangle Diagrams for gg \rightarrow Spin-0 $\rightarrow \gamma \gamma$

- Effective vertices: $\mathcal{L}_{\text{eff}}^{H} = \frac{e}{v} c_{H\gamma\gamma} H F_{\mu\nu} F^{\mu\nu} + \frac{g_s}{v} c_{Hgg} H G_{\mu\nu} G^{\mu\nu}$ $\mathcal{L}_{\text{eff}}^{A} = \frac{e}{v} c_{A\gamma\gamma} A F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{g_s}{v} c_{Agg} A G_{\mu\nu} \tilde{G}^{\mu\nu} ,$
- Decay rates: $\Gamma(\Phi \to gg) = \frac{G_{\mu}\alpha_s^2 M_{\Phi}^3}{64\sqrt{2}\pi^3} \left| \sum_{Q} \hat{g}_{\Phi QQ} A_{1/2}^{\Phi}(\tau_Q) \right|^2$,

$$\begin{split} &\Gamma(\Phi \to \gamma \gamma) = \frac{G_{\mu} \alpha^2 M_{\Phi}^3}{128 \sqrt{2} \pi^3} \bigg| \sum_F \hat{g}_{\Phi FF} N_c e_F^2 A_{1/2}^{\Phi}(\tau_F) \bigg|^2 \\ &\text{Vertex} \\ &\text{form factors:} \quad A_{1/2}^H(\tau) = 2 \left[\tau + (\tau - 1) f(\tau) \right] \tau^{-2}, \quad A_{1/2}^A(\tau) = 2 \tau^{-1} f(\tau) \\ &f(\tau) = \begin{cases} \arccos^2 \sqrt{\tau} & \text{for } \tau \leq 1, \\ -\frac{1}{4} \left[\log \frac{1 + \sqrt{1 - \tau^{-1}}}{1 - \sqrt{1 - \tau^{-1}}} - i\pi \right]^2 & \text{for } \tau > 1. \end{cases} \end{split}$$

Vanish for fermion mass << spin-0 mass

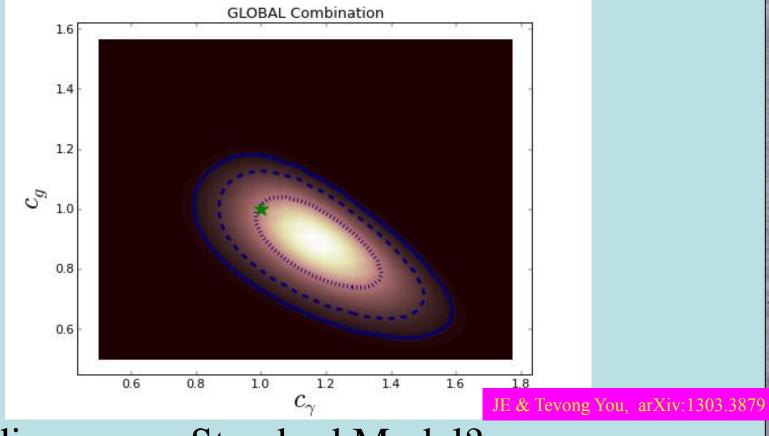
Triangle Diagrams for gg \rightarrow Spin-0 $\rightarrow \gamma \gamma$



• Vanish for fermion mass << spin-0 mass

Loop Corrections ?

• Combination of data on γγ, gluon-gluon couplings



• Loop diagrams ~ Standard Model?

What is it?

- Does it have spin 0 or 2?
 Spin 2 strongly disfavoured
- Is it scalar or pseudoscalar?
 - Pseudoscalar disfavoured

Н0

J = 0

Mass $m = 125.09 \pm 0.24$ GeV

H^0 Signal Strengths in Different Channels

See Listings for the latest unpublished results.

Combined Final States = 1.17 ± 0.17 (S = 1.2) $WW^* = 0.81 \pm 0.16$ $ZZ^* = 1.15^{+0.27}_{-0.23}$ (S = 1.2) $\gamma\gamma = 1.17^{+0.19}_{-0.17}$ $b\overline{b} = 0.85 \pm 0.29$ $\mu^+\mu^- < 7.0$, CL = 95% $\tau^+\tau^- = 0.79 \pm 0.26$ $Z\gamma < 9.5$, CL = 95% $t\overline{t}H^0$ Production = $2.5^{+0.9}_{-0.8}$

- Is it elementary or composite?
 - No significant deviations from Standard Model
- Does it couple to particle masses?
 - Prima facie evidence that it does
- Quantum (loop) corrections?
 - γγ, gg couplings ~ Standard Model
- What are its self-couplings?

Dixit Swedish Academy

Today we believe that "Beyond any reasonable doubt, it is a Higgs boson." [1] http://www.nobelprize.org/nobel_prizes/physics/laureates/2013/a dvanced-physicsprize2013.pdf

Assuming H(125) is SM-like: Model-independent search for new physics

Standard Model Effective Field Theory

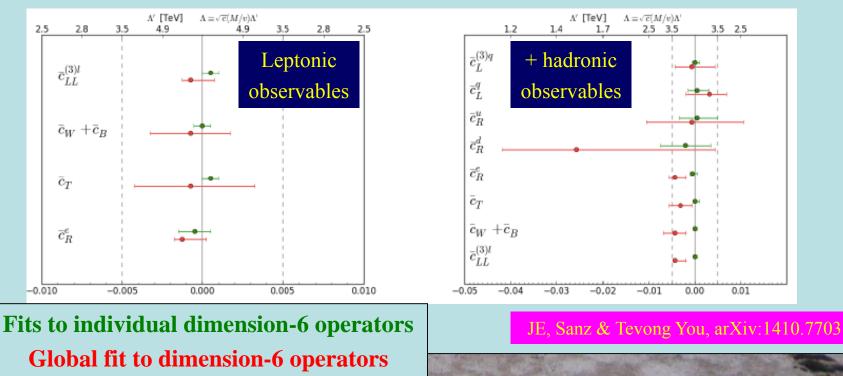
- Higher-dimensional operators as relics of higherenergy physics, e.g., dimension 6: $\mathcal{L}_{eff} = \sum \frac{f_n}{\Lambda^2} \mathcal{O}_n$
- Operators constrained by $SU(2) \times U(1)$ symmetry:
 - $\mathcal{L} \supset \frac{\bar{c}_{H}}{2v^{2}} \partial^{\mu} [\Phi^{\dagger}\Phi] \partial_{\mu} [\Phi^{\dagger}\Phi] + \frac{g'^{2} \bar{c}_{\gamma}}{m_{W}^{2}} \Phi^{\dagger}\Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_{s}^{2} \bar{c}_{g}}{m_{W}^{2}} \Phi^{\dagger}\Phi G_{\mu\nu}^{a} G_{\mu\nu}^{\mu\nu}$ $+ \frac{2ig \bar{c}_{HW}}{m_{W}^{2}} [D^{\mu}\Phi^{\dagger}T_{2k}D^{\nu}\Phi] W_{\mu\nu}^{k} + \frac{ig' \bar{c}_{HB}}{m_{W}^{2}} [D^{\mu}\Phi^{\dagger}D^{\nu}\Phi] B_{\mu\nu}$ $+ \frac{ig \bar{c}_{W}}{m_{W}^{2}} [\Phi^{\dagger}T_{2k}\overleftrightarrow{D}^{\mu}\Phi] D^{\nu} W_{\mu\nu}^{k} + \frac{ig' \bar{c}_{B}}{2m_{W}^{2}} [\Phi^{\dagger}\overleftrightarrow{D}^{\mu}\Phi] \partial^{\nu} B_{\mu\nu}$ $+ \frac{\bar{c}_{t}}{v^{2}} y_{t}\Phi^{\dagger}\Phi \Phi^{\dagger} \cdot \bar{Q}_{L}t_{R} + \frac{\bar{c}_{b}}{v^{2}} y_{b}\Phi^{\dagger}\Phi \Phi \cdot \bar{Q}_{L}b_{R} + \frac{\bar{c}_{\tau}}{v^{2}} y_{\tau} \Phi^{\dagger}\Phi \Phi \cdot \bar{L}_{L}\tau_{R}$
- Constrain with precision EW, Higgs data, TGCs ...

Electroweak Precision Data

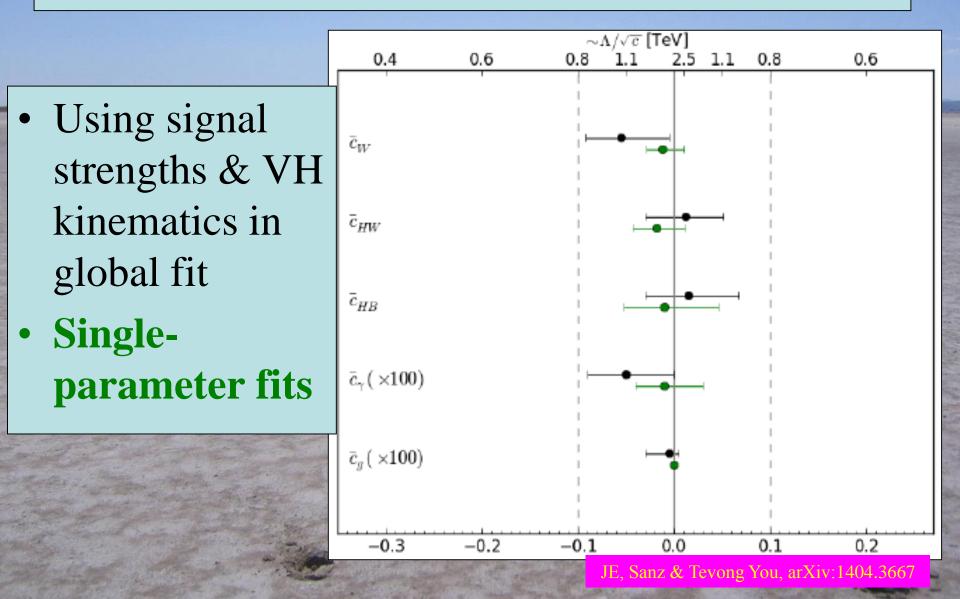
• Operators affecting oblique parameters

$$\mathcal{L}_{\text{dim-6}} \subset \frac{\bar{c}_{WB}}{m_W^2} \mathcal{O}_{WB} + \frac{\bar{c}_W}{m_W^2} \mathcal{O}_W + \frac{\bar{c}_B}{m_W^2} \mathcal{O}_B + \frac{\bar{c}_T}{v^2} \mathcal{O}_T + \frac{\bar{c}_{2W}}{m_W^2} \mathcal{O}_{2W} + \frac{\bar{c}_{2B}}{m_W^2} \mathcal{O}_{2B}$$

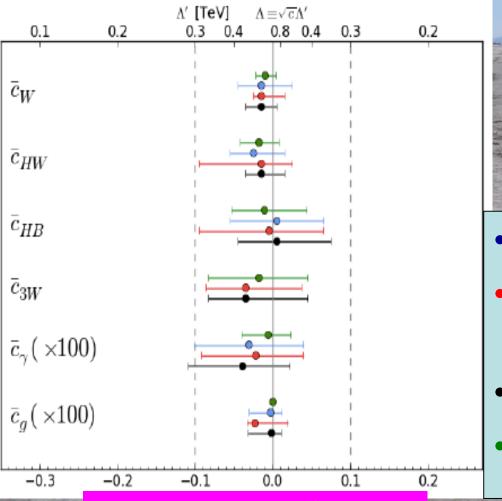
- Also other electroweak tests
- Constraints from LEP et al. data



Fits including Higgs Production

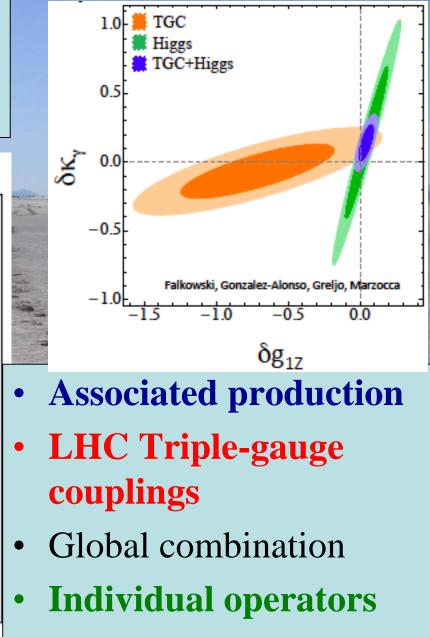


Global Fits including LHC TGCs

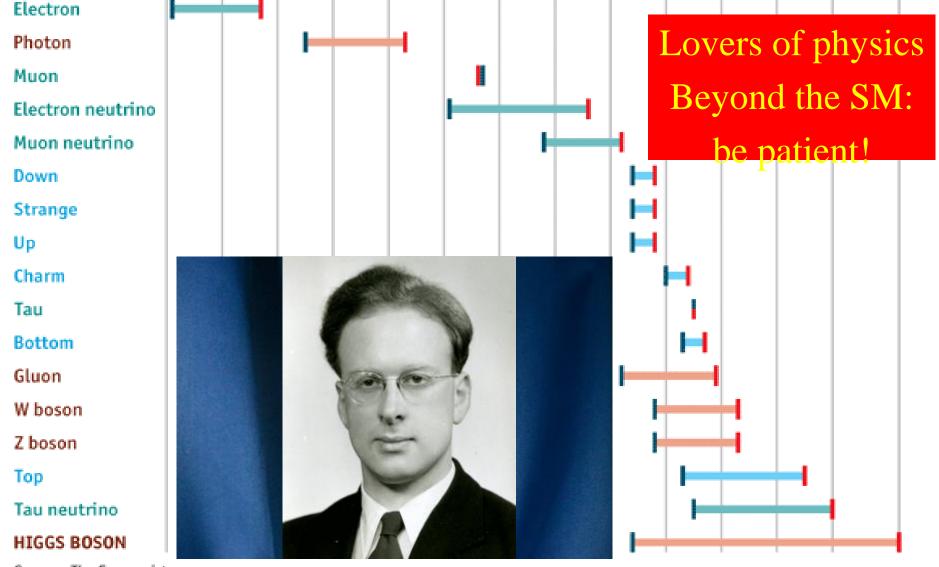


Sanz & Tevong You, arXiv:1410.

JE.



Standard Model Particles: Years from Proposal to Discovery



Source: The Economist



« Empty » space is unsta SUSY

Is Nat EMALIA

- Dark matter
- Origin of matter
- Masses of neutrinos
- Hierarchy problem
- Inflation
- Quantum gravity

SUSY SUSY

SUSY SUSY SUSY

The Standard Model

What lies beyond the Standard Model?

Supersymmetry

Stabilize electroweak vacuum

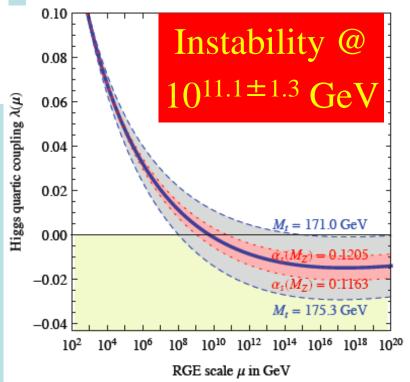
New motivations From LHC Run 1

- Successful prediction for Higgs mass
 Should be < 130 GeV in simple models
- Successful predictions for couplings

 Should be within few % of SM values
 Should be within few % of SM values
- Naturalness, GUTs, string, ..., dark matter

Theoretical Constraints on Higgs Mass

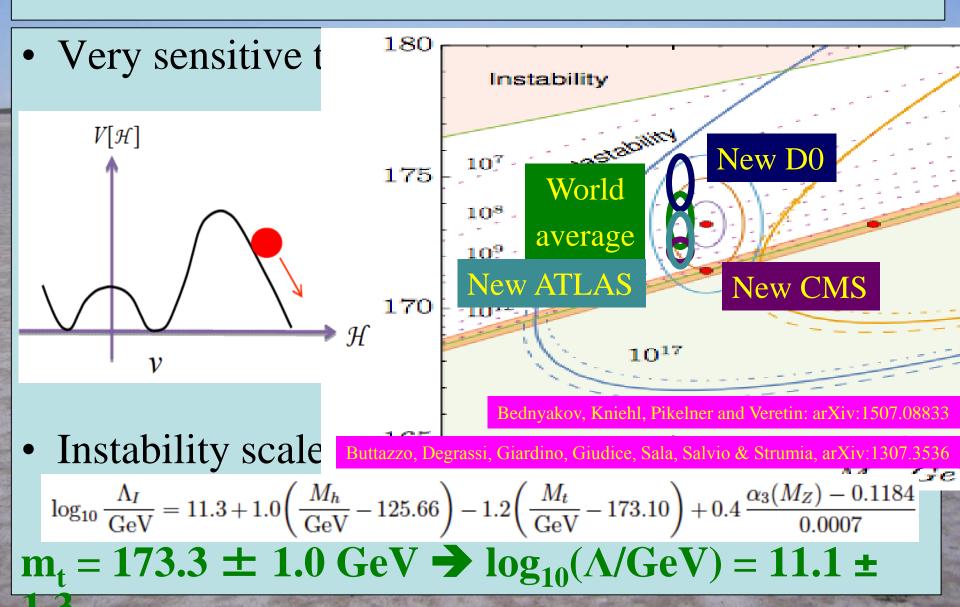
- Large $M_h \rightarrow$ large self-coupling \rightarrow blow up at $\lambda(Q) = \lambda(v) - \frac{3m_t^4}{2\pi^2 v^4} \log \frac{Q}{v} \quad \int_{0.08}^{0.10} \left[\int_{0.08} \text{Instability @} \right]$
- Small: renormalization due to t quark drives quartic coupling < 0 at some scale Λ
 → vacuum unstable



• Vacuum could be stabilized by **Supersymmetry**

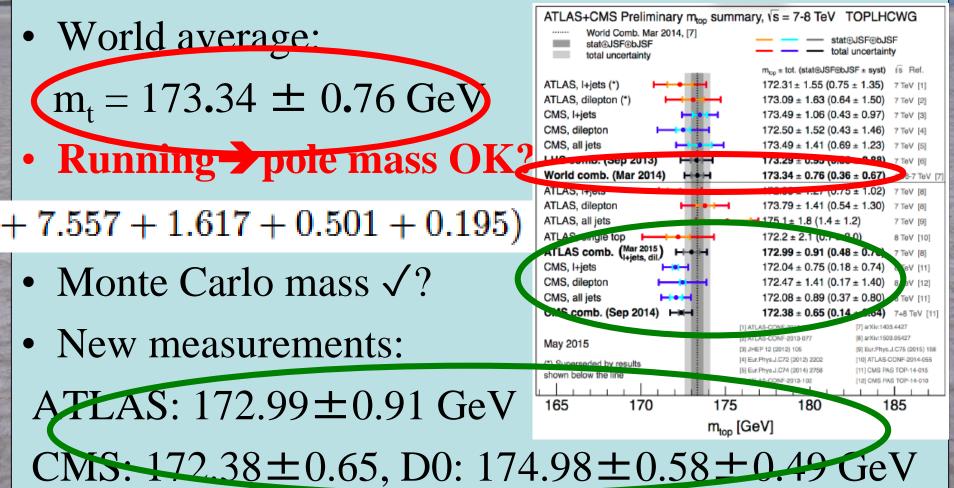
Degrassi, Di Vita, Elias-Miro, Giudice, Isodori & Strumia, arXiv:1205.6497

Vacuum Instability in the Standard Model



Hard QCD: the Top Mass

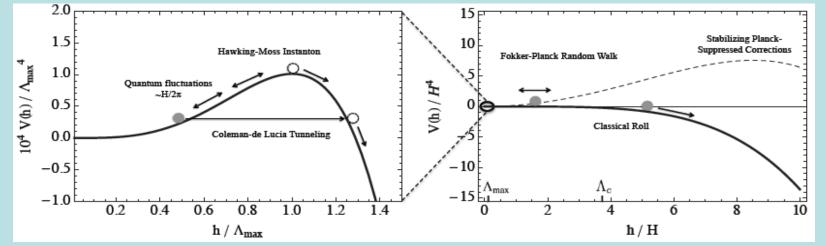
• Basic parameter of SM; stability of EW vacuum?



Instability during Inflation?

Hook, Kearns, Shakya & Zurek: arXiv:1404.5953

Do inflation fluctuations drive us over the hill?



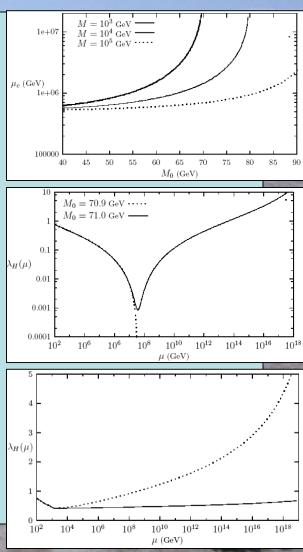
- Then Fokker-Planck evolution
- Do AdS regions eat us?
 - Disaster if so
 - If not, OK if more inflation

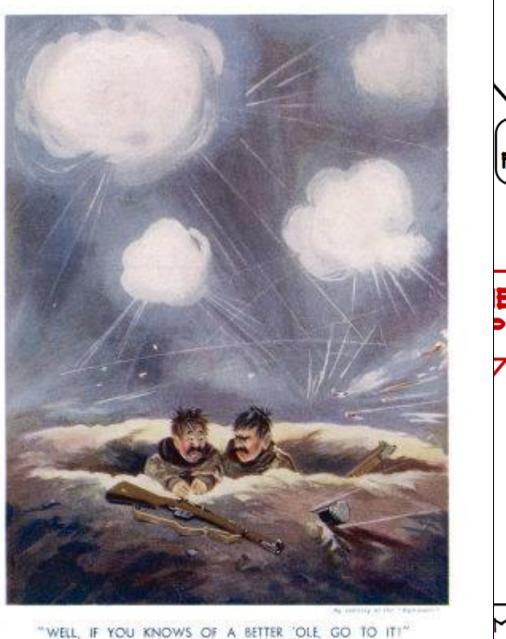
Stabilize vacuum with some physics beyond the SM?

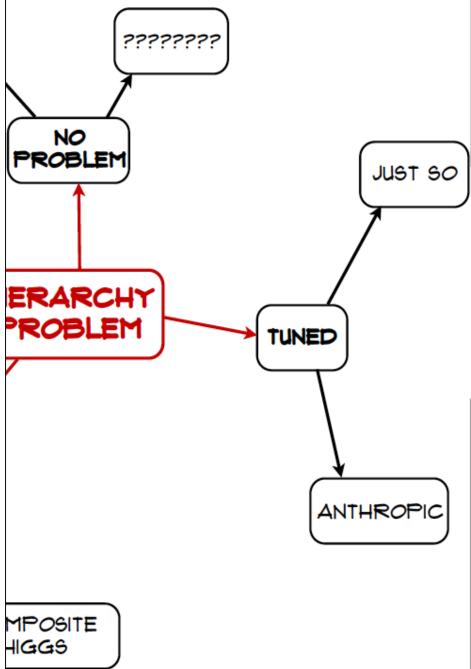
How to Stabilize a Light Higgs Boson?

- Top quark destabilizes potential: introduce stop-like scalar: $\mathcal{L} \supset M^2 |\phi|^2 + \frac{M_0}{v^2} |H|^2 |\phi|^2$
- Can delay collapse of potential:
- But new coupling must be fine-tuned to avoid blow-up:
- Stabilize with new fermions:
 just like Higgsinos
- Very like Supersymmetry!

D Ross







Loop Corrections to Higgs Mass²

• Consider generic fermion and boson loops:

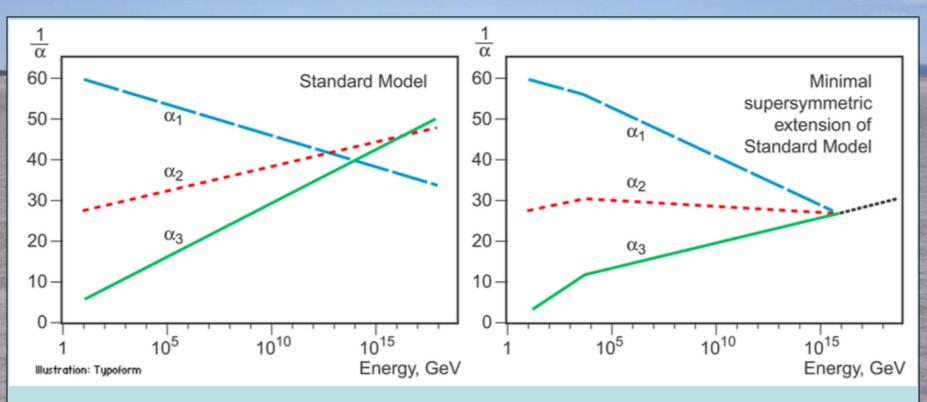


• Each is quadratically divergent: $\int^{\Lambda} d^4k/k^2$

$$\Delta m_H^2 = -\frac{y_f^2}{16\pi^2} [2\Lambda^2 + 6m_f^2 \ln(\Lambda/m_f) + ...]$$
$$\Delta m_H^2 = \frac{\lambda_S}{16\pi^2} [\Lambda^2 - 2m_S^2 \ln(\Lambda/m_S) + ...]$$

• Leading divergence cancelled if $\lambda_S = y_f^2 \ge 2$ Supersymmetry!

Unification of Gauge Couplings



- Impressive!
- Over-ambitious? Hubristic?

Minimal Supersymmetric Extension of Standard Model (MSSM)

• Double up the known particles:

$$\begin{pmatrix} \frac{1}{2} \\ 0 \end{pmatrix} e.g., \ \begin{pmatrix} \ell (lepton) \\ \tilde{\ell} (slepton) \end{pmatrix} or \begin{pmatrix} q (quark) \\ \tilde{q} (squark) \end{pmatrix} \\ \begin{pmatrix} 1 \\ \frac{1}{2} \end{pmatrix} e.g., \ \begin{pmatrix} \gamma (photon) \\ \tilde{\gamma} (photino) \end{pmatrix} or \begin{pmatrix} g (gluon) \\ \tilde{g} (gluino) \end{pmatrix}$$

- Two Higgs doublets
 - 5 physical Higgs bosons:
 - 3 neutral, 2 charged
- Lightest neutral supersymmetric Higgs looks like the single Higgs in the Standard Model

Higgs Bosons in Supersymmetry

- Need 2 complex Higgs doublets (cancel anomalies, form of SUSY couplings)
- 8 3 = 5 physical Higgs bosons
 Scalars h, H; pseudoscalar A; charged H[±]
- Lightest Higgs < MZ at tree level:

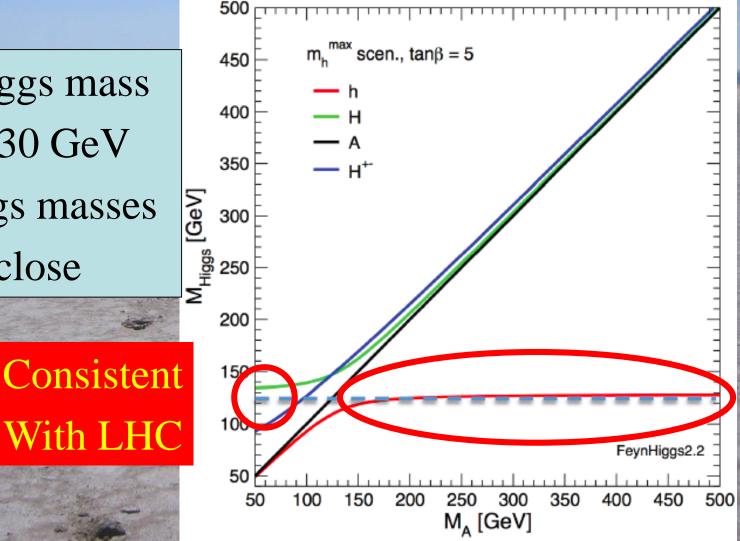
 $M_{\rm H,h}^2 = \frac{1}{2} \left[M_{\rm A}^2 + M_{\rm Z}^2 \pm \sqrt{(M_{\rm A}^2 + M_{\rm Z}^2)^2 - 4M_{\rm Z}^2 M_{\rm A}^2 \cos^2 2\beta} \right]$

• Important radiative corrections to mass:

$$G_{\mu} m_{\mathrm{t}}^{4} \ln \left(\frac{m_{\tilde{\mathrm{t}}_{1}} m_{\tilde{\mathrm{t}}_{2}}}{m_{\mathrm{t}}^{2}} \right) \Delta M_{\mathrm{H}} |_{\mathrm{TH}} \sim 1.5 \mathrm{~GeV}$$

MSSM Higgs Masses & Couplings

Lightest Higgs mass up to ~ 130 GeV Heavy Higgs masses 750? quite close



Lightest Supersymmetric Particle

• Stable in many models because of conservation of R parity:

 $R = (-1)^{2S - L + 3B}$

where S = spin, L = lepton #, B = baryon #

 Particles have R = +1, sparticles R = -1: Sparticles produced in pairs Heavier sparticles → lighter sparticles
 Lightest supersymmetric particle (LSP) stable

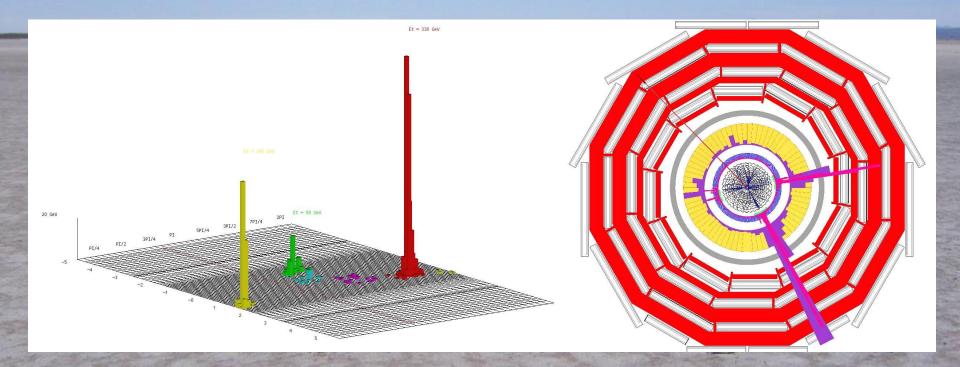
Lightest Sparticle as Dark Matter?

- No strong or electromagnetic interactions Otherwise would bind to matter Detectable as anomalous heavy nucleus
- Possible weakly-interacting scandidates
 Sneutrino

(Excluded by LEP, direct searches) **Lightest neutralino** χ (partner of Z, H, γ) **Gravitino**

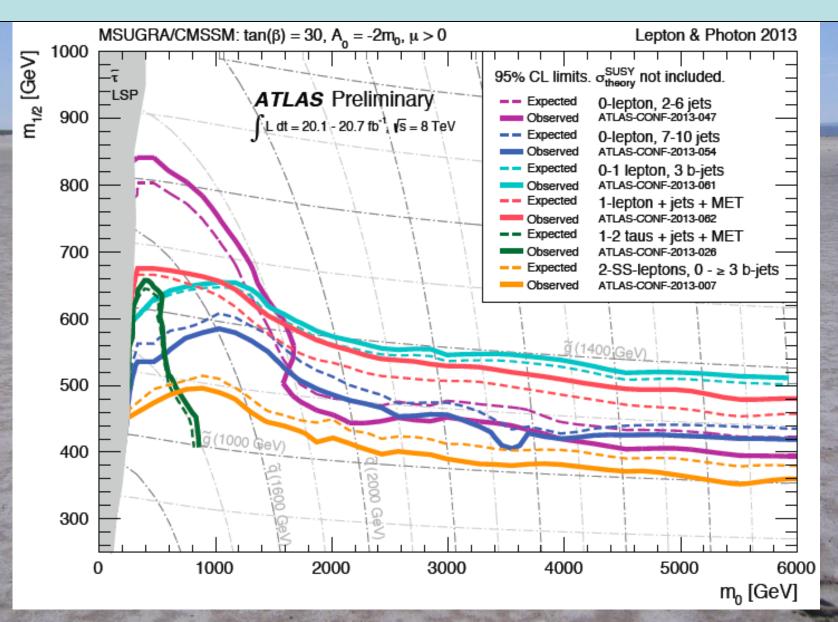
(nightmare for detection)

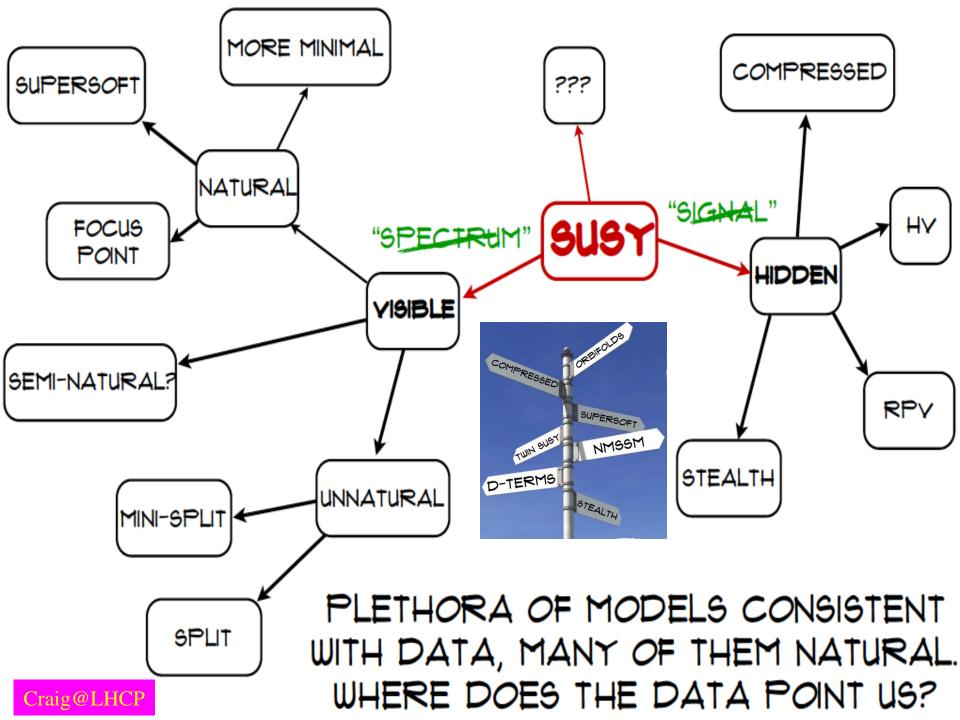
Classic Dark Matter Signature



Missing transverse energy carried away by dark matter particles

Searches with ~ 20/fb @ 8 TeV





(Non-)Universal Scalar Masses?

• Different sfermions with same quantum #s? e.g., d, s squarks? disfavoured by upper limits on flavourchanging neutral interactions • Squarks with different #s, squarks and sleptons? disfavoured in various GUT models e.g., $d_R = e_L$, $d_L = u_L = u_R = e_R$ in SU(5), all in SO(10) • Non-universal susy-breaking masses for Higgses? No reason why not! NUHM

Sample Supersymmetric Models

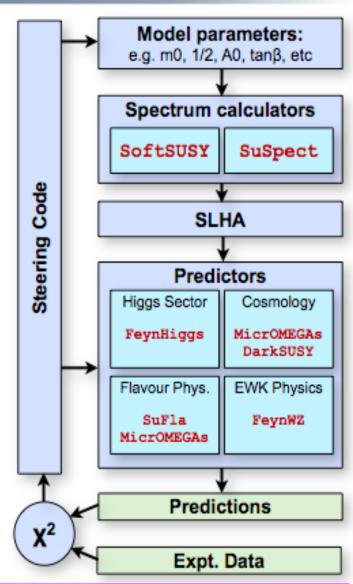
- Universal soft supersymmetry breaking at input GUT scale?
 - For gauginos and all scalars: CMSSM
 - Non-universal Higgs masses: NUHM1,2
- Strong pressure from LHC (p ~ 0.1)
- Treat soft supersymmetry-breaking masses as phenomenological inputs at EW scale
 - pMSSMn (n parameters)
 - With universality motivated by upper limits on flavour-changing neutral interactions: pMSSM10
- Less strongly constrained by LHC (p ~ 0.3)

MasterCode



Combines diverse set of tools

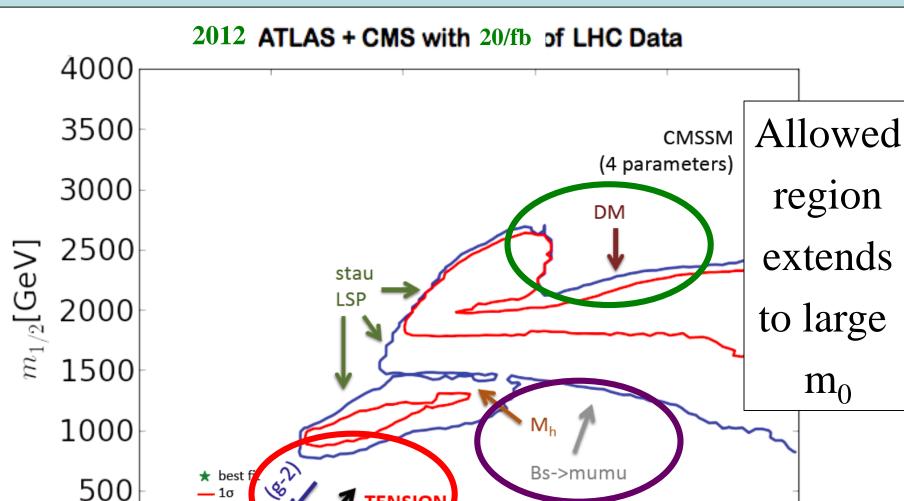
- different codes : all state-of-the-art
 - Electroweak Precision (FeynWZ)
 - Flavour (SuFla, micrOMEGAs)
 - Cold Dark Matter (DarkSUSY, micrOMEGAs)
 - Other low energy (FeynHiggs)
 - Higgs (FeynHiggs)
- different precisions (one-loop, two-loop, etc)
- different languages (Fortran, C++, English, German, Italian, etc)
- different people (theorists, experimentalists)
- Compatibility is crucial! Ensured by
 - close collaboration of tools authors
 - standard interfaces



E. Bagnaschi, M. Borsato, O. Buchmueller, R. Cavanaugh, V. Chobanova, M. Citron, J. Costa, A. De Roeck, M.J. Dolan, J.E., H. Flacher, S. Heinemeyer, G. Isidori, M. Lucio, D. Martinez Santos, K.A. Olive, A. Richards, K. Sakurai, G. Weiglein

			Observable	Source	Constraint
			m_t [GeV]	Th./Ex. [39]	173.2 ± 0.90
	Data		$\Delta \alpha_{\rm had}^{(5)}(m_{\rm Z})$	[39]	173.2 ± 0.90 0.02749 ± 0.00010
			$M_{\rm had}(m_Z)$ $M_Z [{\rm GeV}]$	[40]	91.1875 ± 0.0021
			Γ_Z [GeV]	[24] / [40]	$2.4952 \pm 0.0023 \pm 0.001_{SUSY}$
			$\sigma_{\rm had}^0$ [nb]	[24] / [40]	$\frac{2.4352 \pm 0.0023 \pm 0.0018089}{41.540 \pm 0.037}$
				[24] / [40]	20.767 ± 0.025
			$A_{\rm fb}(\ell)$	[24] / [40]	0.01714 ± 0.00095
			$A_{\ell}(P_{\tau})$	[24] / [40]	0.1465 ± 0.0032
Electroweak precision observables			Rb	[24] / [40]	0.21629 ± 0.00066
			R _c	[24] / [40]	0.1721 ± 0.0030
			$A_{\rm fb}(b)$	[24] / [40]	0.0992 ± 0.0016
			$A_{\rm fb}(c)$	[24] / [40]	0.0707 ± 0.0035
UDSET VADIES			A_b	[24] / [40]	0.923 ± 0.020
			A_c	[24] / [40]	0.670 ± 0.027
Flavour physics observables			$A_{\ell}(\text{SLD})$	[24] / [40]	0.1513 ± 0.0021
			$\sin^2 \theta_{\rm w}^{\ell}(Q_{\rm fb})$	[24] / [40]	0.2324 ± 0.0012
			M_W [GeV]	[24] / [40]	$80.399 \pm 0.023 \pm 0.010_{\rm SUSY}$
			$BR_{b \rightarrow s\gamma}^{EXP}/BR_{b \rightarrow s\gamma}^{SM}$	[41] / [42]	$1.117\pm0.076_{\rm EXP}$
					$\pm 0.082_{\rm SM}\pm 0.050_{\rm SUSY}$
	Deviation from S	ard Model	[27] / [37]	$(< 1.08 \pm 0.02_{\rm SUSY}) \times 10^{-8}$	
	-2	land		[27] / [42]	$1.43\pm0.43_{\rm EXP+TH}$
$\circ g_{\mu} - 2$ Supersymmetry at low \circ			~~1~~~~	[27] / [42]	$< (4.6 \pm 0.01_{\rm SUSY}) \times 10^{-9}$
Supersymmetry at low			scale, or		0.99 ± 0.32
	and model		$BR_{K \to \mu\nu}/BR_{K \to \mu\nu}$	[27] / [44]	$1.008 \pm 0.014_{\rm EXP+TH}$
Higgs mass		1000 C	$BR_{K \to \pi \nu \bar{\nu}}^{EXP} / BR_{K \to \pi \nu \bar{\nu}}^{SM}$	[45]/ [46]	< 4.5
			$\Delta M_{B_s}^{\text{EXP}} / \Delta M_{B_s}^{\text{SM}}$	[45] / [47,48]	$0.97 \pm 0.01_{\rm EXP} \pm 0.27_{\rm SM}$
Dark matter			$\frac{\left(\Delta M_{B_g}^{EXP} / \Delta M_{B_g}^{SM}\right)}{\left(\Delta M_{B_g}^{EXP} / \Delta M_{B_g}^{SM}\right)}$	[27] / [42, 47, 48]	$1.00 \pm 0.01_{\rm EXP} \pm 0.13_{\rm SM}$
			$\Delta \epsilon_K^{\text{EXP}} / \Delta \epsilon_K^{\text{SM}}$	[45] / [47 49]	$1.08 \pm 0.14_{\mathrm{EXP+TH}}$
		a Exp a om	[49] / [38,50]	$(30.2 \pm 8.8 \pm 2.0 \text{sus}) \times 10^{-10}$	
• LHC			M = 125.0		
			$H^{-120.0}$	<u> - 0.21</u> -	$= 0.11 \text{ GeV}_{\pm 0.019}^{.5\text{susy}}$
			σ_p	[23]	$(m_{\gamma}, \frac{SL}{p})$ plane
			jets $+ \not\!\!\!E_T$	[16, 18]	$(m_0, m_{1/2})$ plane
MasterCo	ode: O.Buchmueller, JE et al.		$H/A, H^{\pm}$	[19]	$(M_A, \tan\beta)$ plane
			And a second		

Fit to Constrained MSSM (CMSSM)



p-value of simple models ~ 10% (also SM)

TENSION

1000

Bs->mumu

2000

Buchmueller, JE et al: arXiv:1312.5250

4000

 m_0

3000

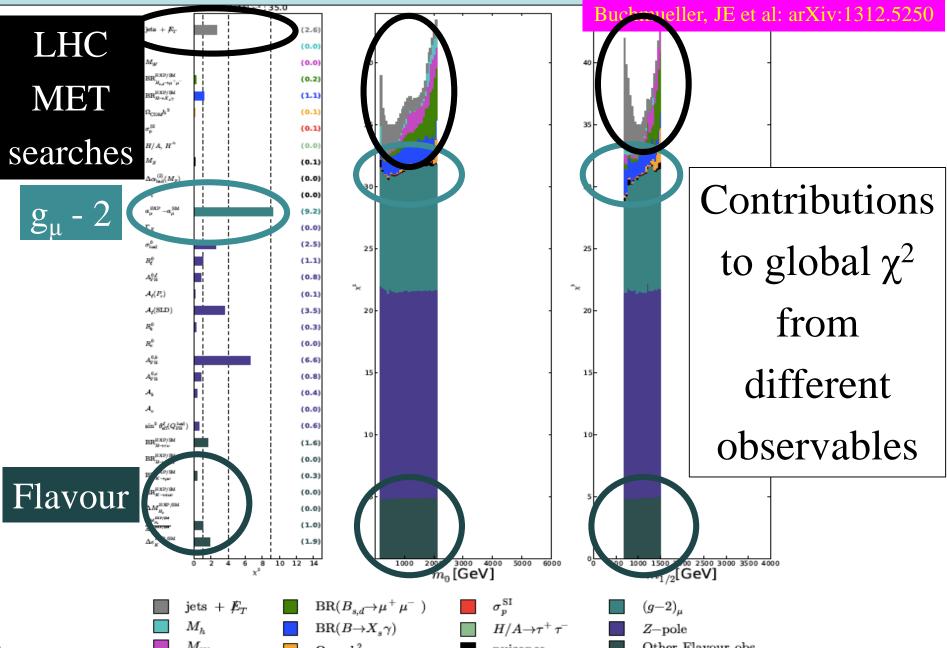
★ best f

- 2σ

500

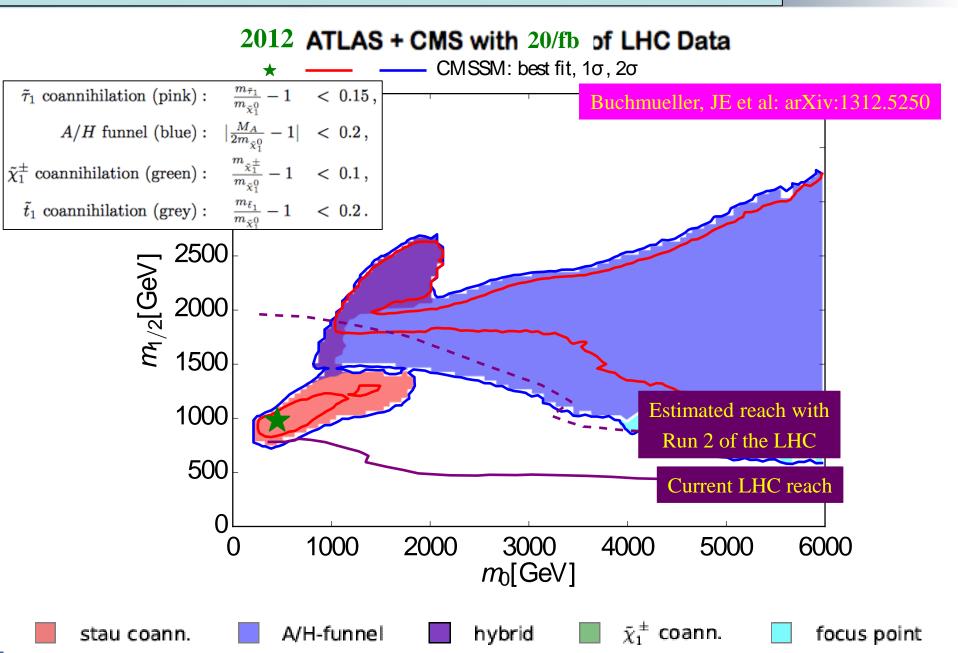
<u>1000</u>

Constrained MSSM (CMSSM)

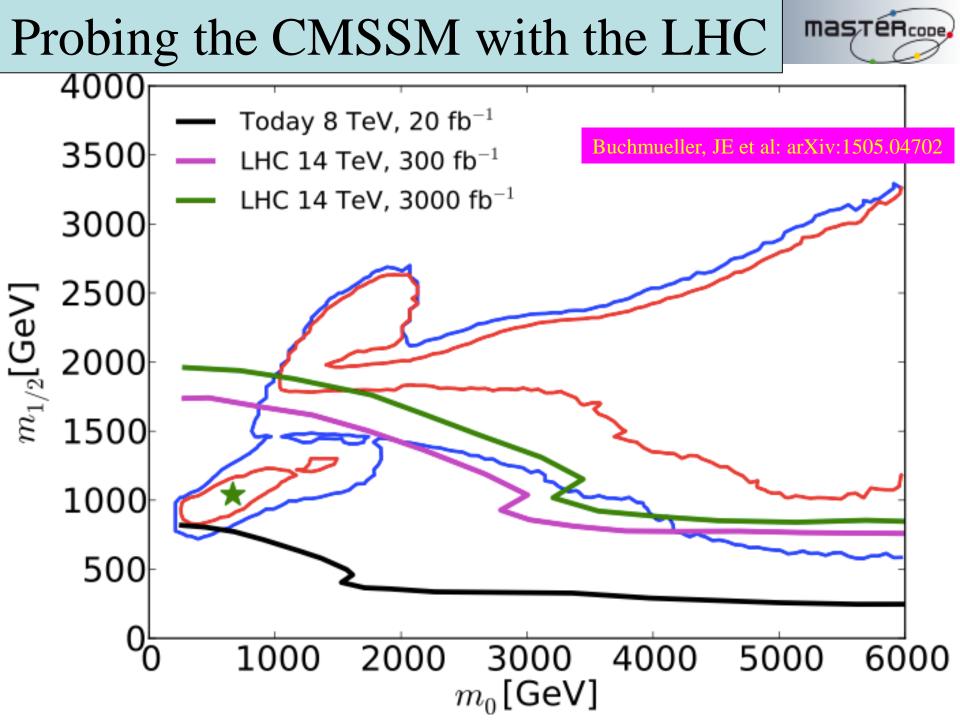


Mas/TeRcope

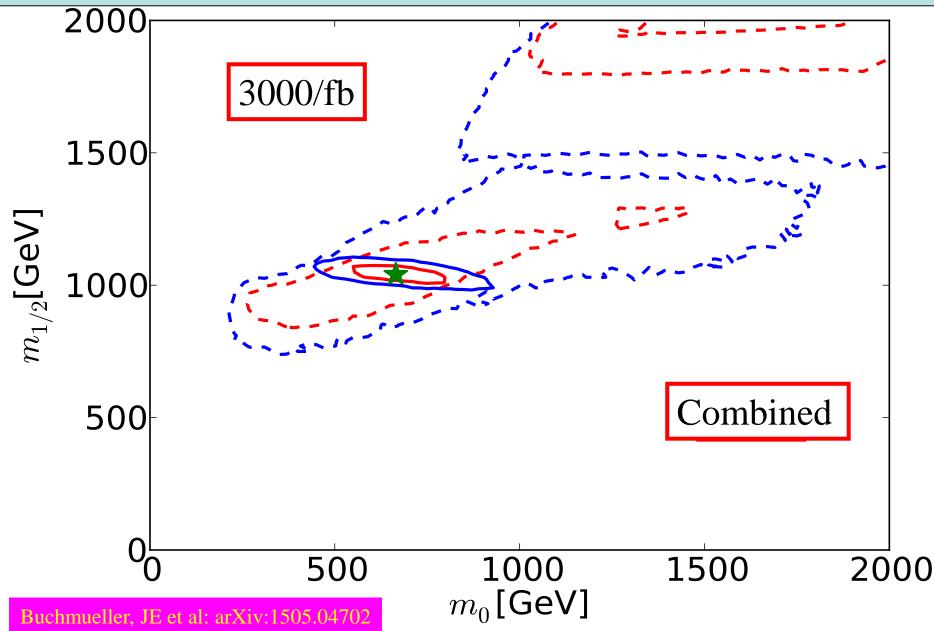
Dark Matter Density Mechanisms



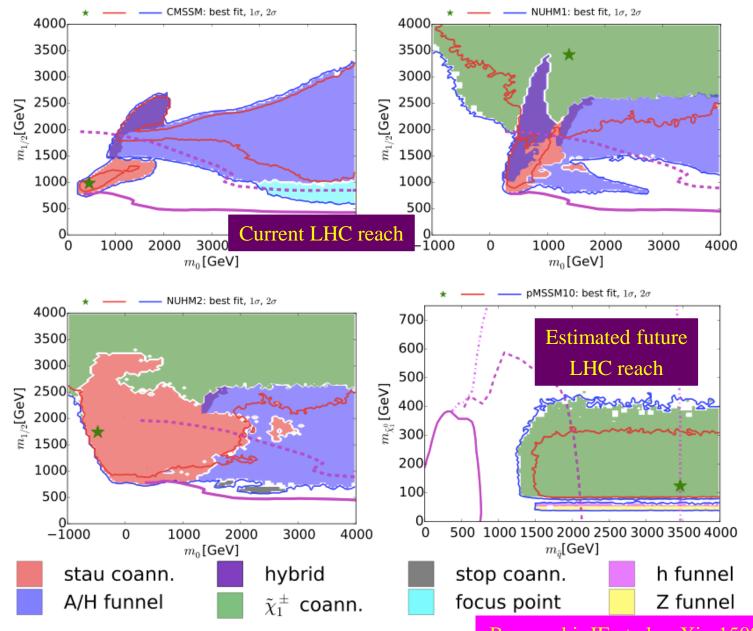
mas Tencore



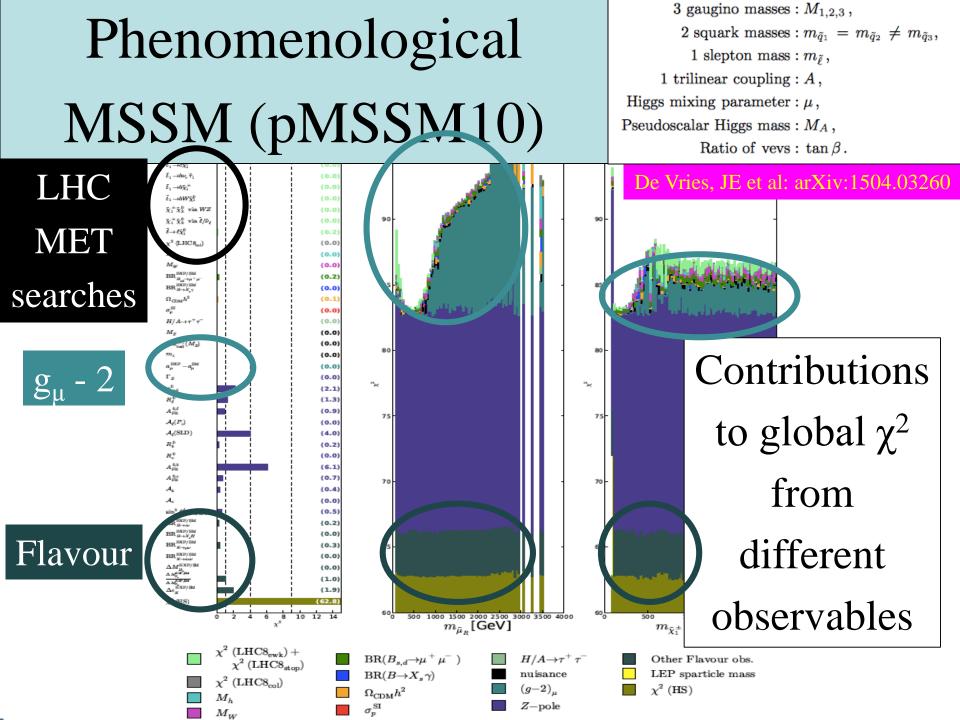
Measuring the CMSSM with the LHC



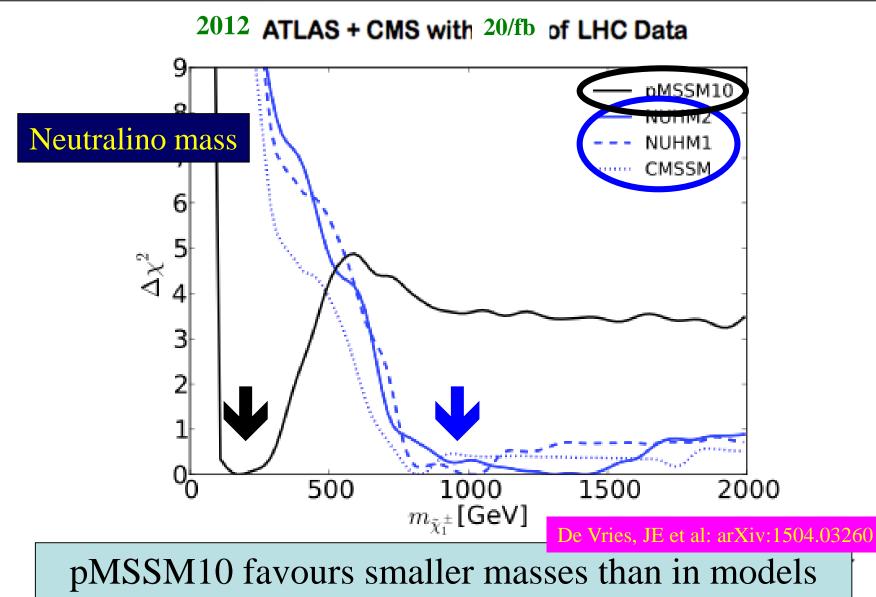
Dark Matter in CMSSM, NUHM1/2, pMSSM10



Bagnaschi, JE et al: arXiv:1508.01173



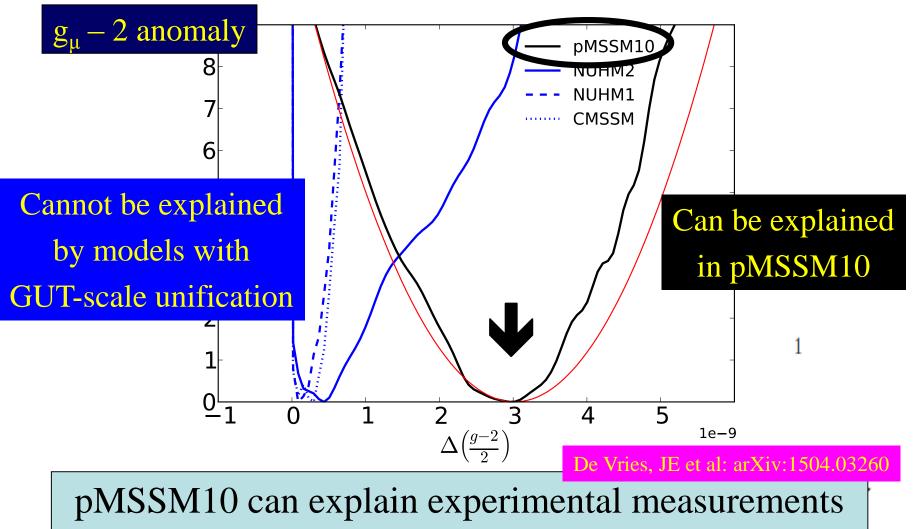
Possible Dark Matter Particle Mass



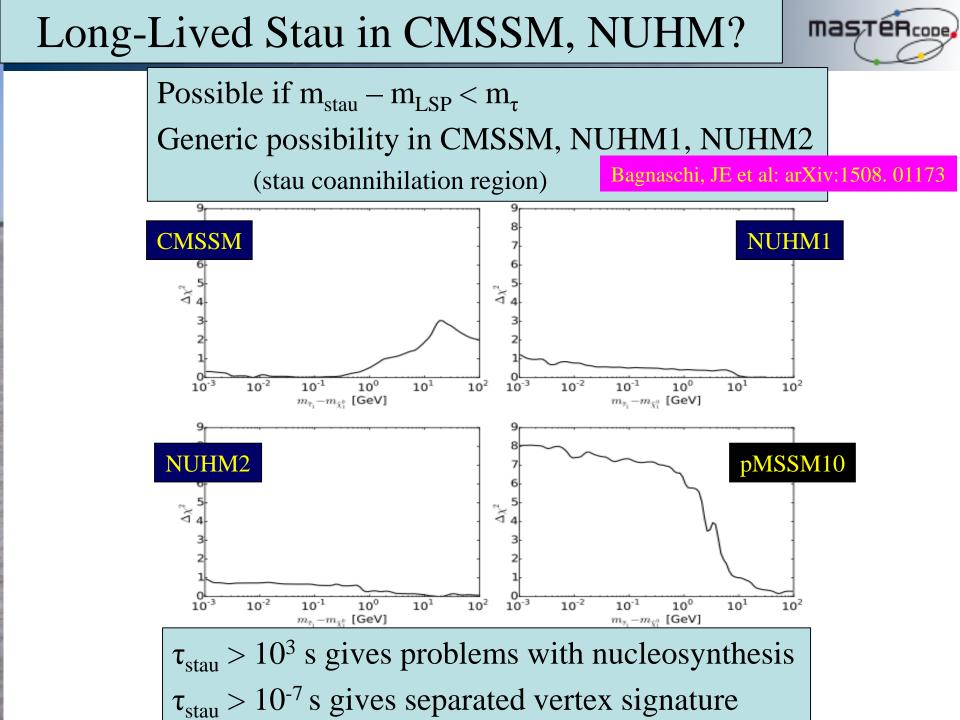
with GUT-scale unification

Anomalous Magnetic Moment of Muon

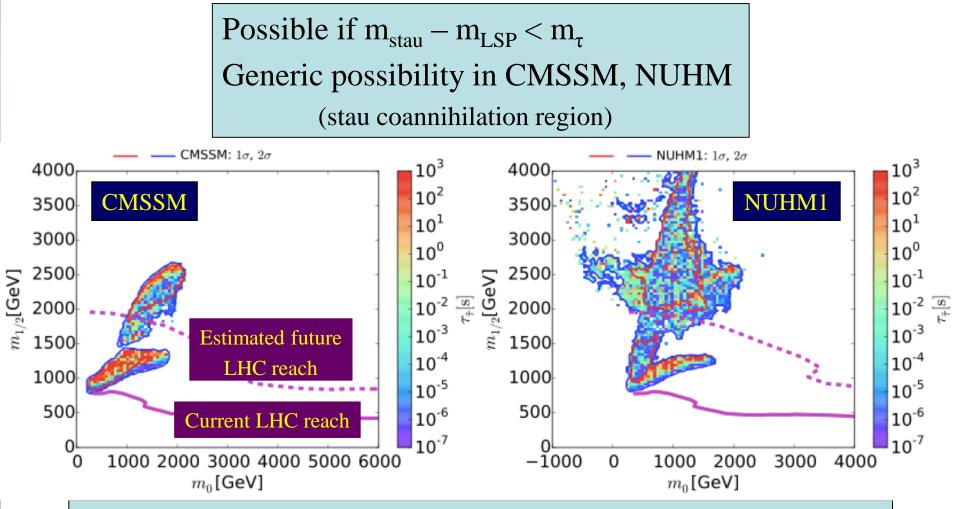
2012 ATLAS + CMS with 20/fb of LHC Data



of
$$g_u - 2$$



Long-Lived Stau in CMSSM, NUHM?



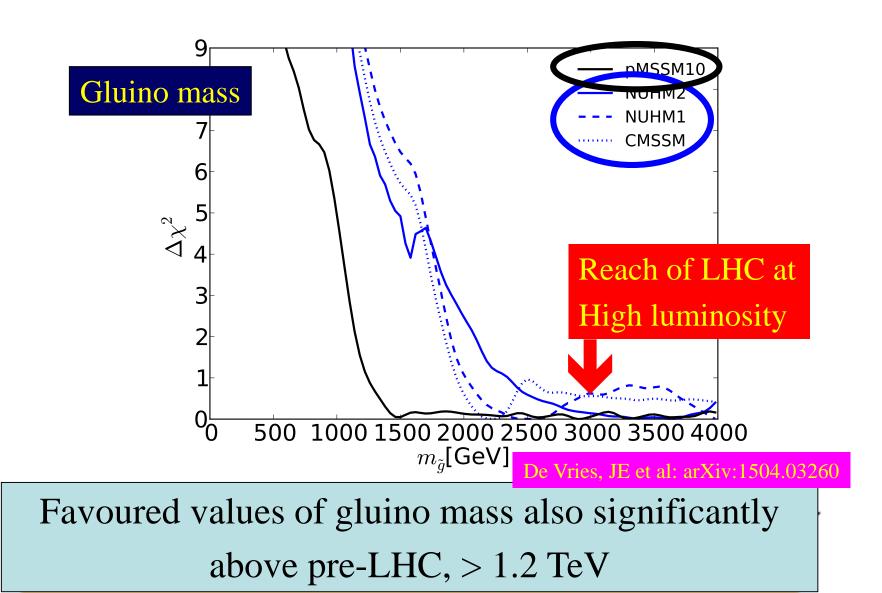
 $\tau_{stau} > 10^3$ s gives problems with nucleosynthesis $\tau_{stau} > 10^{-7}$ s gives separated vertex signature for τ -like decays

Mas/TeRcope

Fits to Supersymmetric Models

2012 ATLAS + CMS with 20/fb of LHC Data

mas/Tencore/

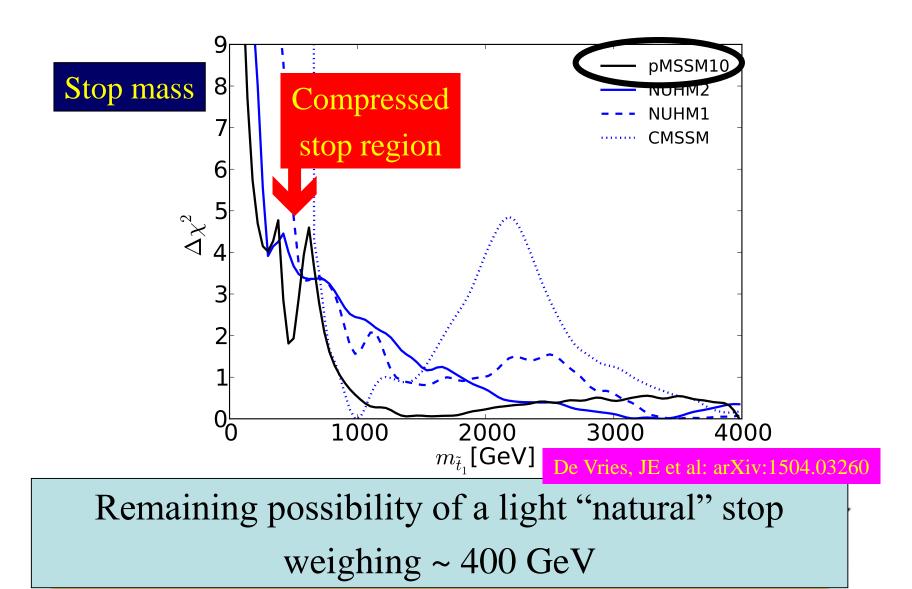


mas/TéRcope Fits to Supersymmetric Models 2012 ATLAS + CMS with 20/fb of LHC Data 9 nMSSM10 Squark mass NUHM1 CMSSM 6 $\Delta\chi^2$ Reach of LHC at High luminosity 0 1000 1500 2000 2500 3000 3500 4000 500 $m_{\tilde{a}}[\text{GeV}]$ De Vries, JE et al: arXiv:1504.03260 Favoured values of squark mass significantly above pre-LHC, ~ 1.5 TeV or more

Fits to Supersymmetric Models

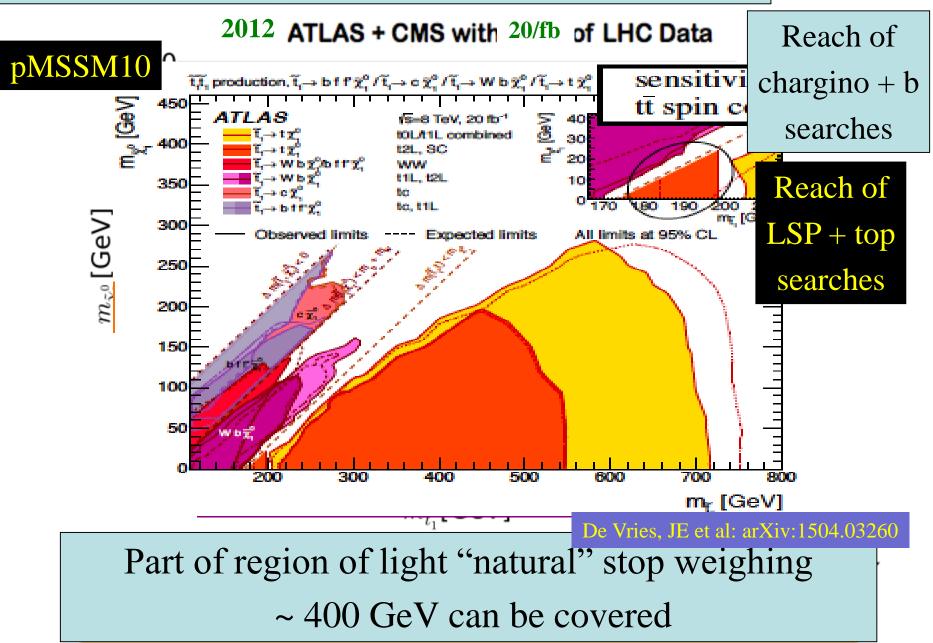


mas/Tencore/



Exploring Light Stops @ Run 2



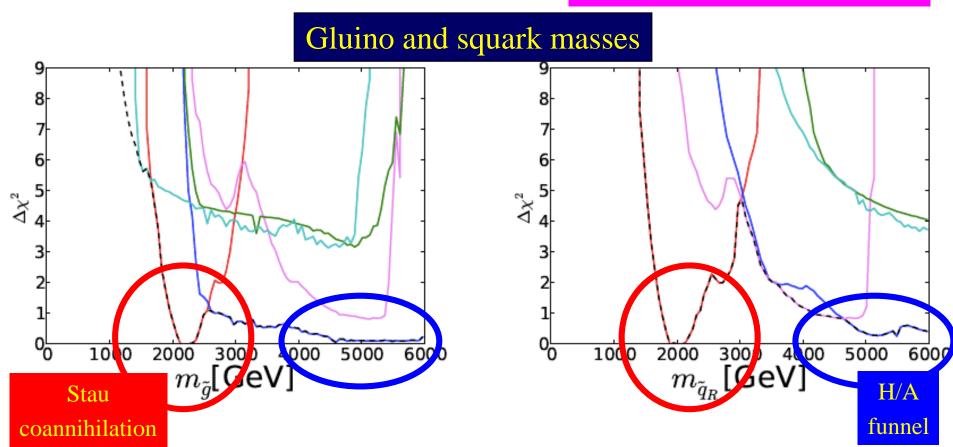


Fits to Supersymmetric Models



2012 ATLAS + CMS with 20/fb of LHC Data

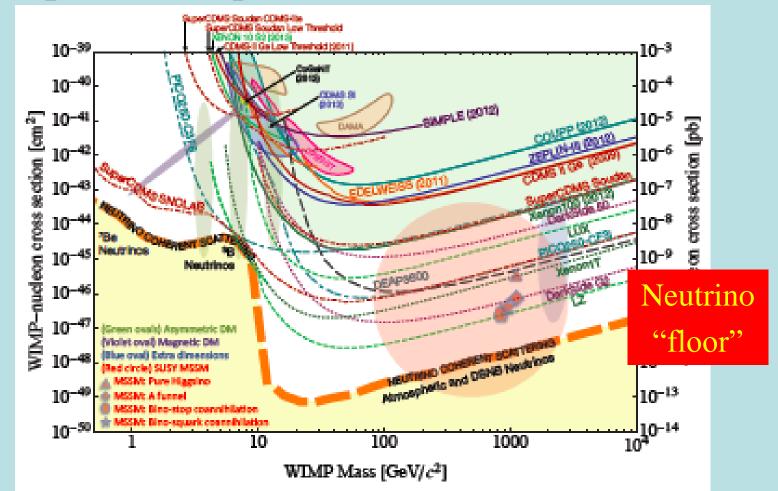
Buchmueller, JE et al: arXiv:1312.5250



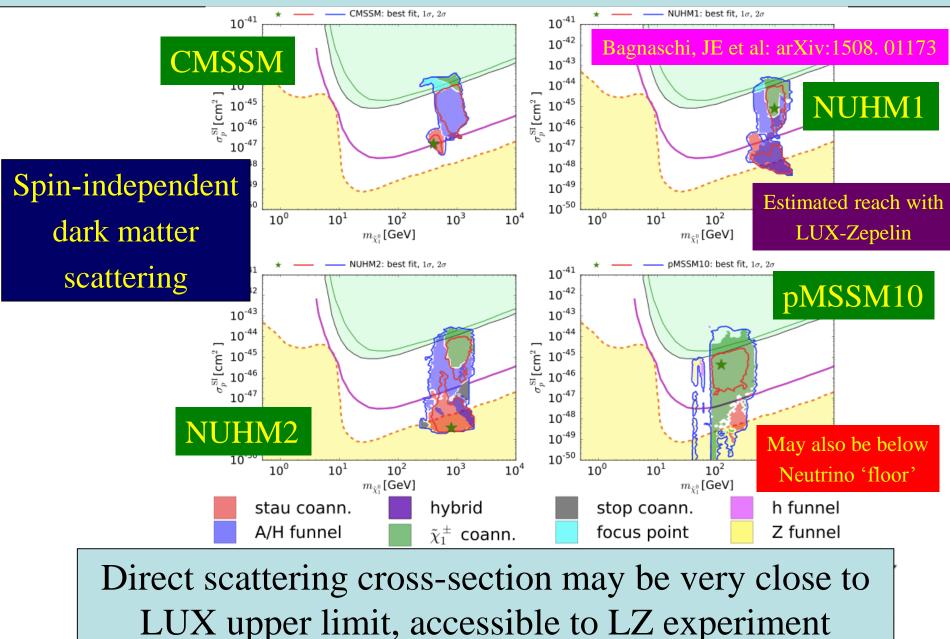
Favoured values of gluino and squark masses significantly above pre-LHC, ~ 2 TeV or more

Direct Dark Matter Searches

• Compilation of present and future sensitivities



Direct Dark Matter Search



Prospects for SUSY Searches

• Different models, various dark matter mechanisms

DM	Exp't	Exp't Models							
mechanism		CMSSM	NUHM1	NUHM2	pMSSM10				
$ ilde{ au_1}$	LHC	$\checkmark E_T, \checkmark LL$	$(\checkmark \not\!\!\!E_T, \checkmark \text{LL})$	$(\checkmark \not\!\!\!E_T, \checkmark \mathrm{LL})$	$(\checkmark \not\!\!\!E_T), \times \mathrm{LL}$				
coann.	DM	(√)	(√)	×	×				
$ ilde{\chi}_1^\pm$	LHC		×	×	$(\checkmark \not\!\!\!\!/ E_T)$				
coann.	DM		\checkmark	\checkmark	(√)				
$ ilde{t}_1$	LHC		_	$\checkmark E_T$	-				
coann.	DM	_ /	_	×	-				
A/H	LHC	$\checkmark A/H$	$(\checkmark A/H)$	$(\checkmark A/H)$	-				
funnel	DM		\checkmark	(√)	-				
Focus	LHC	$(\checkmark \not\!$	_	_	-				
point	DM	1	_	_	-				
h, Z	LHC		_	_	$(\checkmark \not\!\!\!\!\!/ E_T)$				
funnels	DM		—	-	(√)				

• No guarantees, but good prospects