

Toward spatially uniform pulse compression of top-hat beams at the subpetawatt level of peak powers

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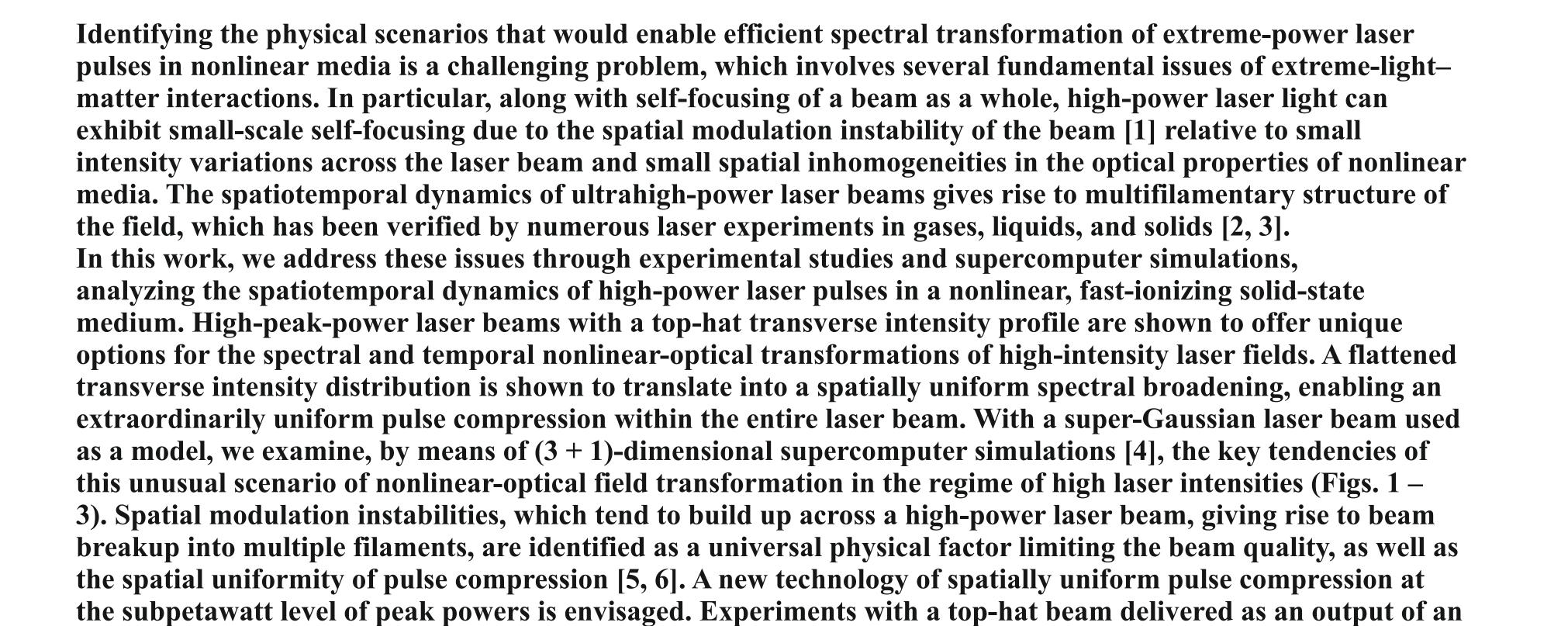


Kurchatov terawatt femtosecond laser system

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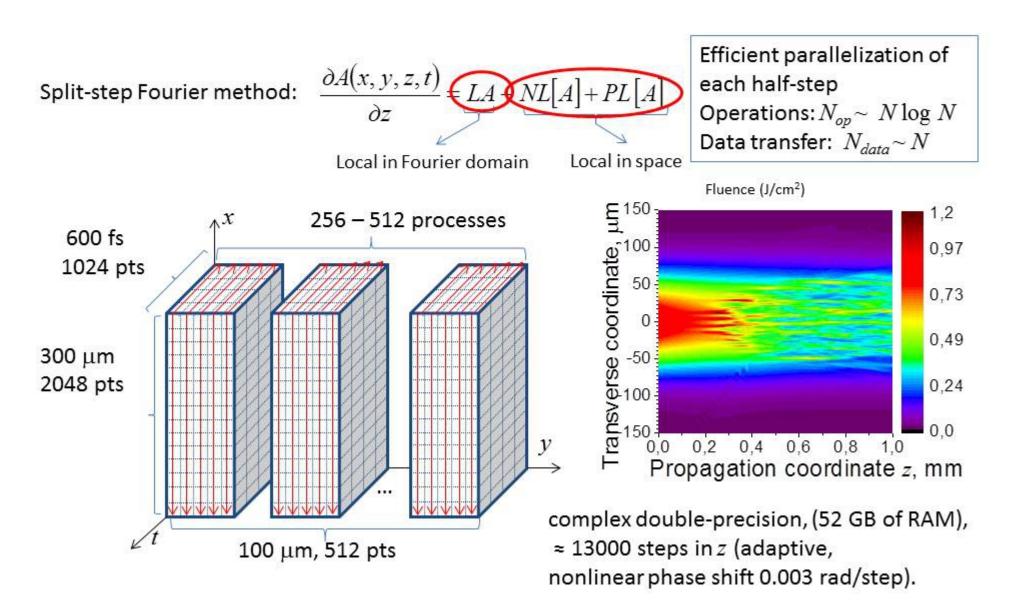
High-peak-power laser beams with a top-hat transverse intensity profile are shown to offer unique options for the spectral and temporal nonlinear-optical transformations of high-intensity laser fields, promising a new technology of spatially uniform pulse compression at the subpetawatt level of peak powers.

> **Kurchatov terawatt femtosecond laser system:** output beam profile

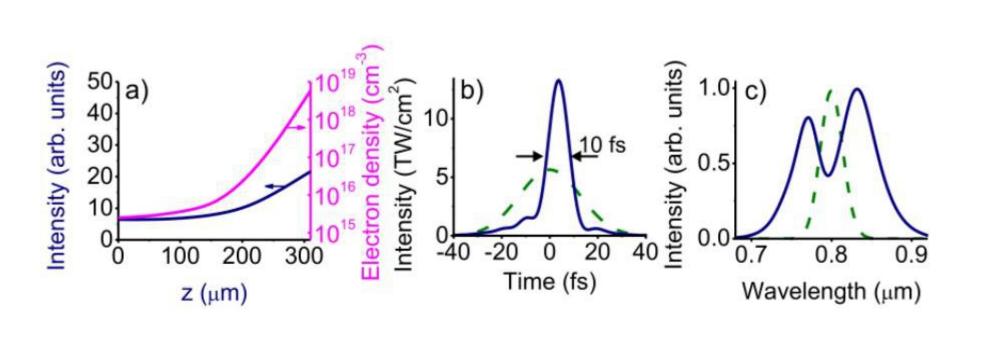


extreme-intensity short-pulse laser source put in operation at Kurchatov Institute are presented.

Numerical calculations



Pulse compression of super-Gaussian beam



Results of (3 + 1)-dimensional supercomputer simulations of nonlinear-optical evolution of a high-intensity fifth-order super-Gaussian beam: (a) the maximum field intensity (blue line) and electron density (red line) across the laser beam, (b) the temporal envelope, (c) the spectrum of the compressed (solid blue line) and input (dashed green line) laser pulse.

References

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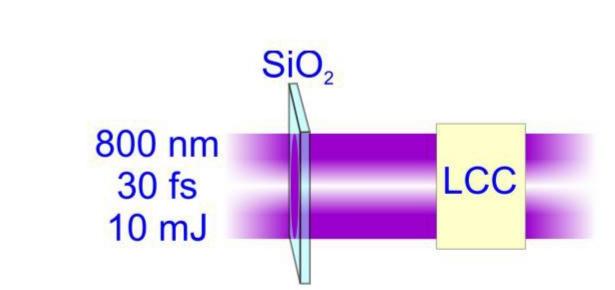
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compression of subpetawatt laser pulses to relativistic-intensity subcycle field waveforms, "Phys. Rev. A 86, 013835 (2012). [6] A. V. Mitrofanov, A. A. Voronin, D. A. Sidorov-Biryukov, S. I. Mitryukovsky, A. B. Fedotov, E. E. Serebryannikov, D. V. Meshchankin, V. Shumakova, S. Ališauskas, A.

Pugžlys, V. Ya. Panchenko, A. Baltuška, and A. M. Zheltikov, "Subterawatt few-cycle

Numerical Experiment



P>>Pcr P=1.7*10⁵ P_{cr} $P_{cr}(SiO_2) = 2MW$

Lomonosov supercomputer $1 - 10 * 10^{15}$ flop

Equation for the electron density

pulse self-steepening

Field evolution equation includes

beam diffraction

dispersion of the medium

susceptibilities of neutral gas

ionization-induced nonlinearities

plasma loss, refraction, and dispersion

photoionization rate W calculated using the Popov--Perelomov--Terentyev version of the Keldysh formalism

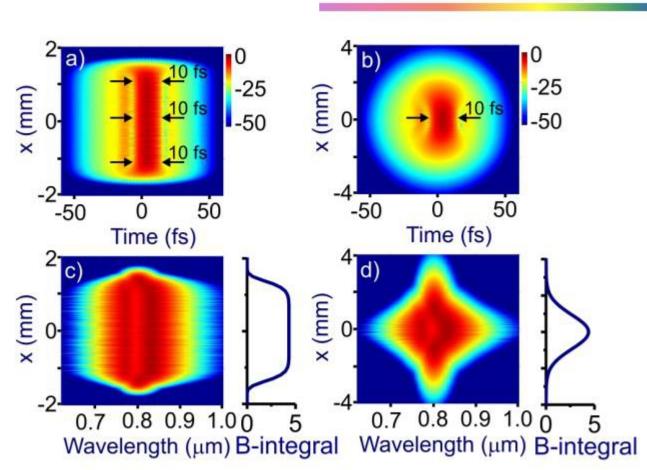
 $\frac{\partial}{\partial z}A(\omega,x,y,z) = \left[\frac{ic}{2\omega n_0}\Delta_{\perp} + i\tilde{D}(\omega)\right]A(\omega,x,y,z) + \tilde{F}\left[i\frac{\omega_0\tilde{T}}{c}n_2IA(t,x,y,z)\right]$

 $-\frac{U_{i}W(I)}{2I}A(t,x,y,z)\left]-\left(\frac{i\omega_{0}^{2}\omega}{2cn_{0}\rho_{c}\left(\omega^{2}+\tau_{c}^{-2}\right)}+\frac{\sigma(\omega)}{2}\right)\tilde{F}\left[\rho(t)A(t,x,y,z)\right]$

optical nonlinearities due to the cubic and quintic

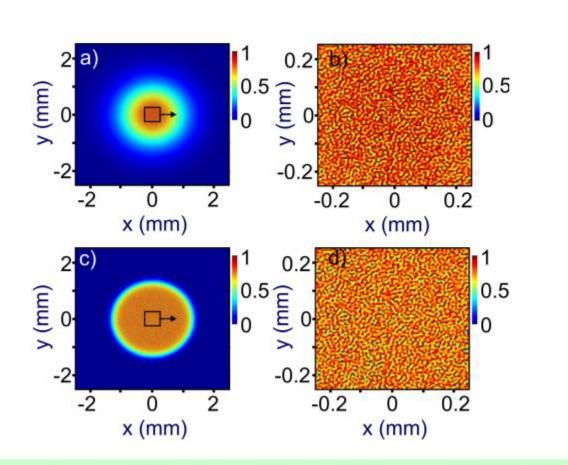
(3 + 1)-dimensional GNSE model

Supercomputer simulations of temporal and spectral evolution in super-Gaussian and a Gaussian beam



Temporal (a, b) and spectral (c, d) evolution in a fifth-order super-Gaussian (a, c) and a Gaussian (b, d) beam. The B integral across the beam is shown as a marginal

Hot-spot generation and beam breakup into multiple filaments for a Gaussian and super-Gaussian beam



-0.1 0.0 0.1 $x/(\lambda f)$ (mm⁻¹)

Results of (3 + 1)-dimensional supercomputer simulations showing hot-spot generation and beam breakup into multiple filaments for a Gaussian (a, b) and a fifth-order super-Gaussian (c, d) beam. The blowup of the central part of the beam is shown in panels (b) and (d).











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midinfrared pulses from a single filament," Optica 3, 299-302 (2016).