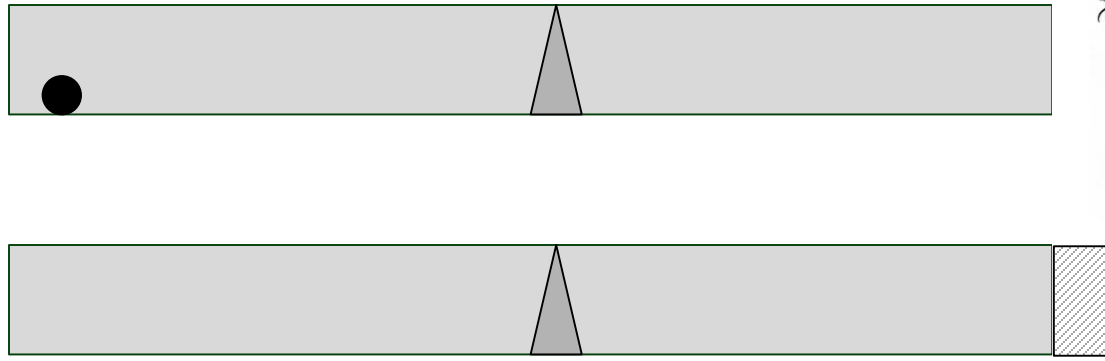
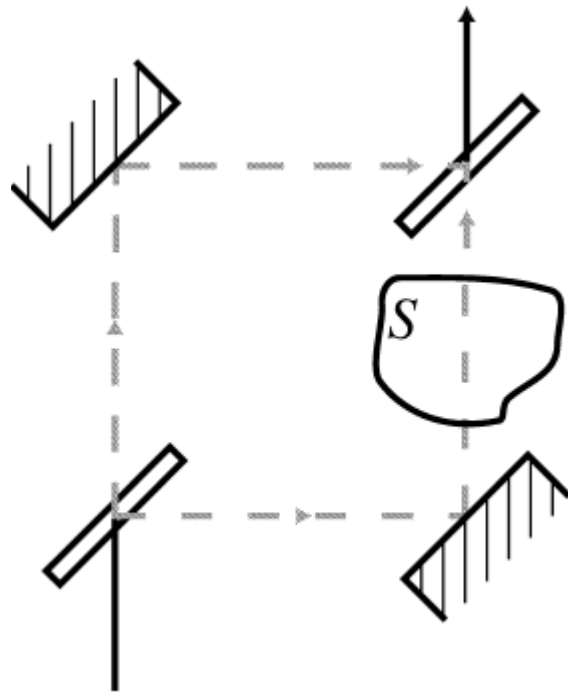




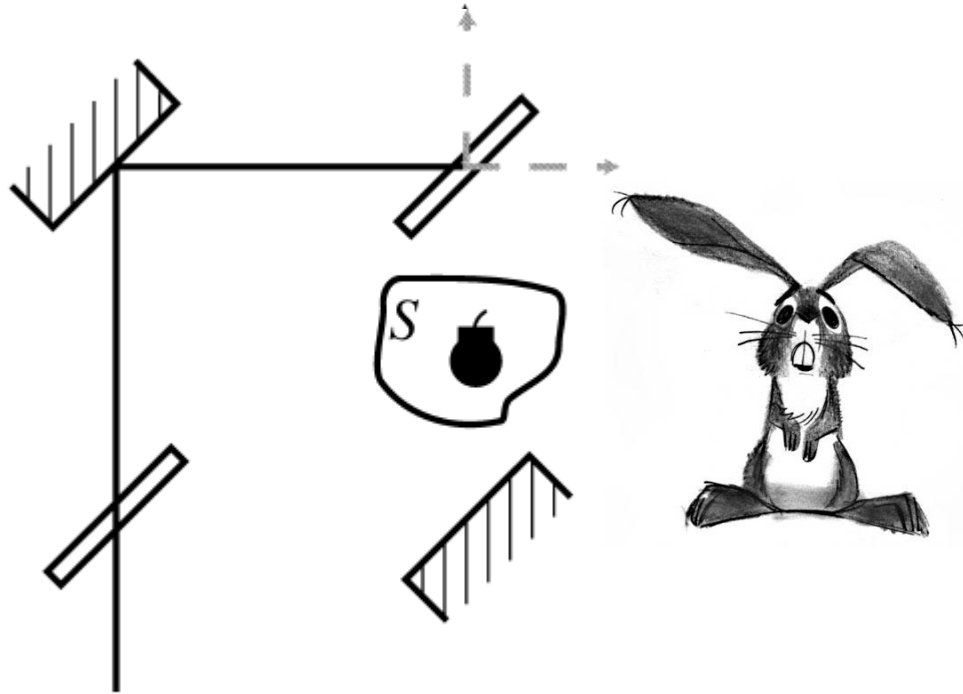
# Locality and nonlocality in “interaction-free” measurements



Y. Aharonov, T. Landsberger and D. Rohrlich (2017)



A. Elitzur and L. Vaidman, *Found. Phys.* **23**, 987 (1993)



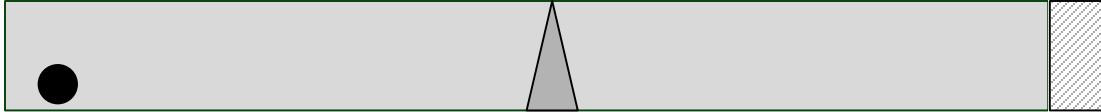
A. Elitzur and L. Vaidman, *Found. Phys.* **23**, 987 (1993)

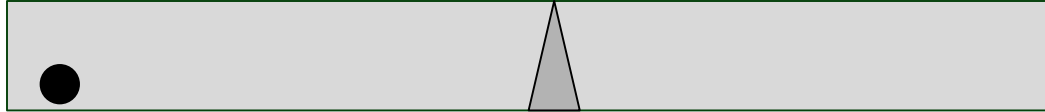
## Outline

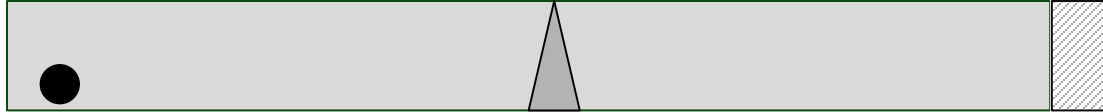
- What is “counterfactual communication”?
- A local “interaction-free measurement”
- “Counterfactual communication” revisited

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- What is “counterfactual communication”?
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- “Counterfactual communication” revisited

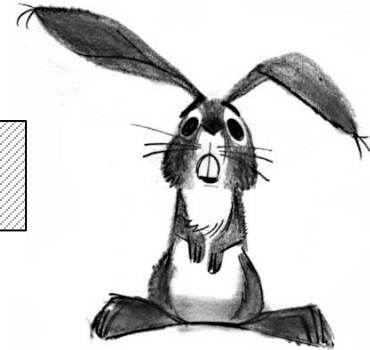
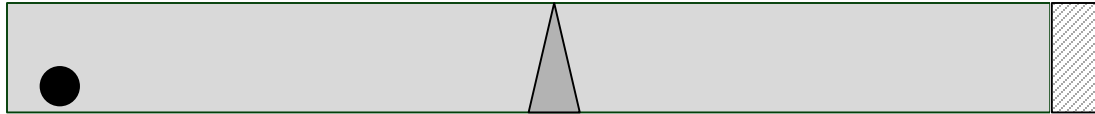






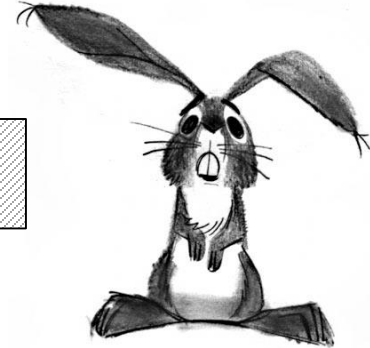
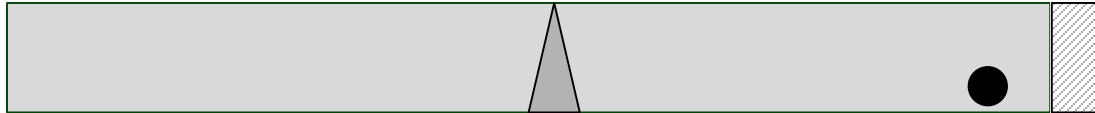
$$U(\epsilon) = \begin{pmatrix} \cos \epsilon & i \sin \epsilon \\ i \sin \epsilon & \cos \epsilon \end{pmatrix}$$





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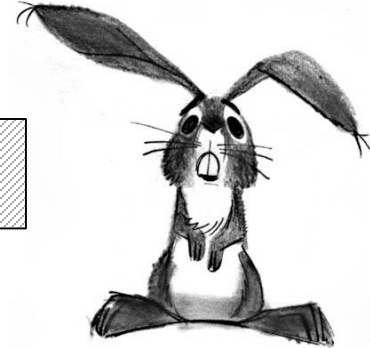
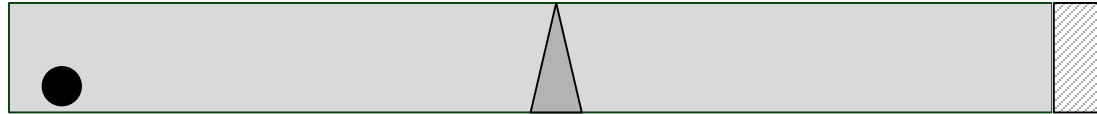
$$U^2(\epsilon) = \begin{pmatrix} \cos 2\epsilon & i \sin 2\epsilon \\ i \sin 2\epsilon & \cos 2\epsilon \end{pmatrix}$$



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Let  $n\epsilon = \pi/2$  at time  $T$ . Then the ball is in Bob's side.

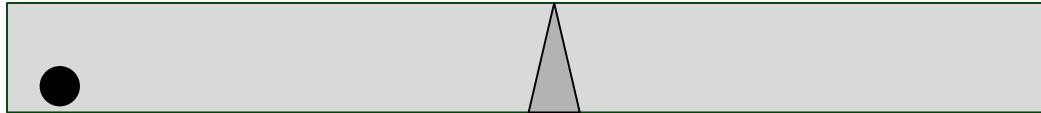


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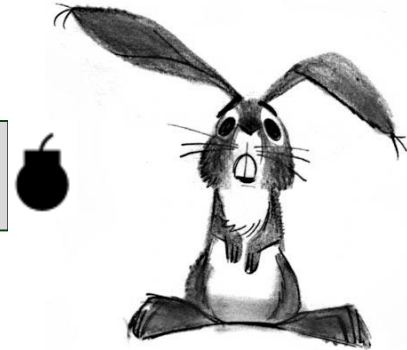
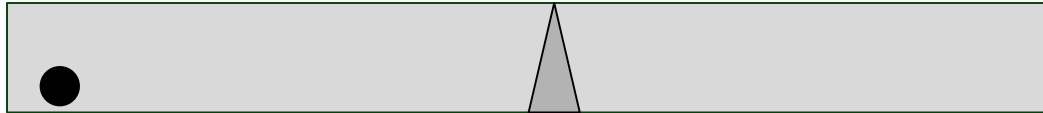
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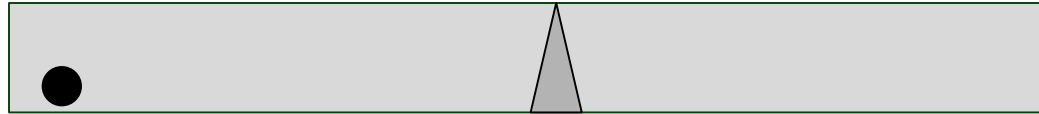
At time  $2T$ , the ball is back on Alice's side, multiplied by  $-1$ .



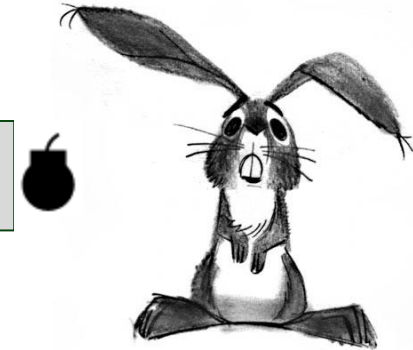
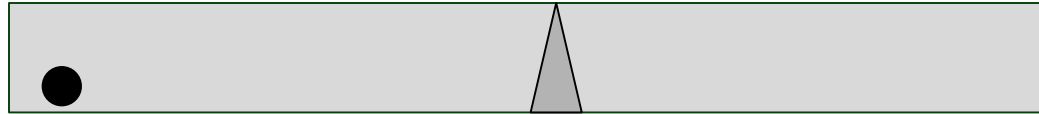
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If Bob's end is open (no mirror) – or if Bob puts a bomb there instead of a mirror –



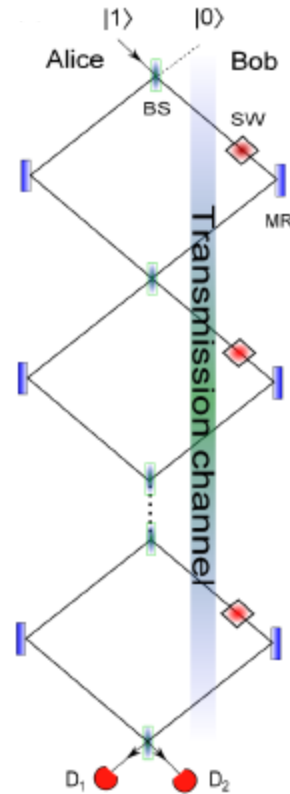
If Bob's end is open (no mirror) – or if Bob puts a bomb there instead of a mirror – then the chance of finding the ball on Bob's side at time  $T$  vanishes (amplitude  $\epsilon$  and probability  $\epsilon^2$  per event, number of events proportional to  $n = \pi/2\epsilon$ ).



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Thus the quantum Zeno effect makes the interaction-free measurement 100% reliable, and Bob can signal to Alice with no exchange of matter.

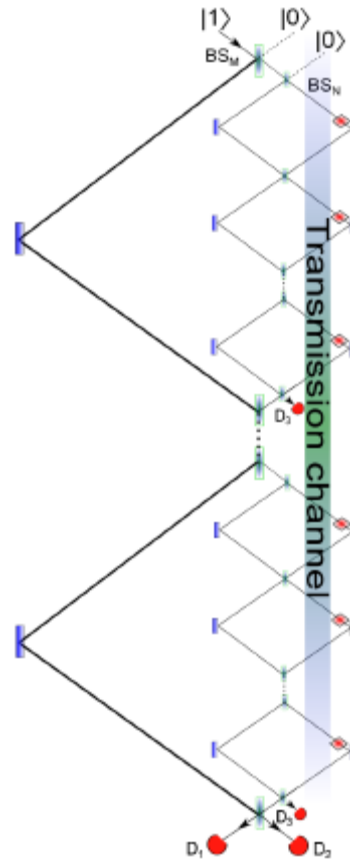
# “Counterfactual communication”



H. Salih *et al.*, *Phys. Rev Lett.* **110**, 170502 (2013)



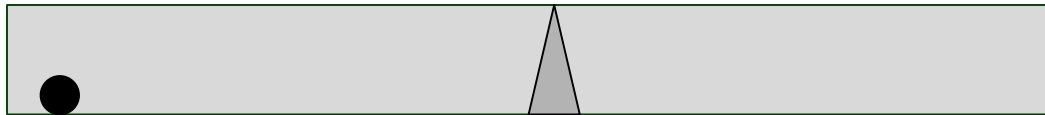
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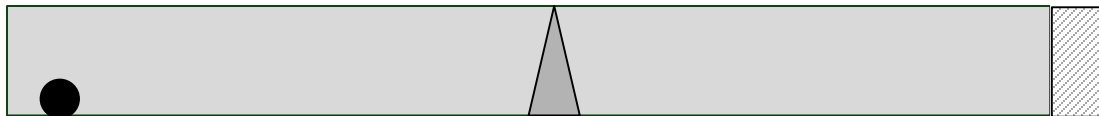
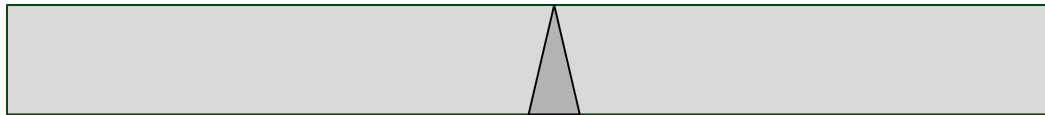


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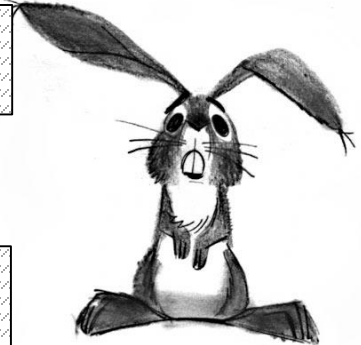
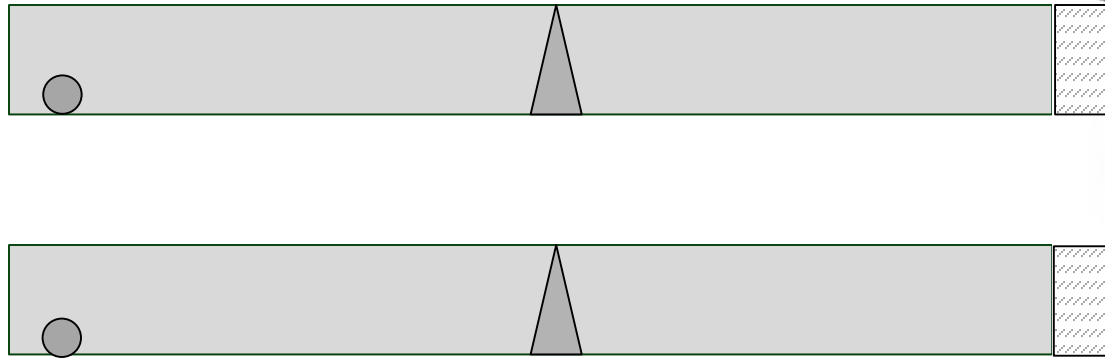
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- What is “counterfactual communication”?
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- “Counterfactual communication” revisited



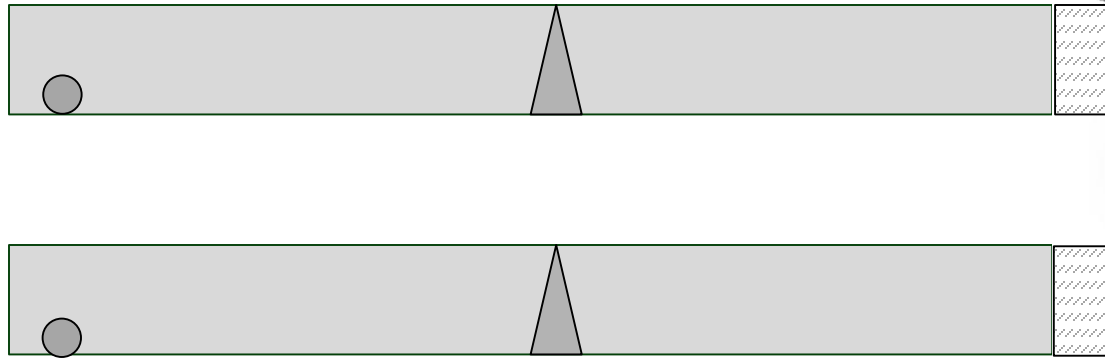


$$\Psi(0) = \frac{1}{2} (|\uparrow\rangle_A + |\downarrow\rangle_A)(|\uparrow\rangle_B + |\downarrow\rangle_B)$$



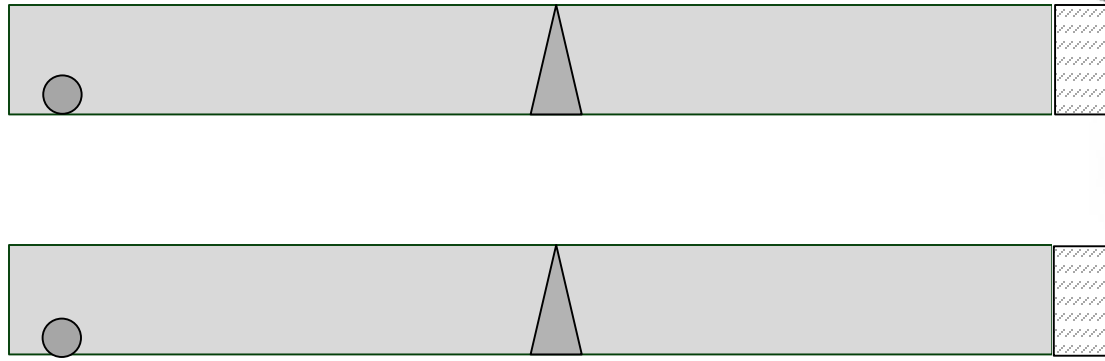
$\Psi(0)$  is an eigenstate at time  $t = 0$  of the “modular angular momenta”  $L^A_{\text{mod}}$  and  $L^B_{\text{mod}}$

$$\Psi(T) = \frac{1}{\sqrt{2}} [|\uparrow\rangle_A |\downarrow\rangle_B + |\downarrow\rangle_A |\uparrow\rangle_B]$$



$\Psi(0)$  is an eigenstate at time  $t = 0$  of the “modular angular momenta”  $L^A_{\text{mod}}$  and  $L^B_{\text{mod}}$ , and then at time  $t = T$  Alice post-selects the ball on her side.

$$\Psi(T) = \frac{1}{\sqrt{2}} [|\uparrow\rangle_B |\downarrow\rangle_M + |\downarrow\rangle_B |\uparrow\rangle_M]$$

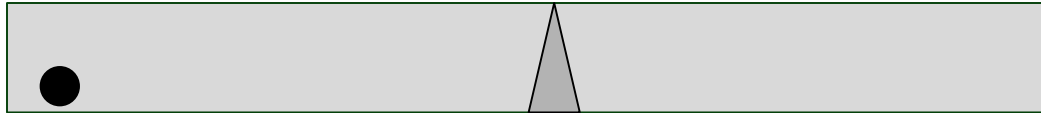


We find a *local* flow of a nonlocal physical quantity – modular angular momentum – on this pre- and post-selected ensemble.

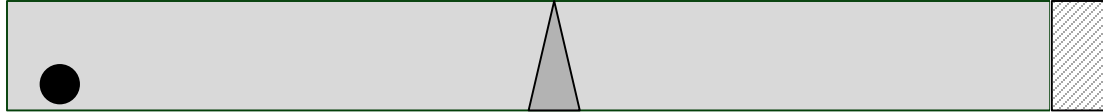
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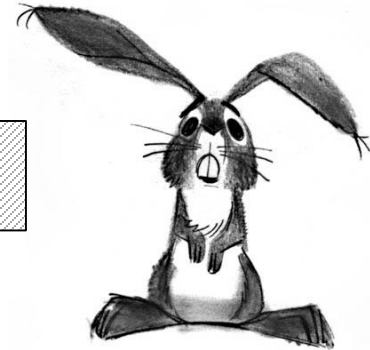
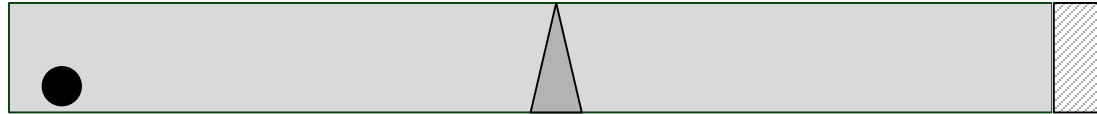




Alice finds the particle at her end after time  $T$ .



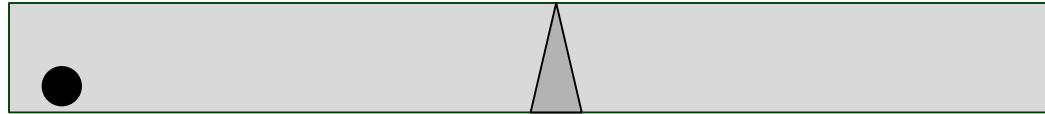
Alice *does not* find the particle at her end after time  $T$ .



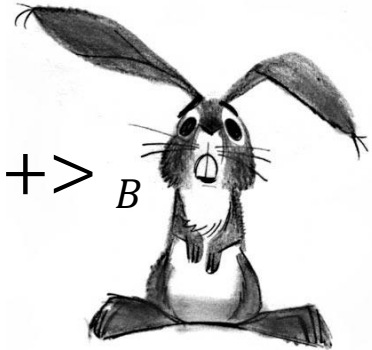
Alice *does not* find the particle at her end after time  $T$ .  
Now let's define states

$$|+\rangle_B = \frac{1}{\sqrt{2}} \left[ \begin{array}{|c|} \hline \text{hatched} \\ \hline \end{array} + \begin{array}{|c|} \hline \text{particle} \\ \hline \end{array} \right] = \frac{1}{\sqrt{2}} \left[ \begin{array}{|c|} \hline \text{hatched} \\ \hline \end{array} + \begin{array}{|c|} \hline \text{bomb} \\ \hline \end{array} \right]$$

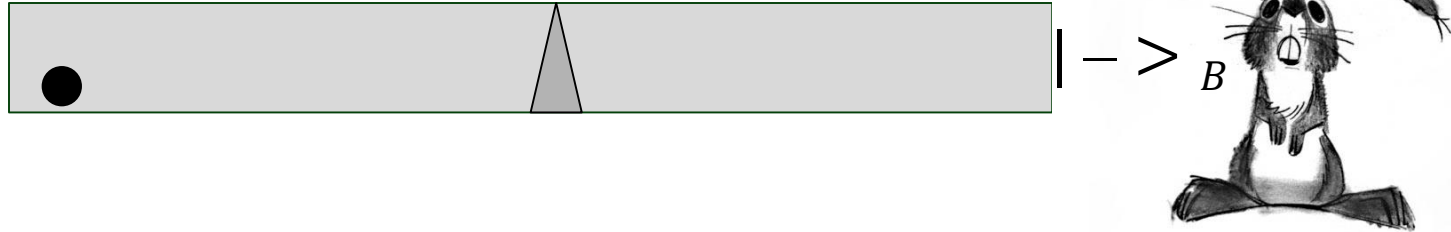
$$|-\rangle_B = \frac{1}{\sqrt{2}} \left[ \begin{array}{|c|} \hline \text{hatched} \\ \hline \end{array} - \begin{array}{|c|} \hline \text{particle} \\ \hline \end{array} \right] = \frac{1}{\sqrt{2}} \left[ \begin{array}{|c|} \hline \text{hatched} \\ \hline \end{array} - \begin{array}{|c|} \hline \text{bomb} \\ \hline \end{array} \right]$$



$| + \rangle_B$



So now if Bob prepares his mirror in the state  $| + \rangle_B$ , after time  $2T$  it turns into  $| - \rangle_B$  and vice versa – not because of “counterfactual communication” but because the ball acquires a *local* phase along the way.



Conclusion: Bob may prepare the mirror *only* in the state  $\frac{1}{\sqrt{2}} [ | + \rangle_B + | - \rangle_B ]$  and thus *only* see “counterfactual communication”, but measurement in any other basis such as  $| \pm \rangle_B$  shows that the ball does reach Bob’s end.