Cosmological Implications of the LARGE Volume String Scenario

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String Scenarios: Some cosmology challenges

- Big-bang singularity X
- Cosmological inflation or alternatives
- After inflation
- Current acceleration
- Consistency with low energy phenomenology

MODULI STABILISATION

4-cycle size: *t* (Kahler moduli)

3-cycle size: U (Complex structure moduli) + Dilaton S

String Scenarios

• IIB (+F-theory) KKLT LVS

Moduli Stabilisation

• IIA

Heterotic

G2 manifolds

GKP Overview

1. Fluxes: GVW $W_0 = \int G \wedge \Omega$ $G_3 = F_3 - iSH_3, \int F_3 = 2\pi M, \int H_3 = -2\pi K$ **Fix CS moduli: z and dilaton: S**

2. Warped throats $z^{1/3} = e^{A} = e^{-\frac{2\pi K}{3g_s M}} \equiv e^{-\alpha}$



KKLT Overview

• Nonperturbative effects: $W_{np} = \sum A_i e^{-a_i T_i}$

SUSY AdS Vacua: DW=0

Anti D3 brane (SUSY breaking+uplift)

$$V_{\text{uplift}} = \frac{D^2}{\left(T + T^*\right)^{\alpha}} = \frac{D^2}{\mathcal{V}^{2\alpha/3}} \quad \begin{cases} \alpha = 3 & \text{KKLT} \\ \alpha = 2 & \text{KKLMMT} \end{cases}$$



LARGE Volume Scenario

Perturbative corrections to K:

$$K = -2\ln\left(\mathcal{V} + \frac{\hat{\xi}}{2}\right)$$

$$V_F \propto \left(\frac{K^{S\bar{S}}|D_SW|^2 + K^{a\bar{b}}D_aW\bar{D}_{\bar{b}}\bar{W}}{\mathcal{V}^2}\right) + \left(\frac{Ae^{-2a\tau}}{\mathcal{V}} - \frac{Be^{-a\tau}W_0}{\mathcal{V}^2} + \frac{C|W_0|^2}{\mathcal{V}^3}\right)$$
$$\mathcal{V} \sim e^{a\tau} \quad \text{with} \quad \tau \sim \text{Re S} \sim 1/g_s > 1.$$

Exponentially large volume for weak coupling (SUSY broken by Fluxes, AdS)



- From F/D terms, hidden matter Burgess et al, Dudas et al, Villadoro-Zwirner, Cicoli et al, T-branes (Cicoli, FQ, Valandro arXiv:1512.04558)
- From non-perturbative effects on hidden brane at singularities



dS Moduli Stabilisation



Relevant ScalesString Scale
$$M_s = \frac{g_s^{1/4} M_P}{\sqrt{4\pi \mathcal{V}}},$$
Kaluza Klein Scale $M_{KK} \simeq \frac{M_P}{\sqrt{4\pi \mathcal{V}^{2/3}}},$

$$M_{KK} \simeq \frac{M_P}{\sqrt{4\pi} \mathcal{V}^{2/3}},$$

Gravitino mass

$$m_{3/2} \simeq \left(rac{g_s^2}{2\sqrt{2\pi}}
ight) rac{W_0 M_P}{\mathcal{V}} \,.$$

Volume modulus mass ¹

$$m_{\mathcal{V}} \simeq m_{3/2} / \sqrt{\mathcal{V}}.$$

Vacuum decay rates

$$\Gamma \sim e^{-\mathcal{V}^3}$$

Constraints on the volume

- Validity of EFT $(m_{3/2} < < M_{kk}) : V >> 10^3$
- **CMP** (m_{volume} > 30 TeV): V<10⁹

Ranges of relevant scales (GeV)

LVS vs KKLT

- $W_0 \sim 0.1 100$
- AdS non SUSY
- Minimum: perturbative in big cycle vs non-perturb. in small cycle
- Uplift:anti D3 branes, Dterms...
- Small parameter = 1/V
- SUSY broken by fluxes
- Many moduli: need $h_{21}>h_{11}>1$ + one blow up, the rest by loop effects/D-terms

- W₀<<1
- AdS SUSY
- Minimum: tree-level vs nonperturbative
- Uplift: anti D3 branes...(no Dterms)
- Small parameter W0
- SUSY broken by uplifting mechanism
- Many moduli: nonperturbative effects for each of them or ...

Revisiting Anti D3 Brane Uplift

Nilpotent Superfields EFT

$$\begin{split} X &= X_0(y) + \sqrt{2}\psi(y)\theta + F(y)\theta\bar{\theta} & \operatorname{Rocek, Komargodski-Seiberg, Dudas et al, Kallosh et al, D'allagata et al...} \\ X_0 &= \frac{\psi\psi}{2F} \\ K &= K_0 X X^* & W = \rho X + W_0 \\ V &= K_0^{-1} \left\| \frac{\partial W}{\partial X} \right\|^2 = \frac{|\rho|^2}{K_0} \ge 0 \\ \mathcal{L} &= -\rho^2 + i\partial_a \bar{\psi} \bar{\sigma}^a \psi + \frac{1}{4\rho^2} \bar{\psi}^2 \partial^2 \psi^2 - \frac{1}{16\rho^6} \psi^2 \bar{\psi}^2 \partial^2 \psi^2 \partial^2 \bar{\psi}^2 \end{split}$$

~ Volkov-Akulov !

Nilpotent Superfields and KKLT

Goldstino: Nilpotent
$$X^2(x, \theta) = 0.$$
 chiral superfield

KKLT $K = -3 \log (T + T^*) + c (T + T^*)^n XX^* + ZCC^* + \cdots$ $Z = (T + T^*)^m + b (T + T^*)^k XX^*$ $W = W_0 + W_{\text{matter}} + W_{np} + \rho X$

Plug into SUGRA expression for V, $V = V_{KKLT} + V_{uplift}$:

$$V_{\text{uplift}} = \frac{|\rho|^2}{c(T+T^*)^{n+3}}$$

(like KKLT, KKLMMT)

Kallosh et al. 2013-15 see also Polchinski @ SUSY 2015

Antibrane uplift from manifestly SUSY EFT!

Anti D3 Brane/O3⁻ Spectrum



Metastability of dS LVS minima

• Brown-Teitelboim (+CdL) D5/NS5 brane nucleation

• AdS: Brane tension>upper bound, so stable in EFT

• dS: Decay rate ~ $exp(-V^3)$

 $P_{dS}/P_{AdS} \sim e^{-V}$ (The larger the volume, more stable!) $P_{dS}/P_{dec} \sim e^{V^2}$

(Also: no evidence for bubble of nothing decay)

S. de Alwis, R. Gupta, E. Hatefi, FQ

arXiv:1308.1222

Inflation (moduli as inflatons)

e.g. 1: Swiss Cheese Calabi-Yau's





e.g.

 $\mathbb{P}_{[1,3,3,3,5]}[15] \qquad \mathbb{P}^4_{1,2,2,10,15}(30) \qquad \mathbb{P}^4_{1,1,2,2,6}(12)/\mathbb{Z}_2 \qquad \mathbf{M}^{(dP_8)^n}_{\mathbf{n}}$

Blumenhagen, et al., Grimm et al., Kreuzer et al. 08

Kähler Moduli Inflation (Blow-up)



$$V \cong V_0 - \frac{4W_0 a_n A_n}{\mathcal{V}^2} \left(\frac{3\mathcal{V}}{4\lambda}\right)^{2/3} (\tau_n^c)^{4/3} \exp\left[-a_n \left(\frac{3\mathcal{V}}{4\lambda}\right)^{2/3} (\tau_n^c)^{4/3}\right].$$

Calabi-Yau: h₂₁>h₁₁>2

Small field inflation (r<<<1) 0.960<n<0.967 Loop corrections??



e.g. 2: Fibre Calabi-Yau

e.g.

$$\mathcal{V} = \lambda_1 t_1 t_2^2 + \lambda_2 t_3^3,$$

 $au_i = \partial \mathcal{V} / \partial t_i,$





String Scenario	n_s	r				
$\mathrm{D}3/\overline{\mathrm{D}3}$ Inflation	$0.966 \le n_s \le 0.972$	$r \le 10^{-5}$				
Inflection Point Inflation	$0.92 \le n_s \le 0.93$	$r \le 10^{-6}$				
DBIInflation	$0.93 \le n_s \le 0.93$	$r \le 10^{-7}$				
Wilson Line Inflation	$0.96 \le n_s \le 0.97$	$r \le 10^{-10}$				
${ m D3/D7}$ Inflation	$0.95 \le n_s \le 0.97$	$10^{-12} \le r \le 10^{-5}$				
Racetrack Inflation	$0.95 \le n_s \le 0.96$	$r \le 10^{-8}$,			
N - flation	$0.93 \le n_s \le 0.95$	$r \le 10^{-3}$				
Axion Monodromy	$0.97 < n_s < 0.98$	0.04 < r < 0.07				
Kahler Moduli Inflation	$0.96 \le n_s \le 0.967$	$r \le 10^{-10}$				
Fibre Inflation	$0.965 \le n_s \le 0.97$	$0.0057 \le r \le 0.007$	-			
Poly – instanton Inflation	$0.95 \le n_s \le 0.97$	$r \le 10^{-5}$	ß			
Ovorall str	ing inflatio	'n	r _{0.002}) 0.20 0.2	•	I	PI. PI.

Overall, string inflation models in good shape after Planck 2013-2015

C. Burgess, M. Cicoli, FQ arXiv:1306.3512



Fibre vs Starobinsky Inflation

• Starobinsky $\alpha = 1$, Fibre $\alpha = 2$.

• Starobinsky from strings? $f(R) = R + \frac{a_2}{M^2}R^2 + \frac{a_3}{M^4}R^3 + \frac{a_4}{M^6}R^4 + \dots$

$$V(\phi) = \frac{M^2}{2\kappa^2} e^{-2\sqrt{\frac{2}{3}}\phi} (U_0 + U_1 e^{\sqrt{\frac{2}{3}}\phi} + U_2 e^{2\sqrt{\frac{2}{3}}\phi} + U_3 e^{3\sqrt{\frac{2}{3}}\phi} + \dots$$

• Fibre very generic: Most known CY are fibrations (Anderson, Gray, et al 2015)

Summary of Fibre Inflation

- String model of inflation with moduli stabilisation incorporated.
- Similar physics but much better rooted than Starobinsky duals (UV completion, tuning, etc.)

- Multi field generalisations (but only small nongaussianities, Burgess et al 2010, 2012)
- Low l effects, α ' inflation, global realisation (Cicoli et al. 2016)

Kahler+Fibre Inflation

Stringy realisation of α-attractors

• α=2 (fibre inflation) Burgess, Cicoli, FQ (2007)





- α=(VInV)⁻¹ (Kahler blow-up inflation)
- Conlon, FQ (2006)



• ...α=(InV)⁻¹ (polyinstanton inflation) Cicoli, Pedro, Tasinato (2011)

Inflation: Fibre+Kahler





After Inflation

General prediction

Axion partner of the volume: mass < exp(-volume) < 10⁻²² eV Dark energy or matter and dark radiation

Cosmological Moduli Problem

- Usually moduli masses = $m_{3/2}$ (de Carlos et al 1993, Scrucca-Gomez-Reino 2008)
- And assume soft terms = $m_{3/2}$
- Identify m_{3/2}=1 TeV
- But LVS is nongeneric scenario
- Volume modulus mass<<m_{3/2}
- So CMP more acute than expected!
- Unless $m_{3/2} >> 1 \text{TeV}$

Cosmological Moduli 'Problem'



Coughlan et al 1983, Banks et al, de Carlos et al 1993

e.g. After Kahler Inflation

Explicit computation of Vacuum misalignement

$$Y = \frac{\delta\varphi}{M_{\rm pl}} = \sqrt{\frac{2}{3}}\delta\phi \simeq 2\sqrt{\frac{2}{3}}R\phi_* \simeq 0.1 - 1$$

M.Cicoli, K. Dutta, A. Maharana, FQ <u>arXiv:1604.08512</u>

Number of efoldings:

$$\begin{split} N_e + \frac{1}{4} N_{\rm mod} + \frac{1}{4} (1 - 3w_{\rm re}) N_{\rm re} &\approx 57 + \frac{1}{4} \ln r + \frac{1}{4} \ln \left(\frac{\rho_*}{\rho_{\rm end}}\right)^{\log(aH)} \\ &\left(55 - \frac{1}{4} N_{\rm mod}\right) \pm 5 \\ N_{\rm mod2} &\approx \frac{2}{3} \ln \left(\frac{16\pi \mathcal{V}^{5/2} (\ln \mathcal{V})^{5/2} Y^4}{10\beta^2}\right) \\ N_e &\simeq 44.65 + \frac{1}{4} \ln \left(\frac{\rho_*}{\rho_{\rm end}}\right) \simeq 45 \qquad n_s \simeq 0.955 \,. \end{split}$$

Thermal History

Alternative History



From S. Watson, SUSY 2013



Volume Reheating*

*Sequestered scenarios

M.Cicoli, J.P. Conlon, FQ arXiv:1208.3562 T. Higaki, F.Takahashi arXiv:1208.3563

$$\begin{split} \Gamma_{\Phi \to a_{b}a_{b}} &= \frac{1}{48\pi} \frac{m_{\Phi}^{3}}{M_{P}^{2}}. & \text{Volume axion } a_{b} \\ \Gamma_{\Phi \to H_{u}H_{d}} &= \frac{2Z^{2}}{48\pi} \frac{m_{\Phi}^{3}}{M_{P}^{2}} & \text{Higgses} \\ \Gamma_{\Phi \to BB} &= \left(\frac{\lambda}{3/2}\right)^{2} \frac{9}{16} \frac{1}{48\pi} \frac{m_{\Phi}^{3}}{M_{P}^{2}} & \text{Closed string axions} \\ \Gamma_{\Phi \to C\bar{C}} &\sim \frac{m_{0}^{2}m_{\Phi}}{M_{P}^{2}} \ll \frac{m_{\Phi}^{3}}{M_{P}^{2}} & \text{Matter scalars C} \\ \hline T_{reheat} &\sim \frac{m_{\Phi}^{3/2}}{M_{Pl}^{1/2}} \sim 0.6 \ \text{GeV} \left(\frac{m_{\Phi}}{10^{6} \text{GeV}}\right)^{3/2}. \end{split}$$

Dark Radiation

Energy density:

$$\rho_{total} = \rho_{\gamma} \left(1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{eff} \right). \label{eq:rhotal}$$

Standard Model N_{eff}=3.04

At CMB: WMAP, ACT, SPT

$$3.12\,\kappa \leq \Delta N_{eff} \leq 3.48\,\kappa$$

 $\kappa = (1+9n_a/16)/n_HZ^2$

Simplest Z=1:

 $1.56 \leq \Delta N_{eff} \leq 1.74$ General: Strong constraints on matter and couplings! _ (

$$0.14 \lesssim \Delta N_{\rm eff} \lesssim 1.6$$

Cicoli+Muia 2016

MSSM: Non-thermal Higgsino DM



Model independent indirect search 300-600 GeV Higgsinos all others multi TeV

CONCLUSIONS

- Several stringy (EFT) de Sitter scenarios
- Several inflationary scenarios with concrete predictions (blow-up: r<<1, low-energyy SUSY, loops? Fibre: r<0.01, too large SUSY scale)
- CMP: after inflation signature of strings (e.g. dark radiation, non-thermal MSSM,...)
- Most known ingredients used (stringy vs simplicity): geometry, fluxes, branes, perturbative, non-perturbative effects

De Sitter 3:

DILATON DEPENDENT NON-PERTURBATIVE EFFECTS (ON LVS)

Hidden Sector on Branes@Singularities



Hidden

Burgess, Cicoli, Maharana, FQ 2012

$$W_{\rm np} = A e^{-a T} + B e^{-b (S+hQ)}$$

$$V = V_{\rm LVS} + V_{\rm up}$$

$$V_{\rm up} \propto h^2 \, \frac{e^{-2b\langle s \rangle}}{\mathcal{V}} \,,$$

Not explicit CY realisation yet

CONCRETE COMPACT CY



Enhancing the value of r?

$$V \simeq \frac{C_2}{\langle \mathcal{V} \rangle^{10/3}} \left[(3-R) - 4\left(1 + \frac{1}{6}R\right) e^{-\kappa\hat{\varphi}/2} + \left(1 + \frac{2}{3}R\right) e^{-2\kappa\hat{\varphi}} + R e^{\kappa\hat{\varphi}} \right]$$

$$R <<1$$

• $n_s = 0.98$ and r=0.01 for $N_e \simeq 50$

• $n_s = 0.99$ and r=0.01 for $N_e \simeq 60$

May need Neff>3.04?

Comments on α-attractors

Kallosh and Linde



e.g. Non-Thermal Dark-Matter (MSSM)

- KKLT: gravitino decay
- KKLT: D7 Higgsino overproduction
- KKLT:D3 small region allowed Higgsino DM

- LVS: Volume decay
- LVS:D7 Higgsino overproduction
- LVS: D3: allowed region to be constrained by 1Ton (Xenon, CTA) and 100TeV (not LHC).

L. Aparicio, M. Cicoli, B Dutta, F. Muia + FQ arXiv:1607.00004