

On accuracy of central mass energy determination for FCCee_z_202_nosol_13.seq

A. Bogomyagkov

Budker Institute of Nuclear Physics
Novosibirsk

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Introduction: different energies

Circumference: Π

Design energy: E_0 magnets fields

Average energy: $\langle E \rangle = \oint E(s) \frac{ds}{\Pi}$

Measured energy: $E_{meas} = f(W)$ function of spin tune

Invariant mass:
(central mass energy) M

Introduction: spin precession frequency

Ω_0 is revolution frequency, W is spin precession frequency.

Gyromagnetic ratio: $q = q_0 + q' = \frac{e}{mc} + q'$.

$$W = \frac{1}{2\pi} \oint \left(\frac{q_0}{\gamma} + q' \right) B_{\perp}(\theta) d\theta = \Omega_0 \cdot \left(1 + \frac{q'}{q_0} \frac{\langle B_{\perp} \rangle}{\langle B_{\perp}/\gamma \rangle} \right)$$
$$\approx \Omega_0 \cdot \left(1 + \langle \gamma \rangle \frac{q'}{q_0} \right),$$

$$\frac{q'}{q_0} = \frac{g-2}{2} = 1.1596521859 \cdot 10^{-3} \pm 3.8 \cdot 10^{-12}.$$

$$E[\text{MeV}] = 440.64843(3) \left(\frac{W}{\Omega_0} - 1 \right).$$

Spin distribution width: synchrotron oscillations

Synchrotron oscillations: $\delta = \Delta E/E_0 = a \cdot \cos(\omega_{\text{synt}} t)$.

$$W = \Omega_0 \left(1 + \nu_0 - \alpha_0 \nu_0 \frac{a^2}{2} \right) + \Omega_0 (\nu_0(1 - \alpha_0) - \alpha_0) \sin(\omega_{\text{synt}} t) + \alpha_0 \Omega_0 \nu_0 \frac{a^2}{2} \cos(2\omega_{\text{synt}} t)$$

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Spin precession frequency distribution shifts and becomes wider by

$$\left\langle \frac{W - \Omega_0(1 + \nu_0)}{\Omega_0(1 + \nu_0)} \right\rangle = \left\langle -\frac{\alpha_0 \nu_0 \frac{a^2}{2}}{1 + \nu_0} \right\rangle = -\frac{\alpha_0 \nu_0 \sigma_\delta^2}{1 + \nu_0} = -2 \cdot 10^{-12}$$

$$\frac{\Delta E}{E_0} = -2 \cdot 10^{-14}$$

Energy dependent momentum compaction

Momentum compaction: $\alpha = \alpha_0 + \alpha_1 \delta$

Synchrotron oscillations: $\ddot{\delta} = -\omega_{syn}^2 \delta - \omega_{syn}^2 \frac{\alpha_1}{\alpha_0} \delta^2$

Average and RMS: $\langle \delta \rangle = -\frac{\alpha_1}{\alpha_0} \sigma^2, \langle \delta^2 \rangle = \sigma^2$

Average W : $\langle W \rangle_\delta = \gamma_0 \Omega_0 \frac{q'}{q_0} \left(1 - \alpha_0 \sigma^2 - \frac{\alpha_1}{\alpha_0} \sigma^2 \right)$

Average energy: $\langle E \rangle = E_0 \left(1 - \frac{\alpha_1}{\alpha_0} \sigma^2 \right)$

Measured energy: $E_{meas} = E_0 \left(1 - \frac{\alpha_1}{\alpha_0} \sigma^2 - \alpha_0 \sigma^2 \right)$

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$$E_0 = 45.6 \text{ GeV}, \alpha_0 = 1.5 \cdot 10^{-5}, \alpha_1 = -9.8 \cdot 10^{-6}, \sigma = 3.8 \cdot 10^{-4}$$

$$\frac{\langle E \rangle - E_{meas}}{E_0} = \alpha_0 \sigma^2 = 2 \cdot 10^{-12}$$

$$\frac{\langle E \rangle - E_0}{E_0} = -\frac{\alpha_1}{\alpha_0} \sigma^2 = 1 \cdot 10^{-7}$$

Longitudinal field compensation

Detector field is $B_0 = 2$ T.

Deviation of compensating field is $\Delta B_c = 0.1$ T.

Length of compensating solenoid is $L_c = 0.75$ m.

$B_\rho = 152.105$ T · m, $E_0 = 45.6$ GeV, $\nu = 103.484$.

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$$\Delta\nu = \frac{\varphi^2}{8\pi} \cot(\pi\nu) \approx \frac{1}{8\pi} \cot(\pi\nu) \left(\frac{\Delta B_c}{B_0} \frac{2B_0 L_c}{B_\rho} \right)^2 \approx 2 \times 10^{-9}.$$

$$\frac{\Delta E}{E_0} = \frac{\Delta\nu \cdot 440.65}{E_0} \approx 2 \times 10^{-11}.$$

Spin distribution width: horizontal betatron oscillations

Ya.S. Derbenev, et al., "Accurate calibration of the beam energy in a storage ring based on measurement of spin precession frequency of polarized particles", Part. Accel. 10 (1980) 177-180

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Sextupole fields introduce additional $B_{\perp} \propto x^2$, $K2 = \frac{1}{B\rho} \frac{\partial^2 B_y}{\partial x^2}$.

Spin precession frequency distribution shifts and becomes wider by

$$\frac{\Delta\nu}{\nu} = -\frac{1}{2\pi} \oint \left(\varepsilon_x \beta_x(s) + \eta_x(s)^2 \sigma_{\delta}^2 \right) K2(s) ds.$$

$$\frac{\Delta\nu}{\nu} = \frac{\Delta E}{E_0} = -2.5 \cdot 10^{-7}.$$

Vertical magnetic fields: horizontal correctors

One corrector with deflection χ : $\frac{\Delta E}{E_0} = -\frac{\chi\eta_x}{\alpha\Pi}$, $\chi = \oint \frac{\Delta B_y}{B\rho} ds$.

RMS of energy shift: $\sigma\left(\frac{\Delta E}{E_0}\right) = \frac{2\sqrt{2}\sin(\pi\nu_x)}{\alpha\Pi} \frac{\langle\eta_x\rangle}{\langle\beta_x\rangle} \sigma_x$.

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$$\sigma\left(\frac{\Delta E}{E_0}\right) = -1.2 \cdot 10^{-3} [m^{-1}] \cdot \sigma_x [m],$$

$\sigma\left(\frac{\Delta E}{E_0}\right) = 10^{-6}$ demands stability of the horizontal orbit between calibrations

$$\sigma_x = 0.8 \text{ mm.}$$

Vertical magnetic fields: quadrupoles

$$\text{Shifted quadrupole: } \frac{\Delta E}{E_0} = -\frac{\chi \eta_x}{\alpha \Pi}, \quad \chi = K1L \cdot \Delta x, \quad K1 = \frac{1}{B\rho} \frac{\partial B_y}{\partial x}.$$

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$\frac{\Delta E}{E_0} = 10^{-6}$ demands stability of quadrupoles position between calibrations (10 min)

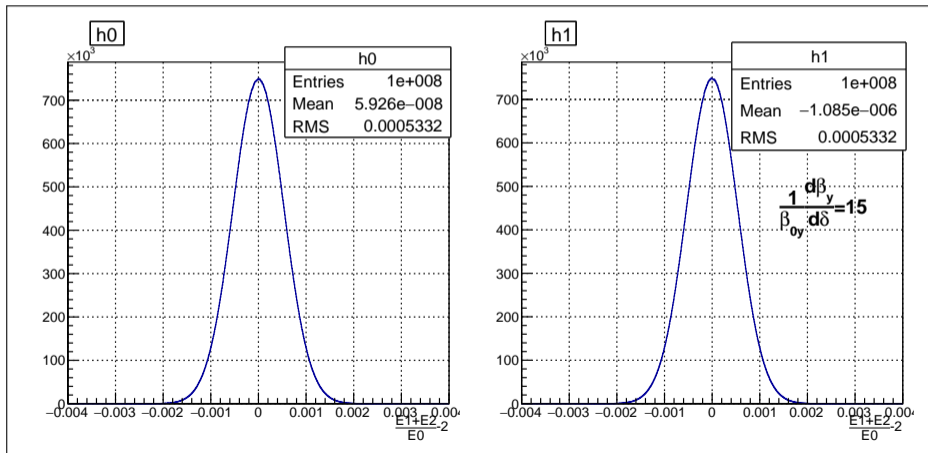
Quadrupole	Δx , m
QC7.1:	$2 \cdot 10^{-4}$
QY2.1:	$7.6 \cdot 10^{-5}$
QFG2.4:	$1.6 \cdot 10^{-4}$
QF4.1:	$1.4 \cdot 10^{-4}$
QG6.1:	$3.5 \cdot 10^{-5}$
QF4:	$\Delta x / \sqrt{720} = 5 \cdot 10^{-6}$

Central mass energy: β chromaticity

Invariant mass: $M^2 = (E_1 + E_2)^2 \cos^2(\theta) + O(m_e^2) + O(\sigma_\alpha^2) + O(\sigma_E^2)$.

Beta function chromaticity at IP: $\beta_{x,y} = \beta_{0x,y} + \beta_{1x,y}\delta$, $\sigma_{x,y}^2 = \varepsilon_{x,y}\beta_{x,y}$.

Particles with energy deviation have higher collision rate.



Central mass energy: β chromaticity

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$\frac{1}{\beta_x} \frac{d\beta_x}{d\delta}$	$\frac{1}{\beta_y} \frac{d\beta_y}{d\delta}$	$\Delta M, \text{ keV}$	$\frac{\Delta M}{E_0}$
0	15	-49 ± 2.4	$-1.1 \cdot 10^{-6} \pm 5 \cdot 10^{-8}$
200	0	-26 ± 2.4	$-5.7 \cdot 10^{-7} \pm 5 \cdot 10^{-8}$
200	15	-75 ± 2.4	$-1.6 \cdot 10^{-6} \pm 5 \cdot 10^{-8}$

Need to measure and adjust $\frac{1}{\beta_{0y}} \frac{d\beta_y}{d\delta}$.

Energy dependence on azimuth: full tapering

Two diametrically opposite RF cavities, U_0 — energy loss per revolution, $E(0)$ — after RF cavity. Full tapering — magnets fields are adjusted to keep design curvature, quadrupole strength etc.

$$\frac{dE}{ds} \propto E^4, \quad E(s) = \frac{E(0)}{(1 + k \cdot s)^{\frac{1}{3}}}, \quad k \approx \frac{3}{\pi} \frac{U_0}{E(0)} + \frac{3}{\pi} \frac{U_0^2}{E(0)^2} + O(U_0^3)$$

Average energy: $\langle E \rangle \approx E(0) - \frac{U_0}{4} - \frac{U_0^2}{12E(0)}$.

Energy at the IP: $E(IP) = E(0) - \frac{U_0}{4}$.

The difference: $\frac{\langle E \rangle - E(IP)}{E(0)} \approx -\frac{1}{12} \frac{U_0^2}{E(0)^2} = 5 \cdot 10^{-8}$, for $E_0 = 45.6$ GeV (Z).

The difference: $\frac{\langle E \rangle - E(IP)}{E(0)} \approx -\frac{1}{12} \frac{U_0^2}{E(0)^2} = 2 \cdot 10^{-7}$, for $E_0 = 80.5$ GeV (WW).

Energy dependence on azimuth: partial tapering

Partial tapering — fields of magnets groups are adjusted to keep approximately design curvature.

Equations of motion (canonical variables)

$$\begin{cases} \sigma' = -K_0 x, \\ p_t' = \left(-\frac{eV_0}{p_0 c} \right) \sin \left(\phi_s + \frac{2\pi}{\lambda_{RF}\sigma} \right) \delta(s - s_0) - \frac{2}{3} \frac{e^2 \gamma^4}{p_0 c} K_0^2 \sigma. \end{cases}$$

Solution: $p_t(s) = p_{0t} - f(s)$.

$$\sigma = 0 = - \int_0^\Pi K_0(s) x(s) ds = -p_{0t} \alpha \Pi + \Pi \langle (K_0 f + \Delta K_0) \eta \rangle_s .$$

$$p_{0t} = \frac{1}{\alpha} \langle (K_0 f + \Delta K_0) \eta \rangle_s .$$

Two RF cavities and symmetrical arcs

$$\begin{cases} \langle p_t \rangle = p_{0t} - \langle f \rangle = p_{0t} - \frac{U_0}{4E_0} = \frac{\langle E \rangle - E_0}{E_0}, \\ p_t(IP) = p_{0t} - f(IP) = p_{0t} - \frac{U_0}{4E_0} = \frac{E_{IP} - E_0}{E_0}, \end{cases}$$
$$\begin{cases} \langle E \rangle = E_0 + E_0 p_{0t} - \frac{U_0}{4}, \\ E_{IP} = E_0 + E_0 p_{0t} - \frac{U_0}{4}. \end{cases}$$

There is no difference between $\langle E \rangle$ and E_{IP} in the first order. Numerical calculations are needed for not symmetrical arcs, magnet misalignments.

Collective field of the own bunch

Electron in the field of own bunch will have potential energy

$$U[\text{eV}] = \frac{N_p e^2 [\text{Gs}]}{\sqrt{2\pi} \sigma_z [\text{cm}]} \left(\gamma_e + \ln(2) - 2 \ln \left(\frac{\sigma_x + \sigma_y}{r} \right) \right) \frac{10^{-7}}{e[\text{C}]},$$

$\gamma_e = 0.577$ Euler constant, $N_p = 4 \cdot 10^{10}$ — bunch population, $r_{ip} = 15$ mm and $r_{arc} = 20$ mm — vacuum chamber radius at IP and in the arcs, $\sigma_{x,IP} = 6.2 \cdot 10^{-6}$ m, $\sigma_{y,IP} = 3.1 \cdot 10^{-8}$ m, $\sigma_{x,arc} = 1.9 \cdot 10^{-4}$ m, $\sigma_{y,arc} = 1.2 \cdot 10^{-5}$ m.

$$\frac{U_{ip}}{E_0} = \frac{192 \text{keV}}{45.6 \text{GeV}} = 4.2 \cdot 10^{-6},$$
$$\frac{U_{arc}}{E_0} = \frac{120 \text{keV}}{45.6 \text{GeV}} = 2.6 \cdot 10^{-6}.$$

Collective field of the opposite bunch

Potential energy at the center of the bunch $\{x, y, s, z = s - ct\} = \{0, 0, 0, 0\}$

$$U(x, y, s, ct) = -\frac{\gamma N_p r_e m c^2}{\sqrt{\pi}} \int_0^\infty dq \frac{\exp \left[-\frac{(x+s \cdot 2\theta)^2}{2\sigma_x^2+q} - \frac{y^2}{2\sigma_y^2+q} - \frac{\gamma^2(s+ct)^2}{2\gamma^2\sigma_s^2+q} \right]}{\sqrt{2\sigma_x^2+q} \sqrt{2\sigma_y^2+q} \sqrt{2\gamma^2\sigma_s^2+q}},$$

$$\frac{U(0, 0, 0, 0)}{E_0} = -\frac{0.4 \text{ MeV}}{45.6 \text{ GeV}} = -9.3 \cdot 10^{-6}.$$

Invariant mass

The four-momentum: $P^\mu = (E - e\varphi, \vec{p}) = (E - e\varphi, \vec{P} - \frac{e}{c}\vec{A})$,

Energy-momentum relation: $(E - e\varphi)^2 = m^2c^4 + c^2(\vec{p})^2$.

Invariant mass:

$$M^2 = (P_1^\mu + P_2^\mu)^2 = 2E_1e_1\varphi + 2E_2e_2\varphi + 2E_1E_2 - (e_1\varphi)^2 - (e_2\varphi)^2 - 2p^{(1)}p^{(2)}.$$

Longitudinal momentum:

$$\begin{aligned} p_{i,s} &= \sqrt{(E_i - e_i\varphi)^2 - p_{i,x}^2 - p_{i,y}^2} \\ &= E_0 \sqrt{(1 + \delta_i - u)^2 - \left(\frac{p_{i,x}}{E_0}\right)^2 - \left(\frac{p_{i,y}}{E_0}\right)^2}. \end{aligned}$$

Invariant mass

Average values

$$\langle M^2 \rangle = 4E_0^2 \cos^2(\theta)(1 - u^2) - 2E_0^2 \sigma_{px}^2 \cos(2\theta) - 2E_0^2 \sigma_{py}^2 \cos(2\theta)$$

$$\langle M \rangle = 2E_0 \cos(\theta) \left(1 - \frac{u^2}{2} \right) - \frac{E_0}{2} \left(\sigma_{\delta}^2 \cos(\theta) + \sigma_{px}^2 \cos(\theta) + \sigma_{py}^2 \frac{\cos(2\theta)}{\cos(\theta)} \right)$$

$$\langle M^2 \rangle - \langle M \rangle^2 = 2E_0^2 \cos^2(\theta) \left(\sigma_{\delta}^2 + \sigma_{px}^2 \tan^2(\theta) \right)$$

Invariant mass shift due to beam potentials

$$\frac{\langle M \rangle - 2E_0 \cos(\theta)}{2E_0 \cos(\theta)} = \left(1 - \frac{(e\varphi)^2}{2E_0^2} \right) \approx 4 \times 10^{-10}$$

What is not estimated?

- 1 Vertical orbit distortions. They will have significant impact.
- 2 Electron positron energy difference due to synchrotron radiation in not identical arcs, energy loss due to not identical impedance of the vacuum chamber in the arcs. Requires impedance estimations.
- 3 Influence of wrong RF cavities (LEP).

Largest corrections and errors

- 1 Beta function chromaticity (correction, tunable) $\sim 2 \times 10^{-6}$.
- 2 Horizontal correctors and shift of quadrupoles (error) $\sim 10^{-6}$ with position stability of arc quadrupoles $\Delta x < 5 \times 10^{-6}$ between calibrations (every 10 minutes).
- 3 Horizontal betatron oscillations and sextupole fields $2.5 \sim 10^{-7}$.

- 1 V.V. Danilov et al., “Longitudinal Beam-Beam Effects for an Ultra-High Luminosity Regime”, proceedings of PAC 1991, p. 526.
- 2 F. Zimmerman and Tor O. Raubenheimer, “Longitudinal space charge in final focus systems for linear colliders”, SLAC-PUB-7304 (1997).
- 3 V.E. Blinov et al., “Analysis of errors and estimation of accuracy in the experiment on precise mass measurement of J/ψ , ψ' mesons and tau lepton on the VEPP-4M collider”, NIM A 494 (2002) 68-74.
- 4 V.E. Blinov et al., “Absolute calibration of particle energy at VEPP-4M ”, NIM A 494 (2002) 81-85.