

Tabletop Probes of Ultra-Low-Mass Bosonic Dark Matter

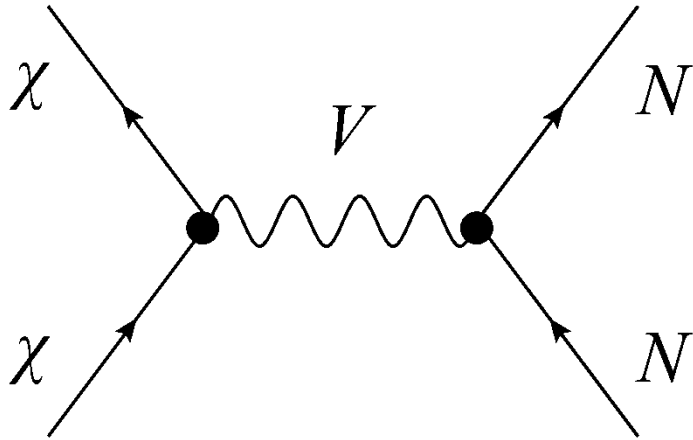
Yevgeny Stadnik

Humboldt Fellow

Johannes Gutenberg University, Mainz, Germany

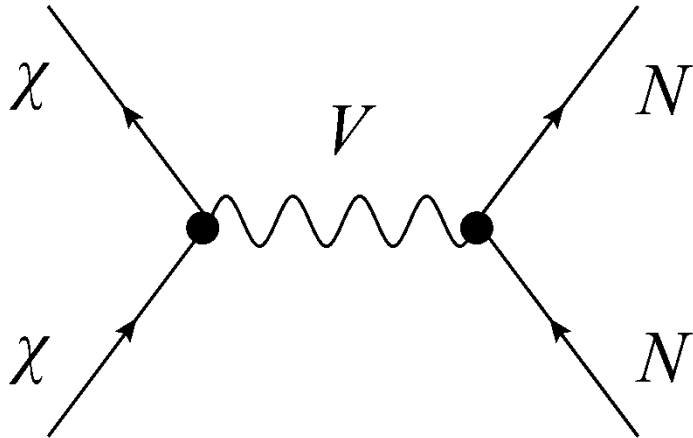
Motivation

Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles ($m_\chi \sim \text{GeV}$) have not yet produced a strong positive result.



Motivation

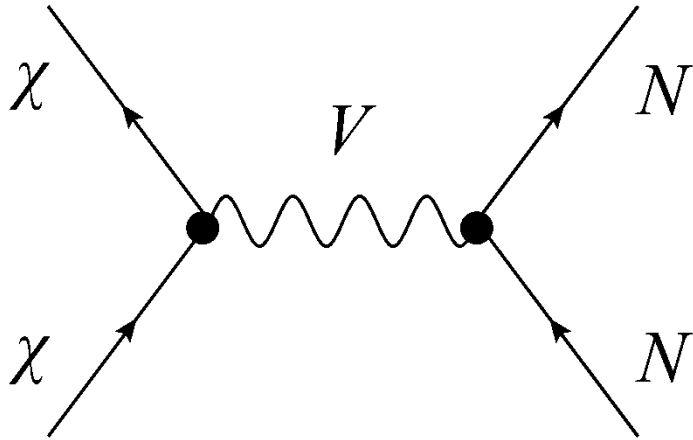
Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles ($m_\chi \sim \text{GeV}$) have not yet produced a strong positive result.



$$\mathcal{M} \propto (e')^2$$

Motivation

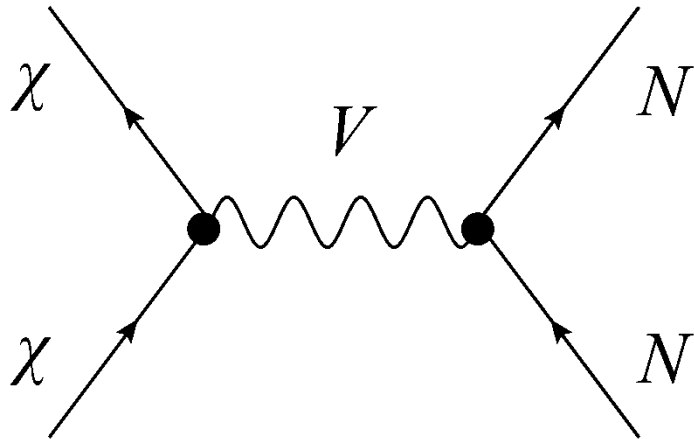
Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles ($m_\chi \sim \text{GeV}$) have not yet produced a strong positive result.



$$\mathcal{M} \propto (e')^2$$
$$\Rightarrow \frac{d\sigma}{d\Omega} \propto |\mathcal{M}|^2 \propto (e')^4$$

Motivation

Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles ($m_\chi \sim \text{GeV}$) have not yet produced a strong positive result.



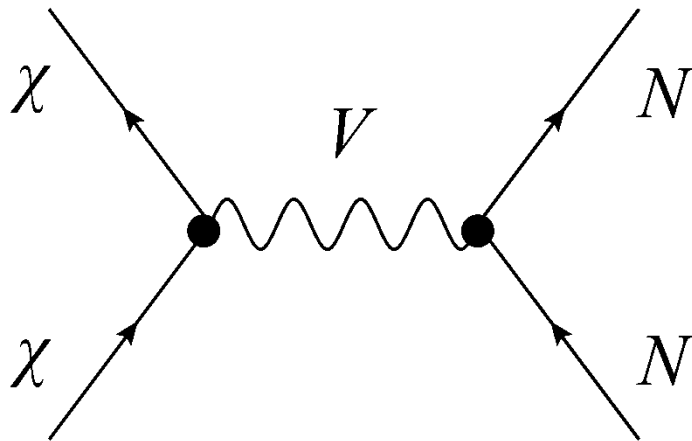
$$\mathcal{M} \propto (e')^2$$

$$\Rightarrow \frac{d\sigma}{d\Omega} \propto |\mathcal{M}|^2 \propto (e')^4$$

Challenge: Observable is **fourth power** in a small interaction constant ($e' \ll 1$)!

Motivation

Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles ($m_\chi \sim \text{GeV}$) have not yet produced a strong positive result.

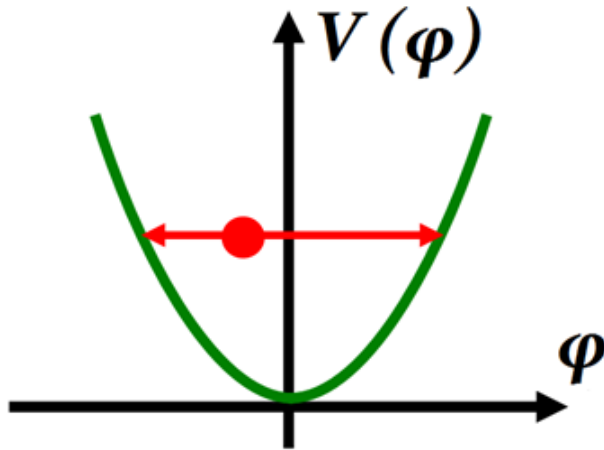


$$\mathcal{M} \propto (e')^2$$
$$\Rightarrow \frac{d\sigma}{d\Omega} \propto |\mathcal{M}|^2 \propto (e')^4$$

Question: *Can we instead look for effects of dark matter that are **first power** in the interaction constant?*

Low-mass Spin-0 Dark Matter

- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)



$$V(\phi) = \frac{m_\phi^2 \phi^2}{2}$$

$$\tau_{\text{coh}} \sim \frac{2\pi}{m_\phi \langle v_{\text{DM}}^2 \rangle} \sim 10^6 T_{\text{osc}}$$

Low-mass Spin-0 Dark Matter

- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- Coherently oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)

Low-mass Spin-0 Dark Matter

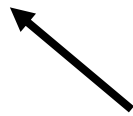
- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- Coherently oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- Classical field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$

Low-mass Spin-0 Dark Matter

- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- Coherently oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- Classical field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$
- Coherent + classical DM field = “**Cosmic laser field**”

Low-mass Spin-0 Dark Matter

- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- Coherently oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- Classical field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$
- Coherent + classical DM field = “**Cosmic laser field**”
- $10^{-22} \text{ eV} \lesssim m_\varphi \lesssim 1 \text{ eV} \Leftrightarrow 10^{-8} \text{ Hz} \lesssim f \lesssim 10^{14} \text{ Hz}$

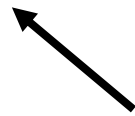


$\lambda_{\text{dB},\varphi} / 2\pi \leq L_{\text{dwarf galaxy}} \sim 1 \text{ kpc}$

Classical field

Low-mass Spin-0 Dark Matter

- *Low-mass spin-0 particles form a coherently oscillating classical field* $\varphi(t) = \varphi_0 \cos(m_\varphi c^2 t / \hbar)$, with energy density $\langle \rho_\varphi \rangle \approx m_\varphi^2 \varphi_0^2 / 2$ ($\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3$)
- Coherently oscillating field, since *cold* ($E_\varphi \approx m_\varphi c^2$)
- Classical field for $m_\varphi \lesssim 1 \text{ eV}$, since $n_\varphi (\lambda_{\text{dB},\varphi} / 2\pi)^3 \gg 1$
- Coherent + classical DM field = “**Cosmic laser field**”
- $10^{-22} \text{ eV} \lesssim m_\varphi \lesssim 1 \text{ eV} \Leftrightarrow 10^{-8} \text{ Hz} \lesssim f \lesssim 10^{14} \text{ Hz}$



$$\lambda_{\text{dB},\varphi} / 2\pi \leq L_{\text{dwarf galaxy}} \sim 1 \text{ kpc}$$

Classical field

- $m_\varphi \sim 10^{-22} \text{ eV} \Leftrightarrow T \sim 1 \text{ year}$

Low-mass Spin-0 Dark Matter

Dark Matter

**Scalars
(Dilaton):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

- Atomic clocks
- Optical cavities
- Fifth-force searches
- Astrophysics (e.g., BBN)

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

- Co-magnetometers
- Nuclear magnetic resonance
 - Torsion pendula

Low-mass Spin-0 Dark Matter

Dark Matter

**Scalars
(Dilatons):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

- Atomic clocks
- Optical cavities
- Fifth-force searches
- Astrophysics (e.g., BBN)

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

- Co-magnetometers
- Nuclear magnetic resonance
 - Torsion pendula

Low-mass Spin-0 Dark Matter

Dark Matter

**Scalars
(Dilatons):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

- Atomic clocks

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

- Co-magnetometers

“Thou shall measure frequency.”

Low-mass Spin-0 Dark Matter

Dark Matter

**Scalars
(Dilatons):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

- Atomic clocks

$$f \sim 10^{15} \text{ Hz}, \Delta f \sim 10^{-3} \text{ Hz}, \Delta f/f \sim 10^{-18}$$

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

- Co-magnetometers

$$f \sim 100 \text{ Hz}, \Delta f \sim 10^{-9} \text{ Hz}, \Delta f/f \sim 10^{-11}$$

Low-mass Spin-0 Dark Matter

Dark Matter

Scalars

(Dilatons):

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

• **Atomic clocks**

• $N \sim 10^5 - 10^{13}$ (or even 1!) [cf. $N \sim 10^{21} - 10^{29}$ (traditional “bulk” detectors)]

Pseudoscalars

(Axions):

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

• **Co-magnetometers**

Low-mass Spin-0 Dark Matter

Dark Matter

**Scalars
(Dilatons):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

• **Atomic clocks**

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

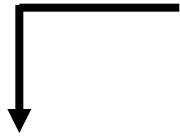
→ **Time-varying spin-
dependent effects**

• **Co-magnetometers**

- $N \sim 10^5 - 10^{13}$ (or even 1!) [cf. $N \sim 10^{21} - 10^{29}$ (traditional “bulk” detectors)]
- Search for *wave-like* signatures [cf. traditional *particle-like* recoil signatures]

Low-mass Spin-0 Dark Matter

Dark Matter



**Scalars
(Dilatons):**

$$\varphi \xrightarrow{P} +\varphi$$

→ **Time-varying
fundamental constants**

- Atomic clocks
- Optical cavities
- Fifth-force searches
- Astrophysics (e.g., BBN)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f$$

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f$$

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t)$$

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.*

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$
$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2}} + \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)}$$

* Linear couplings may be eliminated by a Z_2 symmetry (invariance under $\varphi \rightarrow -\varphi$)

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \frac{\phi_0^2}{2(\Lambda'_f)^2} + \frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)$$

$$\rho_\phi = \frac{m_\phi^2 \phi_0^2}{2} \quad \Rightarrow \quad \phi_0^2 \propto \rho_\phi$$

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2}} + \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)}$$

'Slow' drifts [Astrophysics
(high ρ_{DM}): BBN, CMB]

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2}} + \boxed{\frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)}$$

'Slow' drifts [Astrophysics

(high ρ_{DM}): BBN, CMB]

+ Gradients [Fifth forces]

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)],

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider quadratic couplings of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_\varphi t)$, with SM fields.

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \frac{\phi_0^2}{2(\Lambda'_f)^2} + \frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)$$

'Slow' drifts [Astrophysics
(high ρ_{DM}): BBN, CMB]
+ Gradients [Fifth forces]

Oscillating variations
[Laboratory (high precision)]

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

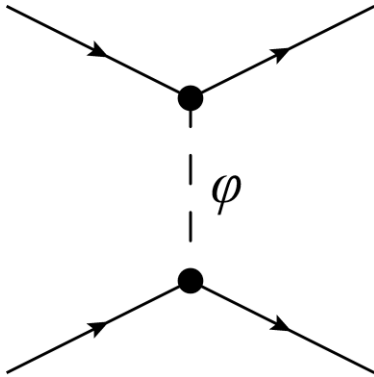
Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\phi\bar{X}X$)



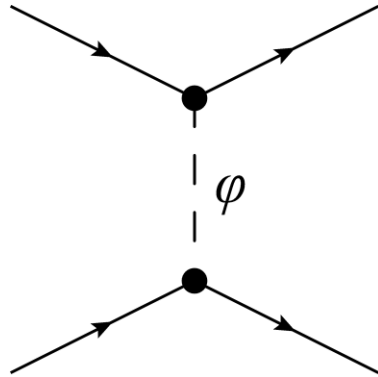
$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

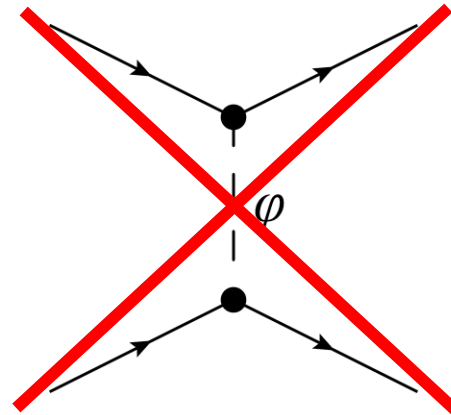
Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\phi\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

Quadratic couplings ($\phi^2\bar{X}X$)

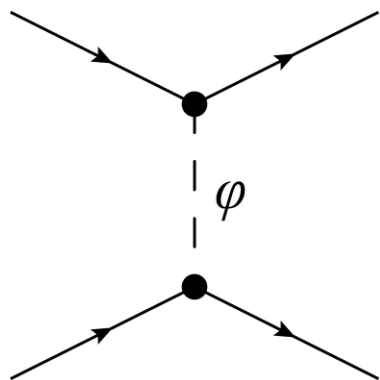


Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

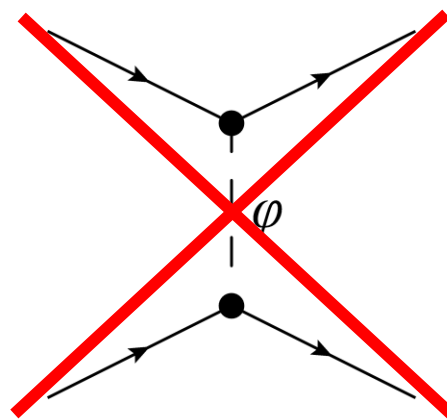
Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\phi\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

Quadratic couplings ($\phi^2\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) \left(1 - \frac{B}{r} \right)$$



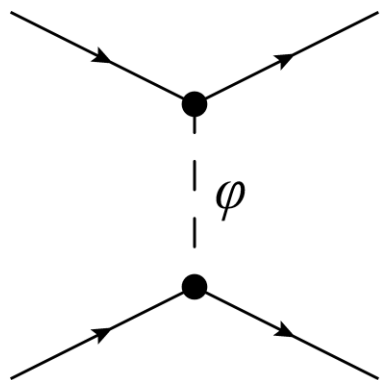
Gradients + screening/amplification

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider the effect of a massive body (e.g., Earth) on the scalar DM field

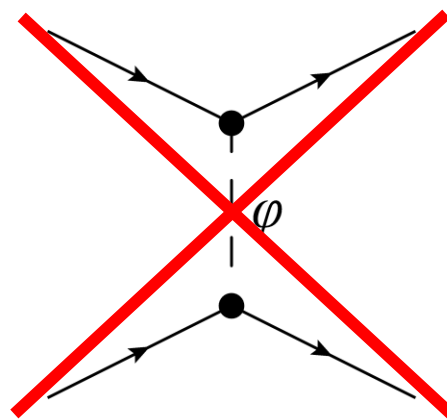
Linear couplings ($\phi\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

**Many different (classical) signatures
in “fifth-force” experiments**

Quadratic couplings ($\phi^2\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) \left(1 - \frac{B}{r} \right)$$



Gradients + screening/amplification

Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, *PRD* **91**, 015015 (2015)], [Stadnik, Flambaum, *PRL* **114**, 161301 (2015)]

$$\frac{\delta(\omega_1/\omega_2)}{\omega_1/\omega_2} \propto \sum_{X=\alpha, m_e/m_p, \dots} (K_{X,1} - K_{X,2}) \cos(\omega t)$$

↑ ↑
Sensitivity coefficients

$\omega = m_\phi$ (linear coupling) or $\omega = 2m_\phi$ (quadratic coupling)

Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, *PRD* **91**, 015015 (2015)], [Stadnik, Flambaum, *PRL* **114**, 161301 (2015)]

$$\frac{\delta(\omega_1/\omega_2)}{\omega_1/\omega_2} \propto \sum_{X=\alpha, m_e/m_p, \dots} (K_{X,1} - K_{X,2}) \cos(\omega t)$$

↑ ↑
Sensitivity coefficients

$\omega = m_\phi$ (linear coupling) or $\omega = 2m_\phi$ (quadratic coupling)

- Precision of optical clocks approaching $\sim 10^{-18}$ fractional level

Atomic Spectroscopy Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Arvanitaki, Huang, Van Tilburg, *PRD* **91**, 015015 (2015)], [Stadnik, Flambaum, *PRL* **114**, 161301 (2015)]

$$\frac{\delta(\omega_1/\omega_2)}{\omega_1/\omega_2} \propto \sum_{X=\alpha, m_e/m_p, \dots} (K_{X,1} - K_{X,2}) \cos(\omega t)$$

↑ ↑
Sensitivity coefficients

$\omega = m_\phi$ (linear coupling) or $\omega = 2m_\phi$ (quadratic coupling)

- Precision of optical clocks approaching $\sim 10^{-18}$ fractional level
- Sensitivity coefficients K_X calculated extensively by Flambaum group and co-workers (1998 – present), see the reviews

[Flambaum, Dzuba, *Can. J. Phys.* **87**, 25 (2009); *Hyperfine Interac.* **236**, 79 (2015)]

Laser Interferometry Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]



**Gravitational-wave
detector (LIGO/Virgo),
 $L \sim 4$ km**



**Small-scale cavity,
 $L \sim 0.2$ m**

Laser Interferometry Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

- Compare $L \sim Na_B$ with λ (or a 2nd L)

Laser Interferometry Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

- Compare $L \sim Na_B$ with λ (or a 2nd L)
- For a “usual” atomic optical transition and in the non-relativistic limit:

$$\Phi = \frac{\omega L}{c} \propto \left(\frac{e^2}{a_B \hbar} \right) \left(\frac{Na_B}{c} \right) = N\alpha \quad \Rightarrow \quad \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

Laser Interferometry Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRA* **93**, 063630 (2016)]

- Compare $L \sim Na_B$ with λ (or a 2nd L)
- For a “usual” atomic optical transition and in the non-relativistic limit:

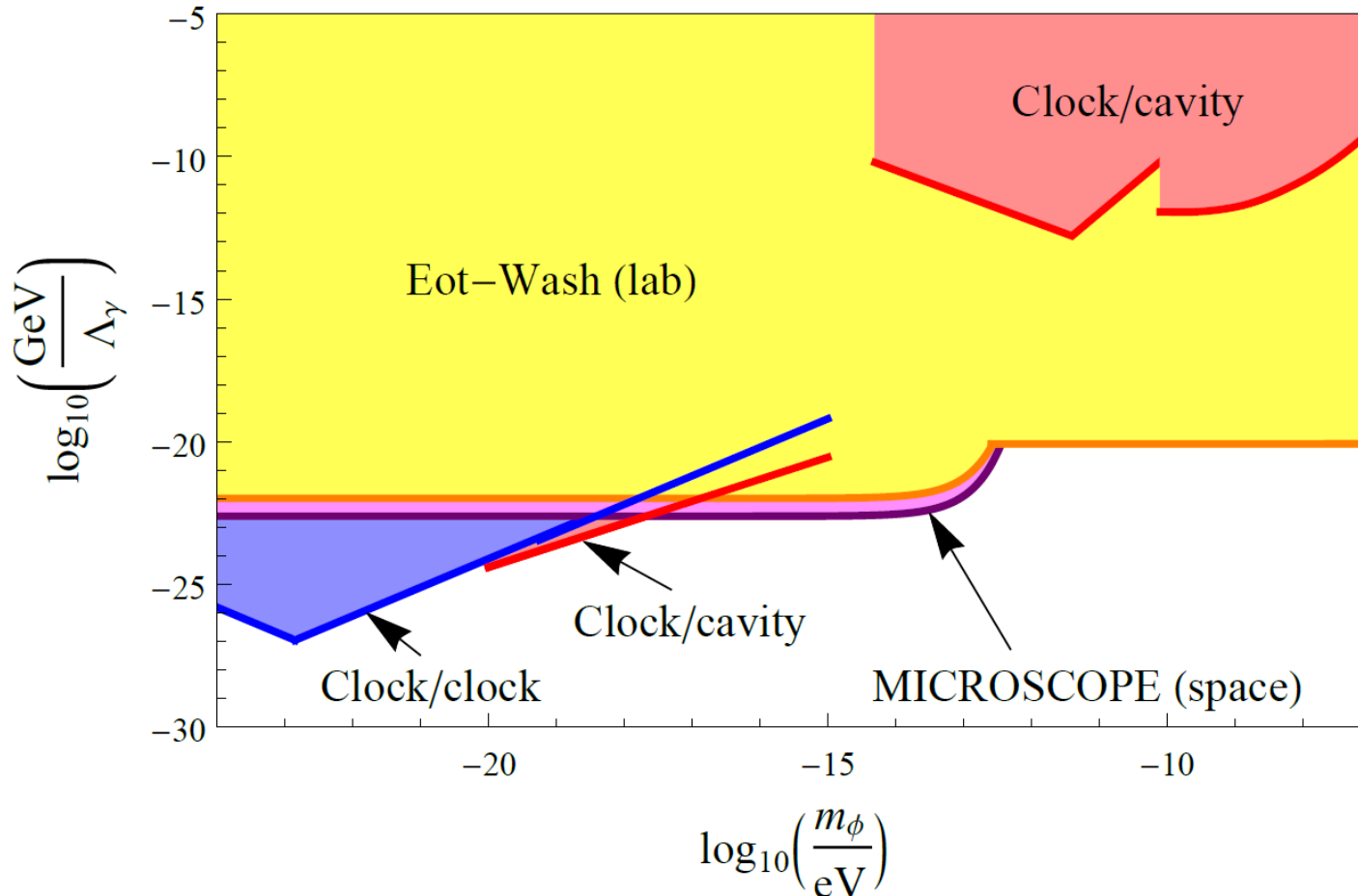
$$\Phi = \frac{\omega L}{c} \propto \left(\frac{e^2}{a_B \hbar} \right) \left(\frac{Na_B}{c} \right) = N\alpha \quad \Rightarrow \quad \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

- Multiple reflections of light beam enhance the effect ($N_{\text{eff}} \sim 10^5$ in small-scale interferometers with highly reflective mirrors; c.f. $N_{\text{eff}} \sim 100$ in LIGO/Virgo)

Constraints on Linear Interaction of Scalar Dark Matter with the Photon

Clock/clock constraints: [Van Tilburg *et al.*, *PRL* **115**, 011802 (2015)], [Hees *et al.*, *PRL* **117**, 061301 (2016)]; **Clock/cavity constraints:** [Robinson, Ye *et al.*, *Bulletin APS*, H06.00005 (2018)], [Aharony *et al.*, arXiv:1902.02788], [Antypas *et al.*, arXiv:1905.02968]

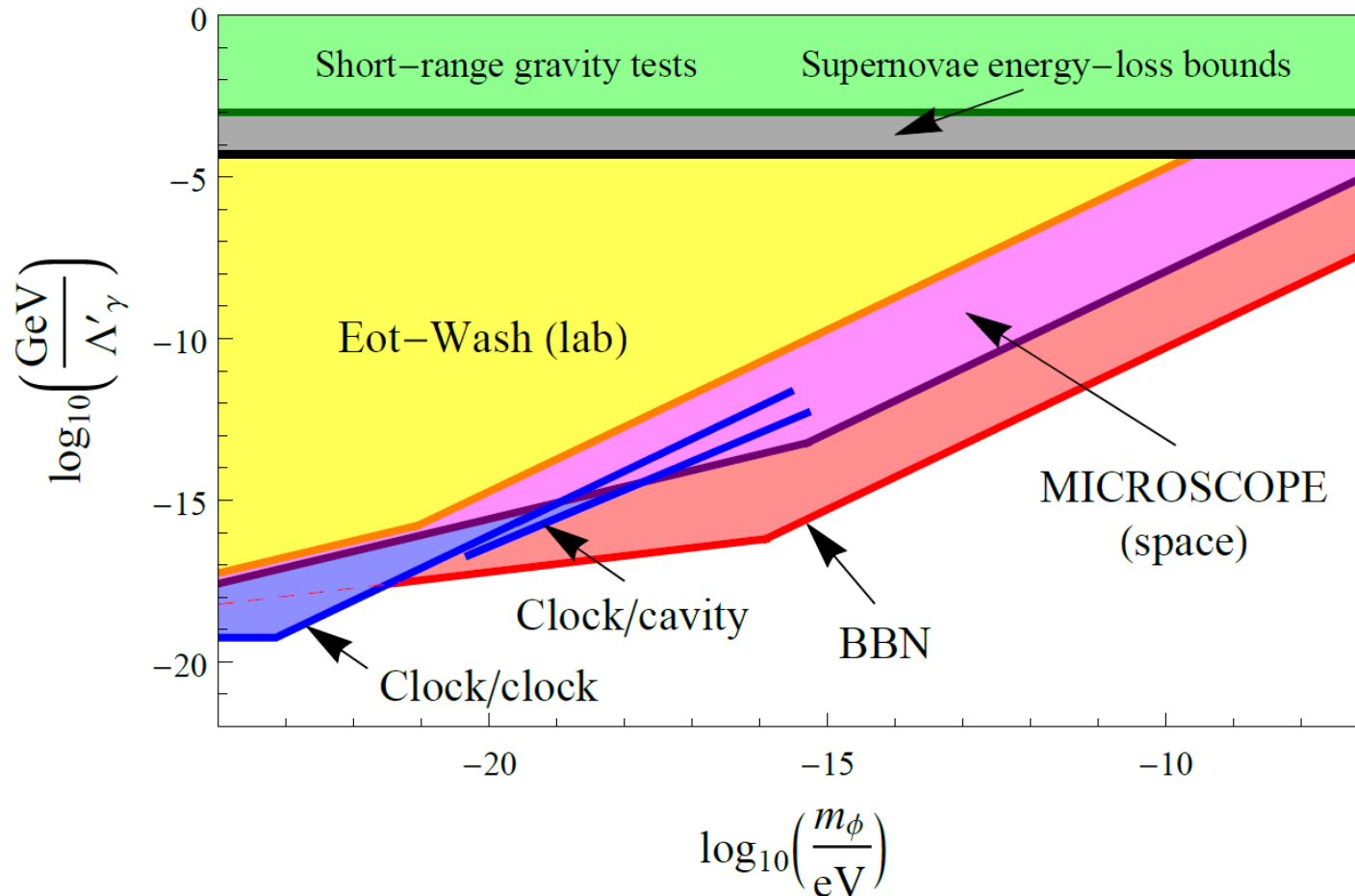
4 orders of magnitude improvement!



Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

Clock/clock + BBN constraints: [Stadnik, Flambaum, *PRL* **115**, 201301 (2015); *PRA* **94**, 022111 (2016)]; **MICROSCOPE + Eöt-Wash constraints:** [Hees *et al.*, *PRD* **98**, 064051 (2018)]

15 orders of magnitude improvement!



Low-mass Spin-0 Dark Matter

Dark Matter



QCD axion resolves
strong CP problem

**Pseudoscalars
(Axions):**

$$\varphi \xrightarrow{P} -\varphi$$

→ **Time-varying spin-
dependent effects**

- Co-magnetometers
- Nuclear magnetic resonance
 - Torsion pendula

“Axion Wind” Spin-Precession Effect

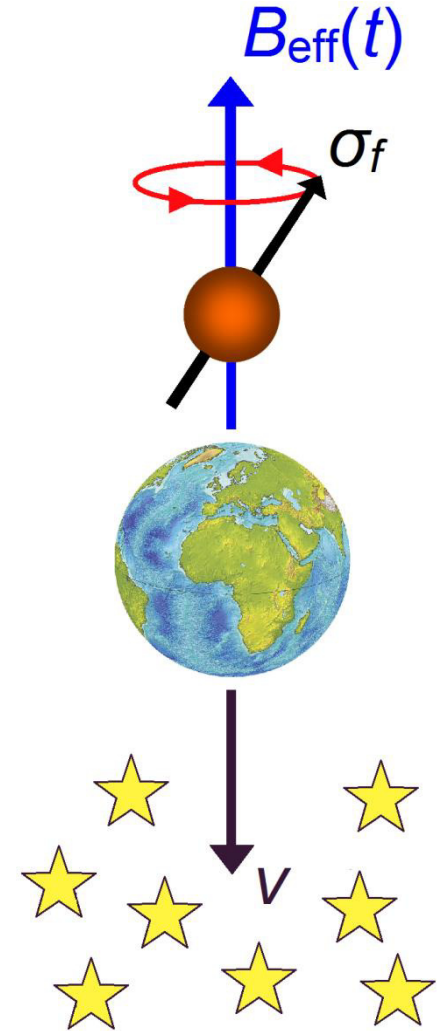
[Flambaum, talk at *Patras Workshop*, 2013], [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aff} = -\frac{C_f}{2f_a} \partial_i [a_0 \cos(\varepsilon_a t - \mathbf{p}_a \cdot \mathbf{x})] \bar{f} \gamma^i \gamma^5 f$$

$$\Rightarrow H_{\text{eff}}(t) \simeq \boldsymbol{\sigma}_f \cdot \mathbf{B}_{\text{eff}} \sin(m_a t)$$

Pseudo-magnetic field*

$$\mathbf{B}_{\text{eff}} \propto \mathbf{v}$$



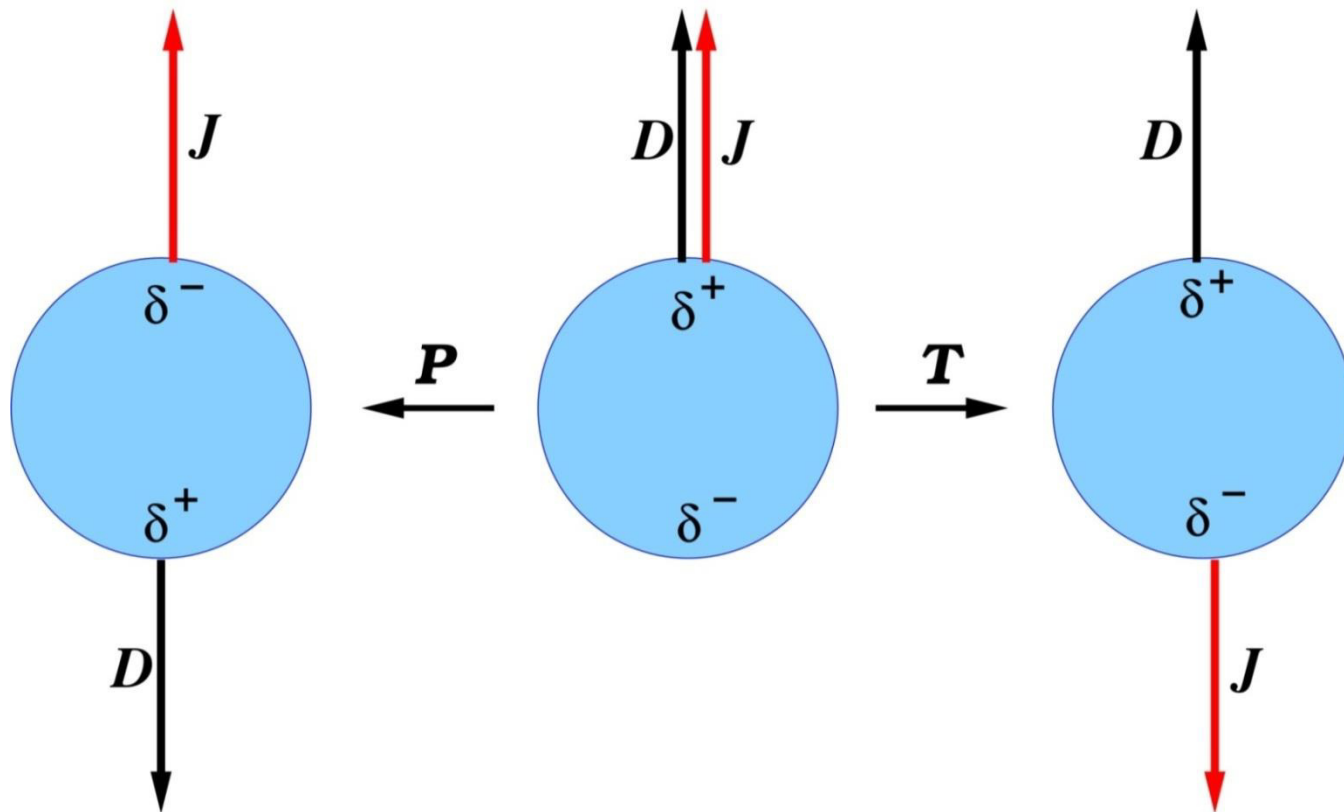
* Compare with usual magnetic field: $H = -\boldsymbol{\mu}_f \cdot \mathbf{B}$

Oscillating Electric Dipole Moments

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)]

Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Electric Dipole Moment (EDM) = parity (P) and time-reversal-invariance (T) violating electric moment



Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Searching for Spin-Dependent Effects

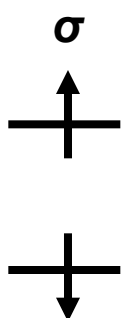
Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\text{Hg}}} = \left| \frac{\mu_n B}{\mu_{\text{Hg}} B} \right| + R(t)$$

Energy



B-field
effect

Axion DM
effect

Searching for Spin-Dependent Effects

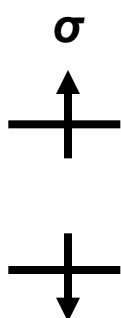
Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\text{Hg}}} = \left| \frac{\mu_n B}{\mu_{\text{Hg}} B} \right| + R(t)$$

Energy



B-field
effect

Axion DM
effect

Searching for Spin-Dependent Effects

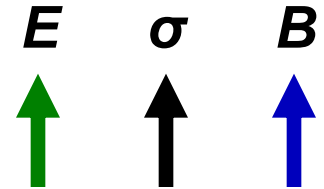
Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\text{Hg}}} = \left| \frac{\mu_n \cancel{B}}{\mu_{\text{Hg}} \cancel{B}} \right| + R(t)$$

$$R_{\text{EDM}}(t) \propto \cos(m_a t)$$



Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
 Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

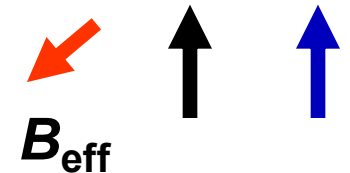
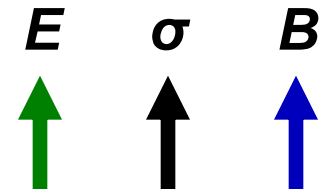
Use *spin-polarised sources*: Atomic magnetometers,
 ultracold neutrons, torsion pendula

Experiment (n/Hg): [nEDM collaboration, *PRX* **7**, 041034 (2017)]

$$\frac{\nu_n}{\nu_{\text{Hg}}} = \left| \frac{\mu_n \cancel{B}}{\mu_{\text{Hg}} \cancel{B}} \right| + R(t)$$

$$R_{\text{EDM}}(t) \propto \cos(m_a t)$$

$$R_{\text{wind}}(t) \propto \sum_{i=1,2,3} A_i \sin(\omega_i t)$$



$$\omega_1 = m_a, \quad \omega_2 = m_a + \Omega_{\text{sidereal}}, \quad \omega_3 = |m_a - \Omega_{\text{sidereal}}|$$

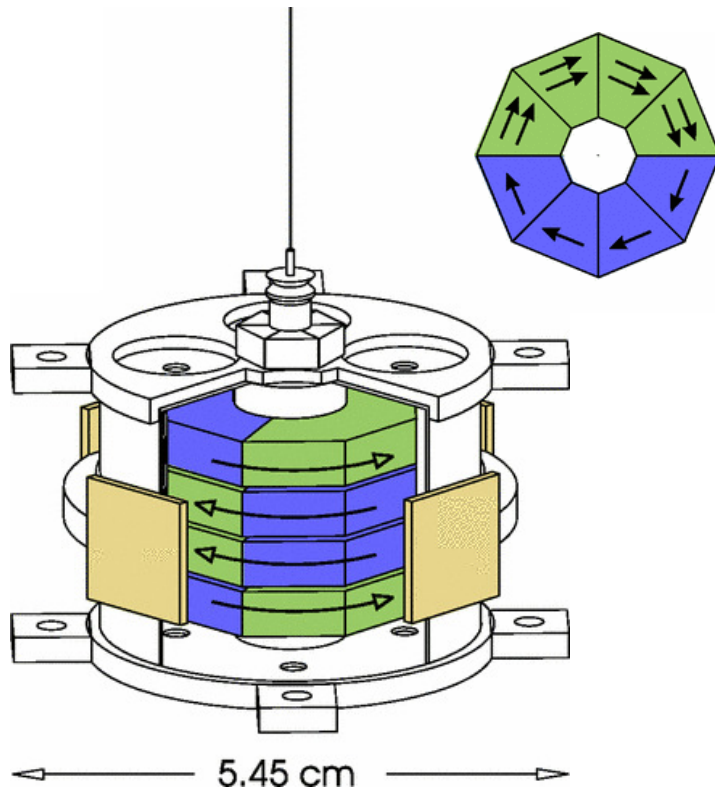


Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (Alnico/SmCo₅): [Terrano *et al.*, arXiv:1902.04246; *PRL* (In press)]

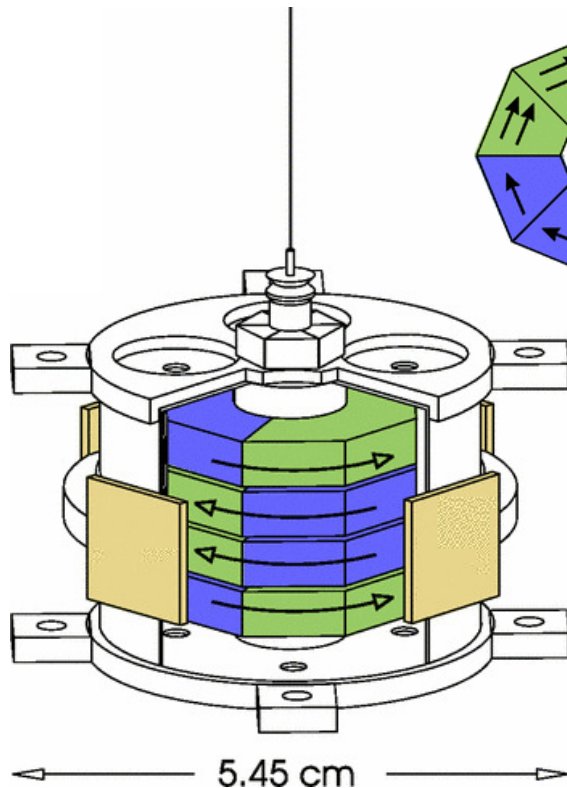


Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (Alnico/SmCo₅): [Terrano *et al.*, arXiv:1902.04246; *PRL* (In press)]



$$\mu_{\text{pendulum}} \approx 0$$

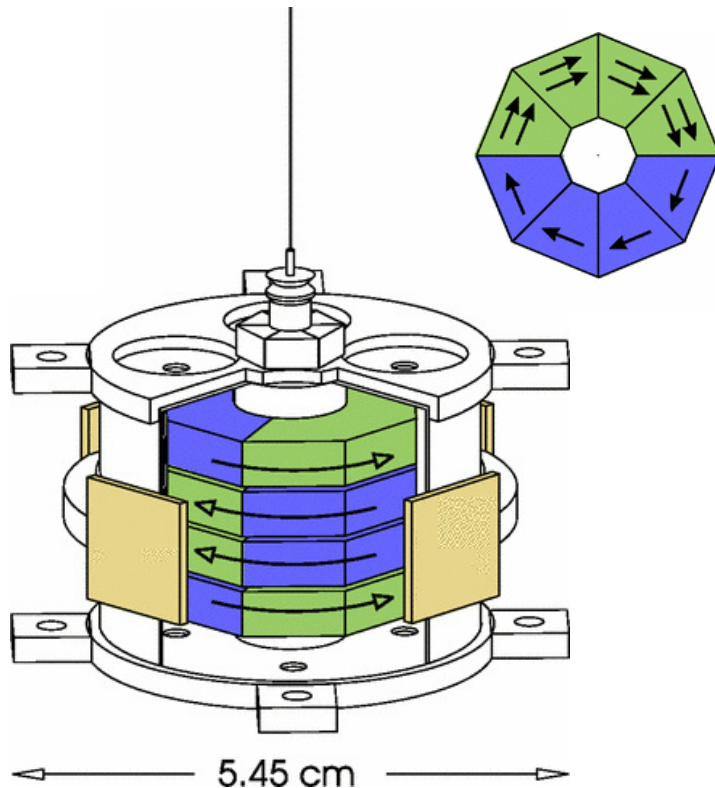
$$(\sigma_e)_{\text{pendulum}} \neq 0$$

Searching for Spin-Dependent Effects

Proposals: [Flambaum, talk at *Patras Workshop*, 2013;
Stadnik, Flambaum, *PRD* **89**, 043522 (2014); Stadnik, thesis (Springer, 2017)]

Use *spin-polarised sources*: Atomic magnetometers,
ultracold neutrons, torsion pendula

Experiment (Alnico/SmCo₅): [Terrano *et al.*, arXiv:1902.04246; *PRL* (In press)]



$$\mu_{\text{pendulum}} \approx \mathbf{0}$$

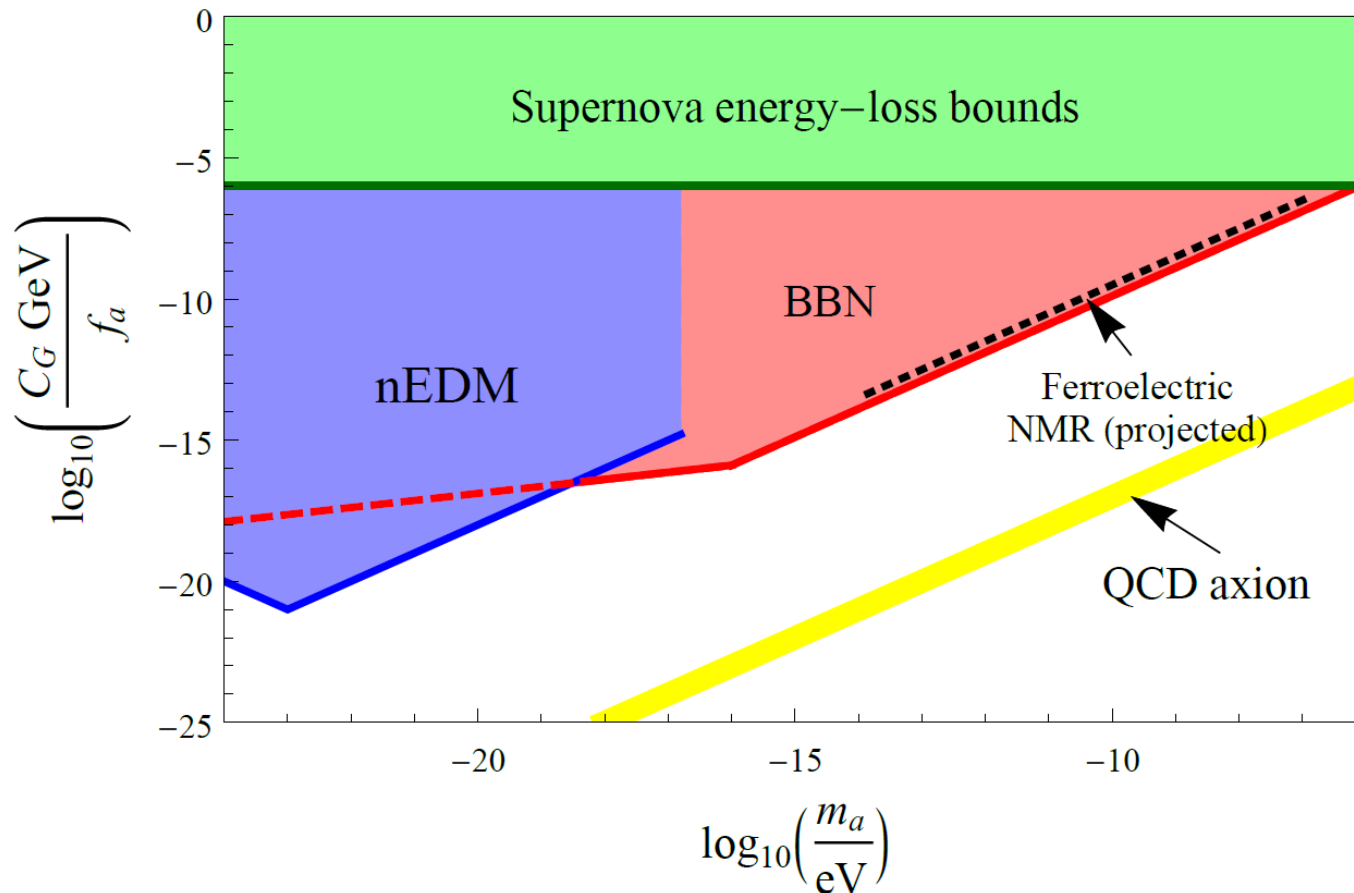
$$(\sigma_e)_{\text{pendulum}} \neq \mathbf{0}$$

$$\tau(t) \propto (\sigma_e)_{\text{pendulum}} \times B_{\text{eff}}(t)$$

Constraints on Interaction of Axion Dark Matter with Gluons

nEDM constraints: [nEDM collaboration, *PRX* 7, 041034 (2017)]

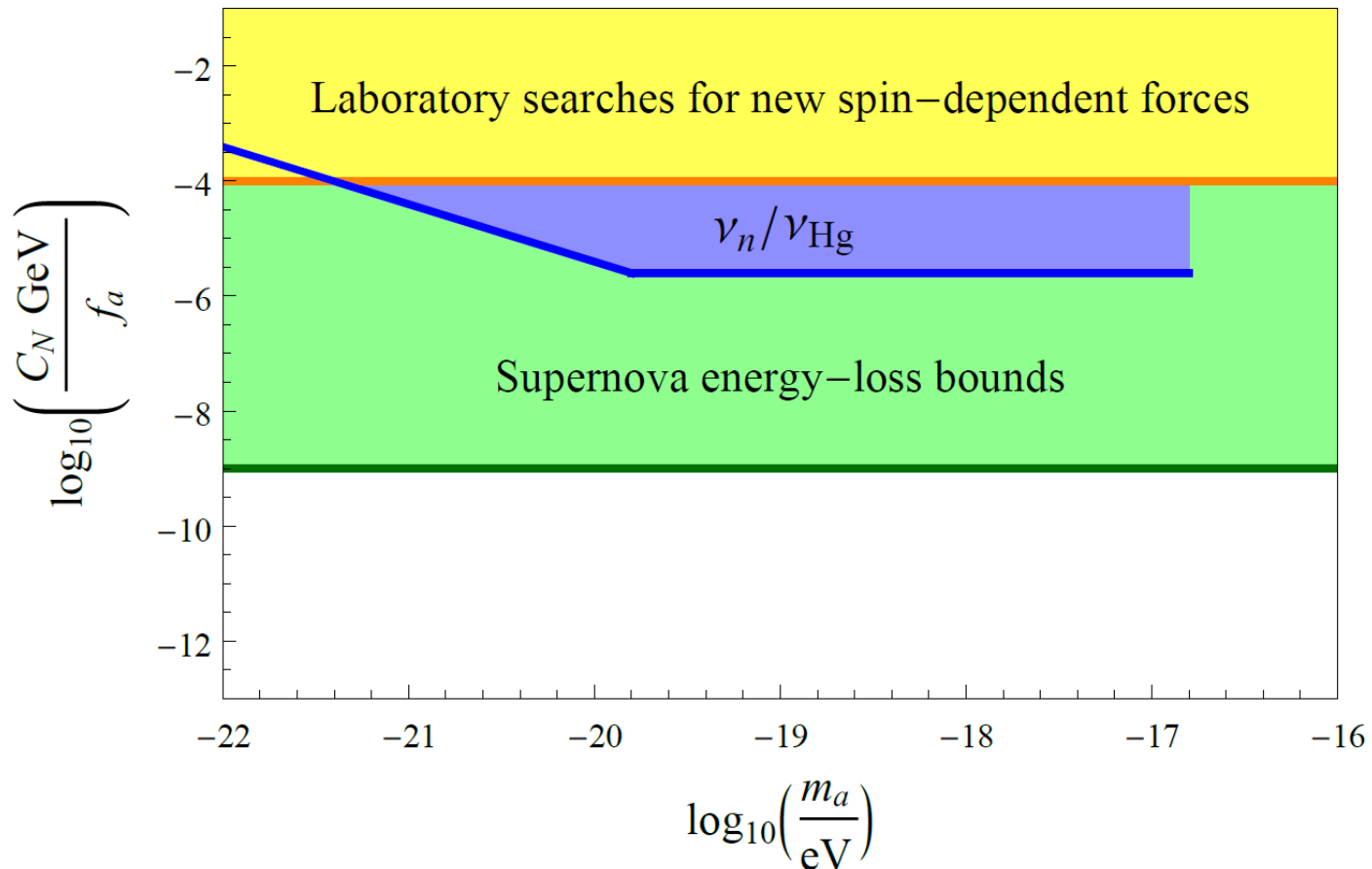
3 orders of magnitude improvement!



Constraints on Interaction of Axion Dark Matter with Nucleons

ν_n/ν_{Hg} constraints: [nEDM collaboration, *PRX* 7, 041034 (2017)]

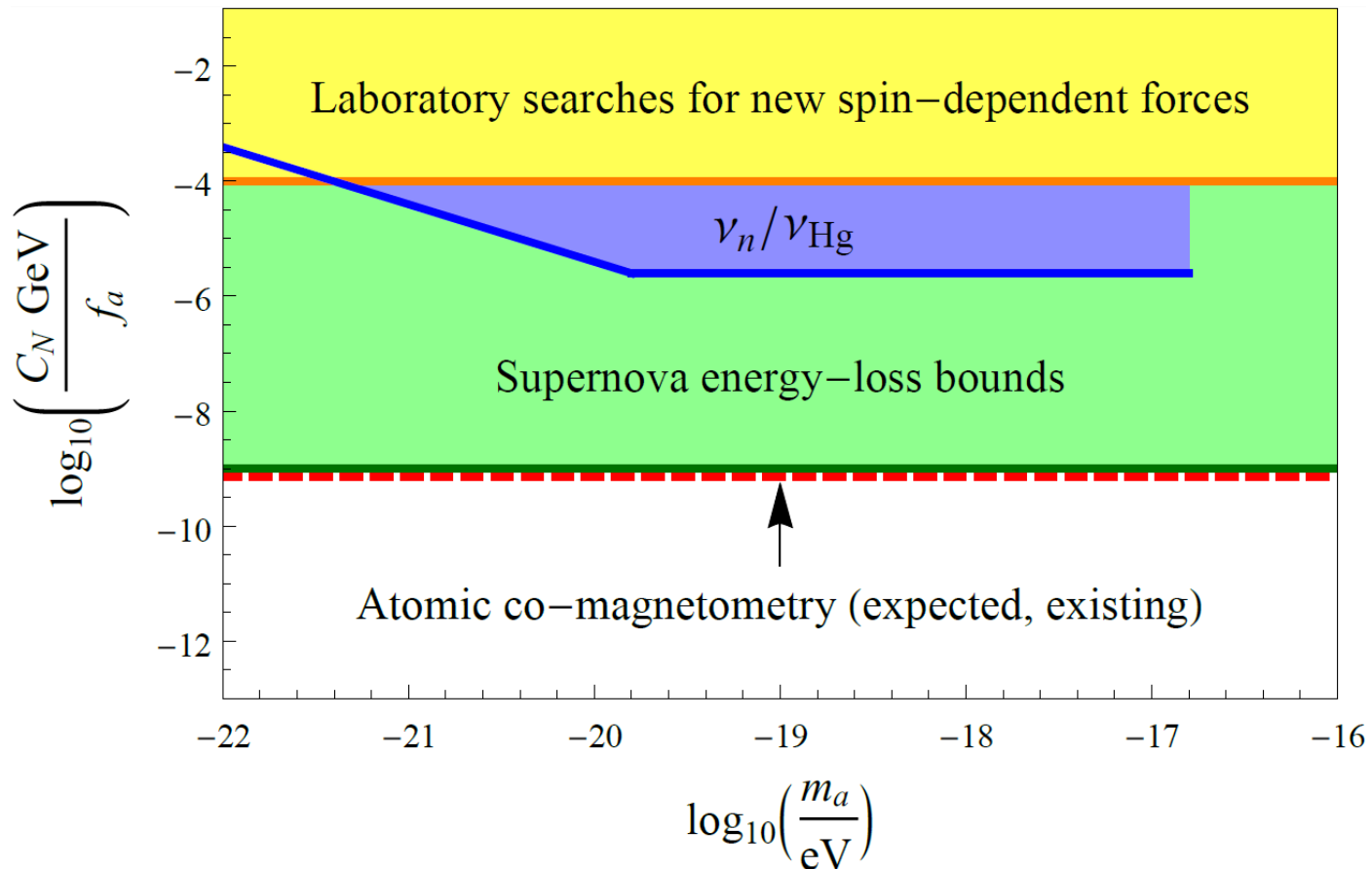
40-fold improvement (laboratory bounds)!



Constraints on Interaction of Axion Dark Matter with Nucleons

ν_n/ν_{Hg} constraints: [nEDM collaboration, *PRX* 7, 041034 (2017)]

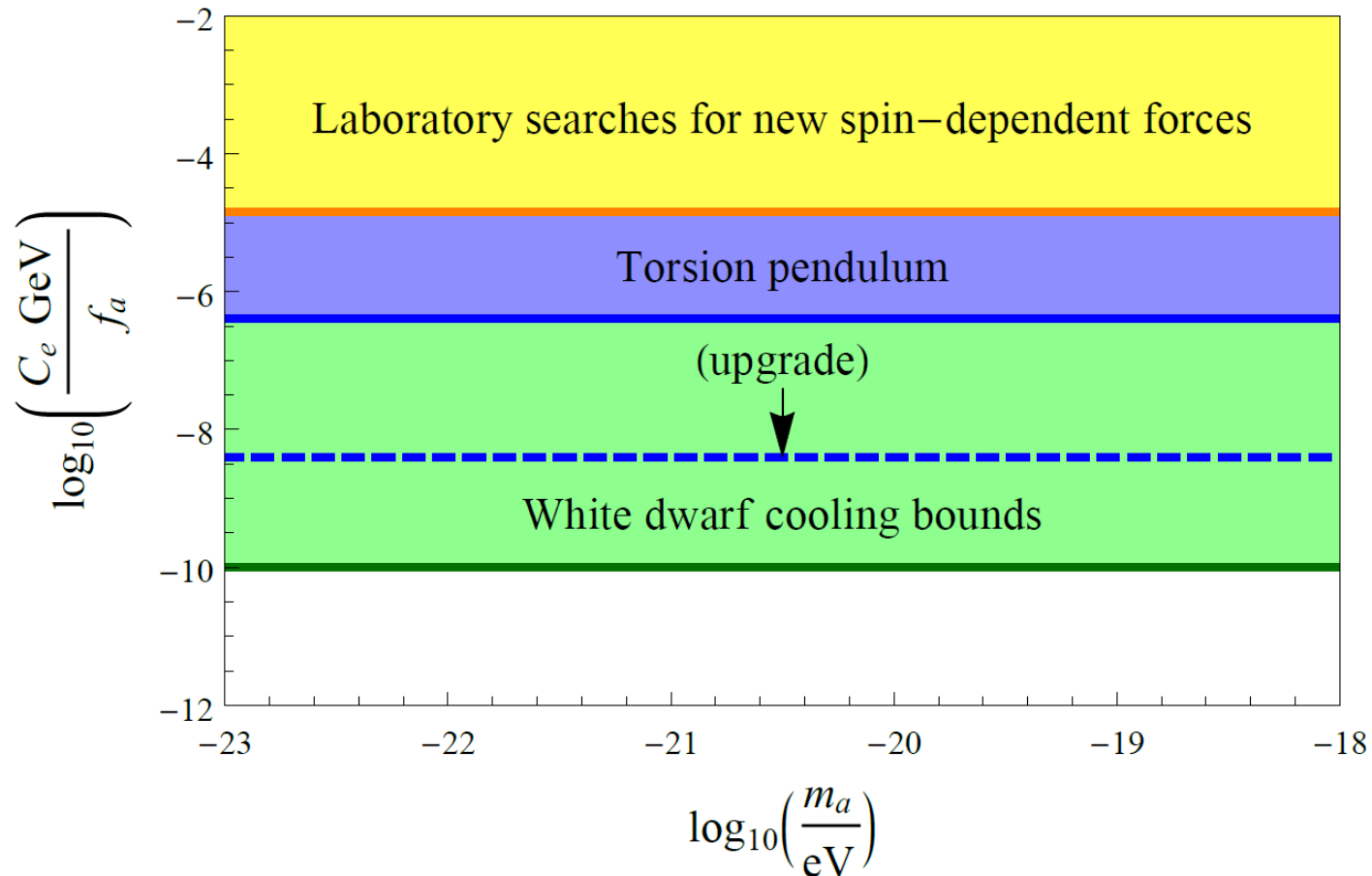
40-fold improvement (laboratory bounds)!



Constraints on Interaction of Axion Dark Matter with the Electron

Torsion pendulum constraints: [Terrano *et al.*, arXiv:1902.04246; *PRL* (In press)]

35-fold improvement (laboratory bounds)!



Summary

- New classes of dark-matter effects that are **first power** in the underlying interaction constant
=> Up to **15 orders of magnitude improvement** with precision, low-energy, table-top experiments:
 - Spectroscopy (clocks)
 - Cavities and interferometry
 - Magnetometry
 - Torsion pendula

⋮

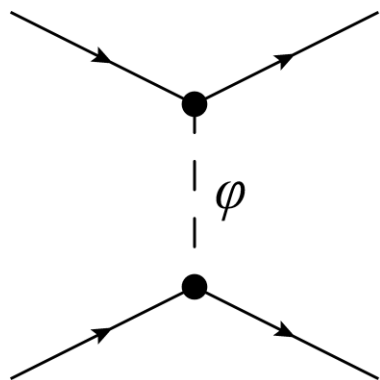
Back-up Slides

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

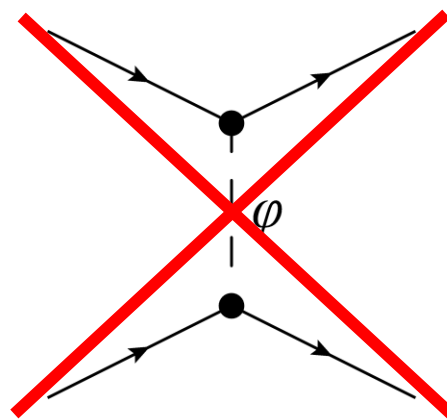
Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\phi\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

Quadratic couplings ($\phi^2\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) \left(1 - \frac{B}{r} \right)$$



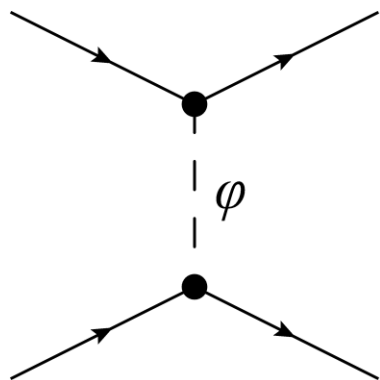
Gradients + screening/amplification

Fifth Forces: Linear vs Quadratic Couplings

[Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

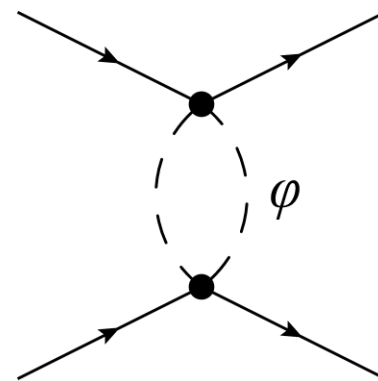
Consider the effect of a massive body (e.g., Earth) on the scalar DM field

Linear couplings ($\phi\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r}$$

Quadratic couplings ($\phi^2\bar{X}X$)



$$\phi = \phi_0 \cos(m_\phi t) \left(1 - \frac{B}{r} \right) - C \frac{e^{-2m_\phi r}}{r^3}$$



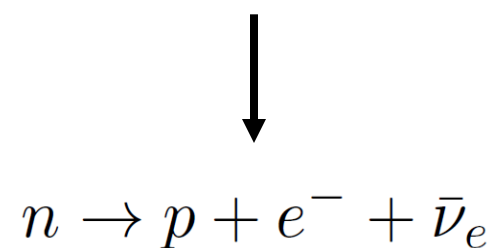
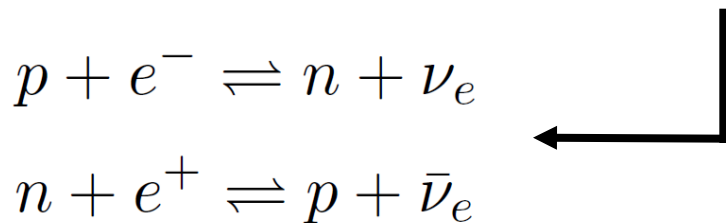
Gradients + screening/amplification

BBN Constraints on 'Slow' Drifts in Fundamental Constants due to Dark Matter

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

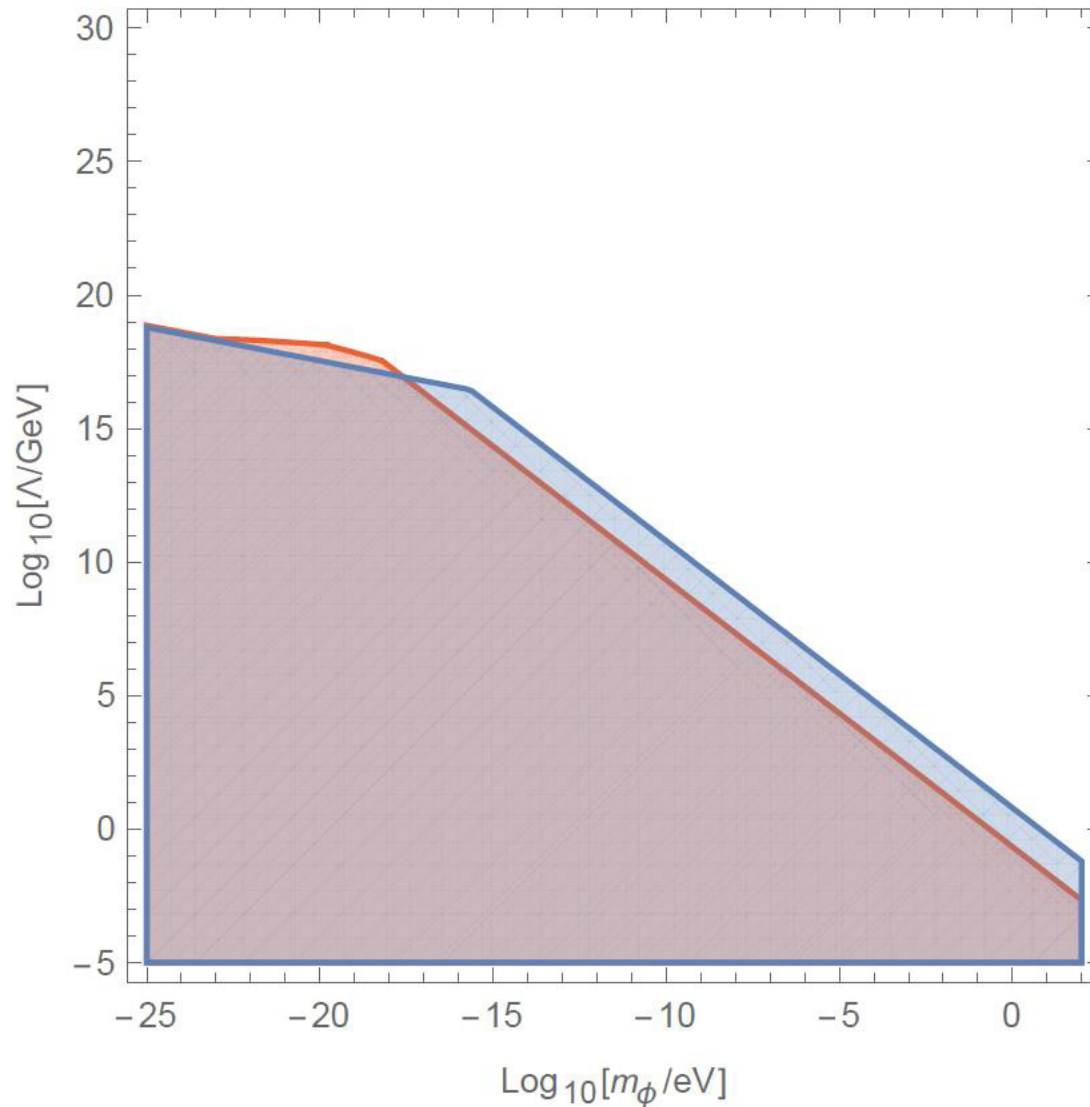
- Largest effects of DM in early Universe (highest ρ_{DM})
- Big Bang nucleosynthesis ($t_{\text{weak}} \approx 1\text{s} - t_{\text{BBN}} \approx 3\text{ min}$)
- Primordial ${}^4\text{He}$ abundance sensitive to n/p ratio
(almost all neutrons bound in ${}^4\text{He}$ after BBN)

$$\frac{\Delta Y_p({}^4\text{He})}{Y_p({}^4\text{He})} \approx \frac{\Delta(n/p)_{\text{weak}}}{(n/p)_{\text{weak}}} - \Delta \left[\int_{t_{\text{weak}}}^{t_{\text{BBN}}} \Gamma_n(t) dt \right]$$



Back-Reaction Effects in BBN

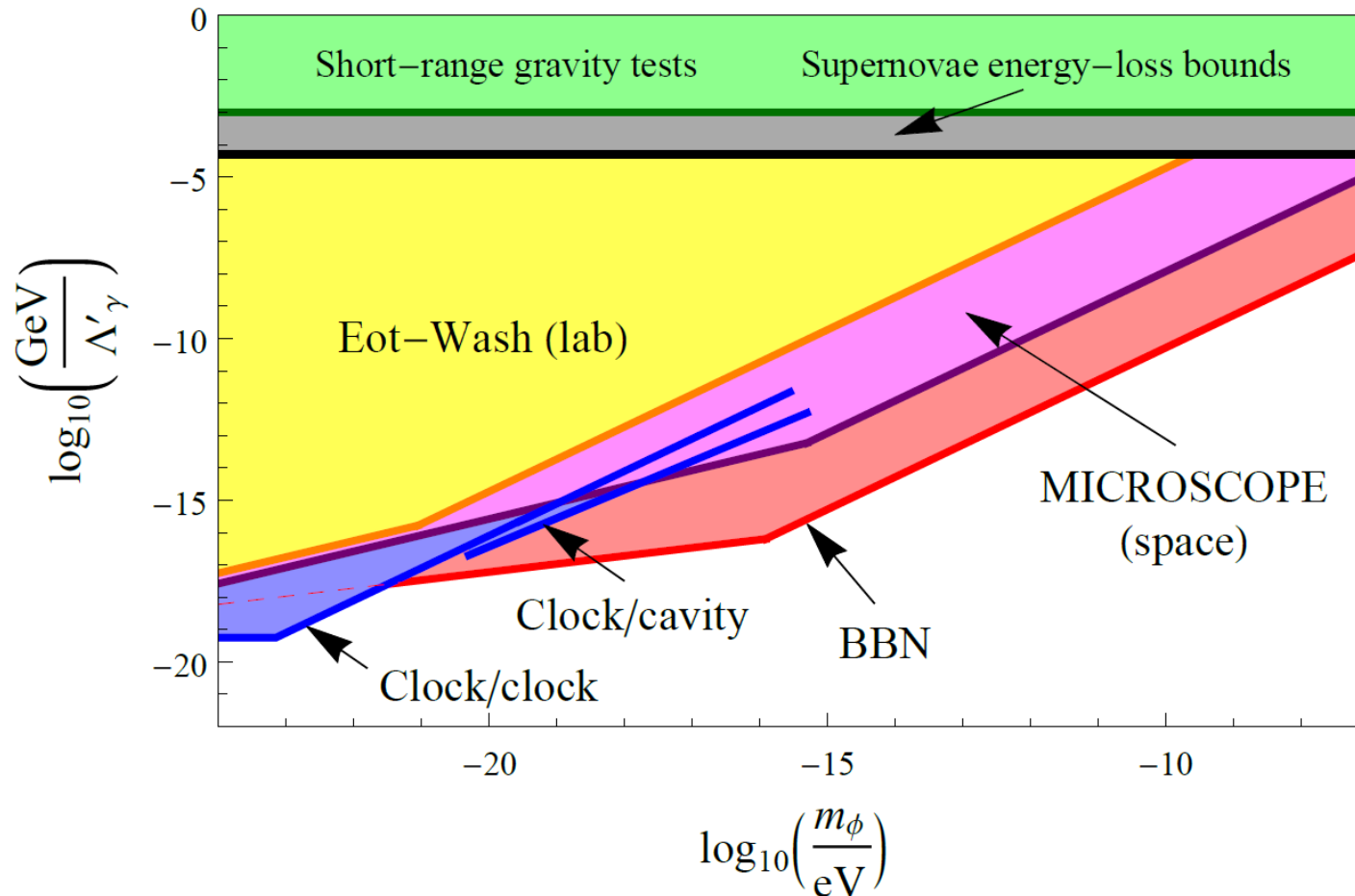
[Sörensen, Sibiryakov, Yu, PRELIMINARY – In preparation]



Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

Clock/clock + BBN constraints: [Stadnik, Flambaum, *PRL* **115**, 201301 (2015); *PRA* **94**, 022111 (2016)]; MICROSCOPE + Eöt-Wash constraints: [Hees *et al.*, *PRD* **98**, 064051 (2018)]

15 orders of magnitude improvement!

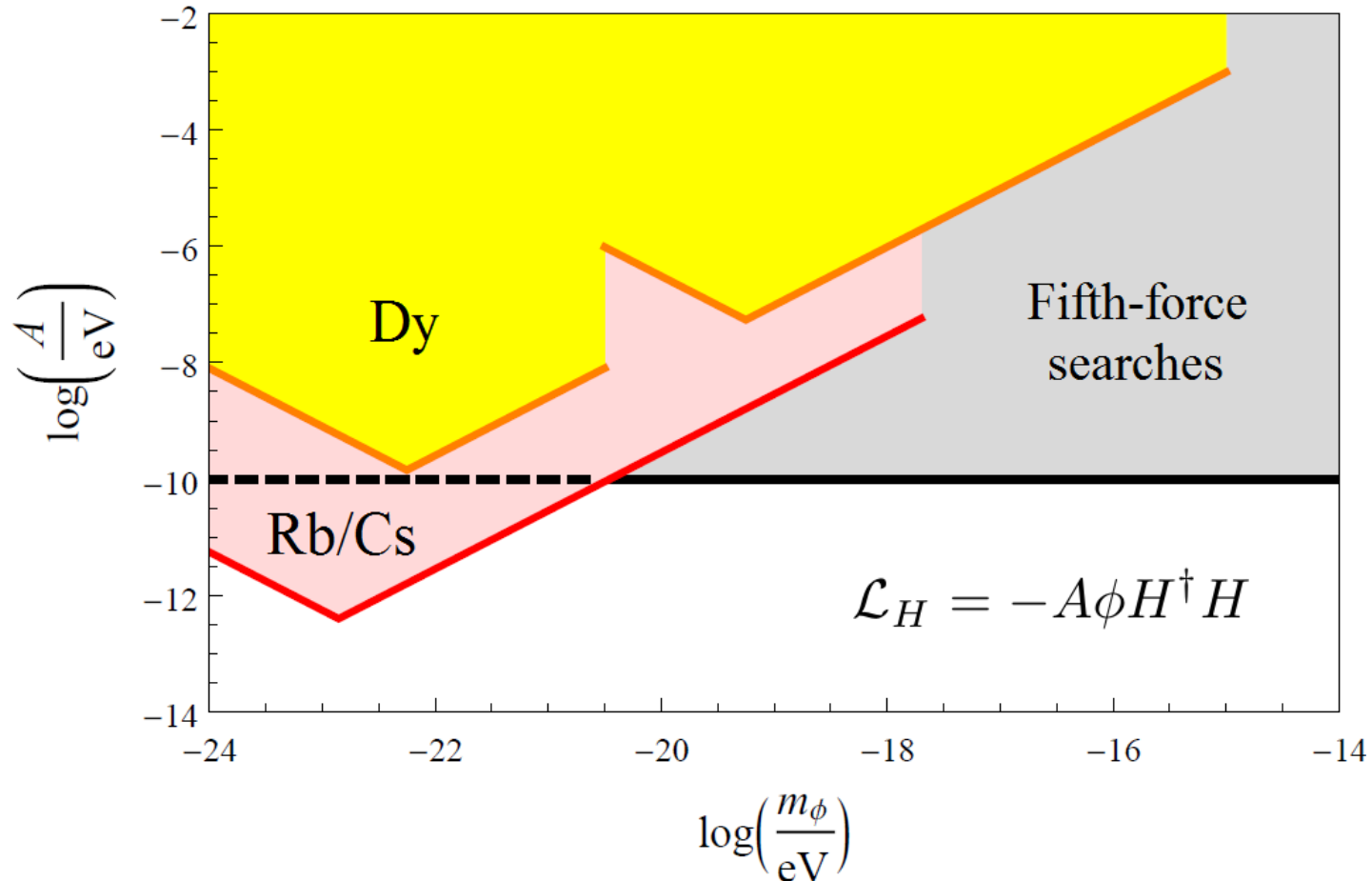


Constraints on Linear Interaction of Scalar Dark Matter with the Higgs Boson

Rb/Cs constraints:

[Stadnik, Flambaum, *PRA* **94**, 022111 (2016)]

2 – 3 orders of magnitude improvement!



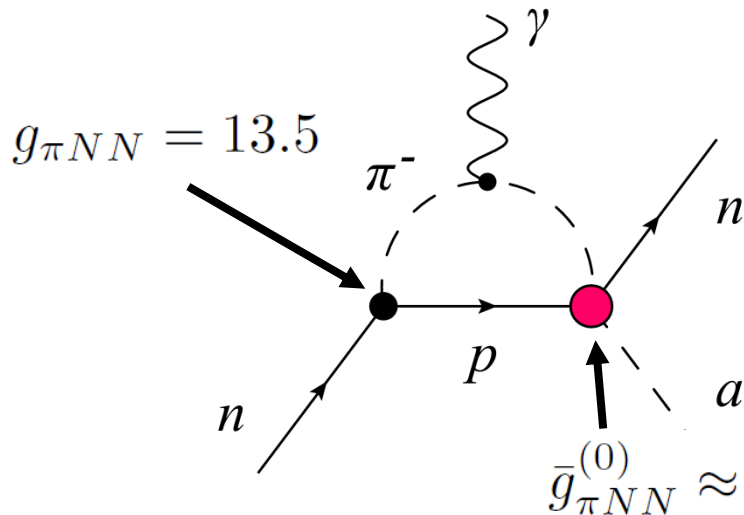
Oscillating Electric Dipole Moments

Nucleons: [Graham, Rajendran, *PRD* **84**, 055013 (2011)]

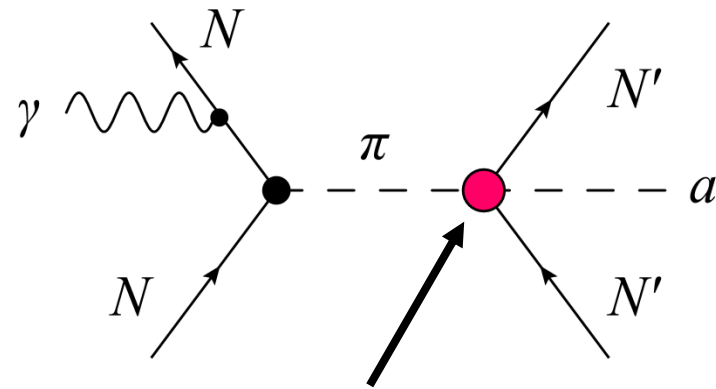
Atoms and molecules: [Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

$$\mathcal{L}_{aGG} = \frac{C_G a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G_{\mu\nu}^a \tilde{G}^{a\mu\nu}$$

Nucleon EDMs



***CP*-violating intranuclear forces**



In nuclei, tree-level *CP*-violating intranuclear forces dominate over loop-induced nucleon EDMs (loop factor = $1/(8\pi^2)$).