

Supersymmetric AdS backgrounds and their moduli spaces

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Motivation

1. **AdS backgrounds often appear as intermediate step (before “uplifting”) in models of string phenomenology**
2. **AdS backgrounds feature prominently in AdS/CFT correspondence**

In both cases their moduli space \mathcal{C} also is of interest

1. **moduli stabilization/values of physical parameters**
2. **corresponds to conformal manifold of dual CFT**

⇒ study (supersymmetric) AdS backgrounds
and their moduli spaces in supergravity
in arbitrary space-time dimensions
with arbitrary number of supercharges

$D = 4, N = 1$ – basic formul

chiral multiplet: $(\phi, \chi)^i$, gravitational multiplet: $(g_{\mu\nu}, \psi_\mu)$

$$\mathcal{L}_{\text{bosonic}} = \frac{1}{2}R - \mathbf{G}_{i\bar{j}} \partial_\mu \phi^i \partial_\mu \bar{\phi}^{\bar{j}} - \mathbf{V} , \quad i, \bar{j} = 1, \dots, n_c ,$$

where

$$\mathbf{G}_{i\bar{j}} = \partial_i \partial_{\bar{j}} \mathbf{K} , \quad \mathbf{V} = |\mathbf{A}_{1i}|^2 - 3|\mathbf{A}_0|^2 ,$$

$$\mathbf{A}_0 = e^{\frac{1}{2}\mathbf{K}} \mathbf{W}(\phi) , \quad \mathbf{A}_{1i} = \mathbf{D}_i \mathbf{A}_0 , \quad \mathbf{D}_i \mathbf{A}_0 = \partial_i \mathbf{A}_0 + \frac{1}{2}(\partial_i \mathbf{K}) \mathbf{A}_0$$

Supersymmetry transformation of the fermions

$$\delta \chi_i = \mathbf{A}_{1i} \epsilon + \dots , \quad \delta \psi_\mu = D_\mu \epsilon + \mathbf{A}_0 \gamma_\mu \epsilon + \dots$$

Supersymmetric AdS backgrounds: $\langle \delta \psi_\mu \rangle = \langle \delta \chi_i \rangle = 0$

$$\Rightarrow \langle \mathbf{A}_{1i} \rangle = 0 , \quad \langle \mathbf{A}_0 \rangle \neq 0 , \quad \langle \mathbf{A}_0 \rangle \sim \text{cos. const.} / \text{AdS-radius}$$

Classification for $N = 1$ difficult/impossible (?)

$D = 4, N = 1$ – supersymmetric moduli space

1. Minkowski background ($\langle \mathbf{A}_0 \rangle = \mathbf{0}$)

generic W : no moduli space

special W : holomorphic $W(\phi)$ implies:
 moduli space \mathcal{C} is Kähler submanifold
 of original Kähler field space \mathcal{M}
 (and protected by supersymmetry)

2. AdS background $\langle \mathbf{A}_0 \rangle \neq 0$

- $\bar{\phi}$ cannot drop out of $\mathbf{A}_{1\phi} = \partial_\phi \mathbf{A}_0 + \frac{1}{2}(\partial_\phi K)\mathbf{A}_0$ and $\langle V \rangle$
- no protection against quantum corrections
- moduli space \mathcal{C} is real submanifold of original Kähler manifold \mathcal{M} , at best with $\dim(\mathcal{C}) = \frac{1}{2}\dim(\mathcal{M})$

$D = 4, N = 1$ – supersymmetric moduli space

necessary condition:

$$\delta\langle\mathbf{A}_0\rangle = \langle\mathbf{A}_{1i}\rangle\delta\phi^i + \langle\bar{\mathbf{A}}_{1i}\rangle\delta\bar{\phi}^{\bar{i}} = 0, \quad \delta\langle\bar{\mathbf{A}}_0\rangle = \dots = 0$$

$$\delta\langle\mathbf{A}_{1i}\rangle = \langle\partial_j\mathbf{A}_{1i}\rangle\delta\phi^j + \frac{1}{2}\langle K_{i\bar{j}}\mathbf{A}_0\rangle\delta\bar{\phi}^{\bar{j}} = 0, \quad \delta\langle\bar{\mathbf{A}}_{1i}\rangle = \dots = 0$$

(1st eq. identically satisfied in supersymmetric backgrounds)

Rewrite 2nd condition:

$$\mathbb{M} \begin{pmatrix} \delta\phi^j \\ \delta\bar{\phi}^{\bar{j}} \end{pmatrix} = 0, \quad \text{with} \quad \mathbb{M} = \begin{pmatrix} \mathbf{m}_{ij} & \langle\mathbf{K}_{i\bar{j}}\mathbf{A}_0\rangle \\ \langle\mathbf{K}_{i\bar{j}}\bar{\mathbf{A}}_0\rangle & \bar{\mathbf{m}}_{\bar{i}\bar{j}} \end{pmatrix},$$

where $\mathbf{m}_{ij} = \langle\nabla_i\mathbf{A}_{1j}\rangle =$ fermionic mass matrix

⇨ generically \mathbb{M} has full rank $\Rightarrow \delta\phi^j = 0 \Rightarrow$ no moduli space

⇨ for special (tuned) \mathbf{W}/\mathbf{A}_0 : moduli space possible [DLMTW, ...]

Maximally supersymmetric backgrounds in arbitrary dim. D

Distinguish two cases: backgrounds **with/without** fluxes

1. **backgrounds without fluxes**

Maximally supersymmetric backgrounds in arbitrary dim. D

Distinguish two cases: backgrounds **with/without** fluxes

1. **backgrounds without fluxes**

Supersymmetry transformation of fermions in gauged supergravities

$$\delta\psi_\mu^I = D_\mu\epsilon^I + \mathbf{A}_{0J}^I(\phi)\gamma_\mu\epsilon^J + \dots, \quad \delta\chi^i = \mathbf{A}_{1J}^i(\phi)\epsilon^J + \dots$$

$$\text{with } V \sim |\mathbf{A}_1|^2 - c(d, q)|\mathbf{A}_0|^2$$

Supersymmetric backgrounds preserving all q supercharges:

$$\langle D_\mu\epsilon^I + \mathbf{A}_{0J}^I(\phi)\gamma_\mu\epsilon^J \rangle = 0, \quad \langle \mathbf{A}_1 \rangle = 0, \quad \forall I, J$$

Consistency condition: $R_{\mu\nu\rho\sigma} \sim \text{tr} \mathbf{A}_0^2 (g_{\mu\rho}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\rho})$

Solutions: AdS_D or $M_d \times T^{(D-d)}$, $1 \leq d \leq D$

(note: $\text{AdS}_d \times Y^{(D-d)}$ not possible)

Supersymmetric moduli spaces \mathcal{C} of AdS_D

- supersymmetric moduli spaces $\mathcal{C} \subset \mathcal{M}$
- necessary condition: $\delta_\phi \langle \mathbf{A}_0 \rangle = 0 = \delta_\phi \langle \mathbf{A}_1 \rangle$
- Gradient-flow equations: [d'Auria, Ferrara]

$$\delta \langle \mathbf{A}_0 \rangle = \langle \mathbf{A}_1 \rangle \delta \vec{\phi} = 0, \quad \delta \langle \mathbf{A}_1 \rangle = \mathbb{M} \delta \vec{\phi}, \quad \mathbb{M} := \mathbf{m}_f + \langle \mathbf{A}_0 \rangle$$

- necessary condition for moduli space: \mathbb{M} has null eigenspace

$$\langle \mathbf{A}_0 \rangle \begin{cases} = 0 & \mathcal{C} \subset \mathcal{M} \text{ has same* geometry as } \mathcal{M} \\ \neq 0 & \mathcal{C} \subset \mathcal{M} \text{ has different geometry as } \mathcal{M} \end{cases}$$

- $q = 4$: can tune superpotential/ \mathbf{m}_f
but \mathcal{C} generically destroyed by quantum corrections
- $q = 8, 16, 32$: no superpotential, can choose spectrum and gauged isometries of \mathcal{M} , supersymmetric protection

Example: $d = 5, q = 16$ supercharges [JL, Triendl, Zagermann]

⇒ multiplets and scalar geometry

- gravity $(g_{\mu\nu}, \psi_{\mu}^i, A_{\mu}^{0,m}, \chi^i, \Sigma)$, $m = 1, \dots, 5 \Rightarrow 6$ graviphotons
- vector $(A_{\mu}^a, \lambda^a, \phi^{am})$, $a = 1, \dots, n_v$

scalar field space: $\mathcal{M}(\Sigma, \phi) = \mathbb{R}^+ \times \frac{\mathbf{SO}(5, n_v)}{\mathbf{SO}(5) \times \mathbf{SO}(n_v)}$

⇒ supersymmetric AdS backgrounds can be “classified”

gauge group:

$$\mathbf{G} = \mathbf{U}(1) \times \hat{\mathbf{G}} \subset \mathbf{SO}(5, n_v)$$

where $\hat{\mathbf{G}} \subset \mathbf{SO}(3, n_v) \rightarrow \mathbf{SO}(3) \times \mathbf{H}$, $\mathbf{H} =$ compact subgroup

⇒ moduli space possible: [Günaydin, Pilch, Warner, Zagermann]

$$\mathcal{C} = \frac{\mathbf{SU}(1, p)}{\mathbf{U}(1) \times \mathbf{SU}(p)} \subset \frac{\mathbf{SO}(2, 2p)}{\mathbf{SO}(2) \times \mathbf{SO}(2p)} \subset \frac{\mathbf{SO}(5, n_v)}{\mathbf{SO}(5) \times \mathbf{SO}(n_v)}$$

General story: AdS moduli space for $\mathcal{M} = G/H$

[JL, S. Lüster, Rüter]

- ⇨ gauging isometries of \mathcal{M} induces gauging of R-symmetry H_R
- ⇨ Maximally supersymmetric AdS backgrounds:
 - $H_R^g \subset H_R$ gauged only by (subset of) graviphotons A_μ
- ⇨ A_0 and A_μ transform in specific rep. of H_R depending on D, q
- ⇨ $\langle A_0 \rangle \neq 0$ implies $H_R \rightarrow H_R^l$ (little group)
- ⇨ additional constraint:
 - $H_R^g \subset H_R^l \subset H_R$ of AdS-vacuum has to have A_μ in its adj. rep.
- ⇨ flat directions of V :
 - H_R^g charged Goldstone bosons + H_R^g -singlets = moduli

Example: $d = 7, q = 32$ supercharges

[**JL, S. Lüst, P. Rüter**]

⇨ gravity multiplet: $(g_{\mu\nu}, \psi_{\mu}^I, A_{\mu}^{[MN]}, \chi^{IJK}, \phi)$

$$I, J, K = 1, \dots, 4, \quad M, N = 1, \dots, 5$$

⇨ scalar field space: $\mathcal{M}(\phi) = \frac{\mathbf{SL}(5)}{\mathbf{SO}(5)}$

⇨ $A_0 \in \mathbf{1} \oplus \mathbf{5}, \quad A_1 \in \mathbf{5} \oplus \mathbf{14} \oplus \mathbf{35}$ of $H_R = SO(5) \sim USp(4)$

$$\Rightarrow H_R = H_R^l = H_R^g = USp(4)$$

⇨ $\delta\phi \in \mathbf{14} \Rightarrow$ no GB, no singlets \Rightarrow no moduli

Results so far:

$\text{AdS}_{(d,q)}$	\mathcal{M}	\mathcal{C}	[]
(4,4)	Kähler	Real	[DLMTW]
(4,8)	SK \times QK	Real \times Kähler	[DMLTW]
(5,8)	SR \times QK	Kähler	[Tachikawa, LM]
(6,8)	$\frac{O(1, n_T)}{O(n_T)} \times$ QK	–	[?]

SK: Special Kähler

QK: Quaternionic Kähler

SR: Special Real

[DMLTW]=[de Alwis, JL, McAllister, Triendl, Westphal]

[LM]=[JL, Muranaka]

Results so far:

$\text{AdS}_{(d,q)}$	\mathcal{M}	H_R^g	\mathcal{C}	[]
(4,16)	$\frac{SO(6,6+n_v)}{SO(6) \times SO(6+n_v)} \times \frac{SU(1,1)}{U(1)}$	$SO(4)$	AX	[LT]
(5,16)	$\frac{SO(5,5+n_v)}{SO(5) \times SO(5+n_v)} \times \mathbb{R}^+$	$U(2)$	$\frac{SU(1,p)}{SU(p) \times U(1)}$	[GPWZ,LTZ]
(6,16)	$\frac{SO(4,4+n_v)}{SO(4) \times SO(4+n_v)} \times \mathbb{R}^+$	$SU(2) \times SU(2)$	AX	[KL]
(7,16)	$\frac{SO(3,3+n_v)}{SO(3) \times SO(3+n_v)} \times \mathbb{R}^+$	$USp(2)$	AX	[LL]
(4,32)	$\frac{E_{7,7}}{SU(8)}$	$SO(8)$	AX	[LLR]
(5,32)	$\frac{E_{6,6}}{USp(8)}$	$SU(4)$	$\frac{SU(1,1)}{U(1)}$	[LLR]
(6,32)	$\frac{SO(5,5)}{SO(5) \times SO(5)}$	XX	XX	[LLR]
(7,32)	$\frac{SL(5)}{SO(5)}$	$USp(4)$	AX	[LLR]

AX: AdS background exists and can be classified but no moduli space,

XX: AdS background does not exist

[LT]=[JL,Triendl]

[LLR]=[JL, Lüst, Rüter]

[KL]=[Karndumri,JL]

AdS/CFT correspondence

AdS_d bulk with q supercharges

\Leftrightarrow

SCFT_{d-1} on AdS_d boundary with $\frac{q}{2} + \frac{q}{2}$ super+sconformal charges

AdS

\Leftrightarrow

SCFT

(broken) gauge group

scalars with mass m

moduli

moduli space \mathcal{C}

(anomalous) flavour group

couplings of gauge invariant operators

exactly marginal couplings

conformal manifold \mathcal{C}

$$S = S_0 + \sum_i \int \varphi^i \mathbf{O}_i$$

$$g_{ij}(\varphi) = x^{2d} \langle \mathbf{O}_i(\mathbf{x}) \mathbf{O}_j(\mathbf{0}) \rangle_S \text{ [Zamolodchikov]}$$

Compare AdS with SCFT-results

$\text{AdS}_{(d,q)}$	\mathcal{C}	[]	$\text{SCFT}_{(d-1,\frac{q}{2})}$	$\mathcal{C}_{\text{SCFT}}$	[]
(4,4)	Real	[DLMTW]	(3,2)	Real ¹	[CDI]
(4,8)	Real \times Kähler	[DMLTW]	(3,4)	Kähler	[GKSTW]
(5,8)	Kähler	[T,LM]	(4,4)	Kähler	[Ansin]
(6,8)	–	[?]	(5,4)	SXX	

SX: SCFT exists but no moduli space,

SXX: SCFT does not exist

[CDI]=[Cordova,Dumitrescu,Intriligator]

[GKSTW]=[Green,Komargodski,Seiberg,Tachikawa,Wecht]

¹ none known, no susy protection

Compare AdS with SCFT-results (Note: no SCFT in $d > 6$)

AdS _(d,q)	\mathcal{C}	[]	SCFT _(d-1, $\frac{q}{2}$)	$\mathcal{C}_{\text{SCFT}}$	[]
(4,16)	AX	[LT]	(3,8)	SX	[CDI]
(5,16)	$\frac{SU(1,p)}{SU(p) \times U(1)}$	[GPWZ,LTZ]	(4,8)	Kähler (?) ²	[P,GGK]
(6,16)	AX	[KL]	(5,8)	SX	[CDI]
(7,16)	AX	[LL]	(6,8)	SX	[LL]
(4,32)	AX	[LLR]	(3,16)	SX	[CDI]
(5,32)	$\frac{SU(1,1)}{U(1)}$	[LLR]	(4,16)	$\frac{SU(1,1)}{U(1)}$	[?]
(6,32)	XX	[LLR]	(5,16)	SXX	[Nahm]
(7,32)	AX	[LLR]	(6,16)	SX	[LL]

[GPWZ]=[Günaydin,Pilch,Warner,Zagermann]

[P,GGK]= [Papadodimas; Gerchkovitz,Gomis,Komargodski]

²[GGK] only show Kähler, [P] shows tt^* -geometry

Supersymmetric backgrounds in arbitrary dimensions D

Distinguish two cases: backgrounds **with/without** fluxes

2. **backgrounds with fluxes**

Supersymmetric backgrounds in arbitrary dimensions D

Distinguish two cases: backgrounds **with/without** fluxes

2. **backgrounds with fluxes** [JL, S. Lüst, ...]

Supersymmetry transformation of the fermions

$$\delta\psi_\mu^I = D_\mu\epsilon^I + (\mathcal{F}_{0\mu})_{\mathbf{J}}^{\mathbf{I}}\epsilon^J + A_{0J}^I\gamma_\mu\epsilon^J, \quad \delta\chi^i = (\mathcal{F}_1)_{\mathbf{J}}^{\mathbf{i}}\epsilon^J + A_{1J}^i\epsilon^J,$$

$$\delta\lambda^A = (\mathcal{F}_2)_{\mathbf{J}}^{\mathbf{A}}\epsilon^J + A_{2J}^A\epsilon^J,$$

$\mathcal{F}_{0,1/2}$ = sum of fluxes in gravitational multiplet/other multiplets

supersymmetric backgrounds: $\mathcal{F}_1 = \mathcal{F}_2 = A_1 = A_2 = 0 \Rightarrow \mathcal{F}_0 = 0$

two exceptions:

1. chiral theories with selfdual fluxes
2. supergravities without χ 's in gravitational multiplet

Supersymmetric backgrounds in arbitrary dimensions D

⇨ candidate supergravities/fluxes:

dimension	supersymmetry	q	possible flux	classification of maximal supersymmetric solutions
$D = 11$	$N = 1$	32	$F^{(4)}$	[Figueroa-O'Farrill,Papadopoulos]
$D = 10$	IIB	32	$F_+^{(5)}$	[Figueroa-O'Farrill,Papadopoulos]
$D = 6$	$N = (2, 0)$	16	$5 \times F_+^{(3)}$	[Chamseddine,Figueroa-O'Farrill,Sabra]
$D = 6$	$N = (1, 0)$	8	$F_+^{(3)}$	[Gutowski,Martelli,Reall]
$D = 5$	$N = 2$	8	$F^{(2)}$	[Gauntlett,Gutowski,Hull,Pakis,Reall]
$D = 4$	$N = 2$	8	$F^{(2)}$	[Tod]

⇨ study this list case by case to find

- $\langle A_0 \rangle = \langle A_1 \rangle = \langle A_2 \rangle = 0$
- no flux for R-symmetry possible

[Hristov,Looyestijn,Vandoren;Gauntlett,Gutowski;Akyol,Papadopoulos]

⇨ background solutions have to coincide with ungauged solutions

they are classified in ungauged case

(see refs. above and [Gran,Gutowski,Papadopoulos])

AdS backgrounds

dim.	SUSY	q	$AdS \times S$	
$D = 11$	$N = 1$	32	$AdS_4 \times S^7$ $AdS_7 \times S^4$	[Freund,Rubin]
$D = 10$	IIB	32	$AdS_5 \times S^5$	[Schwarz,West]
$D = 6$	$N = (2, 0)$	16	$AdS_3 \times S^3$	[Gibbons,Horowitz,Townsend]
	$N = (1, 0)$	8	$AdS_3 \times S^3$	[Gibbons,Horowitz,Townsend]
$D = 5$	$N = 2$	8	$AdS_2 \times S^3$	[Gibbons,Horowitz,Townsend]
			$AdS_3 \times S^2$	[Chamseddine,Ferrara,Gibbons,Kallosh]
$D = 4$	$N = 2$	8	$AdS_2 \times S^2$	[Bertotti,Robinson]

list is exhaustive (cf. [Alonso-Alberca,Ortin])

in addition:

- Hpp-wave solutions [Penrose; Kowalski-Glikman; Blau,Figueroa-O'Farrill,Hull,Papadopoulos]
- further special solutions in $D = 5$
[Cvetic,Larsen; Gauntlett,Meyers,Townsend; Gauntlett,Gutowski,Hull,Pakis,Reall]

Conclusion/Outlook

Studied fully supersymmetric AdS backgrounds in D -dimensional supergravities with q supercharges

1. without fluxes

- classified AdS_D backgrounds with $q = 16, 32$ supercharges and determined their classical moduli spaces
- determined the geometry of classical moduli spaces for $q = 4, 8$

2. with fluxes

- AdS-backgrounds only exist in
 - chiral supergravities
 - supergravities without graviphotinos
- in both cases AdS-solution coincides with solution of corresponding ungauged supergravity which are classified

Results are consistent with dual SCFTs (not yet clear in (4,8))