

Pure state post-selection is universal

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Pure state pre- and postselection

Logical options:

$\mathcal{H}_i, \mathcal{H}_f, \mathcal{H}$ are factors of \mathcal{H}_{tot}

$|\Psi\rangle \in \mathcal{H}_{\text{tot}}$ total initial state

$|i\rangle \in \mathcal{H}_i$ preselection state (projection by $|i\rangle\langle i|$)

$|f\rangle \in \mathcal{H}_f$ postselection state (projection by $|f\rangle\langle f|$)

\mathcal{H} observed state space

Special cases:

$$1) \mathcal{H}_{\text{tot}} = \mathcal{H}_i = \mathcal{H}_f = \mathcal{H}$$

$$2) \mathcal{H}_{\text{tot}} = \mathcal{H}_i = \mathcal{H}_f = \mathcal{H} \otimes \mathcal{H}' \quad (\mathcal{H}': \text{ ancilla space})$$

SilvaGuryanovaBrunnerLindenShortPopescu PRA**89**,012121(2014)

2-time state

$$|i\rangle, |f\rangle \in \mathcal{H}; \quad \sum_{\mu} \hat{A}_{\mu}^{\dagger} \hat{A}_{\mu} = \hat{1} \quad \Longrightarrow \quad p(\mu, succ)$$

$$p(succ) = \sum_{\mu} p(\mu, succ), \quad p(\mu|succ) = \frac{p(\mu, succ)}{\sum_{\mu} p(\mu, succ)}$$

$$\begin{aligned} p(\mu, succ) &= |\langle f | \hat{A}_{\mu} | i \rangle|^2 = \mathbf{tr} (|f\rangle\langle f|) (\hat{A}_{\mu} |i\rangle\langle i| \hat{A}_{\mu}^{\dagger}) \\ &\equiv \mathbf{tr} (\hat{A}_{\mu} \otimes \hat{A}_{\mu}^{\dagger}) \hat{\rho}_{if} \end{aligned}$$

$$\begin{aligned} \text{2-time density : } \hat{\rho}_{if} &= |i\rangle\langle f| \otimes |f\rangle\langle i| \\ &\equiv \hat{\Psi}_{if} \otimes \hat{\Psi}_{if}^{\dagger} \end{aligned}$$

$$\text{2-time pure state : } \hat{\Psi}_{if} = |i\rangle\langle f|$$

So far $\hat{\Psi}_{if} = |i\rangle\langle f|$ is unentangled, $\hat{\rho}_{if} = \hat{\Psi}_{if} \otimes \hat{\Psi}_{if}^{\dagger}$ is not mixed (pure).

General 2-time state

$$p(\mu, \text{succ}) = \text{tr} (\hat{A}_\mu \otimes \hat{A}_\mu^\dagger) \hat{\rho}_{\text{if}}$$

Success becomes independent of measurement if it is weak:

$$p_{\text{WM}}(\text{succ}) = \text{tr} \hat{\rho}_{\text{if}}$$

Definitive math conditions for $\hat{\rho}_{\text{if}}$:

$$1) \text{tr}(\hat{V} \otimes \hat{V}^\dagger) \hat{\rho}_{\text{if}} \geq 0, \quad \forall \hat{V}; \quad 2) \text{tr} \hat{\rho}_{\text{if}} \leq 1$$

Standard form:

$$\hat{\rho}_{\text{if}} = \sum_r \hat{\Psi}_{\text{if}}^r \otimes \hat{\Psi}_{\text{if}}^{r\dagger}, \quad \text{tr} \hat{\Psi}_{\text{if}}^{r\dagger} \hat{\Psi}_{\text{if}}^s = 0 \quad (r \neq s)$$

Can we prepare all $\hat{\rho}_{\text{if}}$? Yes, upto a prefactor since

$$\text{tr} \hat{\rho}_{\text{if}} = p_{\text{WM}}(\text{succ})$$

Preparing entangled 2-time pure state

$$|i\rangle, |f\rangle \in \mathcal{H} \otimes \mathcal{H}'$$

$$p(\mu, \text{succ}) = \text{tr}(|f\rangle\langle f|)(\hat{A}_\mu \otimes \hat{1}')|i\rangle\langle i|(\hat{A}_\mu^\dagger \otimes \hat{1}') \equiv \text{tr}(\hat{A}_\mu \otimes \hat{A}_\mu^\dagger)\hat{\rho}_{if}$$

$$\hat{\rho}_{if} = (\text{tr}'|i\rangle\langle f|) \otimes (\text{tr}'|f\rangle\langle i|) \equiv \hat{\Psi}_{if} \otimes \hat{\Psi}_{if}^\dagger$$

$$\hat{\Psi}_{if} = \text{tr}'|i\rangle\langle f|$$

In coordinates:

$$|i\rangle = \sum_{k,r} c_{kr}^i |k\rangle \otimes |r\rangle', \quad |f\rangle = \sum_{k,r} c_{kr}^f |k\rangle \otimes |r\rangle'$$

$$\hat{\Psi}_{if} = \text{tr}'|i\rangle\langle f| = \text{tr}' \sum_{k,r} c_{kr}^i |k\rangle \otimes |r\rangle' \sum_{l,s} c_{ls}^{f*} \langle l| \otimes \langle s|'$$

$$= \sum_{kl} c_{kl}^{if} |k\rangle\langle l|$$

Can we reach any 2-time amplitudes c_{kl}^{if} ?

Can we prepare any 2-time entangled pure state?

From previous slide:

$$|i\rangle = \sum_{k,r} c_{kr}^i |k\rangle \otimes |r\rangle', \quad |f\rangle = \sum_{k,r} c_{kr}^f |k\rangle \otimes |r\rangle' \quad \implies \hat{\Psi}_{if} = \sum_{kl} c_{kl}^{if} |k\rangle \langle l|$$

$$p_{\text{WM}}(\text{succ}) = |\text{tr} \hat{\Psi}_{if}|^2 = |\text{tr} c^{if}|^2$$

$$c^{if} = c^i c^{f\dagger} \quad (\text{recall: } \text{tr} c^{i\dagger} c^i = \text{tr} c^{f\dagger} c^f = 1)$$

2-time amplitude c^{if} is subnormalized unless $c^i = c^f$ (upto a phase).

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Silva et al. prepare any $\hat{\Psi}$ upto a factor: $c^f = 1/\sqrt{d}$ yields $c^{if} = c^i/\sqrt{d}$

$p_{\text{WM}}(\text{succ}) = |\text{tr} c^i|^2/d$ can be suboptimal.

If $c^{if} \geq 0$, then we can choose $c^i = c^f = \sqrt{c^{if}/\text{tr} c^{if}}$ at

$p_{\text{WM}}(\text{succ}) = 1$.

Open issue: To prepare a given $\hat{\Psi}$ what are the 'closest' $|i\rangle$ and $|f\rangle$ to guarantee the highest $p_{\text{WM}}(\text{succ})$?

Mixed 2-time state

$$|i\rangle, |f\rangle \in \mathcal{H} \otimes \mathcal{H}', \quad \sum_{\mu} \hat{A}_{\mu}^{\dagger} \hat{A}_{\mu} = \hat{1}, \quad \sum_{\mu} \hat{A}'_{\mu}{}^{\dagger} \hat{A}'_{\mu} = \hat{1}'$$

$$p(\mu, r, succ) = \mathbf{tr} (|f\rangle\langle f|)(\hat{A}_{\mu} \otimes \hat{A}'_r) |i\rangle\langle i| (\hat{A}_{\mu}^{\dagger} \otimes \hat{A}'_r{}^{\dagger})$$

$$p(\mu, succ) = \sum_r p(\mu, r, succ) \equiv \mathbf{tr}(\hat{A}_{\mu} \otimes \hat{A}'_{\mu}{}^{\dagger}) \hat{\rho}_{if}$$

$$\hat{\rho}_{if} = \sum_r \left\{ \left(\mathbf{tr}'(\hat{1} \otimes \hat{A}'_r) |i\rangle\langle f| \right) \otimes \left(\mathbf{tr}'(\hat{1} \otimes \hat{A}'_r{}^{\dagger}) |f\rangle\langle i| \right) \right\}$$

$$\equiv \sum_r \hat{\Psi}_{if}^r \otimes \hat{\Psi}_{if}^{r\dagger}$$

Example:

$$|i\rangle = \sum_r \sqrt{p_r} |i; r\rangle \otimes |r\rangle', \quad |f\rangle = \frac{1}{\sqrt{d'}} \sum_r |f; r\rangle \otimes |r\rangle', \quad \hat{A}'_r = |r\rangle'\langle r|'$$

$$\hat{\rho}_{if} = \frac{1}{d'} \sum_r p_r |i; r\rangle\langle f; r| \otimes |f; r\rangle\langle i; r|$$

Issue: Pure state pre/postselection has $1/d'$ -times smaller $p_{\text{WM}}(succ)$ vs Silva et al.

Inferring success rate without tomography

Silva et al.: For mixed 2-states tomography, projective and WM's are insufficient, generalized measurements are needed, they constructed one.

Assume pure 2-states! Both projective and WM's remain insufficient. What WM's are sufficient for?

WMs yield $p_{\text{WM}}(\text{succ})$ without tomography.

Example: AAV spin weak value

$$\hat{\Psi} = |i\rangle\langle f| = |\vec{n}_i\rangle\langle\vec{n}_f|, \quad p_{\text{WM}}(\text{succ}) = |\text{tr}\hat{\Psi}|^2 = \frac{1}{2}(1 + \vec{n}_i\vec{n}_f)$$

$$\vec{\sigma}_W = \text{Re} \frac{\text{tr}\hat{\sigma}\hat{\Psi}}{\text{tr}\hat{\Psi}} = \frac{\vec{n}_i + \vec{n}_f}{1 + \vec{n}_i\vec{n}_f}, \quad |\vec{\sigma}_W|^2 = \frac{1}{p_{\text{WM}}(\text{succ})}$$

$$p_{\text{WM}}(\text{succ}) = \frac{1}{(\sigma_{xW})^2 + (\sigma_{yW})^2 + (\sigma_{zW})^2}$$

Summary

- 2-time density $\hat{\rho}_{\text{if}}$ introduced for $p(\mu, \text{succ}) = \text{tr}(\hat{A}_\mu \otimes \hat{A}_\mu^\dagger) \hat{\rho}_{\text{if}}$

$$p_{\text{WM}}(\text{succ}) = \text{tr} \hat{\rho}_{\text{if}}$$

All $\hat{\rho}_{\text{if}}$ are, upto normalization $p_{\text{WM}}(\text{succ})$, preparable via pure state pre- and postselection.

- $p_{\text{WM}}(\text{succ})$ is a figure of merit of preparation. Optimum preparation protocols remain to be found.
- In the simple VVA case, weak measurements provide $p_{\text{WM}}(\text{succ})$ without tomography:

$$p_{\text{WM}}(\text{succ}) = \frac{1}{(\sigma_{xW})^2 + (\sigma_{yW})^2 + (\sigma_{zW})^2}$$