

Noncommutative Spherically Symmetric Spacetimes

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Quantum Spacetime Hypothesis

Quantum Gravity



Quantum Geometry



Classical Geometry

Outline

- ▶ Differential graded algebra \rightarrow quantum metric
- ▶ Geometry at the semi-classical level
- ▶ Application to spherically symmetric spacetimes

Quantum Differential Algebra

Quantum Algebra

- ▶ (\mathcal{A}, \cdot) unital, associative etc.
- ▶ Not necessarily commutative: $ab \neq ba$

Quantum Differential Algebra

- ▶ Ω^1 is a bimodule of \mathcal{A} so $a((db)c) = (a(db))c$
- ▶ $d : \mathcal{A} \rightarrow \Omega^1$ so $d(ab) = (da)b + a(db)$
- ▶ $\{adb\}$ span Ω^1
- ▶ $\ker d = k$

Extends to Differential Graded Algebra $\Omega = \bigoplus_i \Omega^i$

- ▶ $d^2 = 0$
- ▶ Product $\wedge : \Omega^n(\mathcal{A}) \otimes \Omega^m(\mathcal{A}) \rightarrow \Omega^{n+m}$

Quantum metric

$g \in \Omega^1 \otimes \Omega^1$ with quantum symmetry $\wedge(g) = 0$

Has an inverse:

$$(\cdot, \cdot) : \Omega^1 \otimes \Omega^1 \rightarrow \mathcal{A}$$

Bimodule requirement forces g to be central!

Also require a compatible bimodule connection

$$\nabla : \Omega^1 \rightarrow \Omega^1 \otimes \Omega^1, \quad \sigma : \Omega^1 \otimes \Omega^1 \rightarrow \Omega^1 \otimes \Omega^1$$

- ▶ Metric compatible $\nabla(g) = 0$
- ▶ Torsion free $\wedge \nabla = d$

Quantum Levi-Civita connection!

Example

S. Majid, W. Tao; Phys. Rev. D 91 (2015)

Take Majid-Ruegg model: $[x^i, t] = i\lambda\rho x^i$ & $[t, dt] = i\lambda\rho\alpha dt$

$$g = a_{ij}dx^i \otimes dx^j + b_i(dx^i \otimes dt + dt \otimes dx^i) + cdt \otimes dt$$

$$[g, f] = 0 \quad \forall f \in \mathcal{A}, \quad \wedge(g) = 0$$



$$g = \delta^{-1}d\Omega^2 + ar^{-2}dr \otimes dr + br^{\alpha-1}(dr \otimes dt + dt \otimes dr) + cr^{2\alpha}dt \otimes dt$$

$$\bar{\delta} = \frac{c\alpha^2}{b^2 - ac} > 0 \quad a, b, c, \delta \in \mathbf{R}$$

This is the Bertotti-Robinson metric.

Algebra forces the metric!

Semi-Classical Quantum Gravity

E. Beggs, S.Majid; Class.Quant.Grav. 31 (2014)

Interested in quantization to $\mathcal{O}(\lambda) \rightarrow$ Semiquantization

- ▶ Controlled by Poisson bracket $(\{, \}, C^\infty(\mathcal{M}))$
- ▶ Have $[a, b] = \lambda\{a, b\} + \mathcal{O}(\lambda^2)$
- ▶ $\{, \} \leftrightarrow \omega^{ij}$

For differential structure have

- ▶ For $\eta \in \Omega^1$ have $[a, \eta] = \lambda\nabla_{\hat{a}}\eta + \mathcal{O}(\lambda^2)$
- ▶ $\hat{a} = \{a, \}$ and $\nabla_{\hat{a}}$ is a Poisson (pre)connection defined along a Hamiltonian vector field
- ▶ $d\{a, b\} = \nabla_{\hat{a}}b - \nabla_{\hat{b}}a$

Look at associativity \rightarrow Jacobi identity

$$[a, [db, c]] + [c, [a, db]] + [db, [c, a]] \sim \nabla_{\hat{a}}\nabla_{\hat{b}}db - \nabla_{\hat{a}}\nabla_{\hat{c}}db - \nabla_{\widehat{\{a, c\}}}db$$

Nonflat connection \rightarrow nonassociative calculus at $\mathcal{O}(\lambda^2)$ (but associative functions)

Semiquantization

E. Beggs, S.Majid; J.Geom.Phys. 114 (2017)

- ▶ Inverse problem: Can we construct a quantization that will produce a particular metric g ?
- ▶ Answer: Yes, but only to first order in λ : Semiquantization
- ▶ We map geometric data to algebraic data

Q:(vector bundles and connections) \longrightarrow (bimodules over a deformed algebra)

From the standpoint of physics can view λ as the effective scale of the theory (e.g. Planck scale), so it is reasonable to work only to first order.

Semiquantization

What do we need? Have

- ▶ Metric g
- ▶ Poisson bracket $\{ , \} \Leftrightarrow \sum_{\text{cyclic}(\alpha,\beta,\gamma)} \omega^{\alpha\mu} \omega^{\beta\gamma}{}_{,\mu} = 0$
- ▶ Poisson connection ∇
- ▶ Levi-Civita connection: $\widehat{\nabla} = \nabla + S$
 $S^\alpha{}_{\mu\nu} = \frac{1}{2} g^{\alpha\beta} (T_{\beta\mu\nu} + T_{\mu\beta\nu} + T_{\nu\beta\mu})$

And want

- ▶ Metric: $g \longrightarrow g_1$
- ▶ Connection: $\nabla \longrightarrow \nabla_1$
- ▶ Tensor product: $E \otimes F \longrightarrow Q(E) \otimes_1 Q(F)$
- ▶ Wedge product: $\eta \wedge \xi \longrightarrow Q(\eta) \wedge_1 Q(\xi)$

Semiquantization

We take $\mathcal{A} = C^\infty(\mathcal{M})$ and a deformed product \bullet . Now have

$$a \bullet b = ab + \frac{\lambda}{2} \omega^{\mu\nu} \partial_\mu a \partial_\nu b$$

Extend this to Ω

$$a \bullet dx = adx + \frac{\lambda}{2} \omega^{\mu\nu} a_{,\mu} \nabla_\nu dx$$

$$dx \bullet a = (dx)a + \frac{\lambda}{2} \omega^{\mu\nu} a_{,\nu} \nabla_\mu dx$$

Defines a Bimodule structure $E \otimes_1 A_1 \rightarrow E$

This is the action of $Q!$

Semiquantization

Quantum metric

$$g_1 = g_{\mu\nu} dx^\mu \otimes_1 dx^\nu + \frac{\lambda}{2} \omega^{\alpha\beta} \Gamma_{\mu\alpha\kappa} \Gamma^\kappa_{\beta\nu} dx^\mu \otimes_1 dx^\nu + \frac{\lambda}{2} \mathcal{R}_{\mu\nu} dx^\mu \otimes_1 dx^\nu$$

Quantum Connection

$$\nabla_1 dx^\ell = - \left[\widehat{\Gamma}^\ell_{\mu\nu} + \frac{\lambda}{2} \omega^{\alpha\beta} \left(\widehat{\Gamma}^\ell_{\mu\kappa,\alpha} \Gamma^\kappa_{\beta\nu} - \widehat{\Gamma}^\ell_{\kappa\tau} \Gamma^\kappa_{\alpha\mu} \Gamma^\tau_{\beta\nu} + \widehat{\Gamma}^\ell_{\alpha\kappa} (R^\kappa_{\nu\mu\beta} + \nabla_\beta S^\kappa_{\mu\nu}) \right) \right] dx^\mu \otimes_1 dx^\nu$$

Quantum Laplace operator

$$\square_1 f := (,)_1 \nabla_1 df = \square f + \frac{\lambda}{2} \omega^{\alpha\beta} (\text{Ric}^\gamma_\alpha - S^\gamma_{;\alpha}) (\widehat{\nabla}_\beta df)_\gamma$$

Semiquantization

Generalized Ricci two form

$$\mathcal{R} = g_{\alpha\beta} \omega^{\alpha\gamma} (\nabla_{\gamma} S^{\beta}_{\mu\nu} - R^{\beta}_{\mu\nu\gamma}) dx^{\nu} \wedge dx^{\mu}$$

Then for $\widehat{\nabla} = \nabla + S$ have the conditions

- ▶ Poisson Compatibility $\widehat{\nabla}_{\gamma} \omega^{\alpha\beta} + S^{\alpha}_{\delta\gamma} \omega^{\delta\beta} + S^{\beta}_{\delta\gamma} \omega^{\alpha\delta} = 0$
- ▶ Metric compatibility $\nabla g = 0$
- ▶ Quantum Levi-Civita condition

$$\widehat{\nabla} \mathcal{R} + \omega^{\alpha\beta} g_{\rho\sigma} S^{\sigma}_{\beta\nu} (R^{\rho}_{\mu\gamma\alpha} + \nabla_{\alpha} S^{\rho}_{\gamma\mu}) dx^{\gamma} \otimes dx^{\mu} \wedge dx^{\nu} = 0$$

Application

Take generic spherically symmetric metric

$$g = a^2(r, t)dt \otimes dt + b^2(r, t)dr \otimes dr + c^2(r, t)(d\theta \otimes d\theta + \sin^2(\theta)d\phi \otimes d\phi)$$

Poisson tensor? Ansatz: spherical symmetry

$$\omega^{23} = \frac{f(r, t)}{\sin(\theta)} \quad \omega^{01} = g(r, t)$$

So have Levi-Civita connection $\widehat{\nabla}$.

- ▶ Need ω Poisson: $g\partial_t f = g\partial_r f = 0$
- ▶ Find Poisson connection $\nabla = \widehat{\nabla} - S$
- ▶ Quantum L.C. condition

Uniqueness Theorem

Find $f = k$ and $g = 0$ and

$$S_{022} = c\partial_t c, \quad S_{122} = c\partial_r c, \quad S_{033} = c\partial_t c \sin^2(\theta), \quad S_{133} = c\partial_r c \sin^2(\theta)$$

$$S_{120} = S_{123} = S_{223} = S_{320} = S_{130} = S_{132} = S_{230} = S_{233} = 0$$

Uniquely!

Unique quantization up to $\mathcal{O}(\lambda^2)$ with algebra

$$[z^i, z^j] = \lambda \epsilon^{ij}{}_k z^k, \quad [z^i, dz^j] = \lambda z^i e^j{}_{mn} z^m dz^n$$

$$\sum (z^i)^2 = 1, \quad \sum z^i dz^i = 0$$

Nonassociative fuzzy sphere.

Central (classical) t, r, dt, dr

Quantum Geometry

Recall, we have $g_1, (\cdot, \cdot)_1, \nabla_1, \wedge_1$. Can construct

- ▶ Quantum Riemann tensor

$$\text{Riem}_1(dx^\mu) = (d \otimes_1 \text{id} - (\wedge_1 \otimes_1 \text{id})(\text{id} \otimes_1 \nabla_1))\nabla_1(dx^\mu)$$

- ▶ Quantum Ricci tensor

$$\text{Ricci}_1 = ((\cdot, \cdot)_1 \otimes_1 \text{id} \otimes_1 \otimes \text{id})(\text{id} \otimes_1 i_1 \otimes_1 \text{id})(\text{id} \otimes_1 \text{Riem}_1)(g_1)$$

Where $\wedge \circ i = \text{id}$ and $i : \Omega^2 \rightarrow \Omega \otimes \Omega$

- ▶ Quantum Ricci scalar

$$S_1 = (\cdot, \cdot)_1 \text{Ricci}_1$$

- ▶ Quantum Laplace operator

$$\square_1 f := (\cdot, \cdot)_1 \nabla_1 df = \square f + \frac{\lambda}{2} \omega^{\alpha\beta} (\text{Ric}^\gamma{}_\alpha - S^\gamma{}_{;\alpha})(\widehat{\nabla}_\beta df)_\gamma$$

Metric and Connection

$$g_1 = g_{\mu\nu} dx^\mu \otimes_1 dx^\nu + \frac{\lambda c^2}{2(z^3)^2} \epsilon_{3ij} (z^3 dz^i \otimes_1 dz^j - z^i dz^3 \otimes_1 dz^j)$$

$$\nabla_1(dt) = -\hat{\Gamma}^0_{\mu\nu} dx^\mu \otimes_1 dx^\nu - \frac{\lambda}{2(z^3)^2} \frac{c \partial_t c}{a^2} \epsilon_{3ij} (z^3 dz^i \otimes_1 dz^j - z^i dz^3 \otimes_1 dz^j)$$

$$\nabla_1(dr) = -\hat{\Gamma}^1_{\mu\nu} dx^\mu \otimes_1 dx^\nu + \frac{\lambda}{2(z^3)^2} \frac{c \partial_r c}{b^2} \epsilon_{3ij} (z^3 dz^i \otimes_1 dz^j - z^i dz^3 \otimes_1 dz^j)$$

$$\nabla_1(dz^a) = -\hat{\Gamma}^i_{\mu\nu} dx^\mu \otimes_1 dx^\nu + \frac{\lambda}{2} \left(\epsilon_{ijk} z^k z^a dz^i \otimes_1 dz^j - \frac{1}{(z^3)^2} \epsilon^a{}_{i3} dz^3 \otimes_1 dz^i \right)$$

Riemann Curvature tensor

$$\begin{aligned} \text{Riem}_1(dt) &= -\frac{1}{2}\widehat{R}^0_{\alpha\mu\nu}dx^\mu \wedge dx^\nu \otimes_1 dx^\alpha - \frac{\lambda}{2(z^3)^2} (\\ &\quad \epsilon_{3ij}z^3(F_3dt - F_4dr) \wedge dz^i \otimes_1 dz^j + \\ &\quad + \epsilon_{3ij}z^i(F_3dt - F_4dr) \wedge dz^j \otimes_1 dz^3) \end{aligned}$$

$$\begin{aligned} \text{Riem}_1(dr) &= -\frac{1}{2}\widehat{R}^1_{\alpha\mu\nu}dx^\mu \wedge dx^\nu \otimes_1 dx^\alpha + \frac{\lambda}{2(z^3)^2} (\\ &\quad + \epsilon_{3ij}z^3(F_5dt - F_3dr) \wedge dz^i \otimes_1 dz^j \\ &\quad + \epsilon_{3ij}z^i(F_5dt - F_3dr) \wedge dz^j \otimes_1 dz^3) \end{aligned}$$

$$\text{Riem}_1(dz^\mu) = -\frac{1}{2}\widehat{R}^\mu_{\alpha\mu\nu}dx^\mu \wedge dx^\nu \otimes_1 dx^\alpha + \frac{\lambda F_6}{2(z^3)^2}(1+(z^3)^3)dz^1 \wedge dz^2 \otimes_1 dz^\mu$$

Ricci Curvature tensor and Laplace operator

Ask that $\wedge_1(\text{Ricci}_1) = 0$ then

$$\begin{aligned}\text{Ricci}_1 &= -\frac{1}{2}\widehat{R}_{\mu\nu}dx^\mu \otimes_1 dx^\nu \\ &\quad -\frac{\lambda}{4(z^3)^2}(F_6 + F_5 - F_4)\epsilon_{3ij}(z^3 dz^i \otimes_1 dz^j - z^i dz^3 \otimes_1 dz^j) \\ &\quad -\frac{3\lambda}{4z^3}F_6\epsilon_{3ij}dz^i \otimes_1 dz^j\end{aligned}$$

and

$$S_1 = -\frac{1}{2}\widehat{S}$$

For Laplace operator

$$\square_1 f = g^{\alpha\beta} \left(f_{,\alpha\beta} + f_{,\gamma} \widehat{\Gamma}^\gamma_{\alpha\beta} \right)$$

Any Questions?

$$\begin{aligned}
\text{Riem}(dx^\mu) &= (d \otimes \text{id} - (\wedge \otimes \text{id})(\text{id} \otimes \nabla))(-\Gamma^\mu_{\alpha\beta} dx^\alpha \otimes dx^\beta) \\
&= -\Gamma^\mu_{\alpha\beta,\gamma} dx^\gamma \wedge dx^\alpha \otimes dx^\beta + \Gamma^\mu_{\alpha\beta} \Gamma^\beta_{\gamma\delta} dx^\alpha \wedge dx^\gamma \otimes dx^\delta \\
&= -(\Gamma^\mu_{\beta\gamma,\alpha} - \Gamma^\mu_{\alpha\delta} \Gamma^\delta_{\beta\gamma}) dx^\alpha \wedge dx^\beta \otimes dx^\gamma \\
&= -\frac{1}{2} R^\mu_{\gamma\alpha\beta} dx^\alpha \wedge dx^\beta \otimes dx^\gamma
\end{aligned}$$

$$\begin{aligned}
\text{Ricci} &= ((,) \otimes \text{id})(\text{id} \otimes i \otimes \text{id})(\text{id} \otimes \text{Riem}) g_{\mu\nu} dx^\mu \otimes dx^\nu \\
&= -\frac{1}{2} ((,) \otimes \text{id})(\text{id} \otimes i \otimes \text{id}) g_{\mu\kappa} R^\kappa_{\gamma\alpha\beta} dx^\mu \otimes dx^\alpha \wedge dx^\beta \otimes dx^\gamma \\
&= -\frac{1}{2} ((,) \otimes \text{id}) g_{\mu\kappa} R^\kappa_{\gamma\alpha\beta} dx^\mu \otimes dx^\alpha \otimes dx^\beta \otimes dx^\gamma \\
&= -\frac{1}{2} g_{\mu\kappa} g^{\mu\alpha} R^\kappa_{\gamma\alpha\beta} dx^\beta \otimes dx^\gamma \\
&= -\frac{1}{2} R_{\gamma\beta} dx^\beta \otimes dx^\gamma
\end{aligned}$$