

Development of Far-infrared Spectrophotometers based on Superconducting Tunnel Junction (STJ) for COBAND Experiment



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Yuji Takeuchi (Univ. of Tsukuba)

on behalf of COBAND Collaboration

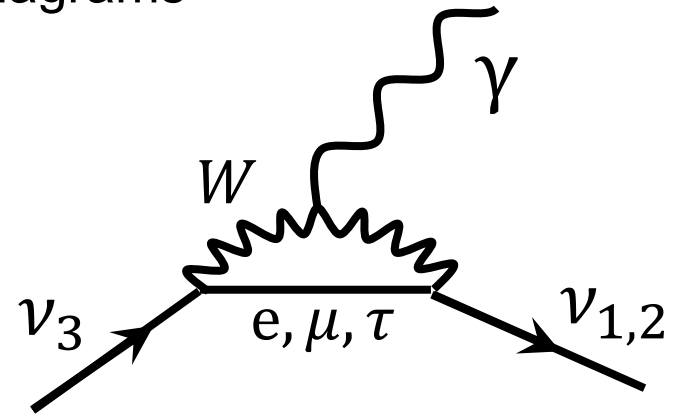
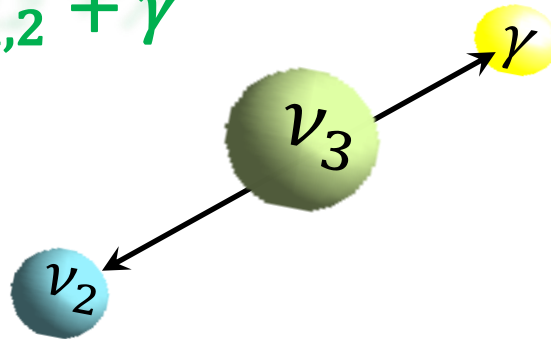
COBAND (COsmic BAckground Neutrino Decay)



□ Heavier neutrinos in mass-eigenstate (ν_2, ν_3) are not stable

– Neutrino can decay through the loop diagrams

– $\nu_3 \rightarrow \nu_{1,2} + \gamma$



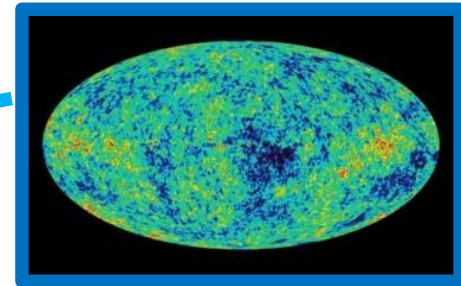
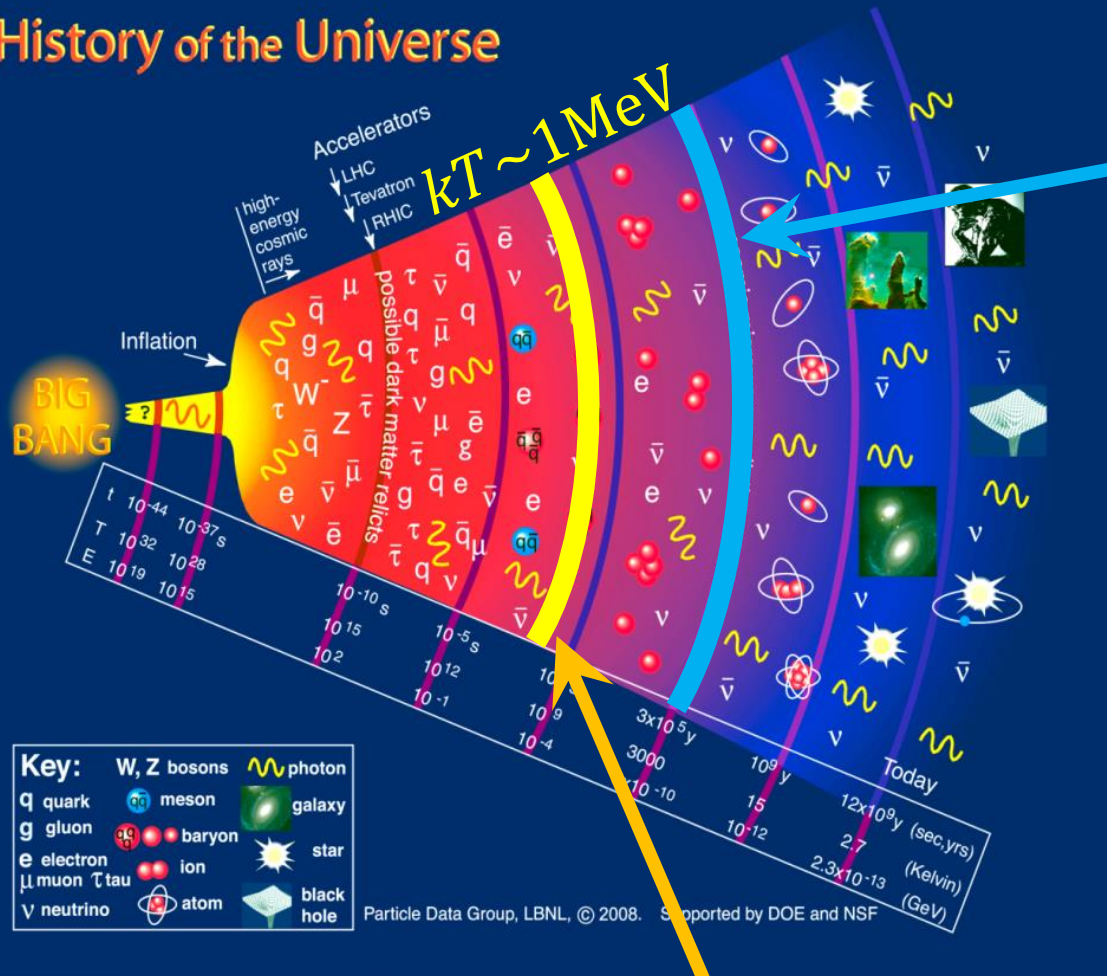
✓ However, the lifetime is expected to be much longer than the age of the universe

→ We search for neutrino decay using Cosmic Background Neutrino (C ν B) as the neutrino source

$$\rho(\nu_3 + \bar{\nu}_3) \sim 110/\text{cm}^3$$

Cosmic background neutrino (CνB)

History of the Universe



CMB

(=Photon decoupling)

$$n_{\gamma} = 411/\text{cm}^3$$

$$T_{\gamma} = 2.73 \text{ K}$$

~380,000yrs after the Big Bang

CνB (=neutrino decoupling)
~1sec after the big bang

COBAND Collaboration Members (As of Nov. 2016)

Shin-Hong Kim, Yuji Takeuchi, Kenichi Takemasa, Kazuki Nagata, Kota Kasahara, Shunsuke Yagi, Rena Wakasa, Yoichi Otsuka (Univ. of Tsukuba), Hirokazu Ikeda, Takehiko Wada, Koichi Nagase (JAXA/ISAS), Shuji Matsuura (Kwansei gakuin Univ), Yasuo Arai, Ikuo Kurachi, Masashi Hazumi (KEK), Takuo Yoshida, Chisa Asano, Takahiro Nakamura, Makoto Sakai (Univ. of Fukui), Satoshi Mima, Kenji Kiuchi (RIKEN), H.Ishino, A.Kibayashi (Okayama Univ.), Yukihiro Kato (Kindai University), Go Fujii, Shigetomo Shiki, Masahiro Ukibe, Masataka Ohkubo (AIST), Shoji Kawahito (Shizuoka Univ.), Erik Ramberg, Paul Rubinov, Dmitri Sergatskov (Fermilab), Soo-Bong Kim (Seoul National University)

Motivation of ν -decay search in $C\nu B$

ν_3 Lifetime

- Standard Model expectation: $\tau = O(10^{43})$ yrs
- Experimental lower limit: $\tau = O(10^{12})$ yrs
- Left-Right symmetric model predicts $\tau = O(10^{17})$ yrs for W_L - W_R mixing angle $|\zeta| \sim 0.02$

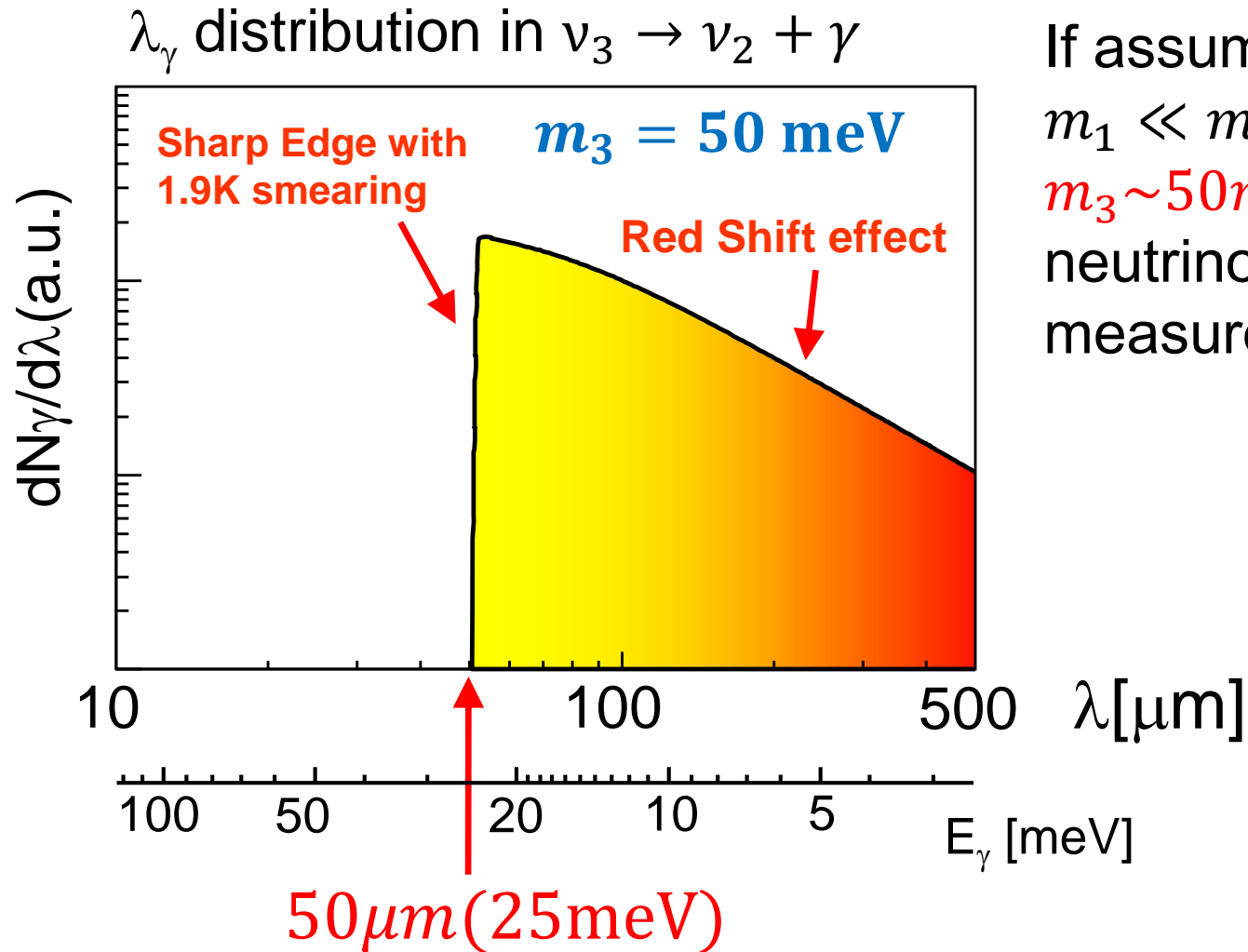
If we observed the neutrino radiative decay at the lifetime much shorter than the SM expectation, it would be

- Physics beyond the Standard Model
- Direct detection of $C\nu B$
- Determination of the neutrino mass

$$- m_3 = (m_3^2 - m_{1,2}^2)/(2E_\gamma)$$

→ Aiming at a sensitivity to ν_3 lifetime in $O(10^{13} - 10^{17})$ yrs

Expected photon wavelength spectrum from CνB decays

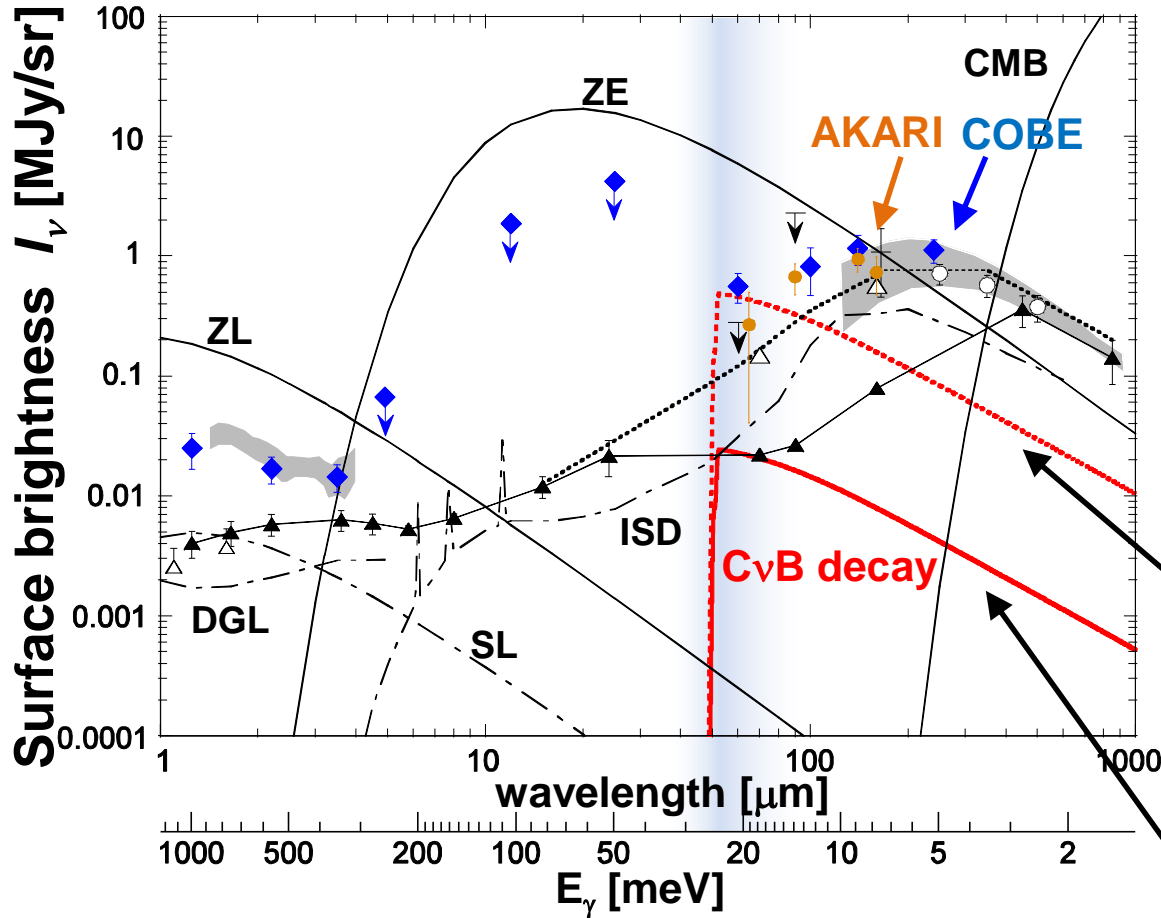


If assume

$m_1 \ll m_2 < m_3$,
 $m_3 \sim 50 \text{ meV}$ from
neutrino oscillation
measurements

No other source has such a sharp edge structure!!

CνB radiative decay and Backgrounds



Zodiacal Emission
 $I_\nu \sim 8 \text{ MJy/sr}$

Cosmic Infrared Background (CIB)
 $I_\nu \sim 0.1 \sim 0.5 \text{ MJy/sr}$

CνB decay
 $\tau = 5 \times 10^{12} \text{ yrs}$
 $I_\nu \sim 0.5 \text{ MJy/sr}$
Excluded (S.H.Kim 2012)

Expected E_γ spectrum
 for $m_3 = 50 \text{ meV}$

$\tau = 1 \times 10^{14} \text{ yrs}$
 $I_\nu \sim 25 \text{ kJy/sr}$

at $\lambda = 50 \mu\text{m}$

$1 \text{ Jy} = 10^{-26} \text{ W/m}^2 \cdot \text{Hz}$

Proposal for COBAND Rocket Experiment

Aiming at a sensitivity to ν lifetime for
 $\tau(\nu_3) = O(10^{14})$ yrs

■ JAXA sounding rocket S-520

- <http://www.isas.jaxa.jp/e/enterp/rockets/sounding/s520.shtml>
- Diameter: 520mm
- Payload: 100kg
- Altitude: 300km

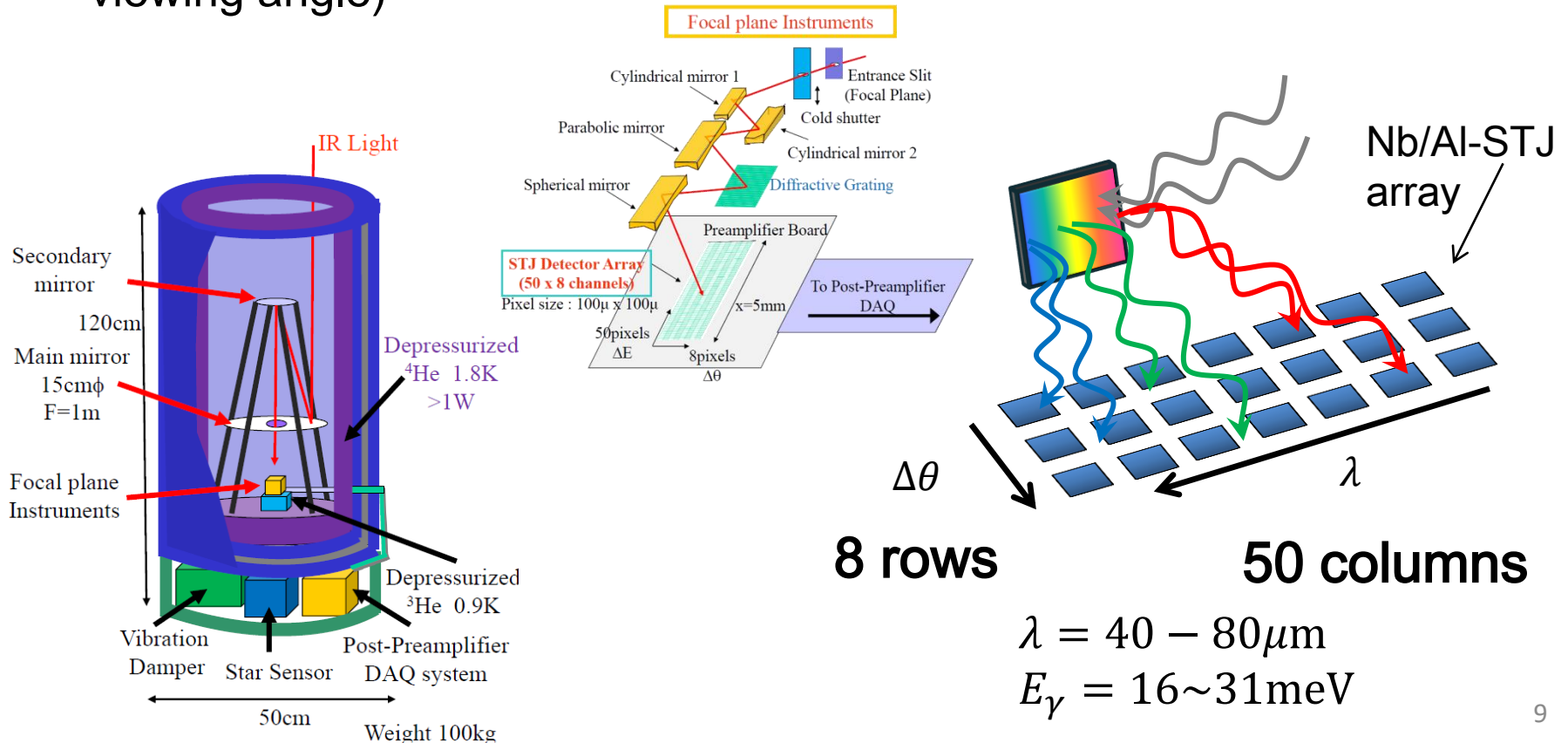


- 200-sec measurement at altitude of 200~300km
 - Telescope with **diameter of 15cm** and **focal length of 1m**
 - All optics (mirrors, filters, shutters and grating) will be cooled at $\sim 1.8\text{K}$

Proposal for COBAND Rocket Experiment

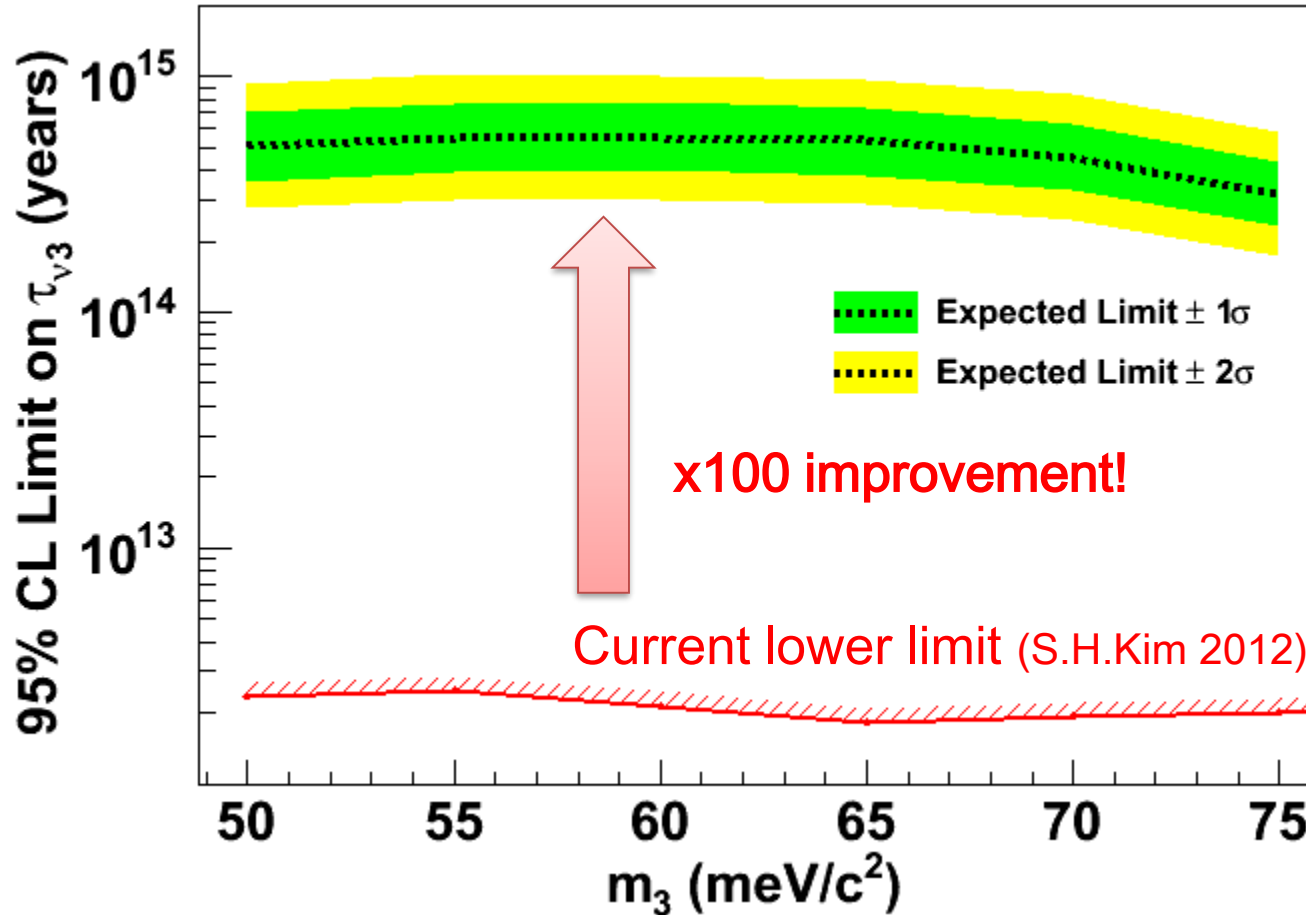
At the focal point, diffraction grating covering $\lambda=40-80\mu\text{m}$ (16-31meV) and array of photo-detector pixels of 50(in wavelength distribution) x 8(in spatial distribution) are placed

- Each pixel **counts FIR photons** in $\lambda=40-80\mu\text{m}$ ($\Delta\lambda = 0.8\mu\text{m}$)
- Sensitive area of $100\mu\text{m} \times 100\mu\text{m}$ for each pixel ($100\mu\text{rad} \times 100\mu\text{rad}$ in viewing angle)



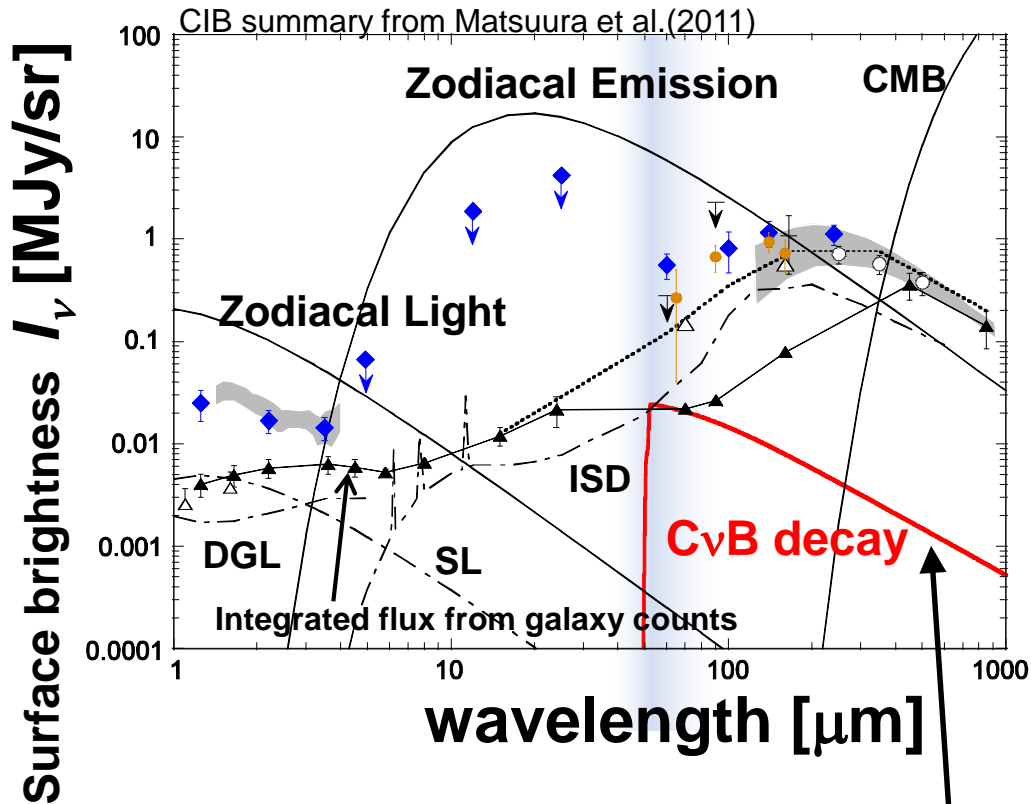
Sensitivity to neutrino decay

Parameters in the rocket experiment simulation are assumed.



- Can set lower limit on ν_3 lifetime at $4-6 \times 10^{14}$ yrs if no neutrino decay
- If ν_3 lifetime were 2×10^{14} yrs, the signal significance would be at 5σ level

Requirements for the photo-detector in the rocket experiment



Zodiacal Emission

- $I_{ZE} = 8 \text{ MJy/sr}$
- $1.1 \times 10^{-17} \text{ W / 8pix}$
- @ $\lambda = 50 \mu\text{m}$

Neutrino Decay

- $I_\nu = 25 \text{ kJy/sr}$
- $3.3 \times 10^{-20} \text{ W / 8pix}$
- @ $\lambda = 50 \mu\text{m}$

$$\tau = 1 \times 10^{14} \text{ yrs}$$

$$m_3 = 50 \text{ meV}$$

Noise Equivalent Power (NEP) Requirements for the photo-detector

- Neutrino decay ($m_3 = 50 \text{ meV}$, $\tau_\nu = 10^{14} \text{ yrs}$): $I_\nu = 25 \text{ kJy/sr}$ @ $\lambda = 50 \mu\text{m}$

$$P_{ND} = 25 \text{ kJy/sr} \times 8 \times 10^{-8} \text{ sr} \times \pi(15 \text{ cm}/2)^2 \times \Delta\nu \\ = 3.3 \times 10^{-20} \text{ W}/8 \text{ pix}$$

- Zodiacal emission: $I_\nu = 8 \text{ MJy/sr}$ @ $\lambda = 50 \mu\text{m}$

$$P_{ZE} = 1.1 \times 10^{-17} \text{ W}/8 \text{ pix}$$

- ◆ Shot noise in P_{ZE} integrated over an interval Δt

– Fluctuation in number of photons with energy ϵ_γ : $\sqrt{\epsilon_\gamma P_{ZE} \Delta t}$

$$\frac{NEP}{\sqrt{2\Delta t}} \times \Delta t \ll \sqrt{\epsilon_\gamma P_{ZE} \Delta t} \ll P_{ND} \Delta t$$

→ $\Delta t > 200 \text{ sec}$

→ $NEP \sim 0(10^{-20}) \text{ W}/\sqrt{\text{Hz}}$ for 1pix

Existing FIR photo-detectors

Detectors	$\lambda(\mu\text{m})$	Operation Temp.	NEP ($\text{W}/\text{Hz}^{1/2}$)	
Monolithic Ge:Ga	50-110	2.2K	$\sim 10^{-17}$	Akari-FIS
Stressed Ge:Ga	60-210	0.3K	$\sim 0.9 \times 10^{-17}$	Herschel-PACS

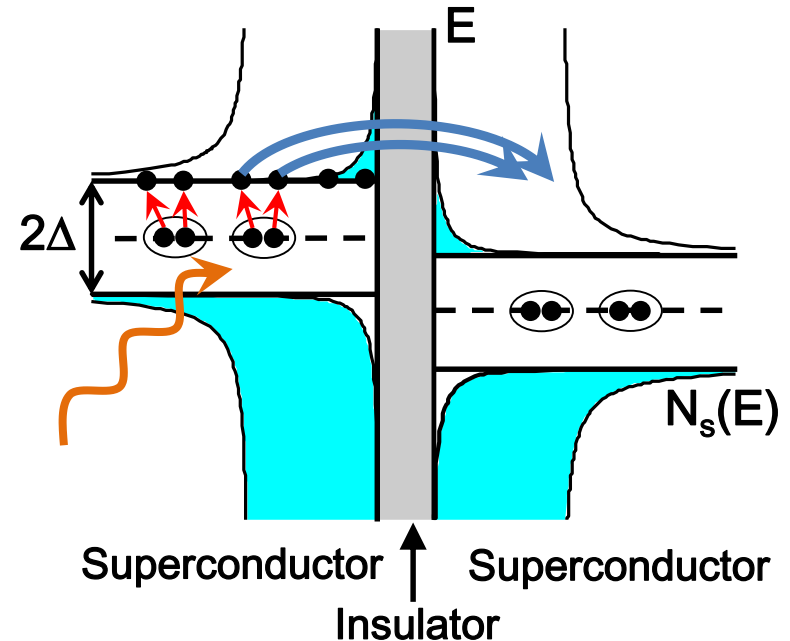
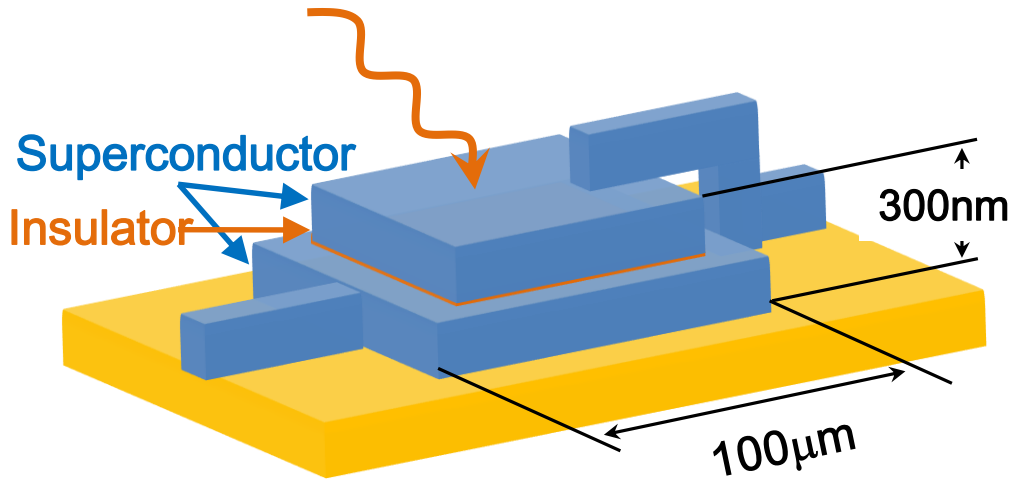
Need **more than 2 orders improvement** from existing photoconductor-based detectors

We are trying to achieve $\text{NEP} \sim 0(10^{-20}) \text{ W}/\sqrt{\text{Hz}}$ by using

- **Superconducting tunneling junction detector**
- **FIR single-photon counting technique**

Superconducting Tunnel Junction (STJ) Detector

- Superconductor / **Insulator** / Superconductor Josephson junction device



Δ : Superconducting gap energy

A constant bias voltage ($|V| < 2\Delta$) is applied across the junction.
A photon absorbed in the superconductor breaks Cooper pairs and creates tunneling current of quasi-particles proportional to the deposited photon energy.

- Much lower gap energy (Δ) than FIR photon → **Can detect FIR photon**
- Faster response ($\sim \mu\text{s}$) → **Suitable for single-photon counting**

STJ energy resolution

Signal = Number of quasi-particles

$$N_{q.p.} = G \frac{E_\gamma}{1.7\Delta}$$

Resolution = Statistical fluctuation in number of quasi-particles

$$\sigma_E/E = \sqrt{(1.7\Delta)F/E}$$

→ Smaller superconducting gap energy Δ yields better energy resolution

Δ : Superconducting gap energy
F: fano factor (~ 0.2 for Nb)
E: Photon energy
G: Back-tunneling gain

	Si	Nb	Al	Hf
Tc[K]		9.23	1.20	0.165
Δ [meV]	1100	1.550	0.172	0.020

Tc :SC critical temperature
Need $\sim 1/10T_c$ for practical operation

STJ candidates

Nb/Al-STJ

- Well-established
 - $\Delta \sim 0.6 \text{ meV}$ by the proximity effect from Al
 - Operation temperature $< 400 \text{ mK}$
 - Back-tunnelling gain $G \sim 10$
- $N_{\text{q.p.}} = 25 \text{ meV} / 1.7\Delta \times 10 \sim 250$ $\sigma_E/E \sim 0.1$ for $E = 25 \text{ meV}$
- 25 meV single-photon detection is feasible in principle
 - Developing for the rocket experiment

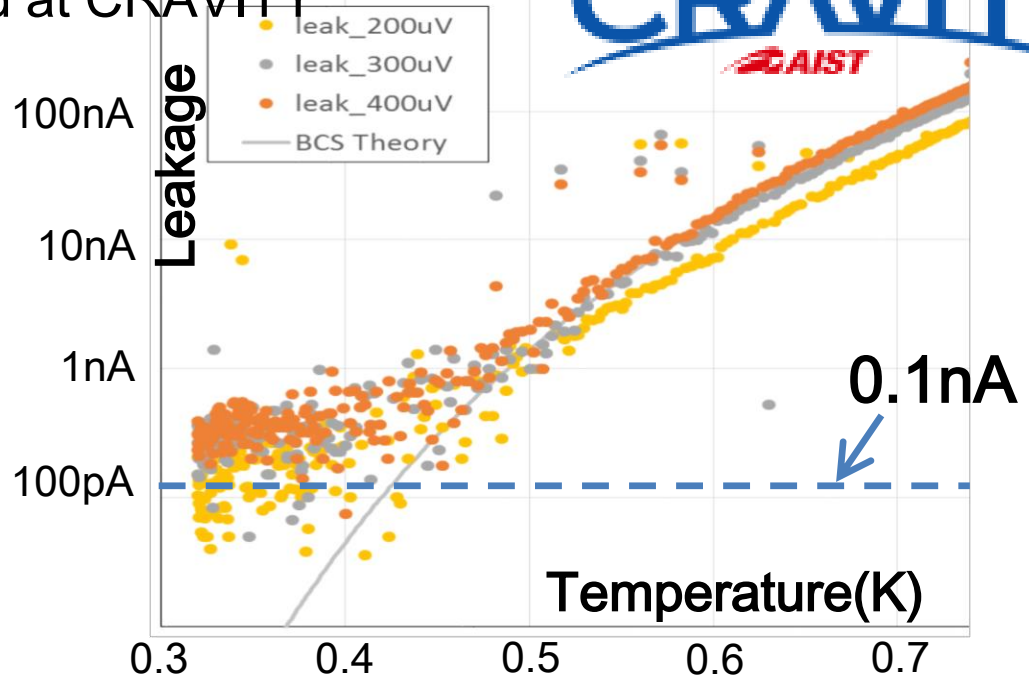
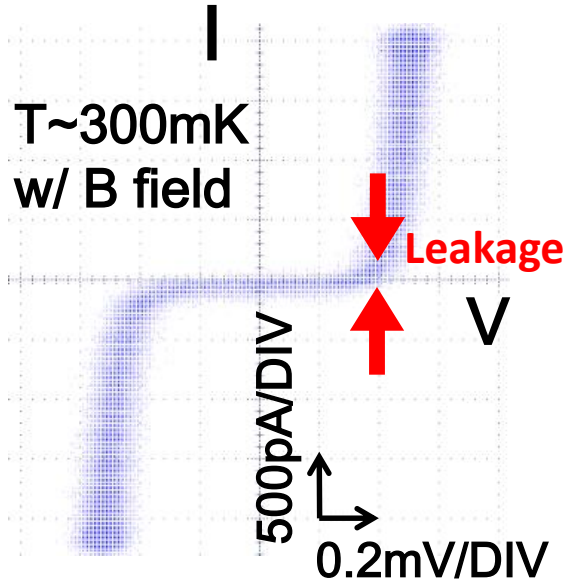
Hf-STJ

- Not established as a practical photo-detector yet by any group
- $N_{\text{q.p.}} = 25 \text{ meV} / 1.7\Delta \sim 735$
- 2% energy resolution for a 25 meV single-photon is achievable if Fano factor < 0.3 for Hf
- Spectrum measurement without a diffraction grating
 - Developing for a future satellite experiment

K.Takemasa “R&D Status of Hf-STJ” on 30th

Nb/Al-STJ development at CRAVITY

50 μ m sq. Nb/Al-STJ fabricated at CRAVITY



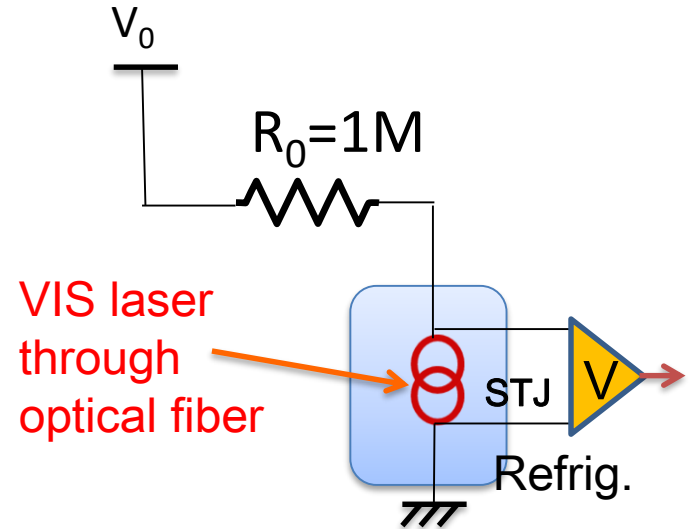
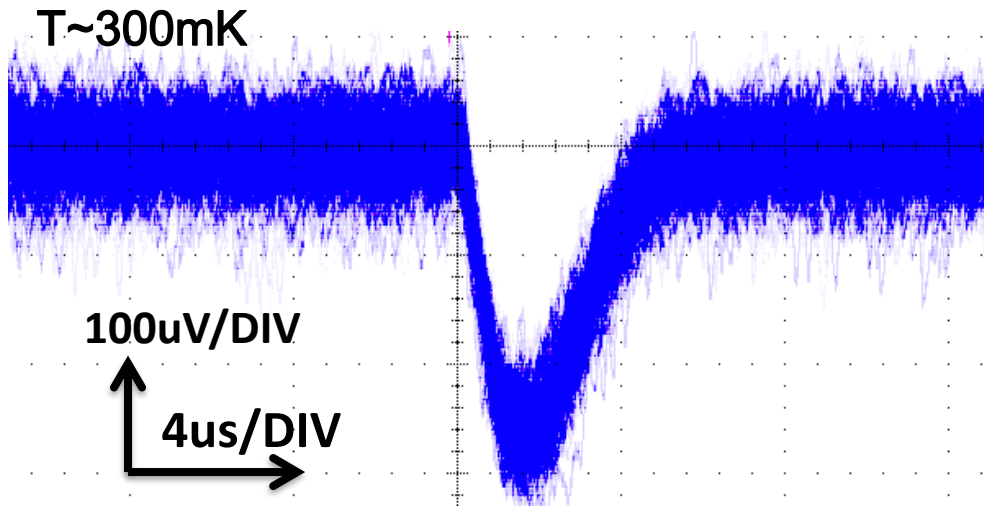
- $I_{leak} \sim 200 \text{ pA}$ for 50 μ m sq. STJ, and achieved 50 pA for 20 μ m sq.
- The shot noise from the leak current in case of $i_{leak} = 50 \text{ pA}$, $\Delta = 0.6 \text{ meV}$ and $G = 10$

$$NEP = \frac{1.7\Delta}{G} \sqrt{\frac{2i_L}{e}} \sim 4 \times 10^{-19} \text{ W}/\sqrt{\text{Hz}} \text{ (Without photon counting)}$$

- Photon counting will yield an improvement by **more than one order**.

Meets our requirements!

STJ response to pulsed laser

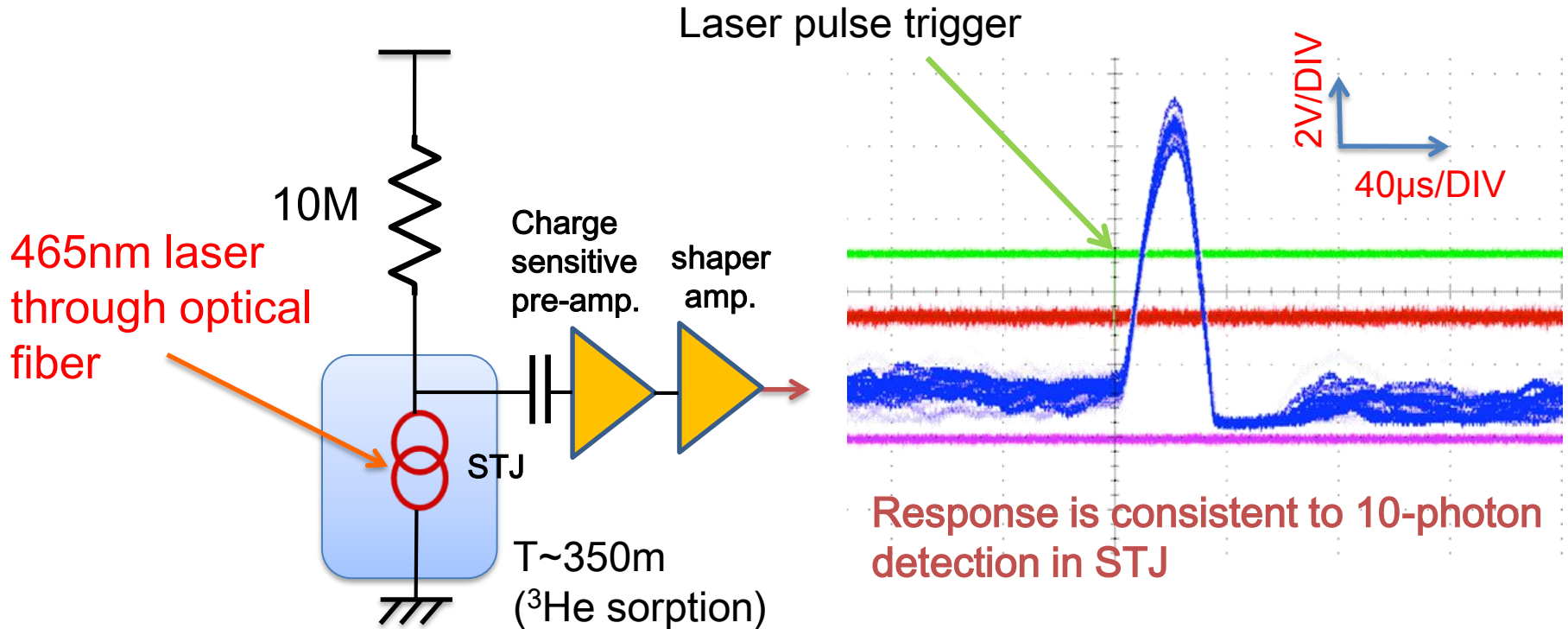


Nb/Al-STJ response to pulsed laser (465nm)
CRAVITY Nb/Al-STJ $100\mu\text{m}$ sq.

Nb/Al-STJ has $\sim 1\mu\text{s}$ response time.

- We can improve NEP by photon counting in $1\mu\text{s}$ integration time
- However we need faster readout system than $f > 1\text{MHz}$

100x100 μm^2 Nb/Al-STJ response to 465nm pulsed laser



We observed NIR-VIS laser pulse **at few-photon level** with a charge-sensitive amplifier placed at the room temperature.

Due to the readout noise, a FIR single-photon detection is not achieved yet.

Need ultra-low noise readout system for STJ signal

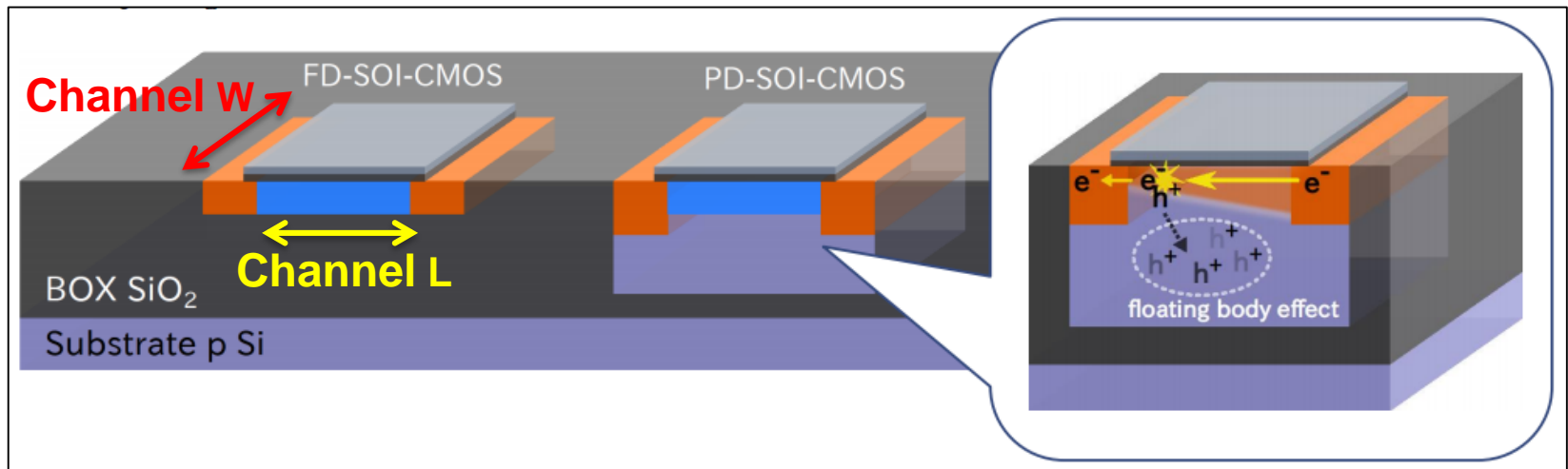
→ Considering a cryogenic pre-amplifier placed close to STJ

FD-SOI-MOSFET at Cryogenic temperature

FD-SOI : Fully Depleted – Silicon On Insulator

- ✓ Very thin channel layer in MOSFET
 - ◆ No floating body effect caused by charge accumulation in the body
- ✓ FD-SOI-MOSFET is reported to work at 4K

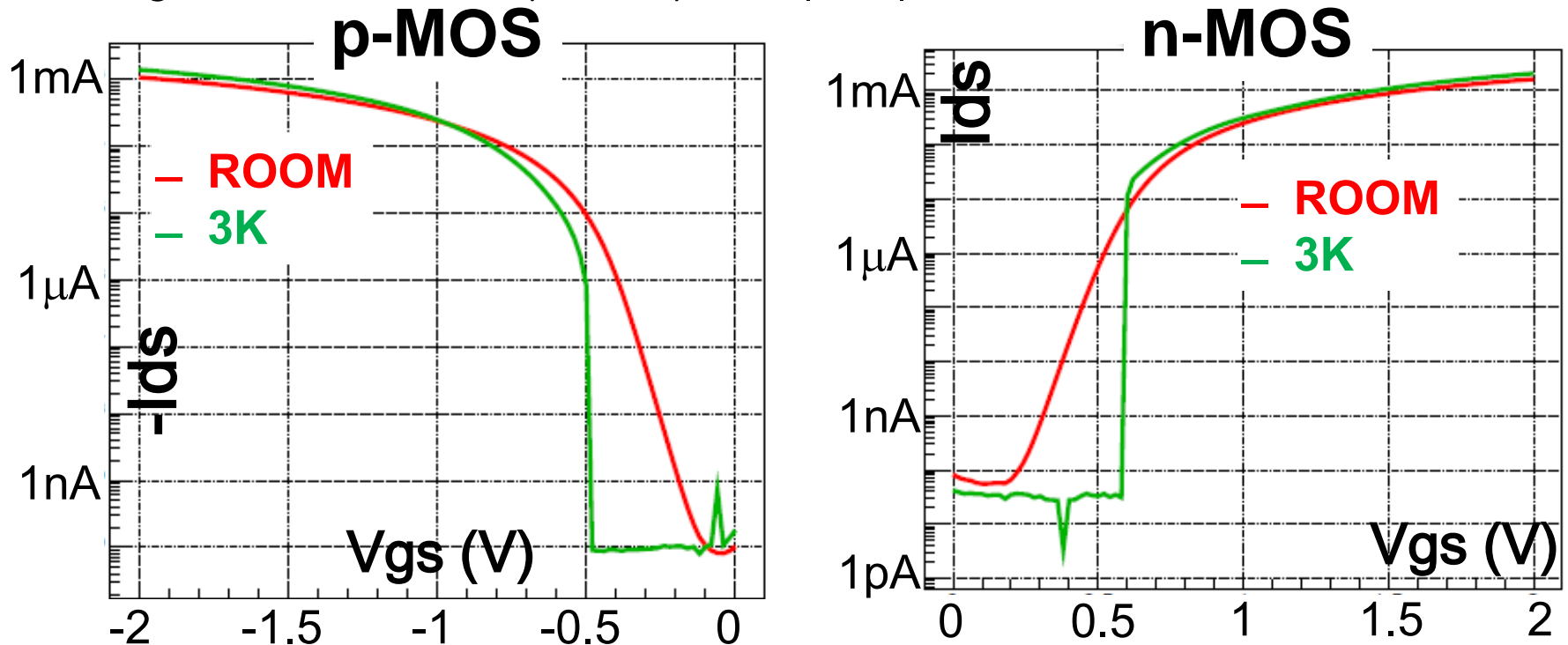
JAXA/ISIS AIPC 1185,286-289(2009)
J Low Temp Phys 167, 602 (2012)



S.Yagi “R&D Status of Nb/Al-STJ with SOI cryogenic preamplifier” on 30th
R.Wakasa “ニュートリノ崩壊光探索に向けた極低温増幅器の開発と現状” on Poster

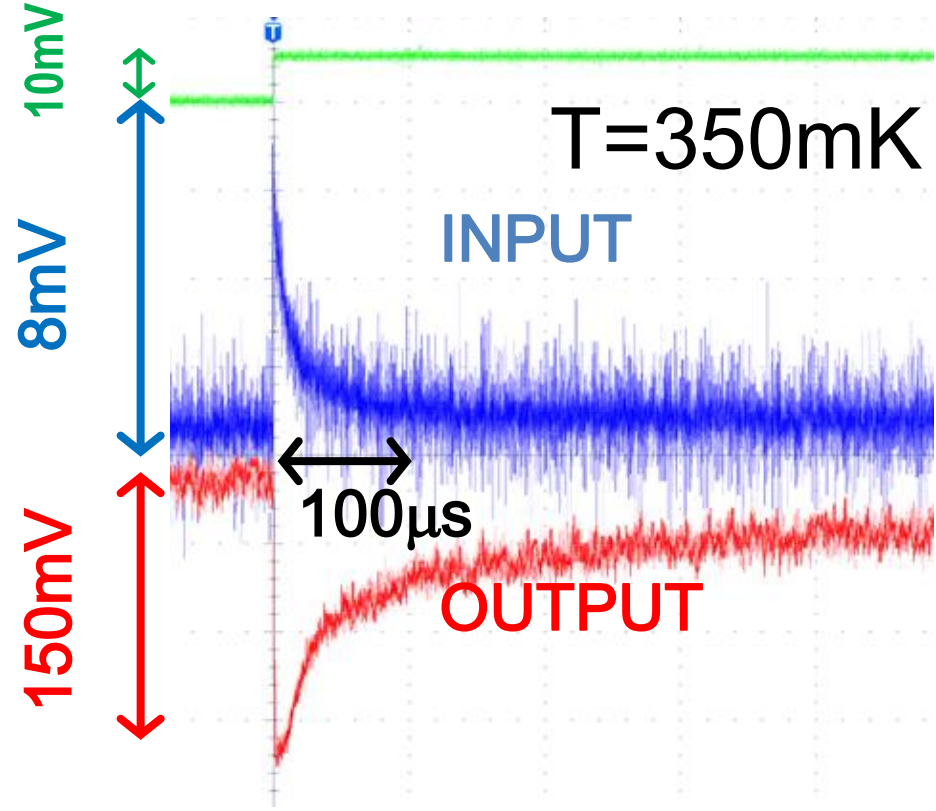
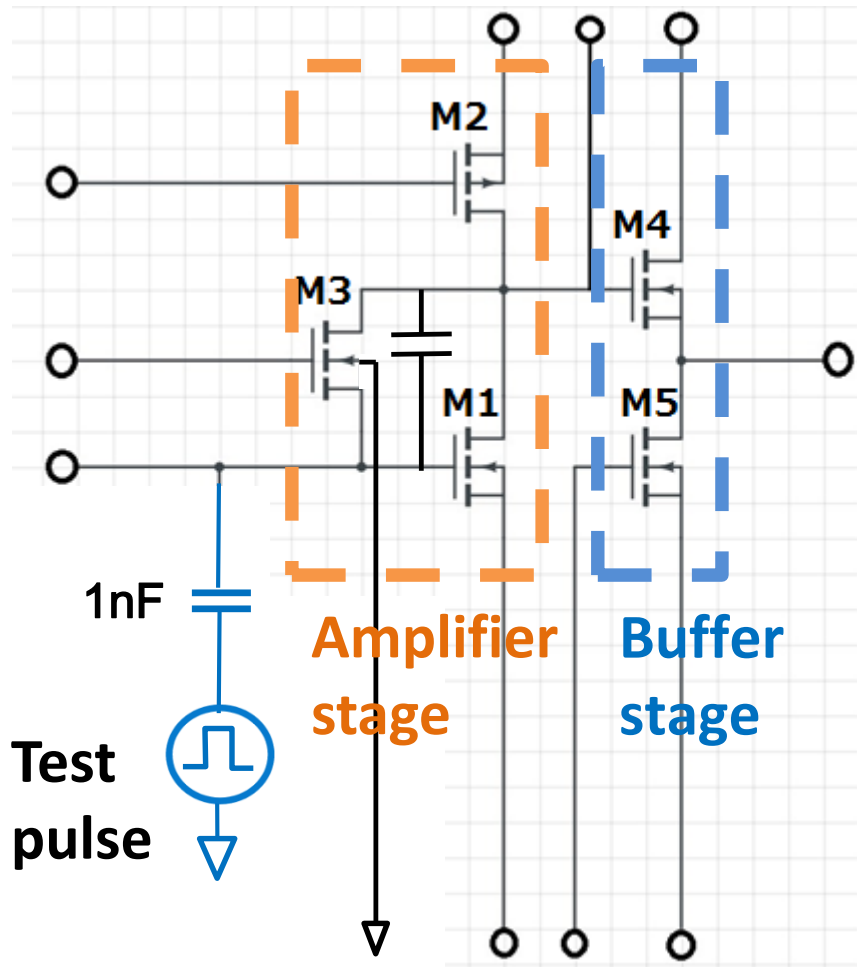
FD-SOI MOSFET I_d - V_g curve

- I_d - V_g curve of $W/L=10\mu\text{m}/0.4\mu\text{m}$ at $|V_{ds}|=1.8\text{V}$



- Both p-MOS and n-MOS show excellent performance at 3K (We confirmed they function down to 300mK).
- Threshold shifts, sub-threshold current suppression and increase of the carrier mobility at low temperature.

SOI prototype amplifier



Test pulse input through $C=1\text{nF}$ capacitance at $T=3\text{K}$ and 350mK

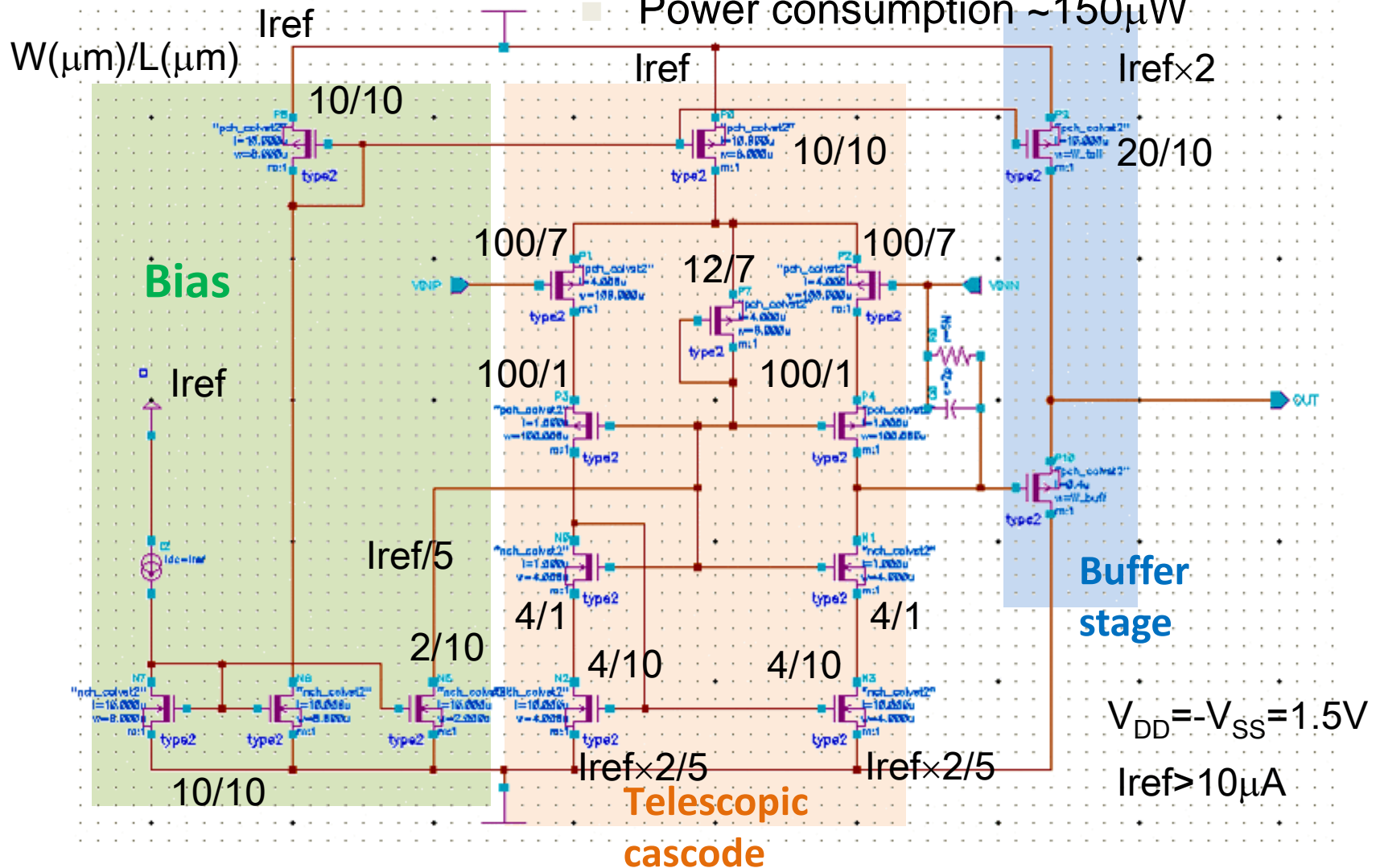
- Power consumption $\sim 100\mu\text{W}$

We can compensate the effect of shifts in the thresholds by adjusting bias voltages.

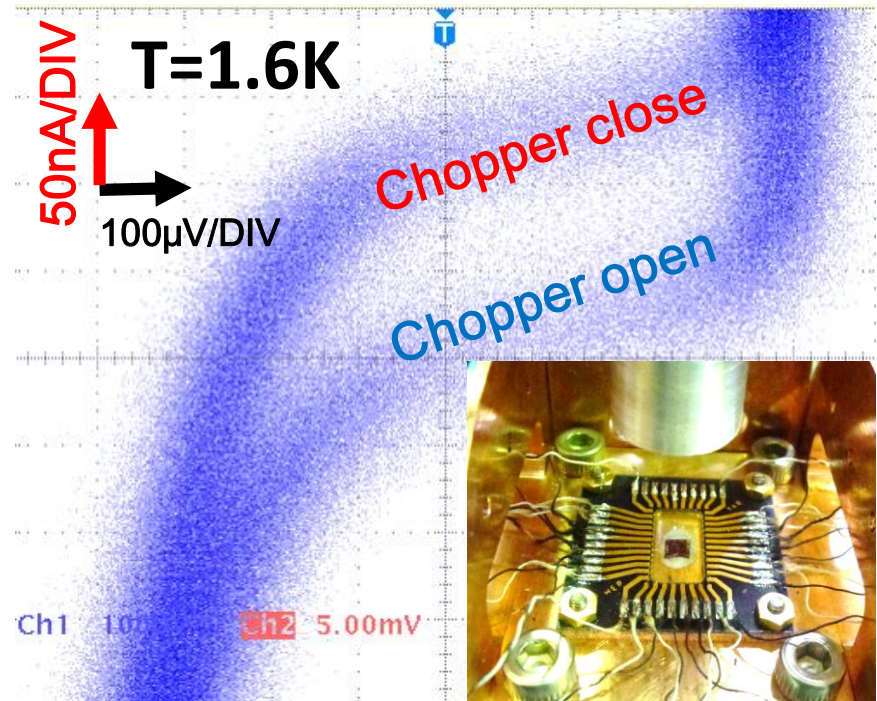
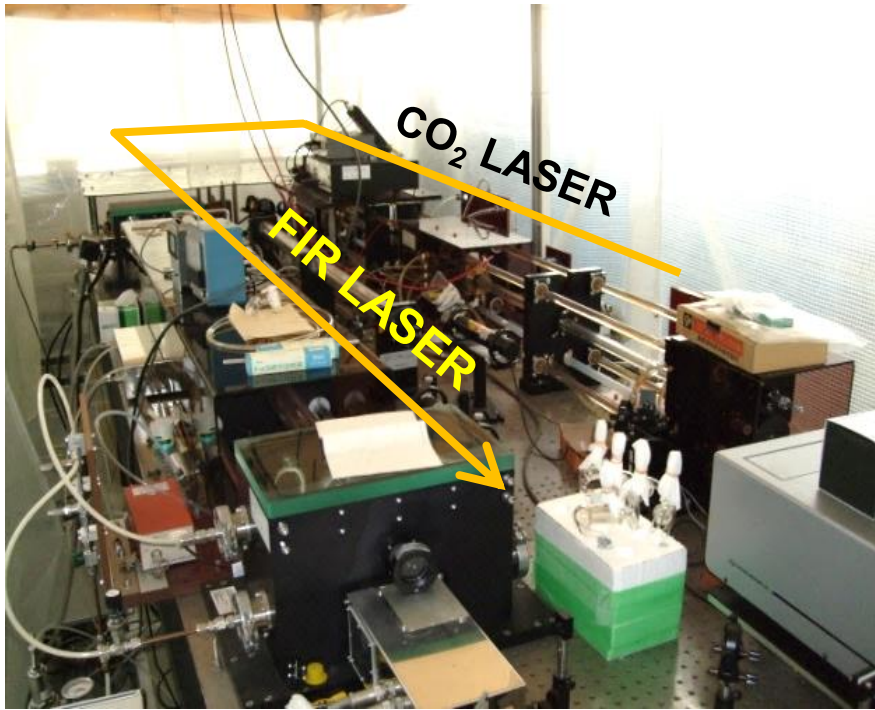
Op-amp Circuit for STJ design

Arriving in this winter!

- telescopic cascode differential amplifier
- Feedback $C=2\text{pF} \times R=5\text{M}\Omega = 10\mu\text{s}$
- Power consumption $\sim 150\mu\text{W}$



Calibration of STJ by Far-infrared Laser



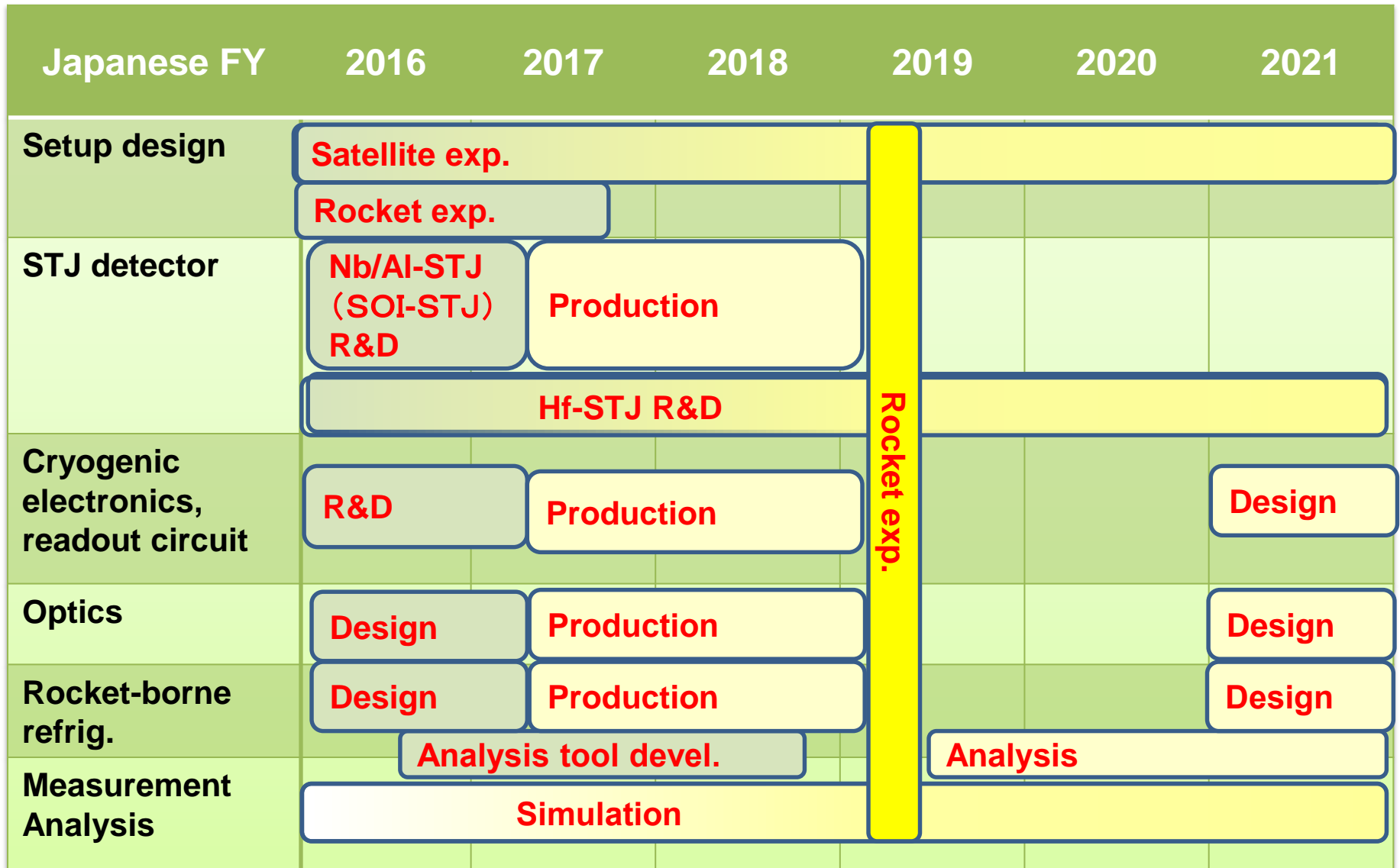
A Nb/Al-STJ is illuminated by FIR laser through a chopper ($f=40\sim 200\text{Hz}$) using a far-infrared molecular laser apparatus at FIR-UF (U. of Fukui)

200 μm sq. Nb/Al-STJ by CRAVITY

- Observed a signal current of $\sim 100\text{nA}$ in response to a $57.2\mu\text{m}$ laser
- FIR source for the STJ calibration is going ready!

C.Asano “COBAND実験におけるSTJ検出器較正用遠赤外光源開発” on Poster

COBAND project



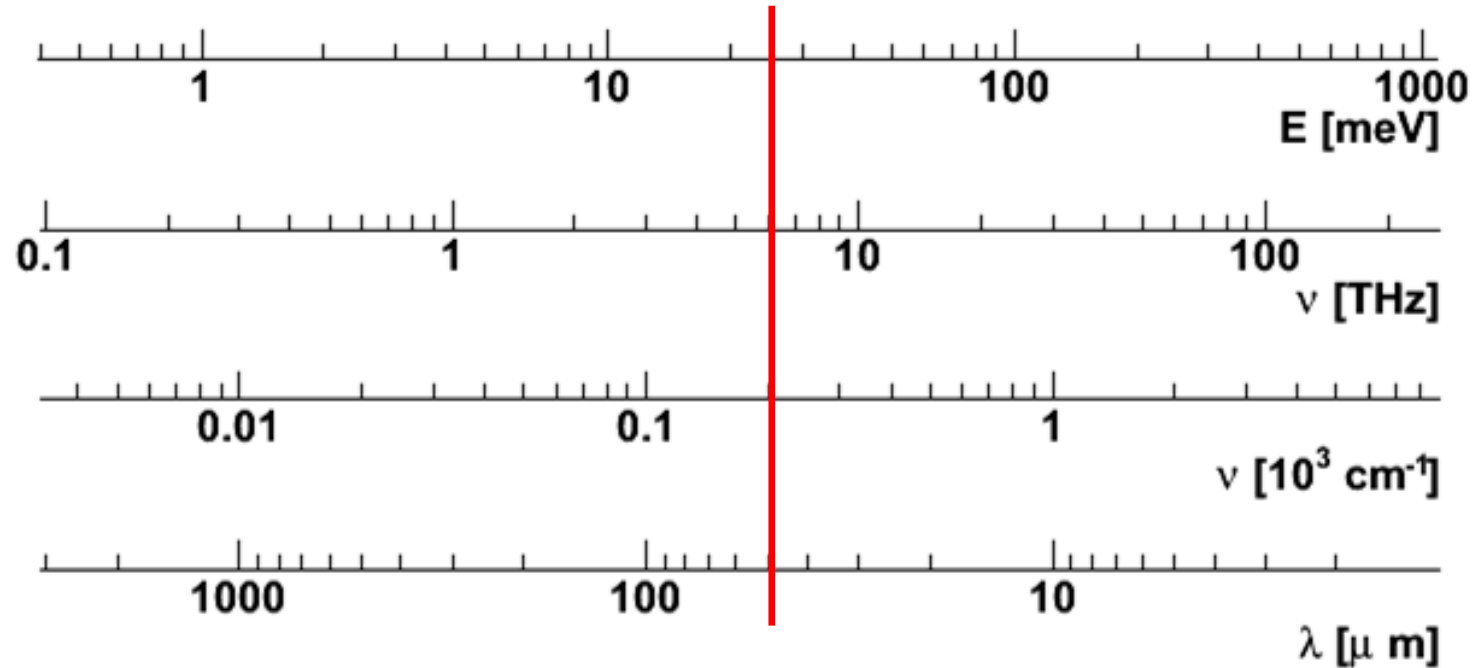
Summary

- We propose a sounding rocket experiment to search for neutrino radiative decay in cosmic neutrino background, and a following satellite experiment .
- Requirements for the detector is a photo-detector with $NEP \sim O(10^{-20})$ W/\sqrt{Hz} .
- Nb/Al-STJ array with a diffractive for the rocket experiment.
 - Nb/Al-STJ fabricated at CRAVITY meets our requirements.
 - FD-SOI readout is under development and almost ready for STJ signal amplification at cryogenic temperature.
 - FIR calibration for STJ is going to be ready.
- Improvement of the neutrino lifetime lower limit up to $O(10^{14} \text{ yrs})$ is feasible for 200-sec measurement in a rocket-borne experiment with the detector.

- **K.Takemasa “R&D Status of Hf-STJ” on 30th**
- **S.Yagi “R&D Status of Nb/Al-STJ with SOI cryogenic preamplifier” on 30th**
- **R.Wakasa “ニュートリノ崩壊光探索に向けた極低温増幅器の開発と現状” on POSTER**
- **C.Asano “COBAND実験におけるSTJ検出器較正用遠赤外光源開発” on POSTER**

Backup

Energy/Wavelength/Frequency



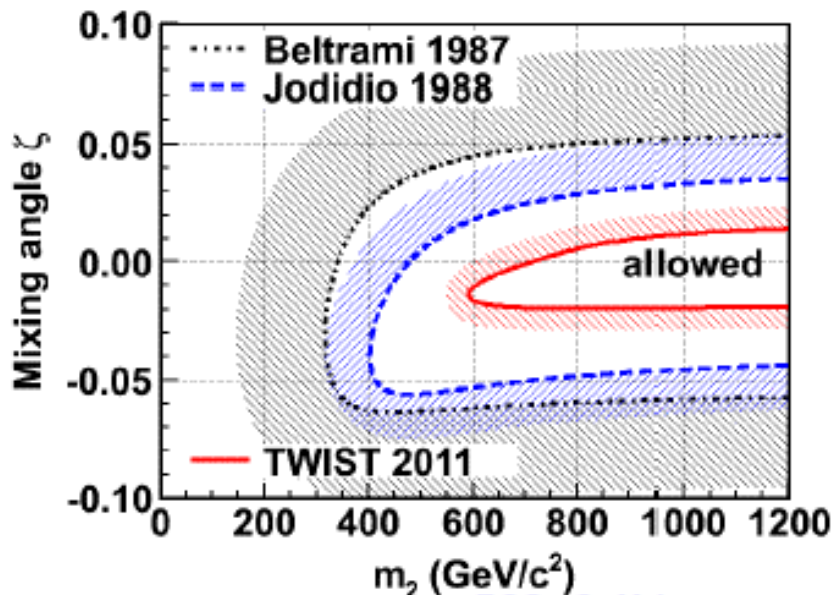
$$E_{\gamma} = 25 \text{ meV}$$

$$\nu = 6 \text{ THz}$$

$$\lambda = 50 \mu\text{m}$$

Limit on LRSM parameters from TWIST

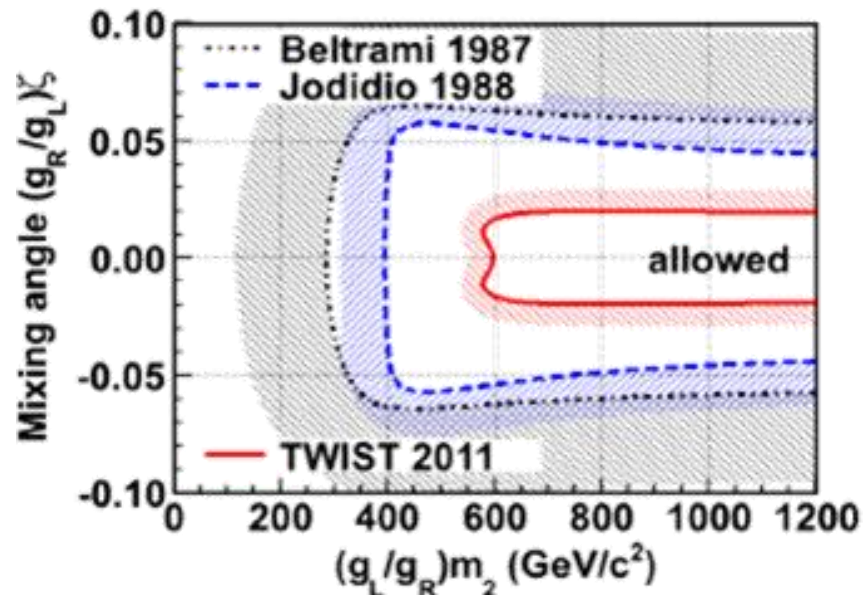
Manifest LRS 90%CL



$$M_{W_2} > 592 \text{ GeV}$$

$$-0.020 < \zeta < +0.017$$

Non-manifest LRS 90%CL



$$(g_L/g_R)M_{W_2} > 578 \text{ GeV}$$

$$-0.020 < (g_L/g_R)\zeta < +0.020$$

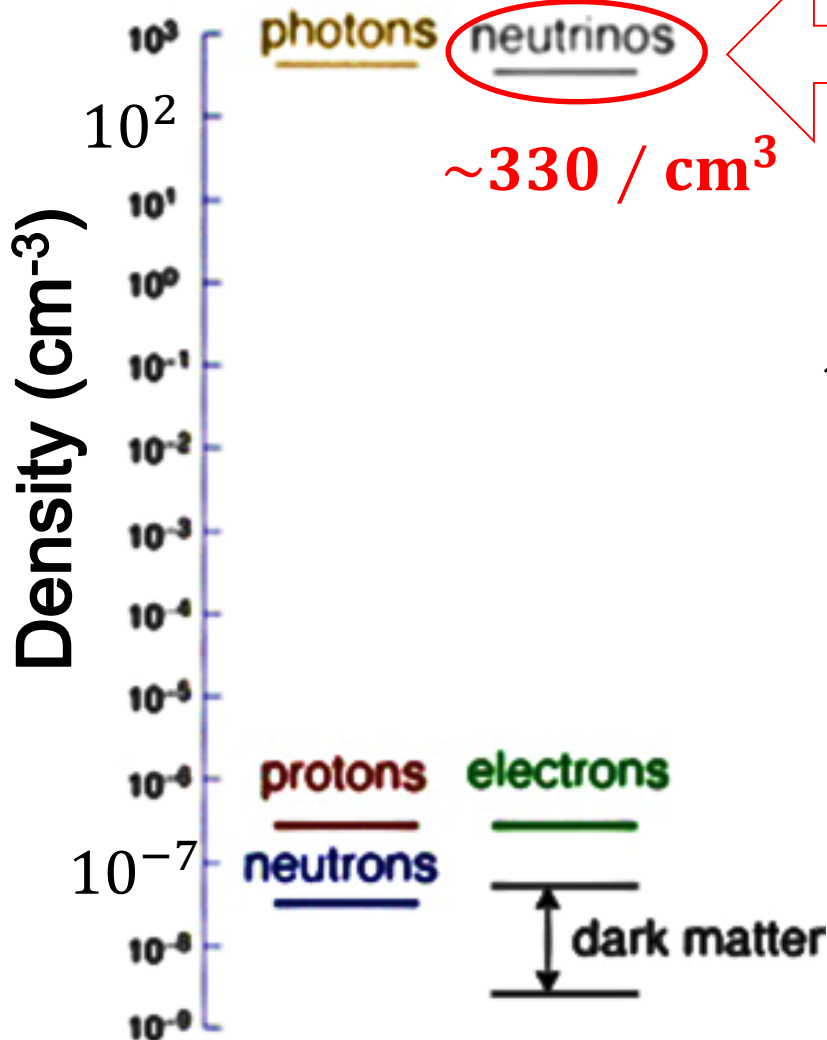
□ Other **W'** direct search mass limits

- ATLAS: $>1.49 \text{ TeV}/c^2$, 95%CL (LLWI11)
- CMS: $>1.58 \text{ TeV}/c^2$, 95%CL (LLWI11)
- CMS: $>1.36 \text{ TeV}/c^2$, 95%CL (2011)
- CDF: $>1.12 \text{ TeV}/c^2$, 95%CL (2011)
- D0: $>1.0 \text{ TeV}/c^2$, 95%CL (2008)

□ Other limits on mixing angle ζ

- Hardy and Towner: <0.0005 (MLRS), <0.04 (generalized)
- K decay: <0.004 (MLRS)

Cosmic background neutrino (CνB)



The universe is filled with neutrinos. However, they have not been detected yet!

$$n_\nu + n_{\bar{\nu}} = \frac{3}{4} \left(\frac{T_\nu}{T_\gamma} \right)^3 n_\gamma$$

$$= 110 / \text{cm}^3 / \text{generation}$$

$$T_\nu = \left(\frac{4}{11} \right)^{\frac{1}{3}} T_\gamma = 1.95 \text{K}$$

$$\langle p_\nu \rangle = 0.5 \text{meV}/c$$

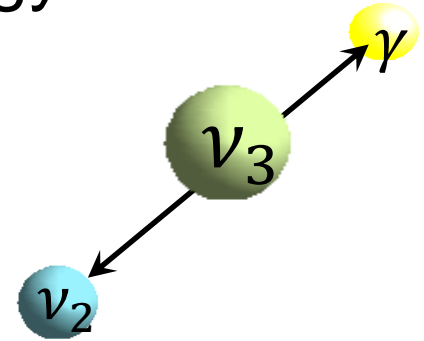
Neutrino Mass and Photon Energy

- From neutrino oscillation
 - $|\Delta m_{23}^2| = |m_3^2 - m_2^2| = 2.4 \times 10^{-3} \text{ eV}^2$
 - $\Delta m_{12}^2 = m_3^2 - m_2^2 = 7.65 \times 10^{-5} \text{ eV}^2$
- CMB (Planck+WP+highL) and BAO
 - $\sum m_i < 0.23 \text{ eV}$

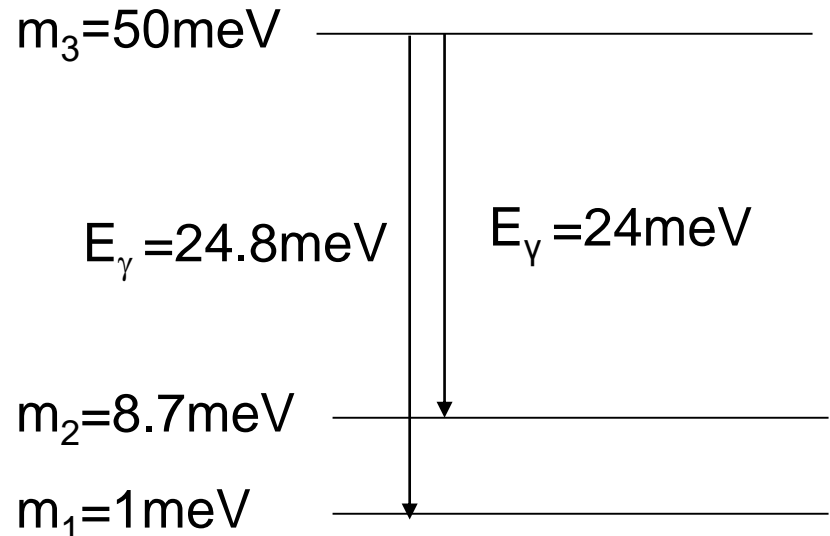
$\rightarrow 50\text{meV} < m_3 < 87\text{meV}$

$E_\gamma = 14 \sim 24\text{meV}$

$\lambda_\gamma = 51 \sim 89\mu\text{m}$



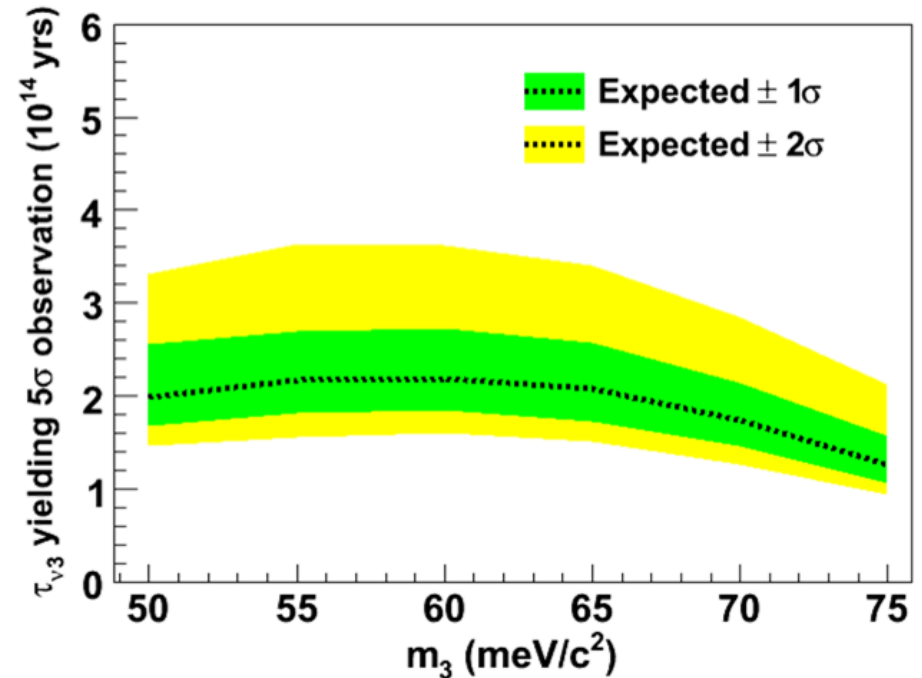
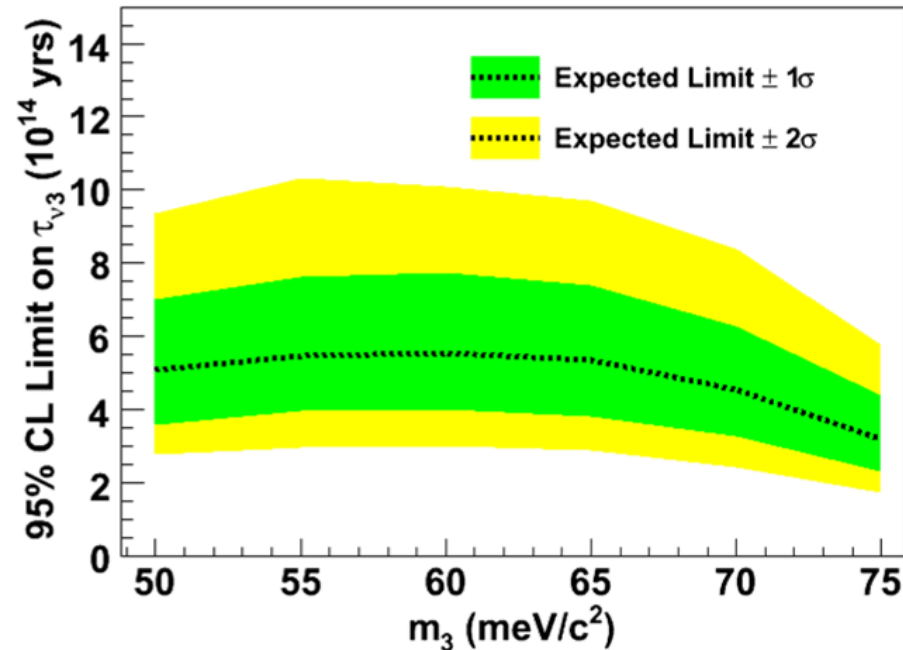
$$E_\gamma = \frac{m_3^2 - m_2^2}{2m_3}$$



Sensitivity to neutrino decay

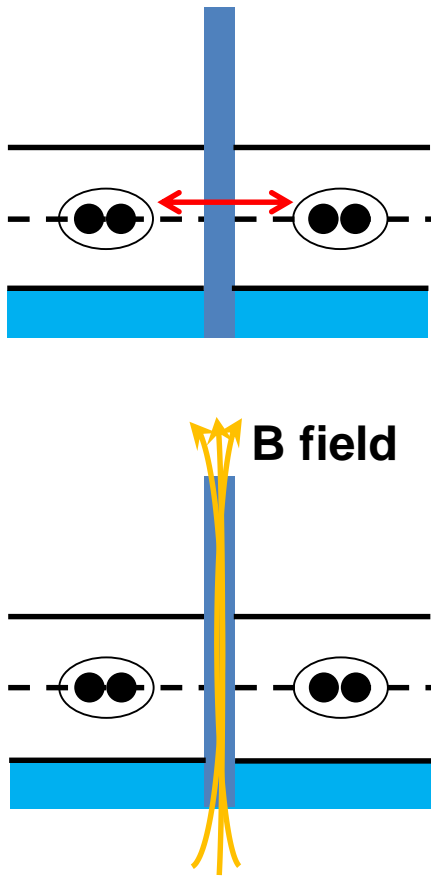
Parameters in the rocket experiment simulation

- telescope dia.: 15cm
- 50-column (λ : 40 μm – 80 μm) \times 8-row array
- Viewing angle per single pixel: 100 μrad \times 100 μrad
- Measurement time: 200 sec.
- Photon detection efficiency: 100%

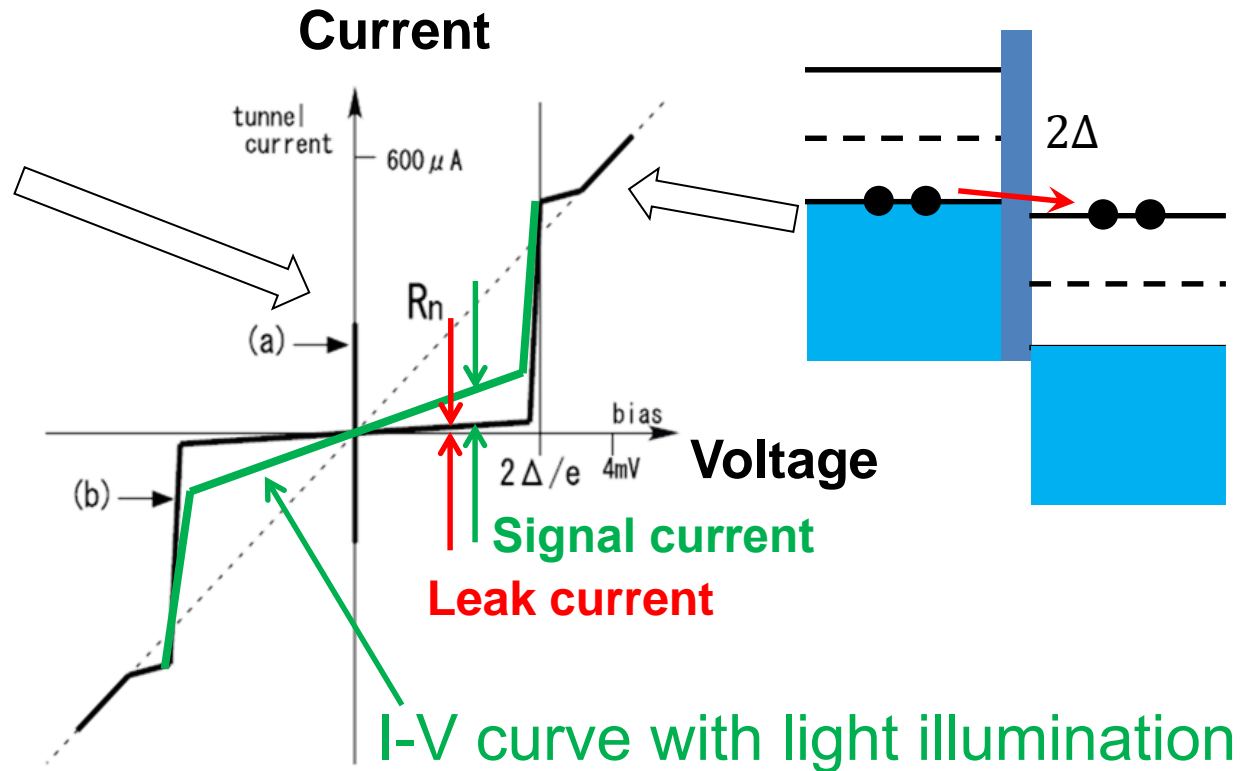


- Can set lower limit on ν_3 lifetime at $4\text{--}6 \times 10^{14}$ yrs if no neutrino decay
- If ν_3 lifetime were 2×10^{14} yrs, the signal significance would be at 5 σ level

STJ current-voltage curve



Tunnel current of Cooper pairs (Josephson current) is suppressed by applying magnetic field

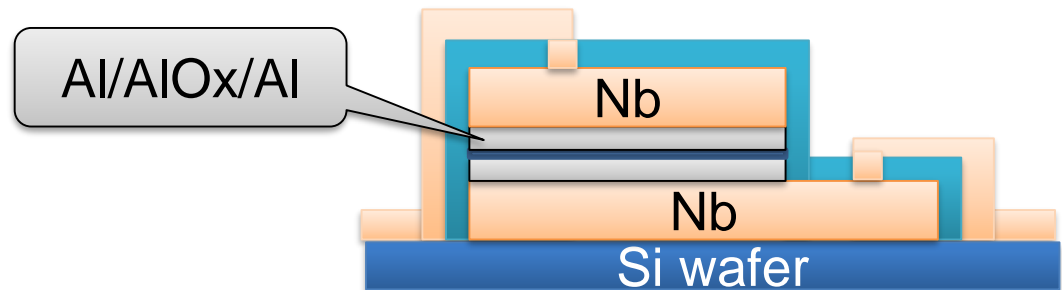


Optical signal readout

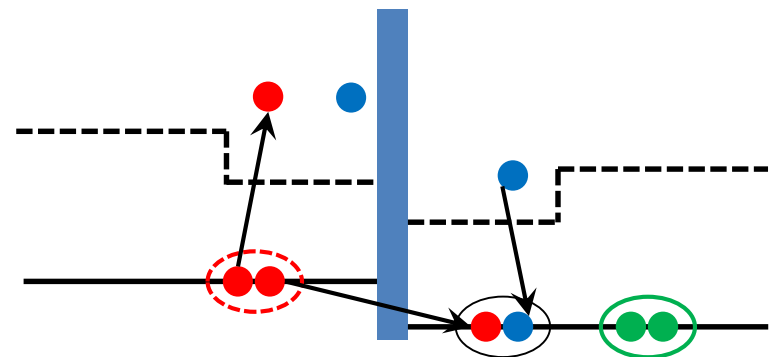
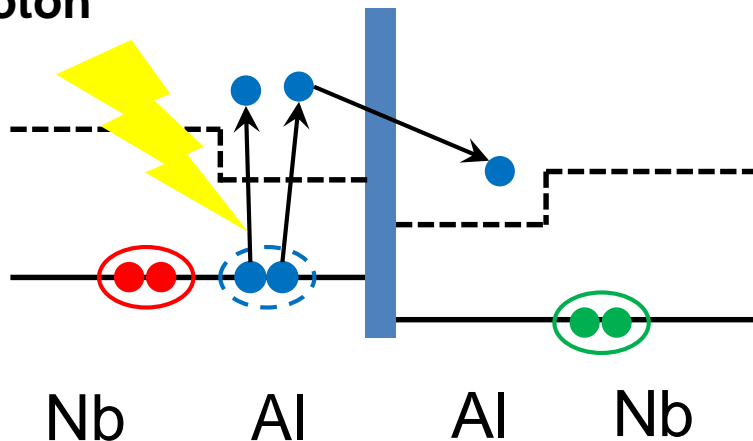
- Apply a constant bias voltage ($|V| < 2\Delta$) across the junction and collect tunneling current of quasi particles created by photons
- ✓ Leak current causes background noise

STJ back-tunneling effect

- Bi-layer fabricated with superconductors of different gaps $\Delta_{\text{Nb}} > \Delta_{\text{Al}}$ to enhance quasi-particle density near the barrier
 - Quasi-particle near the barrier can mediate **multiple Cooper pairs**
- Nb/Al-STJ Nb(200nm)/Al(70nm)/AlOx/Al(70nm)/Nb(200nm)
- Gain: ~ 10



Photon



Detector NEP required for COBAND rocket experiment

What's NEP (Noise Equivalent Power)?

- The signal power yielding unit signal-to-noise ratio in unit bandwidth for a photo-detector

$$NEP = \frac{F}{S/N\sqrt{\Delta f}}$$

F: Radiant flux(W)

S/N: signal-to-noise ratio of the detector output

Δf : detector bandwidth(Hz)

- **A smaller NEP means a more sensitive detector.**
- A detector with $NEP = 10^{-12} W / \sqrt{Hz}$ can detect
 - 1pW signal with S/N=1 after 0.5-sec integration.
 - 0.1pW signal with S/N=1 after 50-sec integration.

Detector NEP required for COBAND rocket experiment

Viewing angle of the photo-detector

- Telescope main mirror: $D=15\text{cm}$, $F=1\text{m}$

➤ $100\mu\text{m} \times 100\mu\text{m} \times 8$ pixels

➔ Viewing angle : $8 \times (100\mu\text{m}/1\text{m})^2 = 8 \times 10^{-8} \text{ sr}$

$\Delta\nu$ per one wavelength division after diffraction grating:

➤ One wavelength division: $(80\mu\text{m} - 40\mu\text{m}) / 50 = 0.8\mu\text{m}$

➤ $\Delta\nu = c/50\mu\text{m} - c/50.8\mu\text{m} = 94\text{GHz}$ @ $\lambda = 50\mu\text{m}$

Hereafter on the assumption of **100% quantum efficiency** at the photo-detector

Detector NEP required for COBAND rocket experiment

- Neutrino decay ($m_3 = 50 \text{ meV}$, $\tau_\nu = 10^{14} \text{ yrs}$): $I_\nu = 25 \text{ kJy/sr @ } \lambda = 50 \mu\text{m}$

$$F_{ND} = 25 \text{ kJy/sr} \times 8 \times 10^{-8} \text{ sr} \times \pi(15 \text{ cm}/2)^2 \times 94 \text{ GHz}$$

$$= 3.3 \times 10^{-20} \text{ W}/8 \text{ pix}$$

- Zodiacal emission: $I_\nu = 8 \text{ MJy/sr @ } \lambda = 50 \mu\text{m}$

$$F_{ZE} = 1.1 \times 10^{-17} \text{ W}/8 \text{ pix}$$

- ◆ The fluctuation in the F_{ZE} integration over Δt interval

- The fluctuation in # of photons: $\epsilon_\gamma \sqrt{F_{ZE} \Delta t / \epsilon_\gamma} = \sqrt{\epsilon_\gamma F_{ZE} \Delta t}$

- ◆ Conditions of the requirements on Δt and NEP

$$\frac{NEP}{\sqrt{2\Delta t}} \times \Delta t \ll \sqrt{\epsilon_\gamma F_{ZE} \Delta t} \ll F_{ND} \Delta t$$

→ $\Delta t > 40 \text{ sec}$ (1σ) $\Delta t > 200 \text{ sec}$ (2.2σ per $\Delta\lambda = 0.8 \mu\text{m}$)

→ $NEP \ll 3 \times 10^{-19} \text{ W}/\sqrt{\text{Hz}}$ for 8 pix

→ $NEP \ll 8.4 \times 10^{-19} \text{ W}/\sqrt{\text{Hz}}$ for 1pix

Shot noise from a detector leak current

- R [A/W] : Detector responsivity
- i_L [A] : leak current

The charge from the leakage after T [sec] integration: $i_L T$ [C]

The number of charge carriers after averaging: $i_L T/e$

The fluctuation in the current: $\frac{e}{T} \sqrt{\frac{i_L T}{e}}$

The fluctuation in the measured incident power: $\frac{1}{R} \sqrt{\frac{e i_L}{T}}$

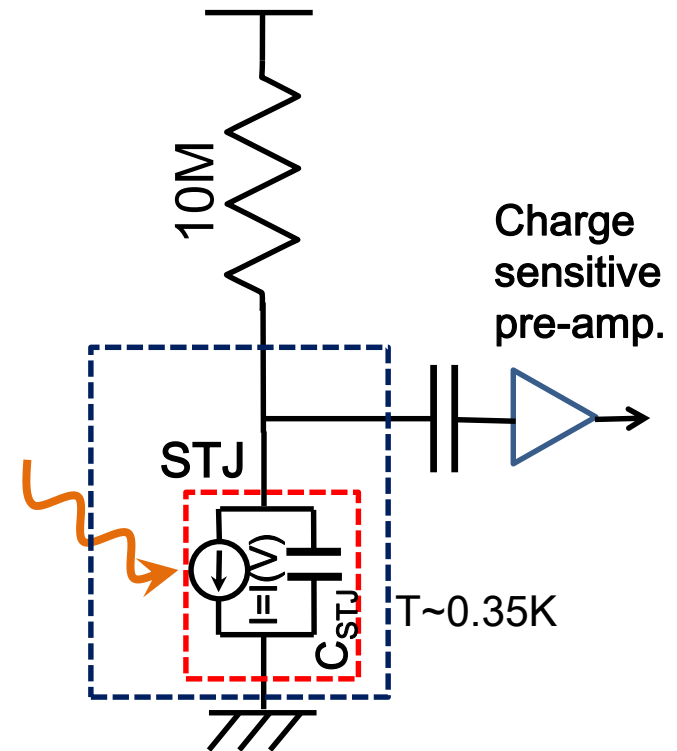
$$\frac{NEP}{\sqrt{2T}} = \frac{1}{R} \sqrt{\frac{e i_L}{T}} \quad \text{i.e.} \quad NEP = \frac{1}{R} \sqrt{2e i_L}$$

$$\text{For STJ: } N_{q.p.} = G \frac{E_\gamma}{1.7\Delta} \quad \text{i.e.} \quad R = G \frac{e}{1.7\Delta}$$

$$NEP = \frac{1.7\Delta}{G} \sqrt{\frac{2i_L}{e}}$$

SOI charge-sensitive pre-amplifier development

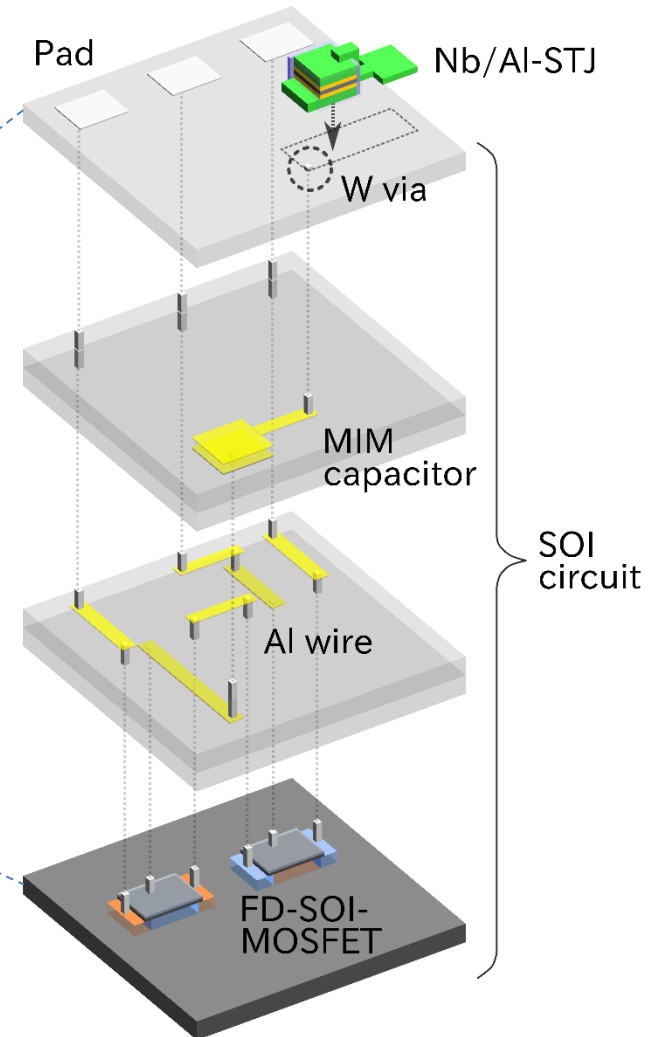
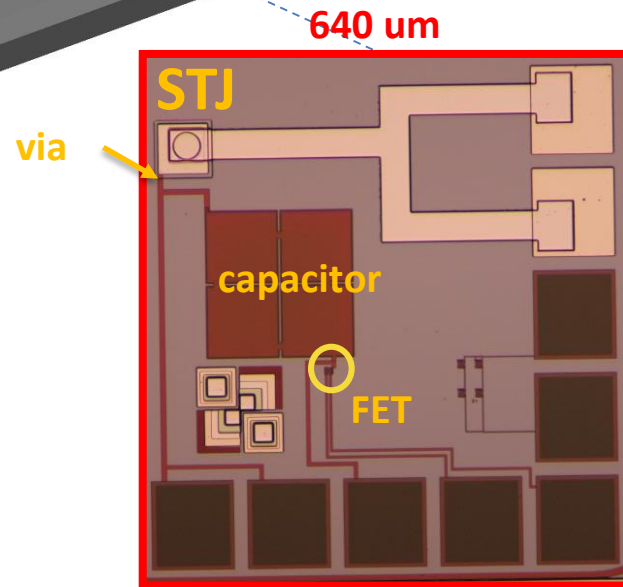
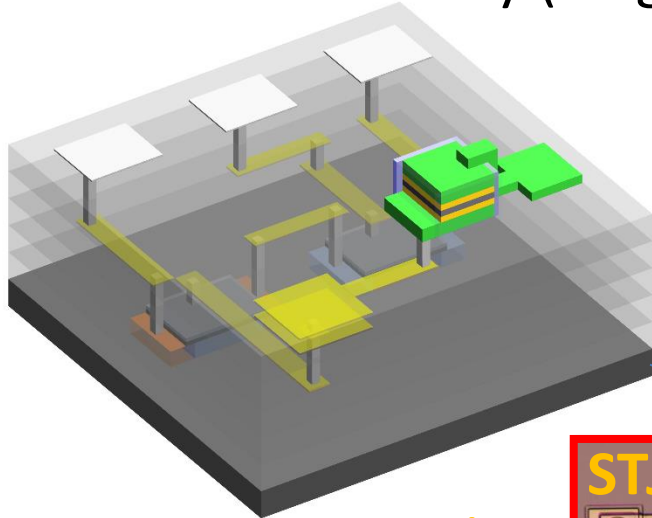
- STJ has comparably large capacitance: $>20\text{pF}$ for $20\mu\text{m}$ sq. STJ.
 - A low input impedance charge-sensitive amplifier is required for STJ single-photon signal readout.
- STJ response time is $\sim 1\mu\text{s}$.
- We designed SOI op-amp which has $>1\text{MHz}$ freq. response, and submitted to the next SOI MPW run. We'll test the amplifier in this winter.



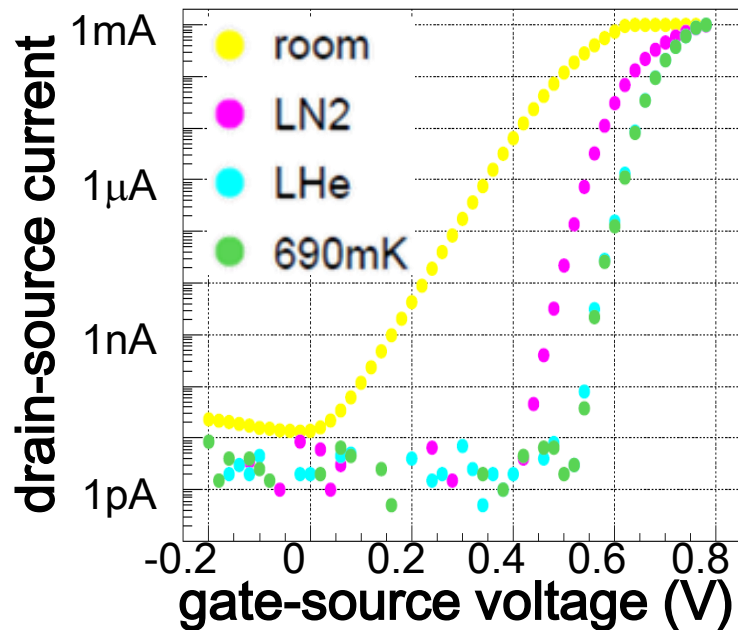
SOI-STJ (STJ directly on SOI) development

- STJ layers are fabricated **directly on** a SOI pre-amplifier board and cooled down together with the STJ

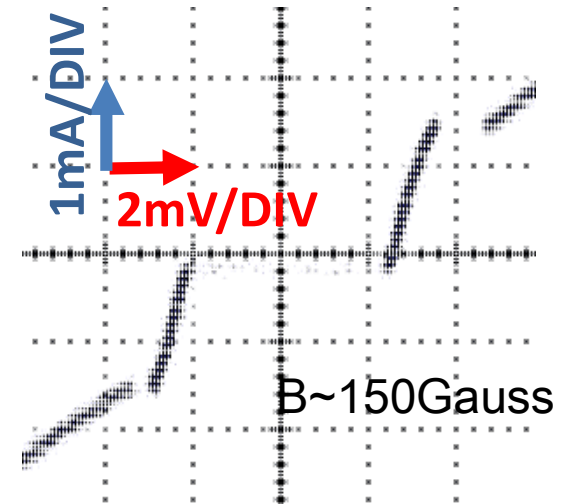
□ Potential to STJ array (Large area STJ)



FD-SOI on which STJ is fabricated



nMOS-FET in FD-SOI wafer on which a STJ is fabricated at KEK



I-V curve of a STJ fabricated at KEK on a FD-SOI wafer

- Both nMOS and pMOS-FET in FD-SOI wafer on which a STJ is fabricated work fine at temperature down below 1K
- Nb/Al-STJ fabricated at KEK on FD-SOI works fine
- We are also developing SOI-STJ where STJ is fabricated at CRAVITY