

# FCC-eh SM and BSM Higgs Studies

**Uta Klein**  
on behalf of  
**the LHeC/FCC-eh Higgs Group**



UNIVERSITY OF  
LIVERPOOL



# SM Higgs Production in ep

## CC : LO SM Higgs Production

e-p (swap charges for e+p)

e- u -> ve h d

e- d~ -> ve h u~

electrons →

$E_T^{\text{miss}}$

WWH

LHC protons →

Fwd jet

around 90-80%

around 10-20%

## NC : LO SM Higgs Production

e-p (swap charges for e+p)

e- d -> e- h d

e- u -> e- h u

electrons →

FS electron

ZZH

LHC protons →

Fwd jet

around 1/3

around 1/3

$\sigma$  (LO QCD CTEQ6L1  $M_H=125$  GeV)

c.m.s. energy	1.3 TeV	3.5 TeV
cross section [fb]		
NC DIS	21	127
CC DIS	109	560
CC DIS polarised cross section [fb] P=-80%	<b>196</b>	<b>1008</b>

→ In ep, direction of quark (FS) is well defined.

- Scale dependencies of the LO calculations are in the range of 5-10%.
- NLO QCD corrections are small, but shape distortions of kinematic distributions up to 20%. QED corrections up to -5%.

[J. Blumlein, G.J. van Oldenborgh, R. Ruckl, Nucl.Phys.B395:35-59,1993]  
[B.Jager, arXiv:1001.3789]

# Analysis Framework

## Event generation

- SM or BSM production
- CC & NC DIS background  
by **MadGraph5/MadEvent**

- Fragmentation
- Hadronization

by **PYTHIA** (modified for ep)

## Fast detector simulation

by **Delphes**  
→ test of FCCeh detector

S/B analysis → cuts or BDT

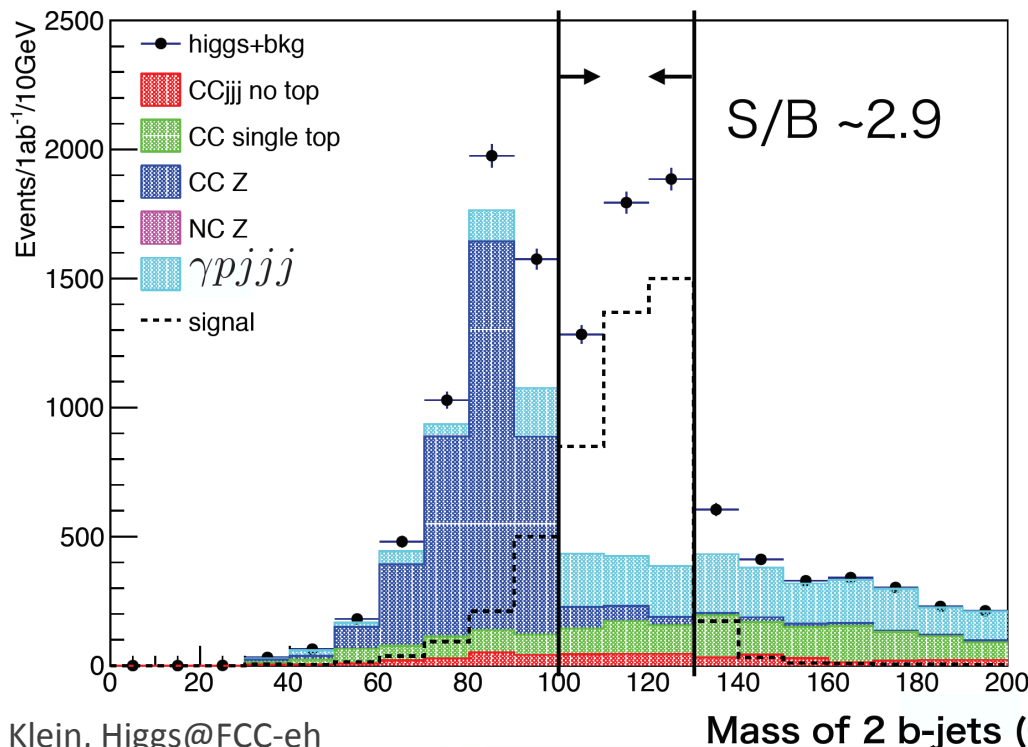
- Calculate cross section with tree-level Feynman diagrams (any UFO) using pT of scattered quark as scale ( $CDR \hat{s}$ ) for ep processes with **MadGraph5**
- Standard HERA tools can NOT to be used !
- **Higgs mass 125 GeV as default**
- Fragmentation & hadronisation uses **ep-customised Pythia**.
- **Delphes 'detector' → displaced vertices and signed impact parameter distributions → studied for LHeC, and used for FCC-eh SM Higgs extrapolations**
- powerful method to optimise detector tuning and S/N for various Higgs, top and BSM decays
- Ongoing : Integration of FCCeh into FCC simulation framework

# Cut-based Results for Hbb @ LHeC

Masahiro Tanaka, Masahiro Kuze

Various studies pursued since the LHeC CDR [ before the Higgs discovery, see <http://cern.ch/lhec> ] focusing on SM 125 GeV Higgs decay into b-quarks

- Assumed  $1000 \text{ fb}^{-1}$  of statistics. ( $\sim 10$  years running for LHeC.)
- Veto efficiency of 90% for photo-production background is assumed, using forward electron tagging.



b-tag performance (cut based)

- b-jet: 75%
- c-jet mis-tagging rate: 5%
- Light-jet mis-tagging rate: 1%

Precision of coupling constant  
(Statistics error only)

$$\kappa = \frac{\sqrt{N_s + N_b}}{2N_s}$$

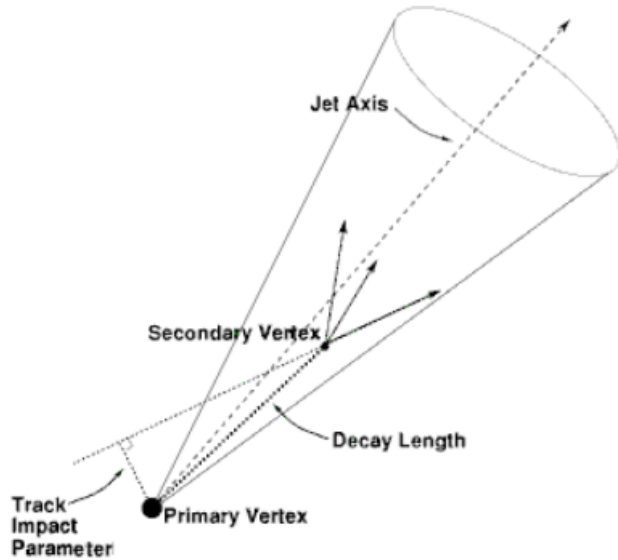
Signal: 3600

Bkg: 1250

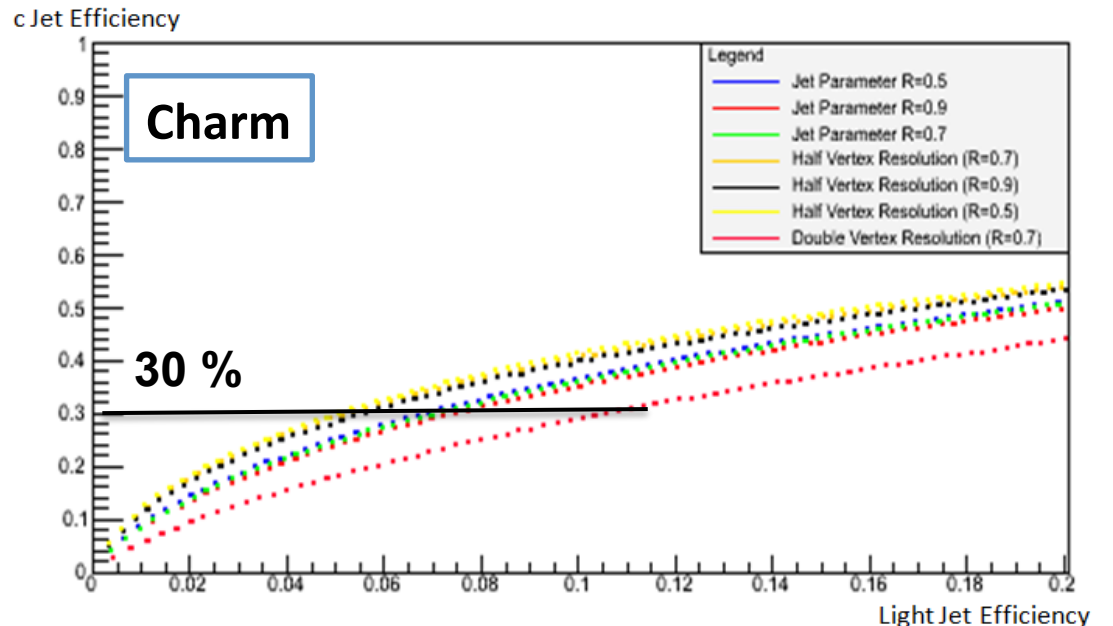
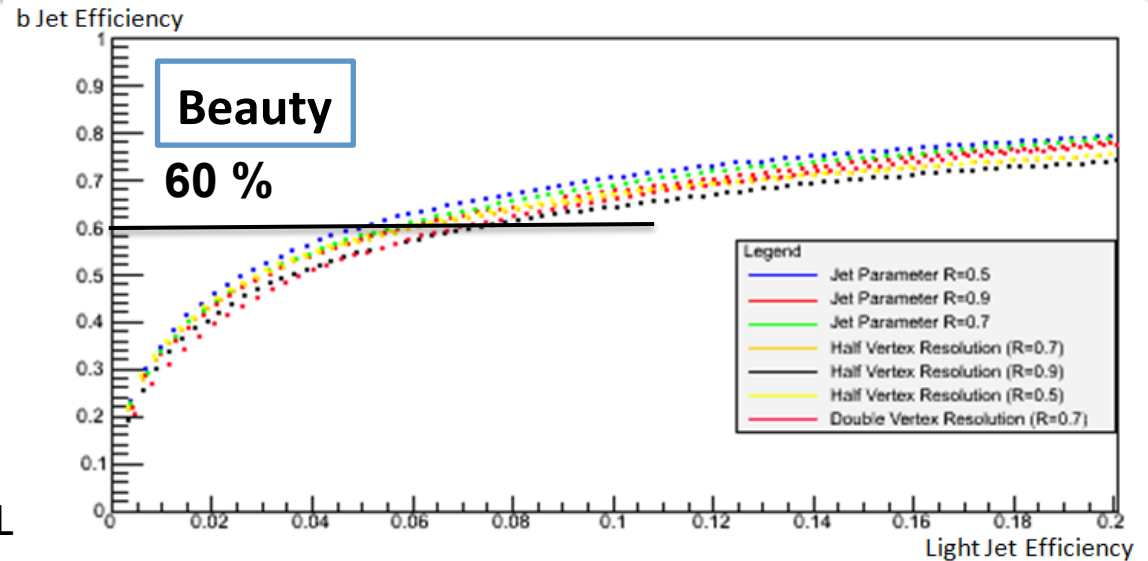
$\kappa(\text{Hbb}) \sim 0.97\%$

# HFL Tagging

Uta Klein &  
Daniel Hampson



- Realistic and conservative HFL tagging within Delphes realised, and dependence on vertex resolution (nominal 10  $\mu\text{m}$ ) and anti-kt jet radius studied
- Light jet rejection very conservative, i.e. factor 10 worse than ATLAS
- **used in full LHeC analysis and for FCC-eh extrapolations**

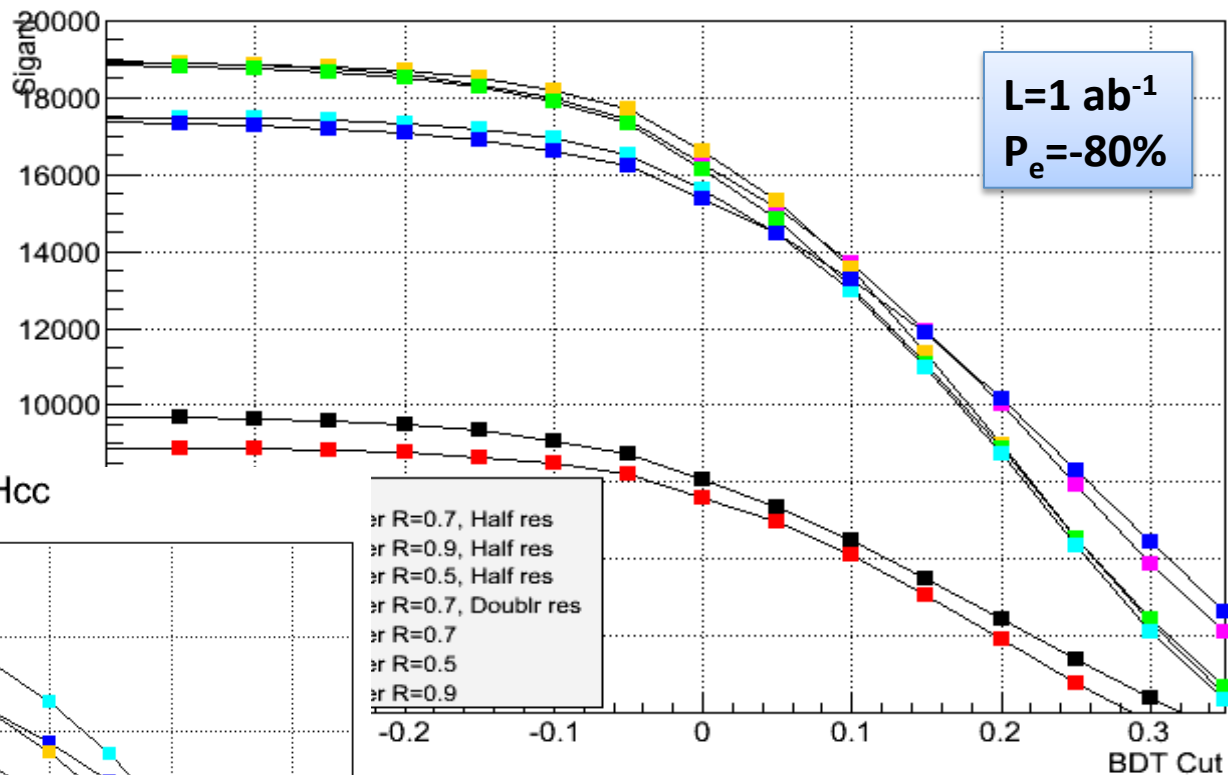


# BDT Results for Higgs @ LHeC

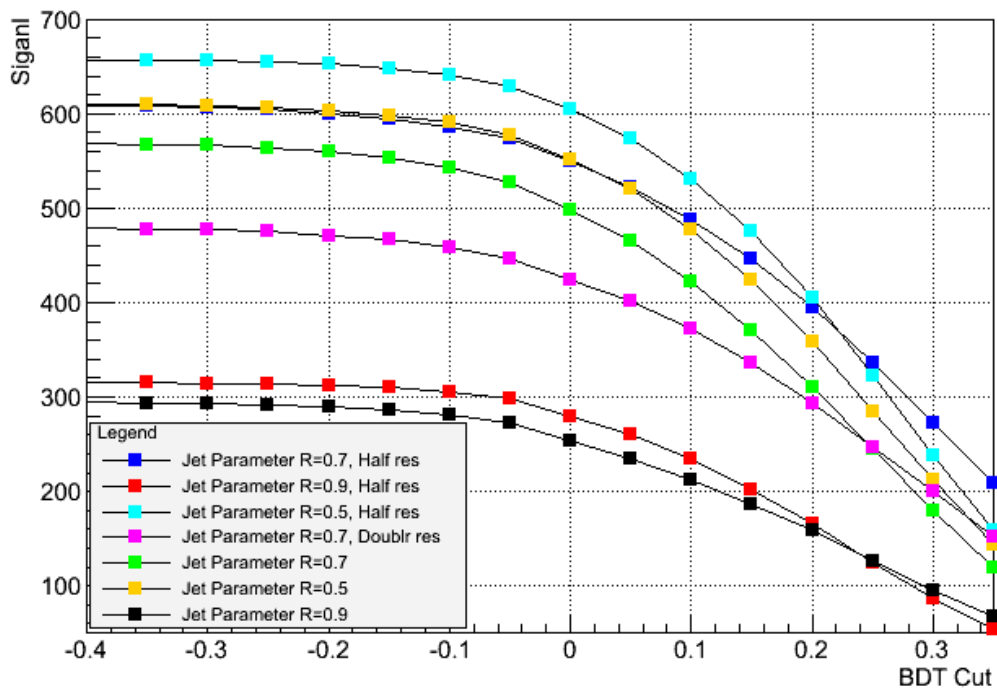
Uta Klein &  
Daniel Hampson

Signal Events Hbb

Hbb : Using same background assumptions as for cut-based analysis, we get factor 5 more Hbb candidates ( $\sim 15000$ ) and a coupling error of 0.6%.



Signal Events Hcc



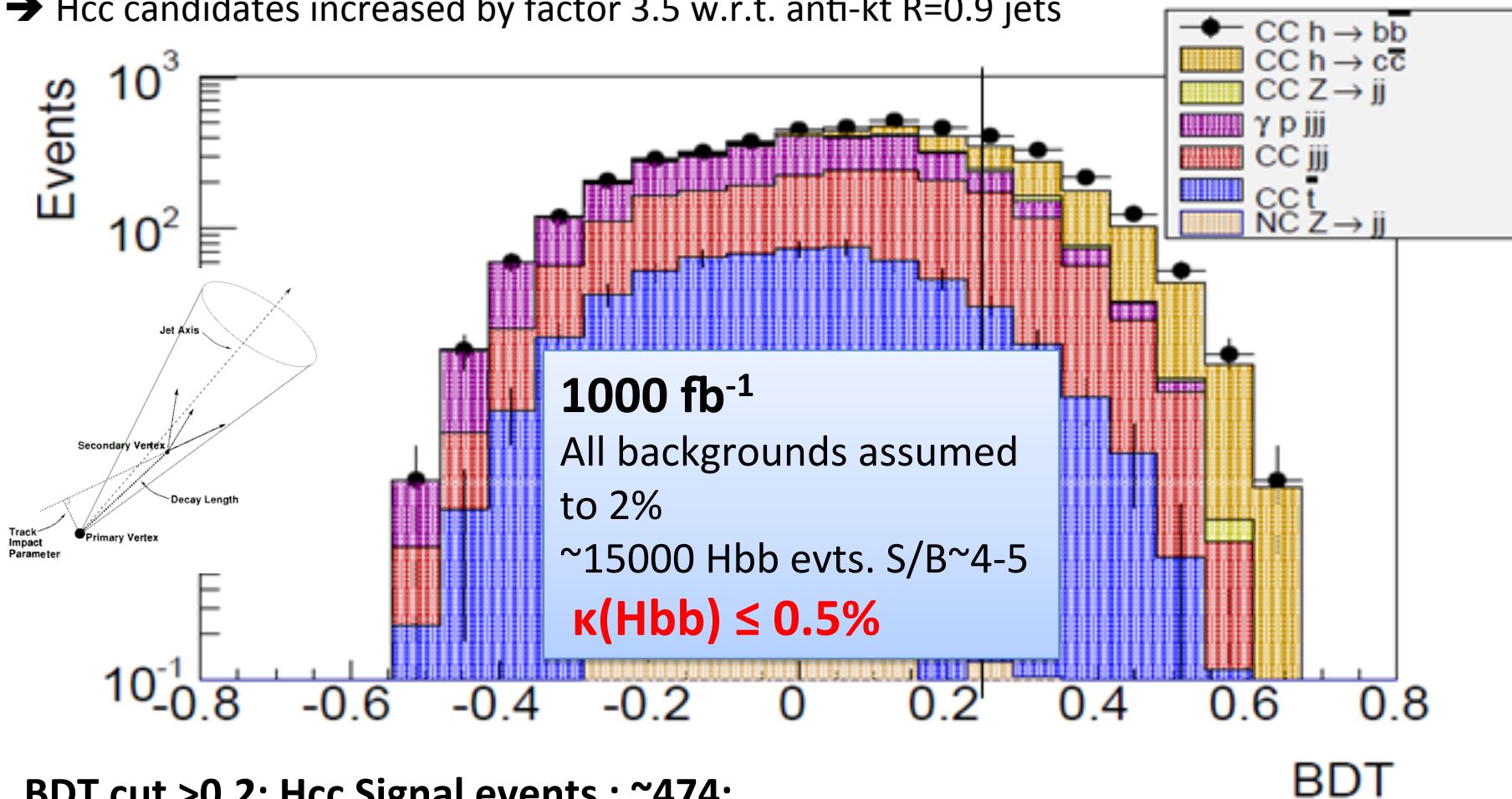
Hcc : High sensitivity to vertex resolution (nominal  $10 \mu\text{m}$ ) and jet radius  
 $\rightarrow$  expect about 400-600 Hcc candidates

# BDT Result for $H \rightarrow cc$

Uta Klein &  
Daniel Hampson

NEW : Using  $R = 0.5$  anti-kt jets and ATLAS IBL vertex resolution ( $5 \mu\text{m}$ )

→ Hcc candidates increased by factor 3.5 w.r.t. anti-kt  $R=0.9$  jets



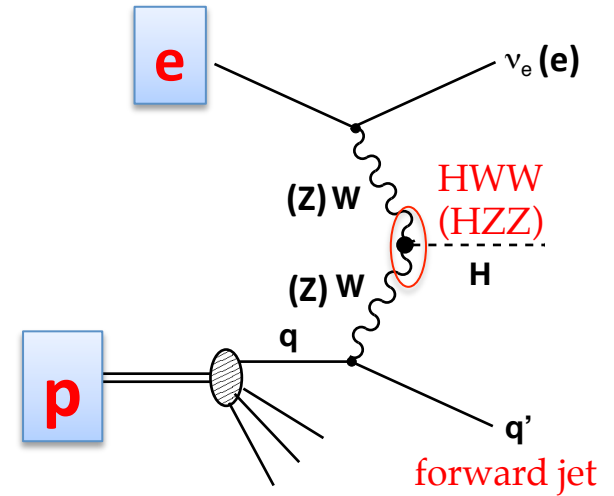
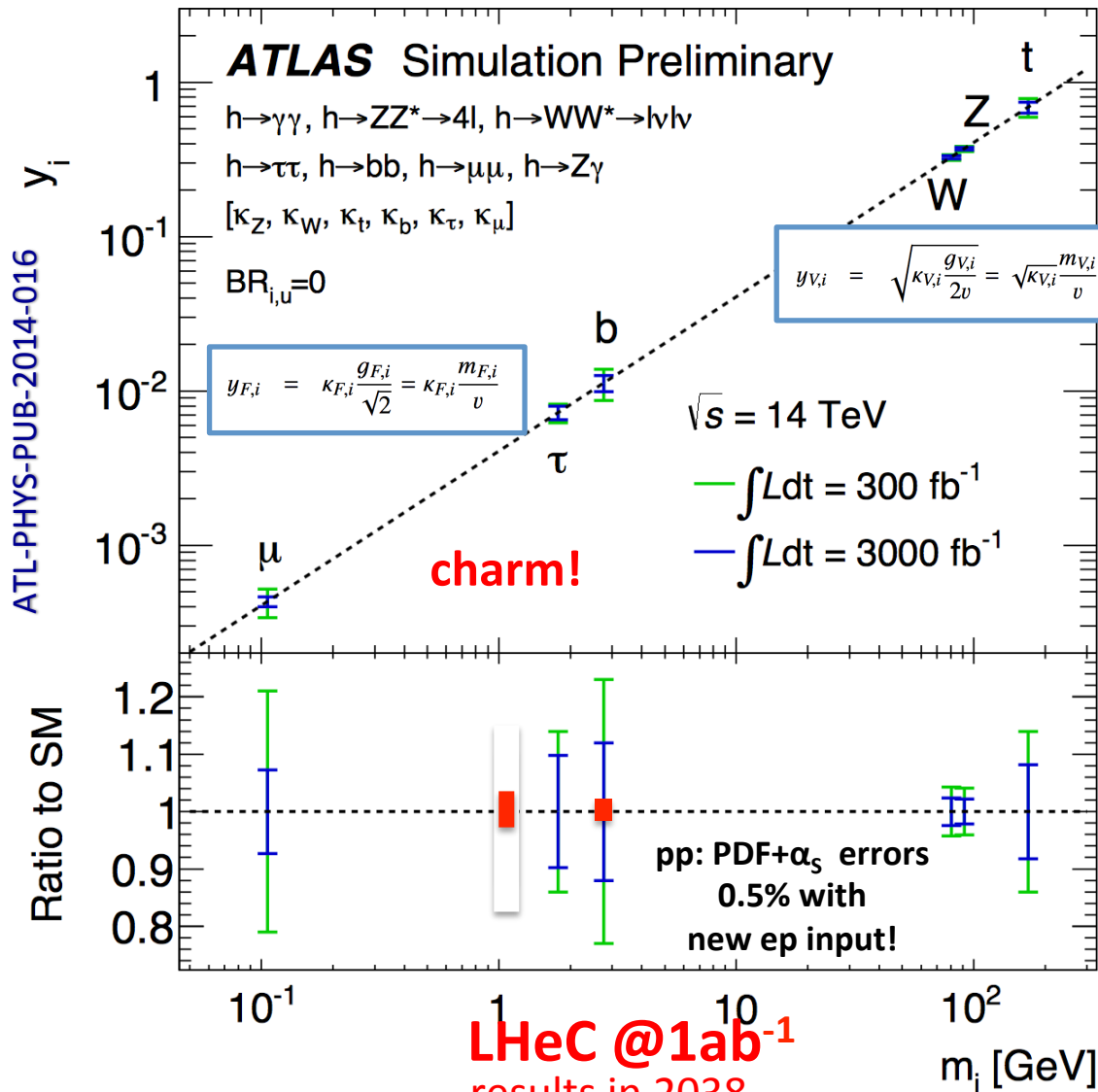
BDT cut  $> 0.2$ : Hcc Signal events :  $\sim 474$ ;

$S/\sqrt{S+B} = 12.8 \rightarrow \kappa(\text{Hcc}) = 4\%$  for  
 $1000 \text{ fb}^{-1}$

Clear potential to access the Higgs to  
charm decay channel in ep.

# Higgs Couplings at pp + ep

running concurrently



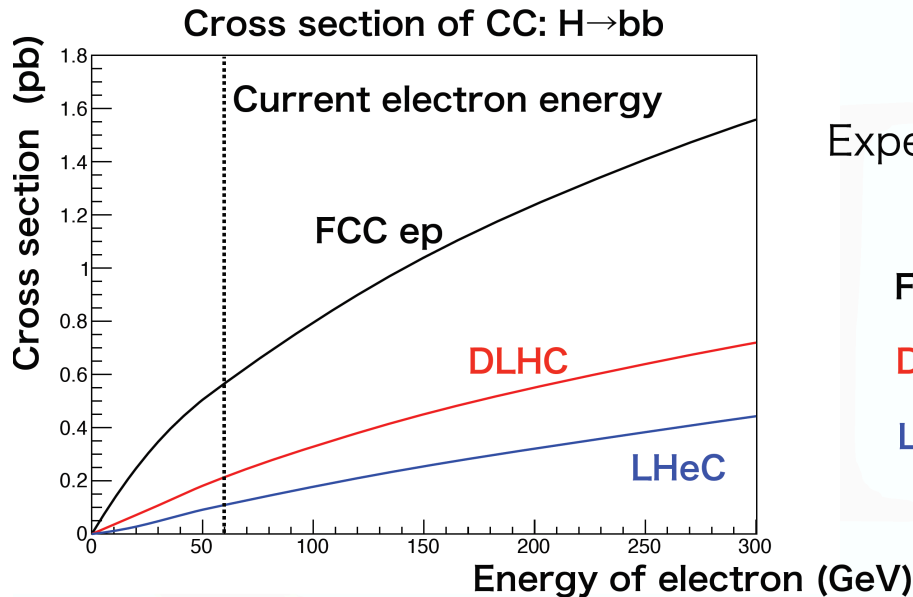
→ use ep as the 'near' detector for pp to beat  $\alpha_s$  and PDF uncertainties to  $< \sim 0.5\%$ ,  $\delta m_b$  to 10 MeV,  $\delta m_{\text{charm}}$  to 3 MeV



# SM Higgs into HFL Summary

- Assume a 60 GeV polarized electron beam and  $1000 \text{ fb}^{-1}$  ( $\sim 10$  years running)
- Expected number of signal events and error of coupling constant from BDT results.
- Background assumed to be known to  $\sim 2\%$

U Klein (Liverpool)



Expected number of signal events  
( $E_e = 60 \text{ GeV}$ )

FCC ep ( $\sim 85,000 H \rightarrow bb$  events)

DLHC ( $\sim 35,000 H \rightarrow bb$  events)

LHeC ( $\sim 15,000 H \rightarrow bb$  events)

	LHeC ( $E_p = 7 \text{ TeV}$ $\sqrt{s} \sim 1.3 \text{ TeV}$ )	DLHC ( $E_p = 14 \text{ TeV}$ $\sqrt{s} \sim 1.8 \text{ TeV}$ )	FCC ep ( $E_p = 50 \text{ TeV}$ $\sqrt{s} \sim 3.5 \text{ TeV}$ )
$\kappa (Hbb)$	0.5%	0.3%	0.2%
$\kappa (Hcc)$	4%	2.8%	1.8%

# Exploring SM EFT & New Physics

M. Trott @ LHeC Workshop 2014

<http://lhec.web.cern.ch>

*In the absence of any explicit new states, or overwhelming theory prejudice, the goal is to systematically study the SM EFT for hints of NP, using all possible future facilities to maximize physics conclusions.*

*What is the SM EFT? A linear realization of gauge symmetry and the new state is a  $0^+$  scalar:*

*Four fermion operators with leptons and quark fields:*

8 : $(\bar{L}L)(\bar{L}L)$		8 : $(\bar{R}R)(\bar{R}R)$		8 : $(\bar{L}L)(\bar{R}R)$	
$Q_{ll}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$Q_{ee}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$Q_{le}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{uu}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{lu}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{dd}$	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{ld}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{eu}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{qe}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{ed}$	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$

8 : $(\bar{L}R)(\bar{R}L) + \text{h.c.}$		8 : $(\bar{L}R)(\bar{L}R) + \text{h.c.}$	
$Q_{ledq}$	$(\bar{l}_p^j e_r)(\bar{d}_s^k q_{tj})$	$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \epsilon_{jk} (\bar{q}_s^k d_t)$
		$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \epsilon_{jk} (\bar{q}_s^k T^A d_t)$
		$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \epsilon_{jk} (\bar{q}_s^k u_t)$
		$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \epsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$

- ➔ 59 operators or 2499 parameters experimentally to constraint!
- ➔ where nearly 50% of the parameters (1053) are sensitive to **lepton-quark interactions** – not just about lepto-quarks

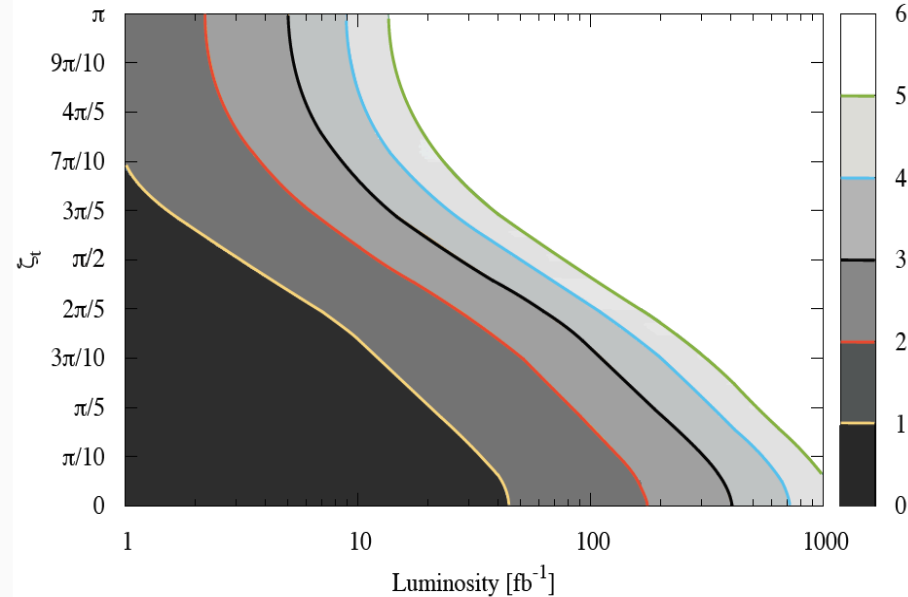
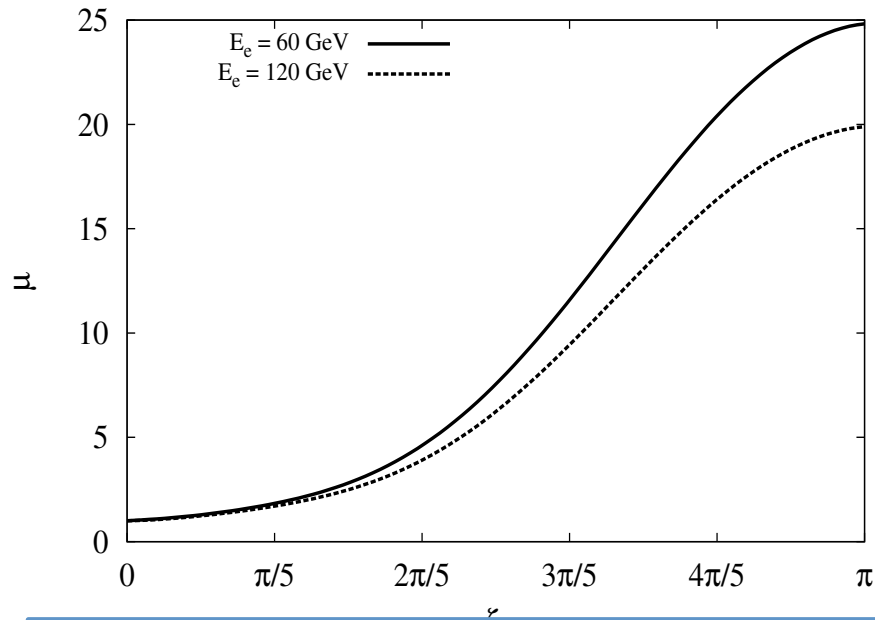
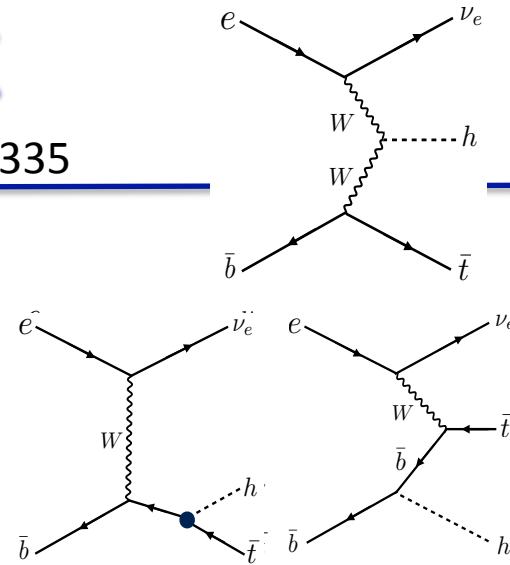
# Top Yukawa Coupling @ LHeC

B.Coleppa, M.Kumar, S.Kumar, B.Mellado, Phys. Lett. B770 (2017) 335

Introduce phase dependent top Yukawa coupling

$$\mathcal{L} = -i \frac{m_t}{v} \bar{t} [\cos \zeta_t + i \gamma_5 \sin \zeta_t] t h$$

Enhancement of the cross-section as a function of phase



Observe/Exclude non-zero phase to better than  $4\sigma$   $\rightarrow$  With Zeor Phase: Measure coupling with 17% accuracy  $\rightarrow$  work ongoing on FCC-eh prospects

# Double Higgs Production at FCC-eh

“Probing anomalous couplings using di-Higgs production in electron-proton collisions” by Mukesh Kumar, Xifeng Ruan, Rashidul Islam, Alan S. Cornell, Max Klein, Uta Klein, Bruce Mellado, *Physics Letters B* 764 (2017) 247-253 [arXiv:1509.04016]

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{hhh}^{(3)} + \mathcal{L}_{hWW}^{(3)} + \mathcal{L}_{hhWW}^{(4)}.$$

FCC-eh  
SM(P=-0.8)  
 $\sigma(\text{HH})=430$  ab  
in VBF!

$$\mathcal{L}_{hhh}^{(3)} = \frac{m_h^2}{2v} (1 - \overset{\text{SM}}{g_{hhh}^{(1)}}) h^3 + \frac{1}{2v} g_{hhh}^{(2)} h \partial_\mu h \partial^\mu h, \quad (2)$$

$$\mathcal{L}_{hWW}^{(3)} = -g \left[ \frac{g_{hWW}^{(1)}}{2m_W} W^{\mu\nu} W_{\mu\nu}^\dagger h + \frac{g_{hWW}^{(2)}}{m_W} (W^\nu \partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.}) + \frac{\tilde{g}_{hWW}}{2m_W} W^{\mu\nu} \tilde{W}_{\mu\nu}^\dagger h \right], \quad (3)$$

$$\mathcal{L}_{hhWW}^{(4)} = -g^2 \left[ \frac{g_{hhWW}^{(1)}}{4m_W^2} W^{\mu\nu} W_{\mu\nu}^\dagger h^2 + \frac{g_{hhWW}^{(2)}}{2m_W^2} (W^\nu \partial^\mu W_{\mu\nu}^\dagger h^2 + \text{h.c.}) + \frac{\tilde{g}_{hhWW}}{4m_W^2} W^{\mu\nu} \tilde{W}_{\mu\nu}^\dagger h^2 \right]. \quad (4)$$

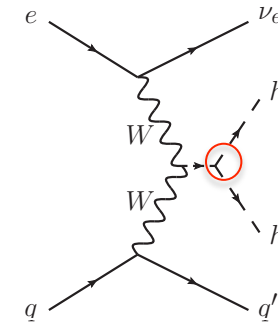
→ All other g coefficients are anomalous couplings to the hhh, hWW and hhWW anomalous vertices → those are 0 in SM

# Effective Vertices

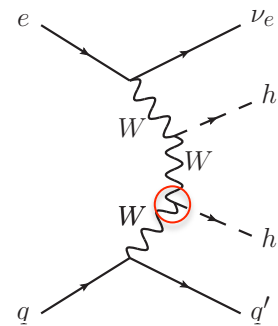
$$\Gamma_{hhh} = -6\lambda v \left[ g_{hhh}^{(1)} + \frac{g_{hhh}^{(2)}}{3m_h^2} (p_1 \cdot p_2 + p_2 \cdot p_3 + p_3 \cdot p_1) \right], \quad (6)$$

$$\Gamma_{hW^-W^+} = gm_W \left[ \left\{ 1 + \frac{g_{hWW}^{(1)}}{m_W^2} p_2 \cdot p_3 + \frac{g_{hWW}^{(2)}}{m_W^2} (p_2^2 + p_3^2) \right\} \eta^{\mu_2\mu_3} \right. \\ \left. - \frac{g_{hWW}^{(1)}}{m_W^2} p_2^{\mu_3} p_3^{\mu_2} - \frac{g_{hWW}^{(2)}}{m_W^2} (p_2^{\mu_2} p_2^{\mu_3} + p_3^{\mu_2} p_3^{\mu_3}) \right. \\ \left. - i \frac{\tilde{g}_{hWW}}{m_W^2} \epsilon_{\mu_2\mu_3\mu\nu} p_2^\mu p_3^\nu \right], \quad (7)$$

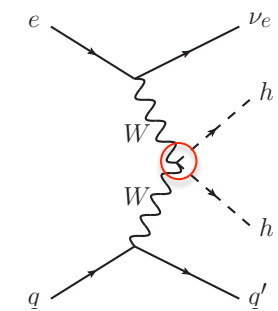
$$\Gamma_{hhW^-W^+} = g^2 \left[ \left\{ \frac{1}{2} + \frac{g_{hhWW}^{(1)}}{m_W^2} p_3 \cdot p_4 + \frac{g_{hhWW}^{(2)}}{m_W^2} (p_3^2 + p_4^2) \right\} \eta^{\mu_3\mu_4} \right. \\ \left. - \frac{g_{hhWW}^{(1)}}{m_W^2} p_3^{\mu_4} p_4^{\mu_3} - \frac{g_{hhWW}^{(2)}}{m_W^2} (p_3^{\mu_3} p_3^{\mu_4} + p_4^{\mu_3} p_4^{\mu_4}) \right. \\ \left. - i \frac{\tilde{g}_{hhWW}}{m_W^2} \epsilon_{\mu_3\mu_4\mu\nu} p_3^\mu p_4^\nu \right]. \quad (8)$$



1,2,3=  
h,h,h



1,2,3 =  
h,W-,W+



1,2,3,4 =  
h,h,W-,W+

**Note the dependence on momenta in non-SM vertices. This induces significant impact on scattering kinematics.**

# Event Selection using $h \rightarrow bb$

$p_e = -0.8$ , Anti-kt jets  $R=0.4$ ,  $E_{\text{miss}} > 40$  GeV,  $\eta(\text{fwd jet}) > 5$ ,

$90 < m_{bb}(1)$ ,  $m_{bb}(2) < 125$  GeV,  $m(4b) > 290$  GeV

b-tagging for  $|\eta| < 5$  assumed to be 70% with misidentifications of 10% for charm and 1% for light quarks /gluons

Delphes detector-level

Cuts / Samples	Signal	4b+jets	2b+jets	Top	ZZ	$b\bar{b}H$	ZH	Total Bkg	Significance
Initial	$2.00 \times 10^3$	$3.21 \times 10^7$	$2.32 \times 10^9$	$7.42 \times 10^6$	$7.70 \times 10^3$	$1.94 \times 10^4$	$6.97 \times 10^3$	$2.36 \times 10^9$	0.04
At least 4b + 1j	$3.11 \times 10^2$	$7.08 \times 10^4$	$2.56 \times 10^4$	$9.87 \times 10^3$	$7.00 \times 10^2$	$6.32 \times 10^2$	$7.23 \times 10^2$	$1.08 \times 10^5$	0.94
Lepton rejection $p_T^\ell > 10$ GeV	$3.11 \times 10^2$	$5.95 \times 10^4$	$9.94 \times 10^3$	$6.44 \times 10^3$	$6.92 \times 10^2$	$2.26 \times 10^2$	$7.16 \times 10^2$	$7.75 \times 10^4$	1.12
Forward jet $\eta_J > 4.0$	233	13007.30	2151.15	307.67	381.04	46.82	503.22	16397.19	1.82
$E_T > 40$ GeV	155	963.20	129.38	85.81	342.18	19.11	388.25	1927.93	3.48
$\Delta\phi_{E_T j} > 0.4$	133	439.79	61.80	63.99	287.10	14.53	337.14	1204.35	3.76
$m_{bb}^1 \in [95, 125]$ , $m_{bb}^2 \in [90, 125]$	54.5	28.69	5.89	6.68	5.14	1.42	17.41	65.23	6.04
$m_{4b} > 290$ GeV	49.2	10.98	1.74	2.90	1.39	1.21	11.01	29.23	7.51

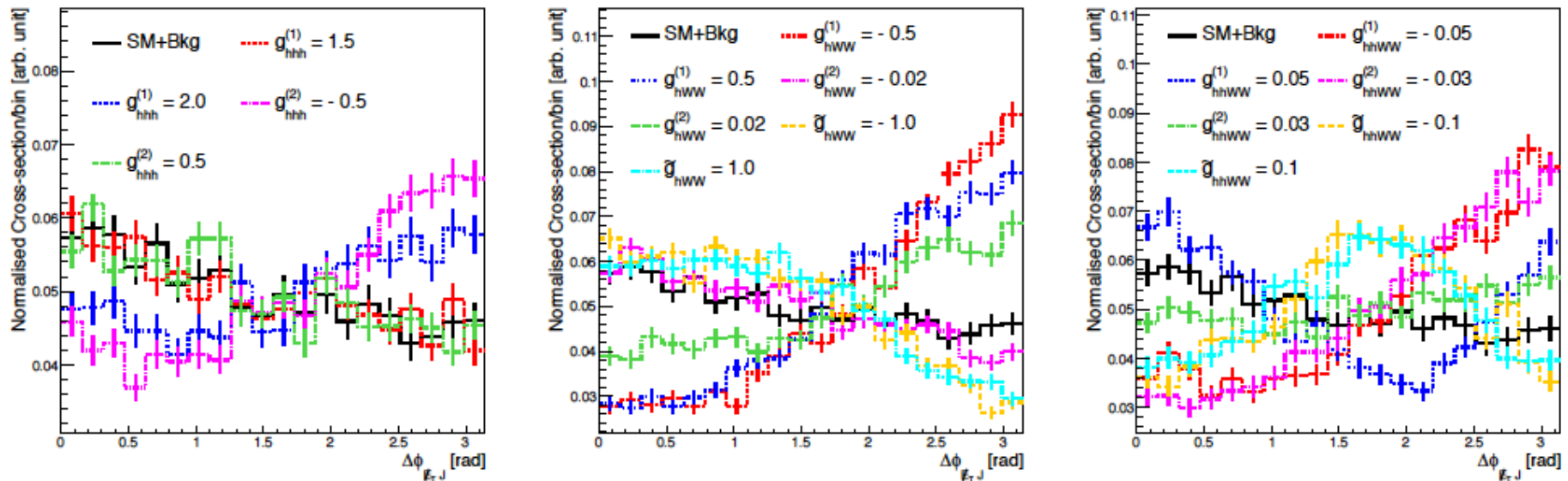
Table 2: A summary table of event selections to optimise the signal with respect to the backgrounds in terms of the weights at  $10 \text{ ab}^{-1}$ . In the first column the selection criteria are given as described in the text. The second column contains the weights of the signal process  $p e^- \rightarrow hh j \nu_e$ , where both the Higgs bosons decay to  $b\bar{b}$  pair. In the next columns the sum of weights of all individual prominent backgrounds in charged current, neutral current and photo-production are given with each selection, whereas in the penultimate column all backgrounds' weights are added. The significance is calculated at each stage of the optimised selection criteria using the formula  $S = \sqrt{2[(S+B)\log(1+S/B) - S]}$ , where  $S$  and  $B$  are the expected signal and background yields at a luminosity of  $10 \text{ ab}^{-1}$  respectively. This optimisation has been performed for  $E_e = 60$  GeV and  $E_n = 50$  TeV.

$$S = \sqrt{2[(S+B)\log(1+S/B) - S]}$$

# Azimuthal Angle Distributions

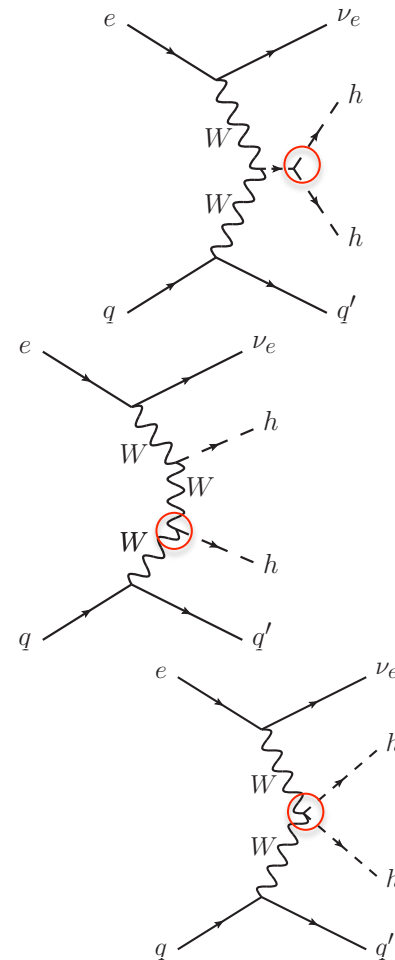
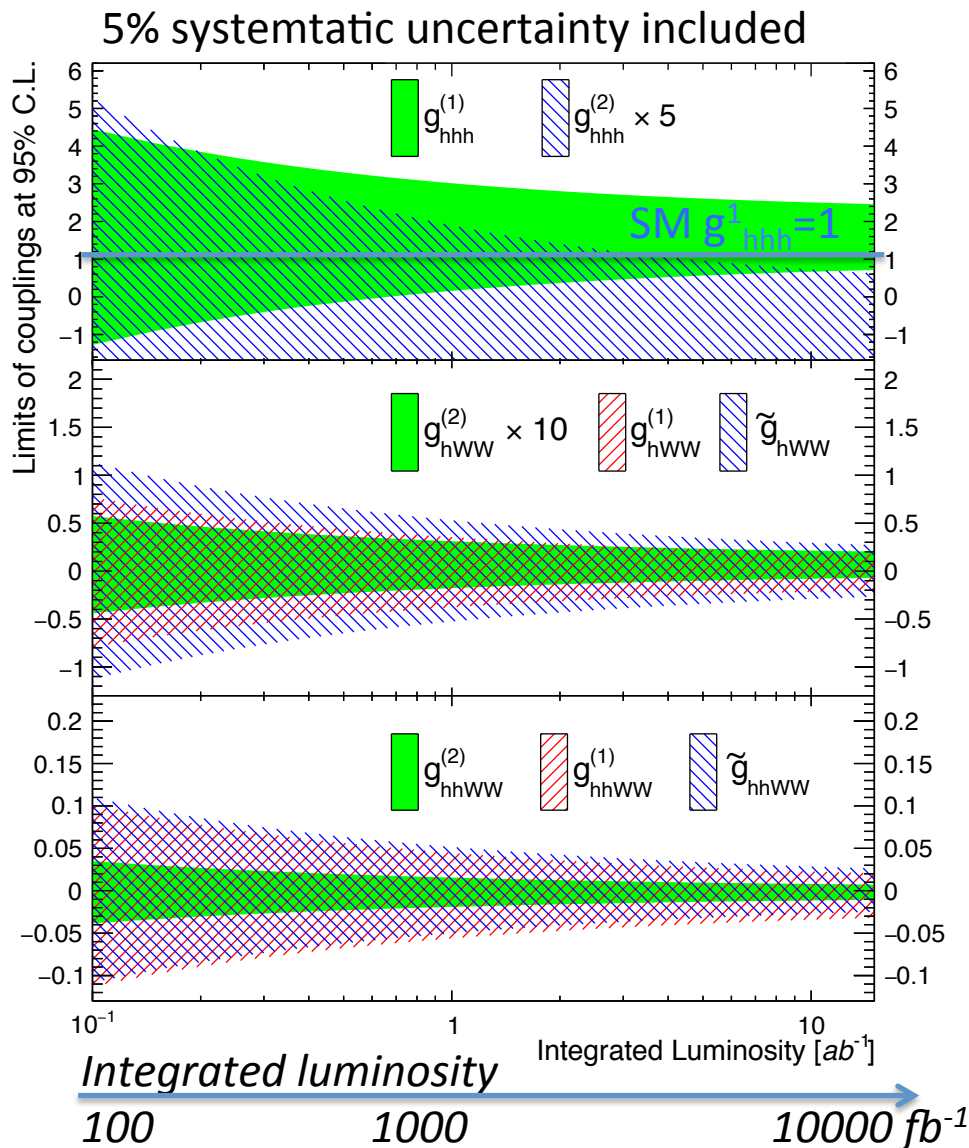
between missing transverse energy and forward jet, at Delphes detector-level, including background :  $bbbbj$ ,  $bbjjj$ ,  $Z(bb)h(bb)j$ ,  $ttj$ ,  $h(bb)bbj$

→ For signal, we consider  $hh \rightarrow bbbb$  decays motivated by  $h \rightarrow bb$  studies.



- normalised DIS cross sections are sensitive to non-BSM vertices
- initial study published for this novel variable
- potential for a deeper analysis and interpretation

# 95% C.L. Exclusion Limits from $\sigma_{\text{fiducial}}$



$1\sigma$  for SM  $hhh$  for  $E_e$   
60 (120) GeV and  $10ab^{-1}$

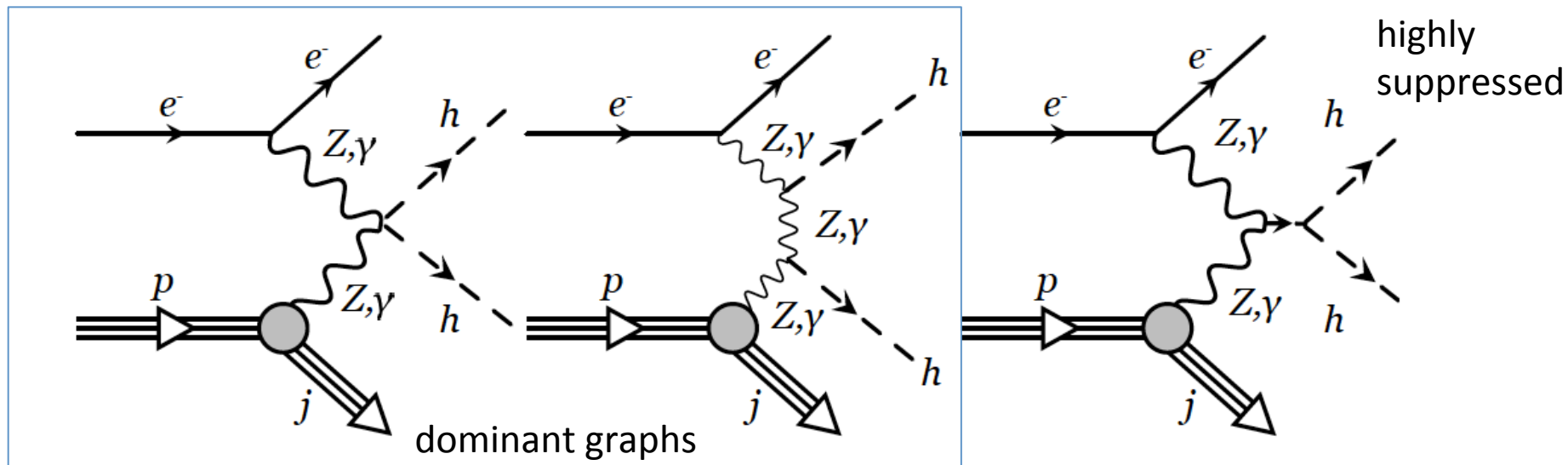
$$g_{hhh}^{(1)} = 1.00^{+0.24(0.14)}_{-0.17(0.12)}$$

Probing anomalous couplings: limits are obtained by scanning one of the non-BSM coupling while keeping other couplings to their SM values.

Here  $g_{(\dots)}^{(i)}$ ,  $i = 1, 2$ , and  $\tilde{g}_{(\dots)}$  are real coefficients corresponding to the CP-even and CP-odd couplings respectively, of the  $hhh$ ,  $hWW$  and  $hhWW$  anomalous vertices.



S. Kuday,<sup>\*</sup> H. Saygı,<sup>†</sup> İ. Hoş,<sup>‡</sup> and F. Çetin<sup>§</sup>



- Vertices for Neutral Current DIS ( $Z, \gamma$ ) and Photoproduction ( $\gamma$ ) studied in Higgs Effective Lagrangian Model : parametrise  $hhZZ$  and  $hh\gamma\gamma$  in 4-point interactions in terms of CP-even and CP-odd Wilson coefficients (and Higgs self coupling and Yukawa coupling)
- Study at Delphes-detector level (FCC-hh) azimuthal dependencies between scattered lepton and forward jet
- $hh$  : 4b final states investigated using a very first version of FCC-hh detector
- Promising sensitivity found while scanning parameter space for Wilson coefficients

# Invisible Higgs@LHeC

relating the Higgs and the 'dark' sectors

Y.-L. Tang et al.,  
arXiv: 1508.01095

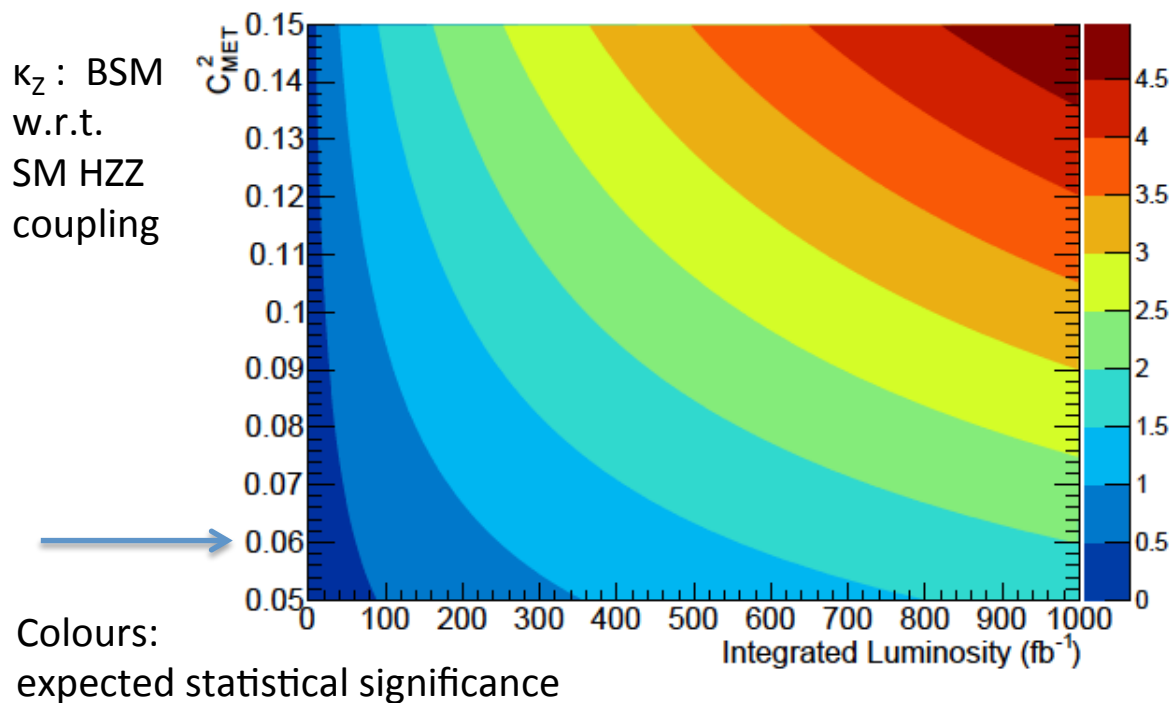
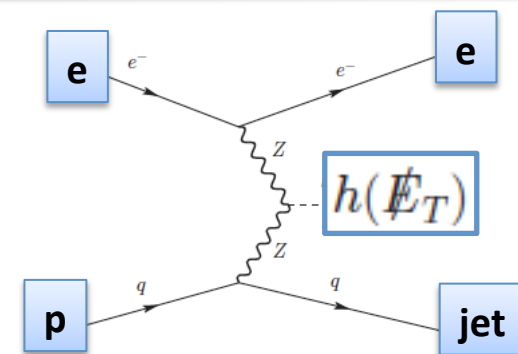
HL-LHC @ 3 ab<sup>-1</sup> [arXiv:1411. 7699]

Br( $h \rightarrow \cancel{E}_T$ ) < 3.5% @95% C.L., MVA based

For LHeC, assume : 1ab<sup>-1</sup>, P<sub>e</sub>=-0.9, cut based

Br( $h \rightarrow \cancel{E}_T$ ) < 6% @ 95 % C.L.

$$C_{\text{MET}}^2 = \kappa_Z^2 \times \text{Br}(h \rightarrow \cancel{E}_T)$$

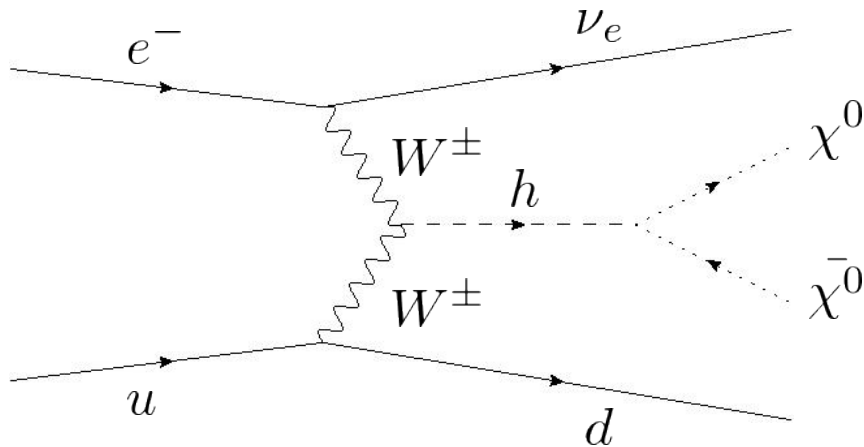


➔ potential much enhanced for FCC-eh @ 3.5 TeV and HE-LHC-eh @ 1.8 TeV  
➔ NEW studies performed on Delphes detector-level using our Madevent framework

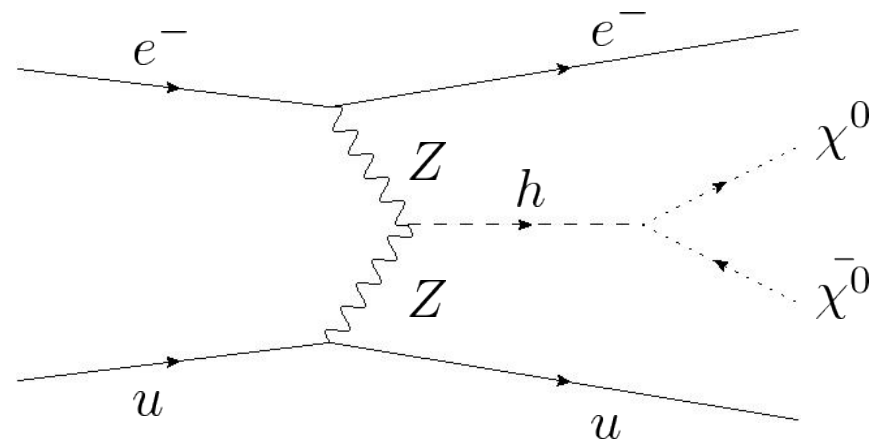
# Invisible Higgs Decay in ep

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Masahiro Kuze  
Tokyo Tech

CC production of an invisible Higgs



NC production of an invisible Higgs



- We focus currently on NC DIS channel: employ that kinematic is over constrained using jet and electron information in the final state
- We use the idea from C. Zhang and Y.-L. Tang : We emulate Higgs to invisible by assuming a branching of 100% for  $H \rightarrow ZZ \rightarrow 4\nu$
- We started to study signals and backgrounds using CMS-style and FCC-eh-style 'Delphes' detectors, using same analysis strategies as developed for LHeC (C. Zhang and BSc thesis S. Kawaguchi)

# Dominant Background

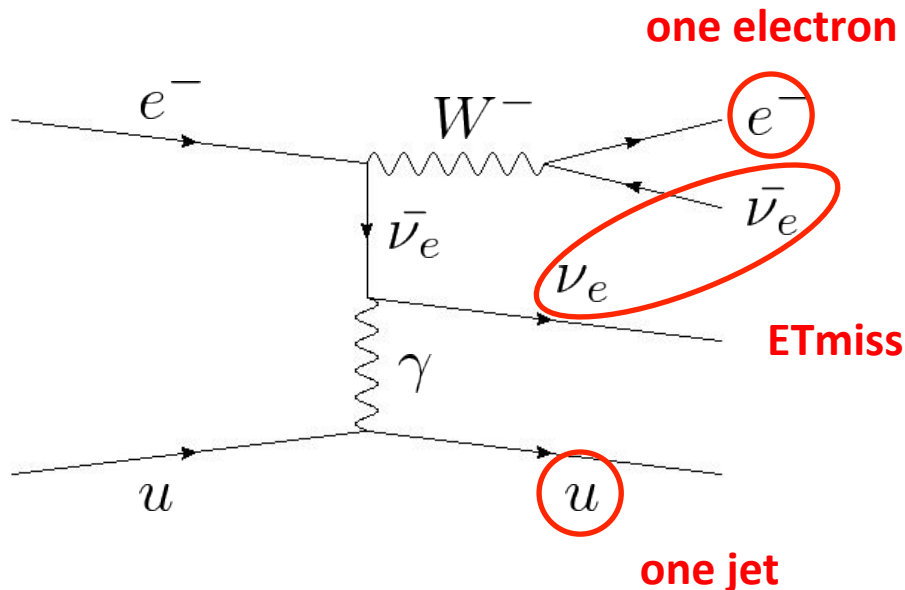
for faking our signal feature:  
one electron, one jet, and missing transverse energy ( $E_T^{\text{miss}}$ )

$W^+j e^-$  and  $W^-j e^-$  backgrounds and

$Wj\nu$  background

$$p + e^- \rightarrow W^- + j + \nu_e$$

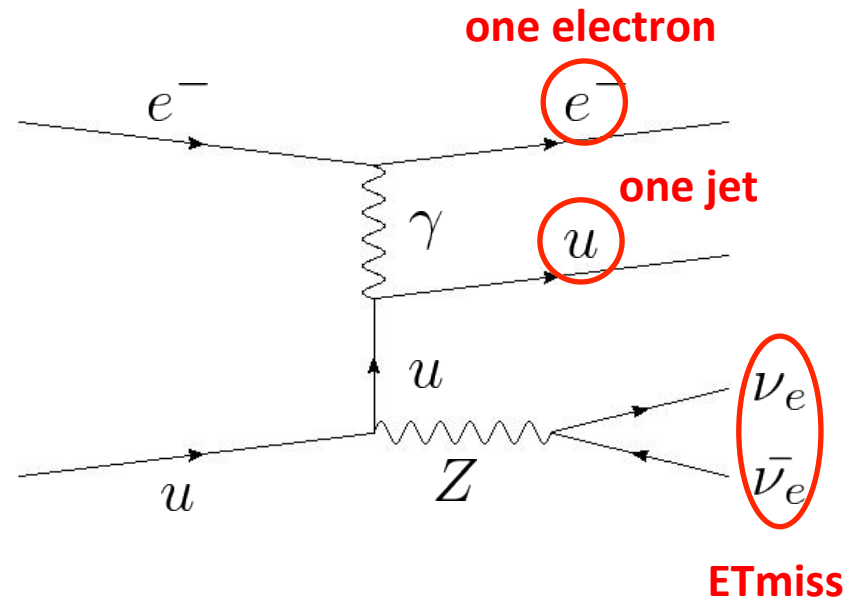
$$W^- \rightarrow e^- + \nu_e$$



$Zje$  background

$$p + e^- \rightarrow Z + j + e^-$$

$$Z \rightarrow \nu + \bar{\nu}$$

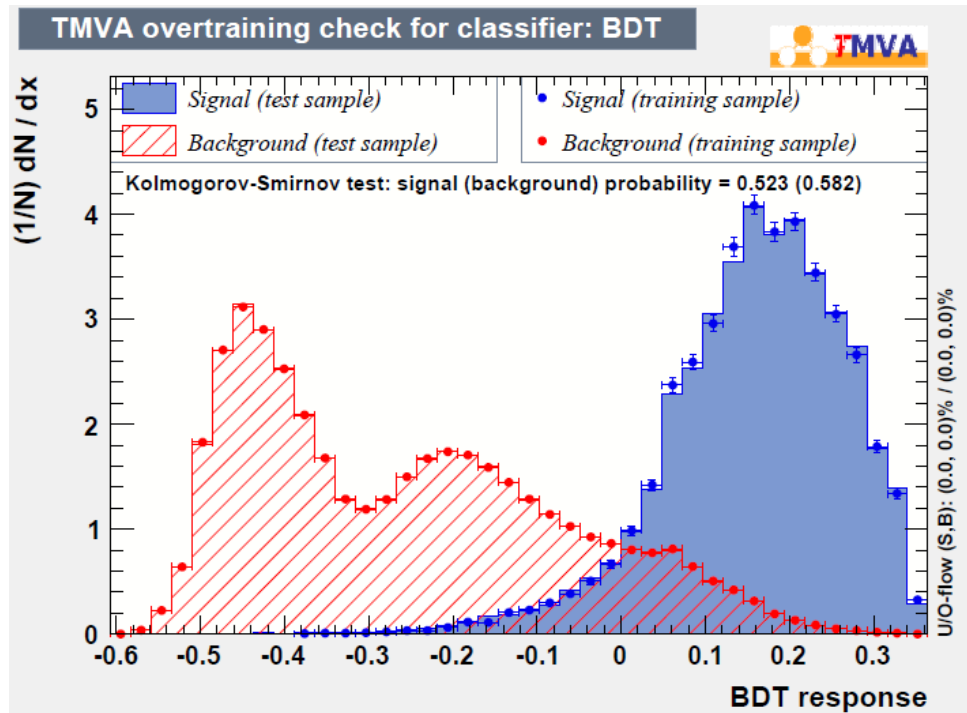


# Results for FCC-eh - Using BDT

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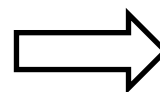
MVA (BDT) using samples with 1 jet and 1e- with high pT, and other variables as a BDT

BDT >	Signal	Background	Z[%]
0	31961	267904	3.25
0.05	29932	176439	2.81
0.1	25686	94138	2.40
0.15	19898	42439	2.08
<b>0.2</b>	<b>13020</b>	<b>15562</b>	<b>1.93</b>
0.25	6998	4969	2.04
0.3	2320	1003	2.82



Branching ratio calculated by  $S/\sqrt{S+B}$  :

$$Z = 13020 \times \text{Br}(h \rightarrow \mu \tau) / \sqrt{13020 \times \text{Br}(h \rightarrow \mu \tau) + 15562}$$



In the case of  $2\sigma$

$$\text{Br}(h \rightarrow \mu \tau) \sim 1.93\%$$

# Branching for invisible Higgs

Values given in case of  $2\sigma$

Delphes detectors	LHeC	DLHeC	FCC-eh
	1.3 TeV	1.8 TeV	3.5 TeV
LHC-style	4.7%	3.2%	1.9%
First 'ep-style'	5.7%		2.6%

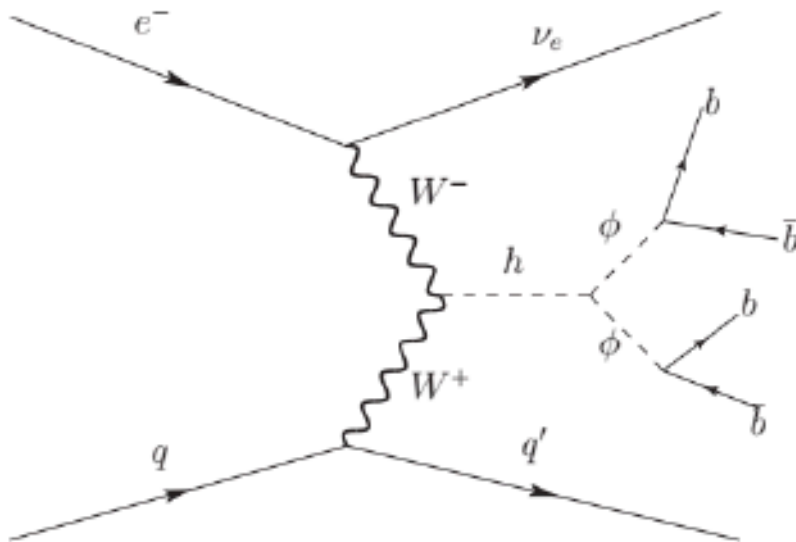
- ✓ **Results look very encouraging for a measurement of the branching of Higgs to invisible in ep down to 2%.**
- ✓ For 2 different detector options we get similar results.
- Certainly : we will use this channel to further optimize analysis strategies (used methods and requirements, e.g. size of jets and electron reconstructions) and to modify our ep-detector
- **employ synergies within FCC study group → detector and analysis details (BDT optimisation) has certainly an impact on results**

# Exotic Higgs Decays

$$h \rightarrow \phi\phi \rightarrow 4b$$

$\phi$ : a spin-0 particle from new physics.

$$eq \rightarrow \nu_e h q' \rightarrow \nu_e \phi\phi q' \rightarrow \nu_e b\bar{b}b\bar{b}q'$$



$$\mathcal{L}_{eff} = \lambda_h v h \phi^2 + \lambda_b \phi \bar{b} b + \mathcal{L}_{\phi \text{ decay, other}}$$

S. Liu, Y. L. Tang, C. Zhang, S. Zhu, 1608.08458

- Well motivated signature in extended Higgs sector.
- Difficult to probe at hadron colliders.
- LHeC signal: here using CC channel.
- Backgrounds: CC multijet, CC  $t/h/W/Z$ +jets, PHP multijet.
- PHP backgrounds assumed to be negligible after MET requirements and electron tagging.
- Current analysis is done at parton level.

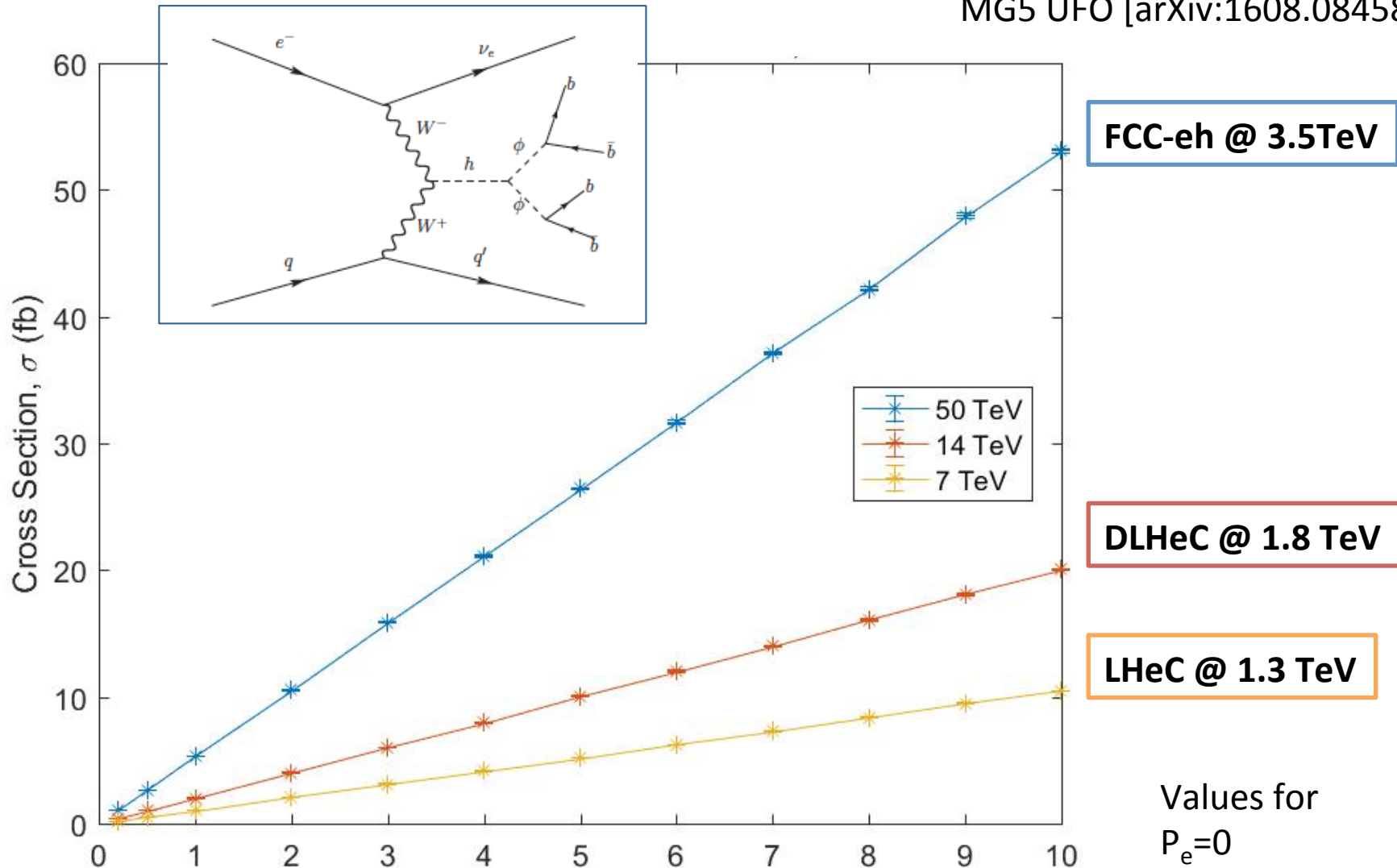
$$C_{4b}^2 = \kappa_V^2 \times \text{Br}(h \rightarrow \phi\phi) \times \text{Br}^2(\phi \rightarrow b\bar{b})$$

@LHeC: 95% C.L. for  $m_\phi$  of 20, 40, 60 GeV is 0.3%, 0.2% and 0.1% for  $C_{4b}^2$

# Exotic Higgs@FCC-eh

Uta Klein  
Michael o'Keefe  
Liverpool

MG5 UFO [arXiv:1608.08458]



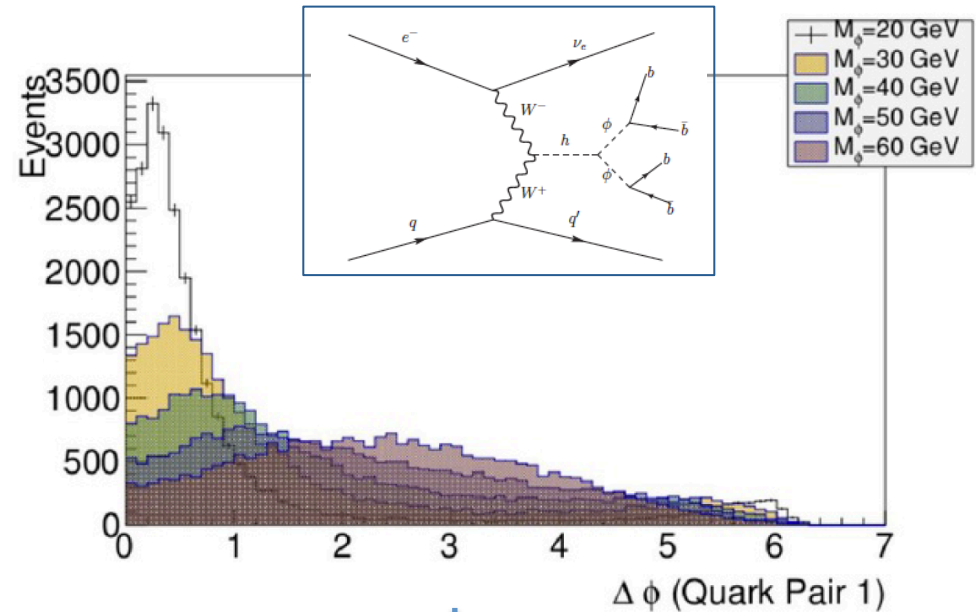
→ reflecting coupling of new scalar to 125 GeV higgs



# Kinematics @ Quark-Level

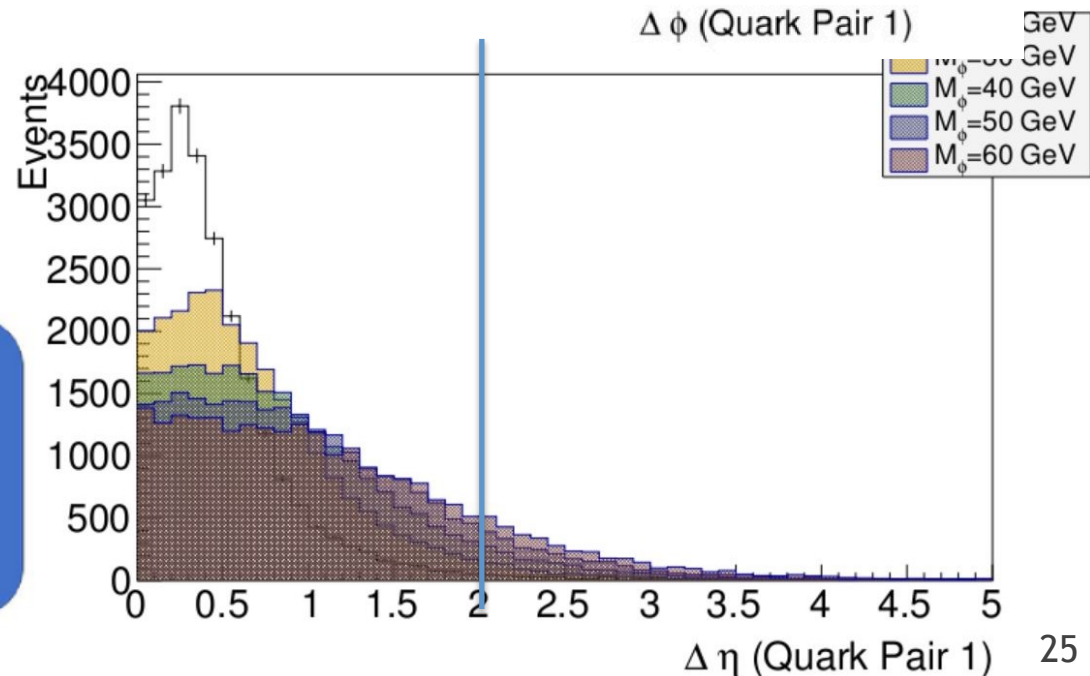
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Michael o'Keefe  
Liverpool

$\Delta\phi$  between b quarks in the scalar (parton level)



$\Delta\eta$  between b quarks in the scalar (parton level)

→ use  $\Delta\eta < 2$  for finding two scalars with mass within  $2m_b$  and  $m_H/2$  looping over N jets minimising  $\Delta m$



Jet A,  $1 < A < N$

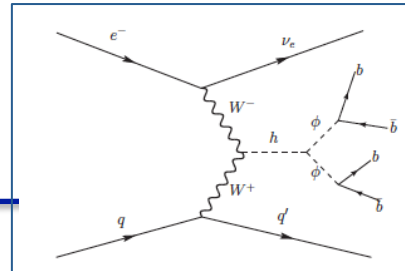
Jet B,  $A+1 < B < N$

Jet C,  $B+1 < C < N$

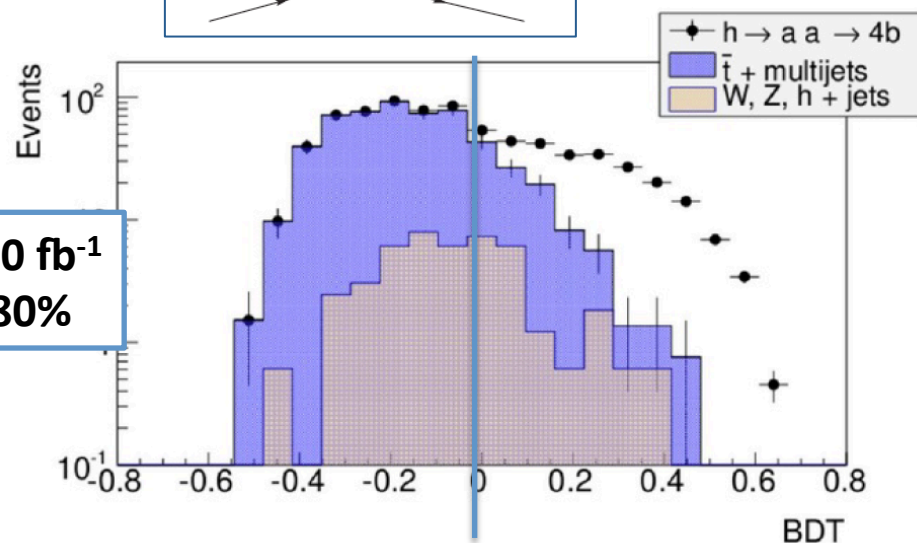
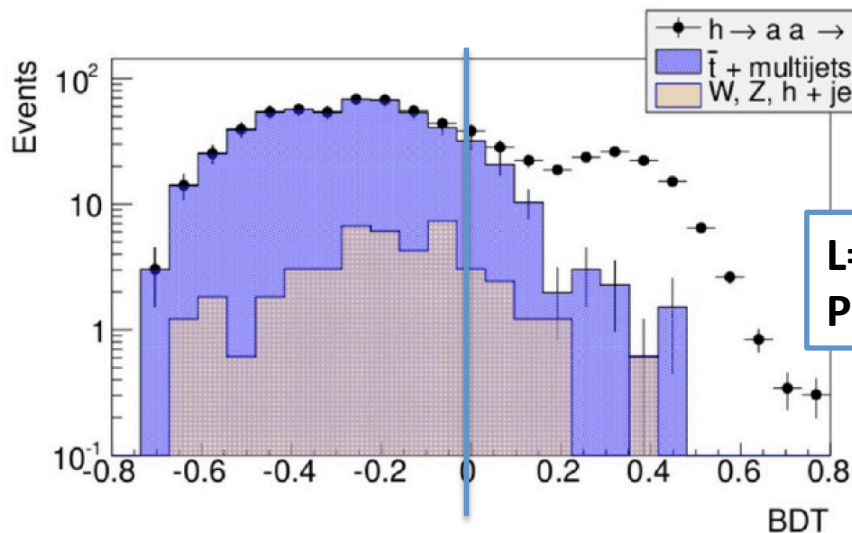
Jet D,  $C+1 < D < N$

# BDT Analysis @ BR=10%

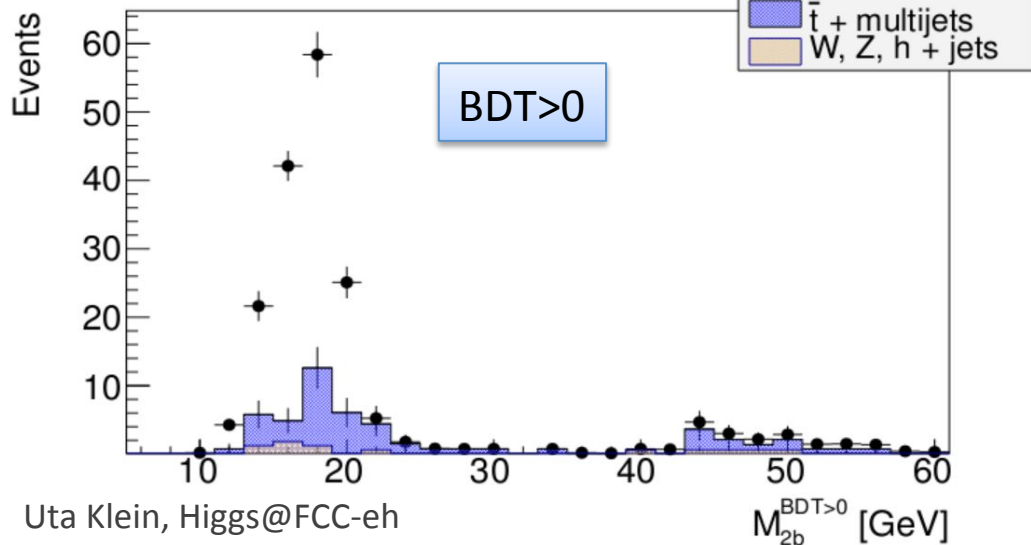
Delphes-detector level with b-tag  $|\eta| < 2.5$



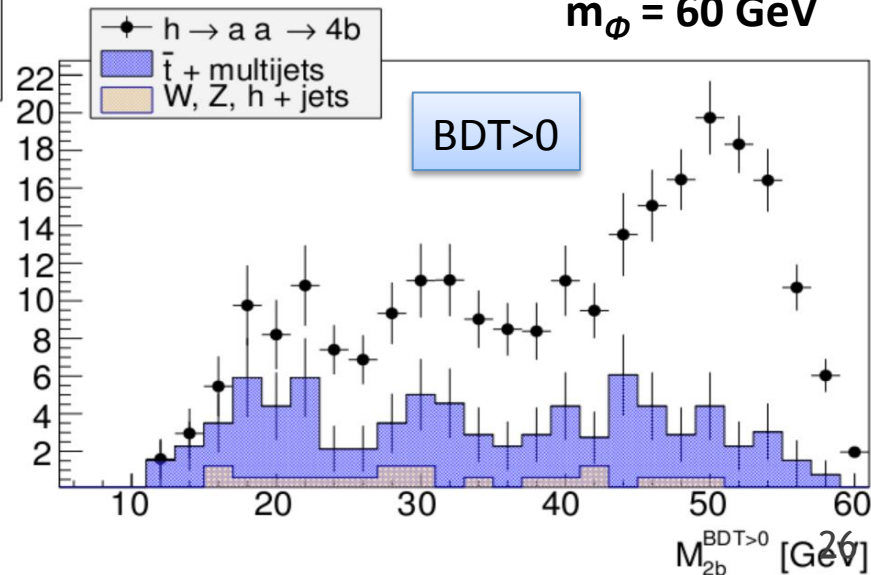
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Michael o'Keefe  
Liverpool



$m_\phi = 20 \text{ GeV}$



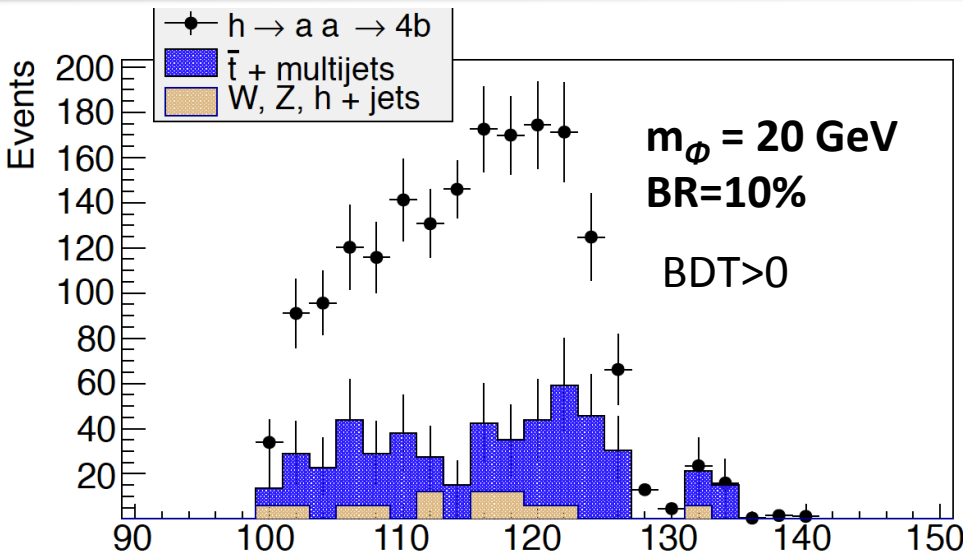
$m_\phi = 60 \text{ GeV}$



# First Results @ FCC-eh

$L=1 \text{ ab}^{-1}$   
 $P_e=-80\%$

Uta Klein  
 Michael o'Keefe  
 Liverpool



Very promising first results to discover an exotic Higgs decay into two new light scalars at FCC-eh down to a BR of 1% for  $1 \text{ ab}^{-1}$ . A BR of 10% could be discovered within 1 year ( $100 \text{ fb}^{-1}$ ).

Values for  $BDT>0$

BR (%)	$M_\phi$ (GeV)						$Z = \sqrt{2 \left[ (S+B) \ln \left( 1 + \frac{S}{B} \right) - S \right]}$
	$\sigma$ (fb)	$\Delta\sigma$ (fb)	Z	$\sigma$ (fb)	$\Delta\sigma$ (fb)	Z	
0.2	0.03	0.02	1.14	0.03	0.03	1.17	
0.4	0.05	0.02	2.27	0.07	0.03	2.33	
0.6	0.08	0.02	3.37	0.10	0.03	3.47	
0.8	0.10	0.02	4.46	0.13	0.03	4.59	
1	0.13	0.03	5.54	0.17	0.03	5.71	

# More New Studies Ongoing

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... and publications in preparation

- “Search for Anomalous HVV couplings at the LHeC and the FCC-ep” by M. Altinli et al.
- “Probing FCNC couplings of Higgs-top at FCC-ep and LHeC” by B. Hacinahinoglu et al.
- “Searching for doubly-charged Higgs bosons in the Georgi-Machacek model at ep colliders “ by H. Sun et al. (see also presentation at DIS2017)

# Conclusions

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- We just started to explore the potential of complementary SM and BSM Higgs searches using concurrent ep collisions at FCC-pp in particular for HH and  $H\phi$  couplings  
→ many more studies ongoing (e.g. anomalous htt coupling) and possible! You are welcome!
- Enhance ep potential further by strengthen analysis techniques and detector developments between p, ep and ee : extended beauty and charm tagging using BDT, jet-substructure, boosted pairs ...
- For the FCC CDR : Quantify the joint potential → combined analysis of pp, ep and ee cross sections to constrain SM/BSM physics scenario's and to design the most powerful and sustainable search complex at the energy frontier.

# Additional Sources & Thanks to

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The LHeC/FCC-eh study group, <http://cern.ch/lhec>.

“On the Relation of the LHeC and the LHC” [arXiv:1211.5102]

Poetic 2016 Workshop, 14.-18.11.2016, Temple University (USA)

[https://phys.cst.temple.edu/poetic-cteq-2016/scientific\\_program.html](https://phys.cst.temple.edu/poetic-cteq-2016/scientific_program.html)

1<sup>st</sup> FCC Physics Workshop, 16.1.-20.1.2017, CERN

<https://indico.cern.ch/event/550509/>

→ see M. Benedikt’s and F. Zimmermann’s and further eh talks given at this workshop

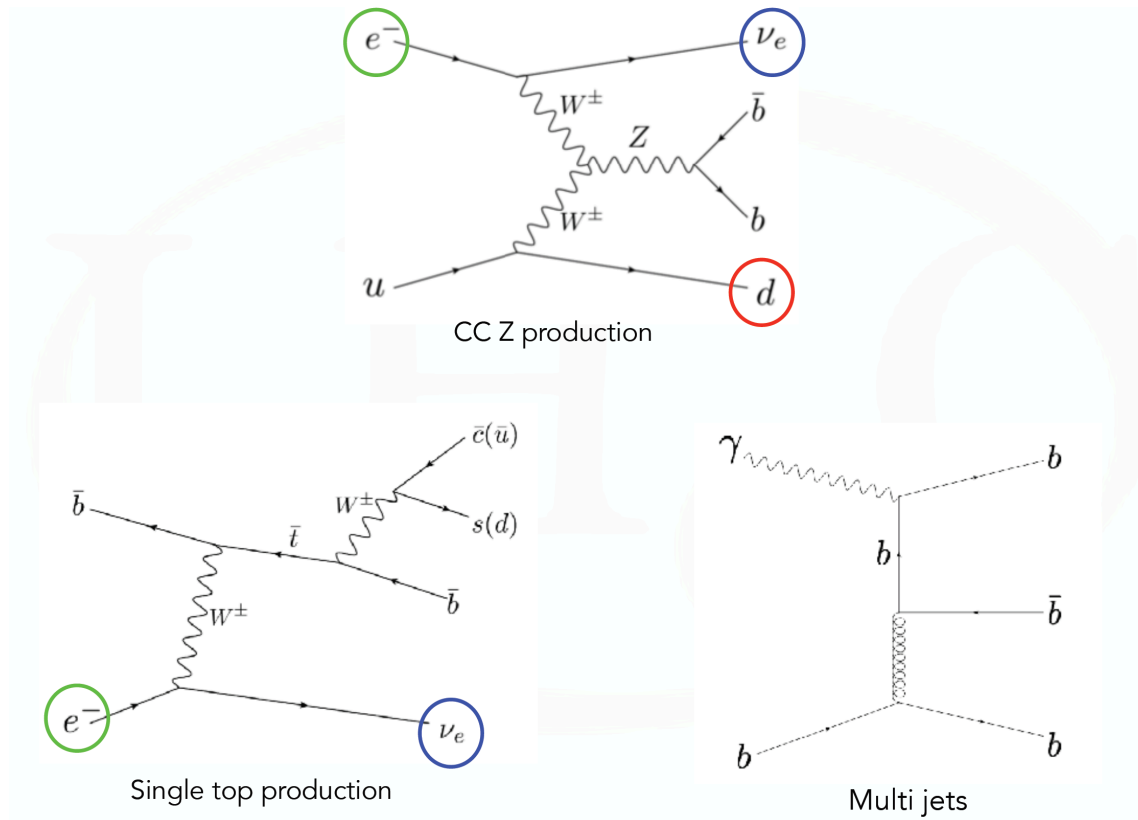
Special thanks to my colleagues in the LHeC/FCC-eh Higgs group, the project leader Max Klein, our detector expert Peter Kostka, and our bi-weekly Higgs-top working group discussions.

# Additional material

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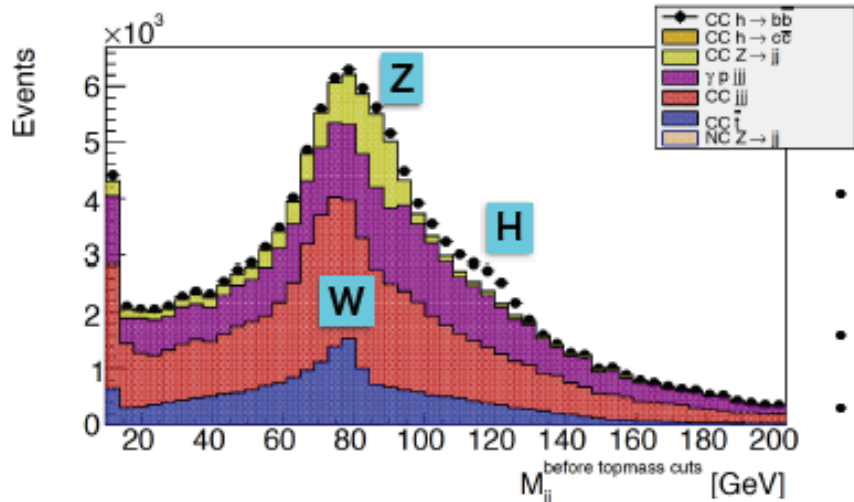
# SM Higgs Decay into b-quarks

- Typical background processes and assumptions about b-tagging for cut-based study

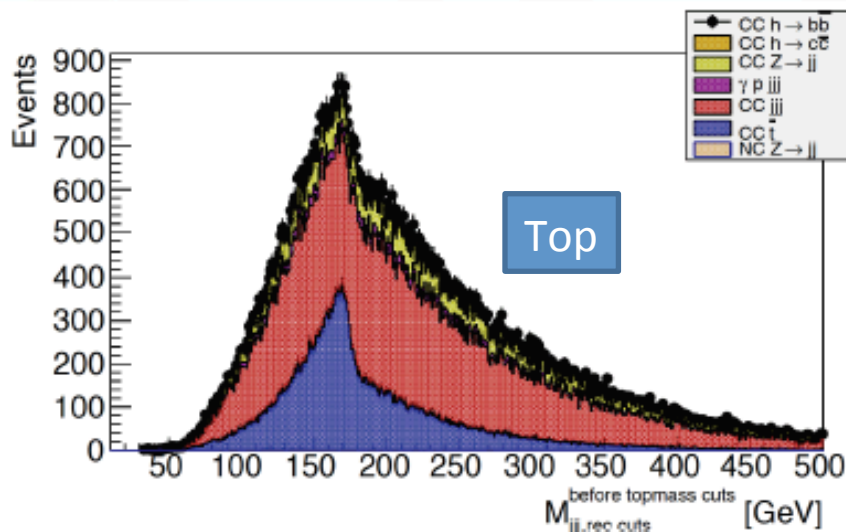




# Higgs w/o HFL tagging



- Invariant Di-jet Mass using 2 lowest eta **un-tagged** jets.
- 100% photo production background.
- No anti top cut.



- Invariant Mass of 3 highest  $p_T$  **un-tagged** jets.
- 10% photo production background.

# Asymmetries

$$\mathcal{A}_{\Delta\phi_{E_T J}} = \frac{|A_{\Delta\phi > \pi/2}| - |A_{\Delta\phi < \pi/2}|}{|A_{\Delta\phi > \pi/2}| + |A_{\Delta\phi < \pi/2}|},$$

[1509.04016]

Samples		$\mathcal{A}_{\Delta\phi_{E_T J}}$	$\sigma(\text{fb})$
SM+Bkg		$0.277 \pm 0.088$	
$g_{hhh}^{(1)}$	= 1.5	$0.279 \pm 0.052$	0.18
	= 2.0	$0.350 \pm 0.053$	0.21
$g_{hhh}^{(2)}$	= -0.5	$0.381 \pm 0.050$	0.19
	= 0.5	$0.274 \pm 0.024$	0.74
$g_{hWW}^{(1)}$	= -0.5	$0.506 \pm 0.022$	0.88
	= 0.5	$0.493 \pm 0.020$	0.94
$g_{hWW}^{(2)}$	= -0.02	$0.257 \pm 0.025$	0.67
	= 0.02	$0.399 \pm 0.040$	0.33
$\tilde{g}_{hWW}$	= -1.0	$0.219 \pm 0.016$	1.53
	= 1.0	$0.228 \pm 0.016$	1.53
$g_{hhWW}^{(1)}$	= -0.05	$0.450 \pm 0.033$	0.52
	= 0.05	$0.254 \pm 0.029$	0.68
$g_{hhWW}^{(2)}$	= -0.03	$0.462 \pm 0.022$	1.22
	= 0.03	$0.333 \pm 0.018$	1.46
$\tilde{g}_{hhWW}$	= -0.1	$0.351 \pm 0.020$	1.60
	= 0.1	$0.345 \pm 0.020$	1.61

Table 3: Estimation of the asymmetry, defined in Eq. (9), and statistical error associated with the kinematic distributions in Fig. 2 at an integrated luminosity of  $10 \text{ ab}^{-1}$ . The cross section ( $\sigma$ ) for the corresponding coupling choice is given in the last column with same parameters as in Table 1.

# Wilson Coefficients in NC DIS

[1702.00185]

Mass Basis	$g_{hh\gamma\gamma}$	$\tilde{g}_{hh\gamma\gamma}, \tilde{g}_{hhzz}$	$g_{hhzz}^{(1)}, g_{hhzz}^{(2)}$	$g_{hhzz}^{(3)}$
Gauge Basis	$-\frac{4\bar{c}_\gamma g^2 s_W^2}{m_W^2}$	$\frac{g}{2m_W} \{ \tilde{g}_{h\gamma\gamma}, \tilde{g}_{hzz} \}$	$\frac{g}{2m_W} \{ g_{hzz}^{(1)}, g_{hzz}^{(2)} \}$	$\frac{g^2}{2c_W^2} [1 - 6\bar{c}_T - \bar{c}_H + 8\bar{c}_\gamma \frac{s_W^4}{c_W^2}]$

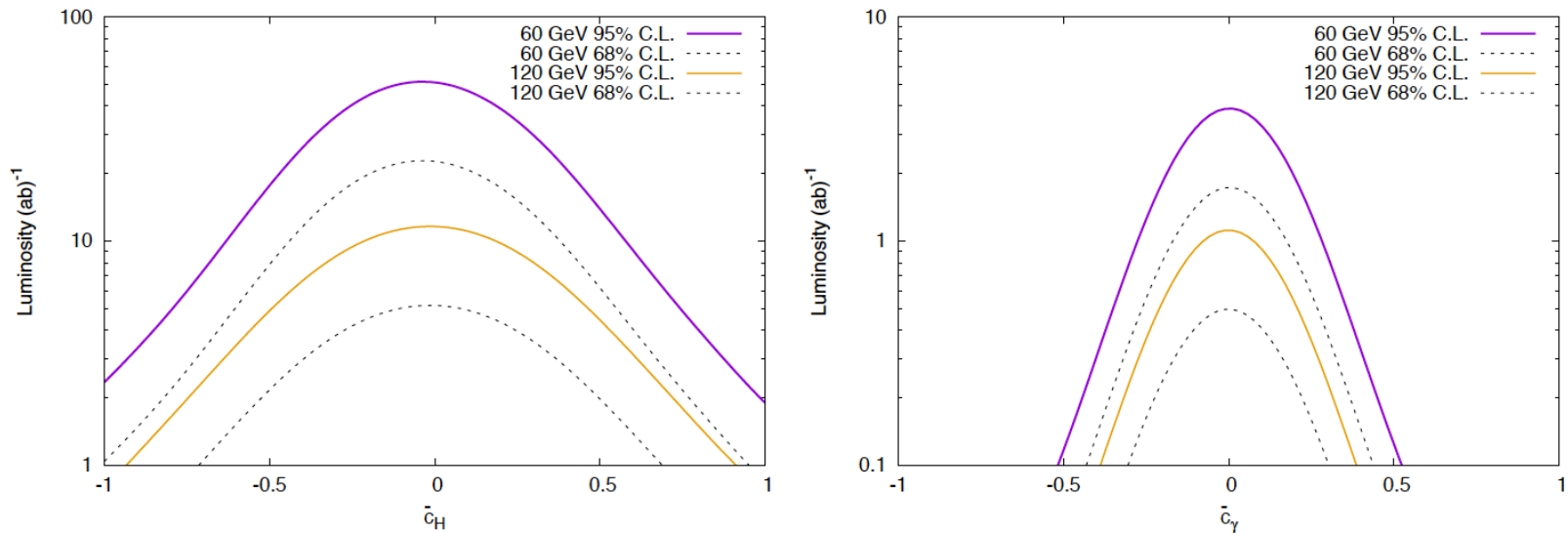


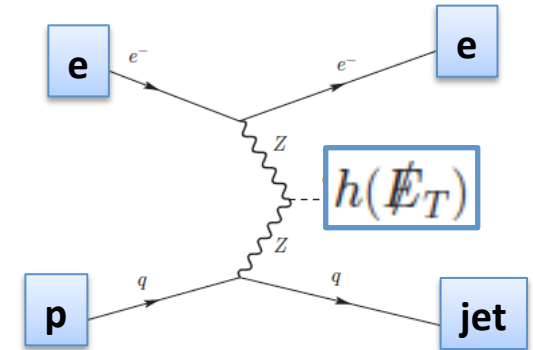
Figure 9: Required integrated luminosities to obtain corresponding Wilson coefficients for FCC-he energy options.

# Selection Requirements

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## Basic cuts (Cut 0)

- $N(\text{jets})$  for the jet and the electron
- $p_T$  for the leading jet and the leading electron
- for the leading jet and the leading electron
- for the leading jet and the leading electron



Cut 1 :  $|\Delta\phi_{\text{jet}, E_{\text{Tmiss}}}| > 1 \text{ rad}$

Cut 2 :  $E_{\text{Tmiss}} > 50 \text{ GeV}$

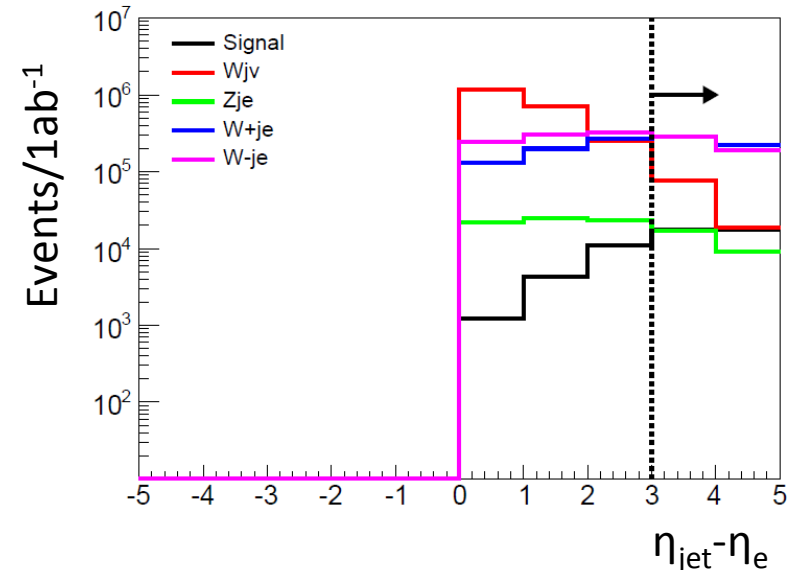
Cut 3 :  $\eta_{\text{jet}} - \eta_e > 3$

Cut 4 :  $\phi_{\text{jet}} - \phi_e < 2.4$

Cut 5 :  $-1.3 < \eta_e < 1.1$

Cut 6 :  $0.08 < \gamma_e < 0.55$

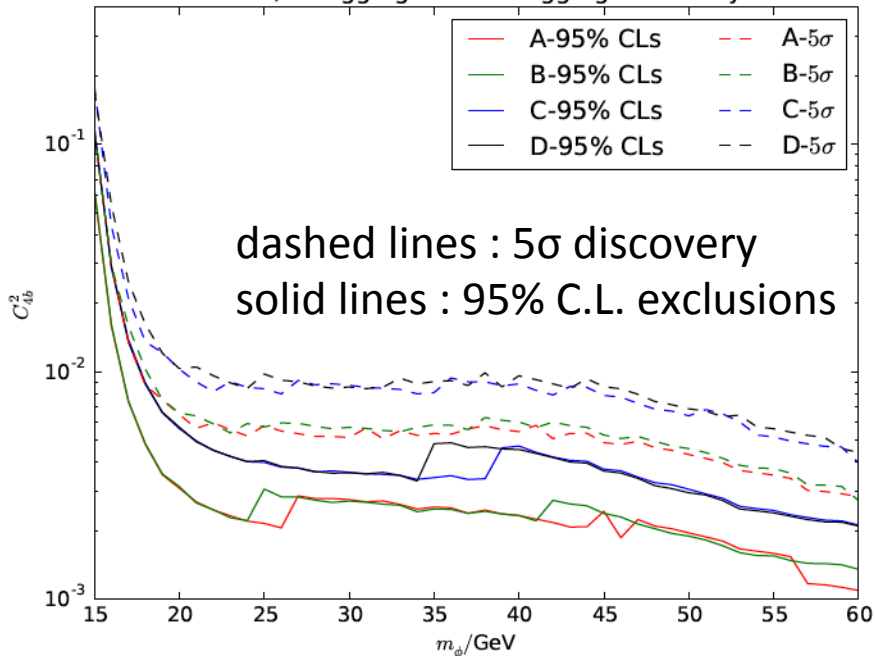
Cut 7 : require 1 electron, 1 jet,  
and veto tau's and muons



b-tag scenarios  
|η| < 5

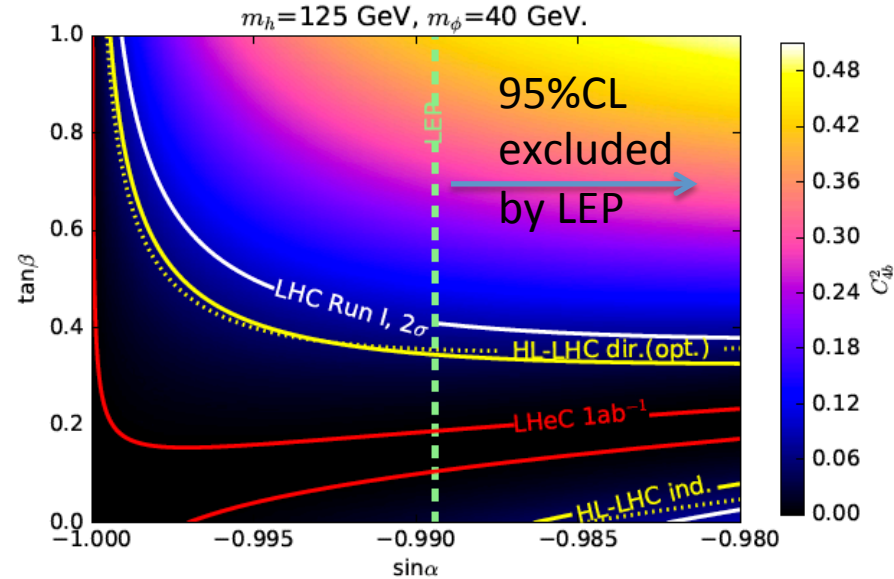
- (A) ε<sub>b</sub> = 70%, ε<sub>c</sub> = 10%, ε<sub>g,u,d,s</sub> = 1%
- (B) ε<sub>b</sub> = 70%, ε<sub>c</sub> = 20%, ε<sub>g,u,d,s</sub> = 1%
- (C) ε<sub>b</sub> = 60%, ε<sub>c</sub> = 10%, ε<sub>g,u,d,s</sub> = 1%
- (D) ε<sub>b</sub> = 60%, ε<sub>c</sub> = 20%, ε<sub>g,u,d,s</sub> = 1%

1ab<sup>-1</sup>, B-tagging and mistagging rates vary.



95% C.L. for m<sub>φ</sub> of 20, 40, 60 GeV for  
 $C_{4b}^2 = \kappa_V^2 \times \text{Br}(h \rightarrow \phi\phi) \times \text{Br}^2(\phi \rightarrow b\bar{b})$   
 is 0.3%, 0.2% and 0.1%

## Sensitivity comparison in Higgs Singlet Model



$$\Phi \equiv \begin{pmatrix} 0 \\ \frac{\tilde{h}+v}{\sqrt{2}} \end{pmatrix}, S \equiv \frac{h'+x}{\sqrt{2}} \quad (12)$$

Here  $v = 246$  GeV ensures the correct mass generation for  $W, Z$  bosons and SM fermions. The gauge eigenstates  $\tilde{h}, h'$  can be related to mass eigenstates  $\phi, h$  via an orthogonal rotation

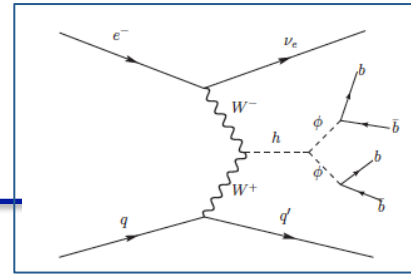
$$\begin{pmatrix} \phi \\ h \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} \tilde{h} \\ h' \end{pmatrix} \quad (13)$$

Now it is convenient to parameterize the model in terms of five more physical quantities: ( $m_\phi, m_h$  are masses of  $\phi$  and  $h$  respectively)

$$m_\phi, m_h, \alpha, v, \tan \beta \equiv \frac{v}{x} \quad (14)$$

# Samples

Focusing on dominant backgrounds



Uta Klein  
Michael o'Keefe  
Liverpool

Sample	Process	Generator Level Constraints	Cross Section (fb)
Signal <sup>1</sup>	$h \rightarrow \phi\phi \rightarrow b b^- b b^-$	PT of Jets/b's/photons/charged leptons > 6.5 GeV $\eta$ of Jets/b's/photons/charged leptons < 6.1 Min. $\Delta R$ between jets = 0.2 Min. Inv. Mass of Jet/ $bb^-$ pair = 8 GeV	51.34
CC Single Top Production	$pe \rightarrow j t^- \nu_l \text{ all } / h$ $(t^- \rightarrow W^- b^-, W^- \rightarrow \text{all all})$	PT of Jets/b's/photons/charged leptons > 6.5 GeV $\eta$ of Jets/b's/photons/charged leptons < 6.1 Min. $\Delta R$ between jets = 0.2 Min. Inv. Mass of Jet/ $bb^-$ pair = 8 GeV	11,347
CC Top+Multijet Sample	all = $g u c d s u^- c^- d^- s^-$ , $\nu_e \nu_m \nu_t \nu_e^- \nu_m^- \nu_t^-$ ta- ta+ b b, $z w^+ w^- h t t^-$ $pe^- \rightarrow b^- \text{ all all } \nu_l$ $pe^- \rightarrow b \text{ all all } \nu_l$	PT of Jets/b's/photons/charged leptons > 6.5 GeV $\eta$ of Jets/b's/photons/charged leptons < 6.1 Min. $\Delta R$ between jets = 0.2 Min. Inv. Mass of Jet/ $bb^-$ pair = 8 GeV Beam Polarisation = 0	9683
CC Inclusive Single W/Z/h Production	$pe^- \rightarrow \nu_l w^- jj / t^- t, w^- \rightarrow jj$ $pe^- \rightarrow \nu_l h jj, h \rightarrow jj$ $pe^- \rightarrow \nu_l z jj, z \rightarrow jj$	$\eta$ of Jets/b's/photons/charged leptons < 6.1 Min. $\Delta R$ between jets = 0.2 Min. Inv. Mass of Jet/ $bb^-$ pair = 8 GeV Beam Polarisation = 0	3566
CC $b b^- + 2j$ Production	all = $g u c d s u^- c^- d^- s^- \nu_e \nu_m$ $\nu_t \nu_e^- \nu_m^- \nu_t^-$ ta- ta+ b $b^-$ $pe^- \rightarrow b^- b \text{ all all } \nu_l$	PT of Jets/b's/photons/charged leptons > 6.5 GeV $\eta$ of Jets/b's/photons/charged leptons < 6.1 Min. $\Delta R$ between jets = 0.2 Min. Inv. Mass of Jet/ $bb^-$ pair = 8 GeV Beam Polarisation = 0	1120

➔ low pT and low dijet mass generation to retain sensitivity for 20 GeV scalar