

New physics?



Bulk

SEARCH 2016

Matt Strassler (Harvard)

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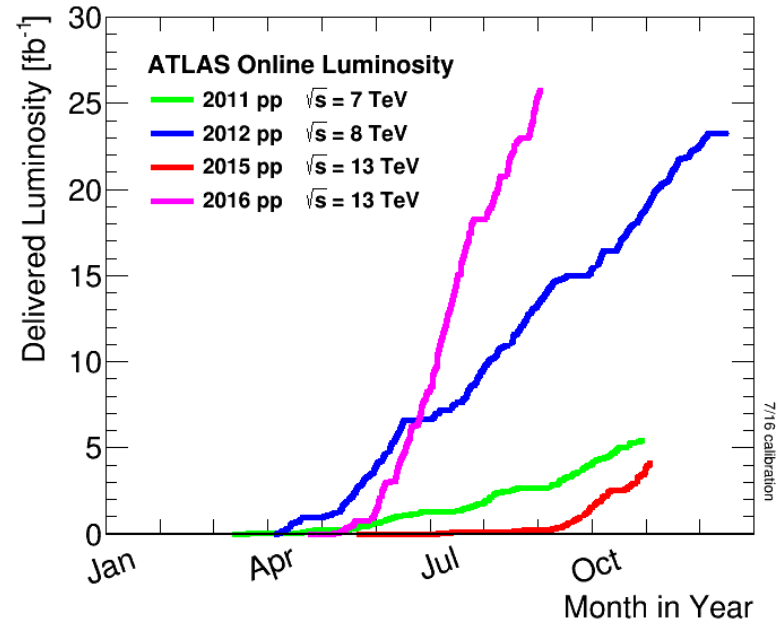
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Slow-Growth Era of LHC

- Statistical improvements will be slow, starting soon... well, now...

How to make rapid progress?

- Unconventional searches
 - *David Curtin's talk*
 - If it's never been done before... (or not since 2011)... who knows?
- Precision (or "high sensitivity")
 - Reducing theoretical & systematic uncertainties on existing searches
 - Planning new searches that weren't previously worth doing



Bulk of a Distribution?

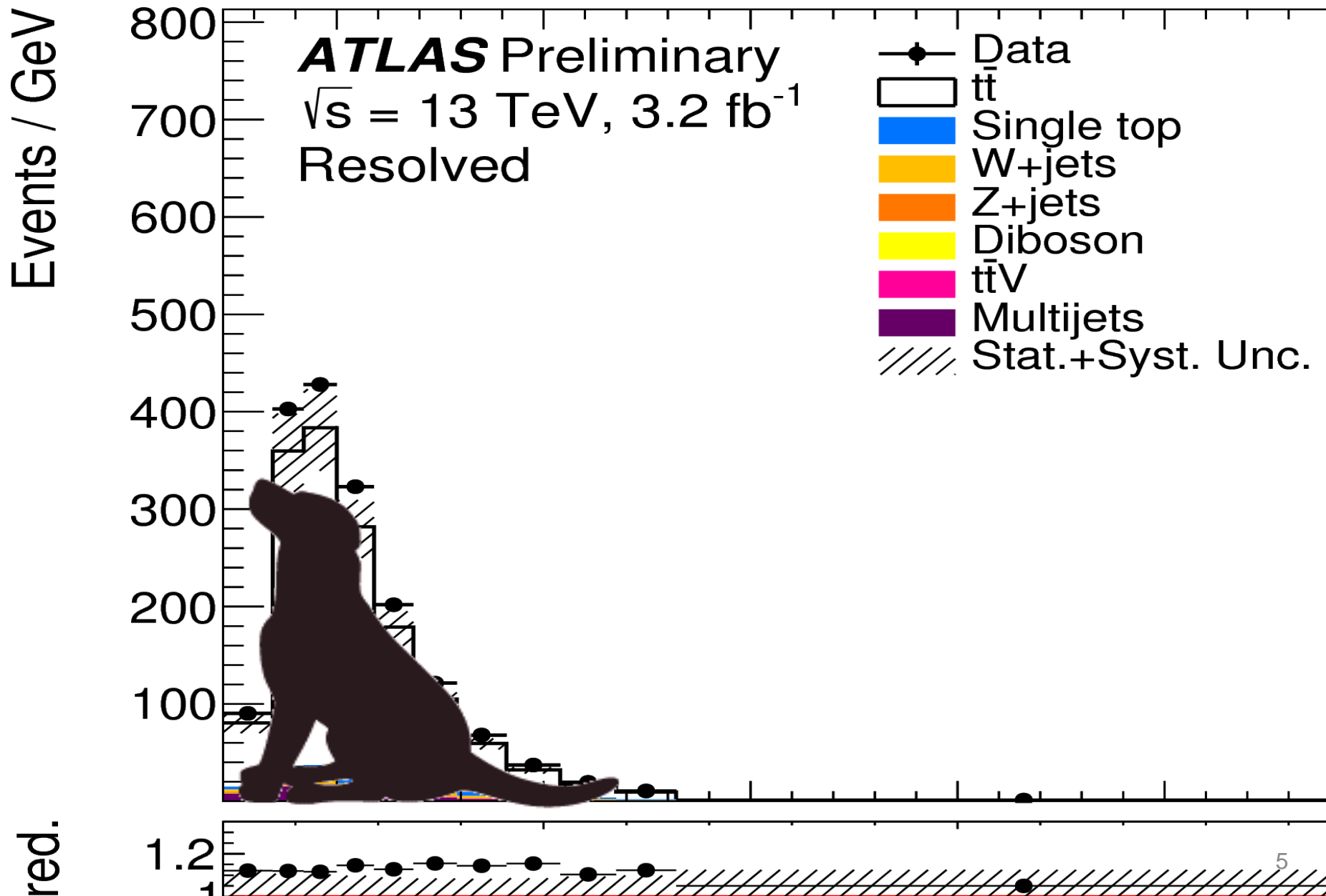
Here:

- Project full space of final states onto a variable or two
- Don't just look on the tail for deviations from prediction

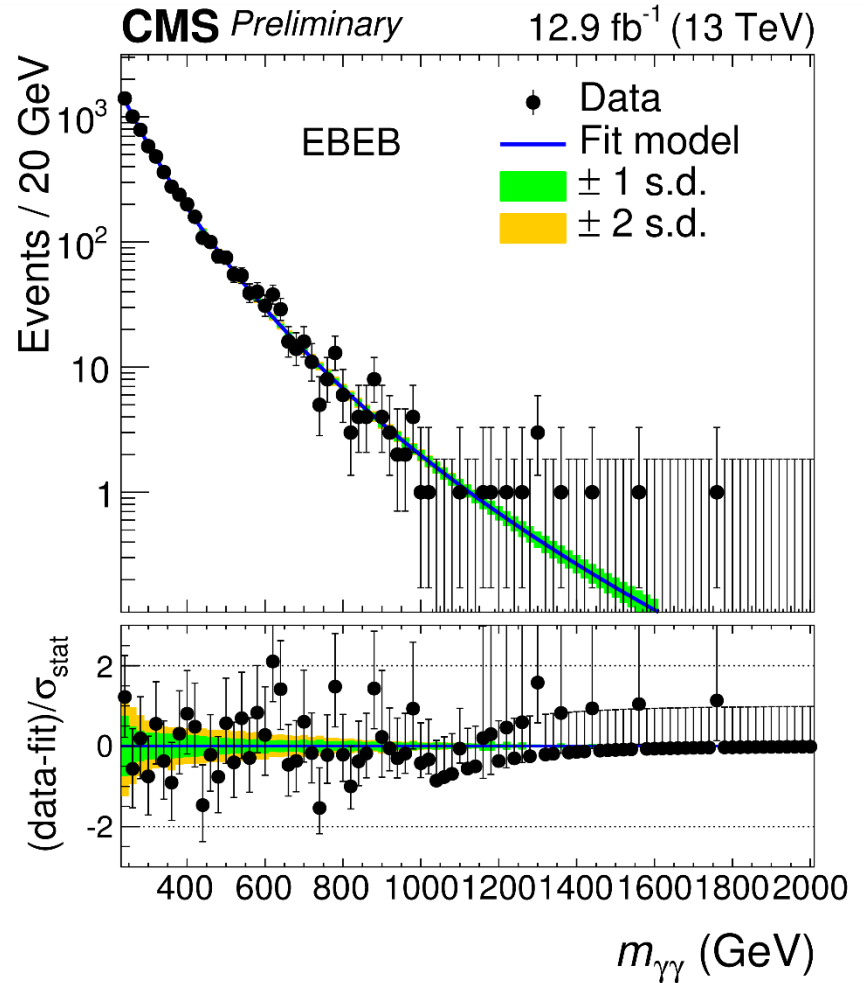
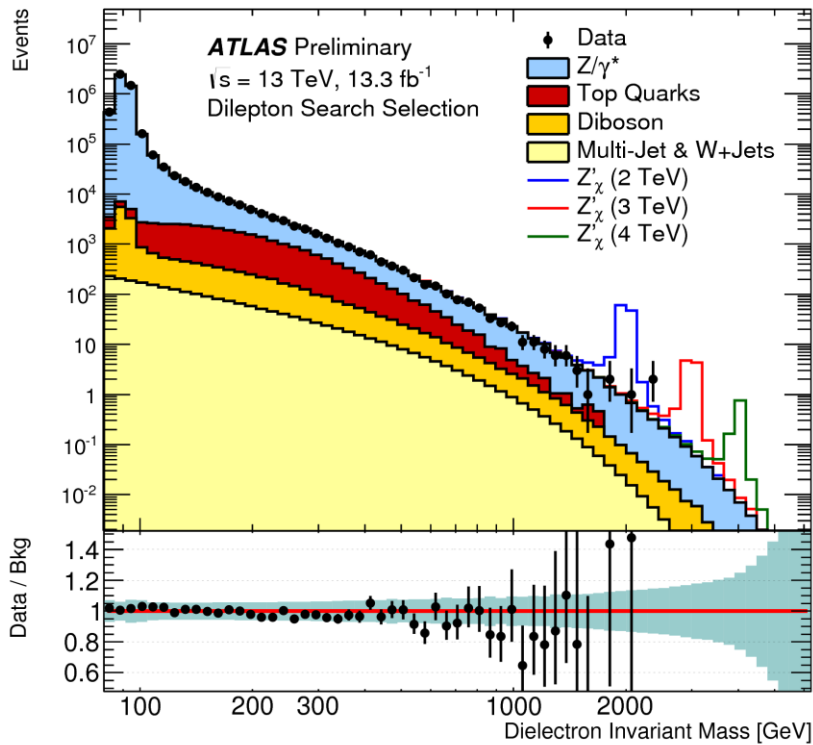
Reasons:

- While there has been a lot of focus for searches on
 - High mass
 - High p_T
 - High METof course many phenomena can show up lower in the plot
- Especially true now as limits on heavy particles begin to max out
 - Precision era means access to EW production, other rare processes
 - SM precision tests (needed for later BSM) demand high statistics

Don't Let the Tail Wag the Dog



Inclusive is not Conclusive



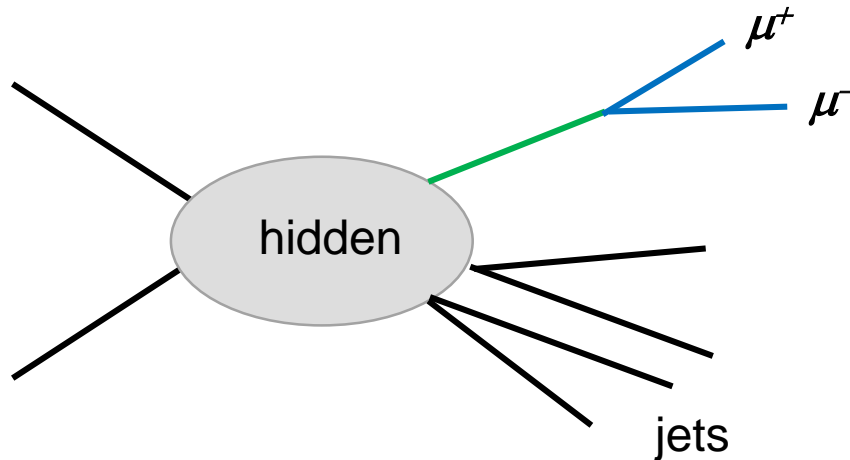
Physics Hidden in Inclusive Distributions

In various models (*common in hidden valleys/dark sectors*)

- Rare or unusual production of new neutral particle
- Bump, edge, endpoint, dip, wiggle is present
 - but swamped in inclusive background

MJS & Zurek '06

Han, Si, Zurek & MJS '07



Rare, prompt, light dilepton resonance along with hard jets

This happens in the SM too

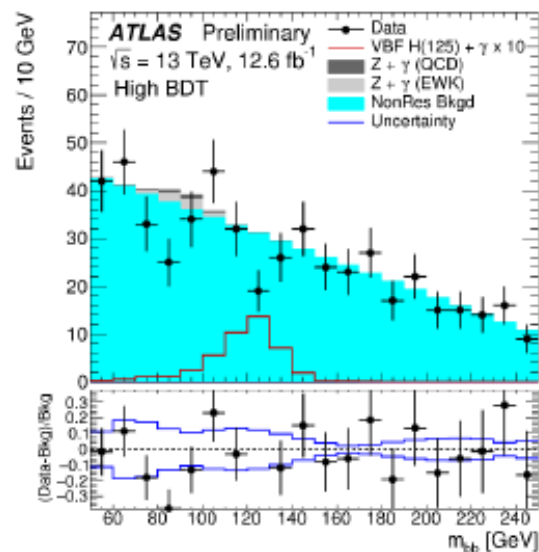
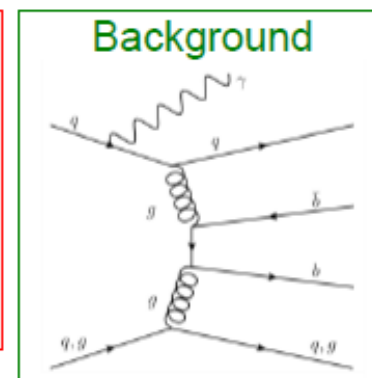
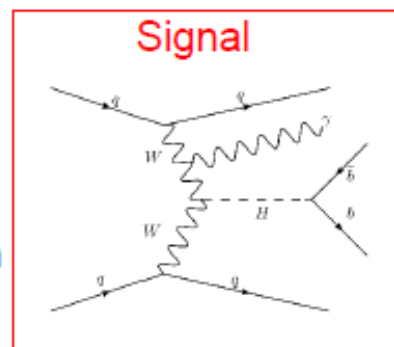
Higgs to bb can't be seen inclusively either

- So do a **semi-exclusive** search... cut background much more than signal
 - VH (require MET/lepton, maybe mild boost)
 - Or... from Guillemin talk

ATLAS-CONF-2016-063

Analysis key points

- $H \rightarrow bb$ VBF analysis (not in association with a photon) performed in Run 1 (sensitivity ~ 5 times the SM)
- Topological 4-jet+ γ trigger signature implemented at Level-1 for Run 2
- Gluon-induced component of the dominant non-resonant $bbj\gamma$ suppressed
- BDT against the non-resonant background: m_{bb} fits in 3 BDT regions



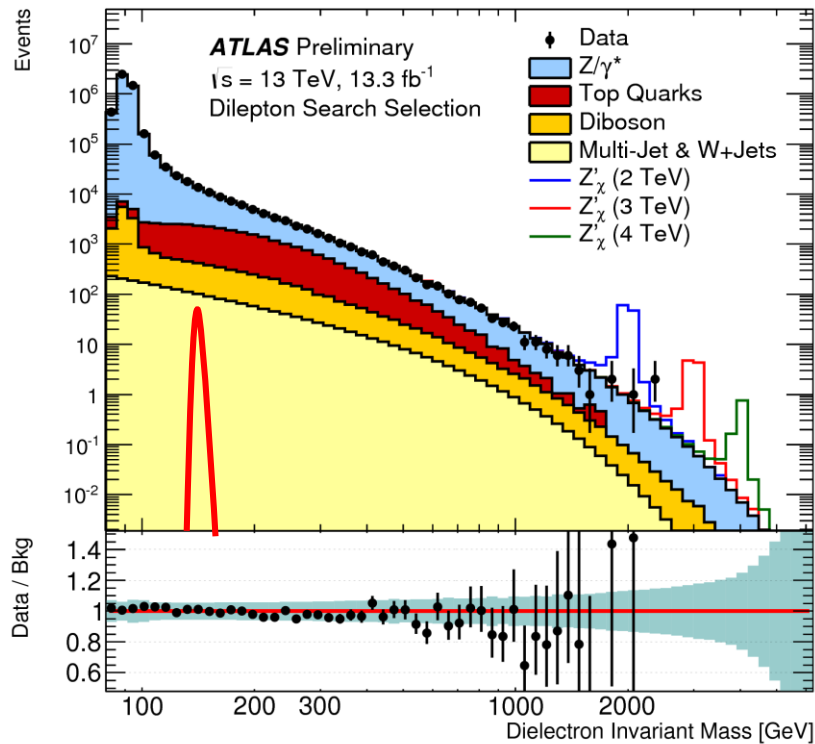
Physics Hidden in Inclusive Distributions

In various models (*common in hidden valleys/dark sectors*)

- Rare or unusual production of new neutral particle
- Bump, edge, endpoint, dip, wiggle is present
 - but swamped in inclusive background
- Point:
 - Higgs \rightarrow bb resonance is invisible in inclusive production
 - But a semi-exclusive search can reveal it
 - The selection criterion reduces background, keeps signal
- We should do this for other resonance searches as a matter of course!

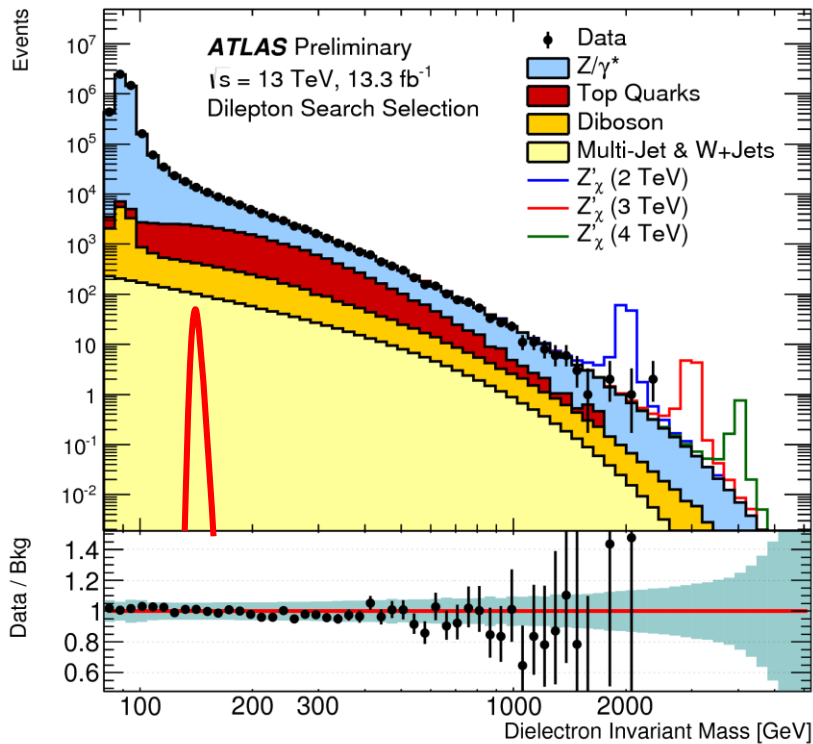
Inclusive is not Conclusive

Inclusive

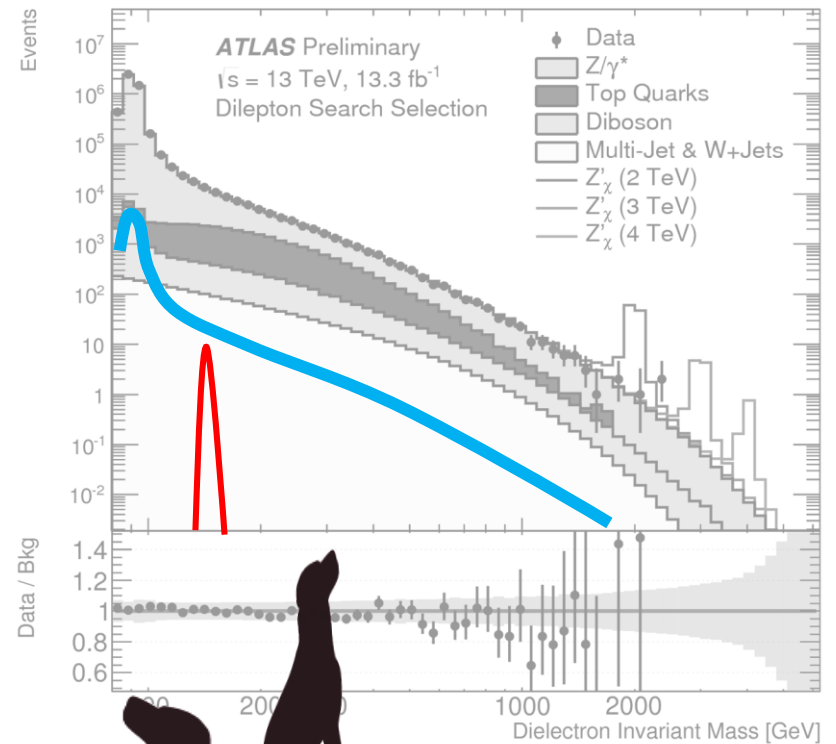


Inclusive is not Conclusive

Inclusive



Require 4 jets and high HT and... presto!



Physics Hidden in Inclusive Distributions

The method of semi-exclusive searches is general:

- Resonances: $X \rightarrow$ diphotons, di-b, di-tau, triphoton, ...
- Edge/Endpoints: $X \rightarrow$ di-object + MET
- Other structures: X interferes and creates wiggle or dip

Any such effect can be searched for

- Inclusively
- Semi-exclusively
 - And Higgs \rightarrow bb shows why we have to do it!
 - Otherwise we are giving up some easy opportunities...

Systematic approach?

Physics Hidden in Total Cross Sections

Some BSM processes might contribute to SM cross-sections

- Electroweakinos contributing to WW Lisanti & Weiner '11; Feigl, Rzehak & Zeppenfeld, '12; Curtin, Jaiswal, Meade, Tien '12,'13,'14; Rolbiecki & Sakurai '13
- Stop \rightarrow top + low mass neutralino near threshold adds to tt Czakon, Mitov, Papucci, Ruderman & Weiler '14
- Stop \rightarrow bino \rightarrow wino contributes to ttW, tth, ... Cf. Huang, Ismail, Low & Wagner '15; Angelescu, Djouadi & Moreau '15; ...

Are we properly cross-correlating all our measurements?

- Sidestep?
 - Spin effects:
 - the SM violates C and P maximally, BSM might be different

e.g. Han, Katz, Krohn & Reece '12

Ratios at different LHC collision energies

$$d\sigma^{\text{jet}} = \sum_{ijkl} dx_1 dx_2 f_i(x_1) f_j(x_2) \frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2} d\Phi_2$$

Hadronic cross-section PDFs for initial state partons Partonic cross-section two-particle phase space

Old strategy:

- Compare physics at two values of s at same value of $\hat{s}/s = x_1 x_2$
 - Pdfs almost the same
 - Partonic process changes (but often SM prediction scales with s)

Mangano & Rojo '12

- Compare physics at two values of s at same value of \hat{s}
 - Partonic process almost the same
 - Pdfs change

Ratios at different LHC collision energies

$$d\sigma^{\text{jet}} = \sum_{ijkl} dx_1 dx_2 f_i(x_1) f_j(x_2) \frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2} d\Phi_2$$

The diagram shows the equation $d\sigma^{\text{jet}} = \sum_{ijkl} dx_1 dx_2 f_i(x_1) f_j(x_2) \frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2} d\Phi_2$ with arrows pointing from labels below to parts of the equation:

- Hadronic cross-section** points to the entire equation.
- PDFs for initial state partons** points to the product $f_i(x_1) f_j(x_2)$.
- Partonic cross-section** points to the term $\frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2}$.
- two-particle phase space** points to the term $d\Phi_2$.

Old strategy:

- Compare physics at two values of s at same value of $\hat{s}/s = x_1 x_2$
 - Pdfs almost the same
 - Partonic process changes (but often SM prediction scales with s)
 - But if a threshold between $x_1 x_2 s_{\text{low}}$ and $x_1 x_2 s_{\text{high}}$ then ...

Mangano & Rojo '12

- Compare physics at two values of s at same value of \hat{s}
 - Partonic process almost the same
 - Pdfs change

Ratios at different LHC collision energies

$$d\sigma^{\text{jet}} = \sum_{ijkl} dx_1 dx_2 f_i(x_1) f_j(x_2) \frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2} d\Phi_2$$

The diagram shows the formula $d\sigma^{\text{jet}} = \sum_{ijkl} dx_1 dx_2 f_i(x_1) f_j(x_2) \frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2} d\Phi_2$ with arrows pointing from labels below to parts of the formula. A green horizontal line is drawn under the PDFs $f_i(x_1) f_j(x_2)$, and a red horizontal line is drawn under the partonic cross-section $\frac{d\hat{\sigma}_{ij \rightarrow kl}}{d\Phi_2}$.

Hadronic cross-section PDFs for initial state partons Partonic cross-section two-particle phase space

Old strategy:

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Mangano & Rojo '12

- Compare physics at two values of s at same value of \hat{s}
 - Partonic process almost the same
 - Pdfs change
 - But if a process subleading at s_{low} comes from different pdf...

14 TeV

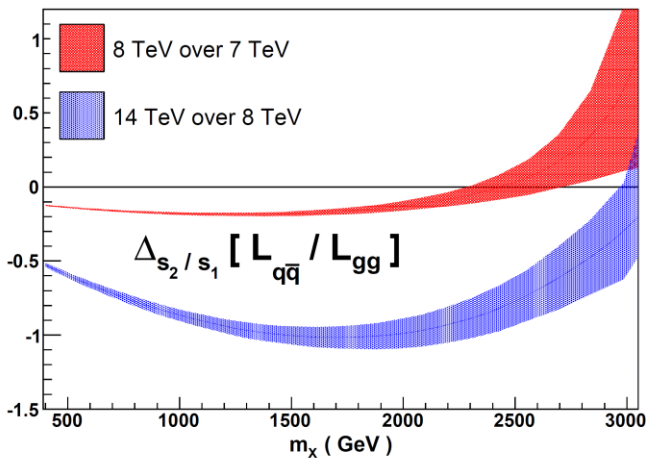
8 TeV

Ratio	R^{nnpdf}	$\delta_{\text{PDF}}(\%)$
$t\bar{t}/Z$	2.12	1.3
$t\bar{t}$	3.90	1.1
Z	1.84	0.7
W^+	1.75	0.7
W^-	1.86	0.6
W^+/W^-	0.94	0.3
W/Z	0.98	0.1
ggH	2.56	0.6
$t\bar{t}(M_{t\bar{t}} \geq 1 \text{ TeV})$	8.18	2.5
$t\bar{t}(M_{t\bar{t}} \geq 2 \text{ TeV})$	24.9	6.3
$\sigma_{\text{jet}}(p_T \geq 1 \text{ TeV})$	15.1	2.1
$\sigma_{\text{jet}}(p_T \geq 2 \text{ TeV})$	181.6	7.7

Mangano-Rojo '15

<1% - 7% variations with pdf, depending on ratio

Small variations among pdfs



Suppose there's a broad resonance at 1.5 TeV producing $t\bar{t}$ -bar

At fixed s -hat, much more gg at 13 TeV than at 8 TeV; so $q\bar{q}$ -bar process suppressed

Thus precise predictions for $t\bar{t}$ bar at 13 vs 8 TeV will fail

$$\Delta_{E_1/E_2}(A) = 1 - \frac{A(E_2)}{A(E_1)}$$

Improvement: cancel luminosity uncertainty using $t\bar{t}$ bar/Z at the two energies.

Charge Asymmetries at the LHC

- Charge asymmetries arise from valence quarks
 - They measure degree of valence quark contribution to a process
 - (Or parity violation – *subtle point...*)
- These change as the center of mass energy changes and the rates (for particular event selections) should be well-predicted
 - Especially if we can normalize them intelligently to W^* or $W+j$
- Useful in checking single top
- Would be useful in diagnosing BSM
- So energy ratios of charge asymms probably good as precision variables

Physics Hidden By Control Regions

- Away from tails, deviations typically **not statistics limited**
- Systematic uncertainties from control-to-signal extrapolation
 - How can these be reduced?
 - Attempts at precision measurements in control regions
 - Using theory at higher order to transfer to signal region
 - Additional questions (data and theory) in validation regions?
 - Large logs can appear in signal regions that are absent in controls?
 - Logs of HT/p_T^{\min}
 - EW Sudakov corrections...
- Case studies??

Ratios on the Backbone

Play off statistical uncertainties versus systematic and theory uncertainties

- Tails give us highest-E sensitivity (but do we understand the tail?)
- Backbones give us highest level of control

For a dimension-six operator, tails, backbones and heads are comparable

- Cf. discussion following Mangano talk – they become complementary.

Tails win in dimension-8

For low-mass physics, backbones may win

- Wide resonance/wiggle
- Off-shell Higgs effects at high mass

Dibosons

with Chris Frye, Marat Freytsis, Jakub Scholtz '15

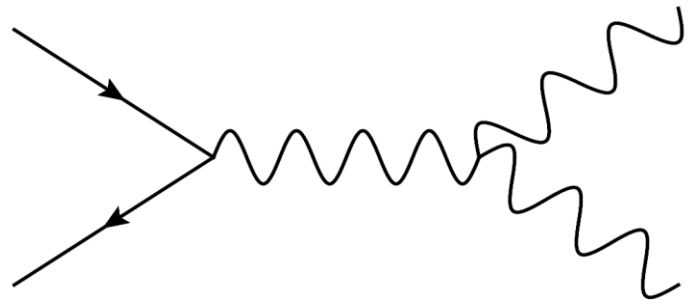
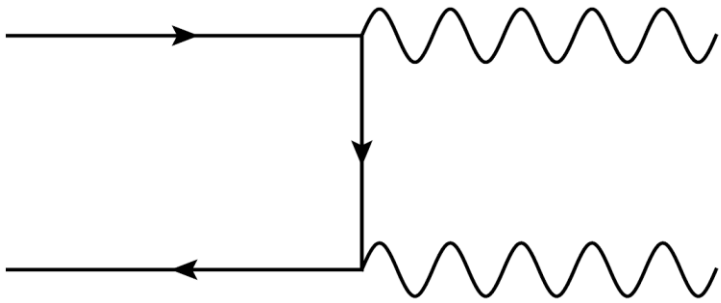
Production of any pair of photon, Z, W^\pm (except same sign)

- Discrepancies have shown up – or not...
- Tails have low statistics; Head has resummation subtleties
- What ratios/variables might help on the backbone?

- Put high-energy $SU(2) \times U(1)$ structure to use
 - Leading-order (tree-level) partonic-level into nicer form
 - Notice useful ratios, show they are still useful in pp collisions

- Proceed to realistic situation for two neutral bosons
 - Show corrections beyond leading order are small at high energy
 - NLO
 - gg-induced NNLO
 - Show remaining uncertainties are small

- *All results below using MCFM Monte Carlo* Campbell, R.K.Ellis



$$a_1 \propto \mathcal{M}(xx) \propto \mathcal{M}(wx) \propto \mathcal{M}(ww_1),$$

t,u

$$a_3 \propto \mathcal{M}(ww_3),$$

s,t,u

$$a_\phi \propto \mathcal{M}(\phi\phi),$$

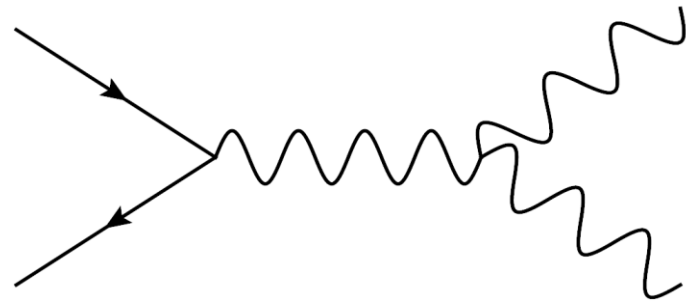
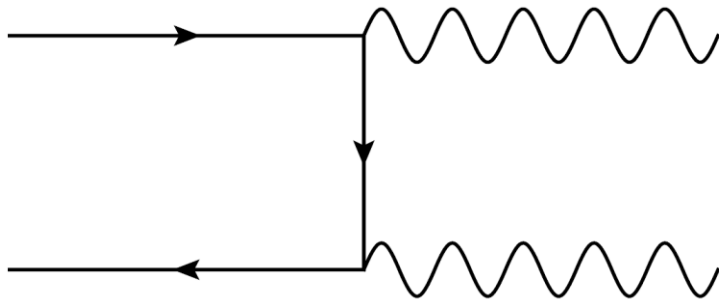
s

$SU(2) w^a (a=1,2,3), U(1) x$

- up to $(m_Z/E)^2$ terms

$$\gamma = c_W x + s_W w^3,$$

$$Z = c_W w^3 - s_W x,$$



$$a_1 \propto \mathcal{M}(xx) \propto \mathcal{M}(wx) \propto \mathcal{M}(ww_1),$$

$$a_3 \propto \mathcal{M}(ww_3),$$

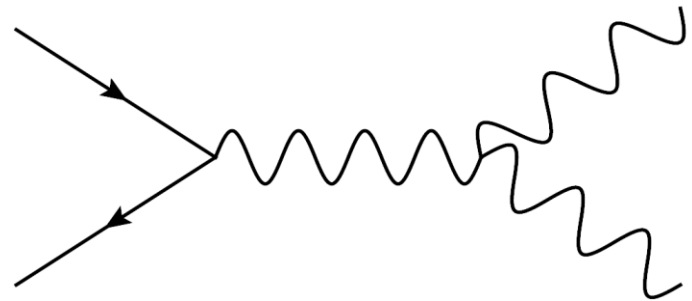
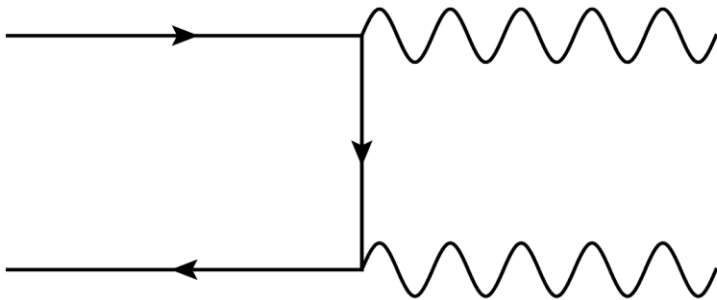
$$a_\phi \propto \mathcal{M}(\phi\phi),$$

$$|a_1|^2 = \frac{\hat{t}}{\hat{u}} + \frac{\hat{u}}{\hat{t}},$$

$$(a_1 a_3) = \left(\frac{\hat{t} - \hat{u}}{2\hat{s}} \right) + \frac{1}{4} \left(\frac{\hat{t}}{\hat{u}} - \frac{\hat{u}}{\hat{t}} \right),$$

$$|a_3|^2 = \frac{\hat{t}\hat{u}}{4\hat{s}^2} - \frac{1}{8} + \frac{1}{32} \left(\frac{\hat{t}}{\hat{u}} + \frac{\hat{u}}{\hat{t}} \right),$$

$$|a_\phi|^2 = \frac{\hat{t}\hat{u}}{4\hat{s}^2}.$$



$ZZ, Z\gamma, \gamma\gamma$ at Leading Order (@LO)

$$|a_1|^2 = \frac{\hat{t}}{\hat{u}} + \frac{\hat{u}}{\hat{t}},$$

$$\frac{d\hat{\sigma}}{d\hat{t}}(q\bar{q} \rightarrow V_1^0 V_2^0) = \frac{C_{12}^q}{\hat{s}^2} |a_1|^2,$$

$$C_{\gamma\gamma}^q = \frac{1}{2} \frac{\pi\alpha_2^2 s_W^4}{N_c} 2Q^4,$$

$$C_{Z\gamma}^q = \frac{\pi\alpha_2^2 s_W^2 c_W^2}{N_c} (L^2 Q^2 + R^2 Q^2),$$

$$C_{ZZ}^q = \frac{1}{2} \frac{\pi\alpha_2^2 c_W^4}{N_c} (L^4 + R^4).$$

Couplings to Z :

$$L = T_3 - Y_L t_W^2, \quad R = -Y_R t_W^2$$

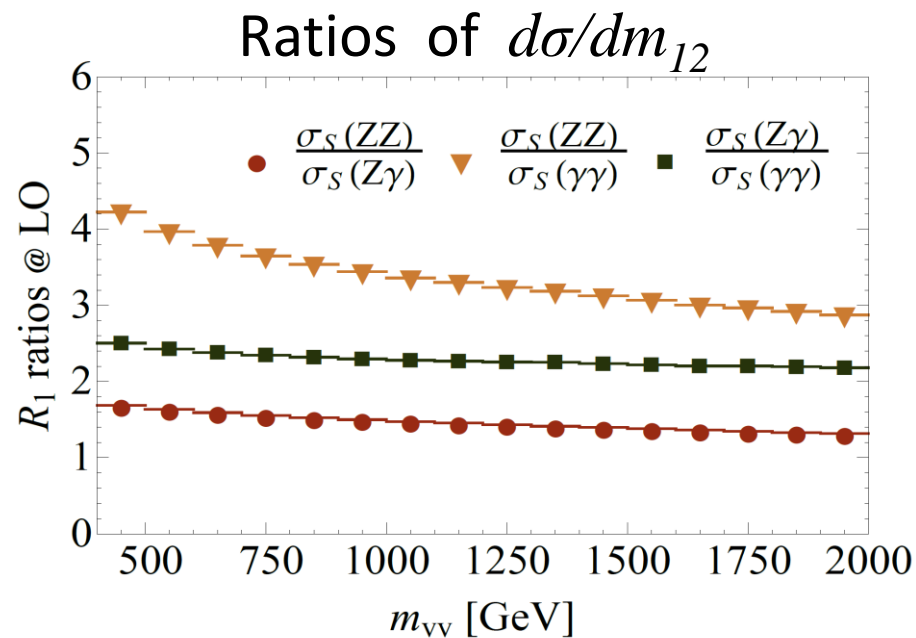
ZZ, Zγ, γγ at Leading Order (@LO)

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Couplings to Z :

$$L = T_3 - Y_L t_W^2, \quad R = -Y_R t_W^2$$

ZZ, Zγ, γγ at Leading Order (@LO)

$V_1^0 V_2^0$	$C_{12}^{uu} \cdot 10^5$	$C_{12}^{dd} \cdot 10^5$
$\gamma\gamma$	1.2	0.07
$Z\gamma$	2.2	0.7
ZZ	1.6	3.3

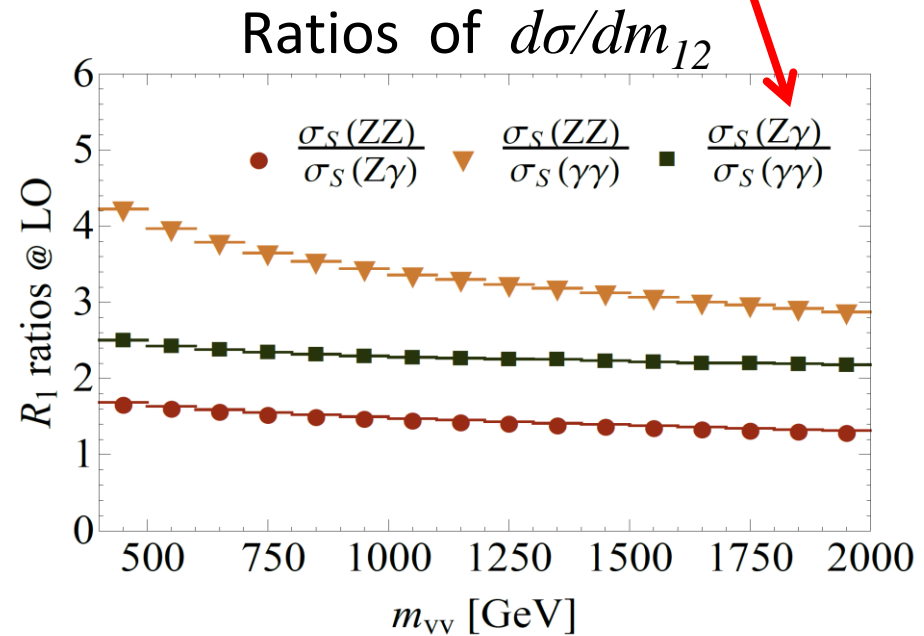
uu dominates;
PDF uncertainties
should cancel

$$\frac{d\hat{\sigma}}{d\hat{t}}(q\bar{q} \rightarrow V_1^0 V_2^0) = \frac{C_{12}^q}{\hat{s}^2} |a_1|^2,$$

$$C_{\gamma\gamma}^q = \frac{1}{2} \frac{\pi\alpha_2^2 s_W^4}{N_c} 2Q^4,$$

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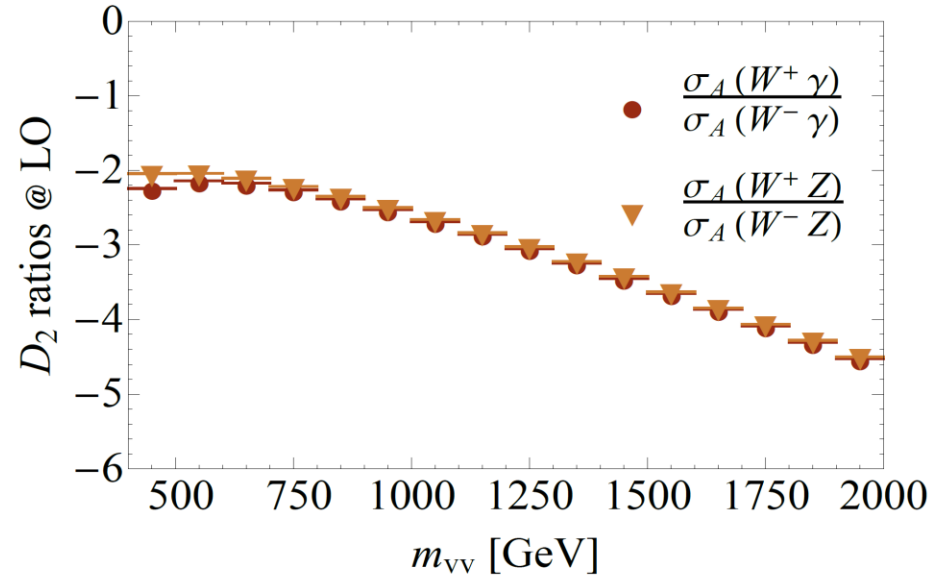
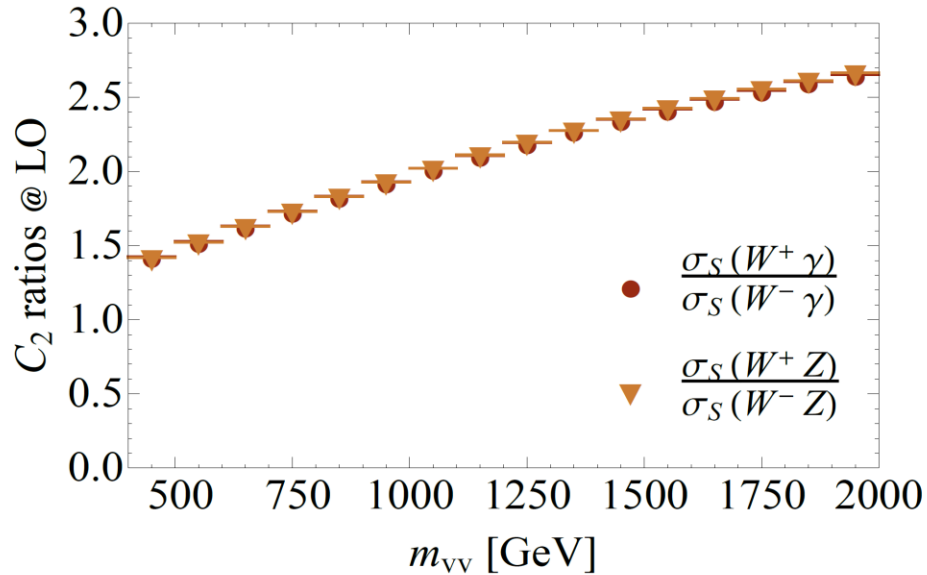


Couplings to Z :

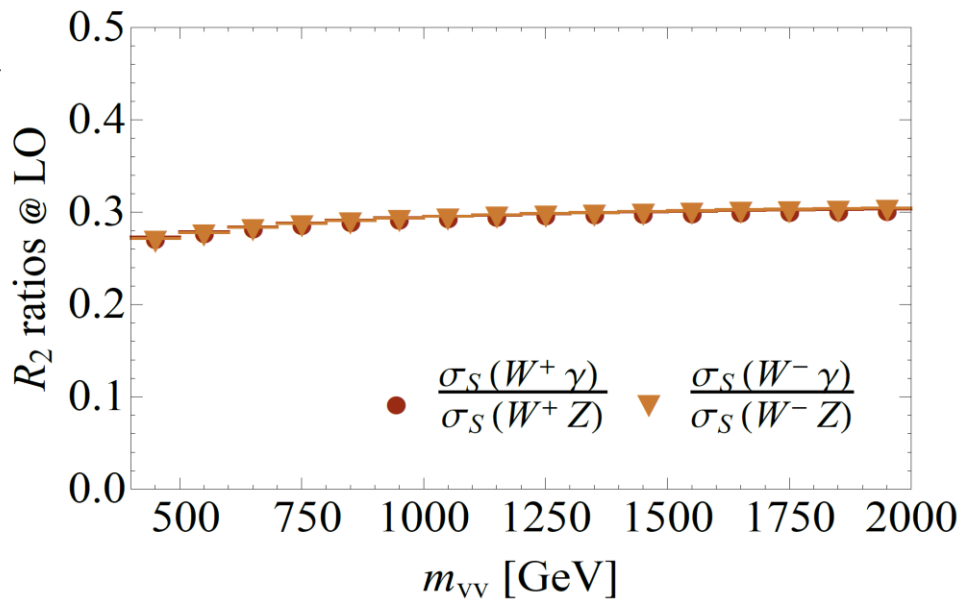
$$L = T_3 - Y_L t_W^2, \quad R = -Y_R t_W^2$$

Charge asymmetries for $W\gamma$, WZ are related

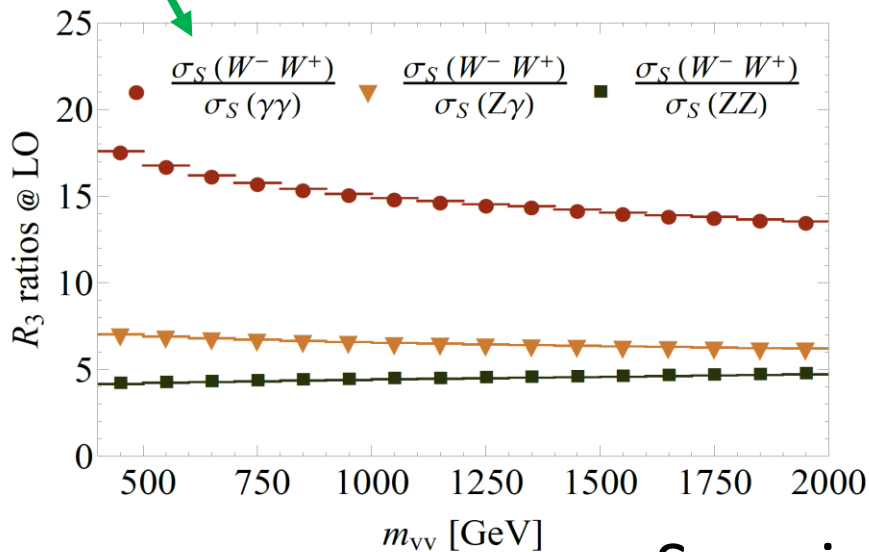
- Determined by the pdfs for both sym, antisym FB quantities



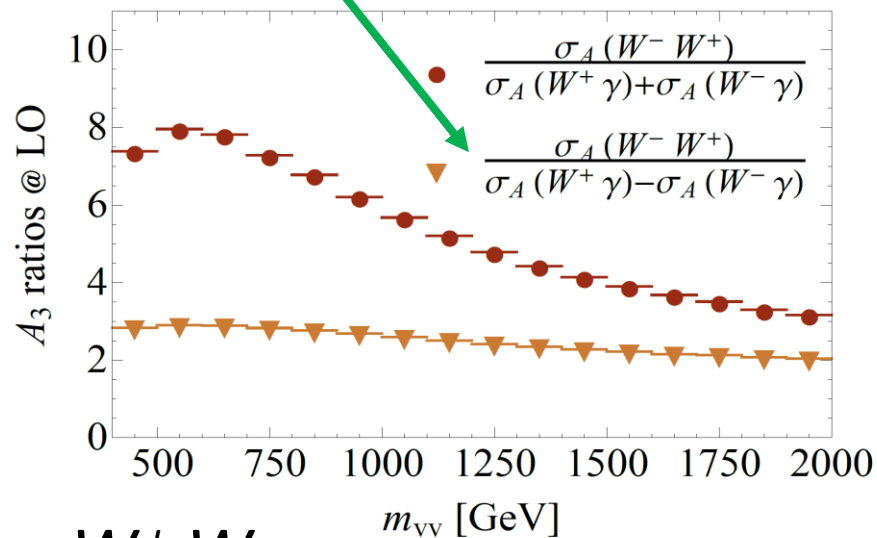
More $W^\pm V$



Best statistics



Best statistics and LO PDF behavior



Surprising $W^+ W^-$

Beyond Leading Order?

- What about higher-order corrections?
 - QCD cancellations?
 - How large are the shifts in the ratios?
 - $SU(2) \times U(1)$ relations should help -- Where do they fail?
 - What uncertainties remain?
 - EW corrections - Partial cancellations?
- Big issue: the radiation zero
 - Where important, LO $SU(2) \times U(1)$ relations may receive large corrections
- Start with $\gamma\gamma$, $Z\gamma$, ZZ
 - No radiation zero
 - Events fully reconstructed ($Z \rightarrow$ leptons ONLY here)
 - Good statistics for first two

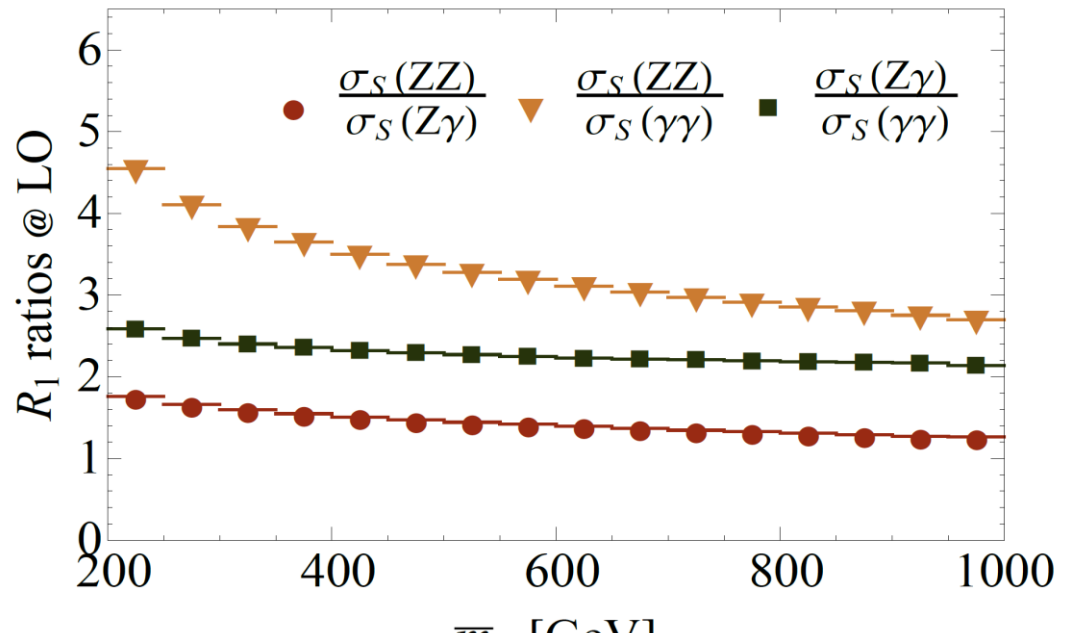
$ZZ, Z\gamma, \gamma\gamma$ at LO \rightarrow NLO

- Must choose observable carefully to avoid large NLO corrections

$$\bar{m}_T = 1/2[m_{T1} + m_{T2}] = \text{min energy at } 90^\circ \text{ scattering}$$

- Radiation cannot reduce this variable
 - so no region of NLO phase space is secretly LO.

Ratios of $d\sigma/dm_T$



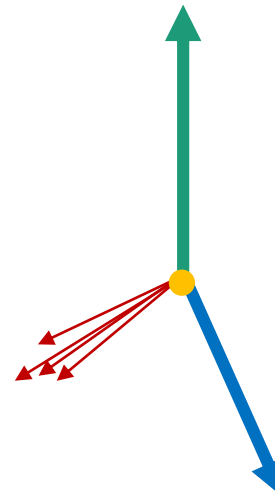
$ZZ, Z\gamma, \gamma\gamma$ at LO \rightarrow NLO

- Need to choose cuts carefully to avoid large NLO corrections
 - Assure cuts select kinematics similar to LO
 - i.e. no vector bosons softer than jets (cf. giant K factors)
 - But do not impose drastic jet veto
- We take

$$p_T^{\text{jet}} < \frac{1}{2} p_T^V|_{\text{min}} ; \frac{1}{2} p_T^V|_{\text{min}} > \frac{1}{2} p_T^V|_{\text{max}}$$

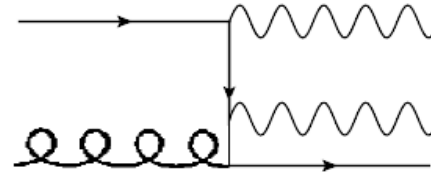
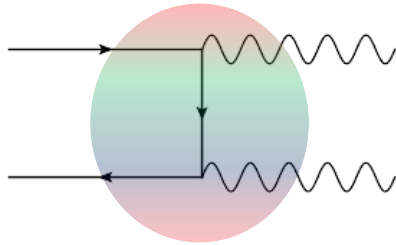


*Notice these cuts scale –
no large logs at high E*



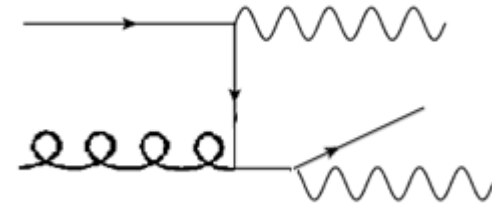
$ZZ, Z\gamma, \gamma\gamma$ at LO \rightarrow NLO

- QCD corrections treat Z, γ identically, largely cancel...



- ...except...

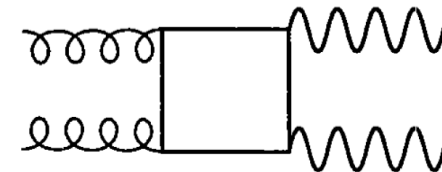
- Collinear quark-boson regime
 - Photon has log enhancement
 - Z has no enhancement



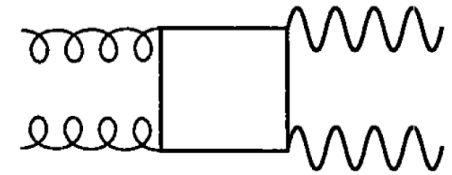
- Gluon fusion process (formally NNLO but numerically large)

- Both of these driven by gluon pdf

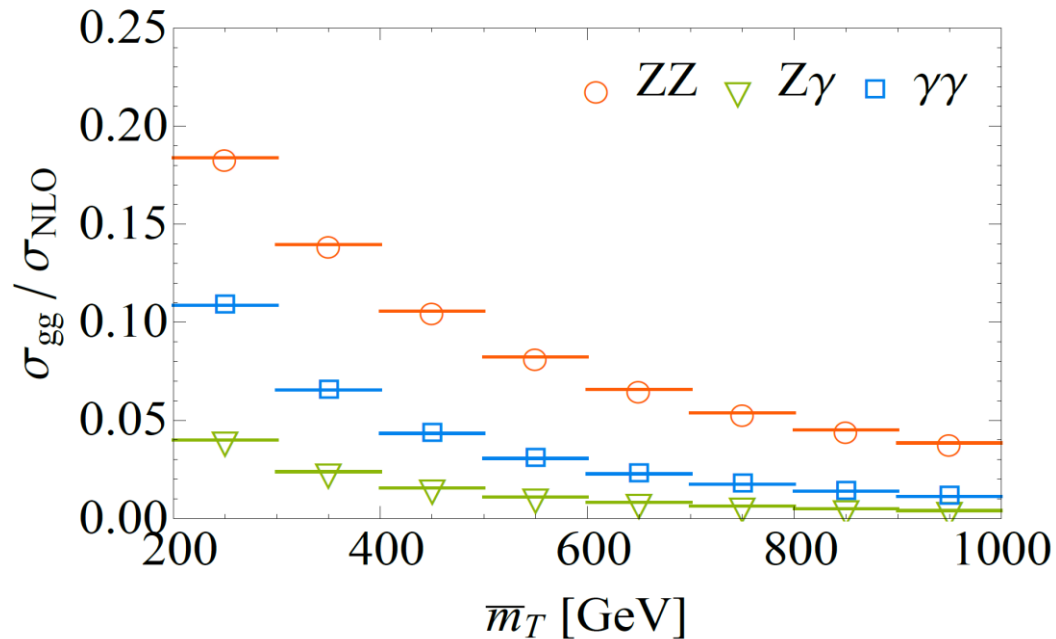
- Both decrease in importance at high energy



NNLO gg / NLO partial K factor

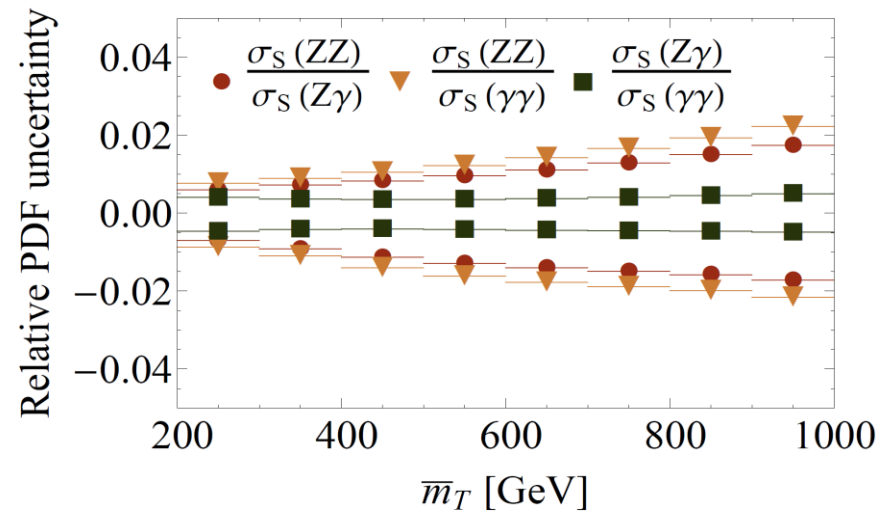
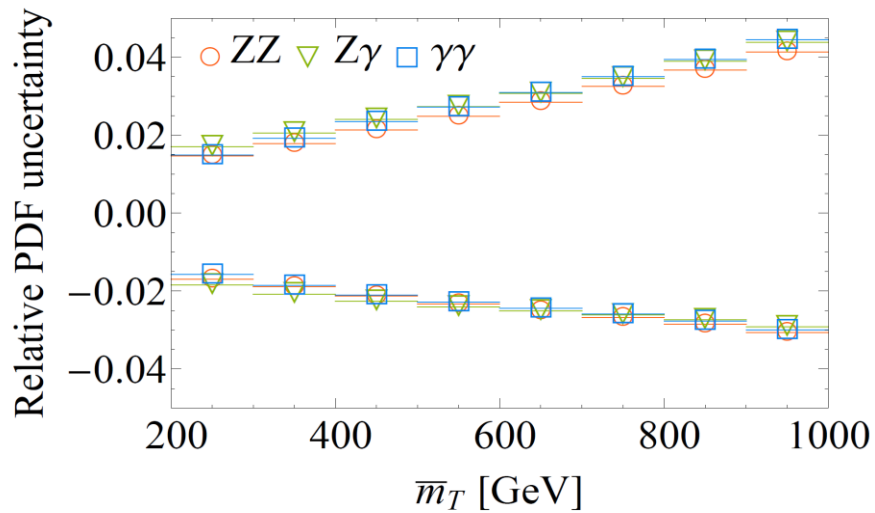


gg gives largest NNLO correction to ratios



- To set scale on gg use partial knowledge of NNNLO gg correction
 - (backup slide)

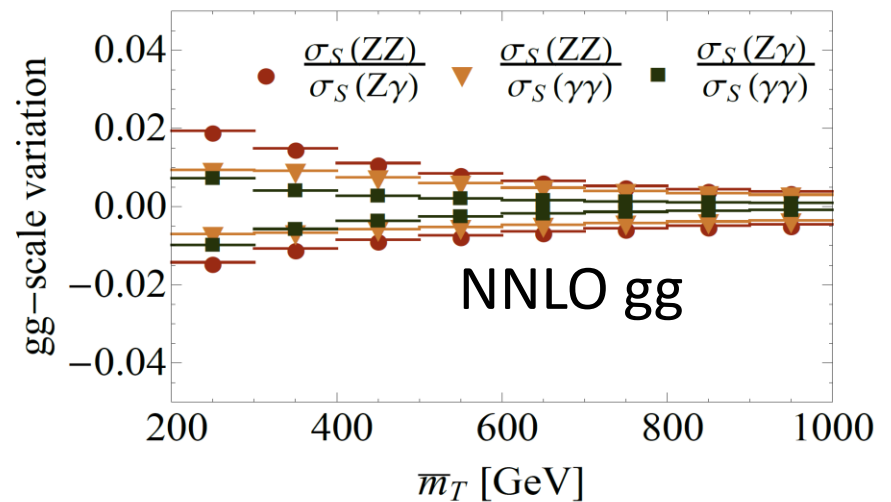
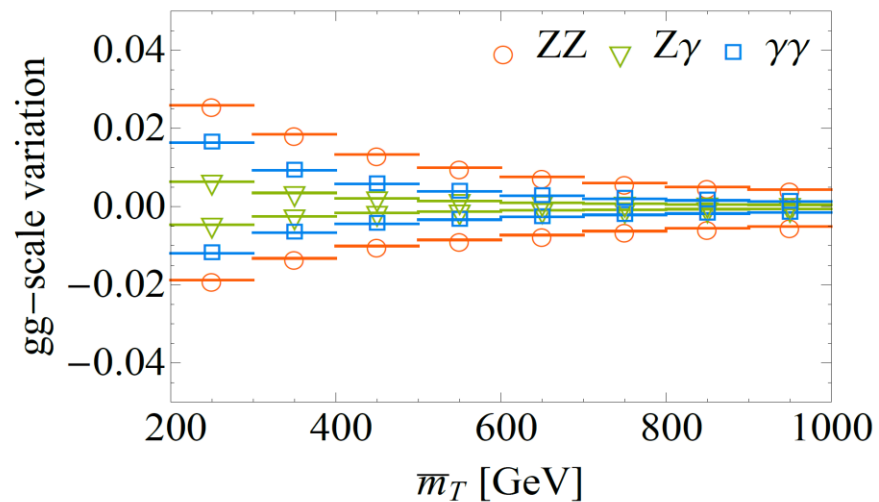
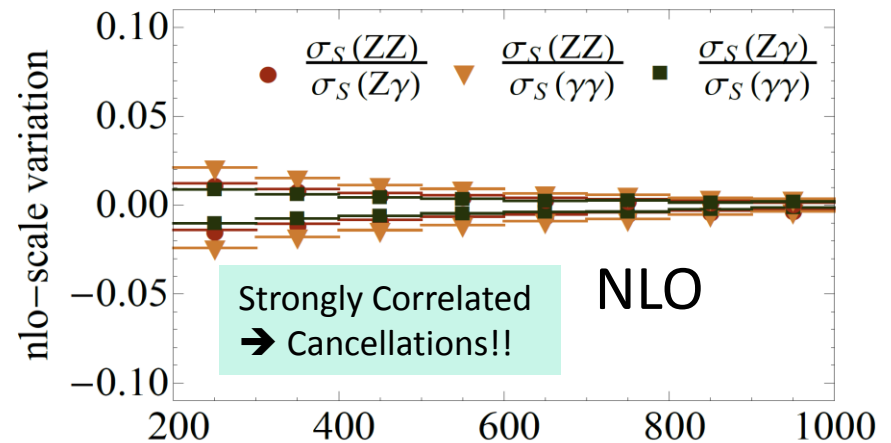
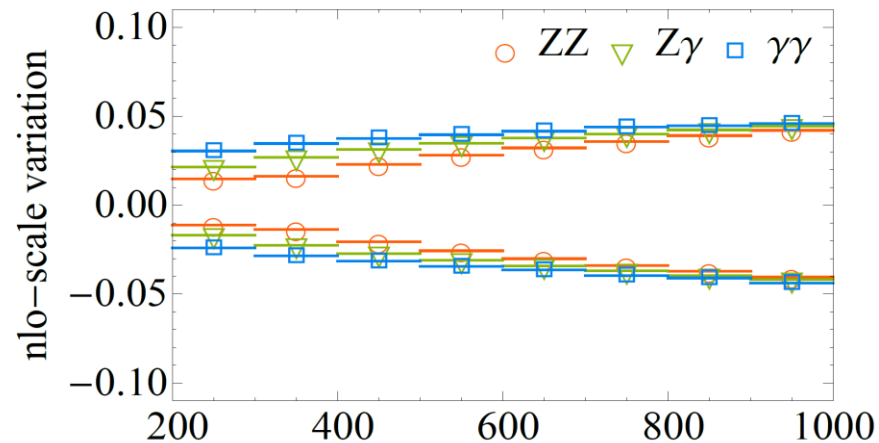
PDF Uncertainties



- Much smaller in ratios
 - 1 – 2 %
 - Especially for $Z\gamma / \gamma\gamma$

Scale [next-order] uncertainties

- Estimates NNLO corrections to what is already present at NLO



- Does not account for new channels (e.g. $q q \rightarrow q q V V \sim 2-3\%$)

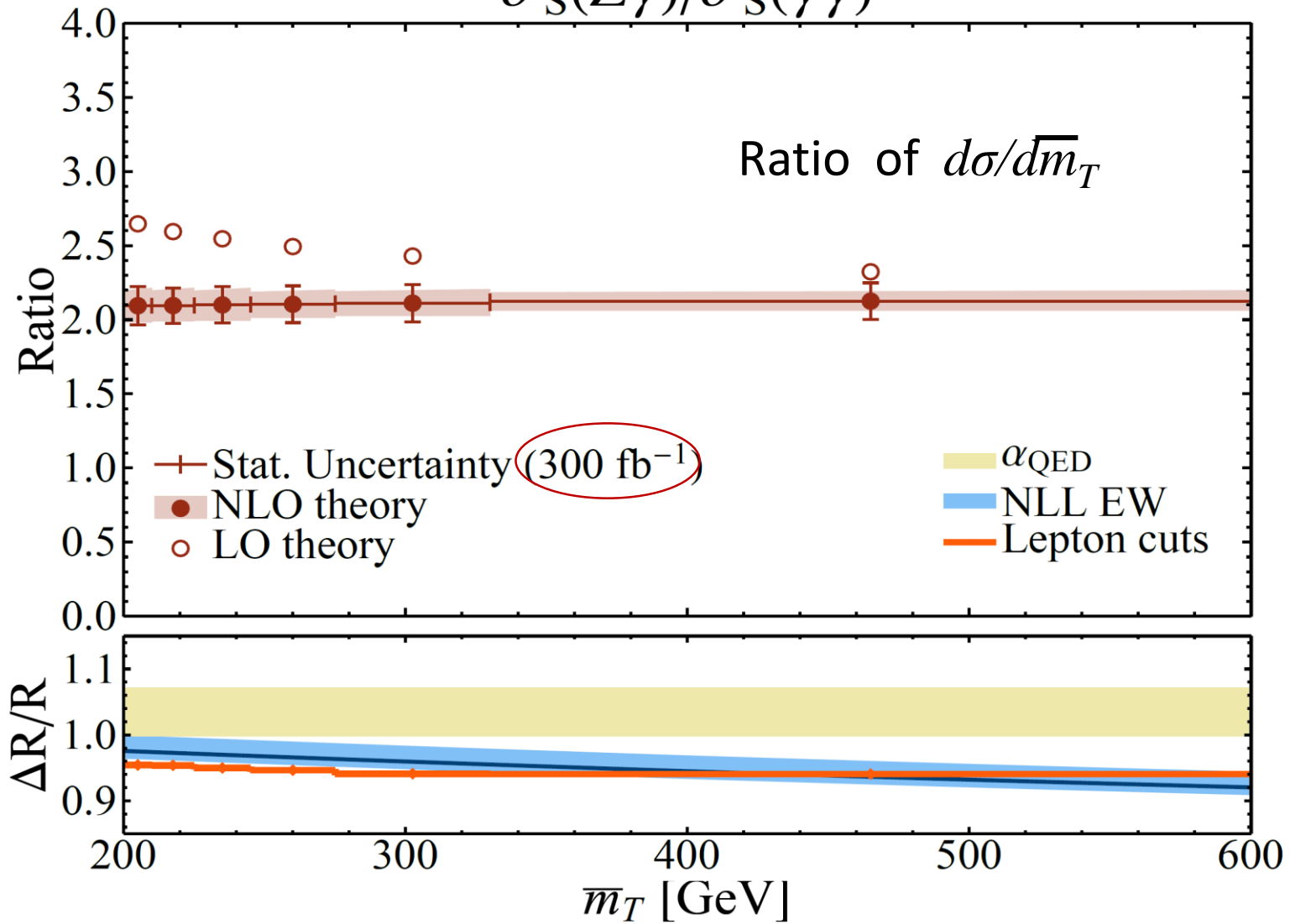
Experimental effects

- Some experimental issues cancel
 - Luminosity
 - Jet energy scale
- Some don't:
 - $Z \rightarrow$ leptons – leptons have their own cuts, acceptance
 - Or \rightarrow neutrinos -- other issues
 - Can be a substantial effect at low p_T
 - But can model, measure with low absolute uncertainty
 - Z – finite width [experimental definition of “Z”]
 - Not large effect
 - Can model

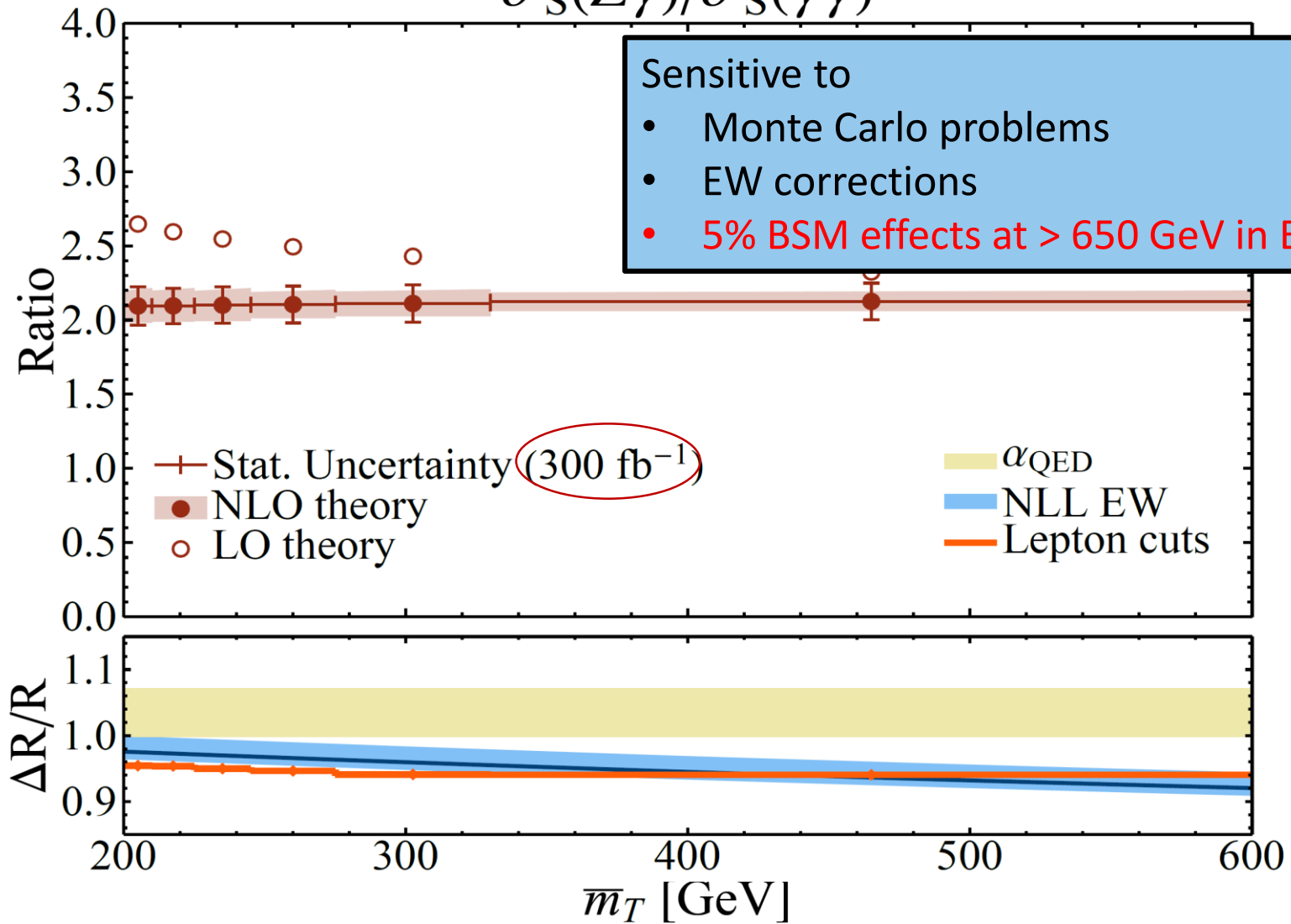
Uncertainty budget

Effect	R_{1a} ($Z\gamma/\gamma\gamma$)	R_{1b} ($ZZ/\gamma\gamma$)	R_{1c} ($ZZ/Z\gamma$)	Comments
$qq \rightarrow VVqq$	2–3%	3–3.5%	1.5–2.5%	extrapolating $p_{T,\min}^j \rightarrow 0$ (Sec. 4.2)
μ_R, μ_F (gg)	0.5–1%	1%	1–2%	uses NLO $gg \rightarrow \gamma\gamma$ (Sec. 4.5)
μ_R, μ_F (NLO)	0.5–1%	1.5–2.5%	1–1.5%	varied independently (Sec. 4.5)
PDF	0.5%	1–1.5%	0.5–1%	MSTW 2008 using MCFM (Sec. 4.5)
NLO EW	+2% –1%	+3% –1%	+2% –1%	EFT scale uncertainty (Sec. 4.4.1)

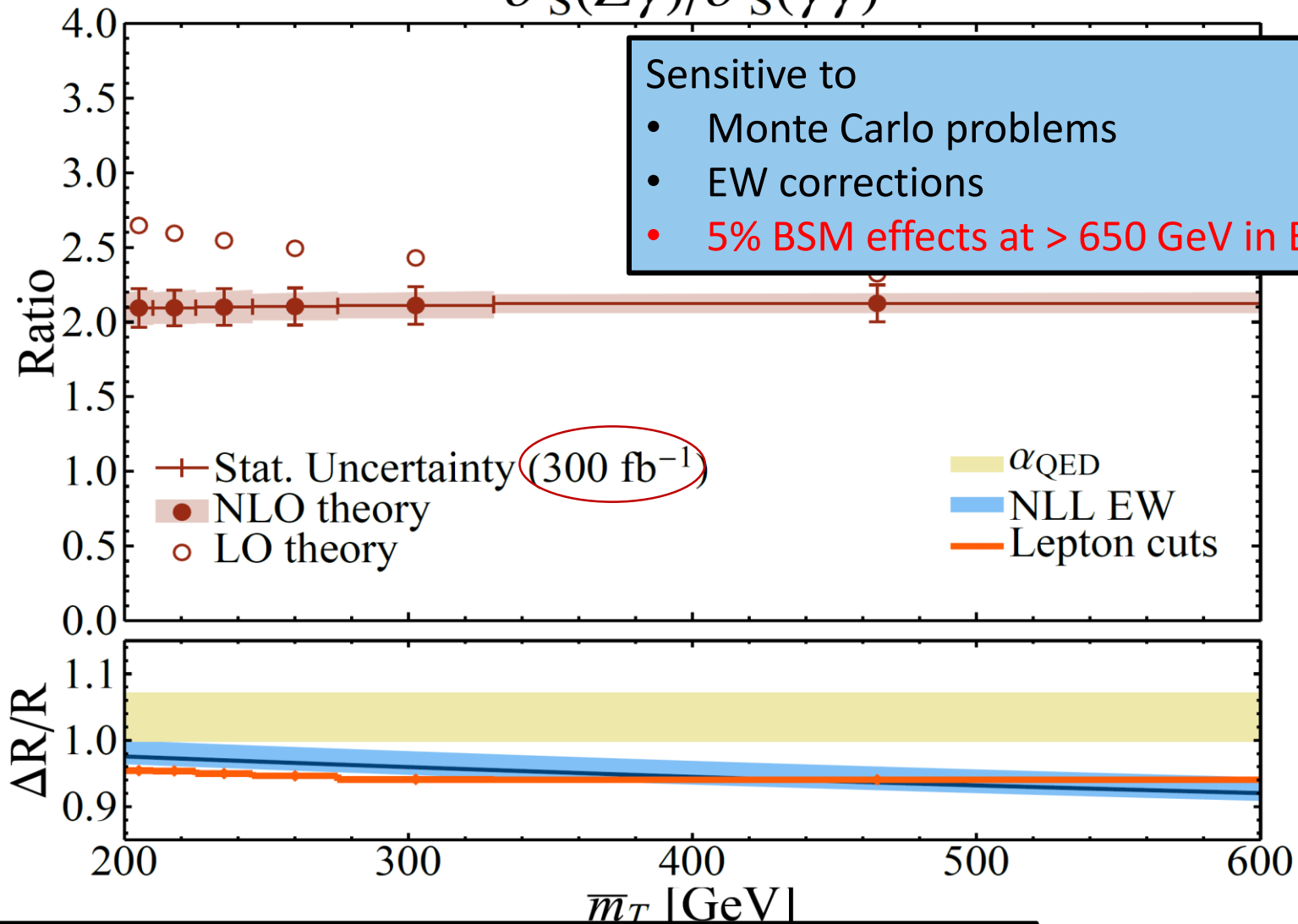
$$\sigma_S(Z\gamma)/\sigma_S(\gamma\gamma)$$



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$$\sigma_S(Z\gamma)/\sigma_S(\gamma\gamma)$$



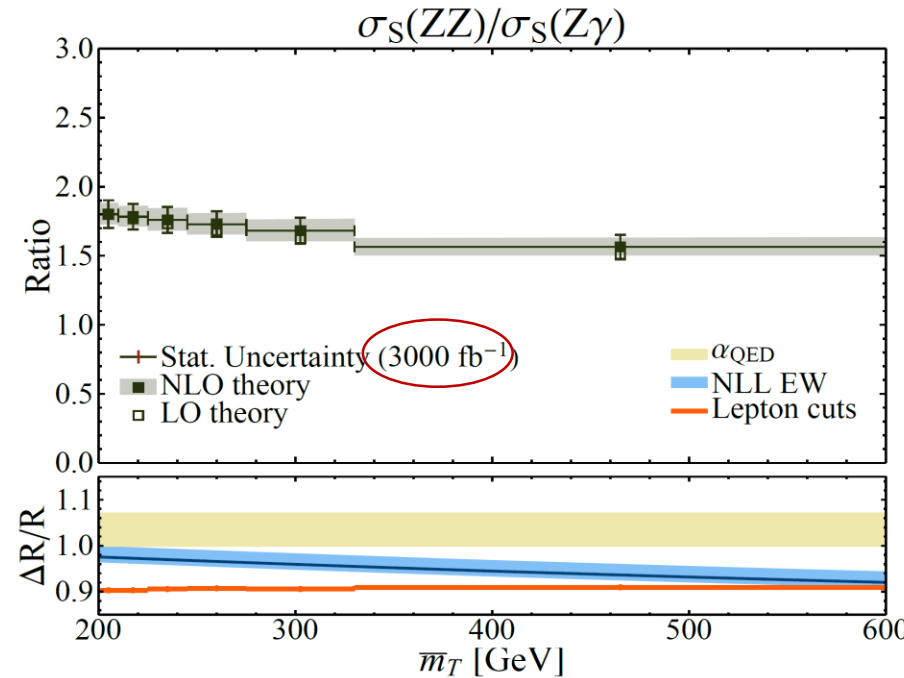
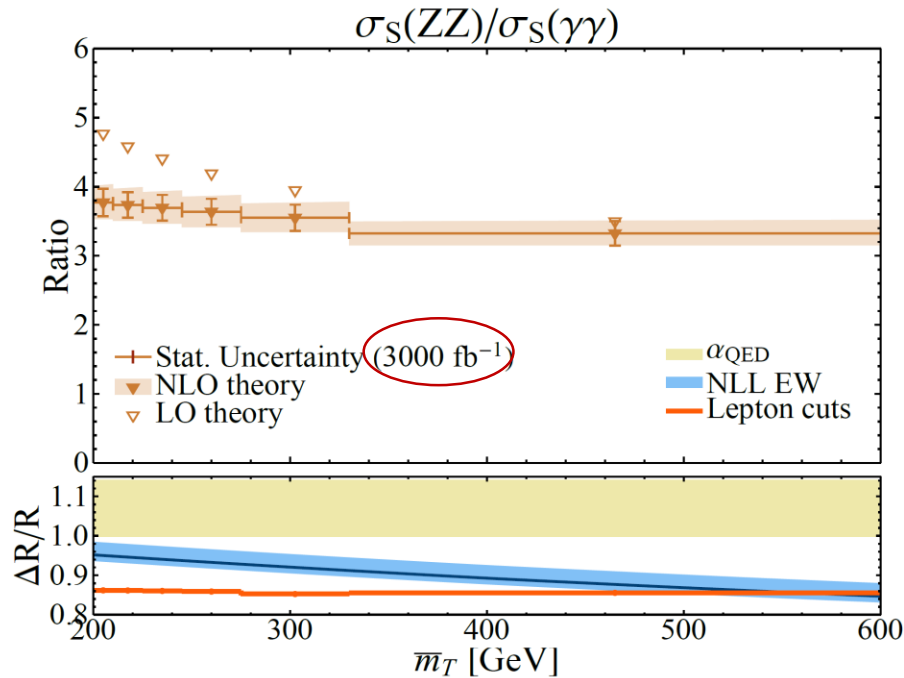
Sensitive to

- Monte Carlo problems
- EW corrections
- 5% BSM effects at $> 650 \text{ GeV}$ in EW sector

Possible Improvements:

- Use $Z \rightarrow$ neutrinos?
- Use $Z \rightarrow$ jets??
- At 3000 fb^{-1} , tens of bins, last bin probes $> 1.2 \text{ TeV}$ at 5%

The other ratios, at 3000 fb⁻¹



Probably want to include $Z \rightarrow$ neutrinos at price of higher theoretical uncertainty.

Compare backbone to tails

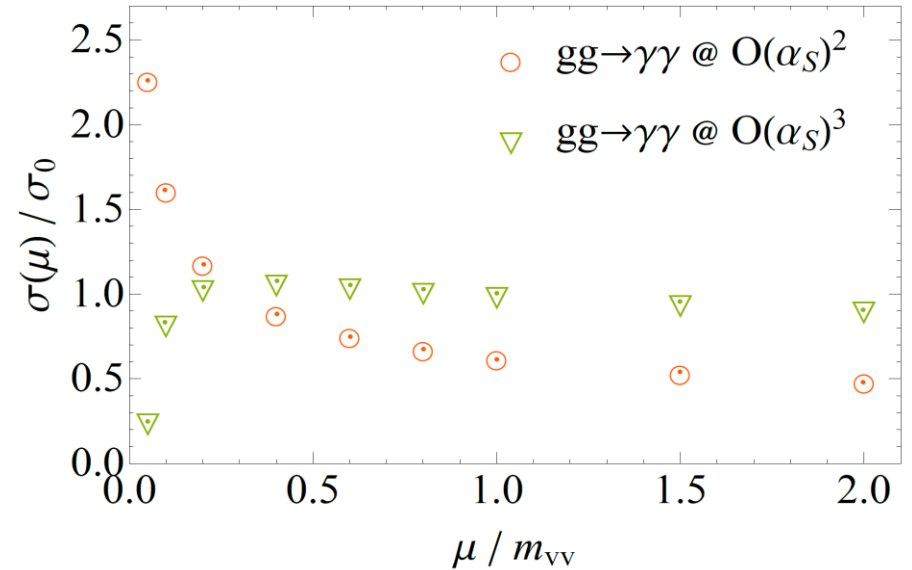
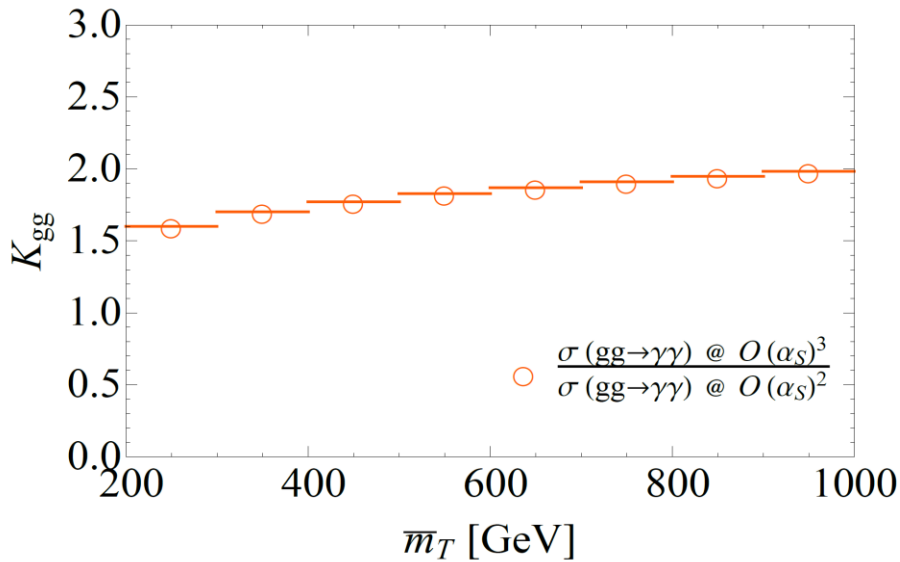
- If you're looking for dimension-6 BSM effect, tail and backbone similar sensitivity.
 - Pdf luminosity roughly falls with a power of s
 - Dim-6 effect roughly grows with power of s
- Unfortunately (*to be confirmed*) anomalous gauge couplings are
 - Dimension six (longitudinal bosons) but suppressed
 - Dimension eight
- But even if this particular effect only on tails,
 - Precise diboson MC can be tested using ratios
 - Electroweak effects (big on tails) can be tested without much QCD pollution
- And other BSM effects can appear on the backbone

Conclusions

- It's not whether to look in the bulk, but how best to do it
 - High statistics can be in your favor, if you can reduce systematics
 - High statistics can be against you; cut wisely (go onto an orthogonal tail)
- Look for buried treasure (*inclusive is not conclusive*)
- Compare 8 TeV and 13 TeV cleverly (*cross-sections; charge asyms?*)
- Control the control regions (*and the transfer to signal regions*)
- New EW/hidden/rare physics requires high precision in bulk
 - Exercise: get high precision in diboson ratios at $p_T \gg M_W$
 - Ratios: small QCD corrections & uncertainties at high energy
 - Certainly good for SM studies (MC and EW)
 - Need still to learn more about how sensitivity to BSM is improved
 - Searches on high-E tails vs. precision at moderate E

Scale setting for gg loops

- For $gg \rightarrow \gamma\gamma$



- For the other processes

$$K_{gg} \equiv \frac{d\sigma_{(3)}(gg \rightarrow \gamma\gamma)}{d\sigma_{(2)}(gg \rightarrow \gamma\gamma)} \approx \frac{d\sigma_{(3)}(gg \rightarrow Z\gamma)}{d\sigma_{(2)}(gg \rightarrow Z\gamma)} \approx \frac{d\sigma_{(3)}(gg \rightarrow ZZ)}{d\sigma_{(2)}(gg \rightarrow ZZ)}$$