

Status of Calculation of $\alpha_{\text{QED}}(M_Z^2)$ and muon $g - 2$

Daisuke Nomura (Nat. Inst. Tech., Kagawa)

talk at 'Precision Theory for Precise Measurements'

at Quy Nhon, Vietnam

September 26, 2016

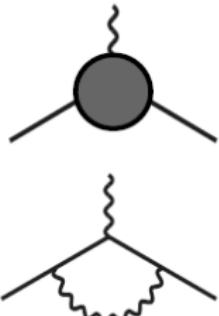
Partially based on:

- K. Hagiwara, A. Keshavarzi, A. D. Martin, DN & T. Teubner,
work in progress
- K. Hagiwara, R. Liao, A. D. Martin, DN & T. Teubner
(HLMNT), J. Phys. **G38** (2011) 085003

Muon $g - 2$: introduction

Lepton magnetic moment $\vec{\mu}$:

$$\boxed{\vec{\mu} = -g \frac{e}{2m} \vec{s}}, \quad (\vec{s} = \frac{1}{2} \vec{\sigma} \text{ (spin)}), \quad g = 2 + 2F_2(0))$$



where

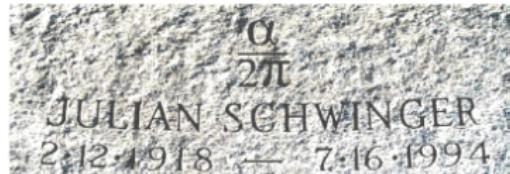
$$\bar{u}(p+q)\Gamma^\mu u(p) = \bar{u}(p+q) \left(\gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu} q_\nu}{2m} F_2(q^2) \right) u(p)$$

Anomalous magnetic moment: $a \equiv (g - 2)/2 (= F_2(0))$

Historically,

- ★ $g = 2$ (tree level, Dirac)
- ★ $a = \alpha/(2\pi)$ (1-loop QED, Schwinger)

Today, still important, since...



- ★ One of the **most precisely measured** quantities:

$$\boxed{a_\mu^{\text{exp}} = 11\ 659\ 208.9(6.3) \times 10^{-10} \quad [0.5\text{ppm}] \quad (\text{Bennett et al})}$$

- ★ **Extremely useful** in probing/constraining physics beyond the SM

Introduction: Standard Model prediction for muon $g - 2$

QED contribution 11 658 471.808 (0.015) Kinoshita & Nio, Aoyama et al

EW contribution 15.4 (0.2) Czarnecki et al

Hadronic contributions

LO hadronic 694.9 (4.3) HLMNT11

NLO hadronic -9.8 (0.1) HLMNT11

light-by-light 10.5 (2.6) Prades, de Rafael & Vainshtein

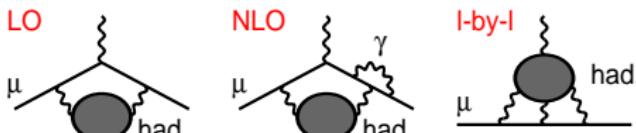
Theory TOTAL 11 659 182.8 (4.9)

Experiment 11 659 208.9 (6.3) world avg

Exp – Theory 26.1 (8.0) 3.3 σ discrepancy

(in units of 10^{-10} . Numbers taken from HLMNT11, arXiv:1105.3149)

n.b.: hadronic contributions:



Improvements in the past few years (1)

- **QED contribution**

5-loop calculation completed:

now $a_\mu(\text{QED}) = 11\ 658\ 471.895(08) \times 10^{-10}$

was $11\ 658\ 471.808(15) \times 10^{-10}$

(numbers are from Aoyama et al (2012, 2007))

- **EW contribution**

Higgs-boson mass just fixed:

now $a_\mu(\text{EW}) = 15.4(1) \times 10^{-10}$ Gnendiger et al '13

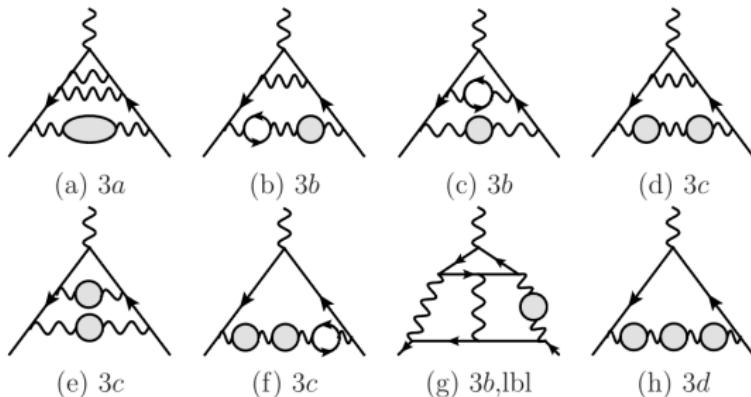
was $15.4(2) \times 10^{-10}$ Czarnecki et al '02

cf: $a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = \begin{cases} (26.1 \pm 8.0) \times 10^{-10} & \text{HLMNT11} \\ (28.7 \pm 8.0) \times 10^{-10} & \text{Davier et al '10} \end{cases}$

Improvements in the past few years (2)

- NNLO hadronic contribution

$$a_\mu(\text{had, NNLO}) = (1.24 \pm 0.01) \times 10^{-10}$$



Numbers and Figs. from
A. Kurz et al, arXiv:1403.6400

$$3a: +0.80 \times 10^{-10}$$

$$3b: -0.41 \times 10^{-10}$$

$$3b,\text{lbl}: +0.91 \times 10^{-10}$$

$$3c: -0.06 \times 10^{-10}$$

$$3d: +0.0005 \times 10^{-10}$$

Figure 2: Sample NNLO Feynman diagrams contributing to a_μ^{had} .

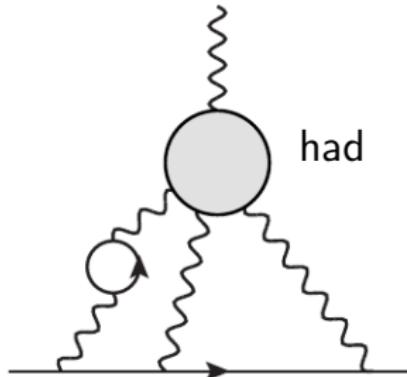
$$\text{cf: } a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = \begin{cases} (26.1 \pm 8.0) \times 10^{-10} & \text{HLMNT11} \\ (28.7 \pm 8.0) \times 10^{-10} & \text{Davier et al '10} \end{cases}$$

$$a_\mu(\text{had, NLO}) = (-9.8 \pm 0.1) \times 10^{-10} \quad \text{HLMNT11}$$

Improvements in the past few years (3)

- **Ibyl-NLO hadronic contribution**

$$a_\mu(\text{had, Ibyl-NLO}) = (0.3 \pm 0.2) \times 10^{-10}$$



Number and Fig. from Colangelo et al,
arXiv:1403.7512

⇒ **Negligible.** But it is always good
to confirm that higher order terms
are really negligible.

cf: $a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = \begin{cases} (26.1 \pm 8.0) \times 10^{-10} & \text{HLMNT11} \\ (28.7 \pm 8.0) \times 10^{-10} & \text{Davier et al '10} \end{cases}$

$$a_\mu^{\text{had,Ibyl}} = \begin{cases} (10.5 \pm 2.6) \times 10^{-10} & \text{Prades et al, '09} \\ (11.6 \pm 3.9) \times 10^{-10} & \text{Jegerlehner+Nyffeler, '09} \end{cases}$$

HLbL scattering: Summary of selected results for $a_\mu^{\text{HLbL}} \times 10^{11}$

Contribution	BPP	HKS, HK	KN	MV	BP, MdRR	PdRV	N, JN
π^0, η, η'	85 ± 13	82.7 ± 6.4	83 ± 12	114 ± 10	—	114 ± 13	99 ± 16
axial vectors	2.5 ± 1.0	1.7 ± 1.7	—	22 ± 5	—	15 ± 10	22 ± 5
scalars	-6.8 ± 2.0	—	—	—	—	-7 ± 7	-7 ± 2
π, K loops	-19 ± 13	-4.5 ± 8.1	—	—	—	-19 ± 19	-19 ± 13
π, K loops +subl. N_C	—	—	—	0 ± 10	—	—	—
quark loops	21 ± 3	9.7 ± 11.1	—	—	—	2.3 (c-quark)	21 ± 3
Total	83 ± 32	89.6 ± 15.4	80 ± 40	136 ± 25	110 ± 40	105 ± 26	116 ± 39

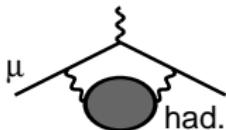
BPP = Bijnens, Pallante, Prades '95, '96, '02; HKS = Hayakawa, Kinoshita, Sanda '95, '96; HK = Hayakawa, Kinoshita '98, '02; KN = Knecht, AN '02; MV = Melnikov, Vainshtein '04; BP = Bijnens, Prades '07; MdRR = Miller, de Rafael, Roberts '07; PdRV = Prades, de Rafael, Vainshtein '09; N = AN '09, JN = Jegerlehner, AN '09

- **Pseudoscalar-exchanges dominate numerically.** Other contributions not negligible. **Cancellation** between π, K -loops and quark loops !
- Note that recent reevaluations of axial vector contribution lead to much smaller estimates than in MV: $a_\mu^{\text{HLbL; axial}} = (8 \pm 3) \times 10^{-11}$ (Pauk, Vanderhaeghen '14; Jegerlehner '14, '15). This would shift central values of compilations downwards: $a_\mu^{\text{HLbL}} = (98 \pm 26) \times 10^{-11}$ (PdRV) and $a_\mu^{\text{HLbL}} = (102 \pm 39) \times 10^{-11}$ (N, JN).
- **PdRV:** Analyzed results obtained by different groups with various models and suggested new estimates for some contributions (shifted central values, enlarged errors). **Do not consider dressed light quark loops as separate contribution. Added all errors in quadrature !**
- **N, JN:** New evaluation of pseudoscalar exchange contribution imposing new short-distance constraint on off-shell form factors. Took over most values from BPP, except axial vectors from MV. **Added all errors linearly.**

A. Nyffeler, talk at Frascati, May 2016

Introduction for $a_\mu^{\text{had},\text{LO}}$

The diagram to be evaluated:

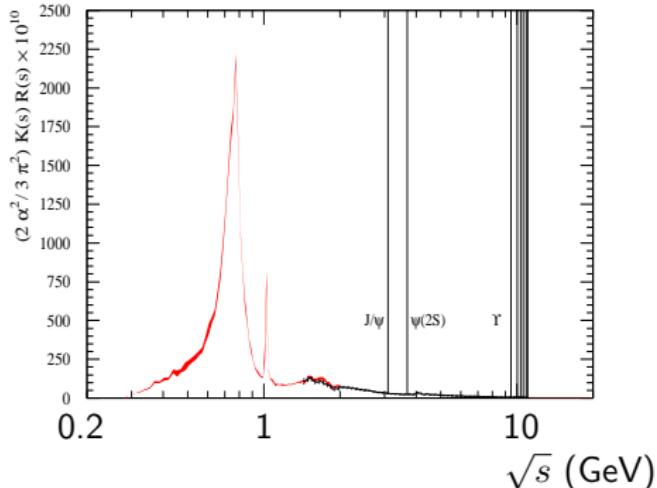


pQCD not useful. Use the **dispersion relation** and the **optical theorem**.

$$\text{had.} = \int \frac{ds}{\pi(s-q^2)} \text{Im } \text{had.}$$

$$2 \text{Im } \text{had.} = \sum_{\text{had.}} \int d\Phi \left| \text{had.} \right|^2$$

$$a_\mu^{\text{had},\text{LO}} = \frac{m_\mu^2}{12\pi^3} \int_{s_{\text{th}}}^\infty ds \frac{1}{s} \hat{K}(s) \sigma_{\text{had}}(s)$$



- Weight function $\hat{K}(s)/s = \mathcal{O}(1)/s$
- ⇒ **Lower energies more important**
- ⇒ $\pi^+\pi^-$ channel: 73% of total $a_\mu^{\text{had},\text{LO}}$

Included Hadronic Final States

Channel	Experiments with References
$\pi^+ \pi^-$	OLYA [16, 17, 18], OLYA-TOF [19], NA7 [20], OLYA and CMD [21, 22], DMI [23], DM2 [24], BCF [25, 26], MEA [27, 28], ORSAY-ACO [29], CMD-2 [10, 11, 30]
$\pi^0 \gamma$	SND [31, 32]
$\eta \gamma$	SND [32, 33], CMD-2 [34, 35, 36]
$\pi^+ \pi^- \pi^0$	ND [22], DMI [37], DM2 [38], CMD-2 [10, 13, 34, 39], SND [40, 41], CMD [42]
$K^+ K^-$	MEA [27], OLYA [43], BCF [26], DMI [44], DM2 [45, 46], CMD [22], CMD-2 [34], SND [47]
$K_S^0 K_L^0$	DMI [48], CMD-2 [10, 14, 49], SND [47]
$\pi^+ \pi^- \pi^0 \pi^0$	M3N [50], DM2 [51], OLYA [52], CMD-2 [53], SND [54], ORSAY-ACO [55], $\gamma\gamma 2$ [56], MEA [57]
$\omega (\rightarrow \pi^0 \gamma) \pi^0$	ND and ARGUS [22], DM2 [51], CMD-2 [53, 58], SND [59, 60], ND [61]
$\pi^+ \pi^- \pi^+ \pi^-$	ND [22], M3N [50], CMD [62], DMI [63, 64], DM2 [51], OLYA [65], $\gamma\gamma 2$ [66], CMD-2 [53, 67, 68], SND [54], ORSAY-ACO [55]
$\pi^+ \pi^- \pi^+ \pi^- \pi^0$	MEA [57], M3N [50], CMD [22, 62], $\gamma\gamma 2$ [56]
$\pi^+ \pi^- \pi^0 \pi^0 \pi^0$	M3N [50]
$\omega (\rightarrow \pi^0 \gamma) \pi^+ \pi^-$	DM2 [38], CMD-2 [69], DMI [70]
$\pi^+ \pi^- \pi^+ \pi^- \pi^+ \pi^-$	M3N [50], CMD [62], DMI [71], DM2 [72]
$\pi^+ \pi^- \pi^+ \pi^- \pi^0 \pi^0$	M3N [50], CMD [62], DM2 [72], $\gamma\gamma 2$ [56], MEA [57]
$\pi^+ \pi^- \pi^0 \pi^0 \pi^0 \pi^0$	isospin-related
$\eta \pi^+ \pi^-$	DM2 [73], CMD-2 [69]
$K^+ K^- \pi^0$	DM2 [74, 75]
$K_S^0 \pi K$	DMI [76], DM2 [74, 75]
$K_S^0 X$	DMI [77]
$\pi^+ \pi^- K^+ K^-$	DM2 [74]
$p\bar{p}$	FENICE [78, 79], DM2 [80, 81], DMI [82]
$n\bar{n}$	FENICE [78, 83]
incl. (< 2 GeV)	$\gamma\gamma 2$ [84], MEA [85], M3N [86], BARYON-ANTIBARYON [87]
incl. (> 2 GeV)	BES [88, 89], Crystal Ball [90, 91, 92], LENA [93], MD-1 [94], DASP [95], CLEO [96], CUSB [97], DHHM [98]

channel	inclusive (1.43,2 GeV) $a_\mu^{\text{had},\text{LO}}$	exclusive (1.43,2 GeV) $a_\mu^{\text{had},\text{LO}}$	inclusive (1.43,2 GeV) $\Delta\alpha_{\text{had}}(M_Z^2)$	exclusive (1.43,2 GeV) $\Delta\alpha_{\text{had}}(M_Z^2)$
$\pi^0 \gamma$ (ChPT)	0.13 ± 0.01	0.00 ± 0.00	0.13 ± 0.01	0.00 ± 0.00
$\pi^0 \gamma$ (data)	4.50 ± 0.15	0.36 ± 0.01	4.50 ± 0.15	0.36 ± 0.01
$\pi^+ \pi^-$ (ChPT)	2.36 ± 0.05	0.04 ± 0.00	2.36 ± 0.05	0.04 ± 0.00
$\pi^+ \pi^-$ (data)	502.78 ± 5.02	34.39 ± 0.29	503.38 ± 5.02	34.59 ± 0.29
$\pi^+ \pi^- \pi^0$ (ChPT)	0.01 ± 0.00	0.00 ± 0.00	0.01 ± 0.00	0.00 ± 0.00
$\pi^+ \pi^- \pi^0$ (data)	46.43 ± 0.90	4.33 ± 0.08	47.04 ± 0.90	4.52 ± 0.08
$\eta \gamma$ (ChPT)	0.00 ± 0.00	0.00 ± 0.00	0.00 ± 0.00	0.00 ± 0.00
$\eta \gamma$ (data)	0.73 ± 0.03	0.09 ± 0.00	0.73 ± 0.03	0.09 ± 0.00
$K^+ K^-$	21.62 ± 0.76	3.01 ± 0.11	22.35 ± 0.77	3.23 ± 0.11
$K_S^0 K_L^0$	13.16 ± 0.31	1.76 ± 0.04	13.30 ± 0.32	1.80 ± 0.04
$2\pi^+ 2\pi^-$	6.16 ± 0.32	1.27 ± 0.07	14.77 ± 0.76	4.04 ± 0.21
$\pi^+ \pi^- 2\pi^0$	9.71 ± 0.63	1.86 ± 0.12	20.55 ± 1.22	5.51 ± 0.35
$2\pi^+ 2\pi^- \pi^0$	0.26 ± 0.04	0.06 ± 0.01	2.85 ± 0.25	0.99 ± 0.09
$\pi^+ \pi^- 3\pi^0$	0.09 ± 0.09	0.02 ± 0.02	1.19 ± 0.33	0.41 ± 0.10
$3\pi^+ \pi^-$	0.00 ± 0.00	0.00 ± 0.00	0.22 ± 0.02	0.09 ± 0.01
$2\pi^+ 2\pi^- 2\pi^0$	0.12 ± 0.03	0.03 ± 0.01	3.32 ± 0.29	1.22 ± 0.11
$\pi^+ \pi^- 4\pi^0$ (isospin)	0.00 ± 0.00	0.00 ± 0.00	0.12 ± 0.12	0.05 ± 0.05
$K^+ K^- \pi^0$	0.00 ± 0.00	0.00 ± 0.00	0.29 ± 0.07	0.10 ± 0.03
$K_S^0 K_L^0 \pi^0$ (isospin)	0.00 ± 0.00	0.00 ± 0.00	0.29 ± 0.07	0.10 ± 0.03
$K_S^0 \pi^\mp K^\pm$	0.05 ± 0.02	0.01 ± 0.00	1.00 ± 0.11	0.33 ± 0.04
$K_L^0 \pi^\mp K^\pm$ (isospin)	0.05 ± 0.02	0.01 ± 0.00	1.00 ± 0.11	0.33 ± 0.04
$K\bar{K} \pi\pi$ (isospin)	0.00 ± 0.00	0.00 ± 0.00	3.63 ± 1.34	1.33 ± 0.48
$\omega (\rightarrow \pi^0 \gamma) \pi^0$	0.64 ± 0.02	0.12 ± 0.00	0.83 ± 0.03	0.17 ± 0.01
$\omega (\rightarrow \pi^0 \gamma) \pi^+ \pi^-$	0.01 ± 0.00	0.00 ± 0.00	0.07 ± 0.01	0.02 ± 0.00
$\eta (\rightarrow \pi^0 \gamma) \pi^+ \pi^-$	0.07 ± 0.01	0.02 ± 0.00	0.49 ± 0.07	0.15 ± 0.02
$\phi (\rightarrow \text{unaccounted})$	0.06 ± 0.06	0.01 ± 0.01	0.06 ± 0.06	0.01 ± 0.01
$p\bar{p}$	0.00 ± 0.00	0.00 ± 0.00	0.04 ± 0.01	0.02 ± 0.00
$n\bar{n}$	0.00 ± 0.00	0.00 ± 0.00	0.07 ± 0.02	0.03 ± 0.01
$J/\psi, \psi'$	7.30 ± 0.43	8.90 ± 0.51	7.30 ± 0.43	8.90 ± 0.51
$\Upsilon [1S - 6S]$	0.10 ± 0.00	1.16 ± 0.04	0.10 ± 0.00	1.16 ± 0.04
inclusive R	73.96 ± 2.68	92.75 ± 1.74	42.05 ± 1.14	81.97 ± 1.53
pQCD	2.11 ± 0.00	125.32 ± 0.15	2.11 ± 0.00	125.32 ± 0.15
sum	692.38 ± 5.88	275.52 ± 1.85	696.15 ± 5.68	276.90 ± 1.77

Table 1: Experiments and references for the $e^+ e^-$ data sets for the different exclusive and the inclusive channels as used in this analysis. The recent re-analysis from CMD-2 [10] supersedes

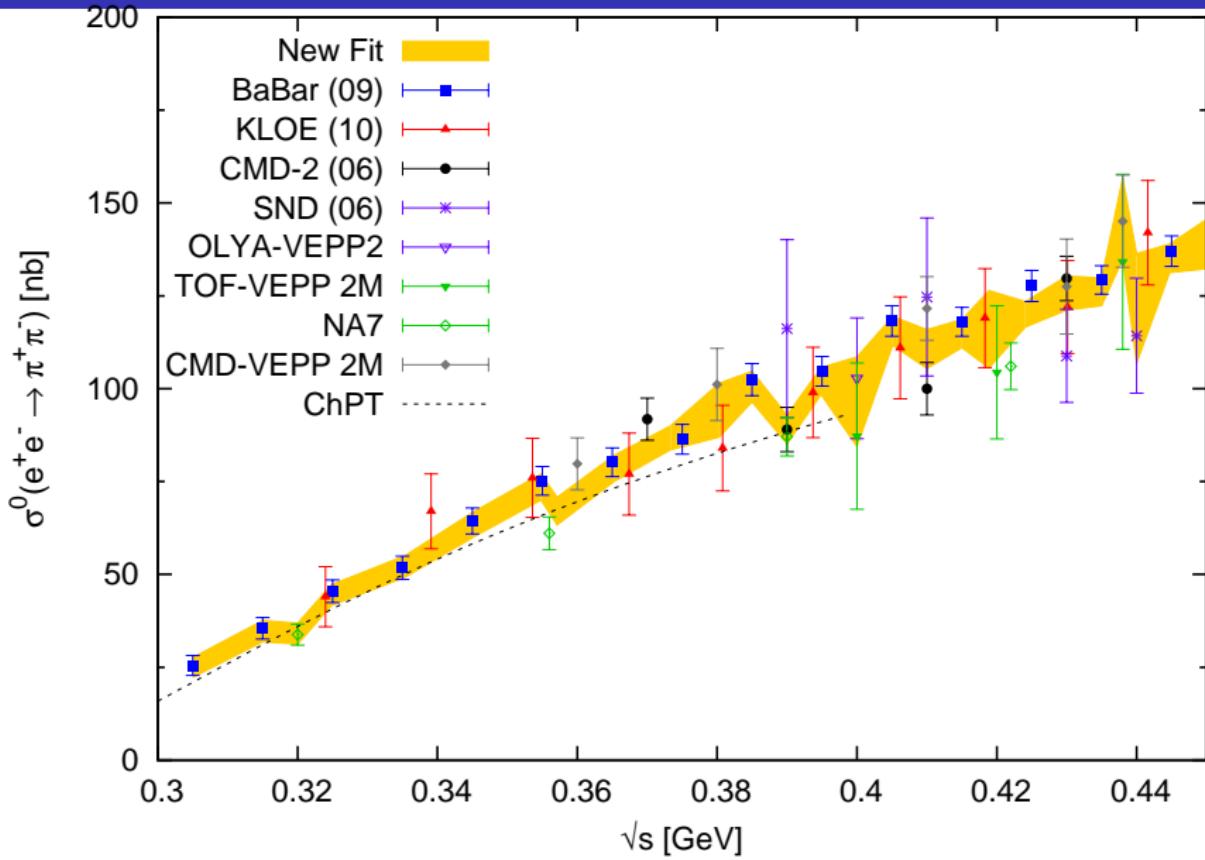
Important Channels

Contributions from various channels to $a_\mu(\text{LO,had})$ for $\sqrt{s} < 1.8\text{GeV}$:

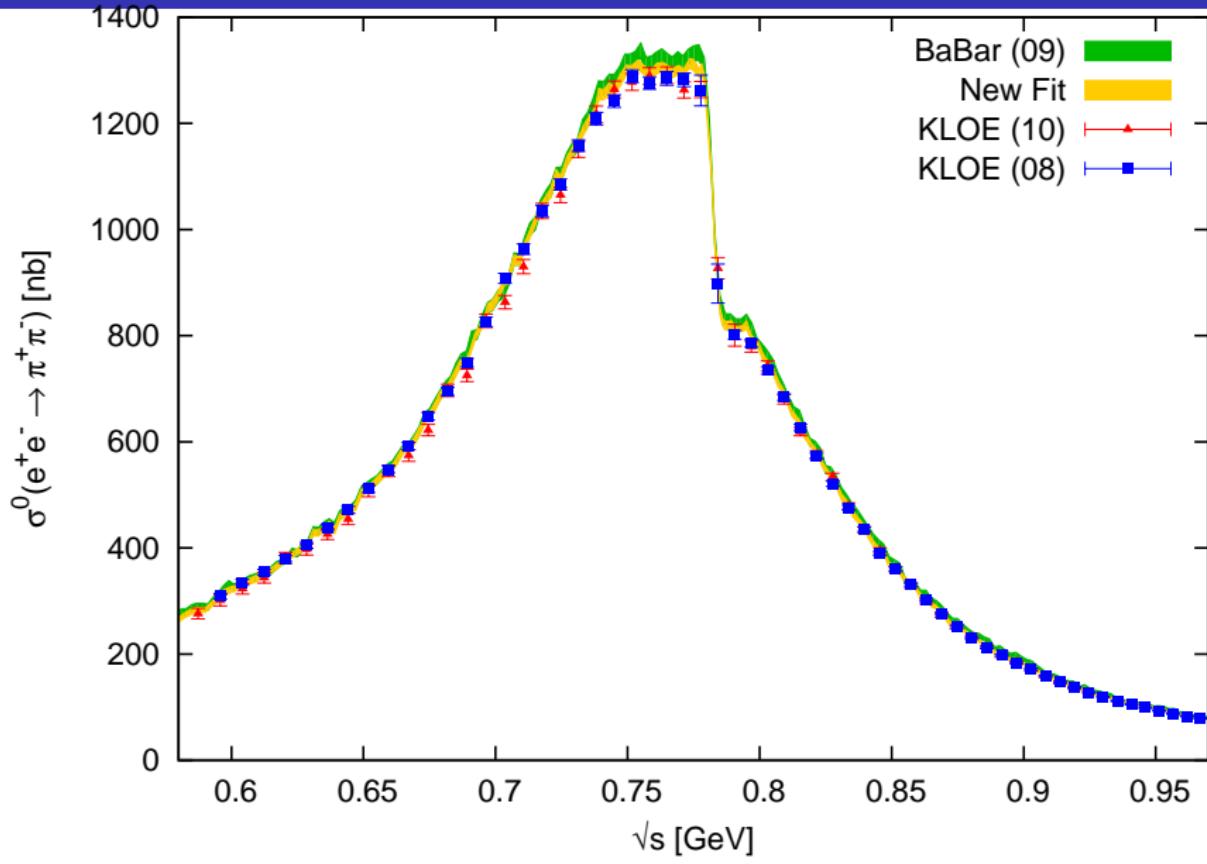
channel	HLMNT11	Davier et al '10	diff
$\pi^+ \pi^-$	505.65 ± 3.09	507.80 ± 2.84	-2.15
$\pi^+ \pi^- \pi^0$	47.38 ± 0.99	46.00 ± 1.48	1.38
$K^+ K^-$	22.09 ± 0.46	21.63 ± 0.73	0.46
$\pi^+ \pi^- 2\pi^0$	18.62 ± 1.15	18.01 ± 1.24	0.61
$2\pi^+ 2\pi^-$	13.50 ± 0.44	13.35 ± 0.53	0.15
$K_S^0 K_L^0$	13.32 ± 0.16	12.96 ± 0.39	0.36
$\pi^0 \gamma$	4.54 ± 0.14	4.42 ± 0.19	0.12
⋮	⋮	⋮	⋮
Sum	634.28 ± 3.53	633.93 ± 3.61	0.35

table taken from HLMNT11

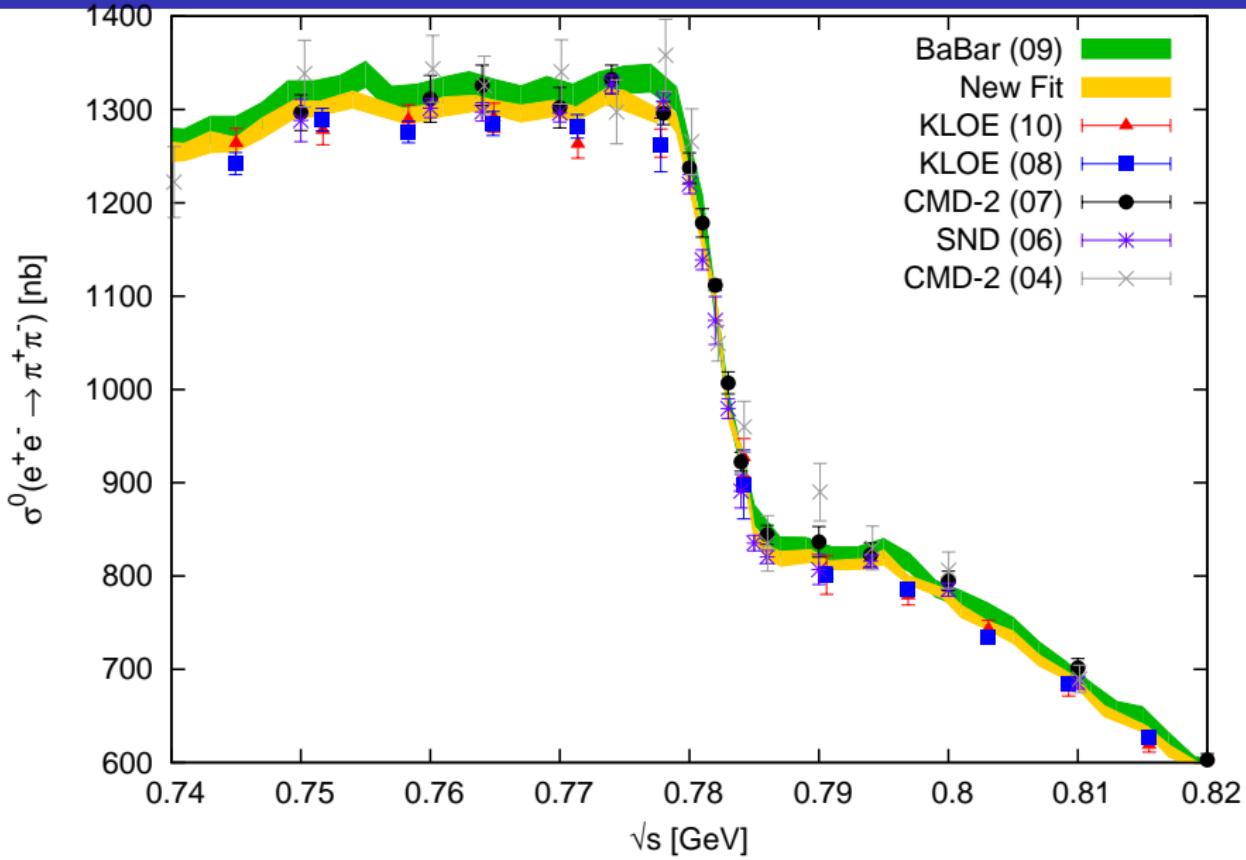
$\pi^+\pi^-$ channel: Low Energy Tail



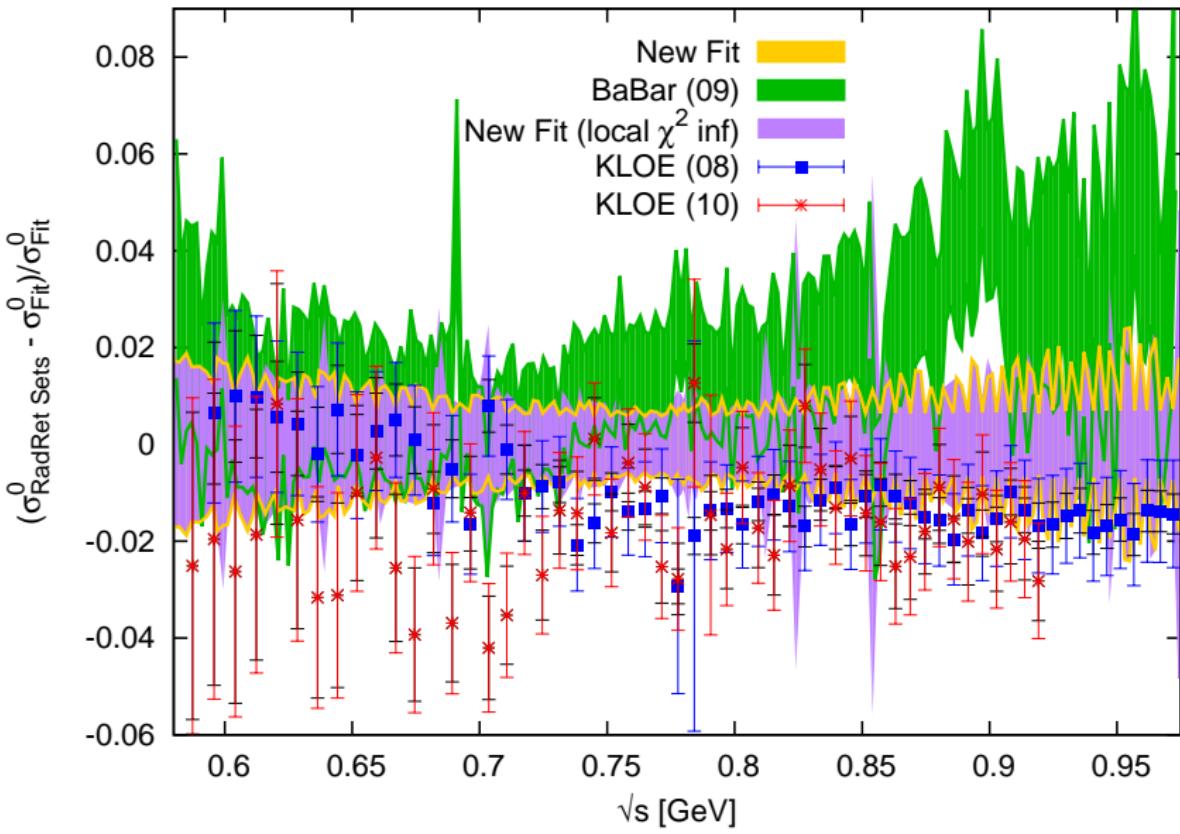
$\pi^+\pi^-$ channel: New Radiative Return Data



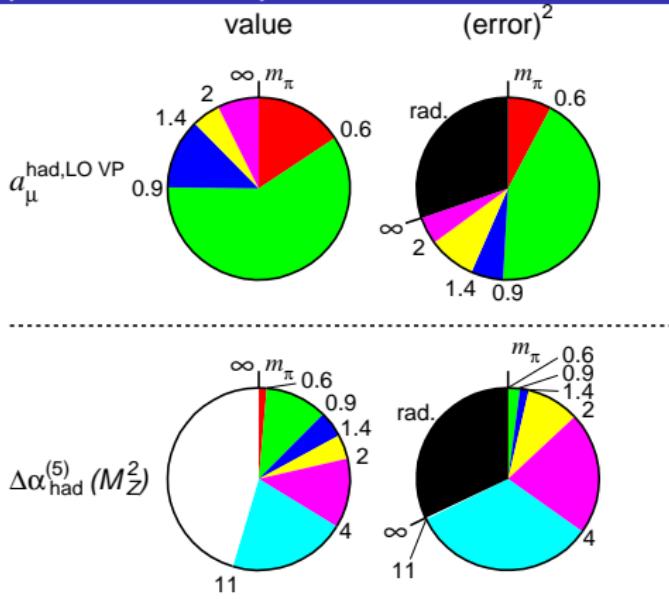
$\pi^+\pi^-$ channel: Zoom-In at ρ - ω Region



Rad. Rtn. Data (for $\pi^+\pi^-$) and Our Combined Result



Results: $a_\mu^{\text{had,LO}}$, $a_\mu^{\text{had,NLO}}$ and $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2)$



$$a_\mu^{\text{had,LO}} = (694.91 \pm 3.72_{\text{exp}} \pm 2.10_{\text{rad}}) \times 10^{-10}$$

$$a_\mu^{\text{had,NLO}} = (-9.84 \pm 0.06_{\text{exp}} \pm 0.04_{\text{rad}}) \times 10^{-10}$$

$$\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = (276.26 \pm 1.16_{\text{exp}} \pm 0.74_{\text{rad}}) \times 10^{-4}$$

Full SM Result and Comparison with Other Groups

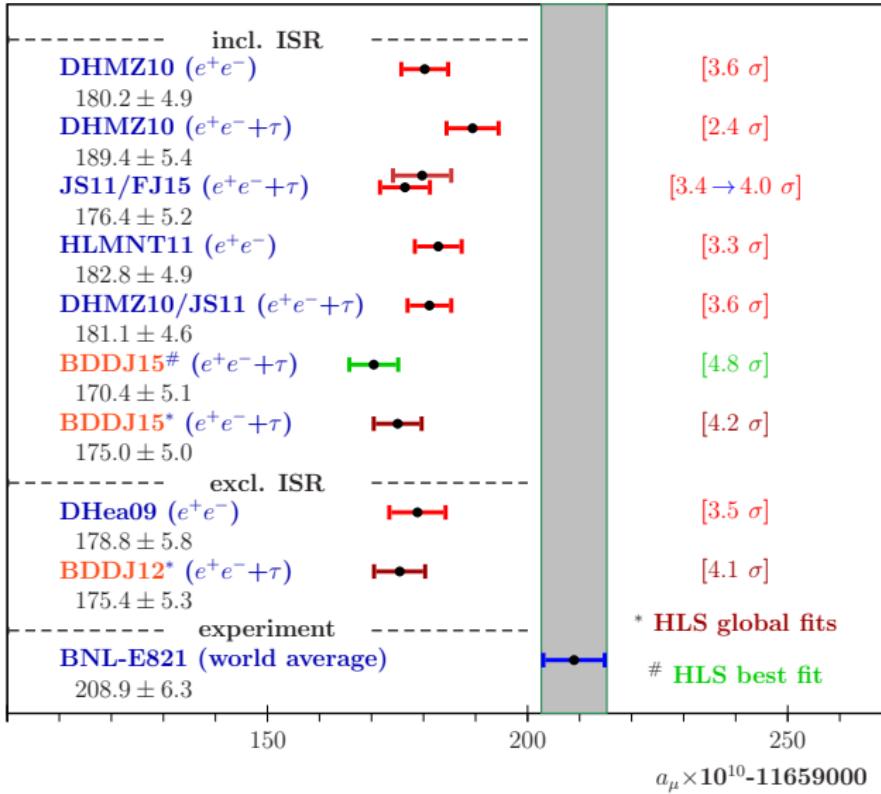
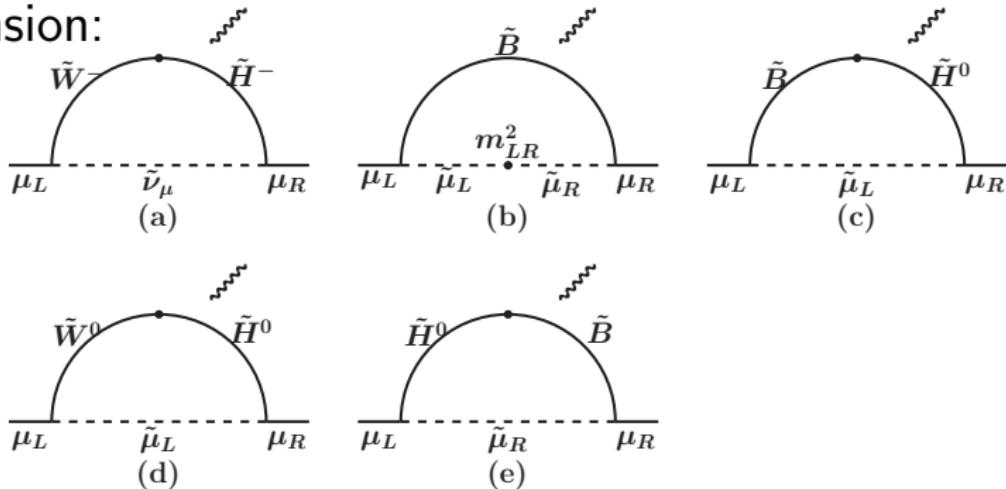


Fig. from F. Jegerlehner, arXiv:1511.04473

SUSY Contributions to Muon $g - 2$

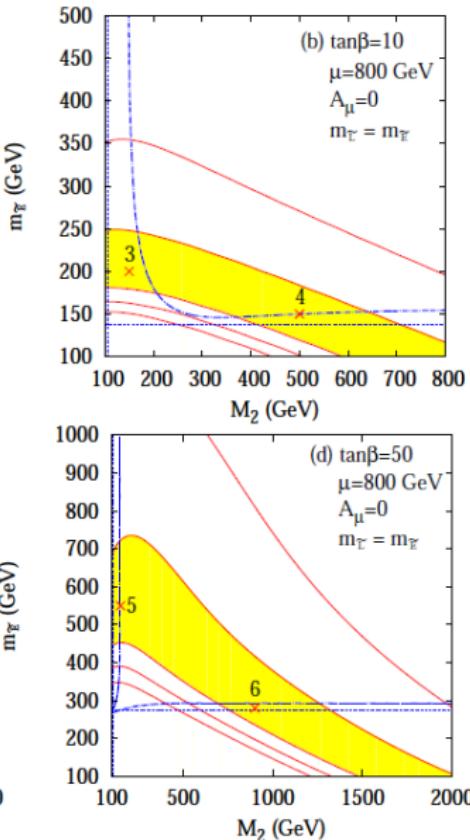
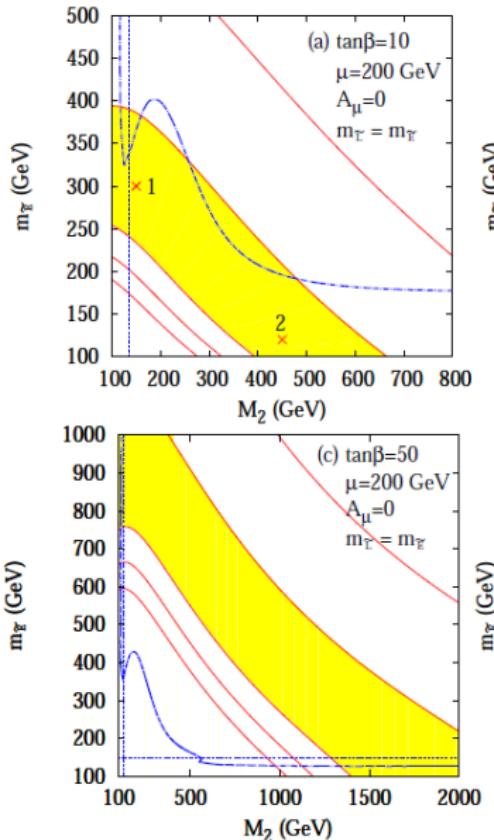
Suppose that the 3.3σ deviation is due to SUSY...

Leading SUSY contributions in the m_Z/m_{SUSY} expansion:



In most cases, the $\tilde{\chi}^\pm$ - $\tilde{\nu}$ diagram (a) and/or the \tilde{B} - $\tilde{\mu}_{L/R}$ diagram (b) dominate. (Lopez-Nanopoulos-Wang, Chattopadhyay-Nath, Moroi, ...)

MSSM Contributions to Muon $g - 2$



x-axis: M_2
(gaugino mass)

y-axis: $m_{\tilde{t}}$
(slepton mass)

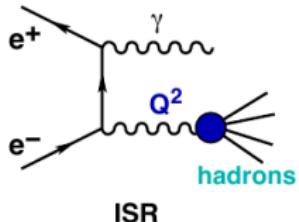
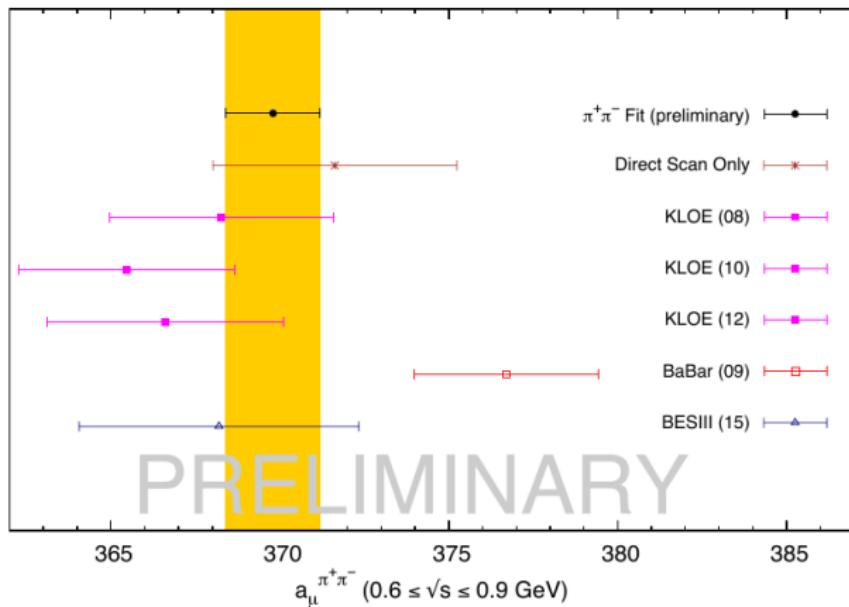
Figs from Cho,
Hagiwara, Matsumoto
and DN

Our preliminary results

HVP: HLMNT \rightarrow HKMNT in preparation

$\pi^+\pi^-$ channel: + KLOE12, + BES III from Rad. Ret.:

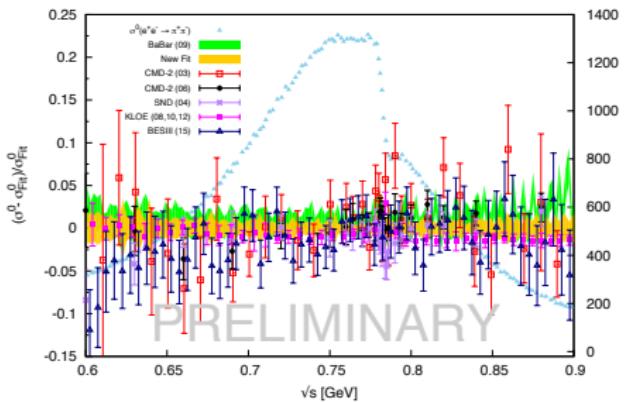
Prel. HKMNT combination w. full cov.-matrices:



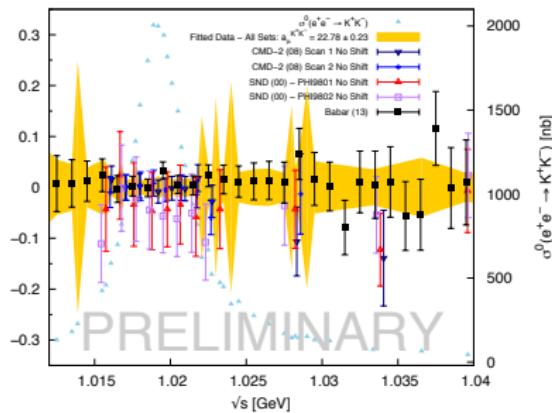
- $\chi^2_{\text{min}}/\text{d.o.f.} = 1.4$
- further improvements expected from CMD-3, more also from BaBar?

HVP: HLMNT -> HKMNT in preparation

$\pi^+\pi^-$ channel

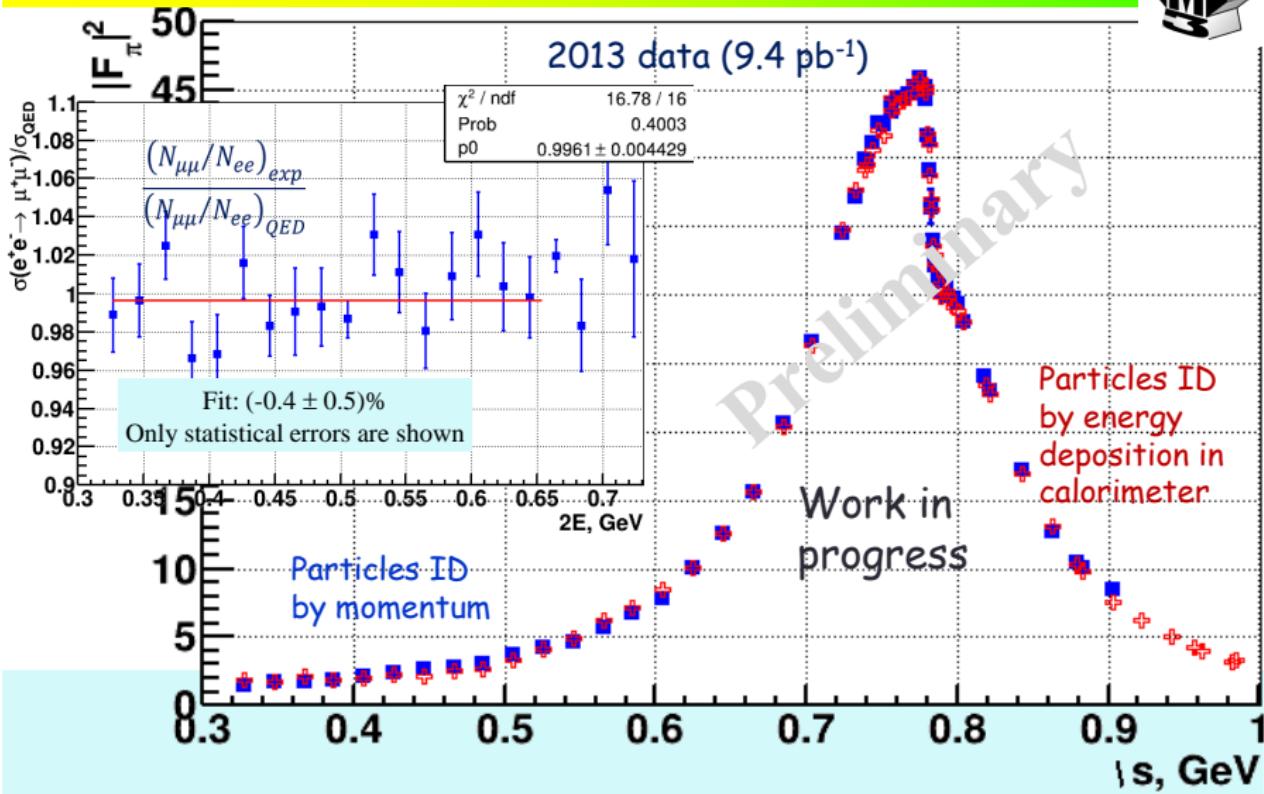


K^+K^- channel with recent BaBar



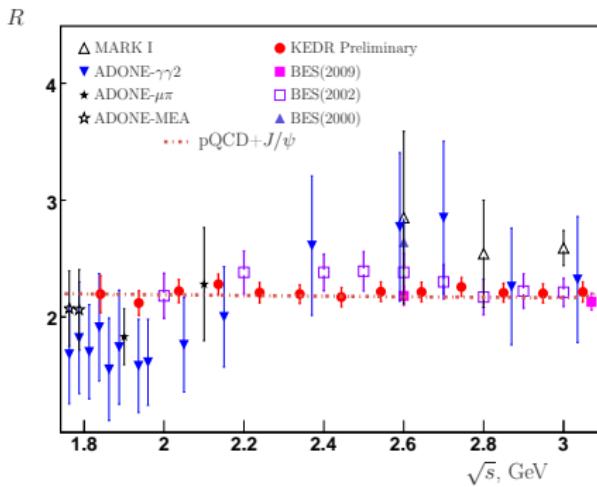
- Many new data sets and an improved combination algorithm, which takes fully into account all available covariance matrices, give significantly reduced errors and a slightly smaller mean value
- Previously sizeable additional (conservative) error from uncertainty in treatment of radiative corrections (VP + FSR), mainly from older data sets, gets reduced
- More exclusive data in multi-pion and K channels reduce uncertainty from estimate based on Iso-spin correlations

Near future prospects



Slide by S. I. Eidelman (BINP), talk at "Tau 2016", Sept. 2016

R Measurement between 1.84 and 3.05 GeV at KEDR - X



$\bar{R} = 2.209 \pm 0.020 \pm 0.046$ agrees with $R_{\text{pQCD}} = 2.18 \pm 0.02$
based on $\alpha_s(m_\tau) = 0.333 \pm 0.013$ derived from hadronic τ decays

Near Future Prospects (Blum et al, arXiv:1311.2198)

Current status and near-future
 ($\sim 5\text{-}7$ yrs) improvements
 in δa_μ (in units of 10^{-11})

Error	[20]	[21]	Future
δa_μ^{SM}	49	50	35
$\delta a_\mu^{\text{HLO}}$	42	43	26
$\delta a_\mu^{\text{HLbL}}$	26	26	25
$\delta(a_\mu^{\text{EXP}} - a_\mu^{\text{SM}})$	80	80	40

[20]: DHMZ

[21]: HLMNT

Near-future improvements in $\delta a_\mu^{\text{HLO}}$
 mainly from **VEPP-2000** and **BES-III**.

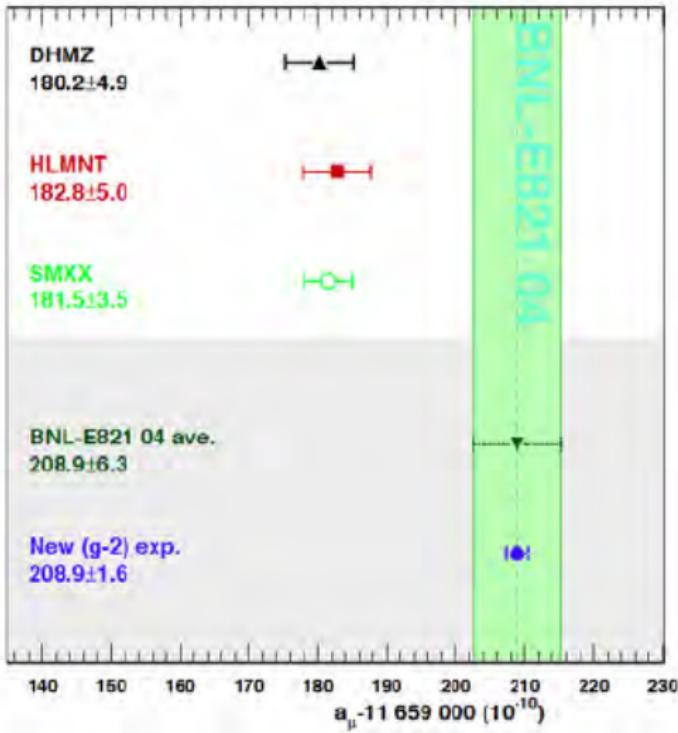


Table and Fig. from T. Blum et al, arXiv:1311.2198

Near Future Prospects, Lattice (Blum et al)

As for the hadronic vacuum polarization (HVP),

"...the lattice-QCD uncertainty on a_μ (HVP), currently at the 5%-level, can be reduced to 1 or 2% within the next few years."

"With increasing experience and computer power, it should be possible to compete with the e^+e^- determination of a_μ (HVP) by the end of the decade, perhaps sooner with additional technical advances."

As for I-by-I,

"... a lattice calculation with even a solid 30% error would already be very interesting. Such a result, ..., is not out of the question during the next 3-5 years."

From T. Blum et al, arXiv:1311.2198

Byproducts

Byproducts: QED coupling at the Z -boson mass

- ★ $\alpha(M_Z^2)$: the least well known among $\{G_\mu, M_Z, \alpha(M_Z^2)\}$, which are used as input to precision electroweak fits.
- ★ Running of α

$$\alpha(M_Z^2) = \frac{\alpha}{1 - \Delta\alpha_{\text{lep}}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha^{\text{top}}(M_Z^2)}$$

where $\Delta\alpha_{\text{lep}}(M_Z^2) = 0.03149769$ (Steinhauser),
 $\Delta\alpha^{\text{top}}(M_Z^2) = -0.0000728(14)$ and $\alpha = 1/137.035999679(94)$ (PDG10).

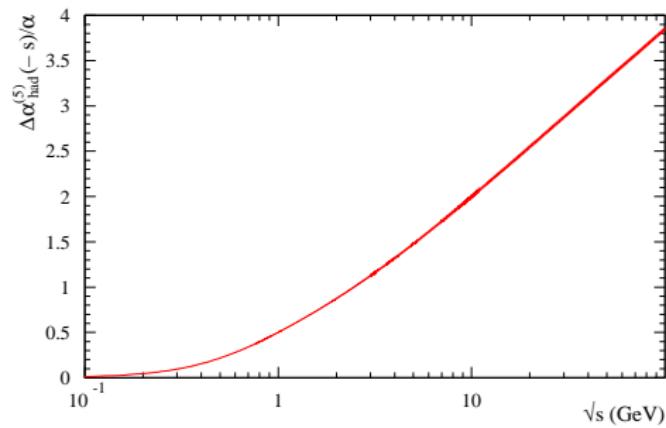
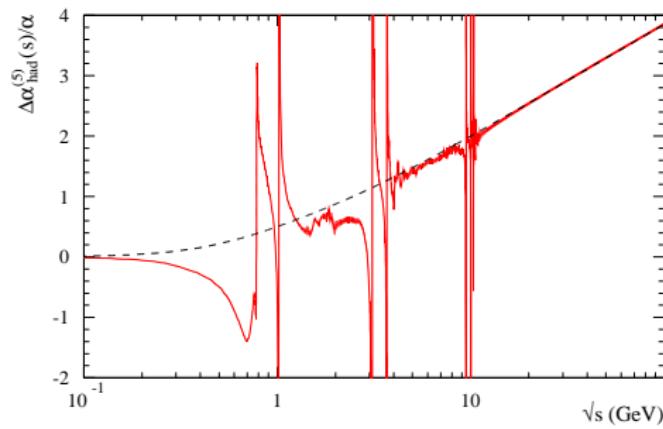
- ★ Similar dispersion relation: (\implies byproduct of $a_\mu^{\text{had}, \text{LO}}$)

$$\Delta\alpha_{\text{had}}^{(5)}(s) = -\frac{\alpha s}{3\pi} P \int \frac{R(s') ds'}{s'(s' - s)}$$

- ★ Our results: $\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) = (276.3 \pm 1.4) \times 10^{-4}$,
 $\alpha(M_Z^2)^{-1} = 128.944 \pm 0.019$.

Byproducts: running QED coupling $\alpha(q^2)$

The hadronic contribution $\Delta\alpha_{\text{had}}^{(5)}(q^2)$ to the running QED coupling for $q^2 > 0$ (left) and $q^2 < 0$ (right)



Fortran subroutine to compute the above is available from us upon request

Byproduct: Standard Model prediction for electron $g - 2$

QED contribution $115\ 965\ 218\ 00.7\ (0.6)_{\text{4loop}}(0.4)_{\text{5loop}}(7.6)\alpha$ (using $\alpha(^{87}\text{Rb})$) Aoyama et al

EW contribution $0.2973\ (0.0052)$ Czarnecki et al

Hadronic contributions

LO hadronic $18.66\ (0.11)$ **DN & TT** (was $18.75(0.18)$) Davier et al '98

NLO hadronic $-2.23\ (0.01)$ **DN & TT** (was $-2.25(0.05)$) Krause '97

light-by-light $0.39\ (0.13)$ Jegerlehner+Nyffeler

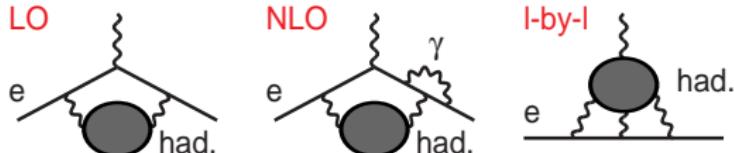
Theory total $115\ 965\ 218\ 17.8\ (0.6)_{\text{4loop}}(0.4)_{\text{5loop}}(0.2)_{\text{had}}(7.6)\alpha$

Experiment $115\ 965\ 218\ 07.3\ (2.8)$ Gabrielse et al

Theory – Exp $10.5\ (8.1)$ 1.3σ discrepancy

(in units of 10^{-13} . Numbers taken from Giudice et al, JHEP11(2012)113)

n.b.: hadronic contributions:



Summary

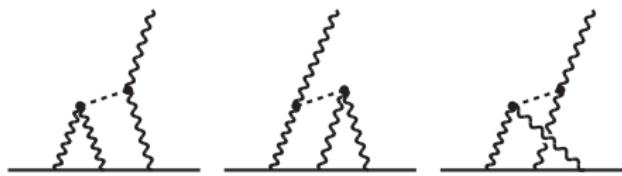
- Hadronic contrib. to the muon $g - 2$: key to improve the Standard Model prediction
- $\gtrsim 3 \sigma$ discrepancy in $(g - 2)_\mu$ between experiment and theory
 - ⇒ **New physics?** “Light” new particles like $\tilde{\mu}$, μ_{KK}, \dots ???
 - ⇒ Worth studying μ -EDM, $\mu \rightarrow e\gamma$, μ - e conv., ...!
 - (\Leftrightarrow No new physics seen at the LHC so far. What does this mean?)
- Two new experiments to measure the muon $g - 2$ planned at J-PARC and Fermilab.
- To establish this discrepancy more firmly, important to resolve the tension between the KLOE and BaBar data
 - ⇒ new data from BES-III appeared recently
 - ⇒ new precise data from VEPP-2000 and SuperKEKB strongly awaited.

Backup Slides

Modern evaluation of l-by-l contribution

(Melnikov & Vainshtein)

1. First, use the large N_C expansion to find that the leading contribution is the pion pole contribution.



2. Choose the momentum-dependence of the $\pi\gamma\gamma$ coupling (form factor) in such a way that it is consistent with a constraint from QCD (OPE) at the momentum region $q_1^2 \sim q_2^2 \gg q_3^2$. Integrate over the loop momenta.
3. Repeat the above for η, η', a_1, \dots . Basically that's all for the LO in $1/N_C$.
4. As for NLO in $1/N_C$, it depends on authors which diagram is numerically important.

For example,

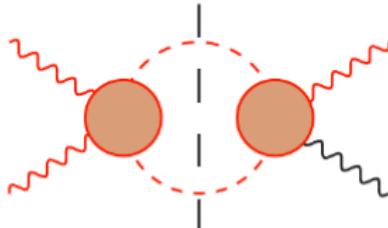
$$a_\mu^{\text{lbyl}} = \begin{cases} (10.5 \pm 2.6) \times 10^{-10} & \text{Prades-de Rafael-Vainshtein, arXiv:0901.0306} \\ (11.6 \pm 4.0) \times 10^{-10} & \text{Nyffeler, arXiv:0901.1172} \end{cases}$$

dispersion relation approaches for a_μ (I)

Hoferichter, Colangelo, Procura, Stoffer (2014)

→ dispersion formalism for $\gamma^* \gamma^* \rightarrow \gamma^* \gamma$

$$\Pi_{\mu\nu\lambda\sigma} = \Pi_{\mu\nu\lambda\sigma}^{\pi^0\text{-pole}} + \Pi_{\mu\nu\lambda\sigma}^{\text{FsQED}} + \bar{\Pi}_{\mu\nu\lambda\sigma} + \dots$$



→ master formula for a_μ

$$a_\mu^{\pi\pi} = e^6 \int \frac{d^4 q_1}{(2\pi)^4} \int \frac{d^4 q_2}{(2\pi)^4} \frac{I^{\pi\pi}}{q_1^2 q_2^2 s ((p+q_1)^2 - m^2) ((p-q_2)^2 - m^2)},$$

$$I^{\pi\pi} = \sum_{l \in \{1,2,3,6,14\}} \left(T_{i,s} I_{l,s} + 2 T_{i,u} I_{l,u} \right) + 2 T_{9,s} I_{9,s} + 2 T_{9,u} I_{9,u} + 2 T_{12,u} I_{12,u}$$

with $I_{l,(s,u)}$ dispersive integrals and $T_{i,(s,u)}$ integration kernels

$$I_{1,s} = \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} \frac{ds'}{s' - s} \left(\frac{1}{s' - s} - \frac{s' - q_1^2 - q_2^2}{\lambda(s', q_1^2, q_2^2)} \right) \text{Im} \bar{h}_{++,++}^0(s'; q_1^2, q_2^2; s, 0),$$

$$I_{6,s} = \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} \frac{ds'}{(s' - q_1^2 - q_2^2)(s' - s)^2} \text{Im} \bar{h}_{+-,+}^2(s'; q_1^2, q_2^2; s, 0) \left(\frac{75}{8} \right)$$

Helicity amplitudes contribute up to $J = 2$ (S and D waves)

Slide by M. Vanderhaeghen, talk at "Lepton Moments 2014", July 2014

Exp. inputs for evaluation of a_μ (had, l-by-l)

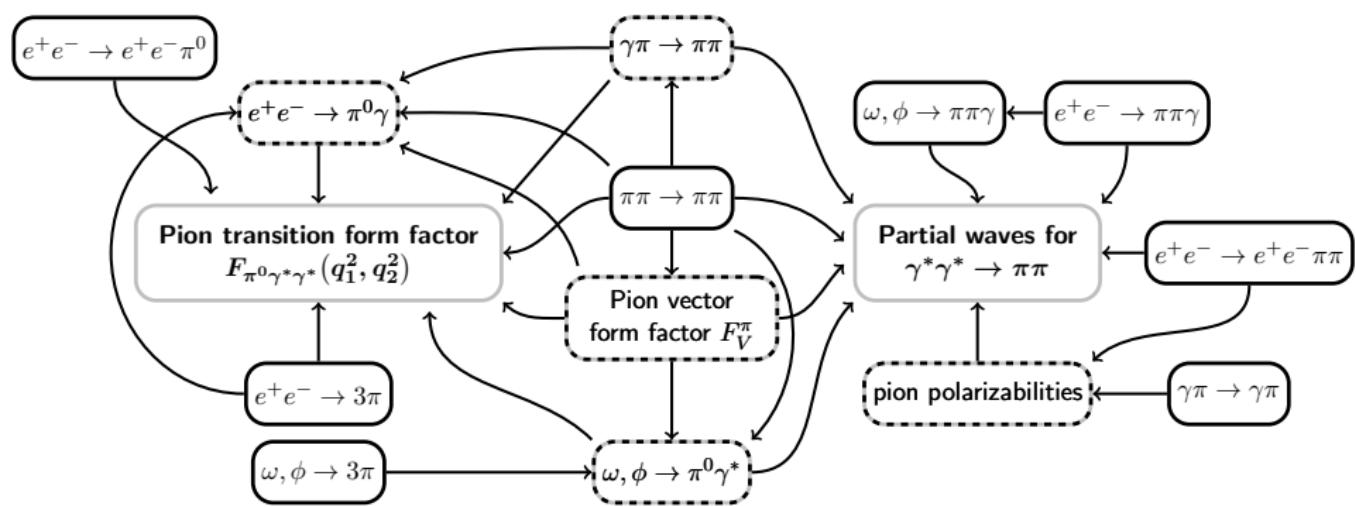
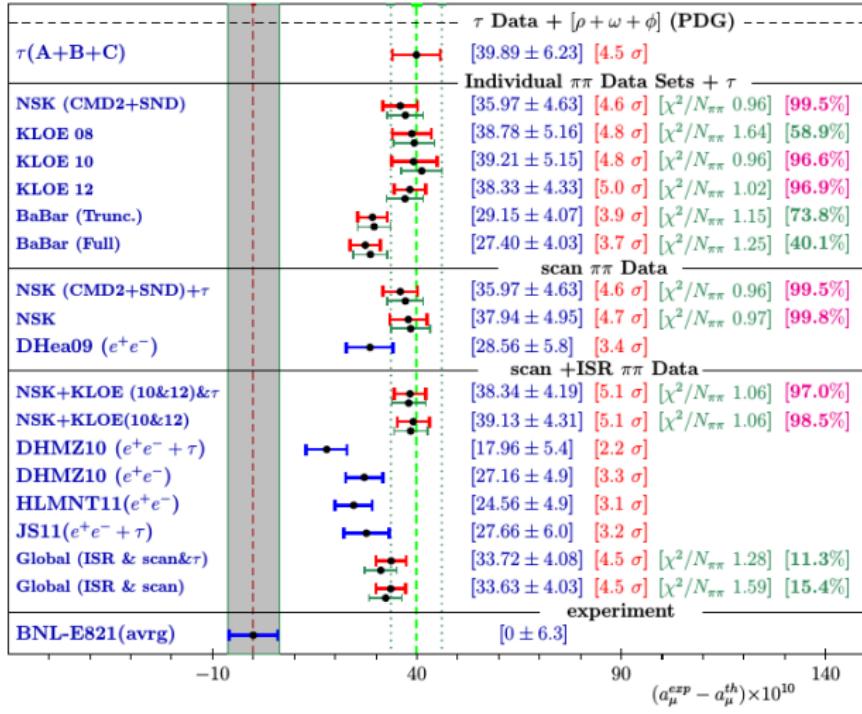


Fig. from G. Colangelo et al, arXiv:1408.2517

Updated anal.
(now with
fitted normal.
factors) in
global fit
based on HLS
model:



← 'preferred'
has 5σ already,
from both
aggressive
errors and
shifts due to
the HLS model
(and the fit?)