

Experimental status on Kappa formalism and EFT approach

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Precision theory for precise measurements
at the LHC and future colliders

Quy Nhon, Vietnam

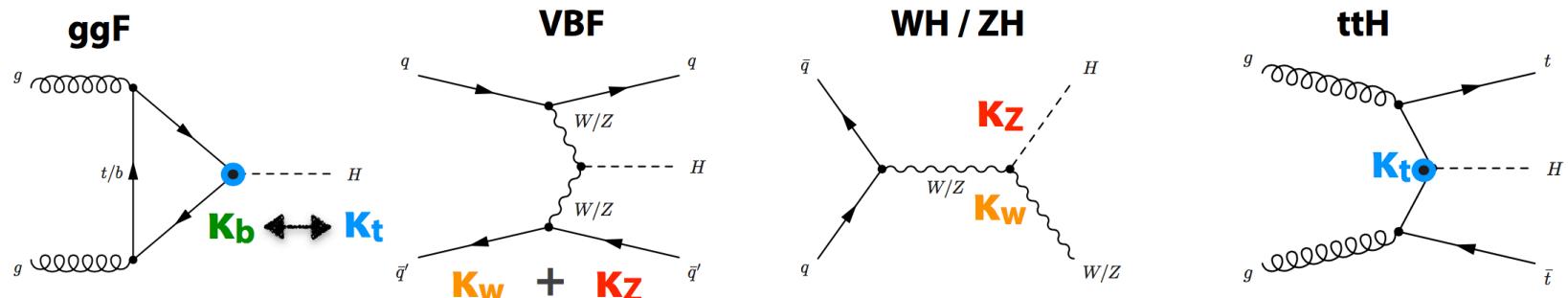
2016.9.25-10.1

Introduction

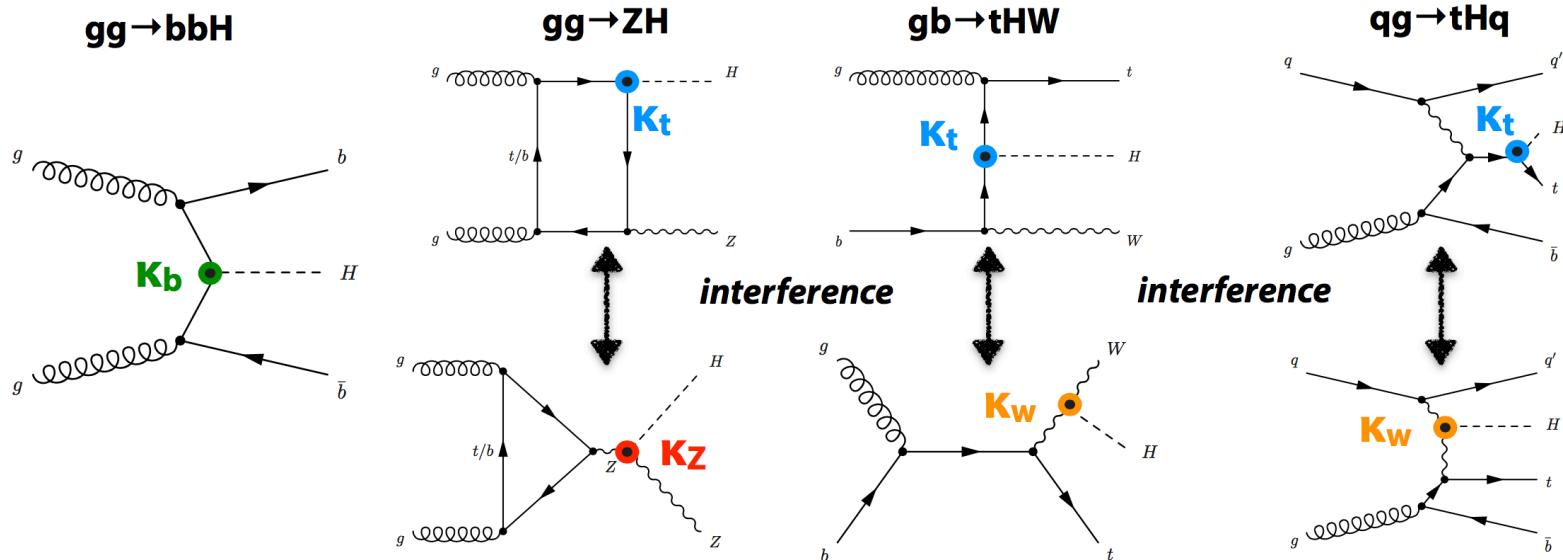
- ATLAS and CMS published in June the Run 1 legacy LHC combination on Higgs rates and couplings
[JHEP 1608 \(2016\) 045](#)
 - Tremendous efforts from both experiments and theorists
 - Close the chapter on Run 1 Higgs experimental results
- Consistency tests performed with the κ -framework based on [LHC HXSWG YR3](#)
 - A LO-motivated framework
 - Stepping from signal strengths to couplings
- Furthermore, first results on differential cross sections and simplified template cross sections have been made with Run 1 and Run 2 data
 - To bridge the data and interpretations (e.g. EFT)

Higgs production processes

- Usual suspects:



- Rare processes:



Signal strengths, μ

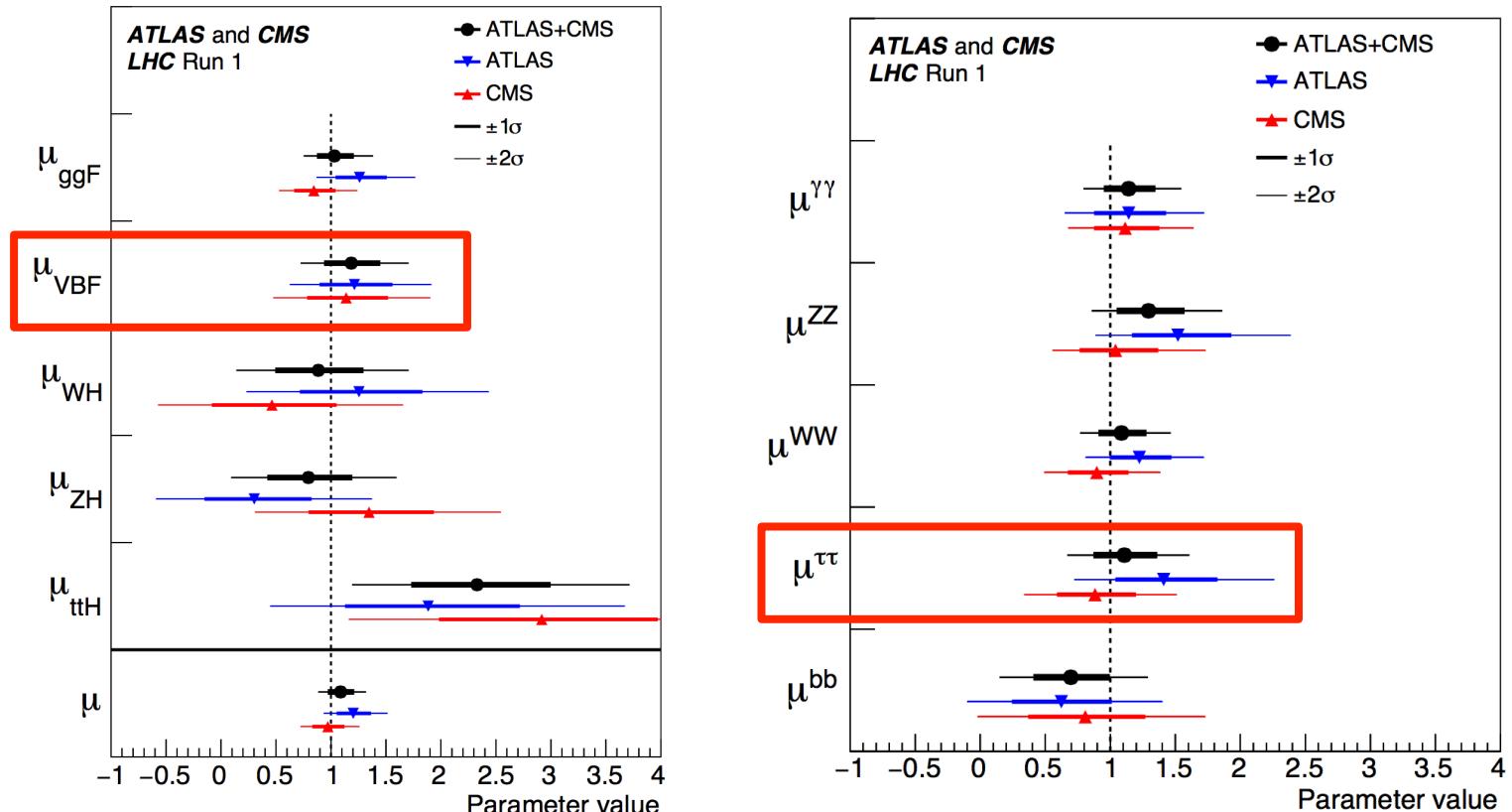
Parameters scale cross sections and BRs relative to SM

$$\mu_i = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \quad \mu^f = \frac{\text{BR}^f}{\text{BR}_\text{SM}^f}$$

Scaling of generic $i \rightarrow H \rightarrow f$ process

$$\mu_i^f \equiv \frac{\sigma_i \cdot \text{BR}^f}{(\sigma_i \cdot \text{BR}^f)_\text{SM}} = \mu_i \times \mu^f$$

Signal strengths, μ

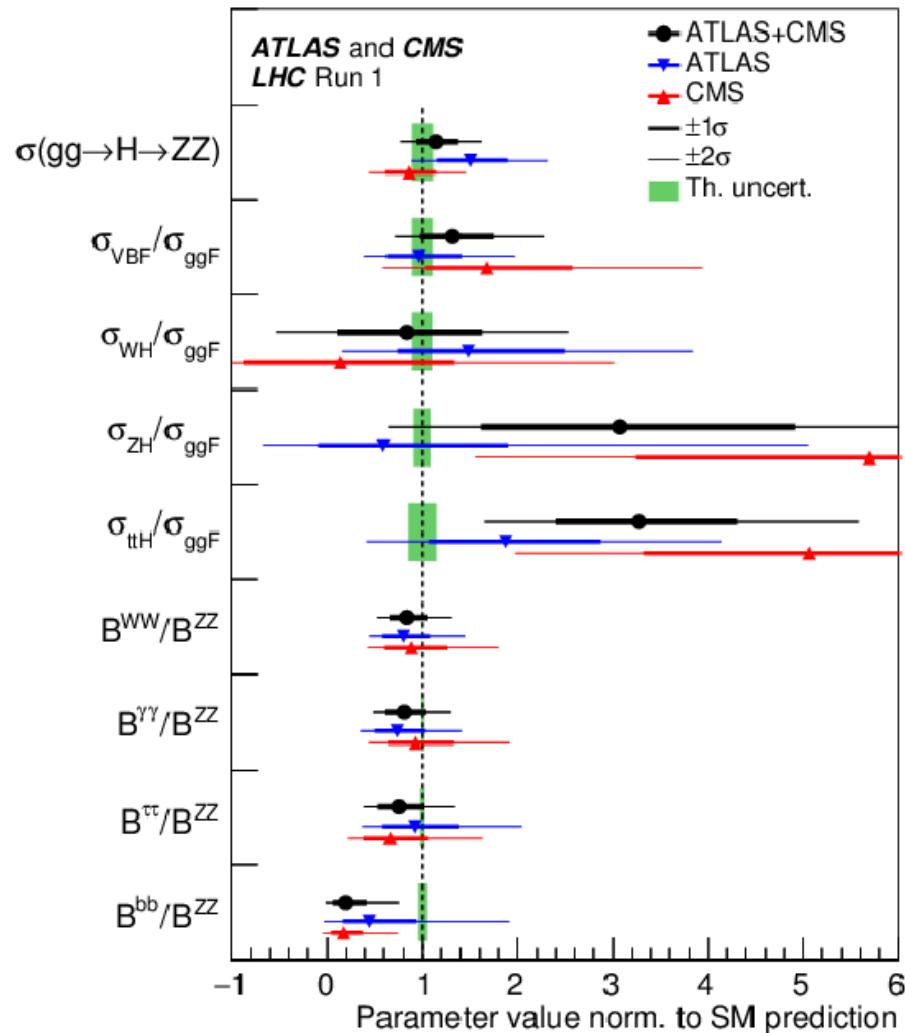


H \rightarrow $\tau\tau$ decay mode and VBF process established through the ATLAS+CMS combination

Ratios of cross-sections and BRs

- Assumptions: **only the 7/8 ratios as in SM**
- Normalize the rate for any particular channel to a reference process using ratios of cross sections and branching ratios
- Motivation:
 - No assumptions on relative cross sections or BRs
 - Measured values independent of SM prediction and inclusive theory uncertainties
 - Cancellation of common systematic uncertainties in ratios
- Choose reference process as one measured with the **smallest systematic uncertainty: $gg \rightarrow H \rightarrow ZZ$**

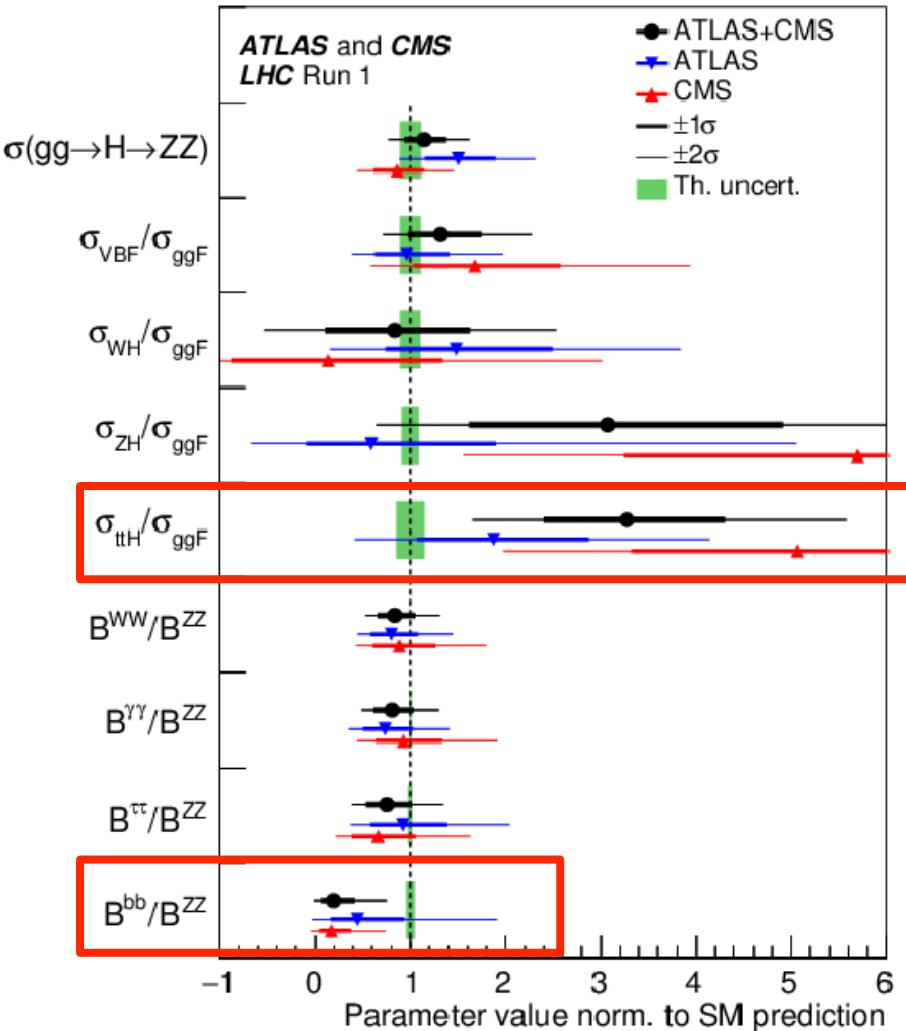
$$\sigma_i \cdot \text{BR}^f = \sigma(gg \rightarrow H \rightarrow ZZ) \times \left(\frac{\sigma_i}{\sigma_{ggF}} \right) \times \left(\frac{\text{BR}^f}{\text{BR}^{ZZ}} \right)$$



Ratios of cross-sections and BRs

- Largest disagreement in $\sigma_{ttH}/\sigma_{ggF}$ ($\sim 3\sigma$)
 - mainly due to the multi-lepton categories
- BR_{bb}/BR_{ZZ} (2.5σ)
 - High observed values σ_{ZH}/σ_{ggF} and $\sigma_{ttH}/\sigma_{ggF}$ induce a low value for $BR(H \rightarrow bb)$
 - $H \rightarrow bb$ decay mode does not contribute to the observed excesses
 - Anti-correlated with above excesses

$$\sigma_i \cdot BR^f = \sigma(gg \rightarrow H \rightarrow ZZ) \times \left(\frac{\sigma_i}{\sigma_{ggF}} \right) \times \left(\frac{BR^f}{BR^{ZZ}} \right)$$



From μ to κ

Signal strengths, μ

Parameters scale cross sections and
BRs relative to SM

$$\mu_i = \frac{\sigma_i}{\sigma_i^{\text{SM}}} \quad \mu^f = \frac{\text{BR}^f}{\text{BR}_{\text{SM}}^f}$$

Scaling of generic $i \rightarrow H \rightarrow f$ process

$$\mu_i^f \equiv \frac{\sigma_i \cdot \text{BR}^f}{(\sigma_i \cdot \text{BR}^f)_{\text{SM}}} = \mu_i \times \mu^f$$

Couplings, κ

Parameters scale cross sections and
partial widths relative to SM

$$\kappa_j^2 = \sigma_j / \sigma_j^{\text{SM}} \quad \kappa_j^2 = \Gamma_j / \Gamma_j^{\text{SM}}$$
$$\sigma_i \cdot \text{BR}^f = \frac{\sigma_i \cdot \Gamma_f}{\Gamma_H},$$

Total width determined as

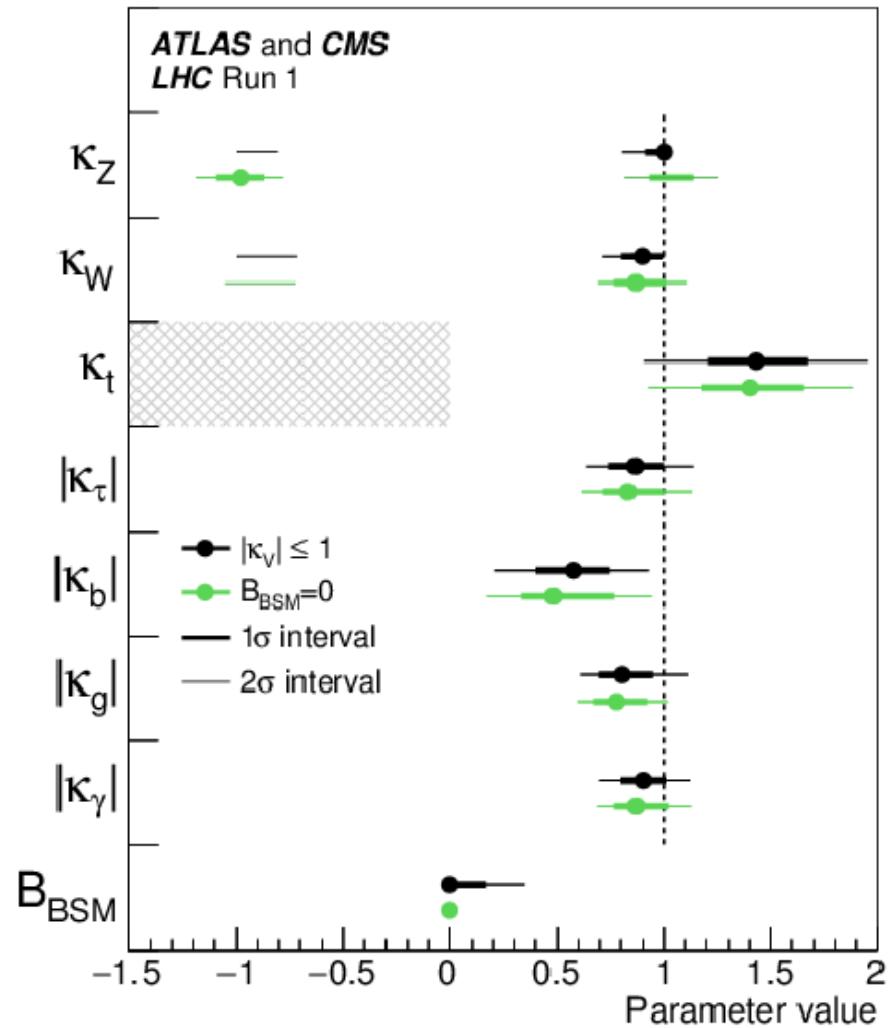
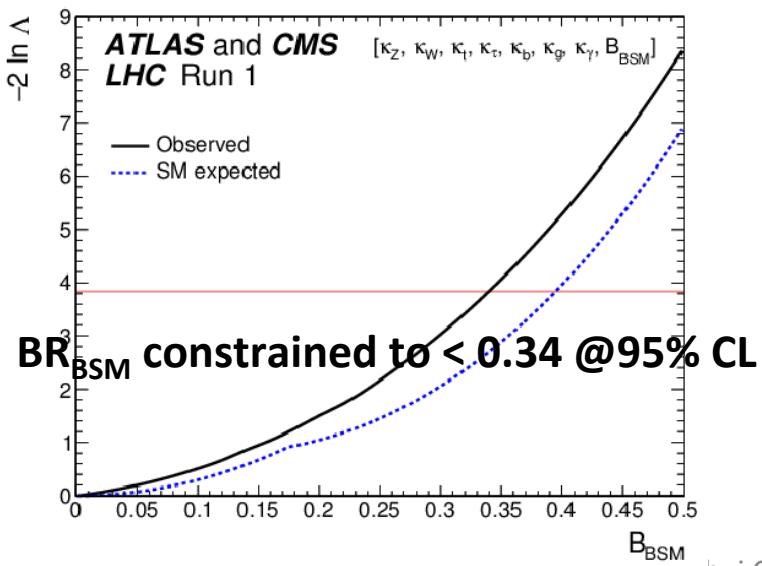
$$\Gamma_H = \frac{\kappa_H^2 \cdot \Gamma_H^{\text{SM}}}{1 - \text{BR}_{\text{BSM}}}$$

Where

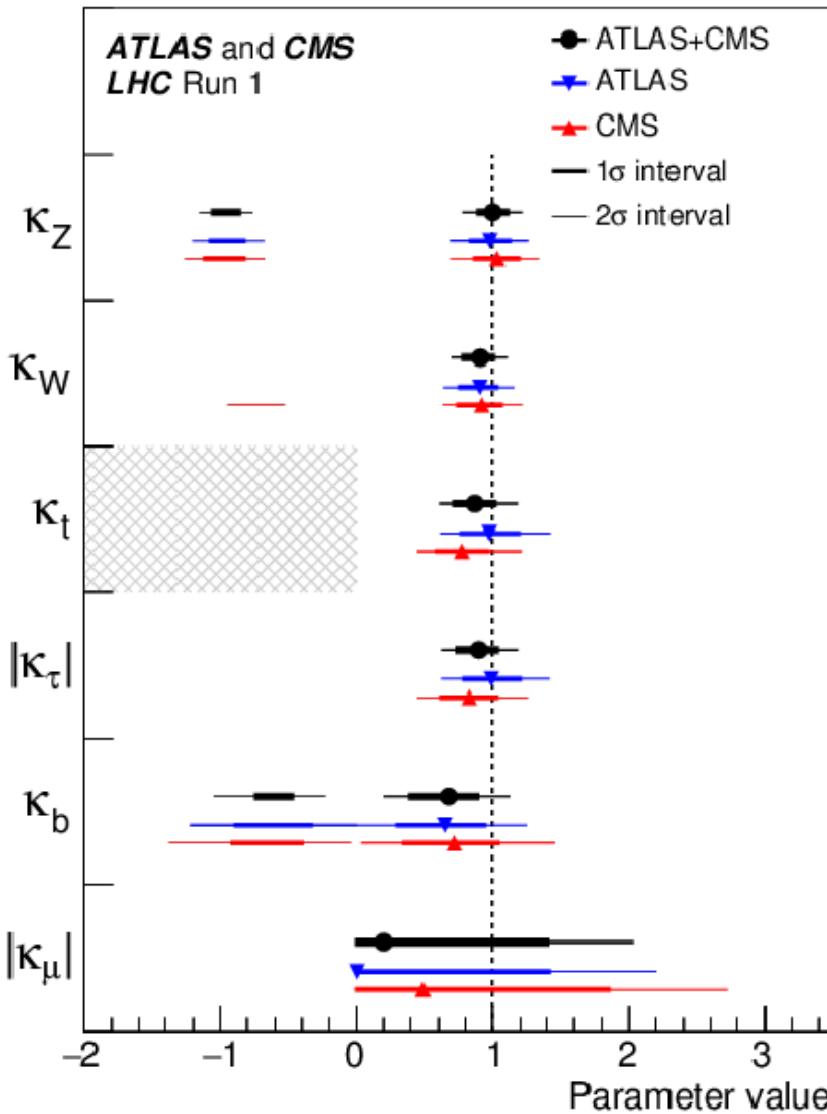
$$\kappa_H^2 = \sum_j \text{BR}_{\text{SM}}^j \kappa_j^2$$

Allow BSM contribution to loop & decay

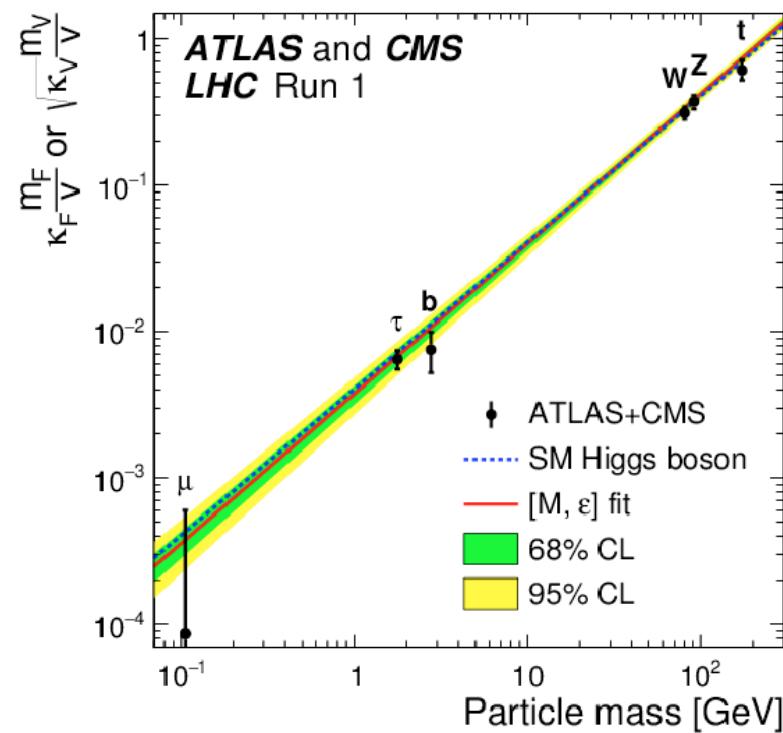
- Use effective couplings for ggH (κ_g) and $H \rightarrow \gamma\gamma$ (κ_γ)
- BR_{BSM} : Higgs decays to BSM particles + non-SM BRs of unmeasured final states (gg, cc, \dots)
- Consider two scenarios:
 - $BR_{BSM} = 0$
 - BR_{BSM} floating, but $|\kappa_v| < 1$



No BSM loop/decay contributions

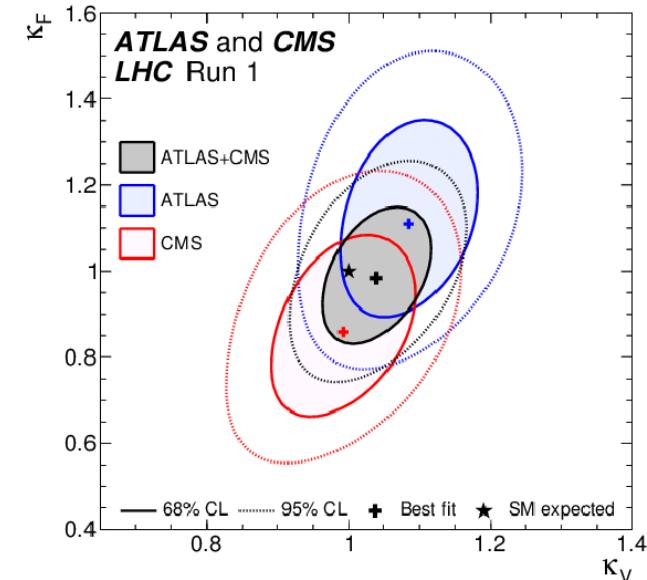
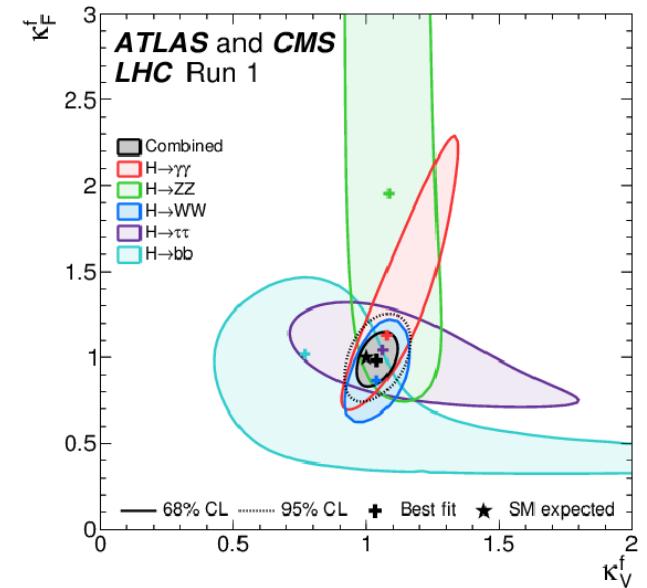


- Resolve ggH (κ_g) and H $\rightarrow\gamma\gamma$ (κ_γ) loops
- Includes H $\rightarrow\mu\mu$ analyses for reduced coupling vs. particle mass



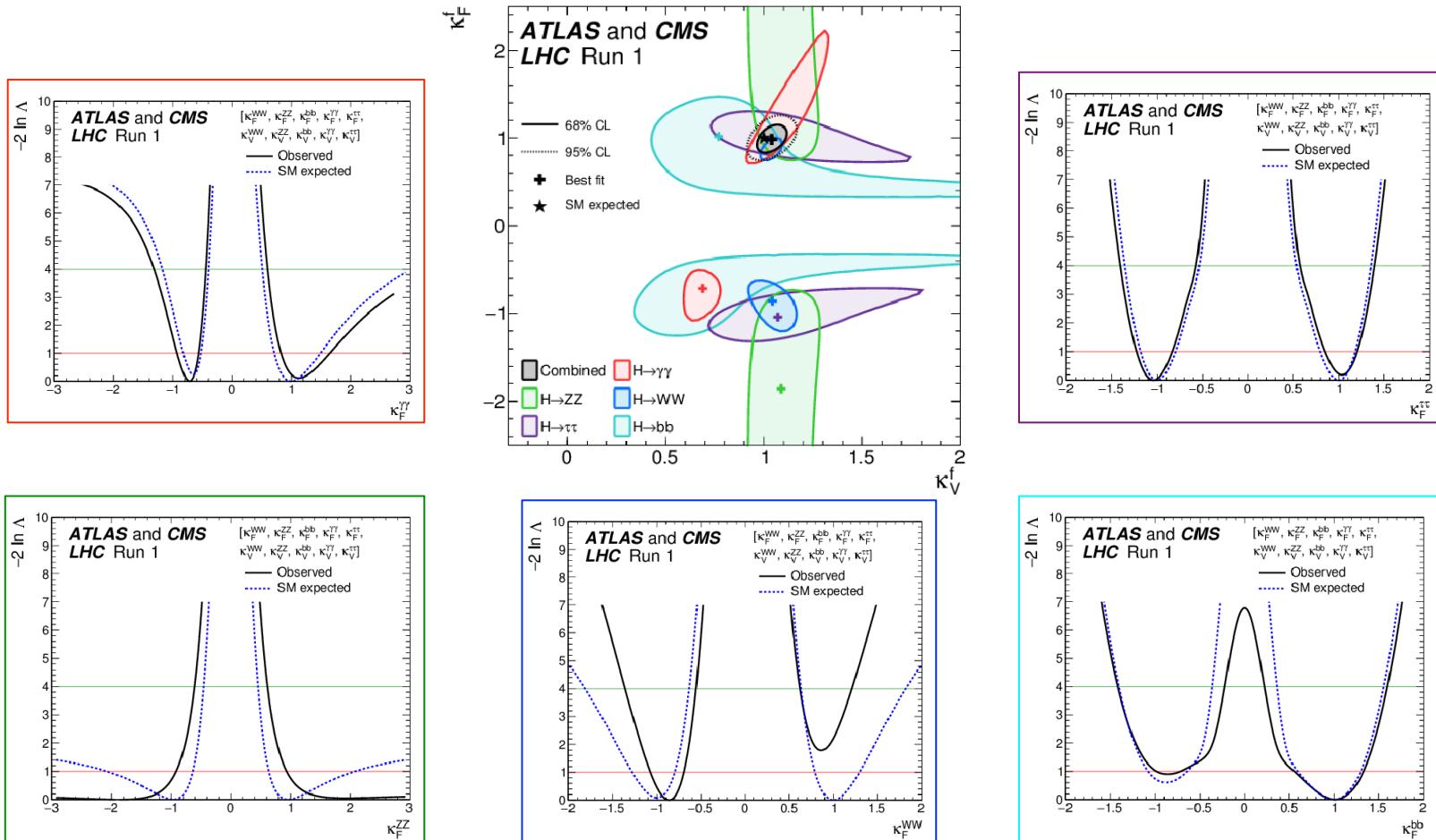
2D scans of κ_V , κ_F

- Simplified model in which
 - $\kappa_V = \kappa_W = \kappa_Z$
 - $\kappa_F = \kappa_t = \kappa_b = k_\tau$
- Different final states give very different constrained region in the $\kappa_V \kappa_F$ (all channels compatible to SM)
- Combination of ATLAS and CMS also compatible within 1σ with SM



2D scans of κ_V , κ_F

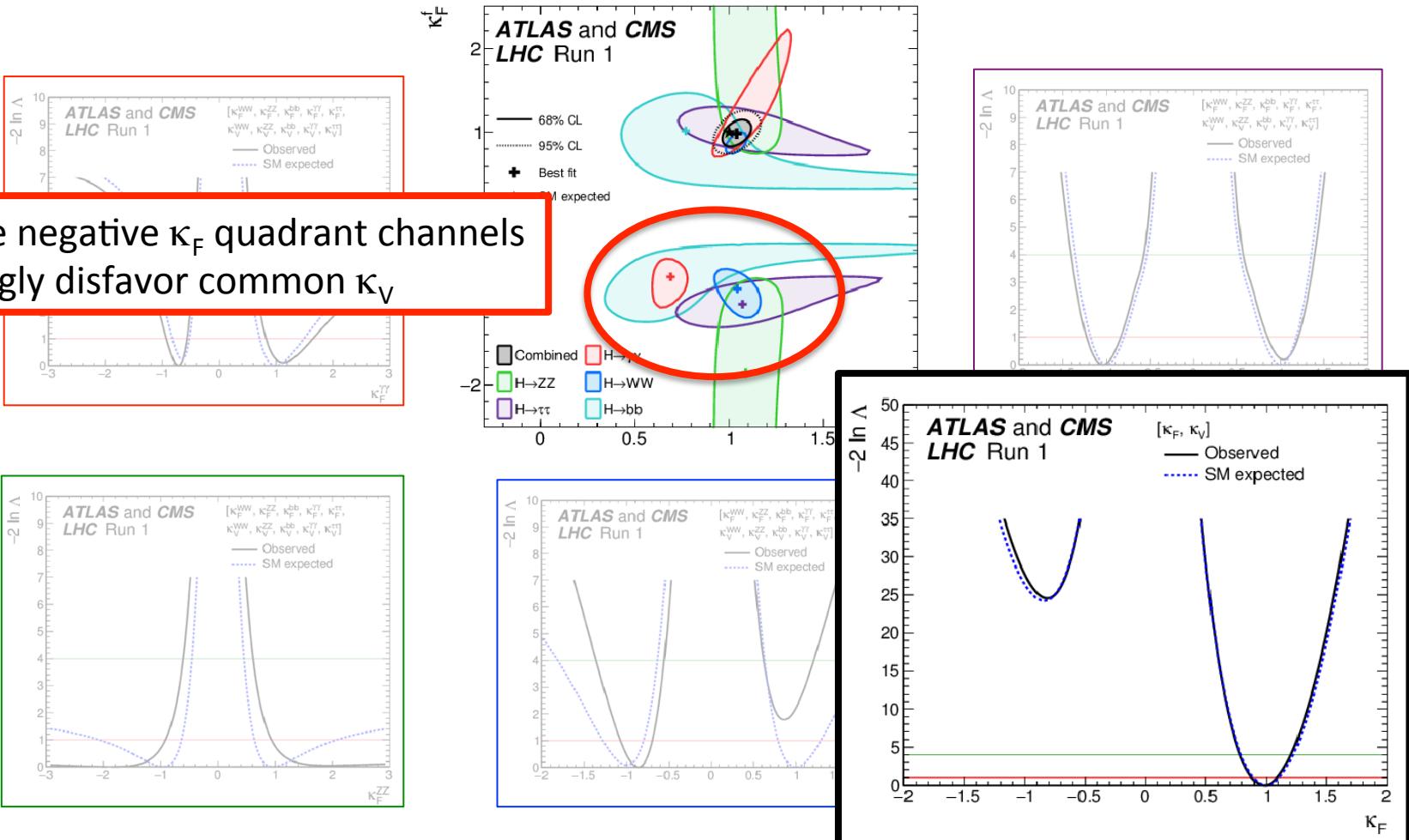
- Most channels nearly degenerate in relative sign of κ_V and κ_F



2D scans of κ_V , κ_F

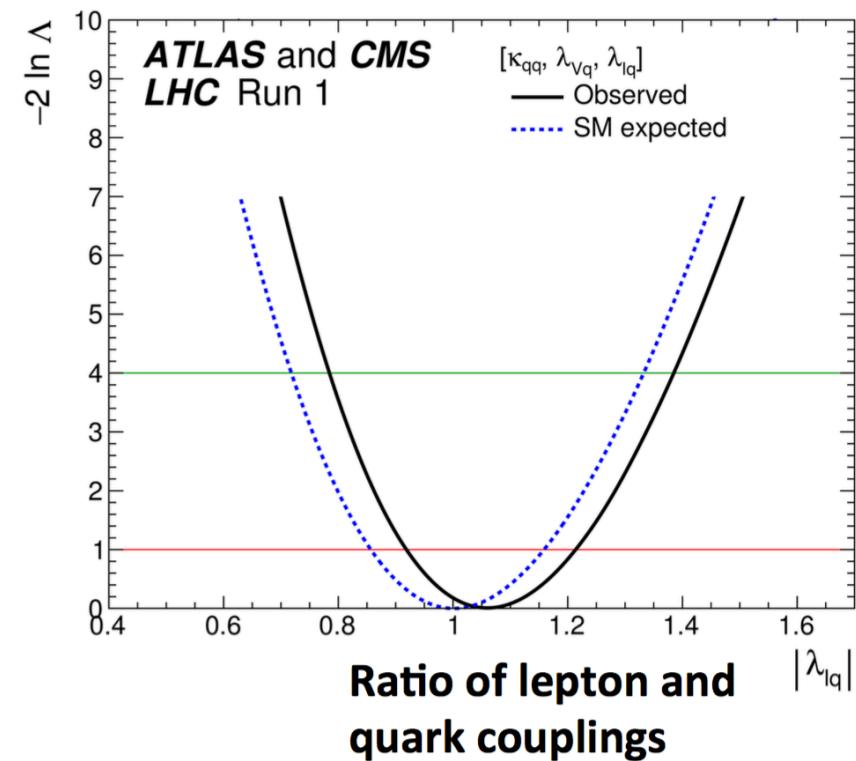
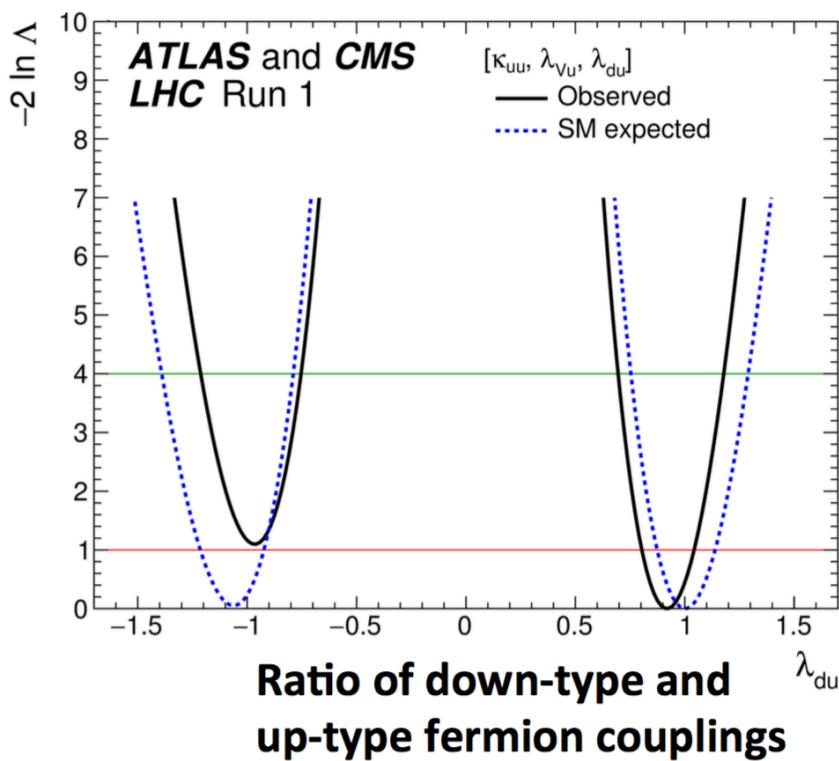
- Most channels nearly degenerate in relative sign of κ_V and κ_F

In the negative κ_F quadrant channels strongly disfavor common κ_V



Fermion Couplings

- In MSSM / 2HDM Type II [$\kappa_v, \kappa_d, \kappa_u$], ratio of down-type (b, τ, μ) and up-type (t) fermion couplings is tested with $\sim 10\%$ precision
- No enhancement observed wrt SM, i.e. consistent with alignment limit
- In 2HDM Lepton-Specific [$\kappa_v, \kappa_l, \kappa_q$], ratio of lepton (τ, μ) and quark couplings (t, b) would be enhanced at large $\tan \beta$
- Also good agreement with SM



Going beyond the κ -framework

- The LO-motivated κ framework has the following assumptions/limitations
 - Single, narrow resonance with $m_H = 125.09 \text{ GeV}$
 - Decay kinematics compatible with SM prediction
 - Tensor structure of interactions is the same as in SM
 - Zero-width approximation in the Higgs boson propagator
- To test kinematics distributions
 - Anomalous couplings
 - EFT interpretation through differential cross section measurements

CMS 13 TeV H4I anomalous couplings

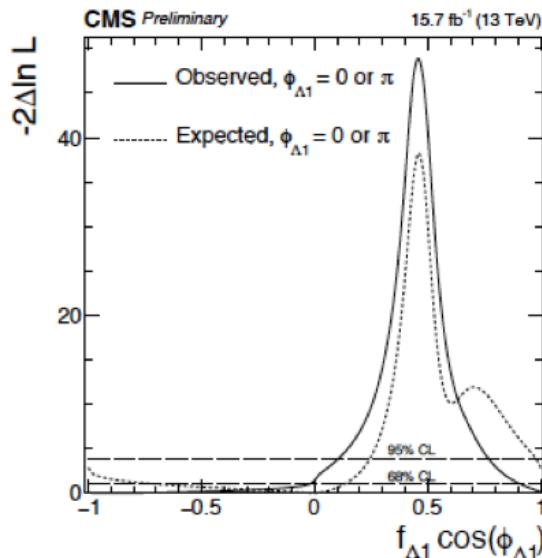
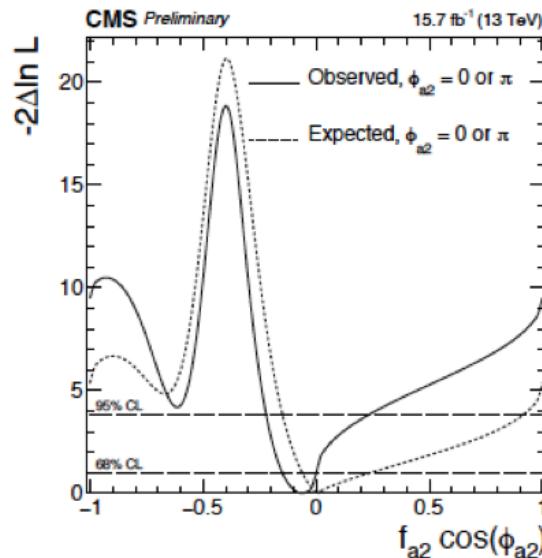
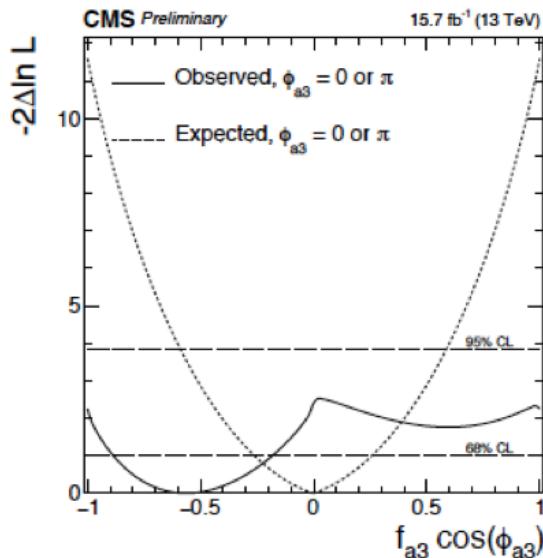
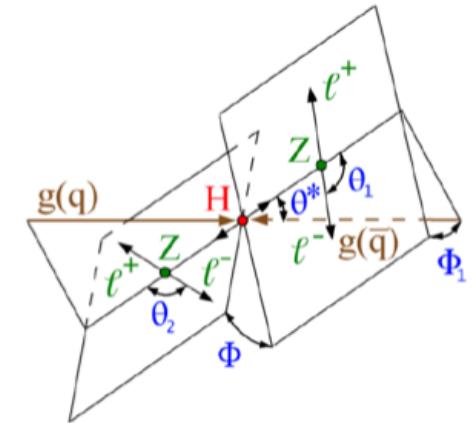
- Characterized through scattering amplitude

$$A(\text{HVV}) \sim \left[a_1^{\text{VV}} + \frac{\kappa_1^{\text{VV}} q_{\text{V}1}^2 + \kappa_2^{\text{VV}} q_{\text{V}2}^2}{(\Lambda_1^{\text{VV}})^2} \right] m_{\text{V}1}^2 \epsilon_{\text{V}1}^* \epsilon_{\text{V}2}^* + a_2^{\text{VV}} f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + a_3^{\text{VV}} f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu}$$

$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4 + \dots}, \quad \phi_{a3} = \arg \left(\frac{a_3}{a_1} \right),$$

$$f_{a2} = \frac{|a_2|^2 \sigma_2}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4 + \dots}, \quad \phi_{a2} = \arg \left(\frac{a_2}{a_1} \right),$$

$$f_{\Lambda 1} = \frac{\tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4 + \dots}, \quad \phi_{\Lambda 1},$$



CMS 13 TeV H4I anomalous couplings

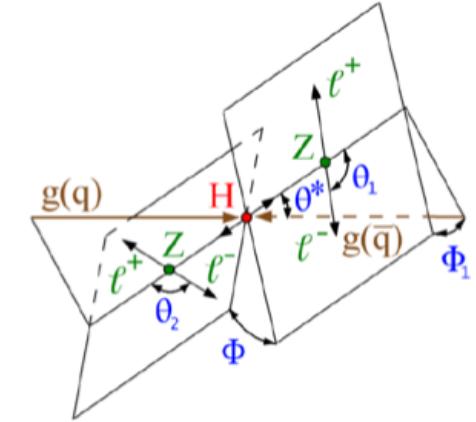
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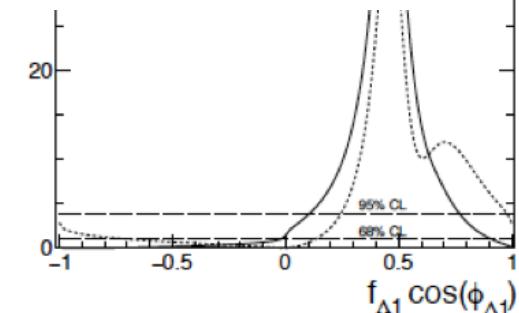
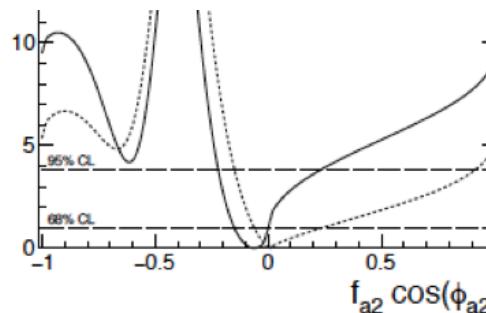
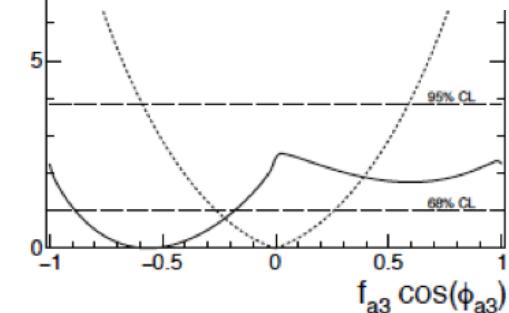
$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4 + \dots}, \quad \phi_{a3} = \arg \left(\frac{a_3}{a_1} \right),$$

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$$f_{\Lambda 1} = \frac{\tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda 1}/(\Lambda_1)^4 + \dots}, \quad \phi_{\Lambda 1},$$



Parameter	Observed	Expected
$f_{a3} \cos(\phi_{a3})$	$-0.56^{+0.38}_{-0.32} [-1.00, 1.00]$	$0.00^{+0.26}_{-0.26} [-0.59, 0.59]$
$f_{a2} \cos(\phi_{a2})$	$-0.06^{+0.06}_{-0.09} [-0.22, 0.24]$	$0.00^{+0.24}_{-0.06} [-0.15, 0.92]$
$f_{\Lambda 1} \cos(\phi_{\Lambda 1})$	$-0.93^{+0.90}_{-0.16} [-1.00, 0.10] \cup [0.77, 1.00]$	$0.00^{+0.13}_{-0.69} [-1.00, 0.24] \cup [0.98, 1.00]$

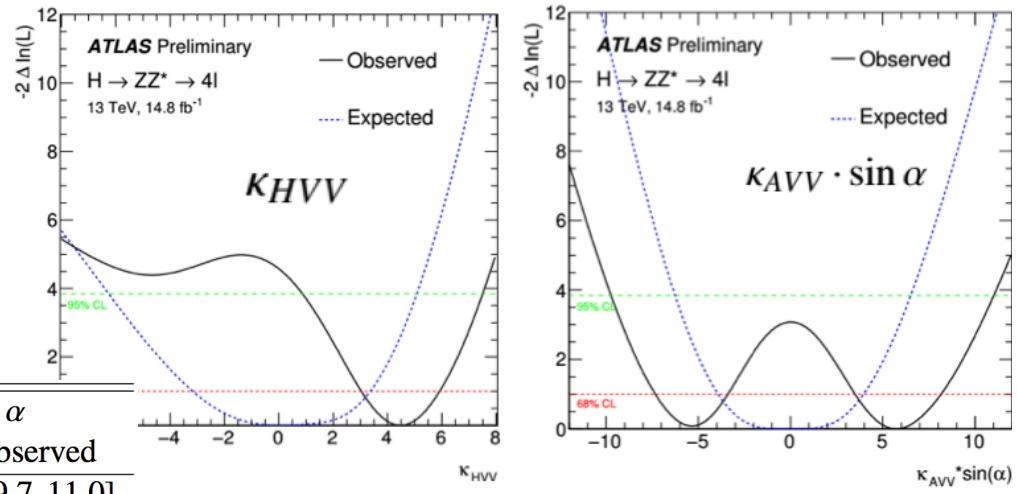


ATLAS 13 TeV H4I

- In terms of EFT couplings defined in Higgs characterization model:

$$\mathcal{L}_0^V = \left\{ \cos(\alpha)\kappa_{\text{SM}} \left[\frac{1}{2}g_{HZZ}Z_\mu Z^\mu + g_{HWW}W_\mu^+ W^{-\mu} \right] \right. \\ - \frac{1}{4}\frac{1}{\Lambda} \left[\cos(\alpha)\kappa_{HZZ}Z_{\mu\nu}Z^{\mu\nu} + \sin(\alpha)\kappa_{AZZ}Z_{\mu\nu}\tilde{Z}^{\mu\nu} \right] \\ \left. - \frac{1}{2}\frac{1}{\Lambda} \left[\cos(\alpha)\kappa_{HWW}W_{\mu\nu}^+ W^{-\mu\nu} + \sin(\alpha)\kappa_{AWW}W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] \right\} X_0$$

J^P	Model	Values of tensor couplings			
		κ_{SM}	κ_{HVV}	κ_{AVV}	α
0^+	SM Higgs boson	1	0	0	0
0_h^+	BSM spin-0 CP-even	0	1	0	0
0^-	BSM spin-0 CP-odd	0	0	1	$\pi/2$



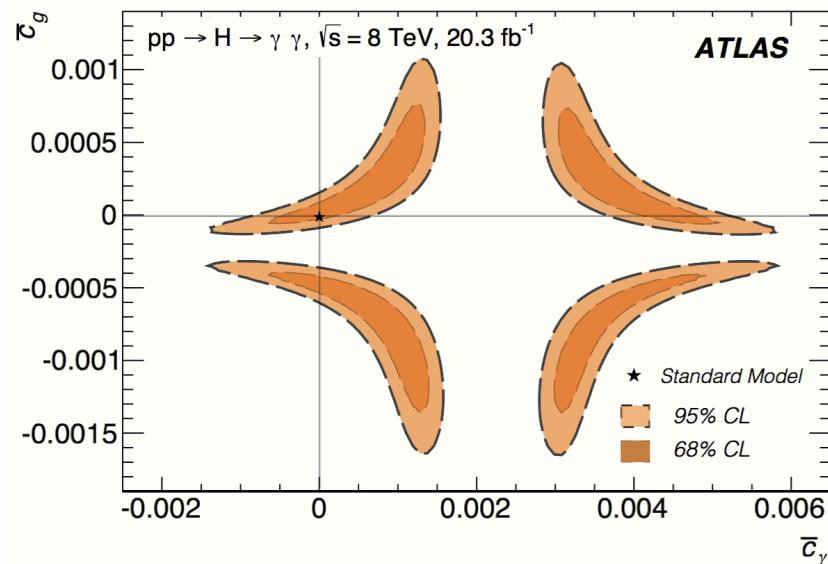
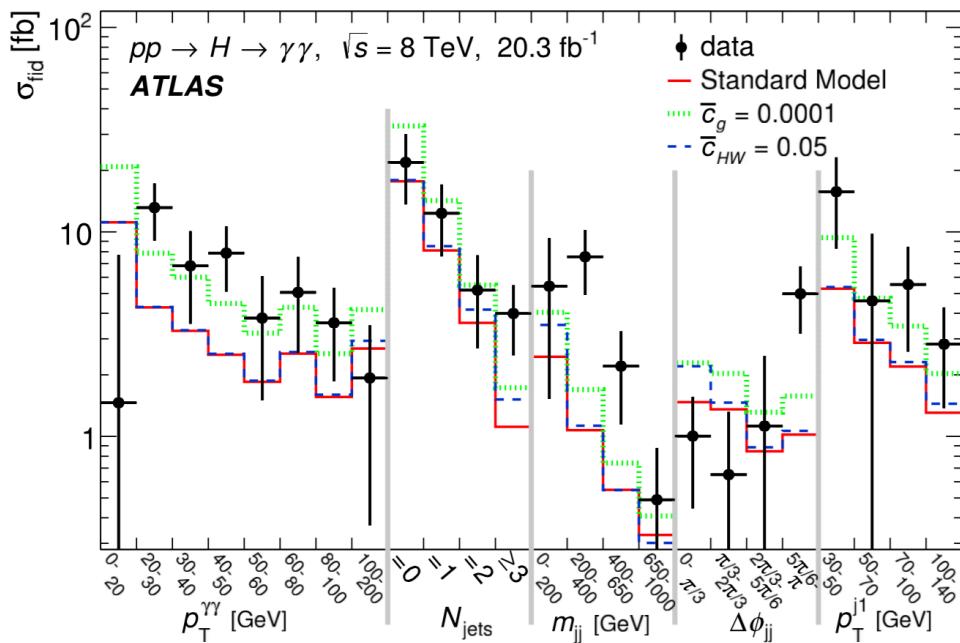
Not excluded range at 95% CL	κ_{HVV}		$\kappa_{AVV} \cdot \sin \alpha$	
	expected	observed	expected	observed
	$[-6.3, 5.1]$	$[0.9, 7.5]$	$[-6.3, 6.5]$	$[-9.7, 11.0]$

EFT with differential cross sections

- Probe BSM CP-even and CP-odd interactions using five differential cross sections in EFT framework and the Strongly Interacting Light Higgs (SILH) formulation:

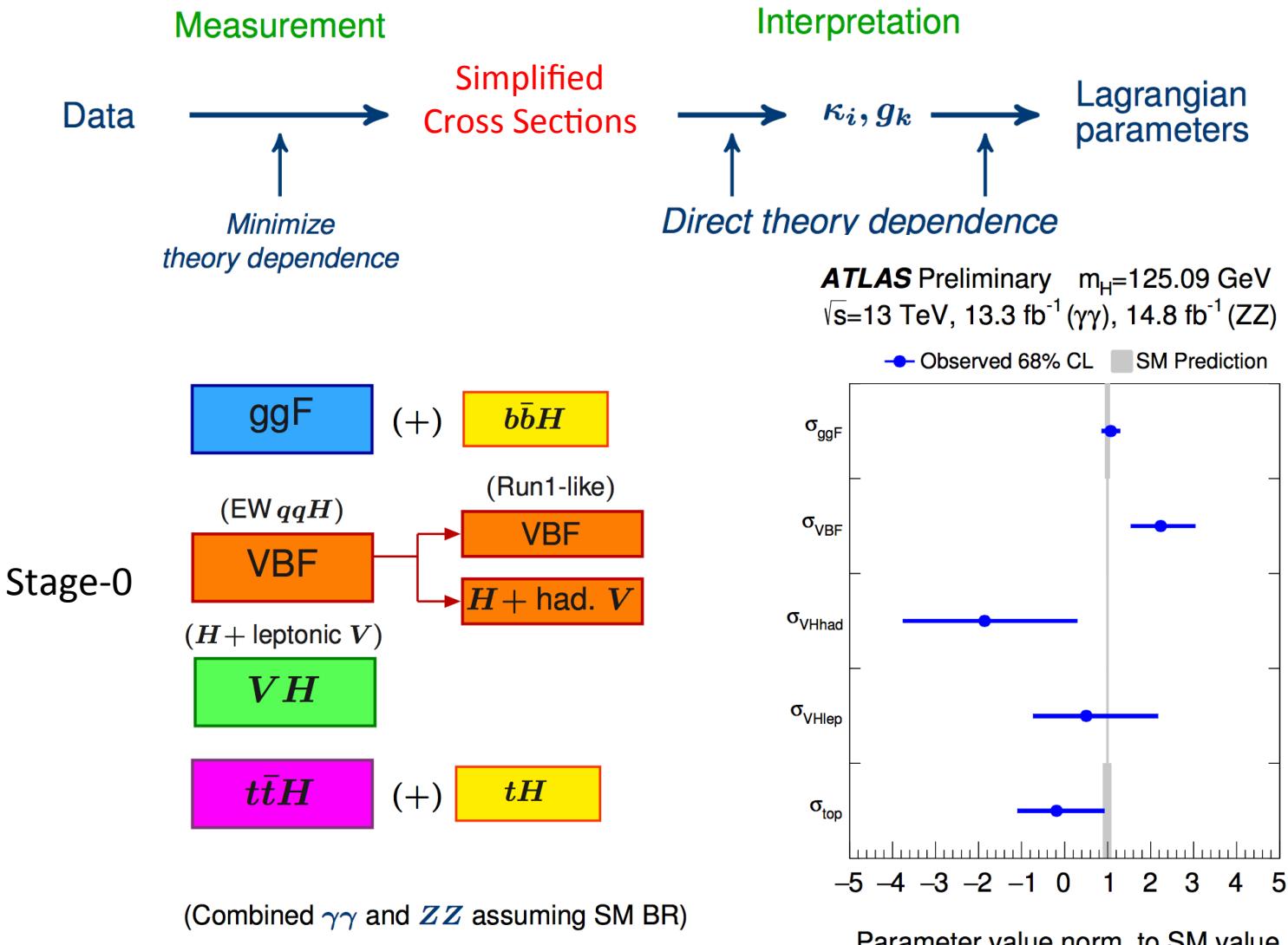
$$\begin{aligned}\mathcal{L} = & \bar{c}_\gamma O_\gamma + \bar{c}_g O_g + \bar{c}_{HW} O_{HW} + \bar{c}_{HB} O_{HB} \\ & + \tilde{c}_\gamma \tilde{O}_\gamma + \tilde{c}_g \tilde{O}_g + \tilde{c}_{HW} \tilde{O}_{HW} + \tilde{c}_{HB} \tilde{O}_{HB}\end{aligned}$$

Wilson coefficients specifying the strength of the new interactions



Simplified Template Cross Sections

- Another complementary approach to connect data and theory interpretations



Summary

- A comprehensive combined measurement of ATLAS and CMS Higgs boson couplings has been performed
 - Precision usually better by $\sim 1/\sqrt{2}$ wrt single experiment
- **First results beyond κ -framework have also been made**
 - **Including the interpretations in EFT approach**
 - **Overall good consistency with SM** expectations
 - Much more data ahead, stay tuned!

The era of Higgs precision measurements has just begun

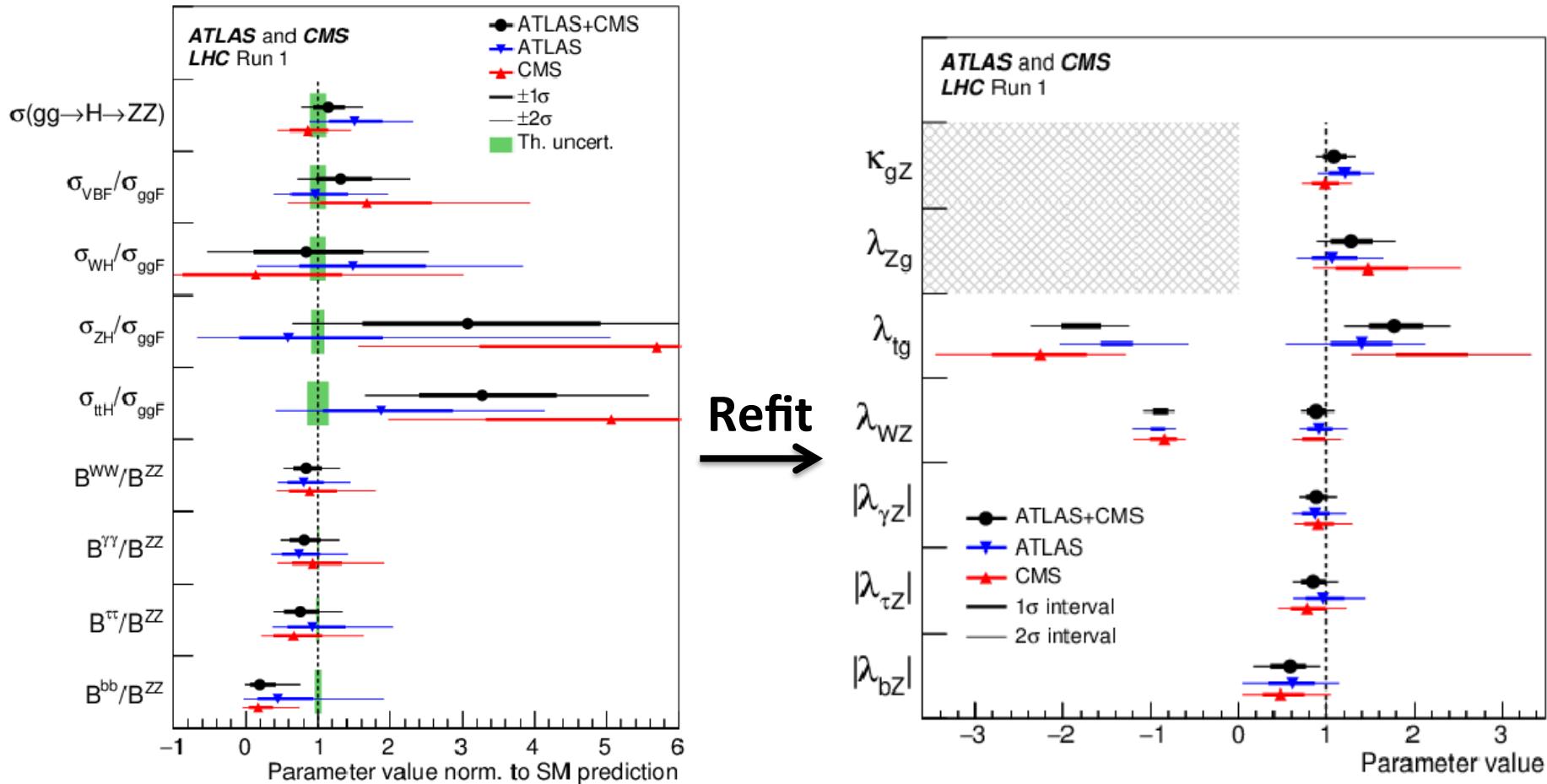


Thanks !

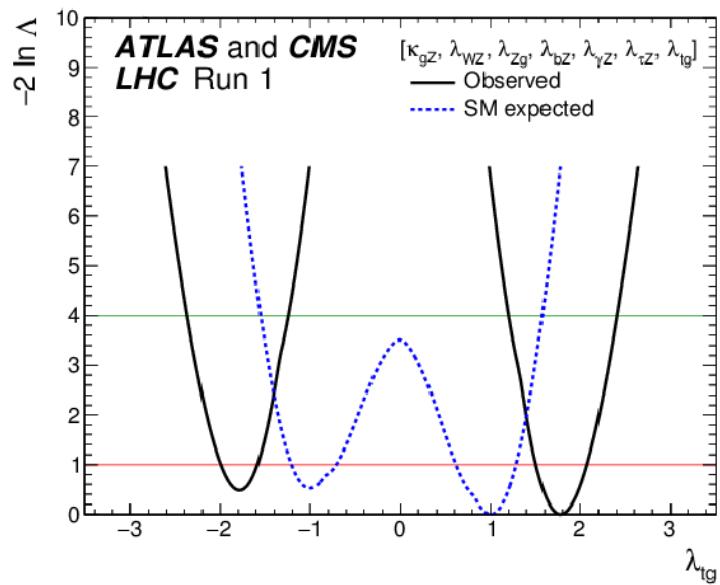
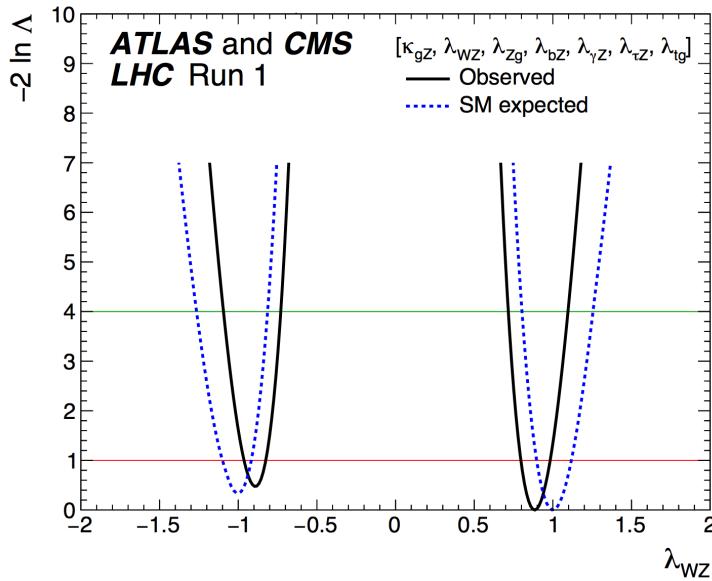
Backup

Ratio of rates → Ratio of couplings

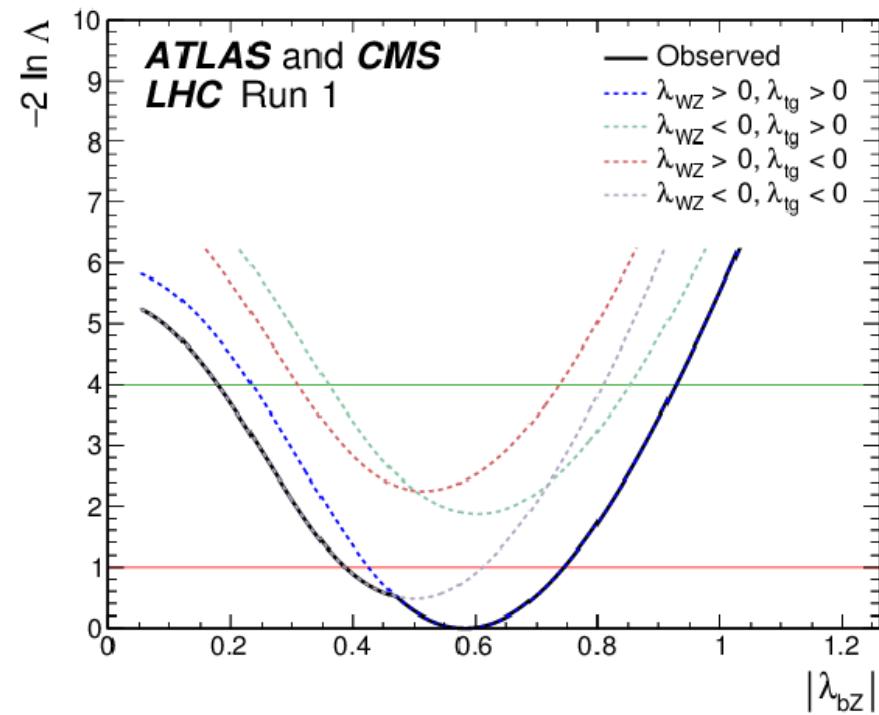
Re-fit fewer couplings to actual particles as ratios to Z and gluon, $\lambda_{ij} = \kappa_i/\kappa_j$



Ratio of couplings



Asymmetry on likelihood due to interference effects in ggZH, tHq/tHW



CP-invariance test: VBF $H \rightarrow \tau\tau$

► HVV couplings as a test of CP-violation / CP-invariance

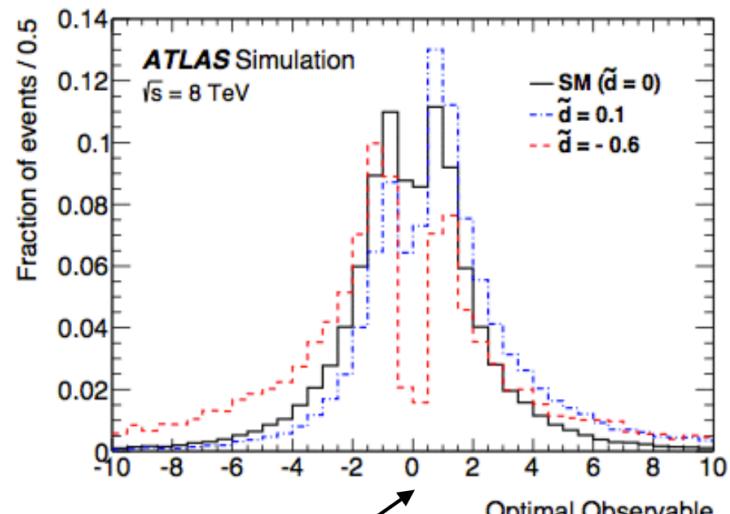
HWW, HZZ decays and $H\gamma\gamma$ differential cross-section
no deviations from Run1 data

Direct test through VBF production ($H \rightarrow \tau\tau$)

- **CP-odd observable:**
 sensitive to interference between SM and CP-odd contributions
- **Optimal observable:** combine multi-dimensional information in a single variable from the VBF production LO matrix-element [independent from H decay mode]

$$OO = \frac{2 \operatorname{Re}(\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{CP-odd}})}{|\mathcal{M}_{\text{SM}}|^2}$$

most sensitive for smallest values



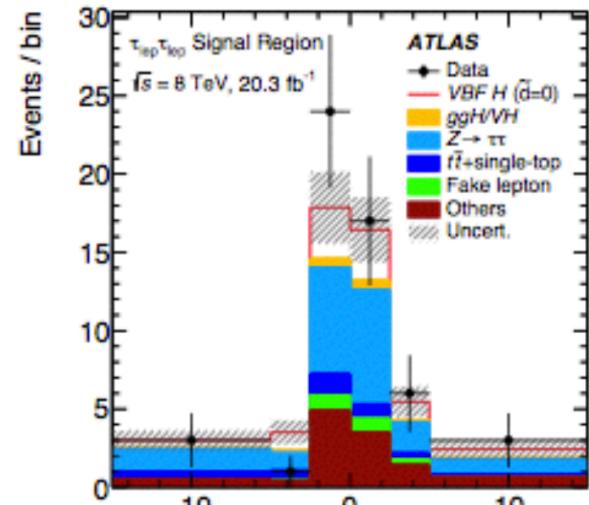
Results interpreted in the Effective Field Theory framework:

CP-violating effects from higher dimension operators on HVV: \tilde{d} parameter

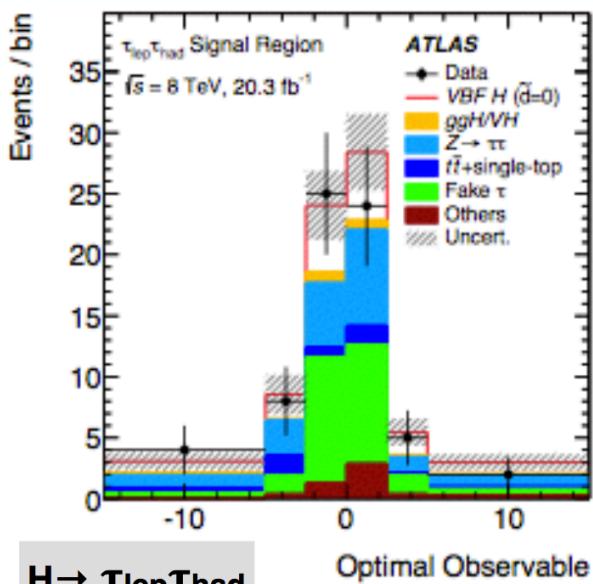
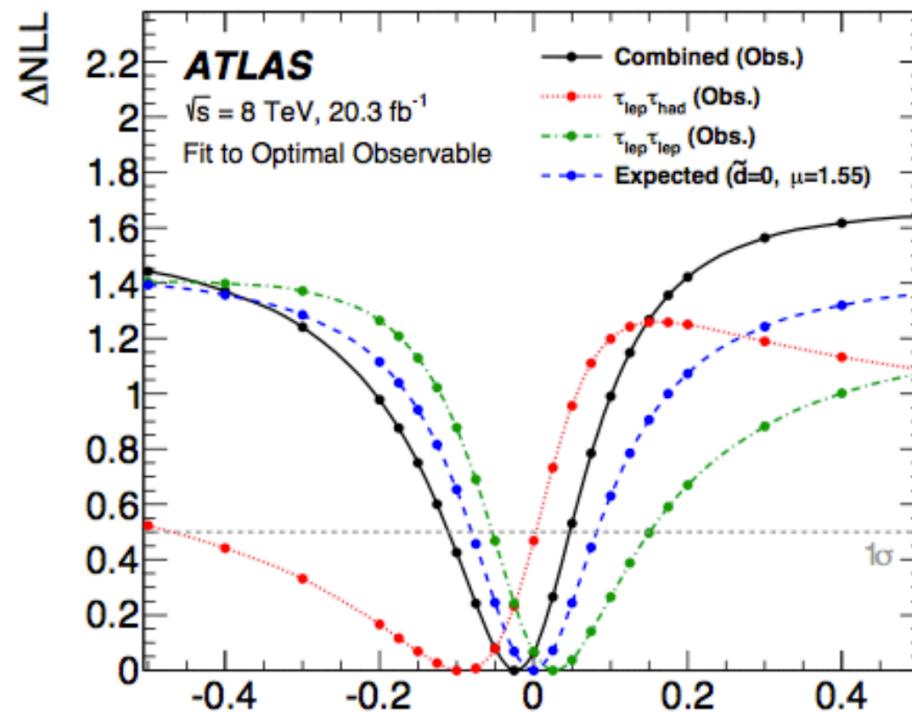
$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \tilde{g}_{HAA} H \tilde{A}_{\mu\nu} A^{\mu\nu} + \tilde{g}_{HAZ} H \tilde{A}_{\mu\nu} Z^{\mu\nu} + \tilde{g}_{HZZ} H \tilde{Z}_{\mu\nu} Z^{\mu\nu} + \tilde{g}_{HWW} H \tilde{W}_{\mu\nu}^+ W^{-\mu\nu}$$

Couplings parametrisation: $\tilde{g}_{HAA} = \tilde{g}_{HZZ} = \frac{1}{2} \tilde{g}_{HWW} = \frac{g}{2m_W} \tilde{d}$

CP-invariance test: VBF $H \rightarrow \tau\tau$



$H \rightarrow \tau_{\text{lep}}\tau_{\text{lep}}$ Optimal Observable



$H \rightarrow \tau_{\text{lep}}\tau_{\text{had}}$ Optimal Observable

No CP-violation ('background-only') minimum at 0 \tilde{d}

- ▶ $[-0.11 < \tilde{d} < 0.05]$ compatible at 68% CL
[competitive with limits from HWW, HZZ]
- ▶ same limit-setting with $\Delta\Phi^{\text{sign}}(\text{jj})$ shows worse results
[azimuthal angle between VBF-tagging jets]

Channel	Fitted value of \tilde{d}
$\tau_{\text{lep}}\tau_{\text{lep}}$	0.3 ± 0.5
$\tau_{\text{lep}}\tau_{\text{had}}$	-0.3 ± 0.4

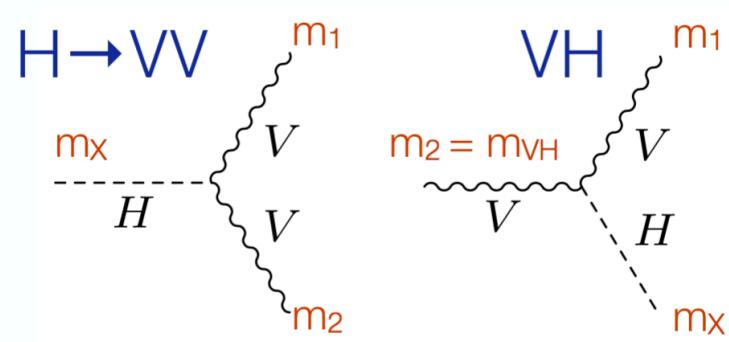
HVV coupling: VHbb

- Generic HVV scattering amplitudes for spin 0 H:

$$A_{00} = -\frac{m_x^2}{v} \left(a_1 \sqrt{1+x} + a_2 \frac{m_1 m_2}{m_x^2} x \right),$$

$$A_{++} = \frac{m_x^2}{v} \left(a_1 + i a_3 \frac{m_1 m_2}{m_x^2} \sqrt{x} \right),$$

$$A_{--} = \frac{m_x^2}{v} \left(a_1 - i a_3 \frac{m_1 m_2}{m_x^2} \sqrt{x} \right),$$



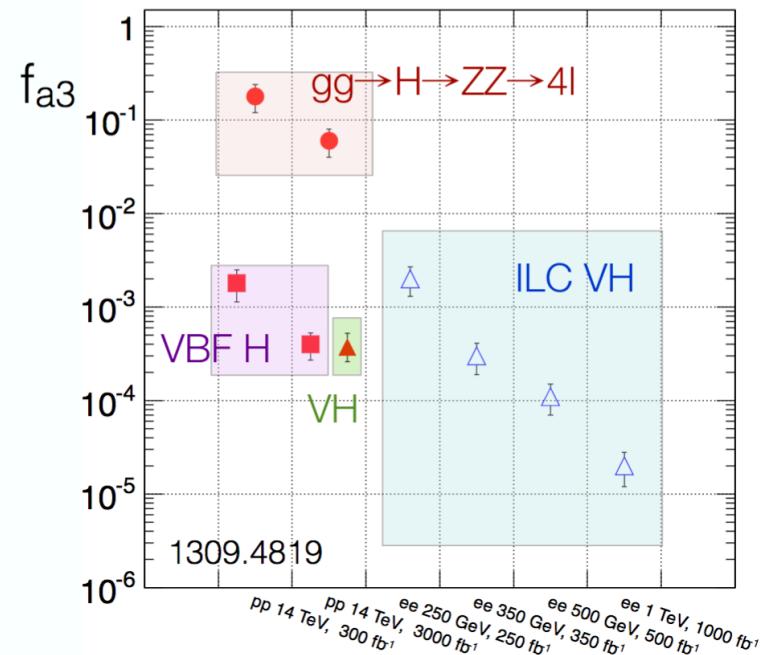
- VH sensitivity from order of magnitude larger $m_1 m_2 / m_x^2$
- Effective fractional cross section to parameterize presence of anomalous pseudoscalar component

$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_3|^2 \sigma_3}$$

where σ_i is the cross section for $a_i=1$ and $a_{j \neq i}=0$

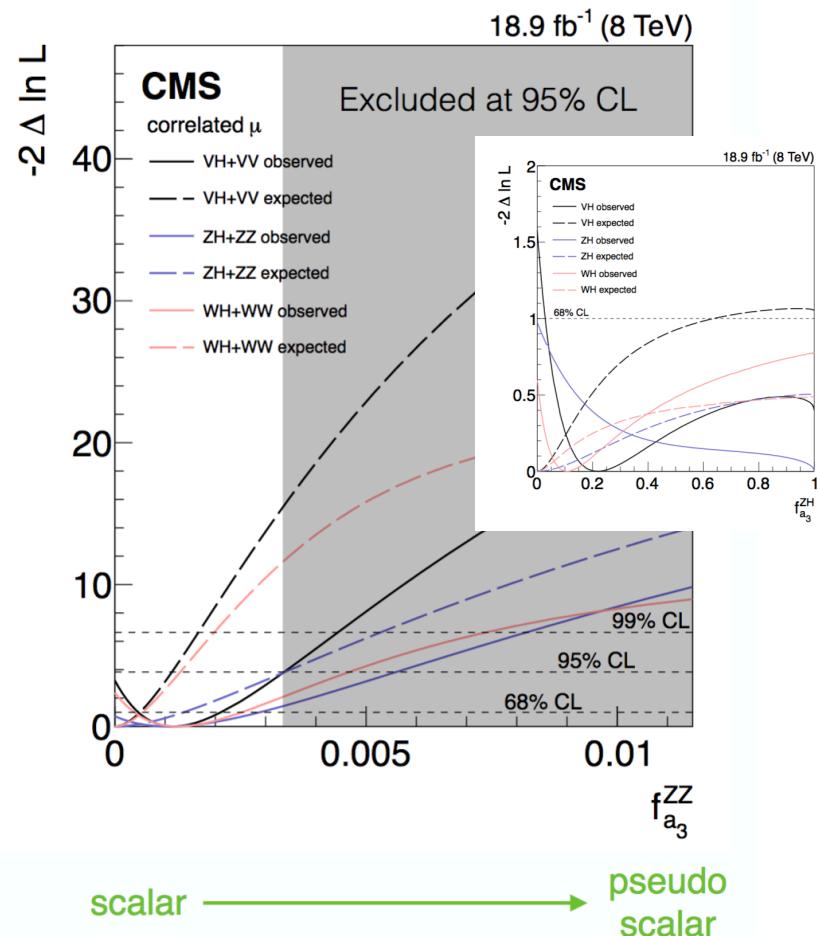
$f_{a3} = 0$ for pure a_1 (SM scalar)

$f_{a3} = 1$ for pure a_3 (pseudoscalar)



HVV coupling: VHbb

- Individual and combined
- Combine with $H \rightarrow VV$
 - Strongest assumption: $a_i^{HZZ} = a_i^{HWW}$ and no additional anomalous couplings, $f_{a_3}^{ZZ} > 0.0034$ excluded at 95% CL
 - Big improvement on $f_{a_3}^{ZZ} > 0.28$ exclusion from $H \rightarrow VV$ alone!



Submitted to PLB in Feb.

HVV anomalous couplings

$$A(HVV) \sim \left[a_1 - e^{i\phi_{\Lambda Q}} \frac{(q_{V1} + q_{V2})^2}{\Lambda_Q^2} - e^{i\phi_{\Lambda 1}} \frac{(q_{V1}^2 + q_{V2}^2)}{\Lambda_1^2} \right] m_V^2 \epsilon_{V1}^* \epsilon_{V2}^*$$

$$+ a_2 f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu} + a_3 f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu}$$

→ Any anomalous coupling can be described with an effective onshell cross sectional fraction and a phase

$$f_{ai} = \frac{|a_i|^2 \sigma_i}{\sum_j |a_j|^2 \sigma_j} \quad \phi_{ai} = \tan^{-1}(a_i/a_1)$$

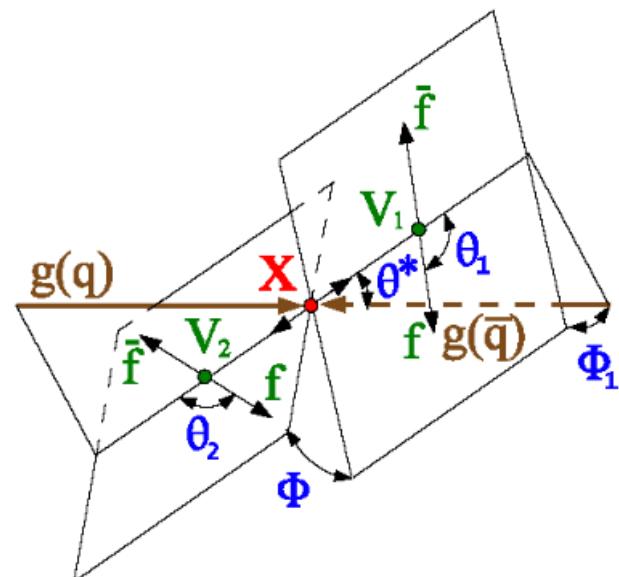
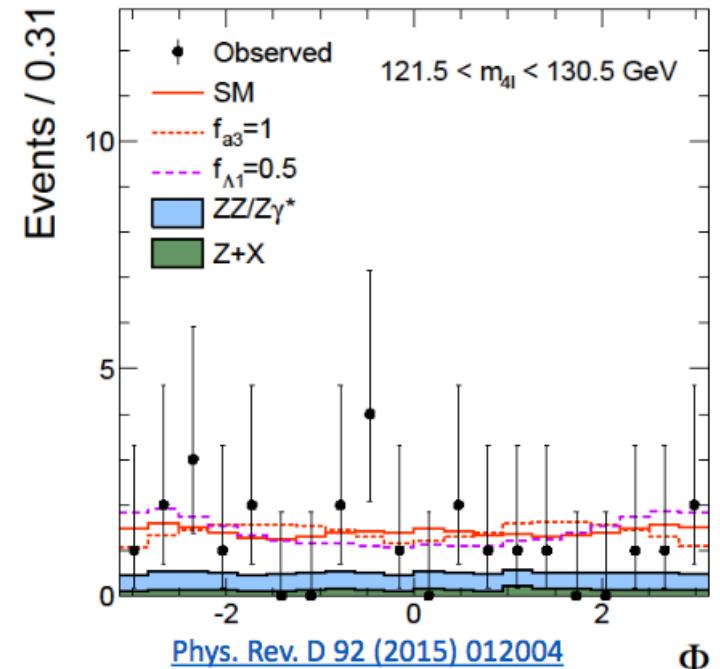
→ $f_{\Lambda Q}$ observable only from offshell, others measurable from either onshell or offshell

$$f_{\Lambda Q} = \frac{m_H^4 / \Lambda_Q^4}{|a_1|^2 + m_H^4 / \Lambda_Q^4}$$

In $4l$, use discriminants based on MELA from MCFM/JHUGen:

$$D_{A vs B} = \frac{P_A}{P_A + P_B}$$

$$D_{A-B \text{ int.}} = \frac{P_{A+B} - P_A - P_B}{P_A + P_B}$$



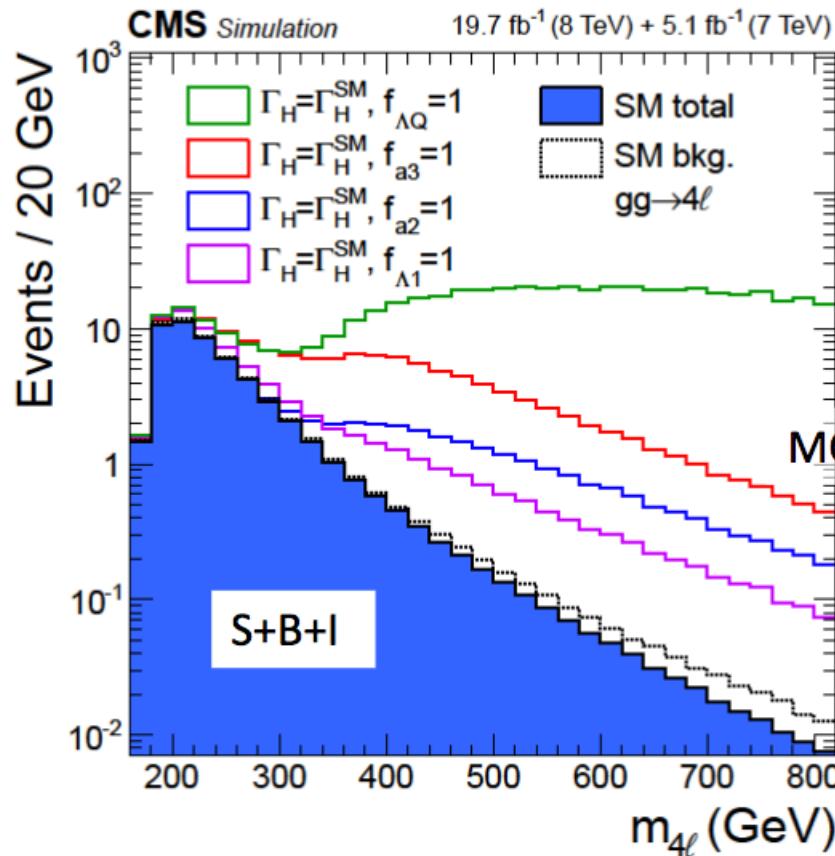
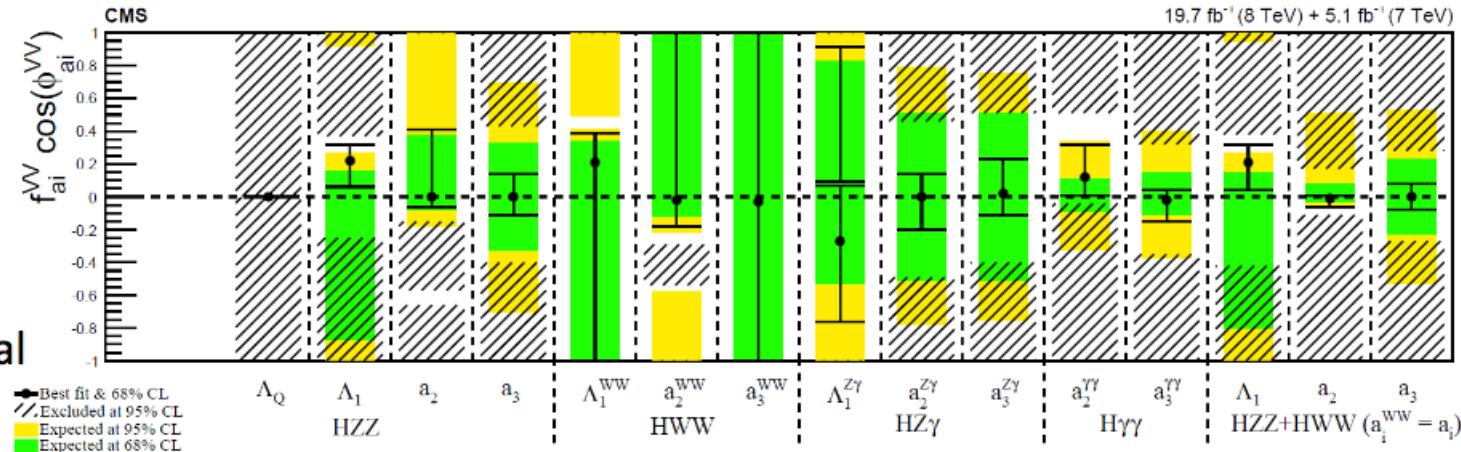
Run I: Anomalous couplings combination

[Phys. Rev. D 92 \(2015\) 012004](#)

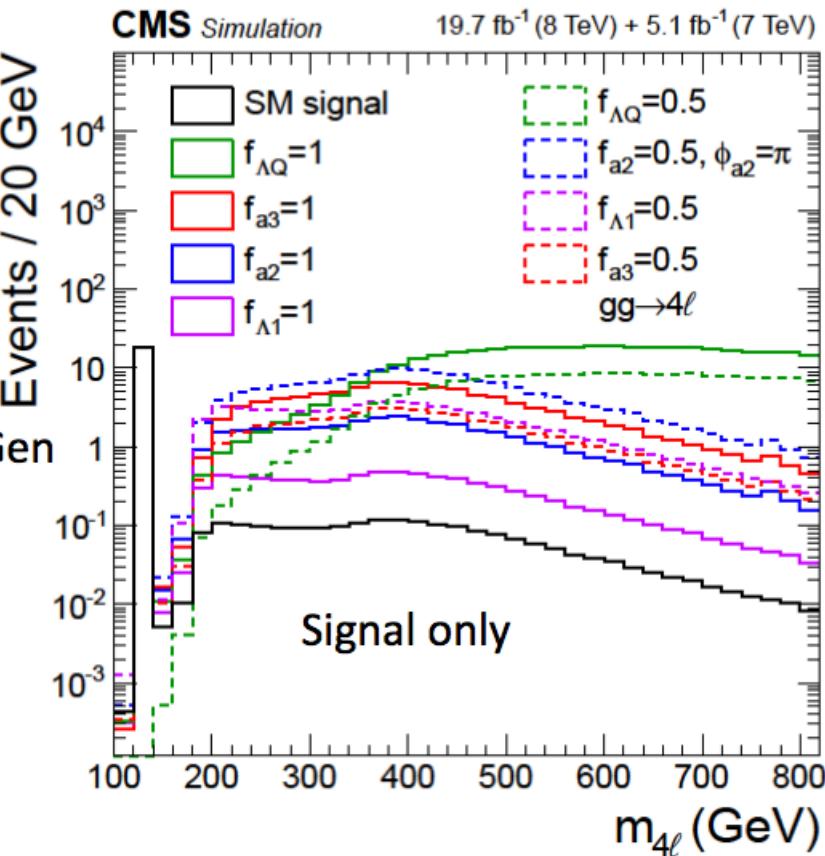
$4l + WW \rightarrow 2l2\nu \rightarrow$

[Phys. Rev. D 92 \(2015\) 072010](#)

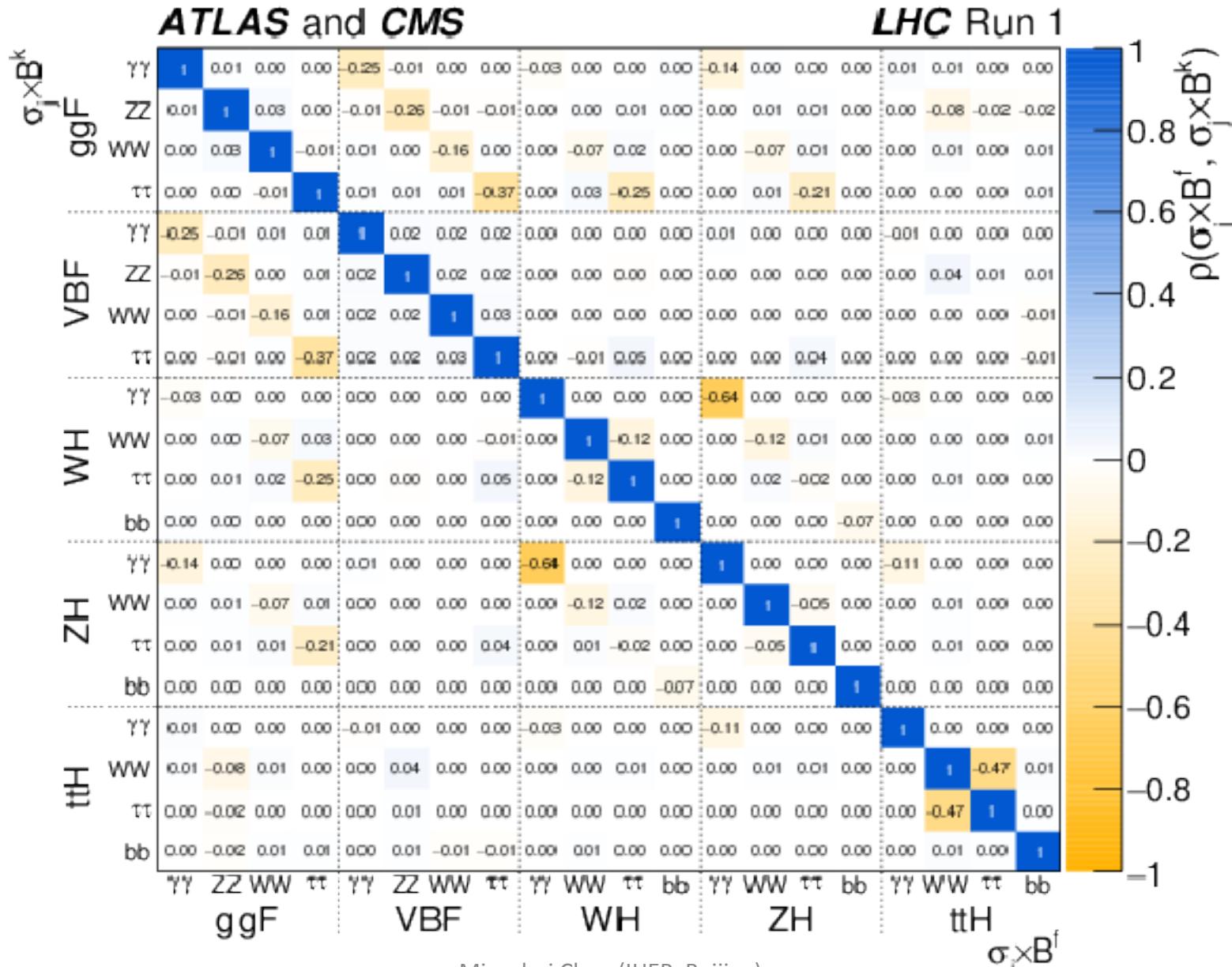
Offshell enhanced for
the same onshell signal
strength ↓



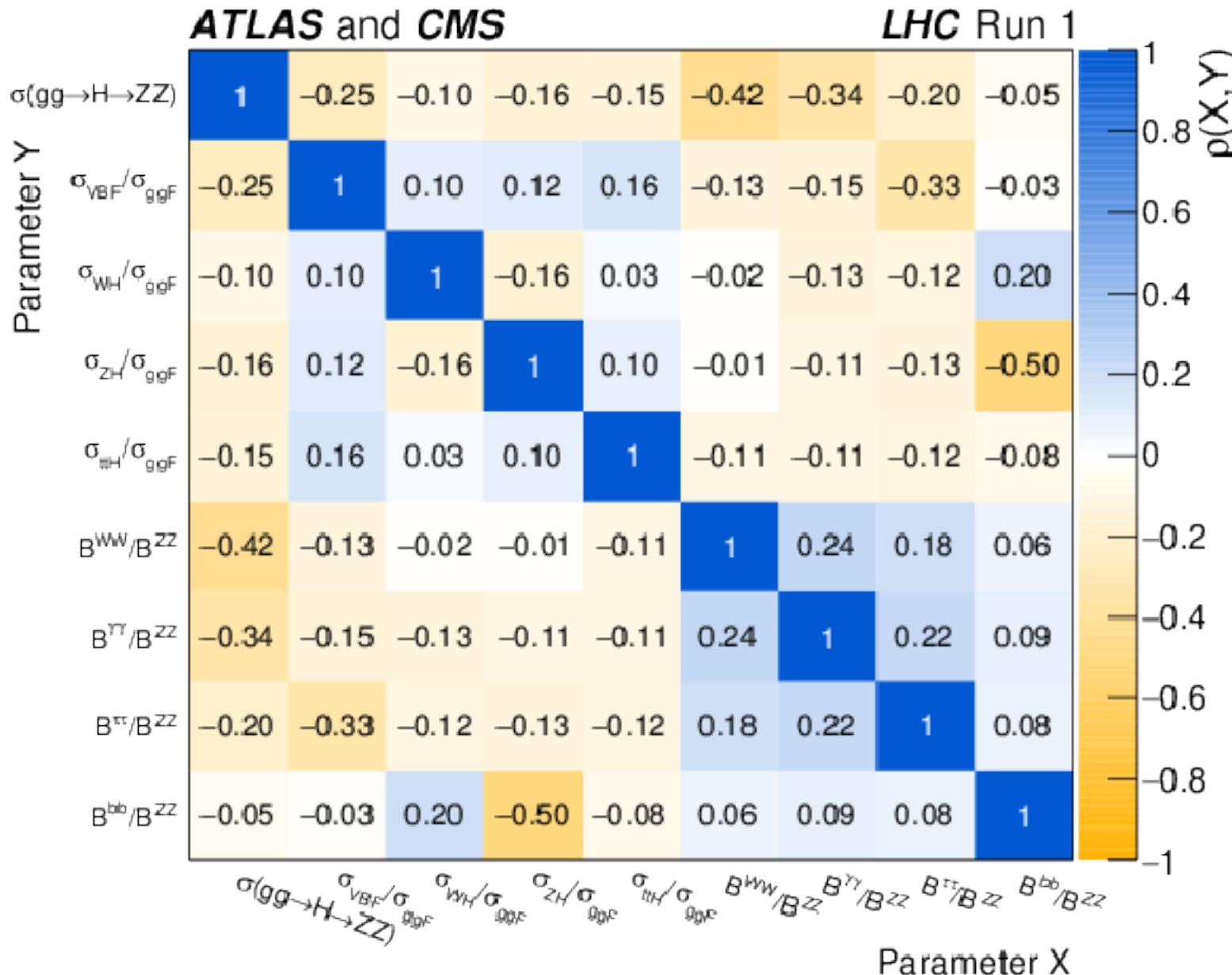
Using
MCFM/JHUGen



Correlation Matrix



Correlation Matrix

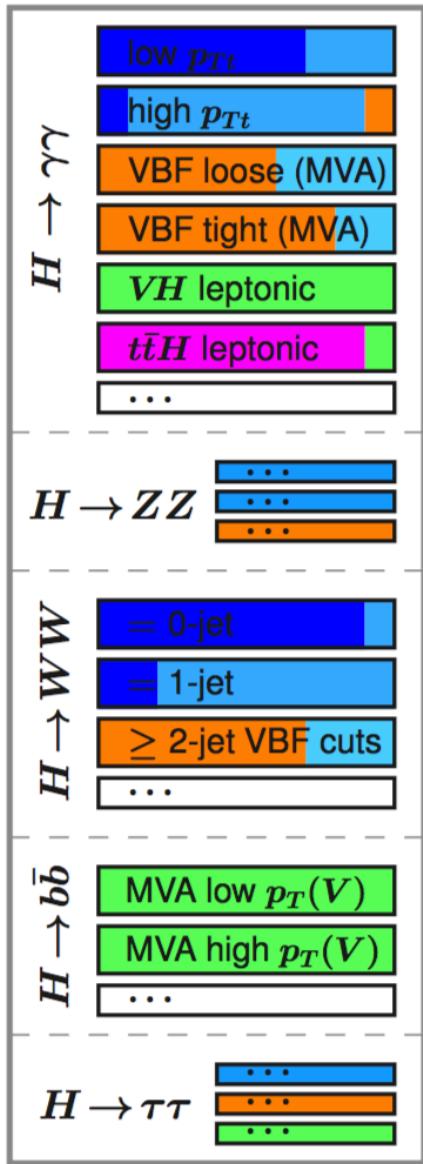


Coupling Parameterization

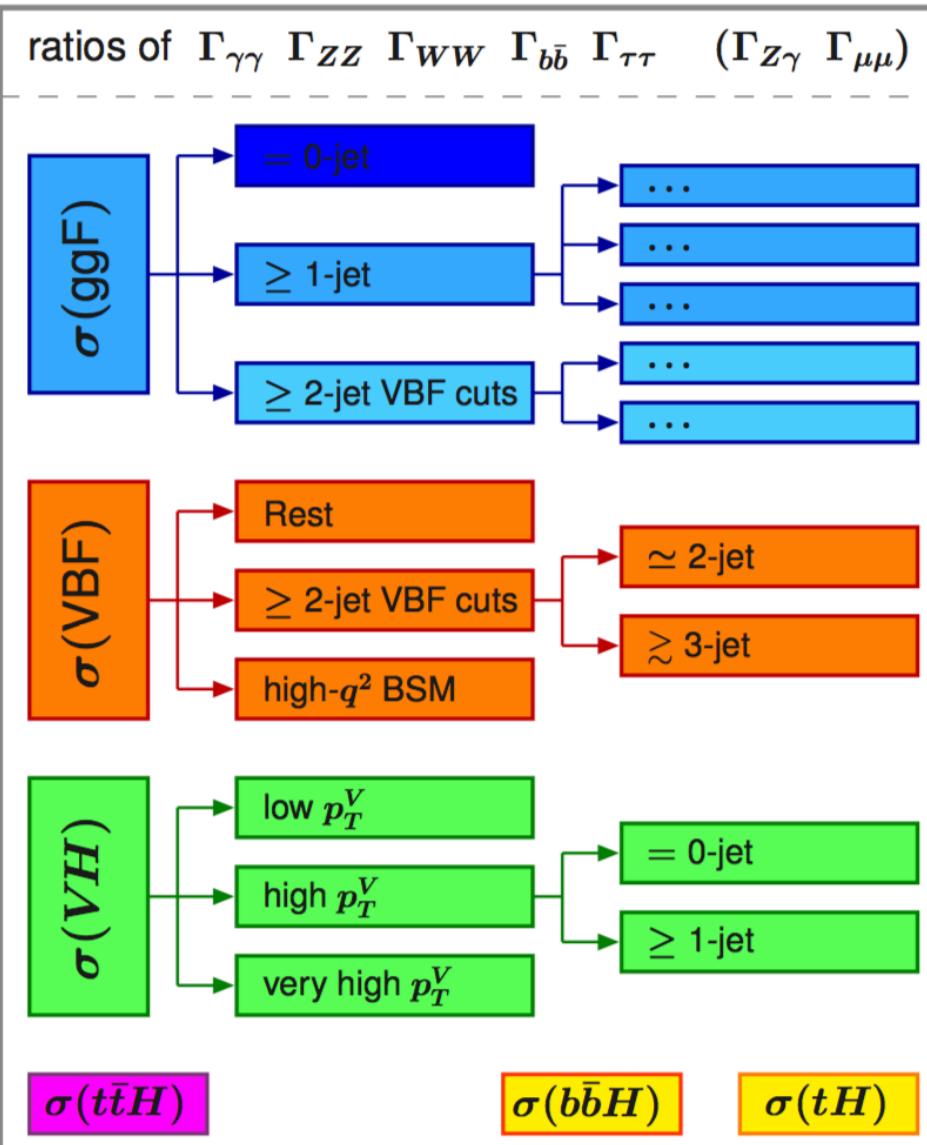
Production	Loops	Interference	Effective	Resolved
			scaling factor	scaling factor
$\sigma(ggF)$	✓	$t-b$	κ_g^2	$1.06 \cdot \kappa_t^2 + 0.01 \cdot \kappa_b^2 - 0.07 \cdot \kappa_t \kappa_b$
$\sigma(VBF)$	—	—		$0.74 \cdot \kappa_W^2 + 0.26 \cdot \kappa_Z^2$
$\sigma(WH)$	—	—		κ_W^2
$\sigma(qq/qg \rightarrow ZH)$	—	—		κ_Z^2
$\sigma(gg \rightarrow ZH)$	✓	$t-Z$		$2.27 \cdot \kappa_Z^2 + 0.37 \cdot \kappa_t^2 - 1.64 \cdot \kappa_Z \kappa_t$
$\sigma(ttH)$	—	—		κ_t^2
$\sigma(gb \rightarrow tHW)$	—	$t-W$		$1.84 \cdot \kappa_t^2 + 1.57 \cdot \kappa_W^2 - 2.41 \cdot \kappa_t \kappa_W$
$\sigma(qq/qb \rightarrow tHq)$	—	$t-W$		$3.40 \cdot \kappa_t^2 + 3.56 \cdot \kappa_W^2 - 5.96 \cdot \kappa_t \kappa_W$
$\sigma(bbH)$	—	—		κ_b^2
Partial decay width				
Γ^{ZZ}	—	—		κ_Z^2
Γ^{WW}	—	—		κ_W^2
$\Gamma^{\gamma\gamma}$	✓	$t-W$	κ_γ^2	$1.59 \cdot \kappa_W^2 + 0.07 \cdot \kappa_t^2 - 0.66 \cdot \kappa_W \kappa_t$
$\Gamma^{\tau\tau}$	—	—		κ_τ^2
Γ^{bb}	—	—		κ_b^2
$\Gamma^{\mu\mu}$	—	—		κ_μ^2
Total width ($B_{BSM} = 0$)				
Γ_H	✓	—	κ_H^2	$0.57 \cdot \kappa_b^2 + 0.22 \cdot \kappa_W^2 + 0.09 \cdot \kappa_g^2 +$ $0.06 \cdot \kappa_\tau^2 + 0.03 \cdot \kappa_Z^2 + 0.03 \cdot \kappa_c^2 +$ $0.0023 \cdot \kappa_\gamma^2 + 0.0016 \cdot \kappa_{(Z\gamma)}^2 +$ $0.0001 \cdot \kappa_s^2 + 0.00022 \cdot \kappa_\mu^2$

Simplified Template Cross Section Framework.

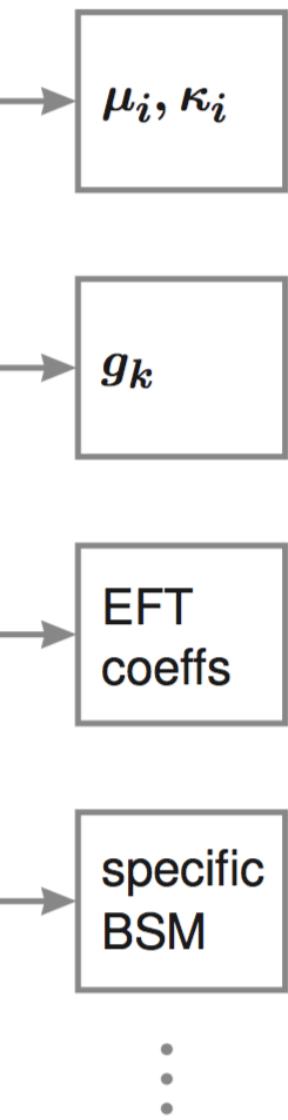
Analysis categories



Simplified Template Cross Sections



Lagrangian parameters



Advantages of FXS and STXS

Fiducial: Optimized for maximal theory independence

- Minimize acceptance corrections
- Simple (rectangular) signal cuts
- “Exact” fiducial volume
- Fiducial in Higgs decay
- Targeted object definitions

Agnostic to production modes

(Single-)differential distributions
(overlapping events)

Only $H \rightarrow \gamma\gamma, ZZ, (WW)$
(by default no combination of channels)

Simplified: Maximize sensitivity while reducing theory dependence

- Allow larger acceptance corrections
- Allow event categories, MVAs, ...
- Abstracted/simplified fiducial volumes
- Inclusive in Higgs decay
- Common idealized object definitions

Xsec split by production mode

Xsec split into mutually exclusive regions of phase space

Explicitly designed for combination of all decay channels