



Logo of German KET Workshop, May 2016

Precision physics at energy frontier e+e- colliders



Precision theory
for precise measurements at
the LHC and future colliders

ICISE
Quy-Nhon, Vietnam
September 25 - October 1, 2016

Precision challenges in theory for data analyses
at the LHC and at Future Colliders

- Higgs boson production and properties
- Electroweak measurements
- Top quark physics
- Heavy flavour physics
- PDFs

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<http://rencontresduvietnam.org/conferences/2016/precision-theory/>

Roman Pöschl
Directeur de Recherche of CNRS



R.P. is indebted to A. Freitas, S. Heinemeyer, F. Richard and F. LeDiberder for useful discussions and to many authors from whom I have reused their material

- Chapter 1: Introduction
- Chapter 2: Electroweak precision tests
- Chapter 3: Top physics
- Chapter 4: BSM physics

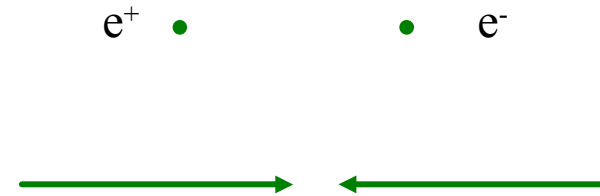
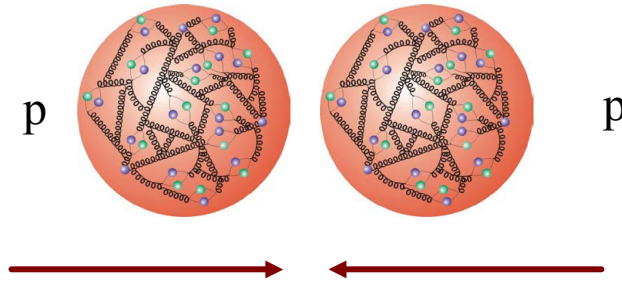
1. Introduction



- 1) Collisions at energies well above the electroweak scale
 - Requires now and in the foreseeable future Hadron colliders
 - Direct production of new particles
 - Produce large number of rare particles and study rare decays
 - First precision measurements of key particles of electroweak theory-> High energy, High luminosity LHC

- 2) **e+e-Collisions at energies at the electroweak scale**
 - Probe the electroweak scale with high precision
 - ... in particular particles that carry the “imprint of the Higgs Field such as W, Z and top”

- 3) e+e- collisions at 'smaller' energies
 - Requires high luminosity to get sensitive to tiny quantum effects-> SuperKEKB

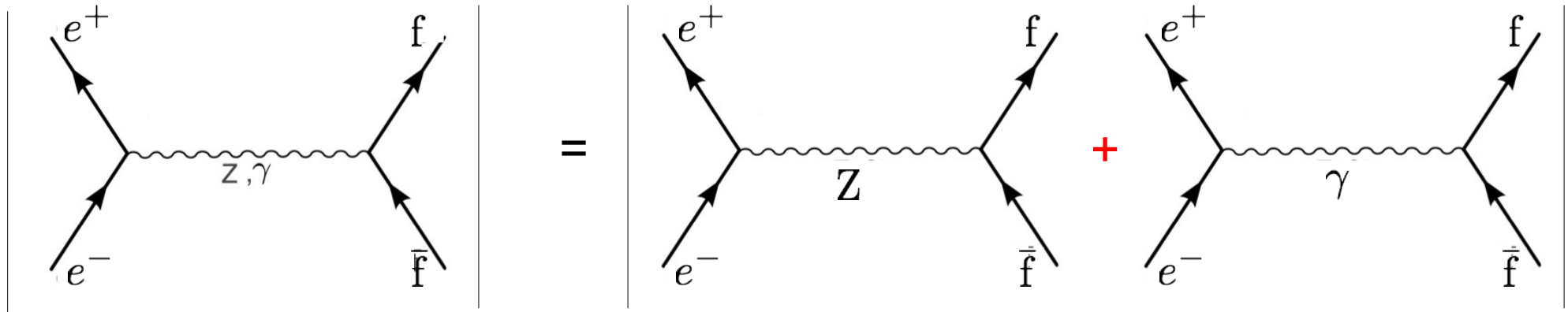


Proton:

- Composed particle (hadron)
- Unknown energy of collision partners
- Parasitic reactions
- Strong interaction
=> Considerable physics background
- Advantage: Scan of energy
Range within one experiment

Electron:

- Elementary particle
- Well known and adjustable
energy of collision partners
- High precision measurements
in a model independent way
with small electroweak backgrounds
- Each energy point needs a
new set of machine parameters



Interference between individual amplitudes of γ and Z exchange

$$\mathcal{M}_Z = -\frac{\sqrt{2}G_F M_Z^2}{s - M_Z^2} \left[\bar{f} \gamma^\rho \left(c_V^f - c_A^f \gamma^5 \right) f \right] g_{\rho\sigma} \left[\bar{e} \gamma^\sigma \left(c_V^e - c_A^e \gamma^5 \right) e \right]$$

$$\mathcal{M}_\gamma = -\frac{e^2}{s} (\bar{f} \gamma^\nu f) g_{\mu\nu} (\bar{e} \gamma^\nu e)$$

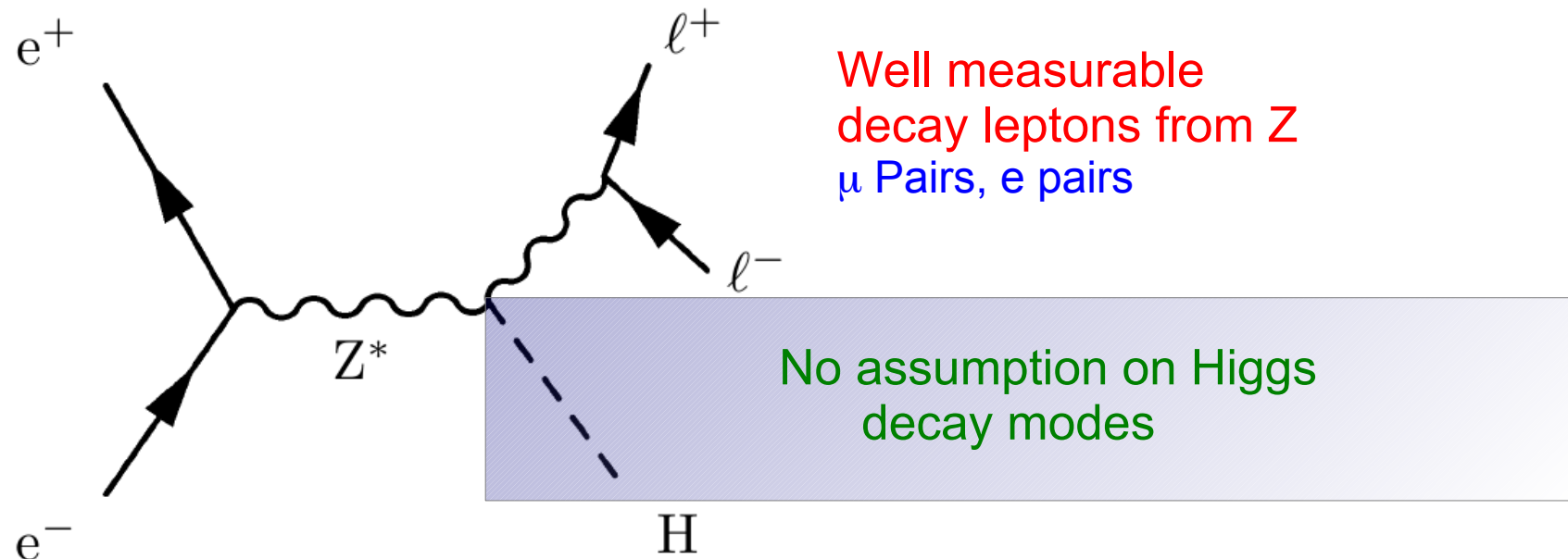
Differential cross section:

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} \left[A_0(1 + \cos^2\theta) + A_1 \cos\theta \right] \left\{ \begin{array}{ll} \sim (1 + \cos^2\theta) & \text{'Usual' Vector current, symmetric in } \cos\theta \\ \sim \cos\theta & \text{Axial Vector current, asymmetric in } \cos\theta \end{array} \right.$$

Weak interaction introduces forward backward asymmetry

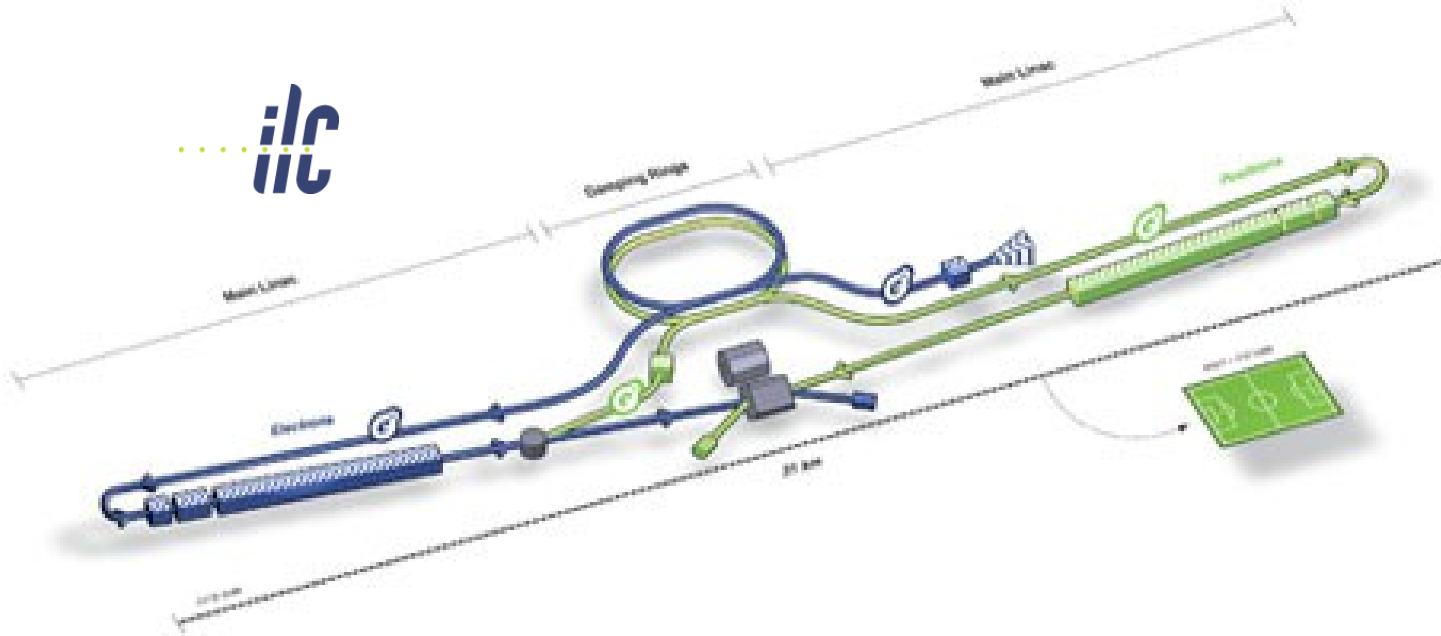
=> Asymmetry is intrinsic to electroweak processes!!!

Higgs Mass and ZZH coupling by **Model Independent** measurement

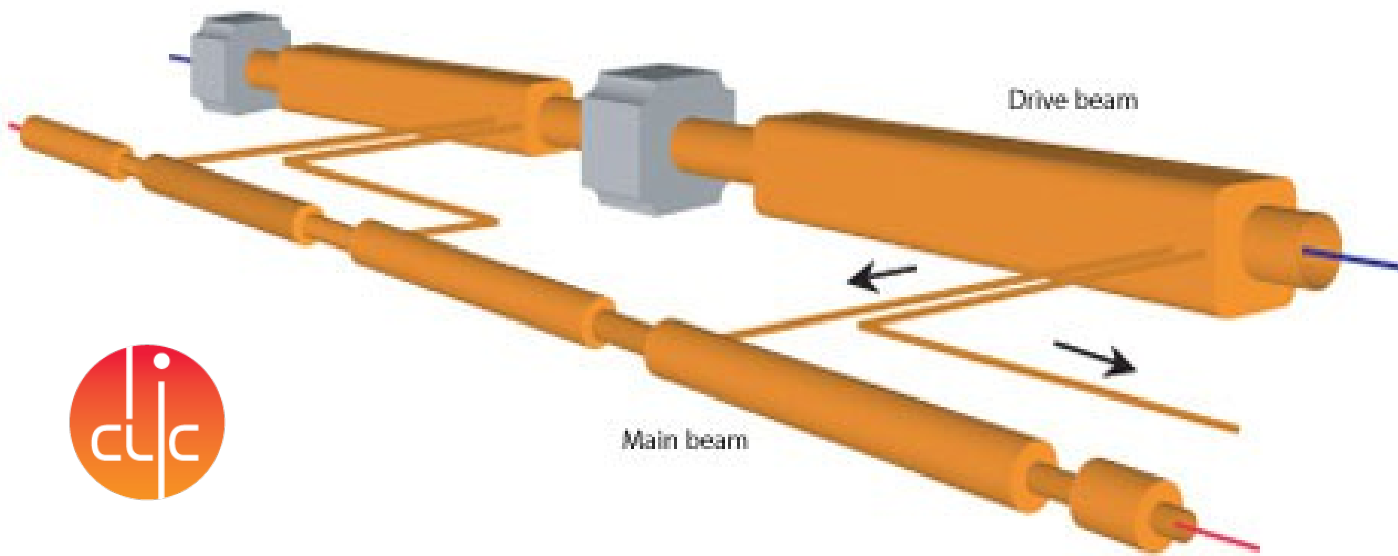


Higgs Recoil Mass:
$$M_h^2 = M_{recoil}^2 = s + M_Z^2 - 2 E_Z \sqrt{s}$$

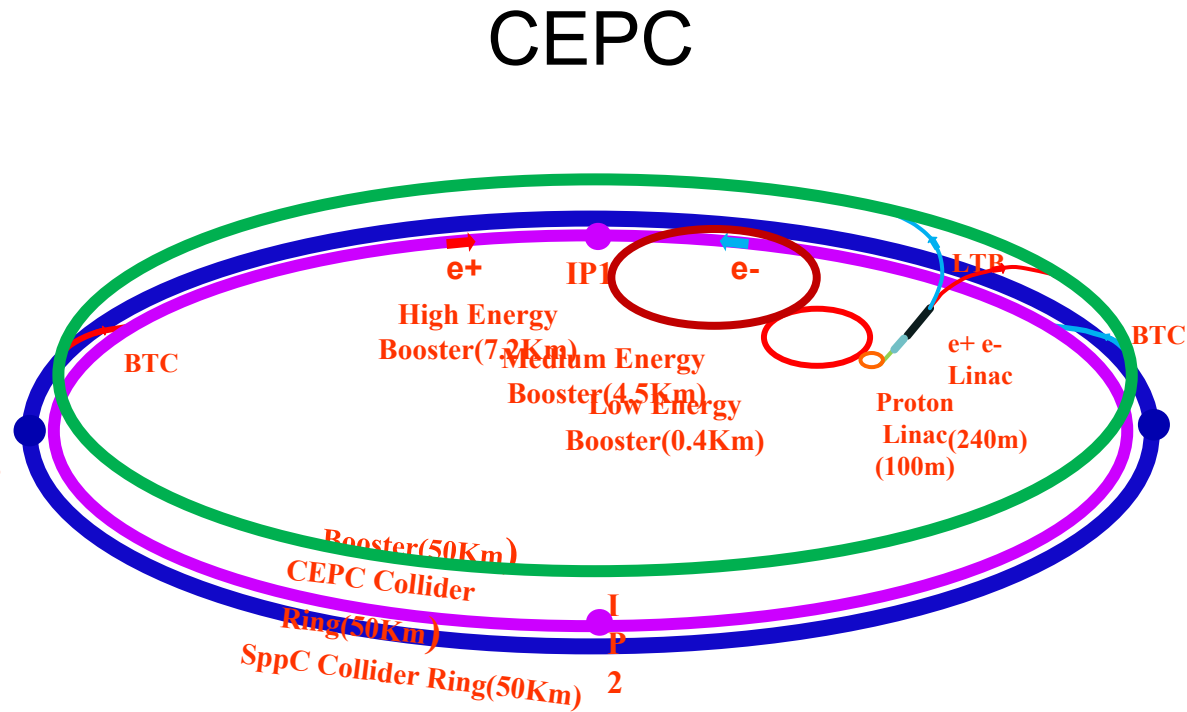
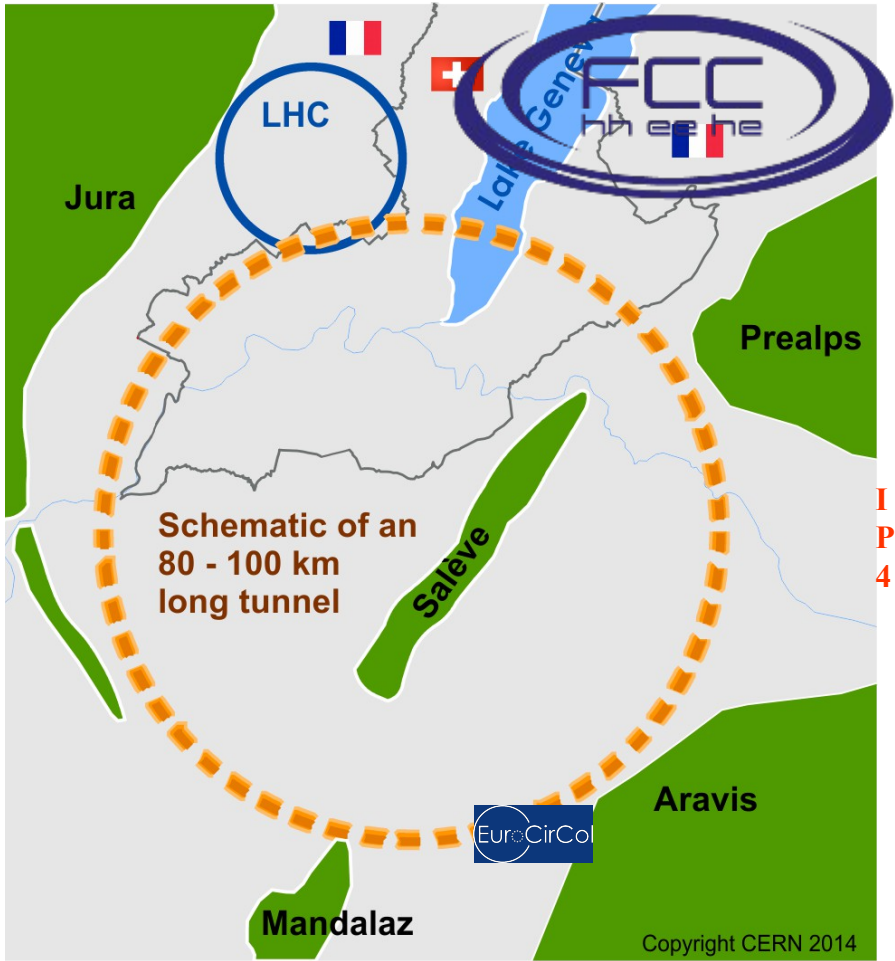
More details on Higgs physics at e+e- colliders see talk by Junping Tian



Energy: 0.1 - 1 TeV
Electron (and positron)
polarisation
TDR in 2013
+ DBD for detectors
 Footprint 31 km

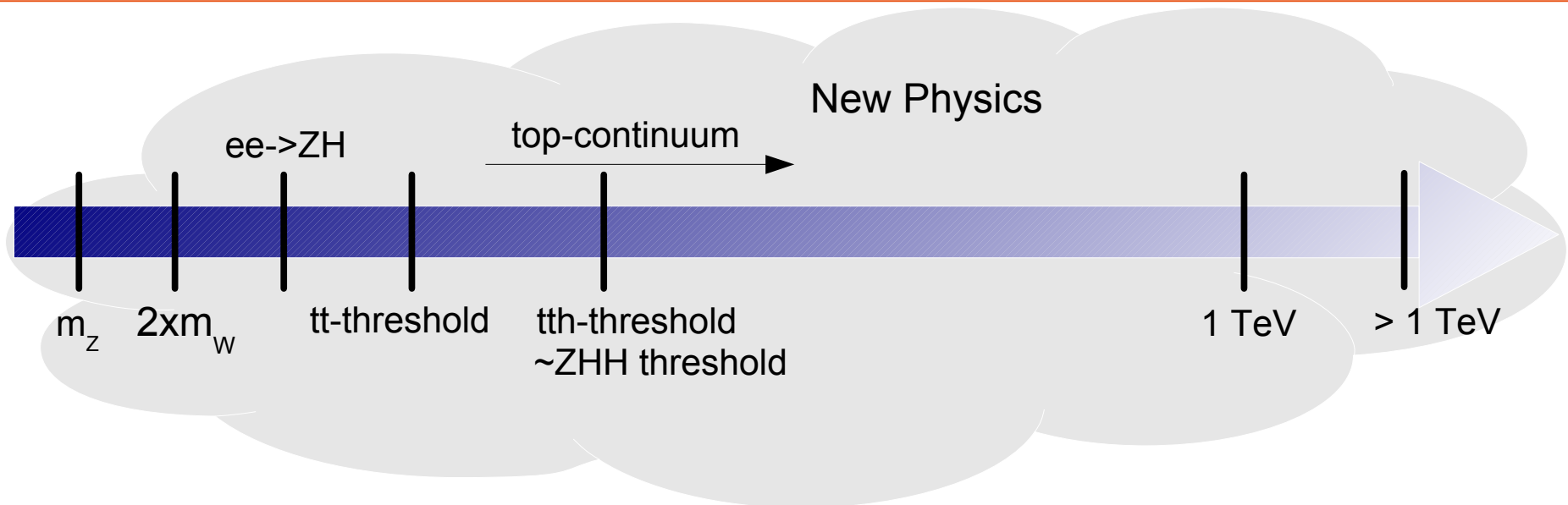


Energy: 0.5 - 3 TeV
CDR in 2012
 Footprint 48km



- ~100 km storage rings
- Coupled to hadron collider proposal
- 90 – 350 GeV cms energy
- No long. beam polarisation
- CDR Phase

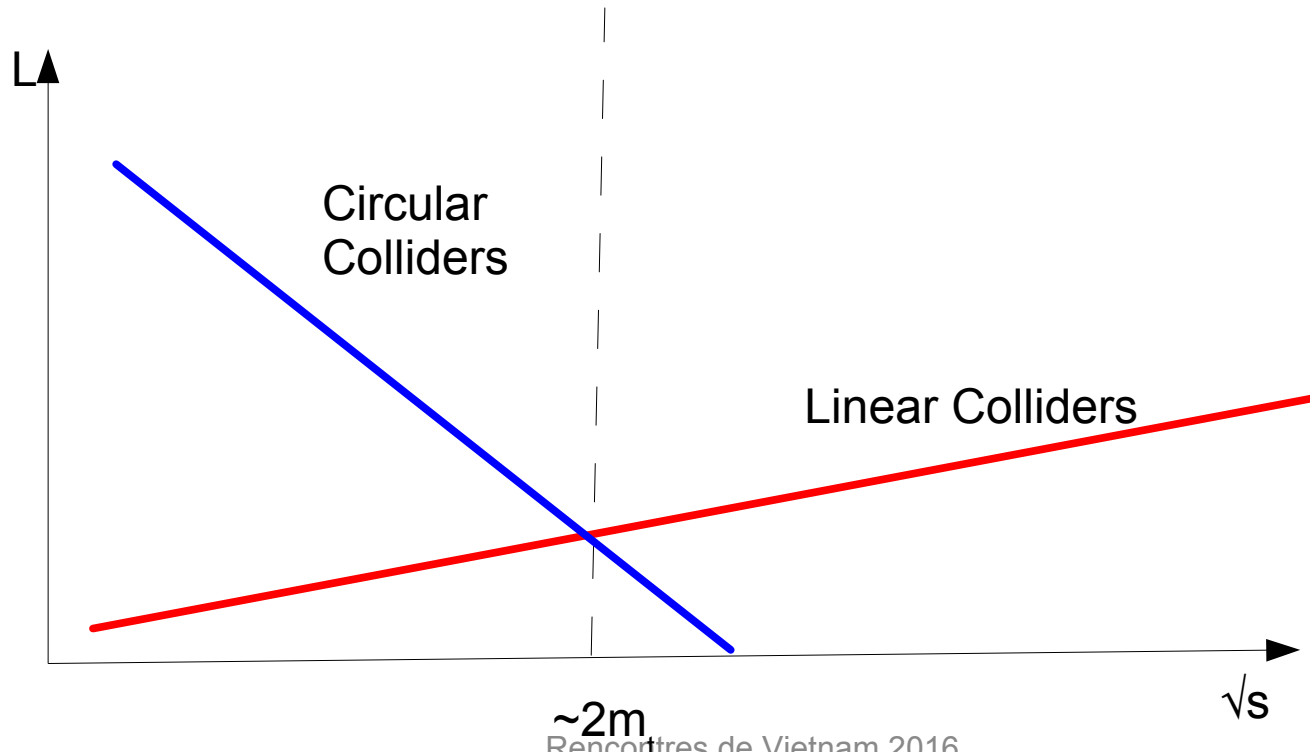
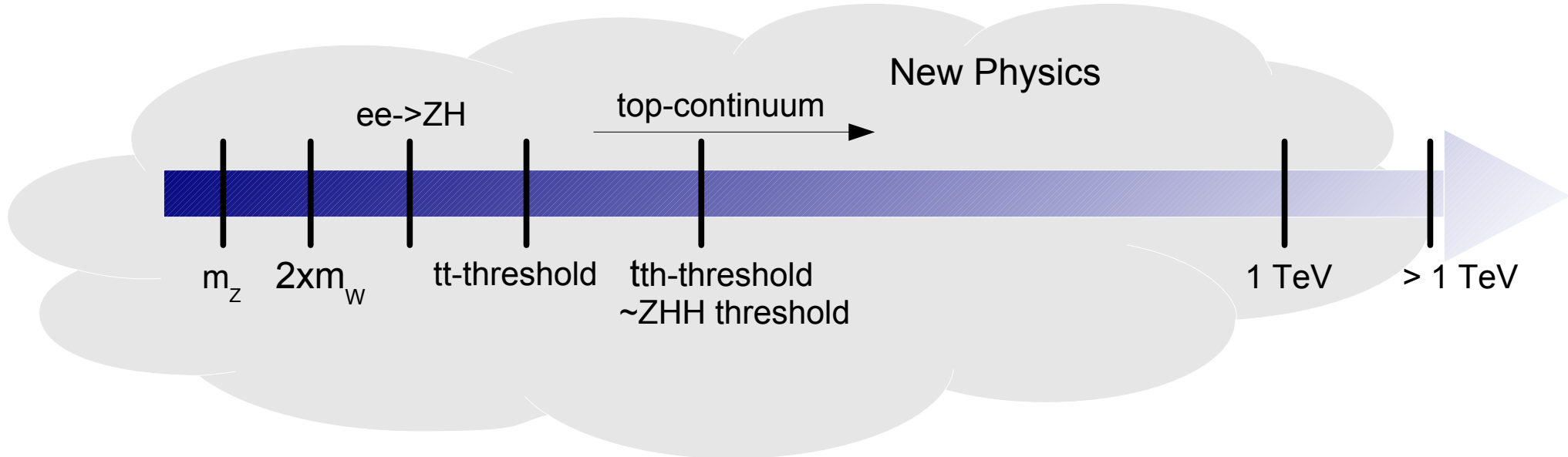
- ~50 km storage rings
- Coupled to hadron collider proposal
- 90 – 240 GeV cms energy
- No long. beam polarisation
- (Pre-)CDR Phase



- All Standard Model particles within reach of planned e+e- colliders
 - High precision tests of Standard Model over wide range to detect onset of New Physics
- Machine settings can be “tailored” for specific processes
 - Centre-of-Mass energy
 - Beam polarisation

$$\sigma_{P,P'} = \frac{1}{4} [(1 - PP')(\sigma_{LR} + \sigma_{RL}) + (P - P')(\sigma_{RL} - \sigma_{LR})]$$

- “Background free” searches for BSM through beam polarisation



Theory

Loop calculations (prospect on 3-loops)

... to distinguish new effects from ordinary effects, how many are needed?

Global Fits

... to assure consistency of results

New models

... as concrete manifestations of new physics

Effective field theory

... for generic effects of new physics

Experiment

Better accelerators

- Higher accelerating gradients
- Beam polarisation
- Nanosize beams

Better detectors

- 4π hermetic
- Highly granular devices

New analysis techniques

... multi-variate techniques
(must however not replace first two points)

Track momentum: $\sigma_{1/p} < 5 \times 10^{-5}/\text{GeV}$ (1/10 x LEP)

(e.g. Measurement of Z boson mass in Higgs Recoil)

Impact parameter: $\sigma_{d_0} < [5 \oplus 10/(p[\text{GeV}]\sin^{3/2}\theta)] \mu\text{m}$ (1/3 x SLD)

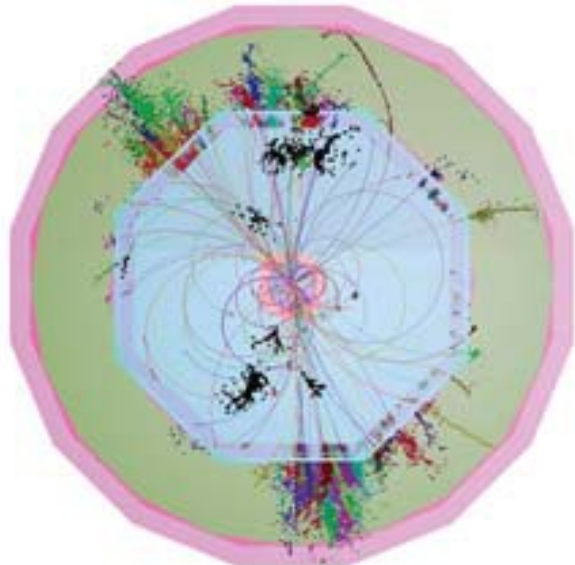
(Quark tagging c/b)

Jet energy resolution : $dE/E = 0.3/(E(\text{GeV}))^{1/2}$ (1/2 x LEP)

(W/Z masses with jets)

Hermeticity : $\theta_{\min} = 5 \text{ mrad}$

(for events with missing energy e.g. SUSY)



Final state will comprise events with a large number of charged tracks and jets(6+)

- High granularity
- Excellent momentum measurement
- High separation power for particles
- Particle Flow Detectors

2. Electroweak precision tests

Copied from deBlas, Higgs-Hunting 2016

- Precise measurements of W&Z properties taken at e+e- colliders

$$M_Z, \Gamma_Z, \sigma_{had}^0, \sin^2 \theta_{eff}^{lept}, P_{\tau}^{Pol}, A_f, A_{FB}^{0,f}, R_f^0$$

Z-Pole observables
SLD/LEP
0.002 - O(1%)

$$M_W, \Gamma_W$$

W-observables
LEP2
0.02 - O(1%)

- Tevatron/LHC **but in future also from e+e- colliders**

$$M_W, \Gamma_W$$

0.02-O(1%)

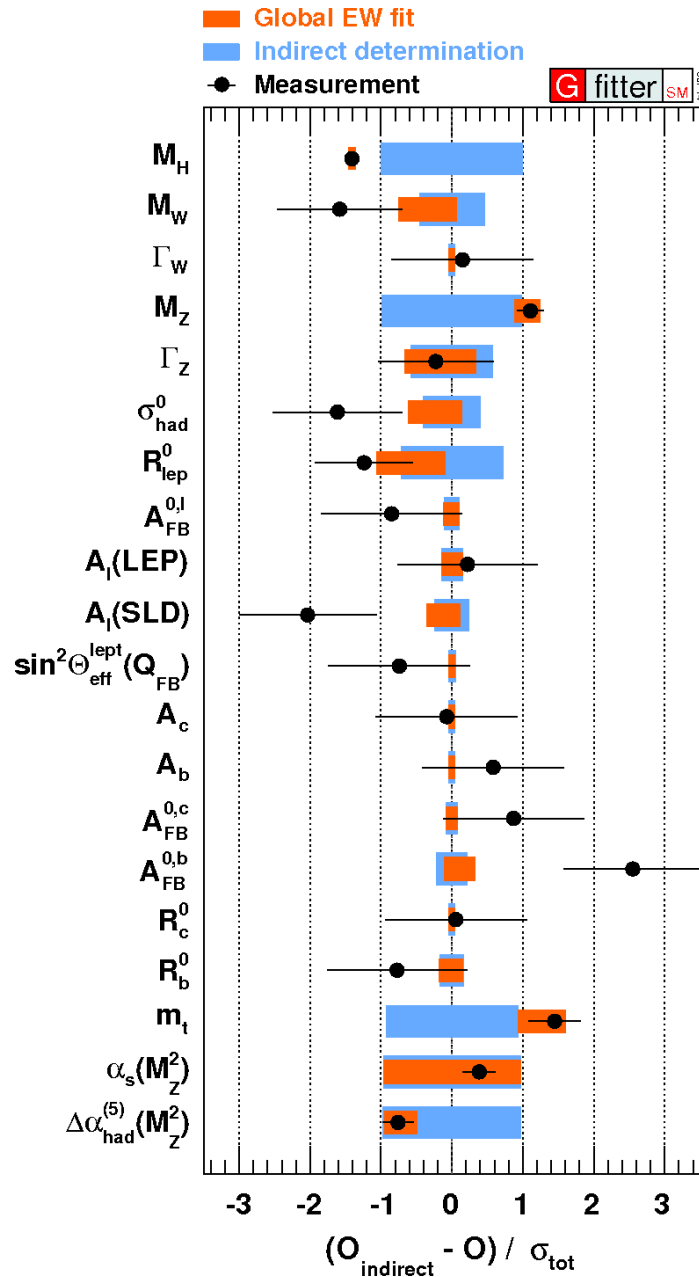
$$m_t$$

0.4%

$$M_H$$

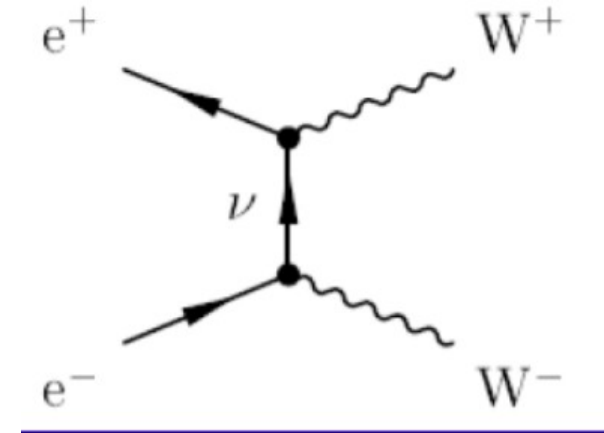
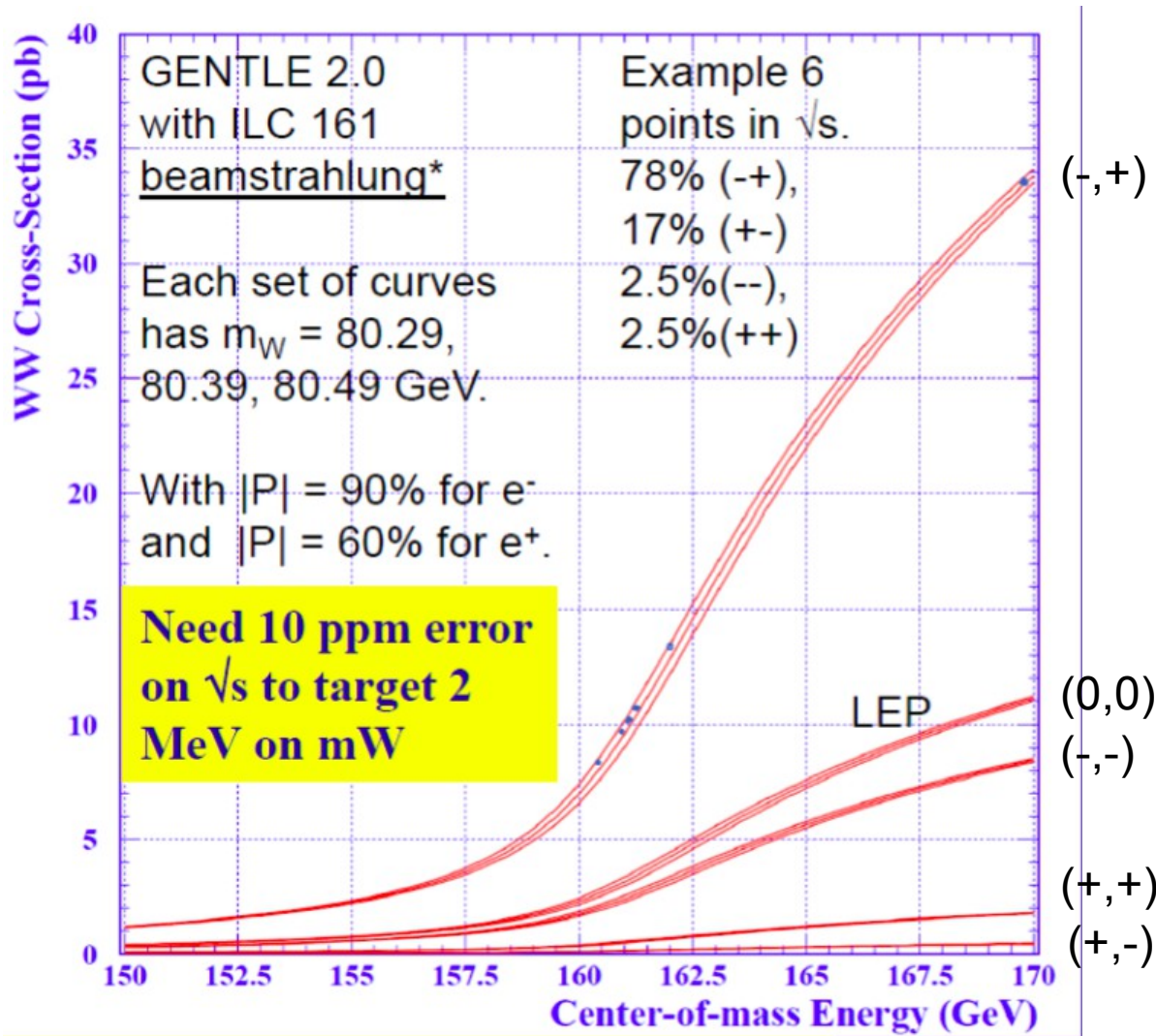
0.2%

Note competitive measurements of $\sin^2 \theta_{eff}^2$ by LHC



- Consistency between fit and measured values for vast majority of precision observables

With one notable exception!!!!
(see later)

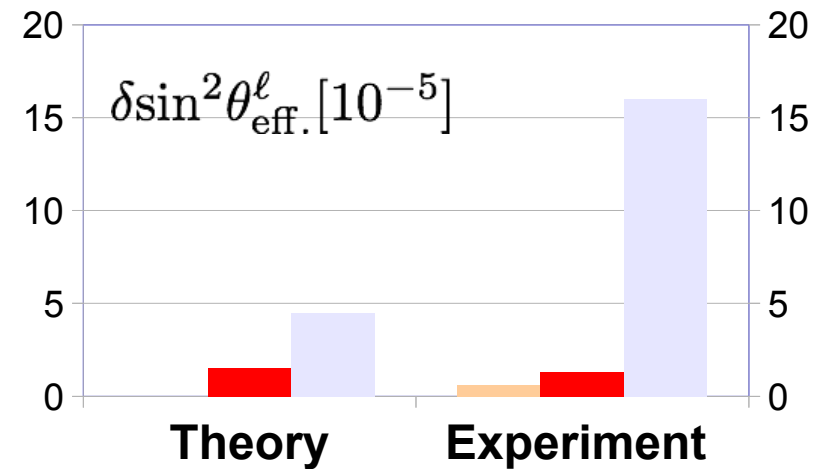
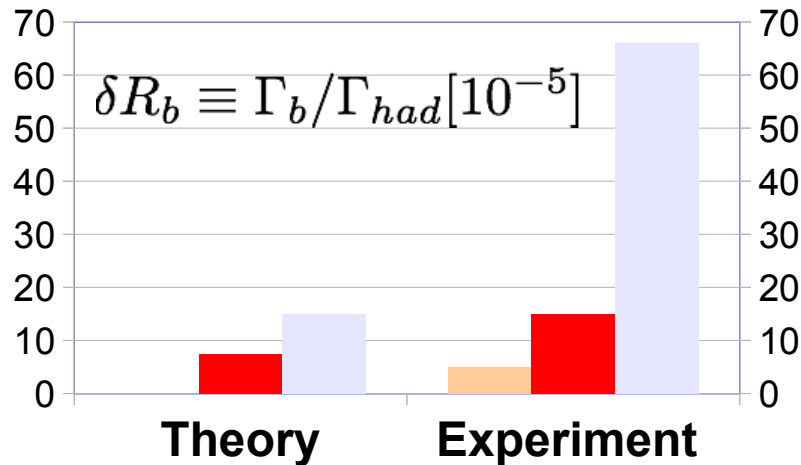
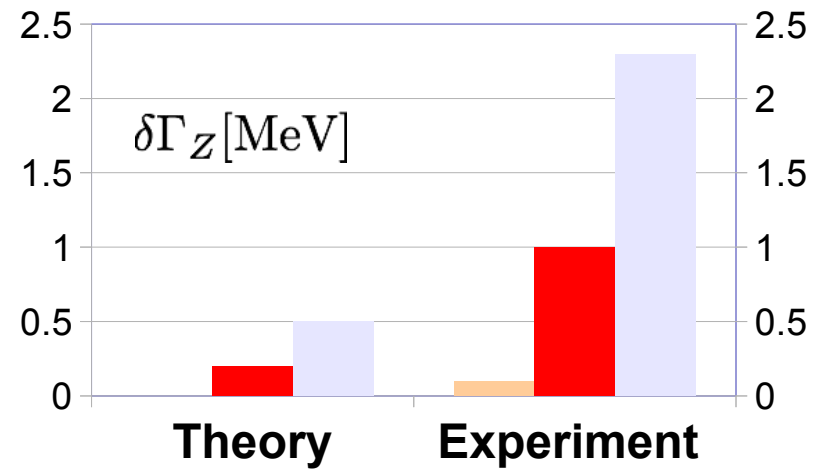
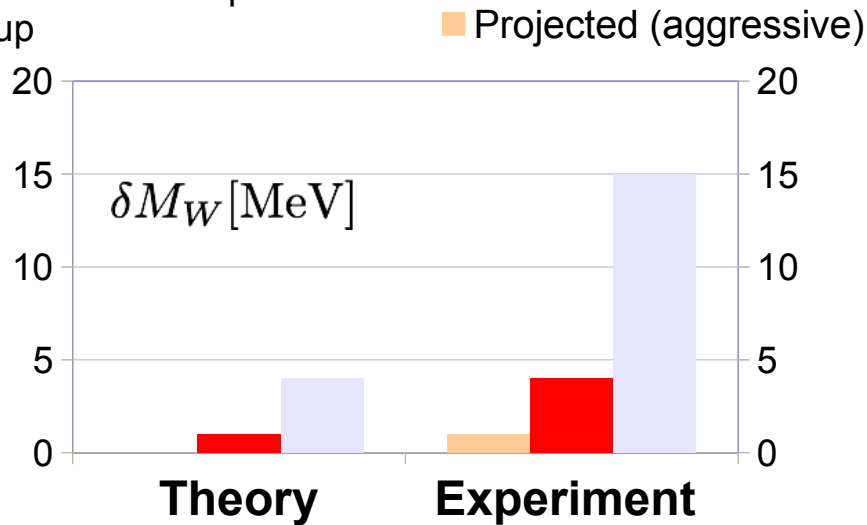


- Robust method
- Beam polarisation essential
- Need extreme good control of beam energy

G. Wilson

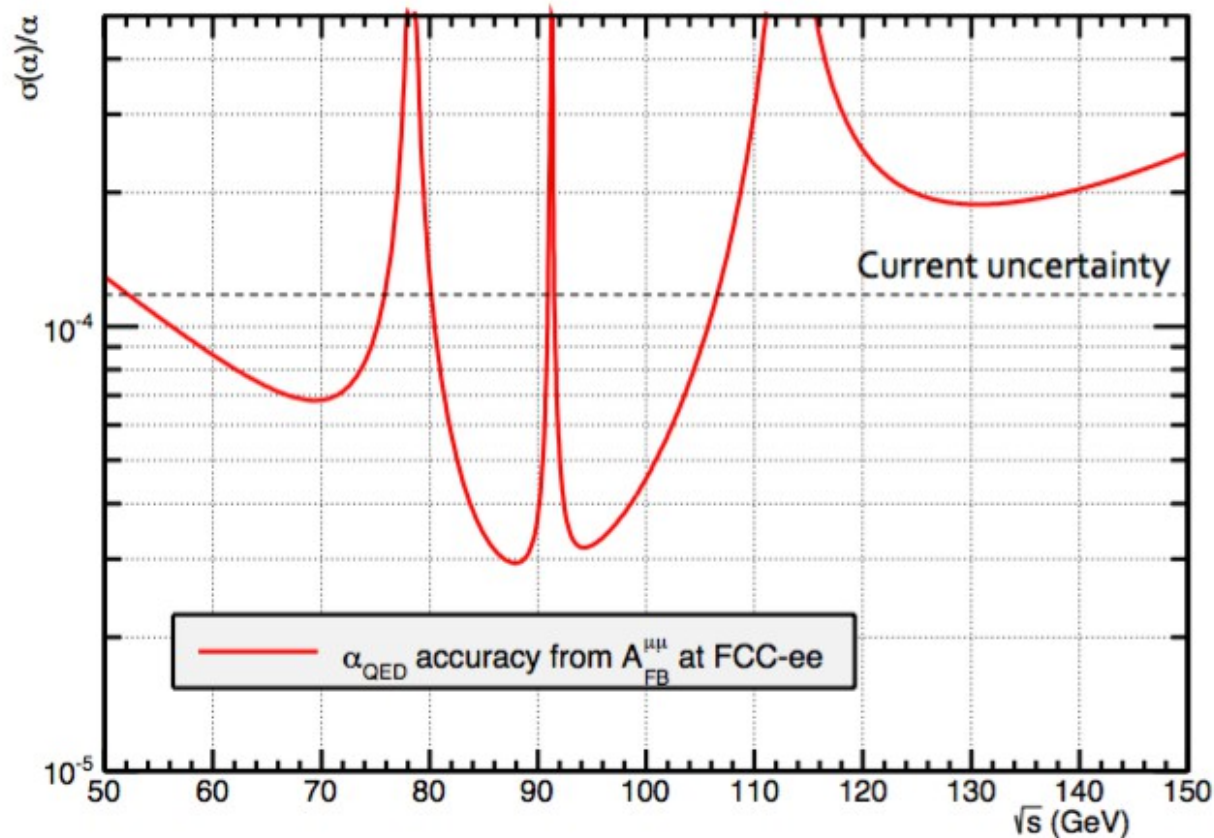
à la Freitas, FCC-ee Workshop Feb. 2016

Details -> backup



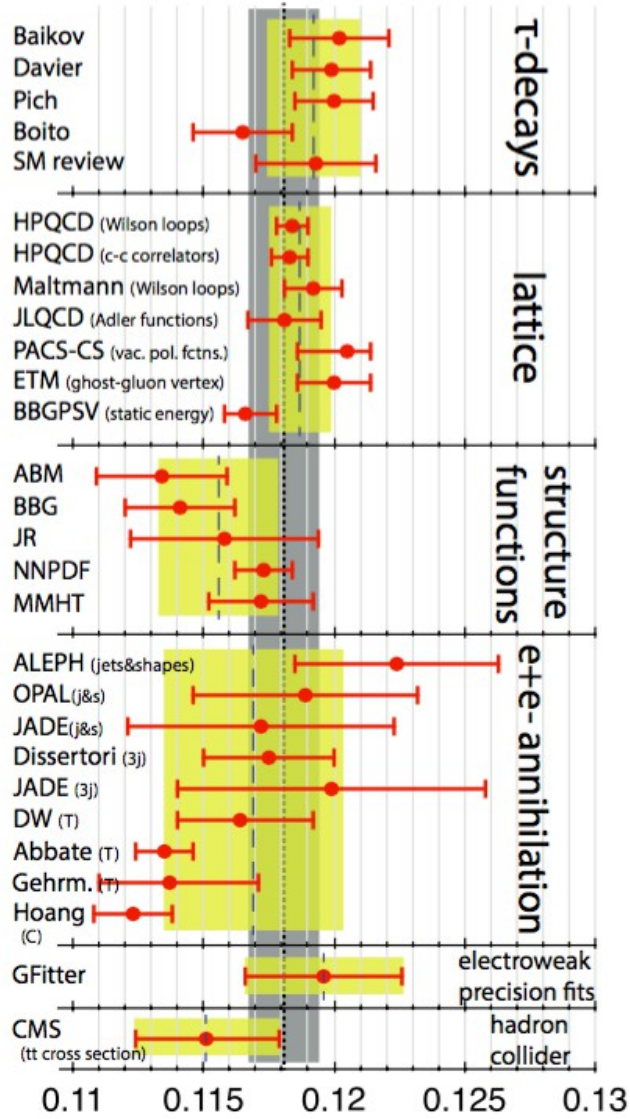
- Optimistic/aggressive scenario will lead to theory dominated errors
- Theory errors may become guide line for planning of future projects

Prospects for α_{QED} in high statistics e^+e^- running



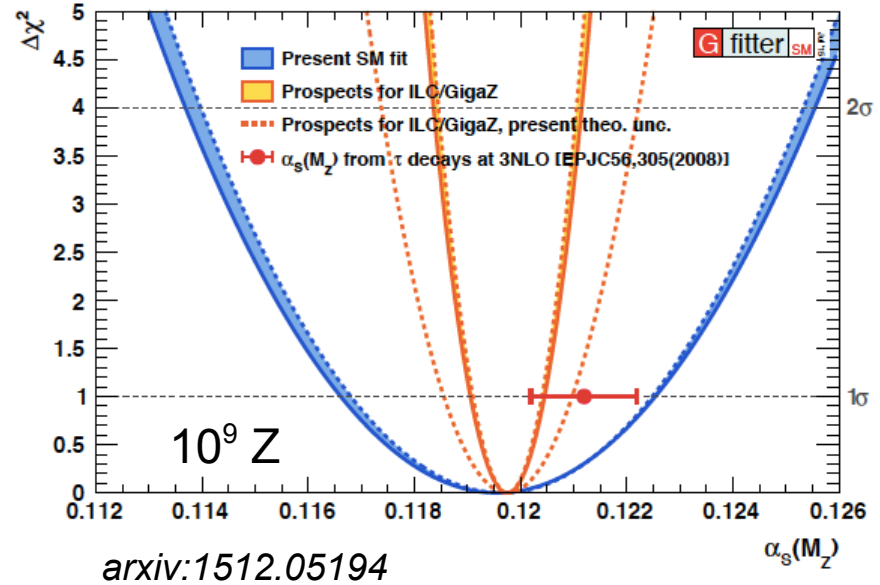
- Combining measurements slightly above and below the Z-Pole to reduce experimental errors
- Most optimistic experimental prospection arrives at $d\alpha/\alpha$ 3×10^{-5}
- Puts challenge on theory, current prospect $\sim 10^{-4}$

Current status



Dominated by lattice QCD

Prospects Z-running



Electroweak fit with updated EWPO:
and theory uncertainties

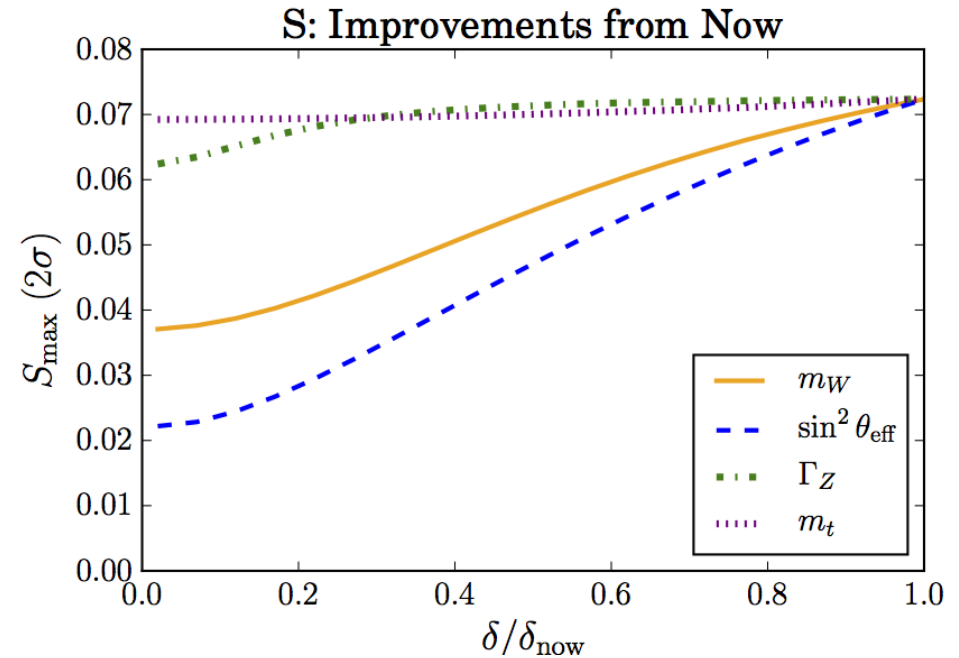
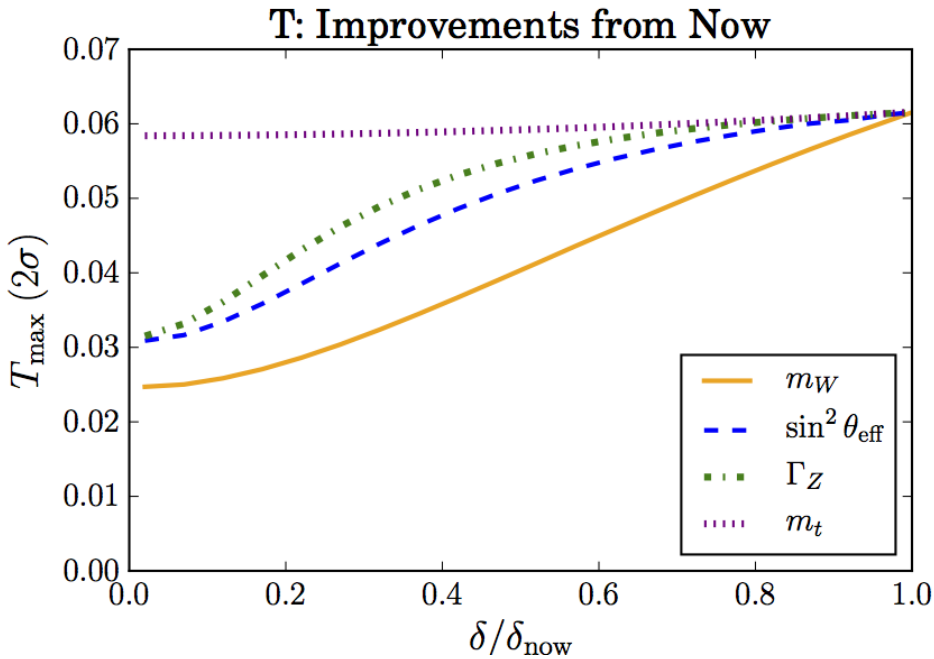
$$\delta \alpha_s(M_Z) \sim 0.0007 \text{ for } 10^9 Z$$

$$\delta \alpha_s(M_Z) \sim 0.0003(16) \text{ for } 10^{12} Z$$

Prospects Lattice

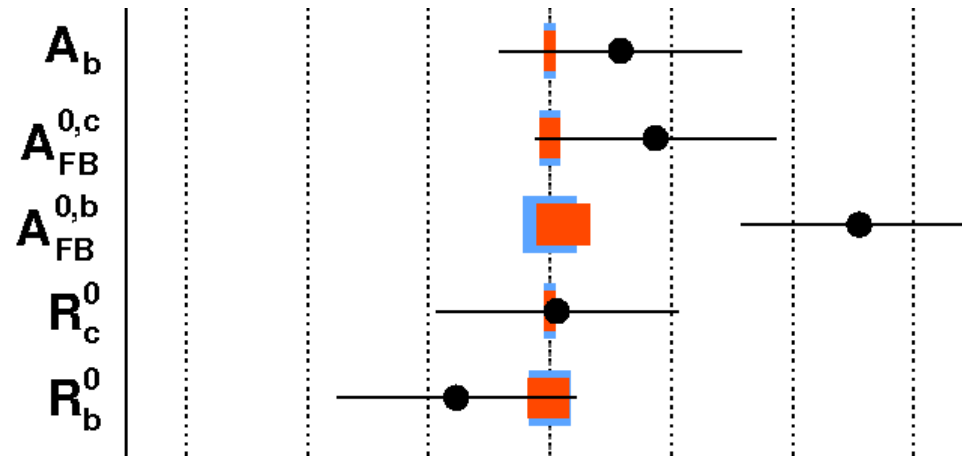
$$\delta \alpha_s(M_Z) \sim 0.0003$$

à la M. Reece 1609.03018

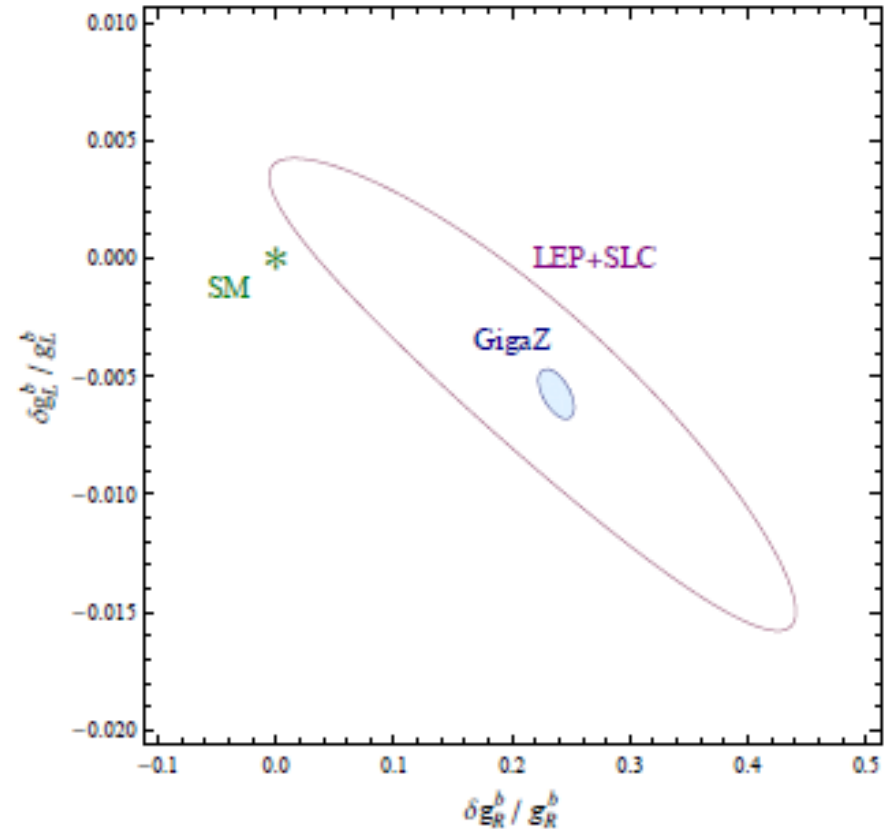


- T depends mainly on M_W
 - Threshold scan around $2M_W$ (with polarised beams), see above
- S depends mainly on $\sin^2 \theta_{\text{eff}}^2$.
 - Ultra-high statistics sample $ee \rightarrow ff$ at Z-mass or smaller sample exploiting beam polarisation (Remember LEP and SLC times)
- Precision on S and T seems to saturate at $\delta/\delta_{\text{now}} \approx 0.1 - 0.2$

$\sim 3\sigma$ in heavy quark observable A_{FB}^b

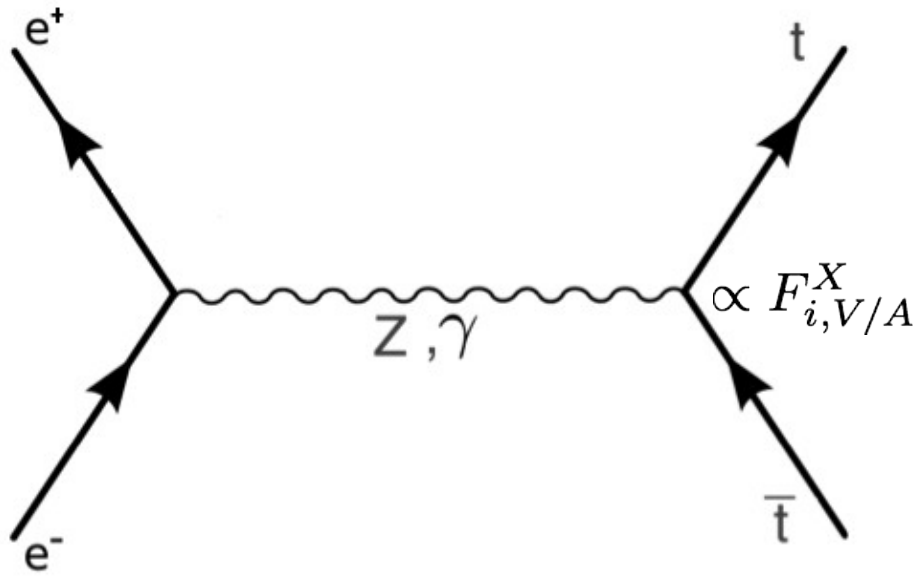


Projection with $10^9 Z$



- Is tension due to underestimation of errors or due to new physics?
- High precision e+e- collider will give final word on anomaly
- In case it will persist polarised beams will allow for discrimination between effects on left and right handed couplings (Remember $Zb_l b_l$ is protected by cross section)

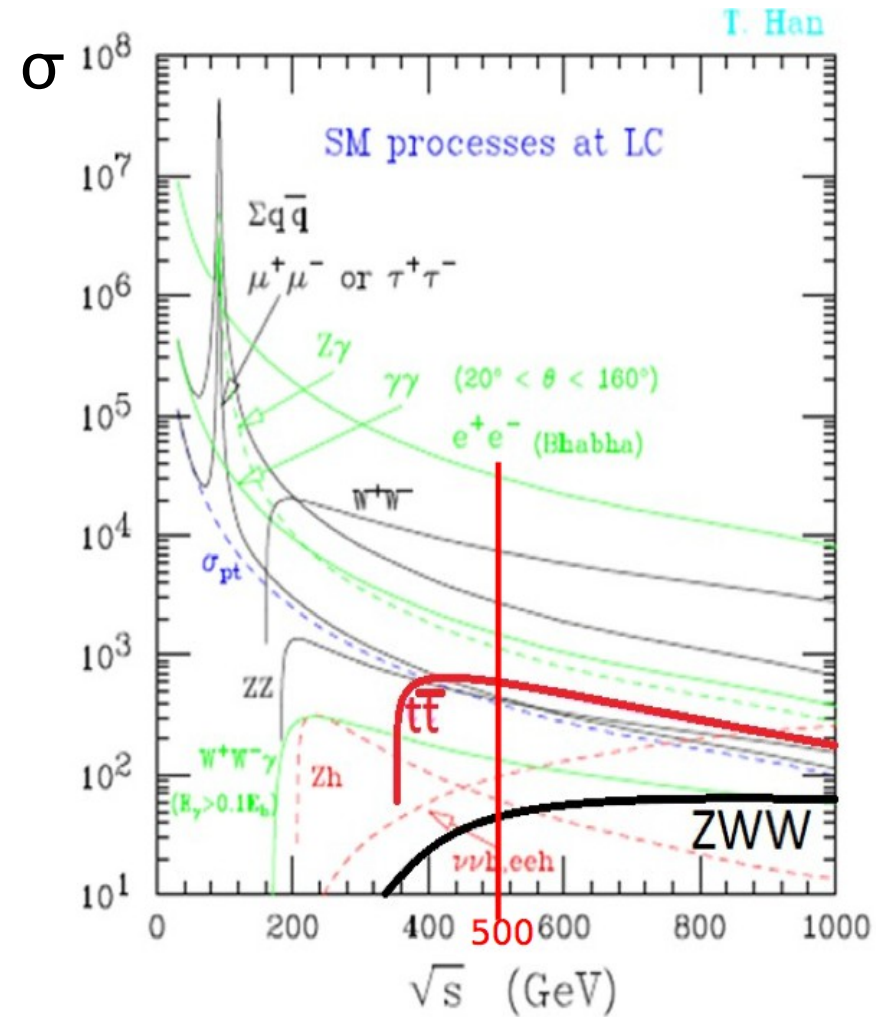
3. Top physics



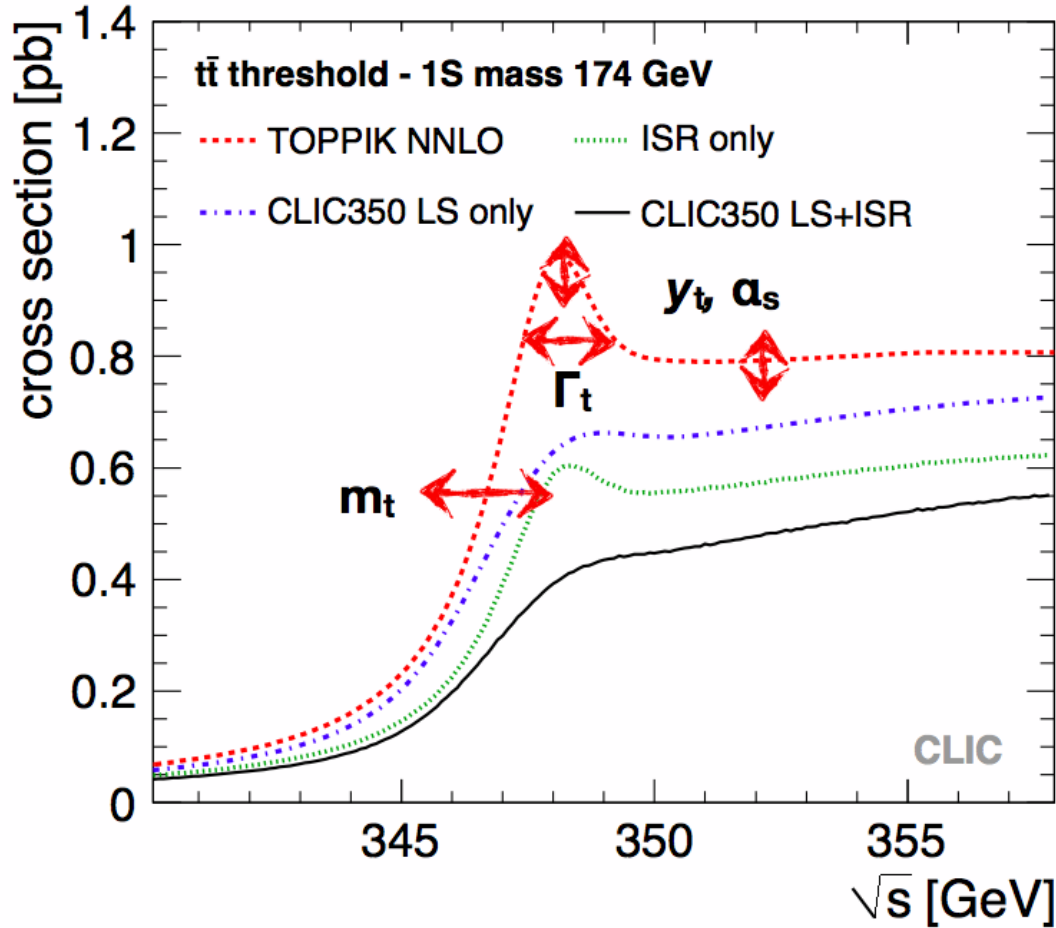
- Top quark production through electroweak processes
no competing QCD production => Small theoretical errors!

- High precision measurements

- Top quark mass at ~ 350 GeV through threshold scan
- Polarised beams allow testing chiral structure at $t\bar{t}X$ vertex
=> Precision on form factors F

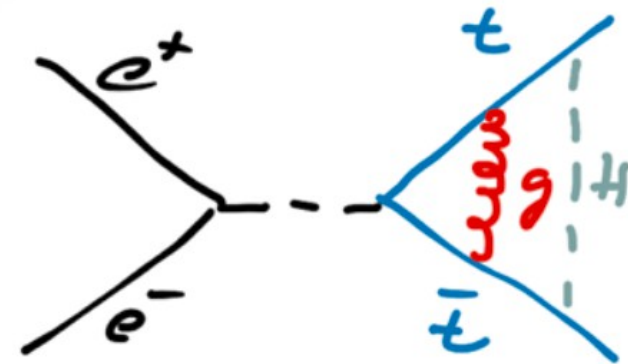


Small size of $t\bar{t}$ “bound state” at threshold ideal premise for precision physics

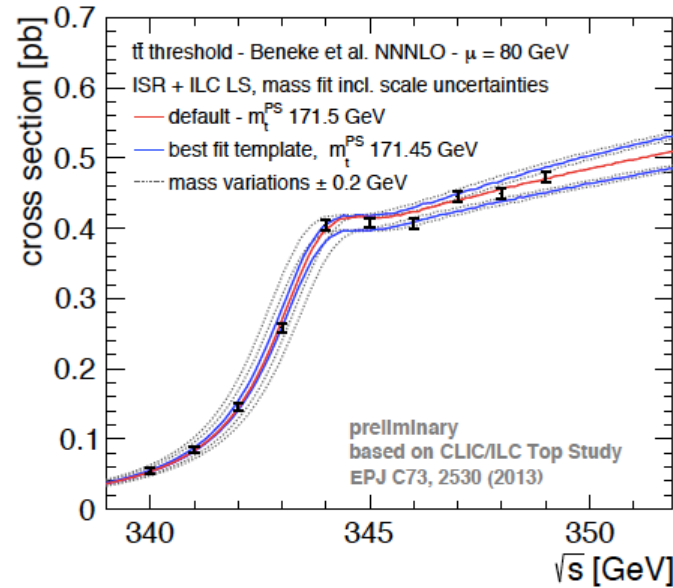


Cross section around threshold is affected by several properties of the top quark and by QCD

- Top mass, width Yukawa coupling
- Strong coupling constant



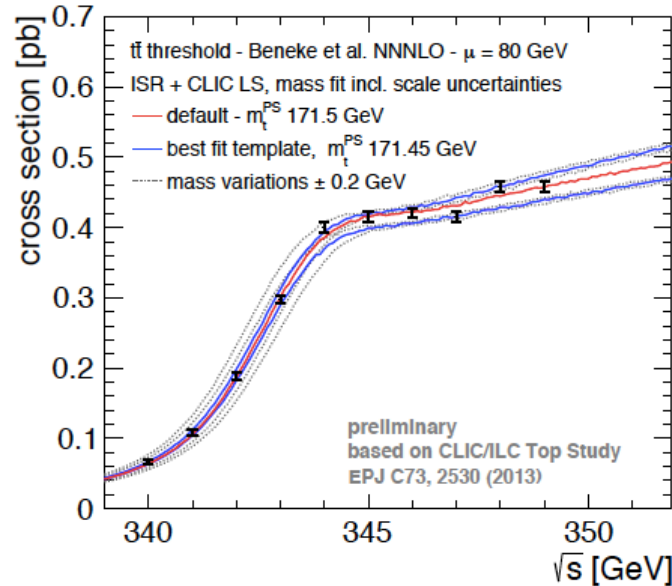
Effects of some parameters are correlated:
Dependence on Yukawa coupling rather weak,
Precise external α_s helps



ILC

Fit uncertainty:
28.5 MeV (18 MeV stat)

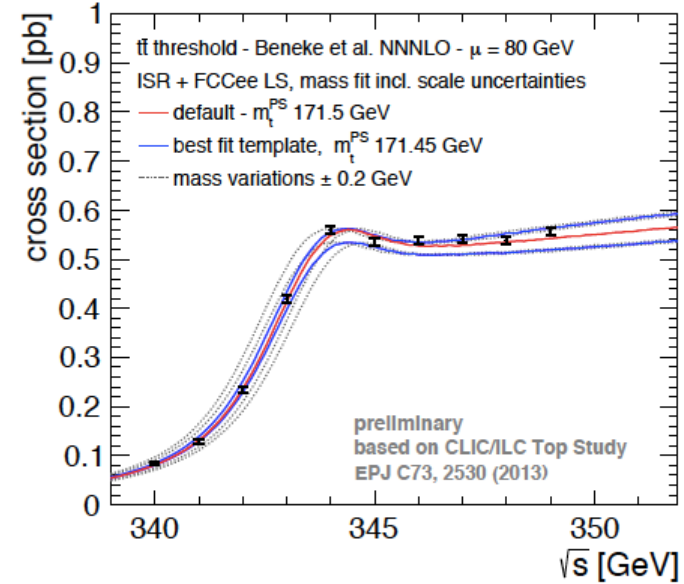
Scale uncertainty:
40 MeV



CLIC

Fit uncertainty:
31 MeV (21 MeV stat)

Scale uncertainty:
42 MeV



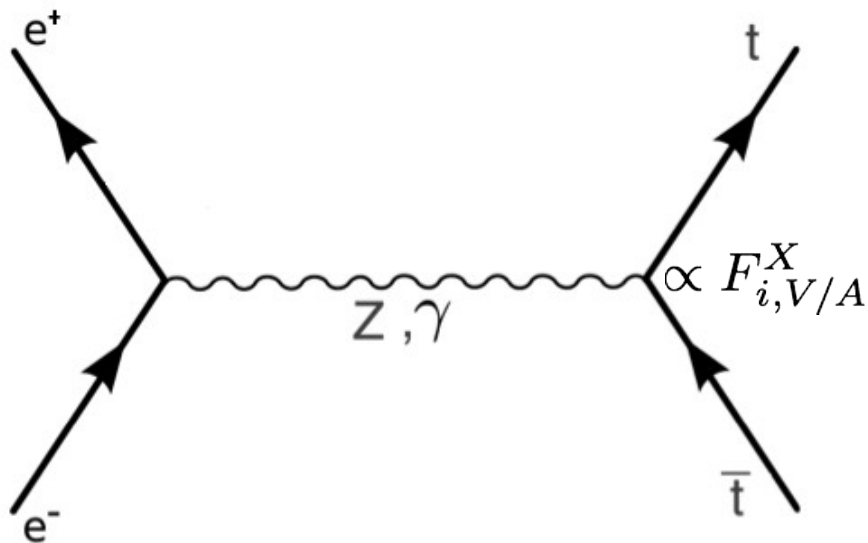
FCC-ee

Fit uncertainty:
27 MeV (15 MeV stat)

Scale uncertainty:
40 MeV

More details on top mass in e+e- collisions by A. Hoang and M. Beneke

- Fermion mass generation closely related to the origin electroweak symmetry breaking
- Expect residual effects for particles with masses closest to symmetry breaking scale



Manifestation of New Physics:

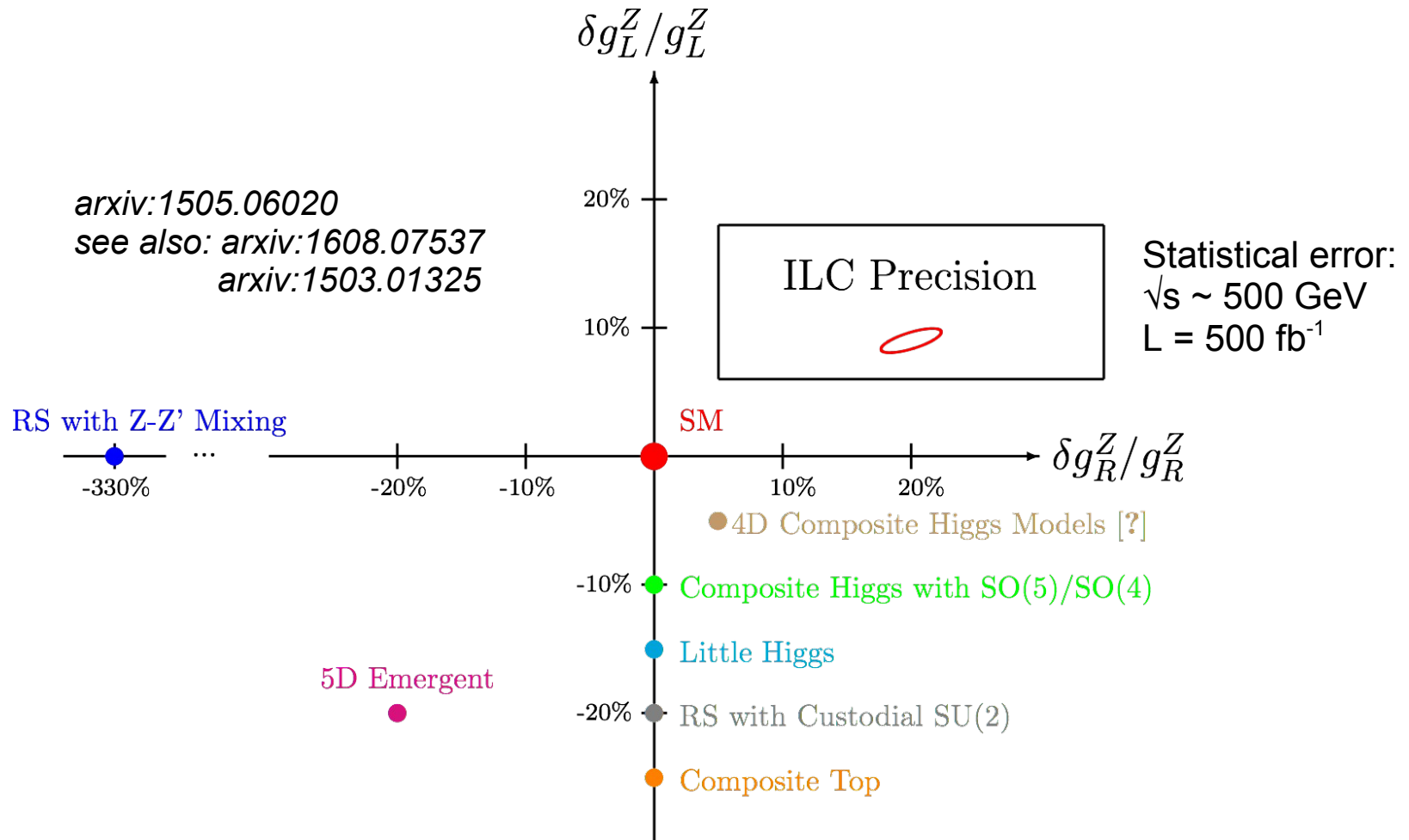
- Modification of Ztt coupling
Mixing between top and partners
Mixing Z/Z'
- s-channel exchange of New Z'
Including interference effects

$$\Gamma_{\mu}^{ttX}(k^2, q, \bar{q}) = -ie \left\{ \gamma_{\mu} (F_{1V}^X(k^2) + \gamma_5 F_{1A}^X(k^2)) + \frac{\sigma_{\mu\nu}}{2m_t} (q + \bar{q})^{\nu} (iF_{2V}^X(k^2) + \gamma_5 F_{2A}^X(k^2)) \right\}, \quad (2)$$

Pure γ or pure Z^0 : $\sigma \sim (F_i)^2 \Rightarrow$ No sensitivity to sign of Form Factors

Z^0/γ interference : $\sigma \sim (F_i) \Rightarrow$ Sensitivity to sign of Form Factors

Top is primary candidate to be a messenger new physics in many BSM models
Incorporating compositeness and/or extra dimensions

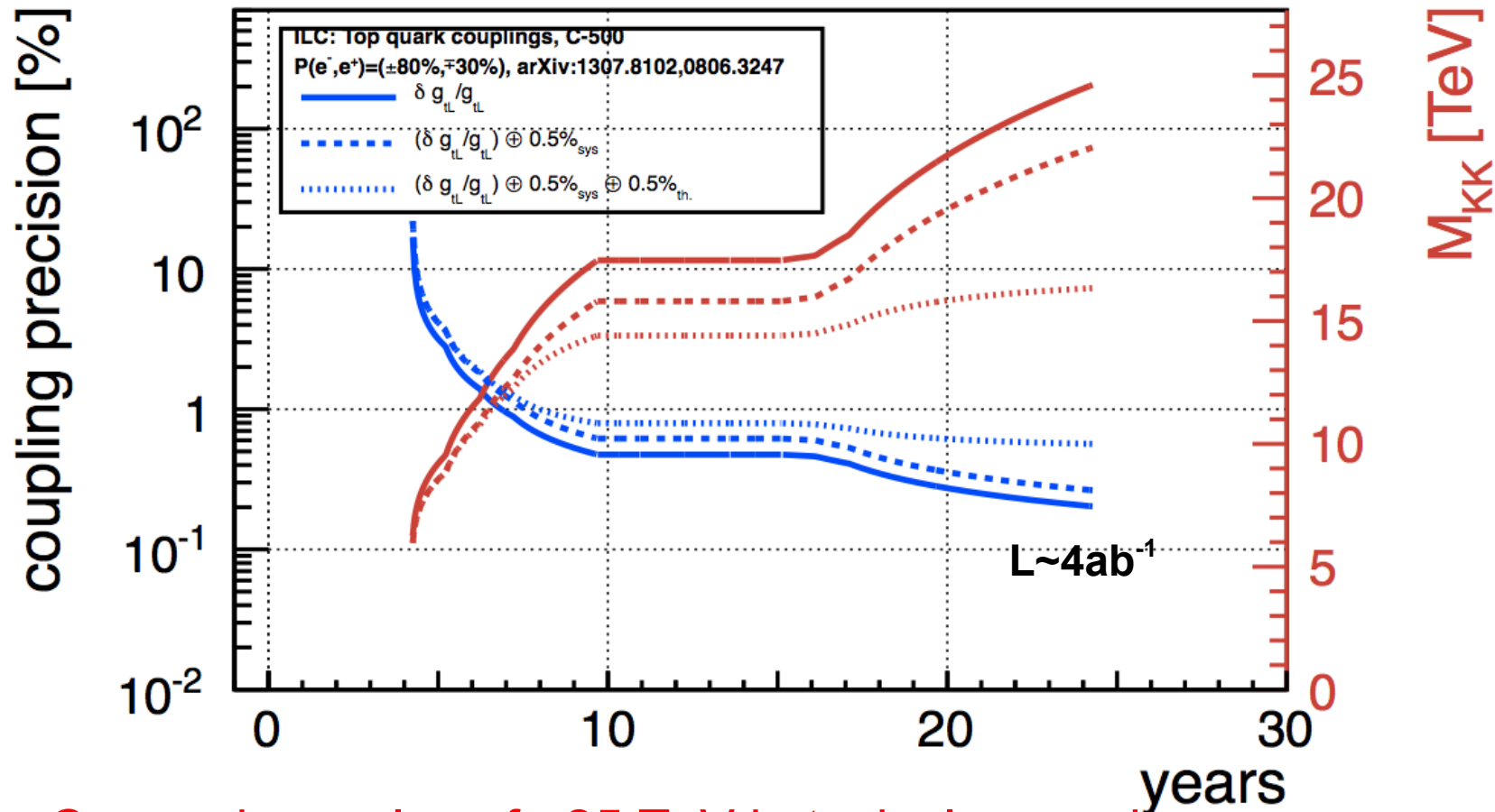


Precision expected for top quark couplings will allow to distinguish between models

Remark: All presented models are compatible with LEP elw. precision data

New physics reach for typical BSM scenarios with composite Higgs/Top and or extra dimensions

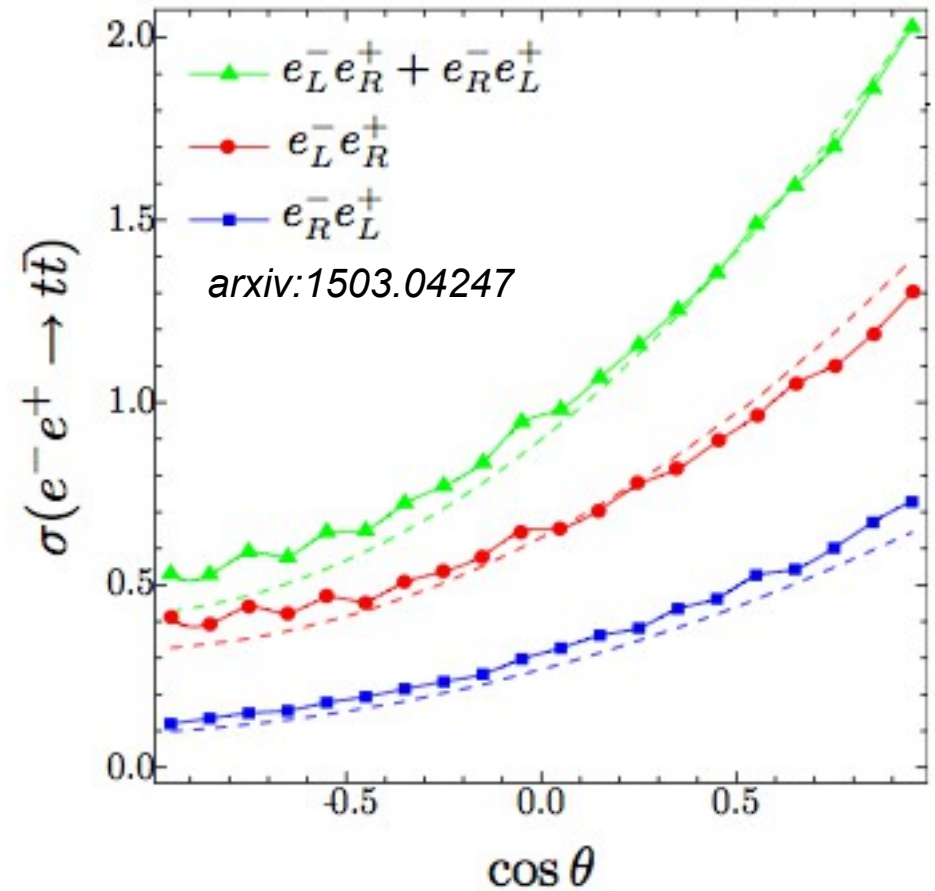
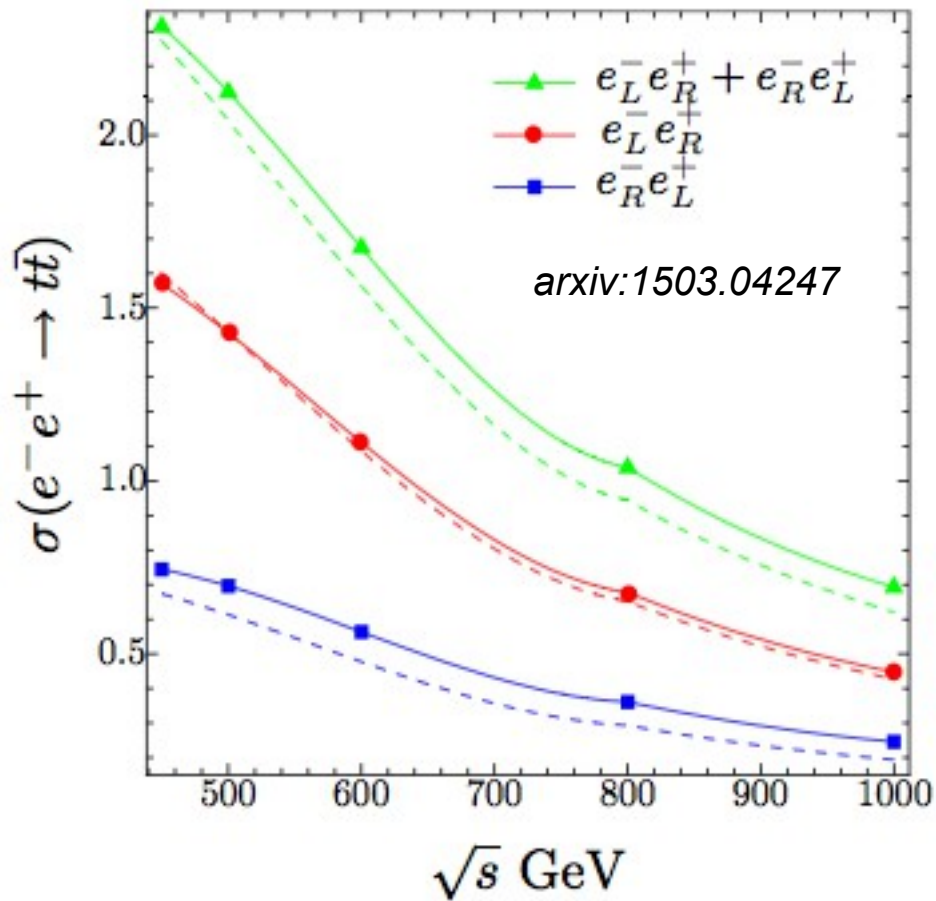
Based on phenomenology described in Pomerol et al. arXiv:0806.3247



Can probe scales of ~ 25 TeV in typical scenarios

(... and up to 80 TeV for extreme scenarios)

=> Important guidance for e.g. 100 TeV pp-collider



- Electroweak corrections manifest themselves differently for different beam polarisations

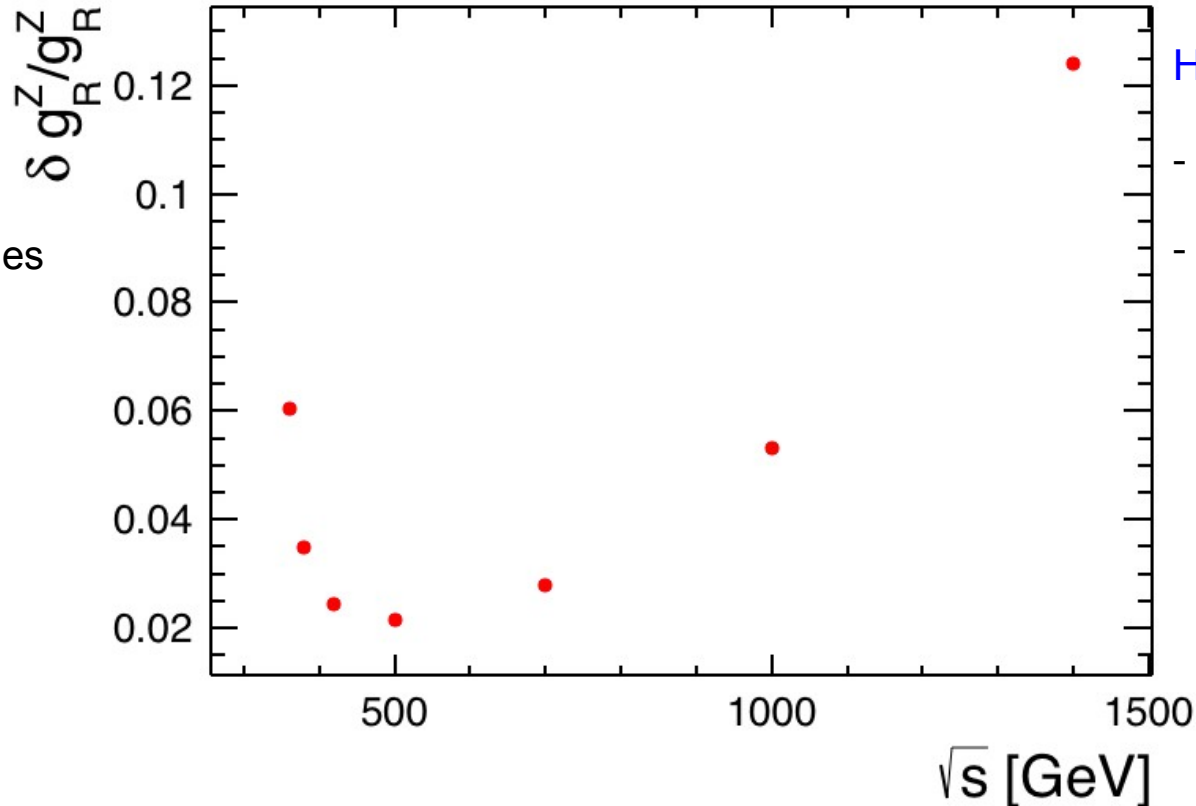
Beam polarisation important asset to disentangle SM and effects of new physics

Configuration $e_R^-e_L^+$ seems to lead to “simpler” corrections

... simplified discussion

Small cms energies:

- Vanishing axial vector coupling
- large QCD uncertainties ... and
- Lumi decreases at linear colliders



High cms energies:

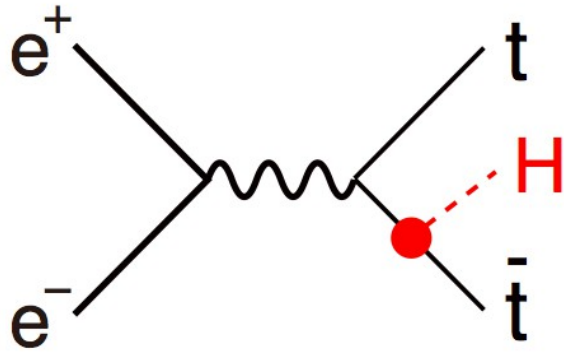
- Quickly decreasing cross section
- ... partially compensated by increasing luminosity

Broad minimum between 400 and 700 GeV

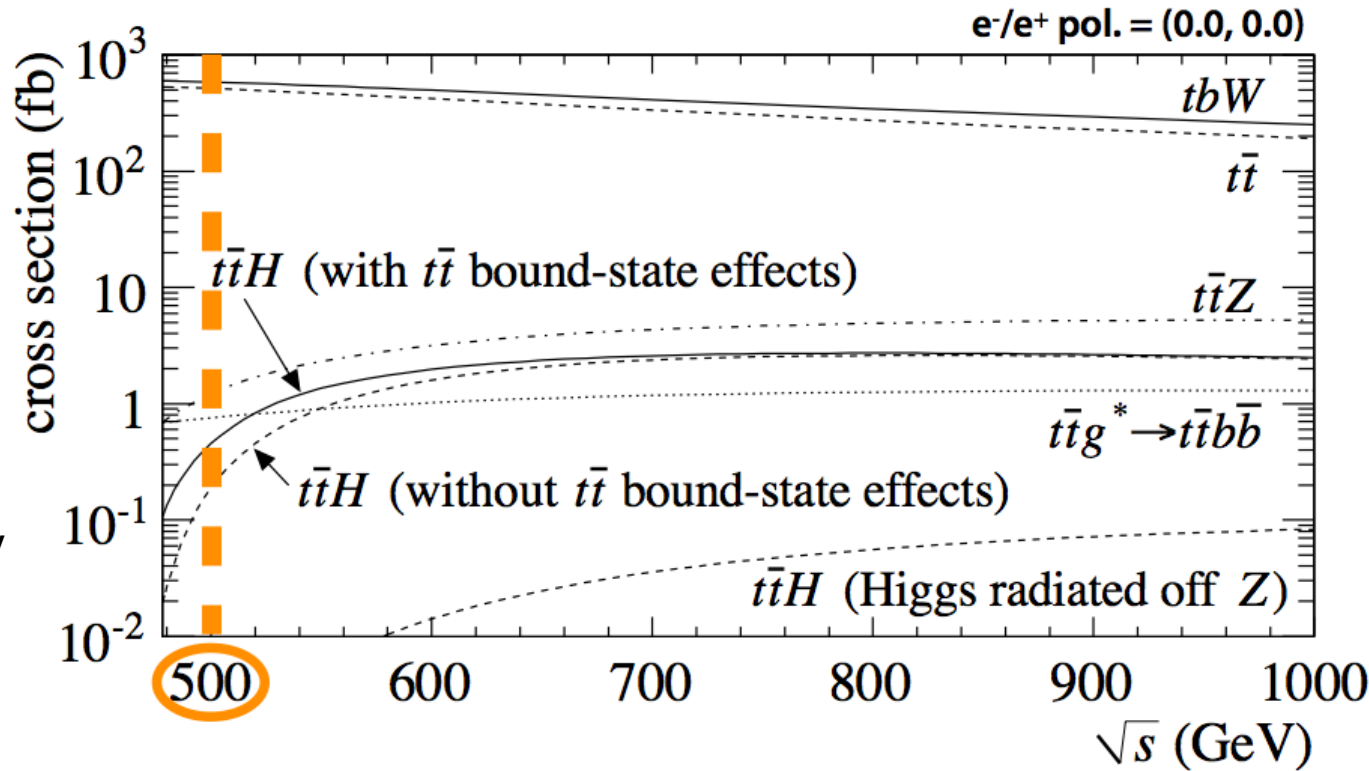
$\sqrt{s} \sim 500$ GeV is “sweet spot” for coupling measurements

However:

- Sensitivity to CP violating Higgs at smaller cms energies
- New physics at higher energies may increase cross section



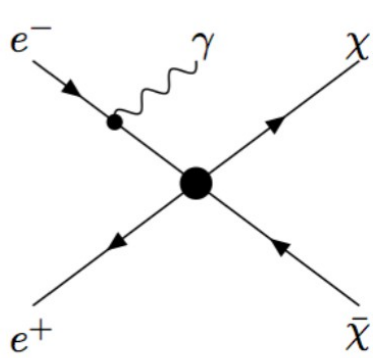
- Coupling of Higgs to heaviest particle known today
- Up to eight final state jets



$\Delta g_{ttH} / g_{ttH}$	500 GeV	+ 1 TeV
Snowmass	7.8%	2.0%
H20	6.3%	1.5%

4. BSM Physics at e^+e^- colliders

Z' with **vector couplings** to Dirac type Dark Matter X and **axial couplings** to ordinary matter

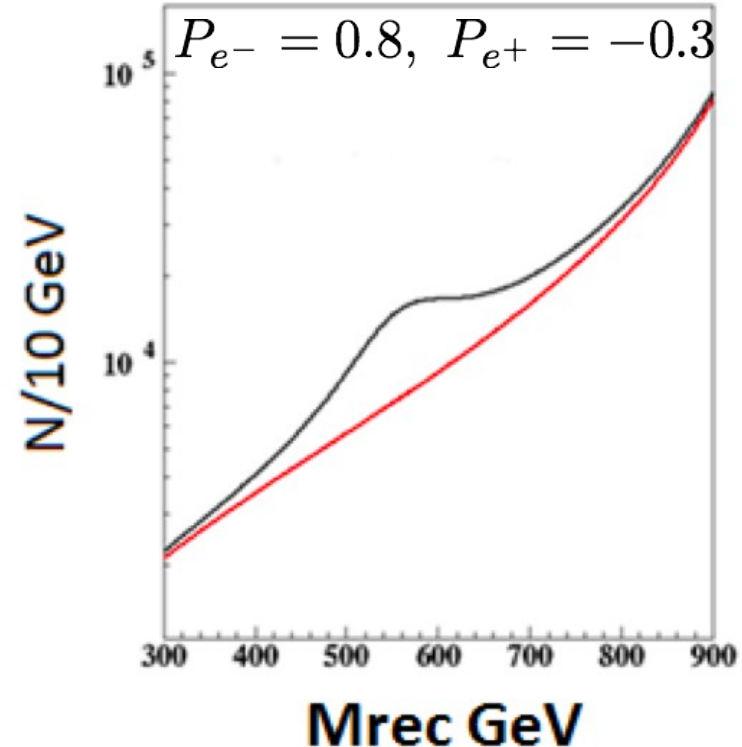
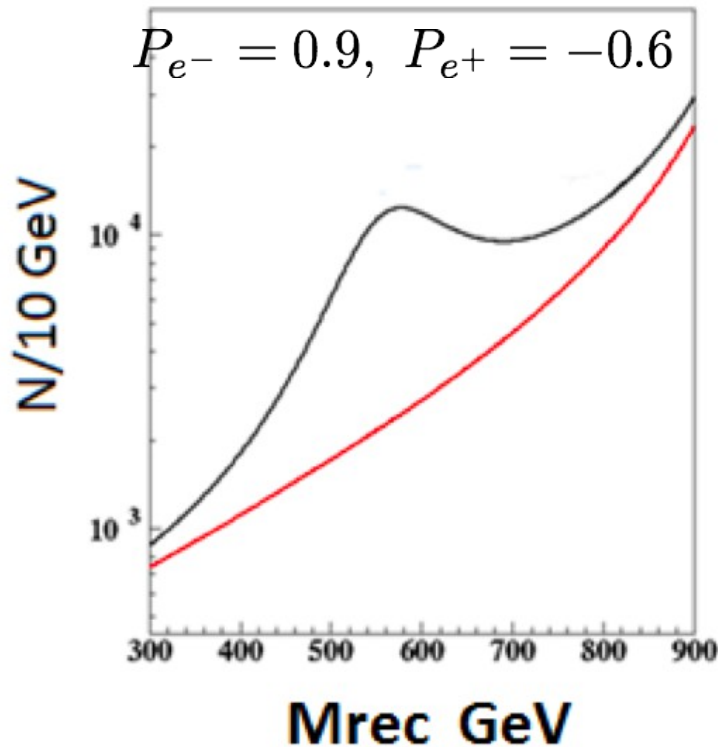


$$\sigma v = \boxed{|g_V^X|^2} K^2 \sum_f n_{cf} \boxed{|g_A^f|^2} \frac{2m_X^2 + s}{12\pi [(s - m_{Z'}^2) + (m_{Z'}\Gamma_{Z'})^2]}$$

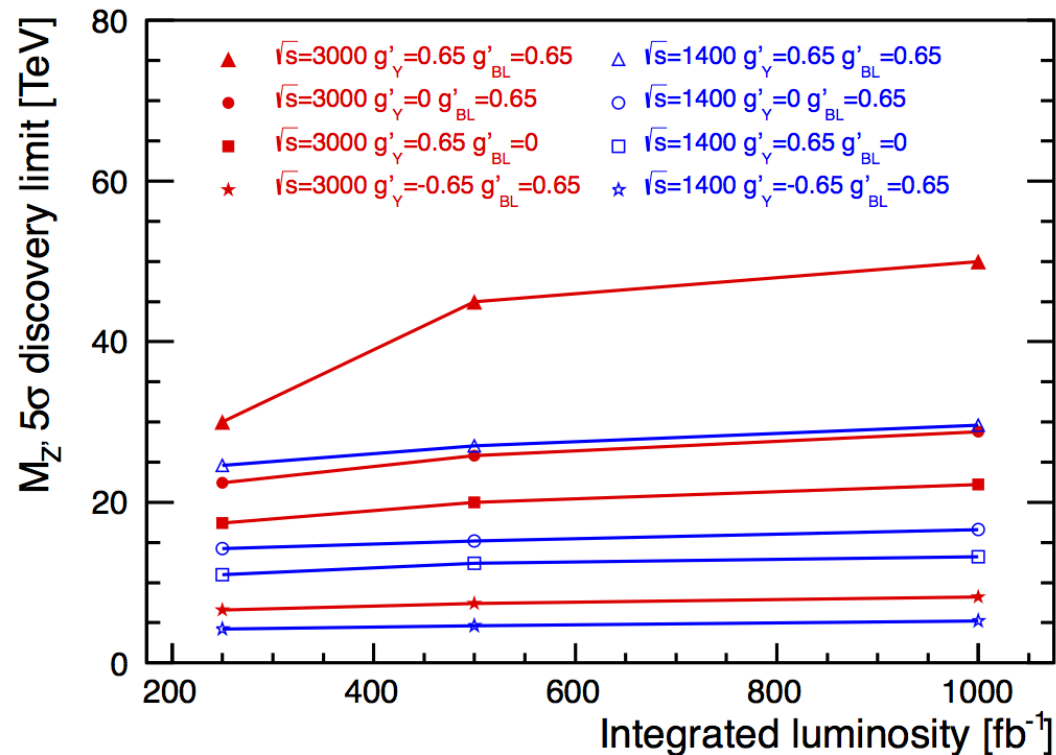
-> Monophoton search

-> Background suppression through polarised beams

$$\sqrt{s} = 1 \text{ TeV}, \quad m_{Z'} = 550 \text{ GeV}, \quad L = 1 \text{ ab}^{-1}, \quad K^2 = 0.1$$



Z' discovery limit in $e^+e^- \rightarrow \mu^+\mu^-$



- Remarkable discovery potential
- Different sensitivities for different scenarios
- Sensitivities to $M_{Z'}$ of up to 50 TeV @ $\sqrt{s} = 3$ TeV

- Versatile machines for precision physics in the range $m_Z - 3$ TeV under study
 - Linear and circular machines would/could be complementary
- LEP/SLC programme can be repeated with much higher precision
- Which precision is reasonable precision?
- (Higgs and) Top quark are physics guaranteed
 - (My conviction) both are messengers to New Physics
 - Electroweak top quark pair production
 - $t\bar{t}h$ coupling in clean environment
- Discovery potential up to $\sqrt{s}/2$
 - Sensitivity complementary to LHC and Dark Matter Experiments
 - Benefits from large lever arm in energy
 - Polarisation to suppress SM background
- The hardest part is to keep all the promises
 - Requires large effort on experimental and theory side and excellent communication

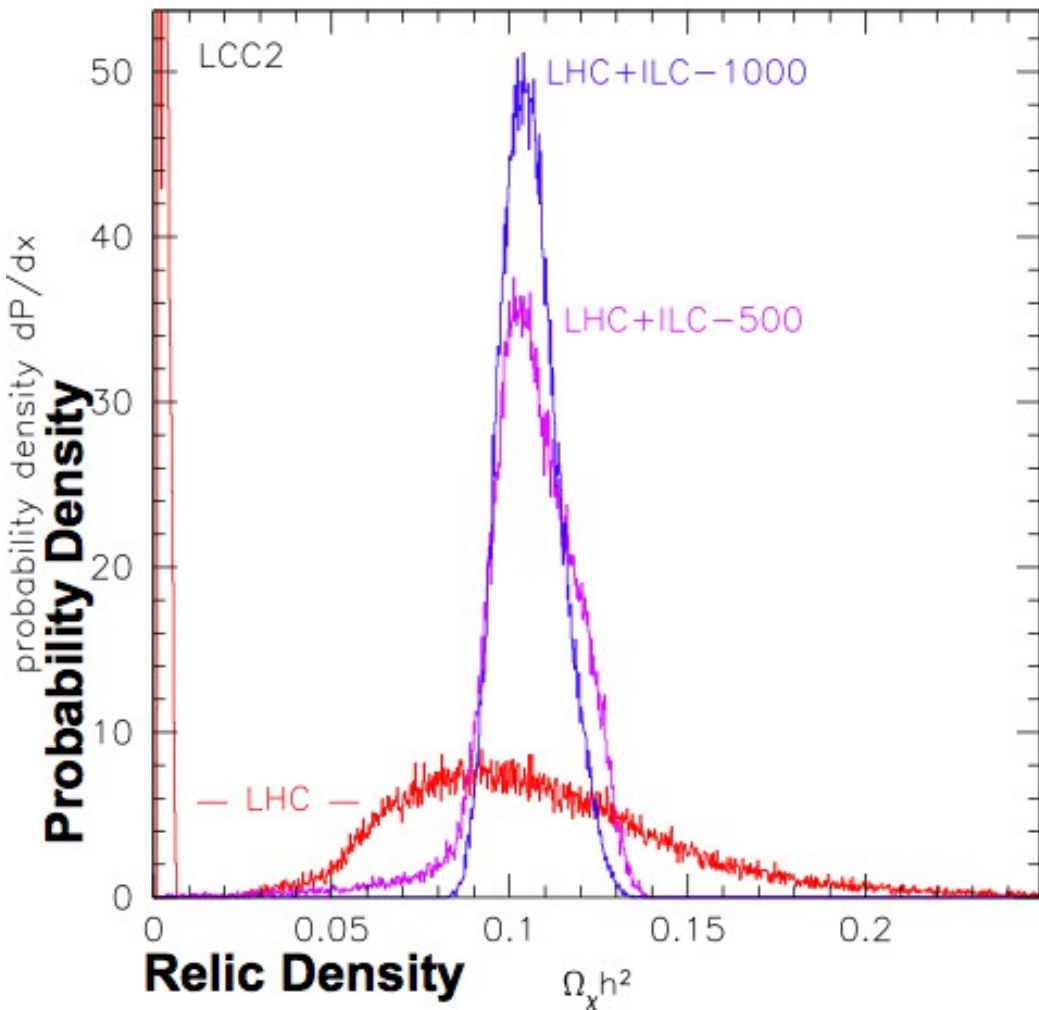
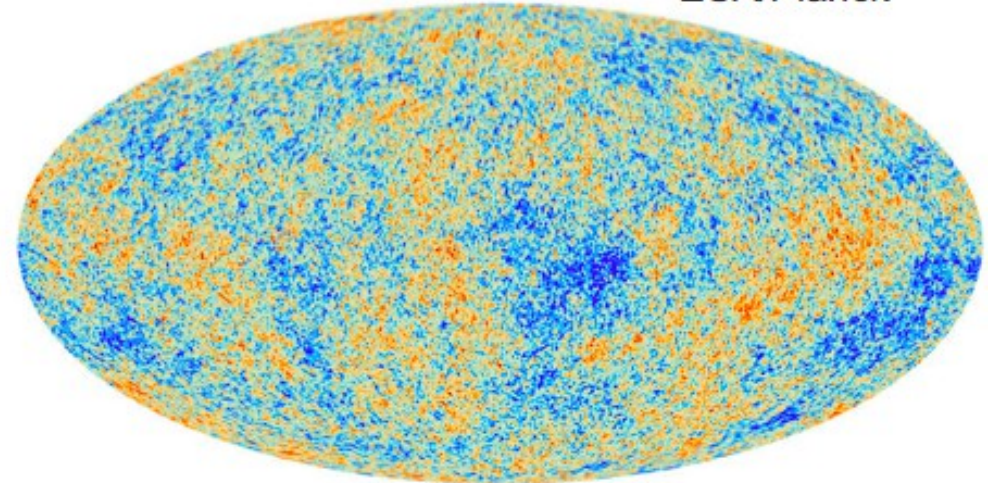
- Versatile machines for precision physics in the range $m_e - 3 \text{ TeV}$ under study
 - Linear and circular machines would/could be complementary
- LEP/SLC programme can be repeated with much higher precision
- Which precision is reasonable precision?
- (Higgs and) Top quark are physics guaranteed
 - (My conviction) both are messengers to New Physics
 - Electroweak top quark pair production
 - $t\bar{t}h$ coupling in clean environment
- Discovery potential up to $\sqrt{s}/2$
 - Sensitivity complementary to LHC and Dark Matter Experiments
 - Benefits from large lever arm in energy
 - Polarisation to suppress SM background
- The hardest part is to keep all the promises
 - Requires large effort on experimental and theory side and excellent communication

Rencontres de Vietnam on e- precision measurements?

Backup

WMAP/Planck (68% CL)
 $\Omega_c h^2 = 0.1196 \pm 0.0027$

ESA/Planck



Once a DM candidate is discovered, need to check the consistency with the measured DM relic abundance.

→ ILC's precise measurements of the mass and cross sections provide crucial input.

Baltz, Battaglia, Peskin, Wizansky
 PRD74 (2006) 103521, arXiv:hep-ph/0602187

Tomohiko Tanabe
 ILD Meeting 2014

Current uncertainties for EWPOs

6/19

	Experiment	Theory error	Main source
M_W	80.385 ± 0.015 MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
Γ_Z	2495.2 ± 2.3 MeV	0.5 MeV	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
σ_{had}^0	41540 ± 37 pb	6 pb	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$R_b \equiv \Gamma_Z^b / \Gamma_Z^{\text{had}}$	0.21629 ± 0.00066	0.00015	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^l$	0.23153 ± 0.00016	4.5×10^{-5}	$\alpha^3, \alpha^2\alpha_s$

Methods for theory error estimates:

- Parametric factors, *i. e.* factors of α, N_c, N_f, \dots
- Geometric progression, *e. g.* $\frac{\mathcal{O}(\alpha^3)}{\mathcal{O}(\alpha^2)} \sim \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)}$
- Renormalization scale dependence (often underestimates error)
- Renormalization scheme dependence (may underestimate error)

Future projections

7/19

	ILC	FCC-ee	perturb. error with 3-loop [†]	Param. error ILC*	Param. error FCC-ee**
M_W [MeV]	3–5	~ 1	1	2.6	1
Γ_Z [MeV]	~ 1	~ 0.1	$\lesssim 0.2$	0.5	0.06
R_b [10^{-5}]	15	$\lesssim 5$	5–10	< 1	< 1
$\sin^2 \theta_{\text{eff}}^l$ [10^{-5}]	1.3	0.6	1.5	2	2

[†] **Theory scenario:** $\mathcal{O}(\alpha\alpha_S^2)$, $\mathcal{O}(N_f\alpha^2\alpha_S)$, $\mathcal{O}(N_f^2\alpha^2\alpha_S)$
 (N_f^n = at least n closed fermion loops)

Parametric inputs:

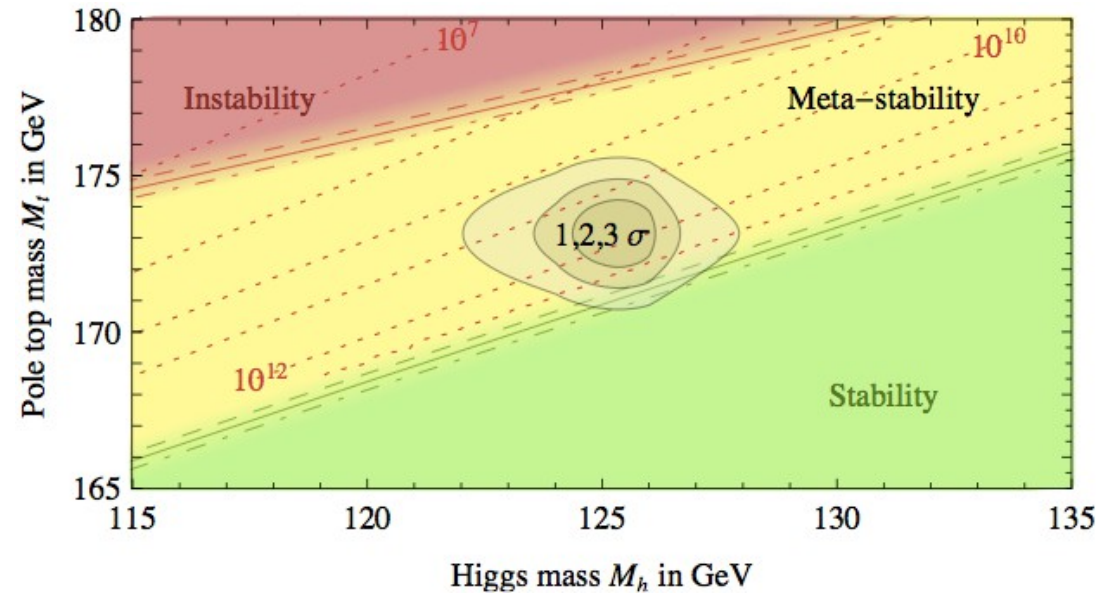
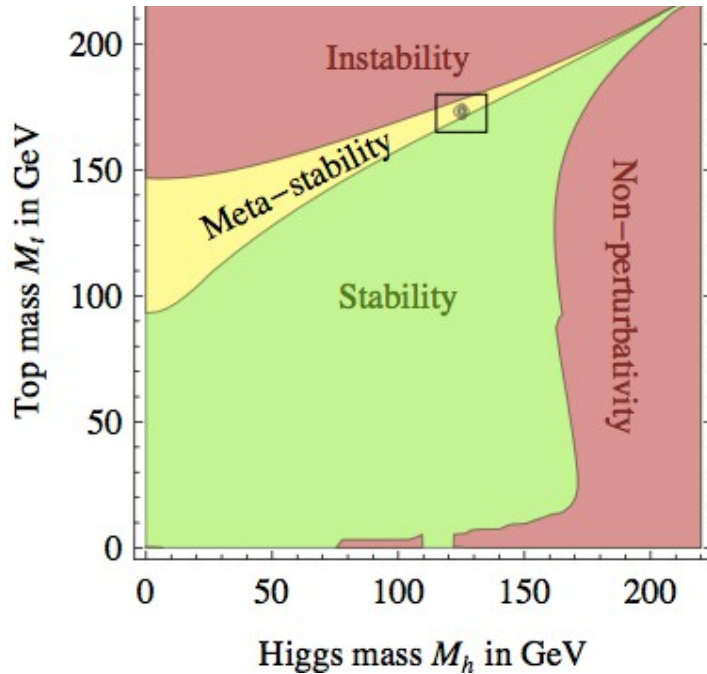
* **ILC:** $\delta m_t = 100$ MeV, $\delta\alpha_S = 0.001$, $\delta M_Z = 2.1$ MeV

****FCC-ee:** $\delta m_t = 50$ MeV, $\delta\alpha_S = 0.0001$, $\delta M_Z = 0.1$ MeV

also: $\delta(\Delta\alpha) \sim 5 \times 10^{-5}$

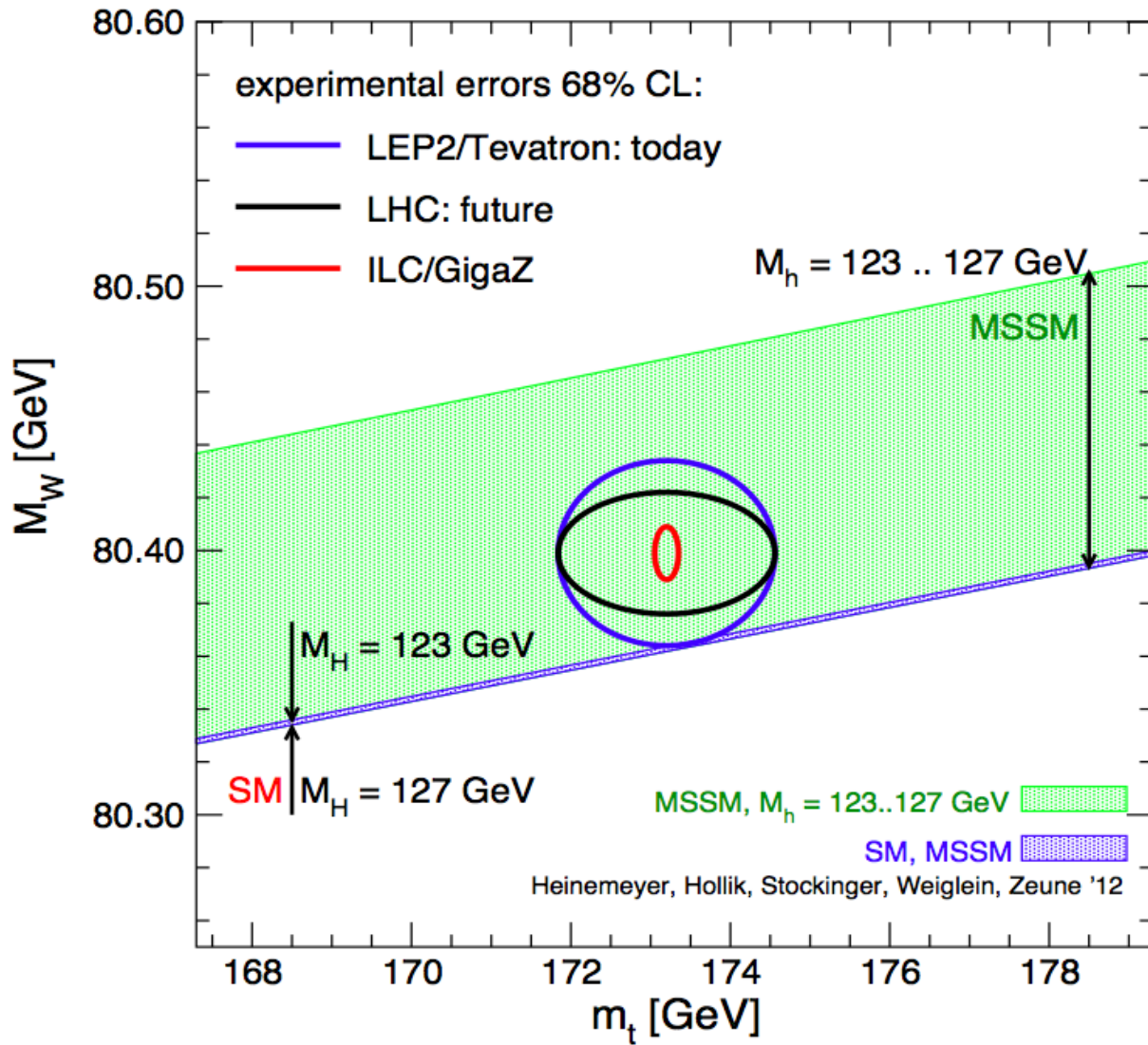


$$M_h \text{ [GeV]} > 129.4 + 1.4 \left(\frac{M_t \text{ [GeV]} - 173.1}{0.7} \right) - 0.5 \left(\frac{\alpha_s(M_Z) - 0.1184}{0.0007} \right) \pm 1.0_{\text{th}} .$$



Type of error	Estimate of the error	Impact on M_h
M_t	experimental uncertainty in M_t	± 1.4 GeV
α_s	experimental uncertainty in α_s	± 0.5 GeV
Experiment	Total combined in quadrature	± 1.5 GeV
λ	scale variation in λ	± 0.7 GeV
y_t	$\mathcal{O}(\Lambda_{\text{QCD}})$ correction to M_t	± 0.6 GeV
y_t	QCD threshold at 4 loops	± 0.3 GeV
RGE	EW at 3 loops + QCD at 4 loops	± 0.2 GeV
Theory	Total combined in quadrature	± 1.0 GeV

Uncertainty on (pole) top quark mass dominates uncertainty on stability conditions



Precise Top (and W) mass crucial to test compatibility of measured Higgs mass

MS might not be sufficient to explain Higgs mass

LHC may not reach sufficient discriminative power

A lepton collider will

Coronation of the Standard Model
and
First step on a road yet largely unexplored
Slightly modified citation of Barbieri arXiv:1309.3447



HIGGS

Particle Physics

HIGGS

Chip Brook, Snowmass Summary Talk

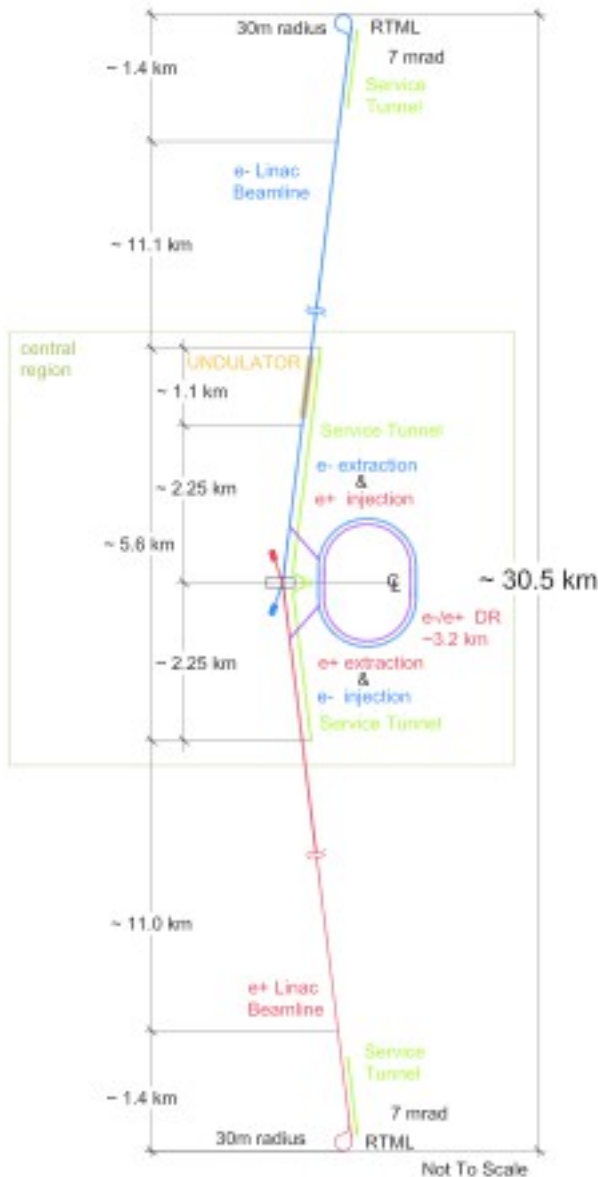
Where do we go from here?

ILC design parameters	
\sqrt{s}	91-500 GeV
\mathcal{L}	$2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
P_{e^-}	>80%
P_{e^+}	upto 30%
Length	~31 km

Comment

- 500 GeV is baseline
Option to upgrade to 1 TeV
- ~Factor 4 technically possible
- Proven by SLC
- ~Conservative estimate
- Current site allows for 50km

- Discussion on possible running scenarios has started
- Luminosity and running time to achieve at a ~25 years research programme
That includes running at 250 GeV, 350 GeV, 500 GeV and 1 TeV
- No official statement yet but integrated luminosities indicated in following transparencies are realistic



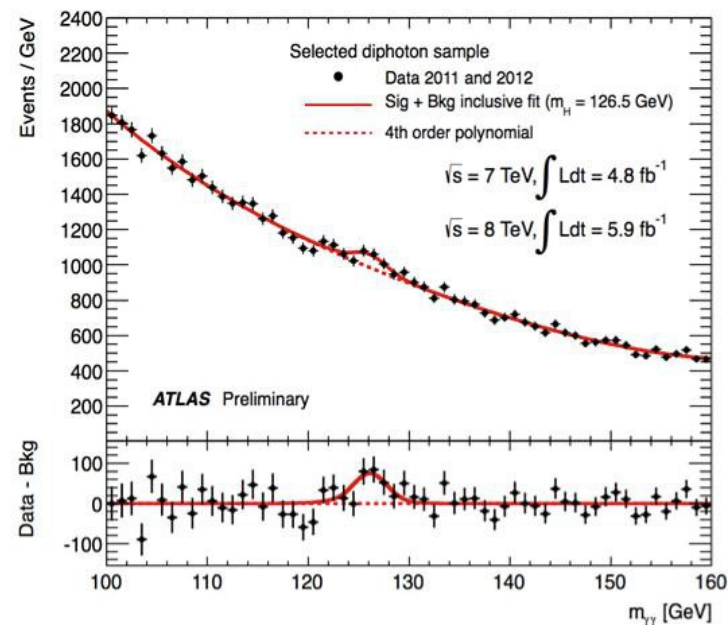
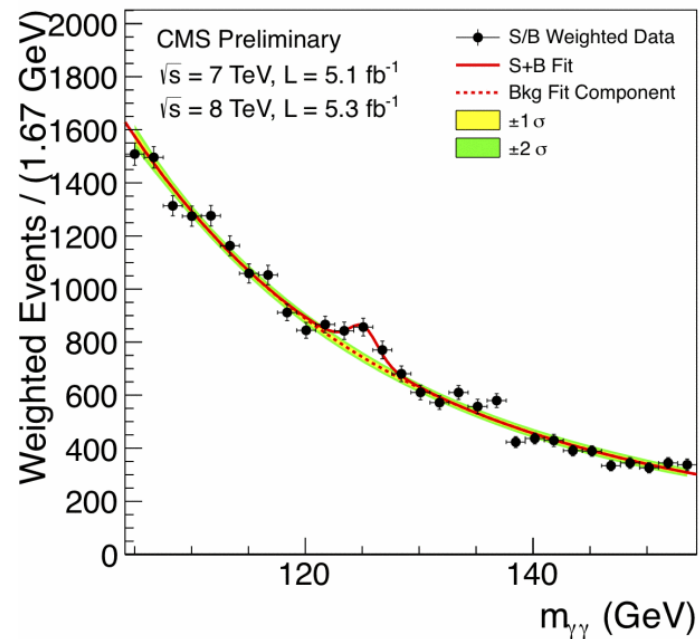
• SCRF Technology

- 1.3GHz SCRF with 31.5 MV/m
- 17,000 cavities
- 1,700 cryomodules
- 2×11 km linacs

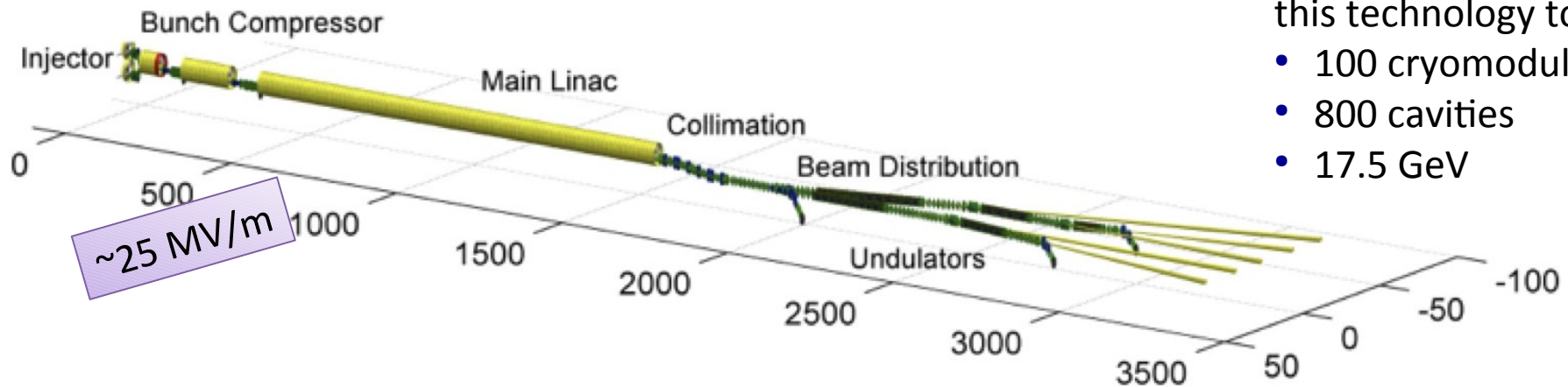
Luminosity

$$L \propto \frac{\eta_{RF} P_{RF}}{E_{cm}} \sqrt{\frac{\delta_{BS}}{\epsilon_{n,y}}} H_D$$

$\eta_{RF} \sim 40\%$ for SCRF technology
-> efficient technology



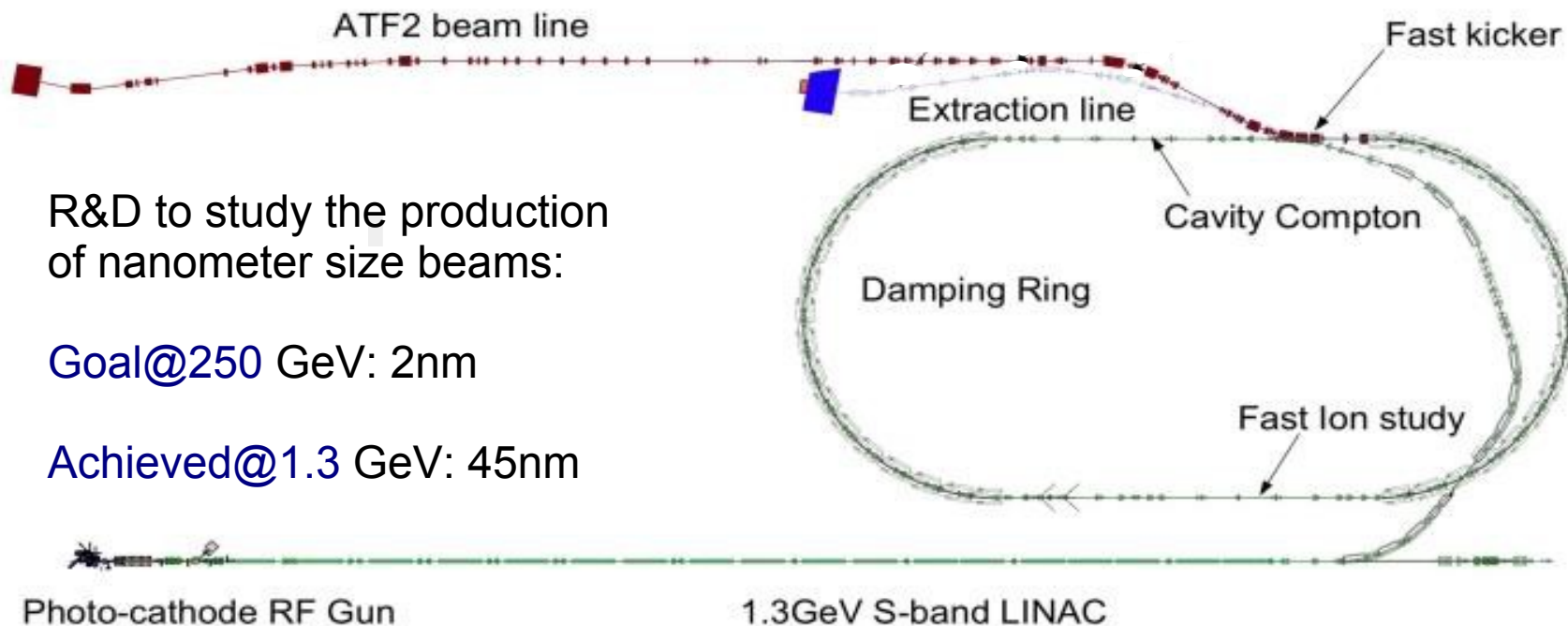
European XFEL Project: Location DESY Hamburg, Start 2015



Largest deployment of this technology to date

- 100 cryomodules
- 800 cavities
- 17.5 GeV

ATF at KEK Japan:

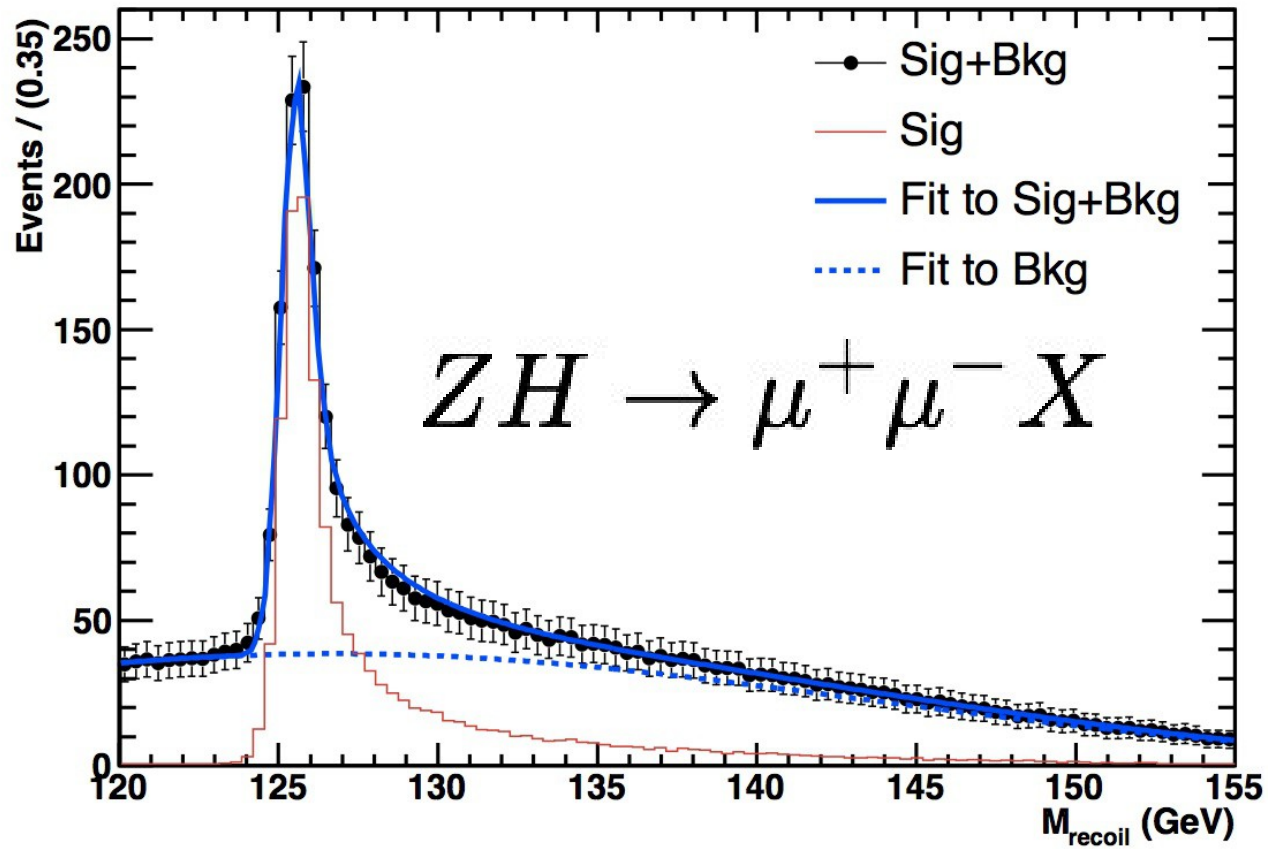


R&D to study the production of nanometer size beams:

Goal@250 GeV: 2nm

Achieved@1.3 GeV: 45nm

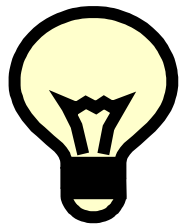
2. Higgs Physics at the ILC



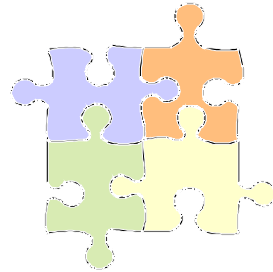
$$M_h = 125.3 \pm 0.03 \text{ GeV}$$

$$\sigma_{ZH} = 10.32 \pm 0.37 \text{ fb, } 3.6\%$$

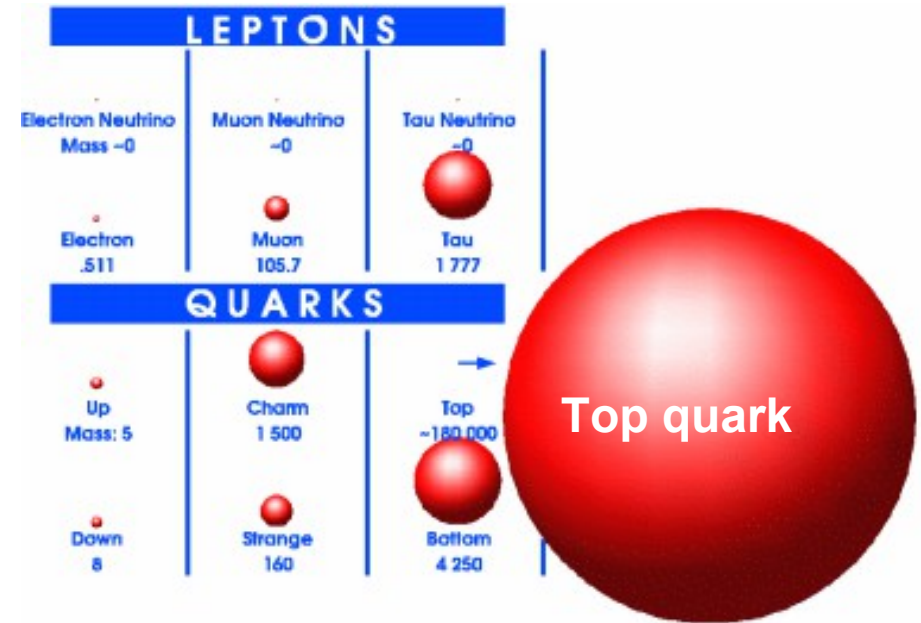
More on Higgs Physics, see talk of Junping



Elementary Scalar?



Composite object?



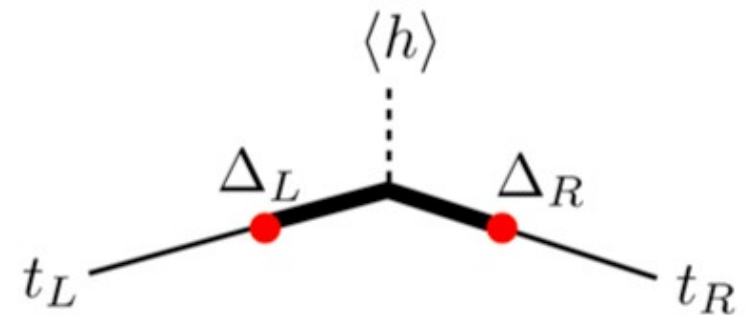
- Higgs and top quark are intimately coupled!

Top Yukawa coupling $O(1)$!

=> Top mass important SM Parameter

- New physics by compositeness?

Higgs and top composite objects?



Courtesy of S. Rychkov

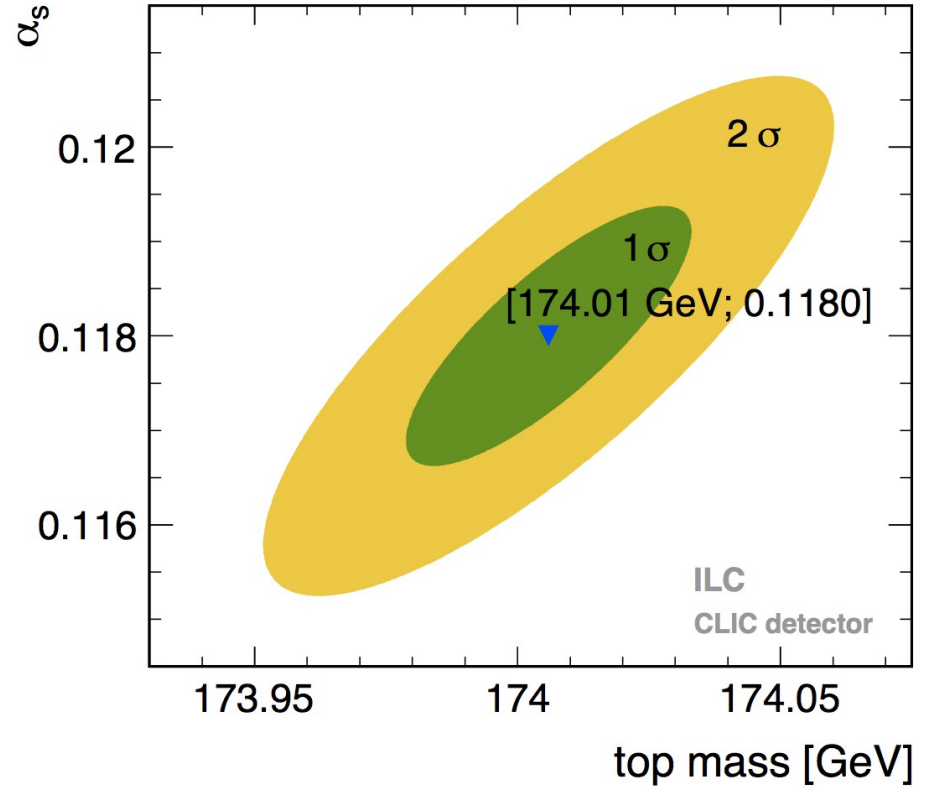
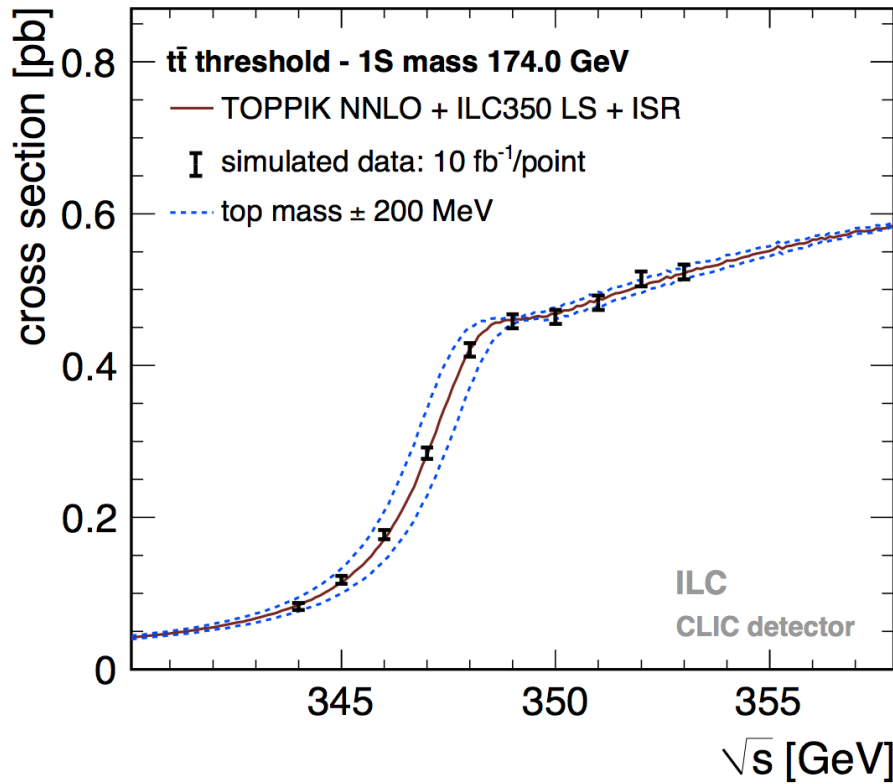
- **e+e- collider perfectly suited to decipher both particles**

The higher the energy the better!!!



Mass and α_s

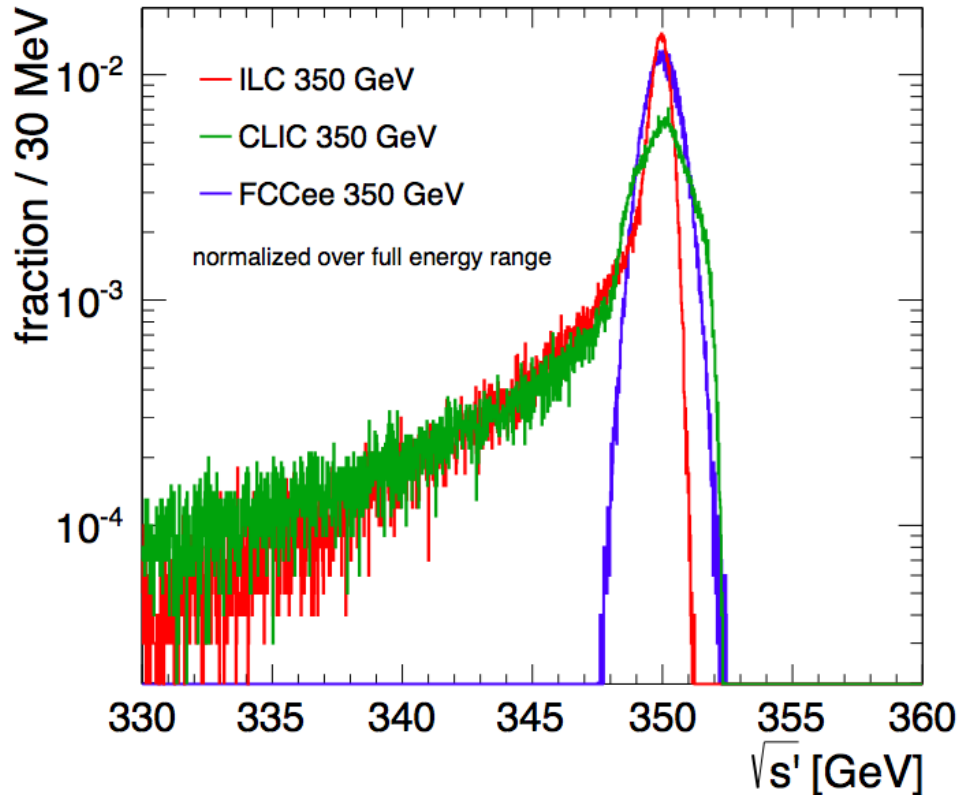
EPJC C73 (2013) 2530



~100 MeV

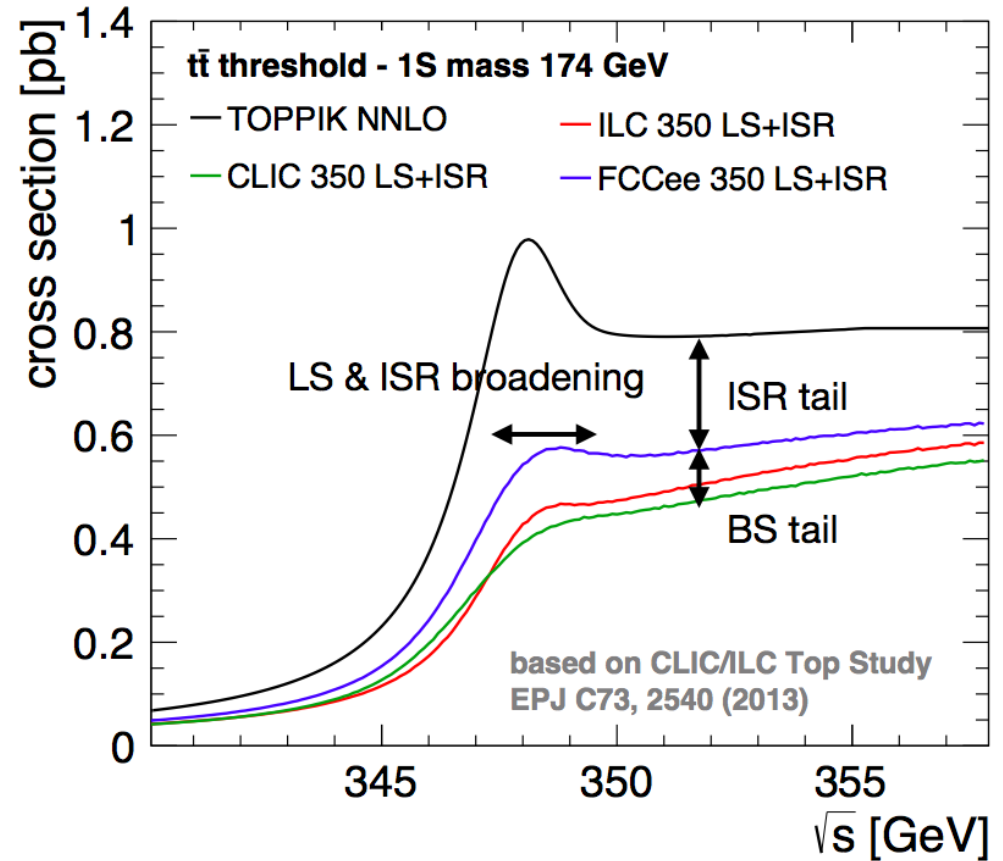
1S top mass and α_s combined 2D fit

m_t stat. error	27 MeV
m_t theory syst. (1%/3%)	5 MeV / 9 MeV
α_s stat. error	0.0008
α_s theory syst. (1%/3%)	0.0007 / 0.0022



- Slight changes in statistics due to cross section, changes in sensitivity due to steepness of threshold turn on
- For 100 fb⁻¹, no polarisation, 1D mass fit
15 MeV → 18 MeV → 21 MeV (stat.)
FCCee ILC CLIC

- Somewhat different luminosity spectra for different machines
- No beamstrahlung in storage ring
- Sharper main peak at ILC broader for CLIC



Higgs Mechanism

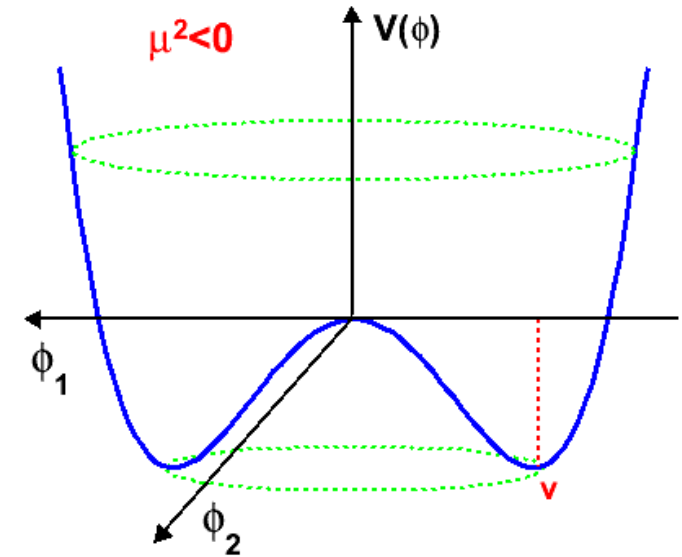
Scalar field which doesn't vanish in the vacuum

Choice in SM:
Doublet Field

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} \quad \text{4 degrees of freedom}$$

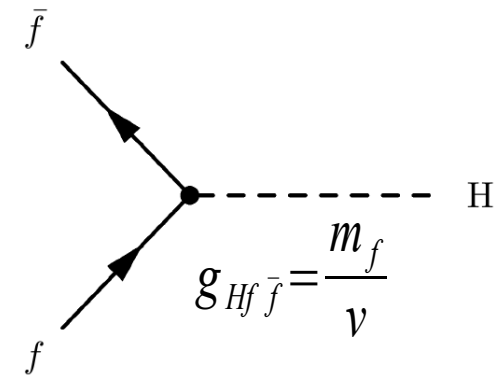
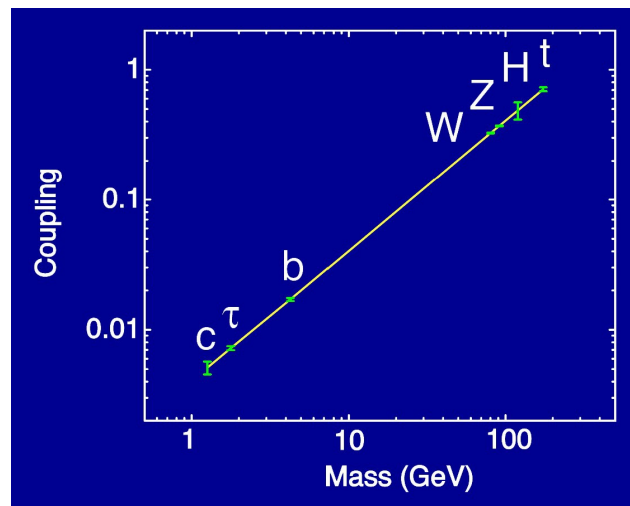
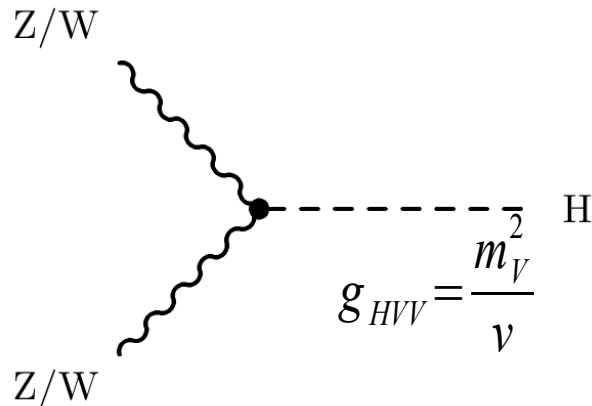
Higgs Boson

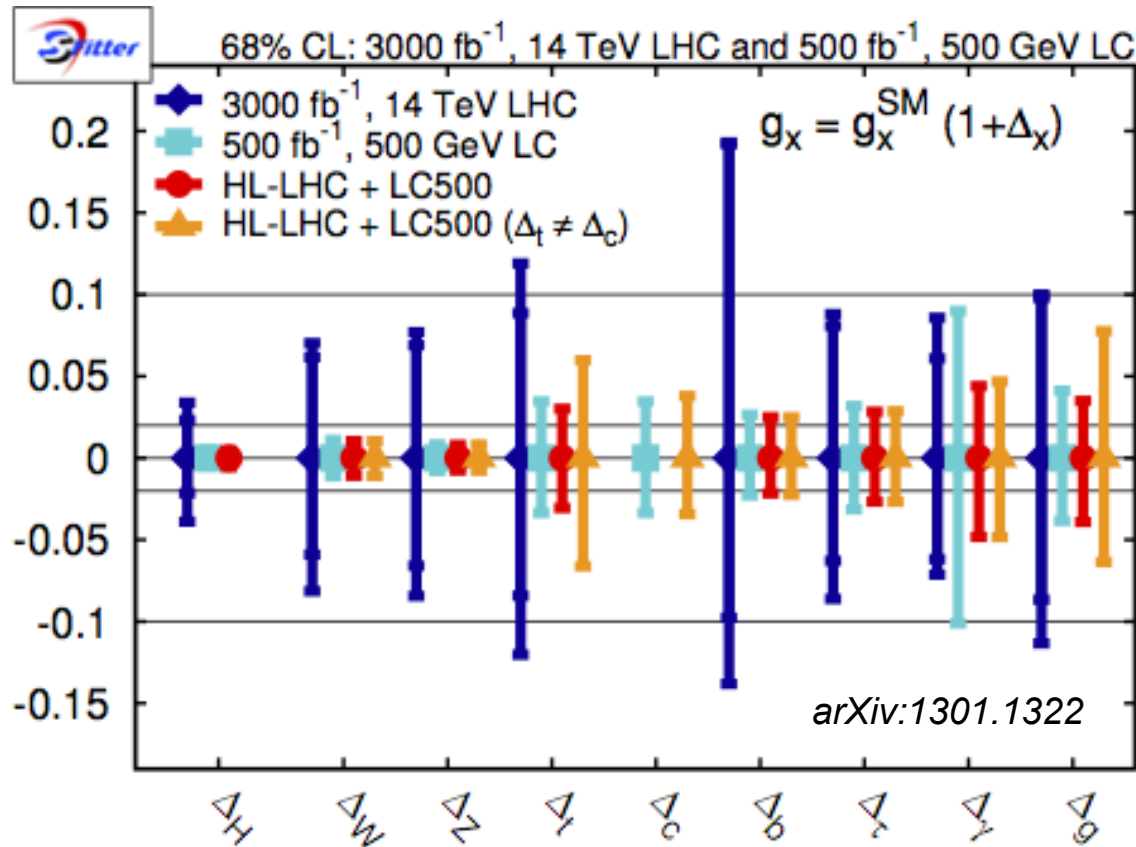
Longitudinally degrees
of W,Z Bosons



Couplings to Higgs Boson in Standard Model

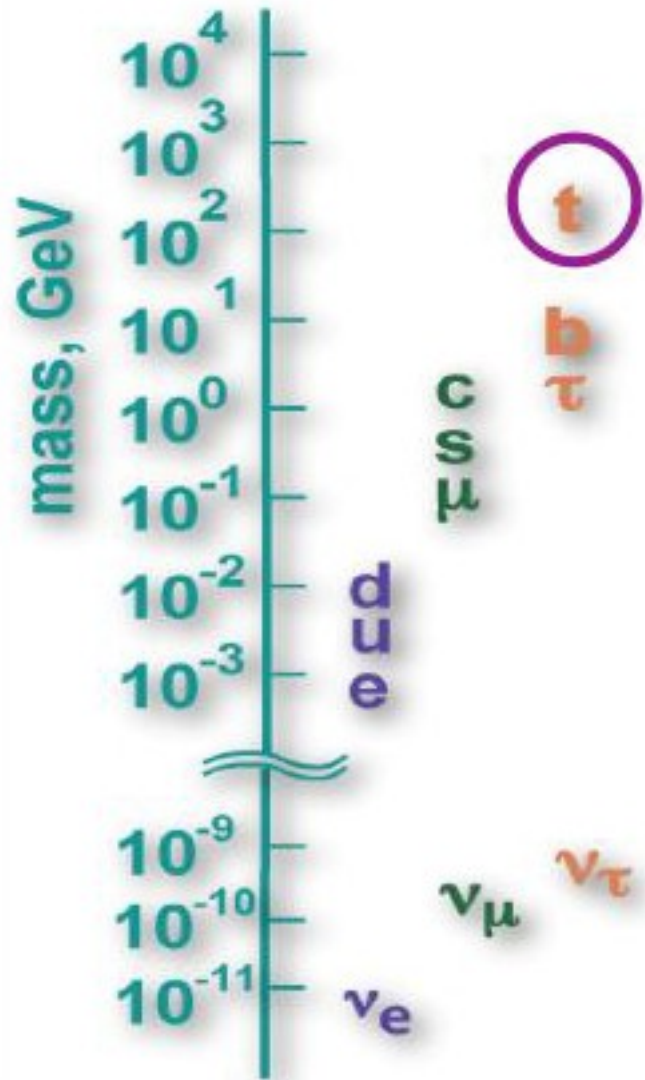
Increase with particle mass





- A e+e- machine (Linear Collider) running at several energies will provide precise measurements of relevant Higgs couplings: Possibility to confirm the Higgs mechanism of the SM
- Precision matters: Detect deviations, for example due to extended Higgs sectors (SUSY, composite, ...): Expected on the 10% - 15% level in fermions, on the few % level in gauge bosons in typical Two-Higgs-Doublet models

- Expected statistical uncertainty **10 – 30 MeV**
- Experimental systematics
 - Beam energy: **~30 MeV** or lower
 - Non-ttbar background, selection efficiencies: **~ 15 MeV**
 - Luminosity spectrum: **10 MeV**
 - Single top contamination: **< 30 MeV**
- Theory uncertainties
 - **~40-45 MeV** due to scale uncertainties in NNNLO calculations
 - When not included in the fit: $\sim 3 \text{ MeV per } 10^{-4} \text{ uncertainty on } \alpha_s \text{ today} \rightarrow \text{~18 MeV}$
 - Conversion from 1S/PS masses to MSbar mass Currently: **~50 MeV**
However conversion now known to N⁴LO (arxiv:1502.01030)
 - Now at point where results become sensitive to effects other than QCD

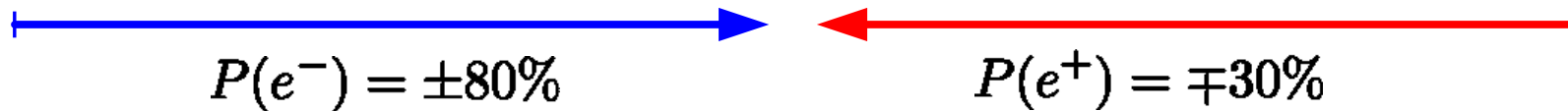


- SM does not provides no explanation for mass spectrum of fermions (and gauge bosons)
- Fermion mass generation closely related to the origin electroweak symmetry breaking
- Expect residual effects for particles with masses closest to symmetry breaking scale
 - A_{FB} anomaly at LEP for b quark

Strong motivation to study chiral structure of top vertex in high energy e^+e^- collisions

At ILC **no** separate access to ttZ or tty vertex, but ...

ILC 'provides' two beam polarisations



There exist a number of observables sensitive to chiral structure, e.g.

$$\sigma_I \quad A_{FB,I}^t = \frac{N(\cos\theta > 0) - N(\cos\theta < 0)}{N(\cos\theta > 0) + N(\cos\theta < 0)} \quad (F_R)_I = \frac{(\sigma_{tR})_I}{\sigma_I}$$

x-section

Forward backward asymmetry

Fraction of right handed top quarks



Extraction of relevant unknowns

$$F_{1V}^\gamma, F_{1V}^Z, F_{1A}^\gamma = 0, F_{1A}^Z \quad \text{or equivalently} \quad g_L^\gamma, g_R^\gamma, g_L^Z, g_R^Z$$

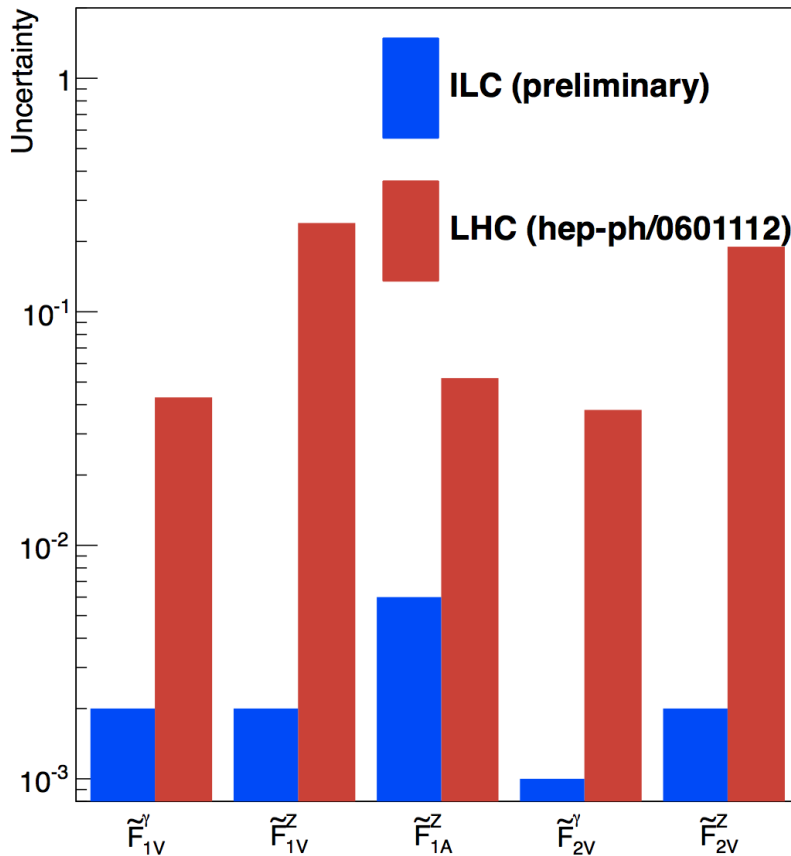
$$F_{2V}^\gamma, F_{2V}^Z$$

Precision: cross section $\sim 0.5\%$,

Precision $A_{FB} \sim 2\%$,

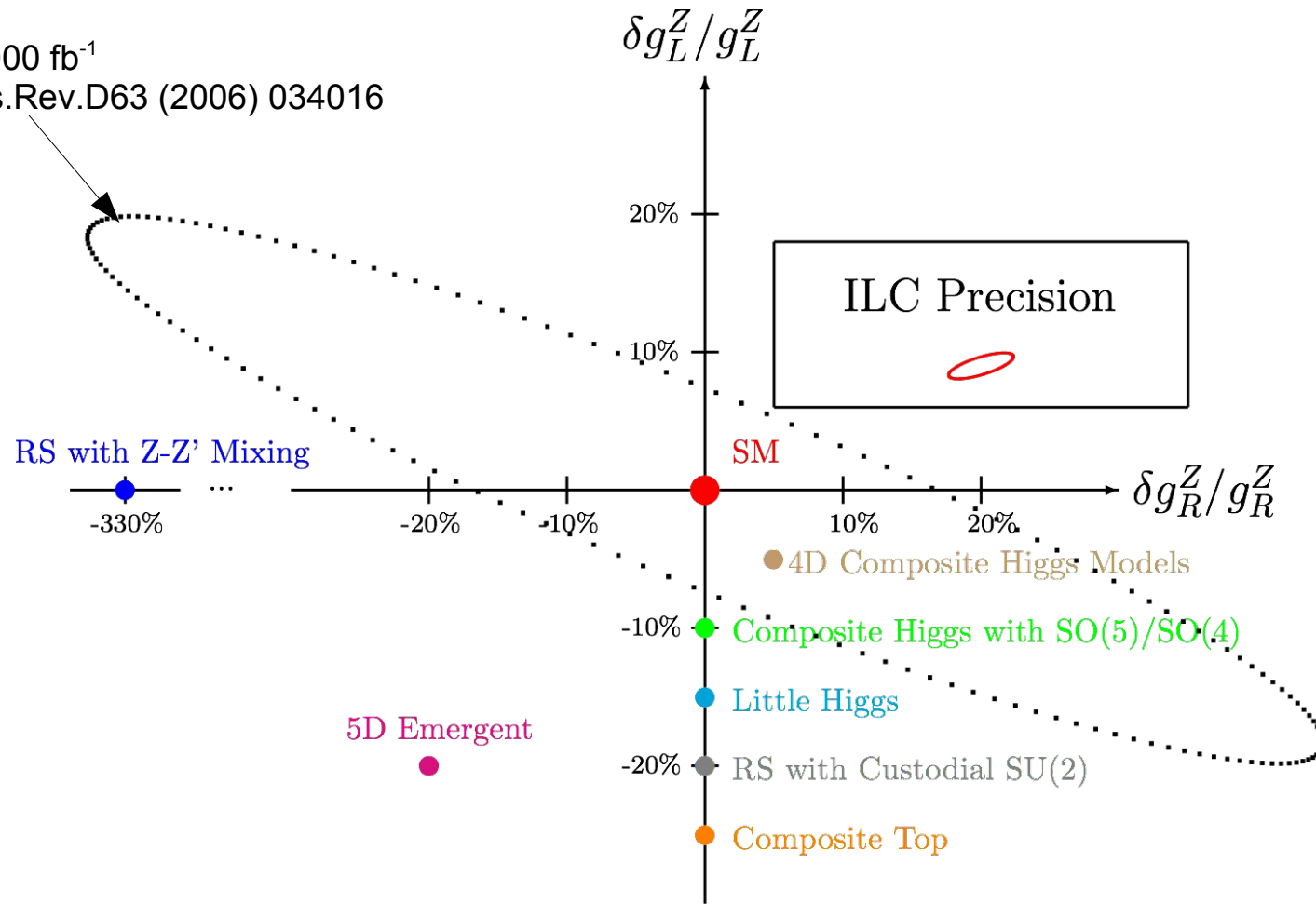
Precision $\lambda_t \sim 3-4\%$

Accuracy on CP conserving couplings



- ILC might be up to two orders of magnitude more precise than LHC ($\sqrt{s} = 14$ TeV, 300 fb^{-1})
Disentangling of couplings for ILC
One variable at a time For LHC
However LHC projections from 8 years old study
- Need to control experimental (e.g. Top angle) and theoretical uncertainties (e.g. Electroweak corrections)
-> Dedicated work has started
- Potential for CP violating couplings at ILC under study

ILC promises to be high precision machine for electroweak top couplings

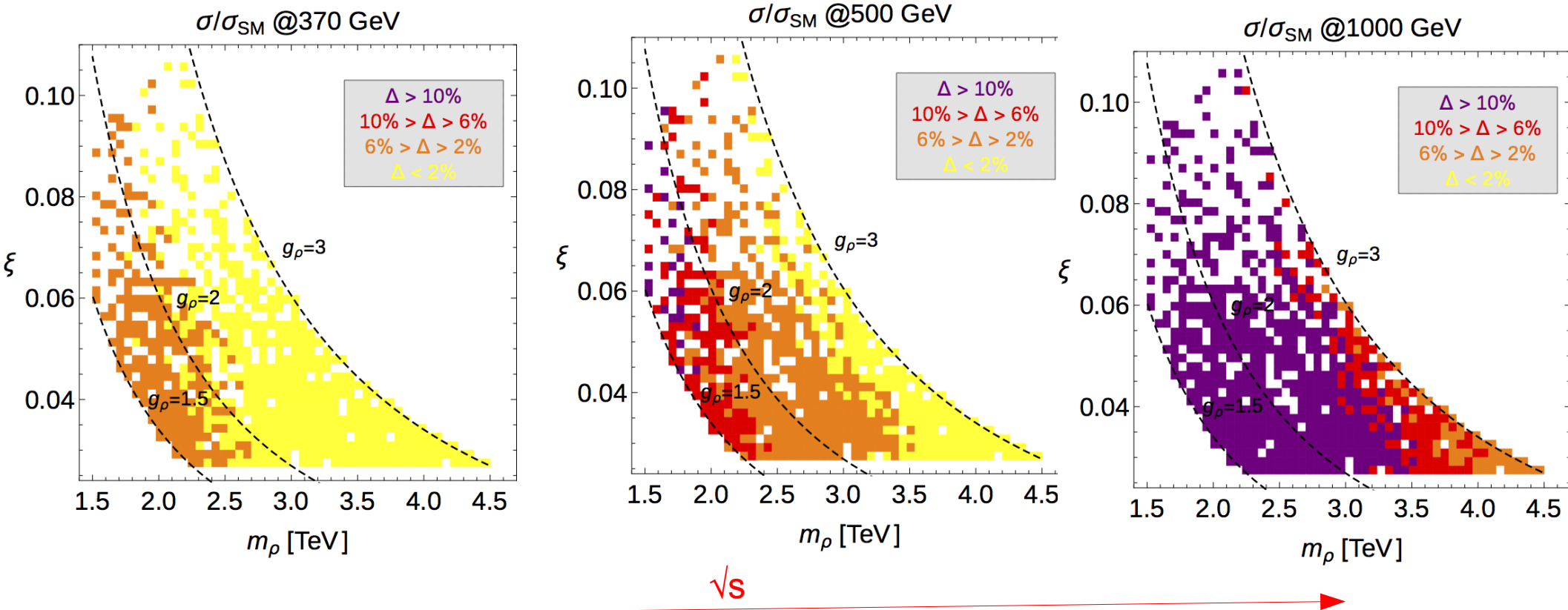


Linear Collider will outperform LHC results

- Particular poor constraint on g_R (this holds also for flavor physics results)
- LHC LO QCD analysis, ~30% improvement through NLO QCD
- LHC may still be capable to exclude models

Example: Sensitivity to $M_{Z'} = M_\rho$ in 4D Higgs Composite Model, arxiv: 1504.05407

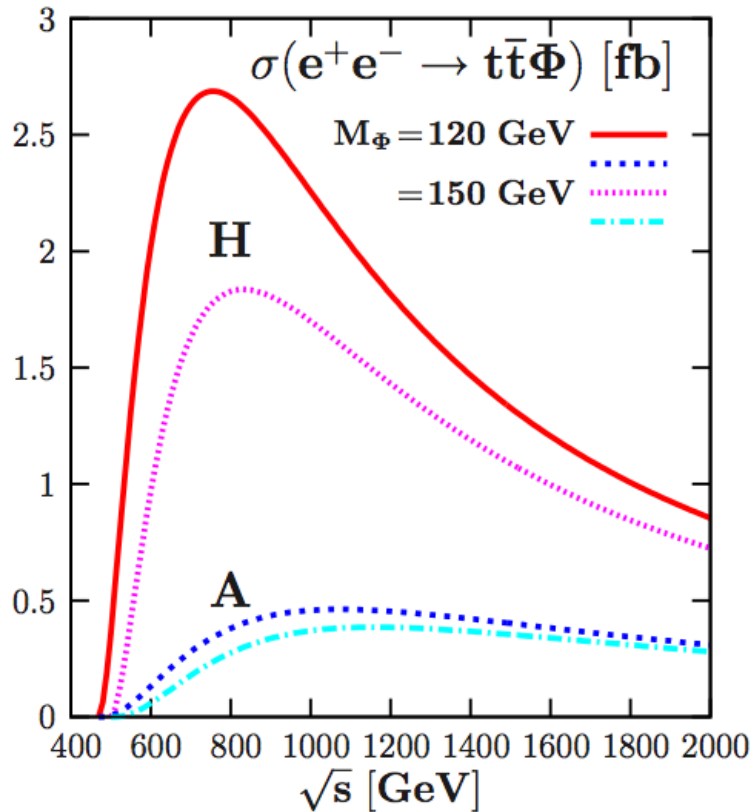
$$\frac{\delta g_I}{g_I} \sim \xi \sim \left(\frac{v g_\rho}{M_\rho} \right)^2$$



Effects observed at smaller energies may be amplified at higher energies

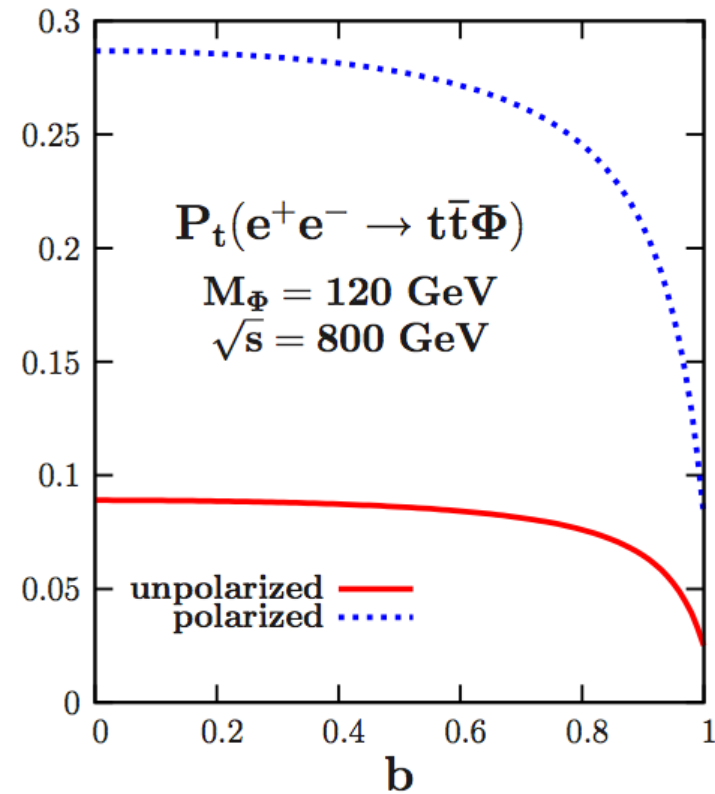
Direct coupling of top quark to CP odd and CP even scalar

Cross section



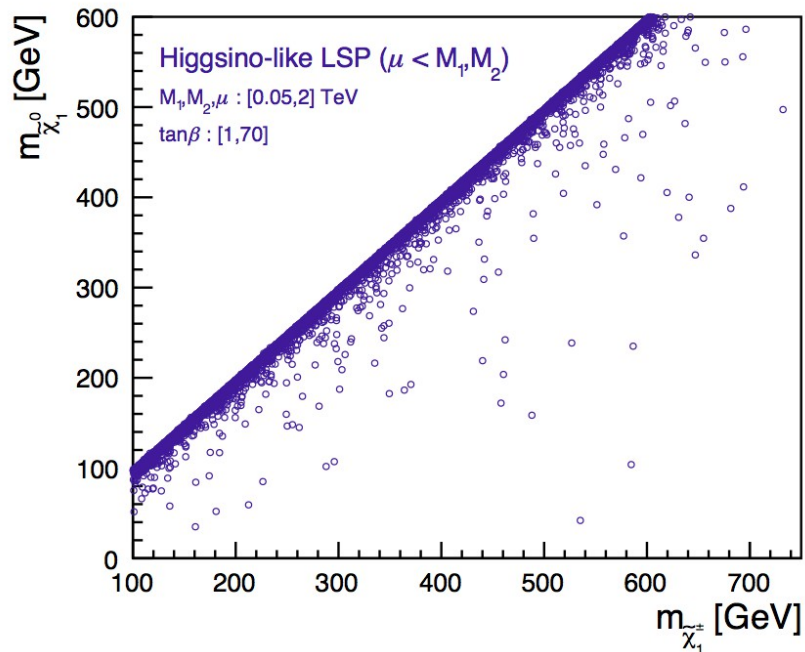
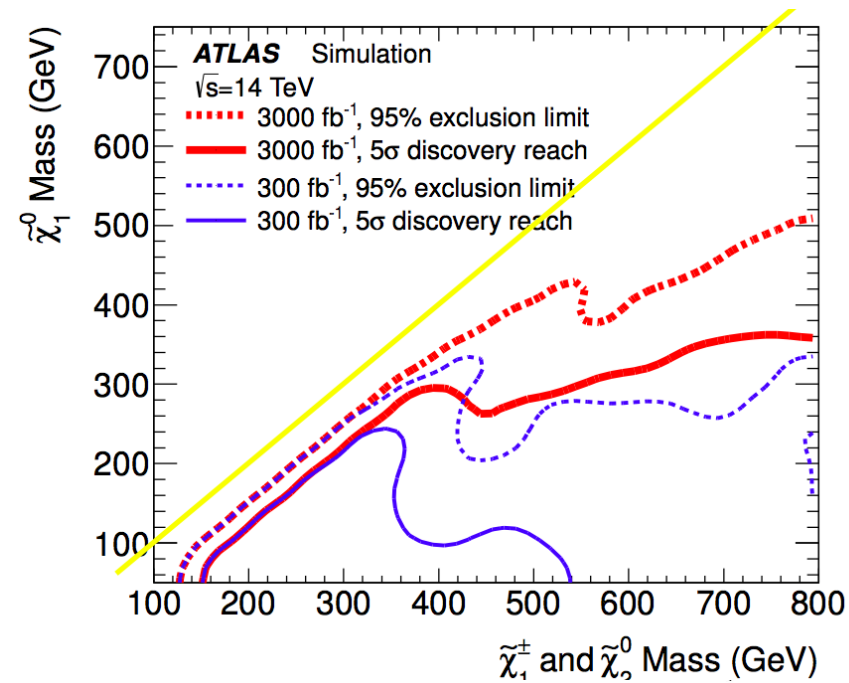
Dramatic differences for
CP odd and CP even scalar

Top quark polarisation



Sensitivity to CP odd admixture b
Merit of beam polarisation

Determination of CP nature of scalar boson in an unambiguous way

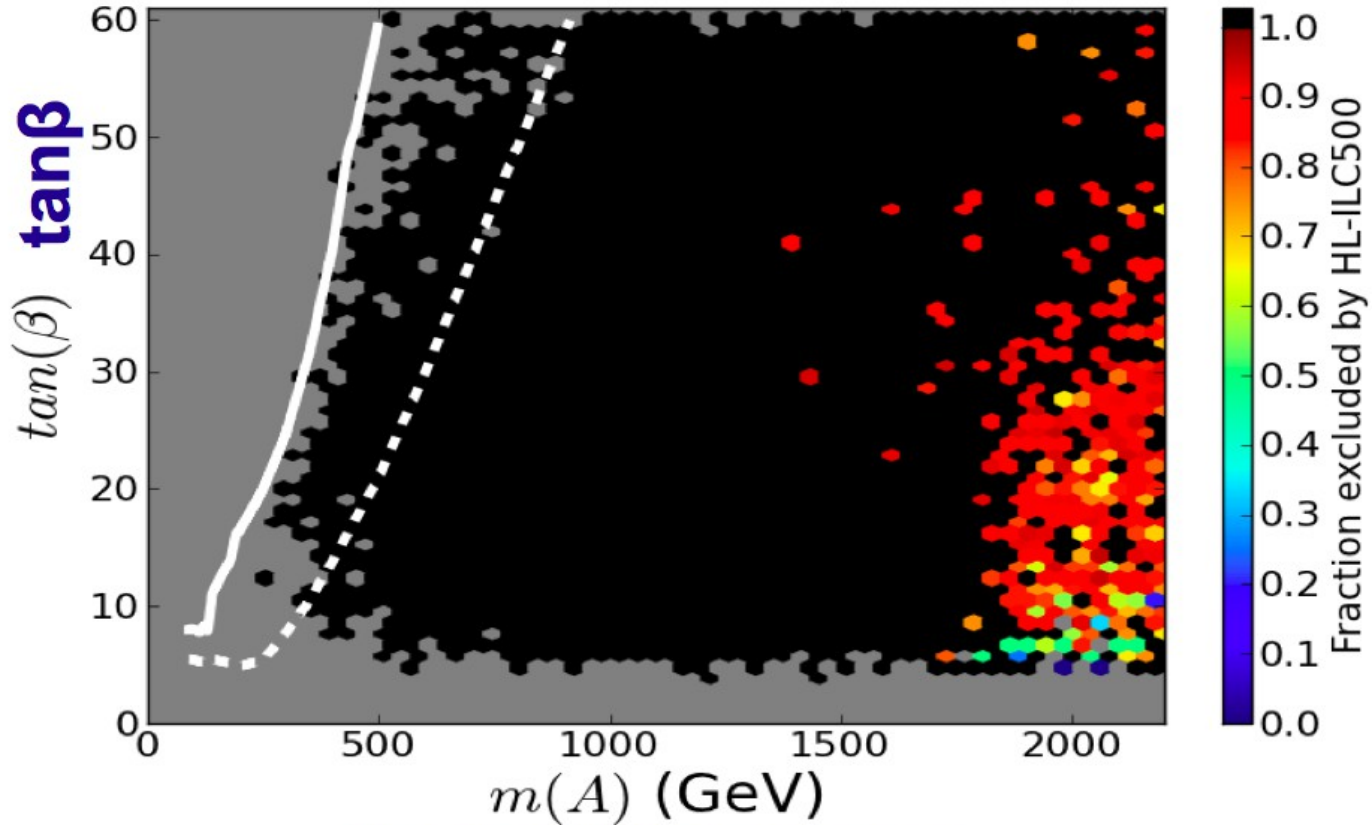
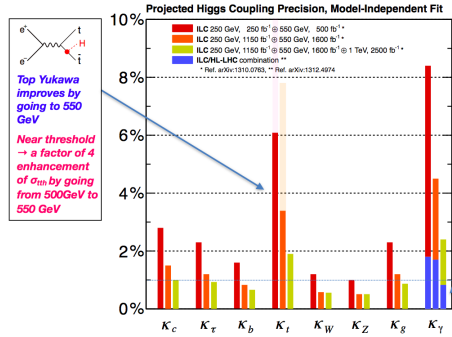


- Hadron colliders have a great potential to discover supersymmetric particles
 - coloured and neutral
- Hadron colliders cannot exclude low mass SUSY with light neutralino and chargino(s) Degenerated in mass
- Example: scenario with light higgsinos $\mu \sim O(v)$

Exclusion of pMSSM points via Higgs Couplings – arXiv 1407.7021

ILC (1150 fb⁻¹@250 GeV & 1600 fb⁻¹@500 GeV)

See Chapter 2



Heavy Higgs mass

Precision Higgs coupling measurements are sensitive probe for heavy Higgs Bosons $m_A \sim 2$ TeV reach for any $\tan\beta$ in high energy e+e- collisions

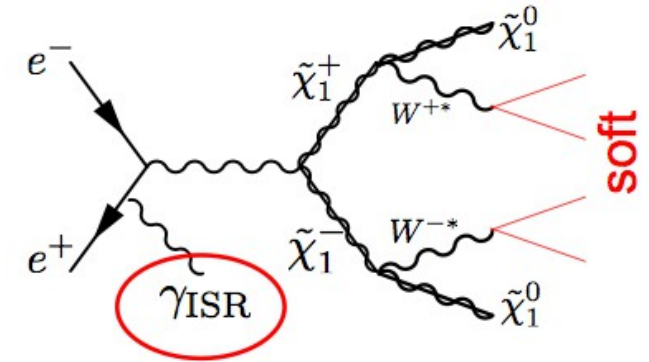
Study of Higgsino pair production, with ISR tag

Benchmark models with

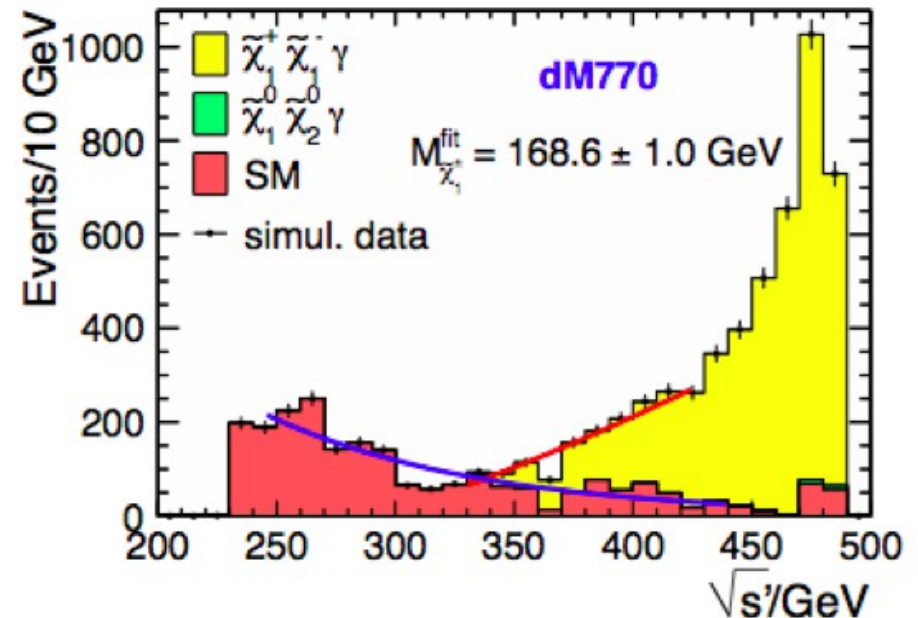
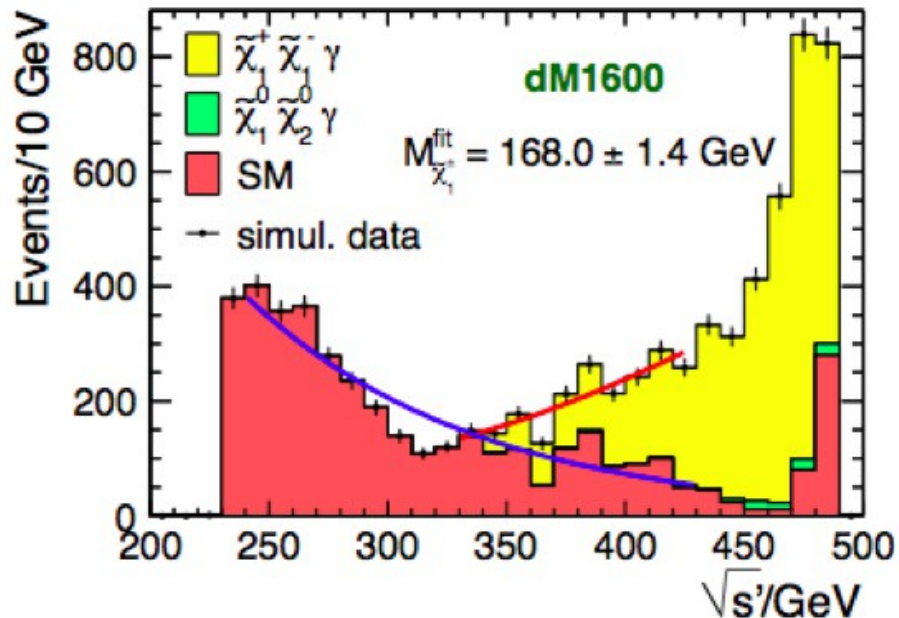
$m(\text{NLSP}) - M(\text{LSP}) = 1.6 \text{ GeV}$ and 0.8 GeV

$$\sigma(e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^-) = 78.7 \text{ (77.0) fb}$$

$$\Delta M = 1.60 \text{ (0.77) GeV}$$



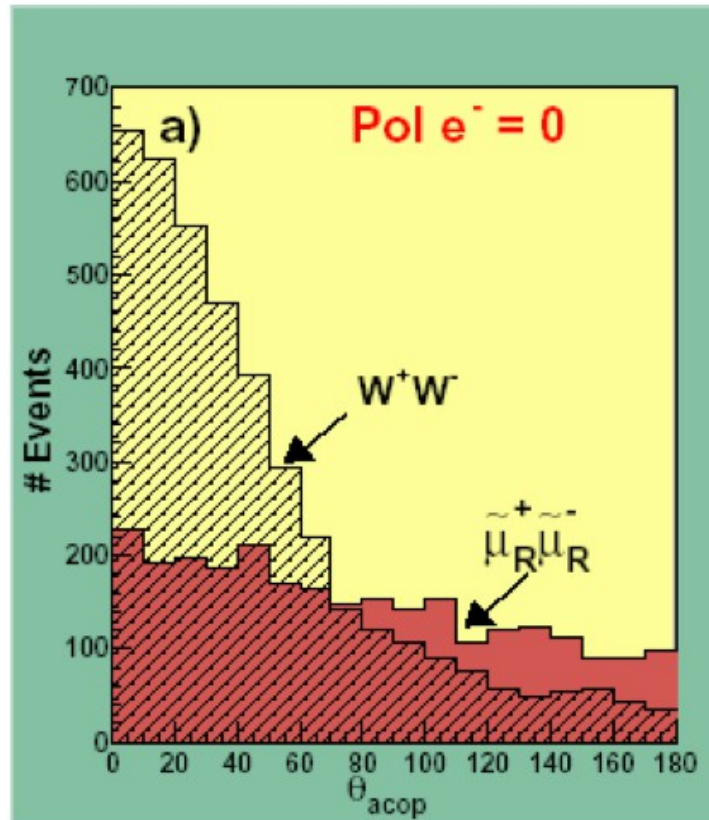
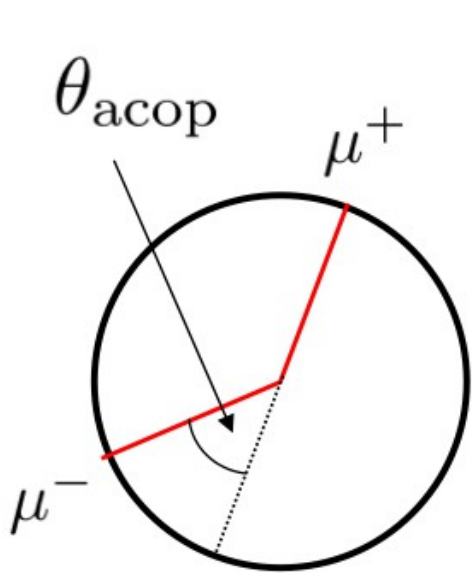
Berggren, Bruemmer, List, Moortgat-Pick, Robens, Rolbiecki, Sert, EPJ C73 (2013) 2660 [arXiv:1307.3566]



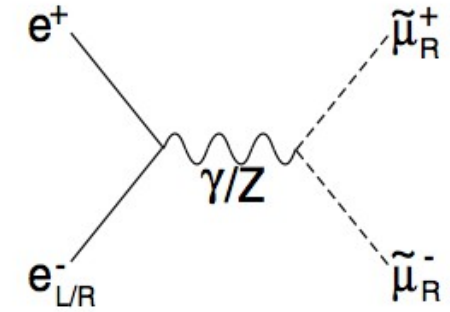
$\sqrt{s}=500 \text{ GeV}$, $\text{Lumi}=500 \text{ fb}^{-1}$, $P(e^-,e^+)=(-0.8,+0.3) \rightarrow \text{LSP mass resolution } \sim 1\%$

Clear signal => ILC covers important corner of phase space for SUSY Searches

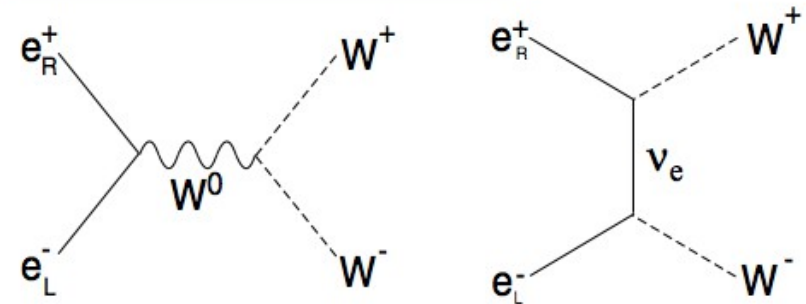
Example: Smuon pair production



signal



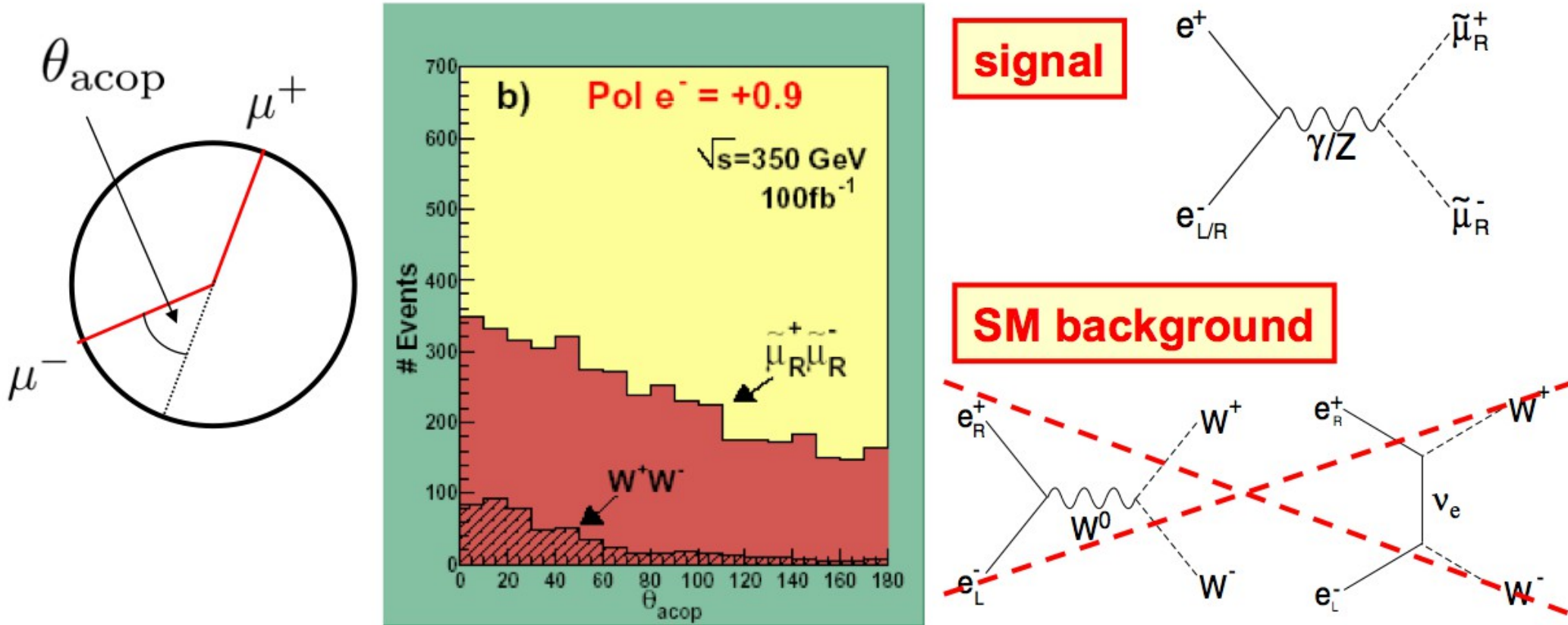
SM background



M. Thomson, IoP Meeting 2007

Strong SM Background

Example: Smuon pair production

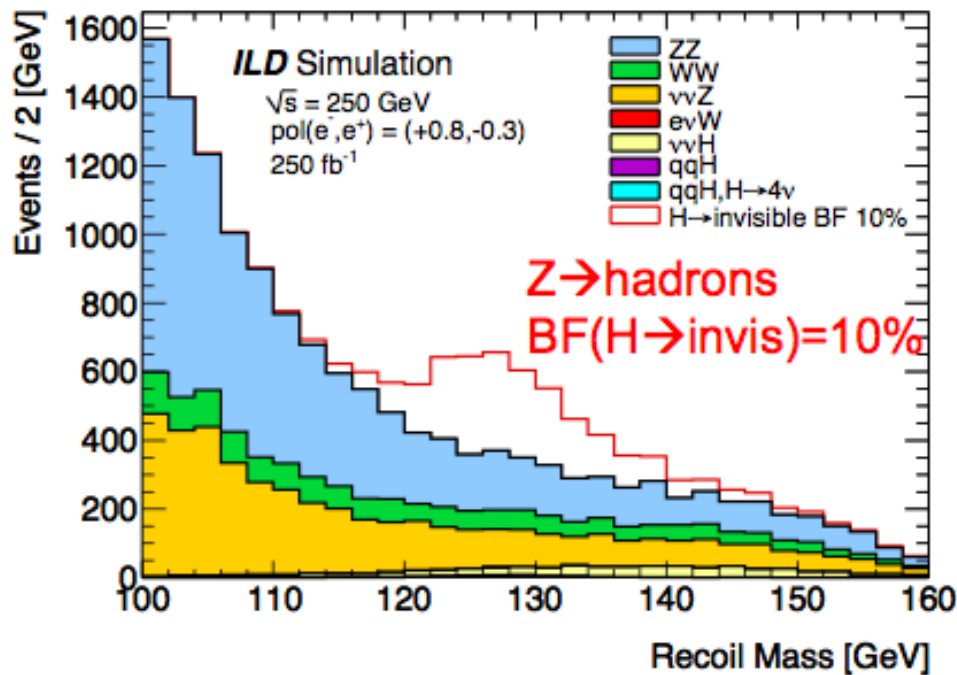


M. Thomson, IoP Meeting 2007

Strong background suppression through beam polarisation

WIMP searches at colliders are complementary to direct/indirect searches.
Examples at the ILC:

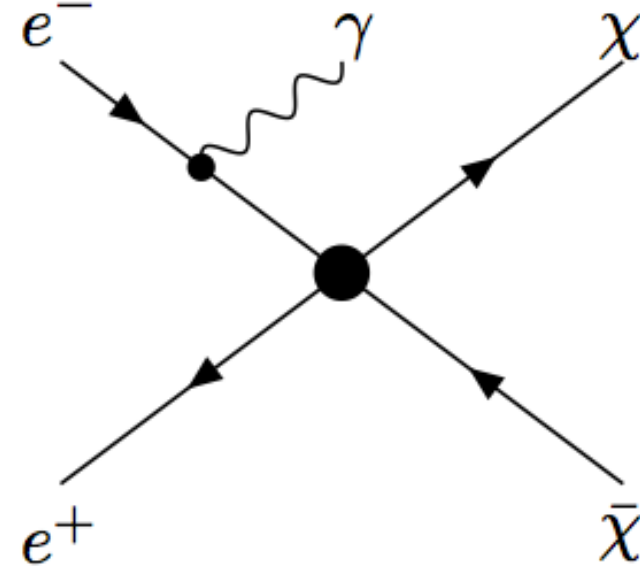
Higgs Invisible Decays



$BR(H \rightarrow \text{invis.}) < 0.4\%$ at 250 GeV, 1150 fb^{-1}

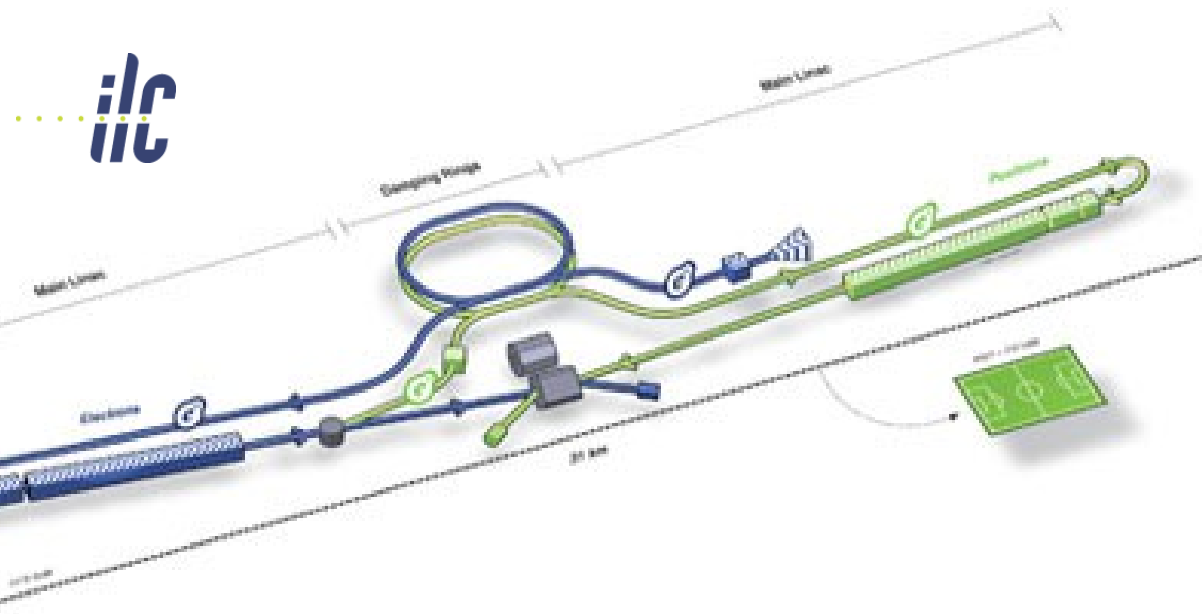
Impact of jet energy resolution

Monophoton Searches



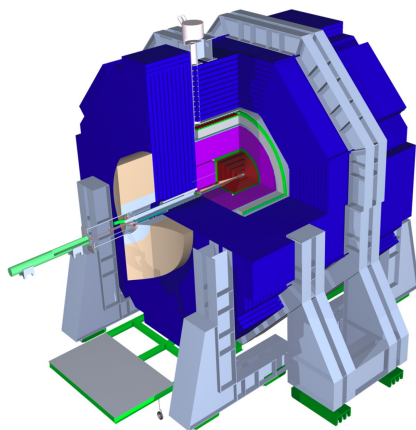
→ DM mass sensitivity nearly half \sqrt{s}

Soft photons, forward detectors

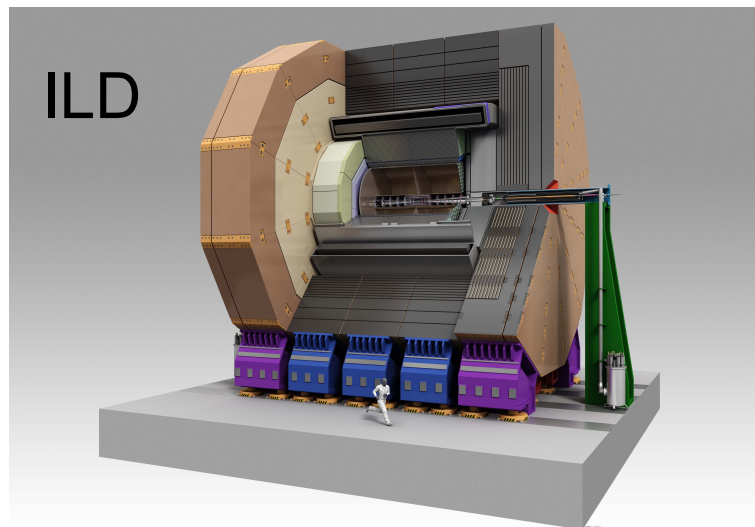


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\mathcal{L}	$2 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
P_{e^-}	>80%
P_{e^+}	~30%
Length	~31 km

-> Talk by Steinar Stapnes



SiD



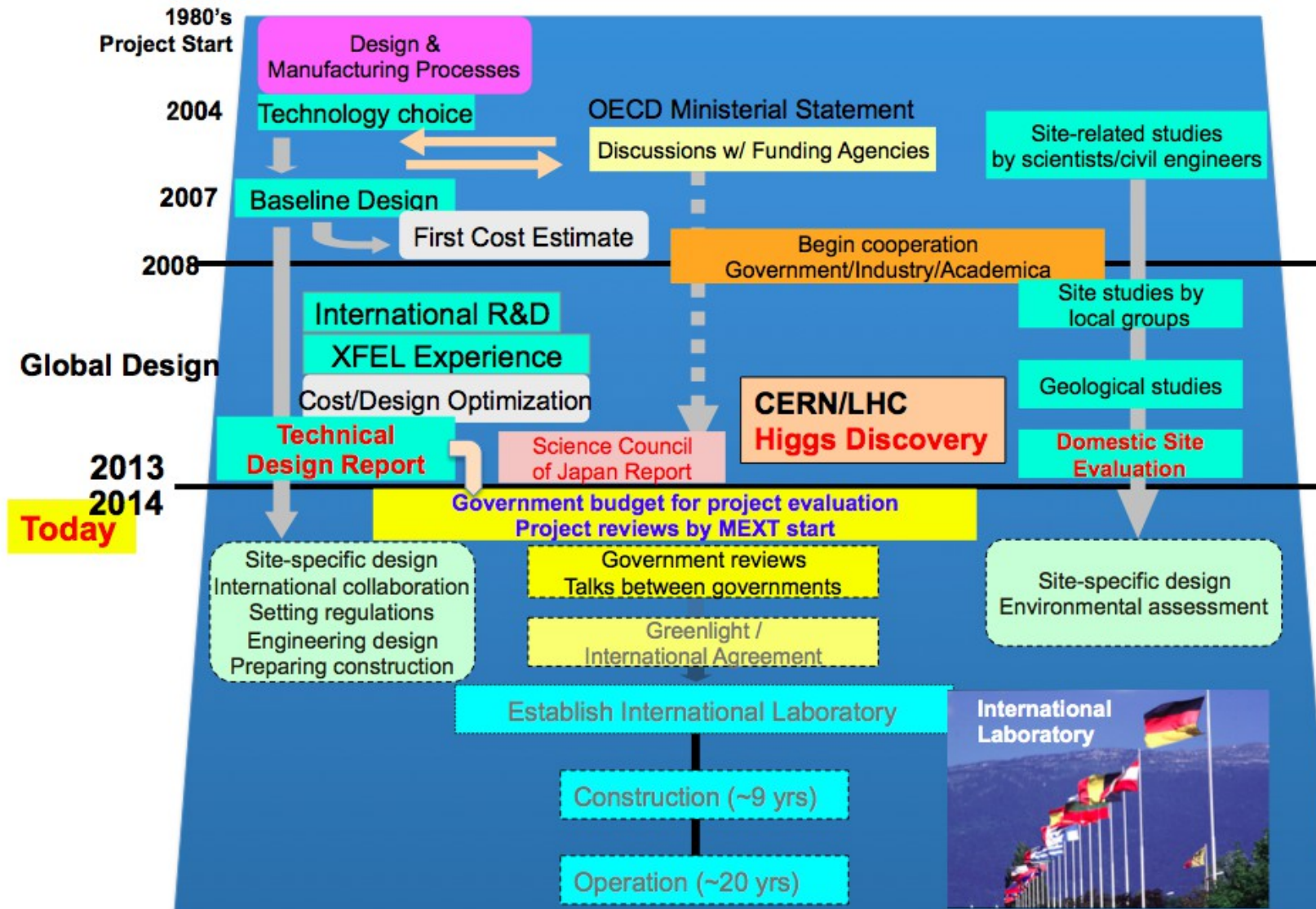
ILD

Talks by:

Imad Laktineh
 Frank Simon
 Marek Idzik
 Lucie Linssen

Machine TDR in 2013 + DBD for detectors

Timeline of ILC



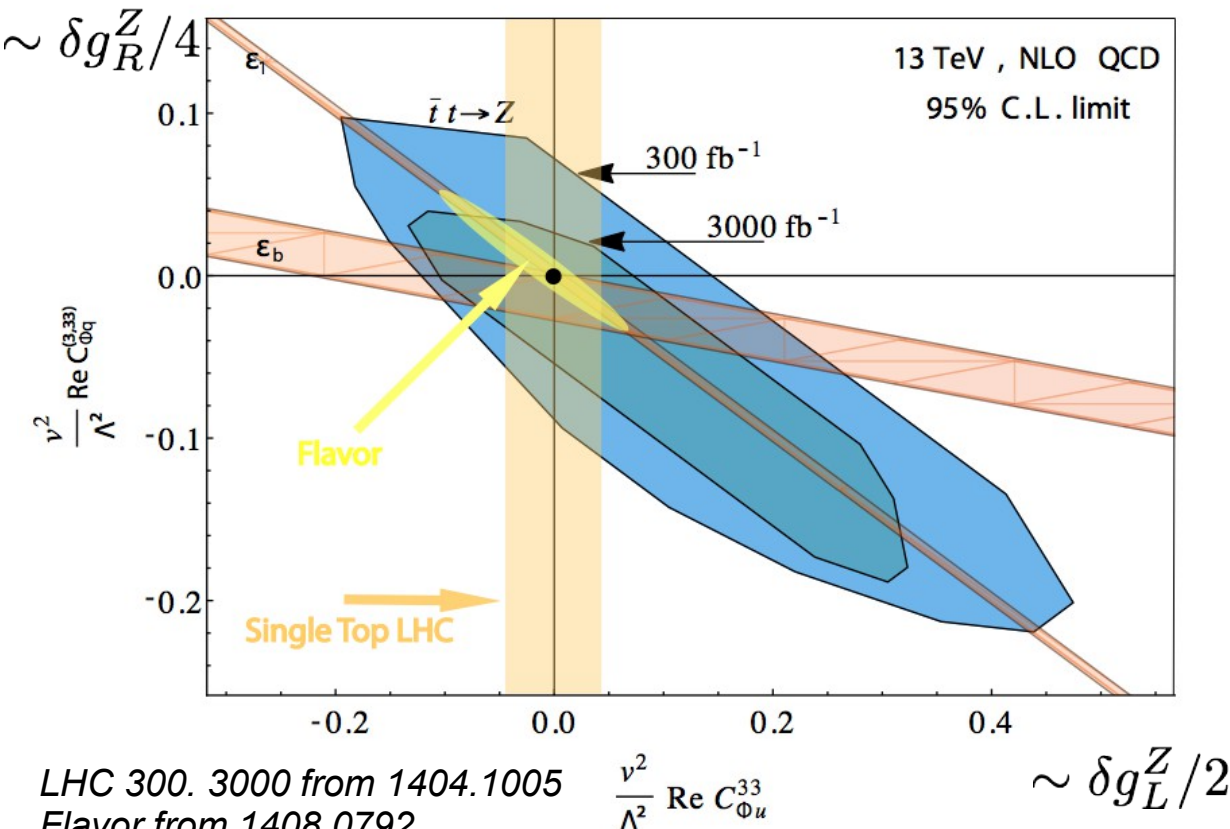
Remark R.P.: MEXT report in March 2016

Precision cross section $\sim 0.5\%$,

Precision $A_{FB} \sim 2\%$,

Precision $\lambda_t \sim 3-4\%$

Accuracy on SM Z couplings compared with other experiments



LHC 300. 3000 from 1404.1005
Flavor from 1408.0792
LHC Single top added by F. Richard

- ILC with polarised beams outperforms all present and future experiments (Stringent limits only from LEP)
- Before ILC single top at LHC and B factories can deliver complementary information
- In particular g_R can only be constrained by ILC!
- Maintaining this high level still requires substantial experimental and theoretical work

ILC promises to be high precision machine for electroweak top couplings