

# New Strategy to Explore CP Violation with $B_s \rightarrow K^- K^+$

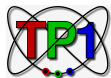
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based on [arXiv:1608.00901](https://arxiv.org/abs/1608.00901)



Theor. Physik 1



# Motivation

- $B$  meson decays offer an interesting laboratory to search for new physics
  - ▶ New sources of CP violation
- Unprecedented precision required for LHCb upgrade and Belle II
- Theoretical precision limited by strong interactions

## $U$ -spin symmetry

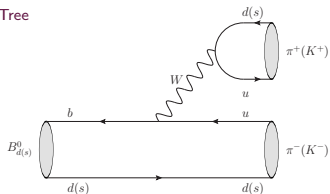
Flavor symmetries provide valuable insights into hadronic non-perturbative parameters

- Precision limited by symmetry-breaking corrections
- *New strategy*

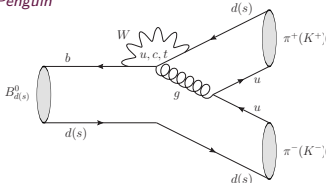
# $B_s^0 \rightarrow K^- K^+$ decay topologies

R. Fleischer [1999]

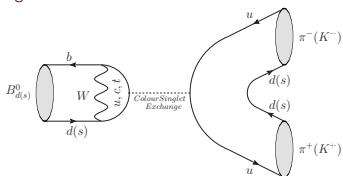
Tree



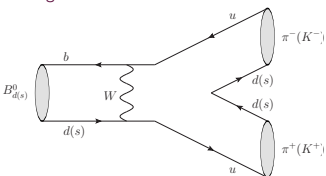
Penguin



Penguin Annihilation



Exchange



- Penguin topologies dominate
  - ▶ Sensitive to physics beyond the Standard Model
- Related to  $B_d^0 \rightarrow \pi^- \pi^+$  via  $U$ -spin symmetry ( $s$ -quark  $\leftrightarrow$   $d$ -quark)

$$B_s^0 \rightarrow K^- K^+ \text{ and } B_d^0 \rightarrow \pi^- \pi^+$$

R. Fleischer [1999]

$$A(B_s^0 \rightarrow K^- K^+) = \sqrt{\epsilon} e^{i\gamma} C' \left[ 1 + \frac{1}{\epsilon} d' e^{i\theta'} e^{-i\gamma} \right]$$

$$A(B_d^0 \rightarrow \pi^- \pi^+) = e^{i\gamma} C \left[ 1 - d e^{i\theta} e^{-i\gamma} \right]$$

$$C' \propto T' + P^{(ut)'} + E' + PA^{(ut)'} \quad d' e^{i\theta'} \propto \frac{P^{(ct)'} + PA^{(ct)'}}{T' + P^{(ut)'} + E' + PA^{(ut)'}}$$

- Penguin dominated  $\epsilon \simeq 0.05$
- Weak phase  $\gamma$  of Unitarity Triangle
- $C$  and  $d$  analogous to  $C'$  and  $d'$

### ***U-spin symmetry***

$$d e^{i\theta} = d' e^{i\theta'} \text{ and } C = C'$$

- Direct and Mixing induced CP violation

$$\begin{aligned} \mathcal{A}_{\text{CP}}(t) &\equiv \frac{\Gamma(B_q^0(t) \rightarrow f) - \Gamma(\bar{B}_q^0(t) \rightarrow \bar{f})}{\Gamma(B_q^0(t) \rightarrow f) + \Gamma(\bar{B}_q^0(t) \rightarrow \bar{f})} \\ &\propto \mathcal{A}_{\text{CP}}^{\text{dir}}(B_q \rightarrow f) \cos(\Delta M_q t) + \mathcal{A}_{\text{CP}}^{\text{mix}}(B_q \rightarrow f) \sin(\Delta M_q t) \end{aligned}$$

- ▶ Depend on  $\gamma$  and hadronic parameters  $d$  and  $\theta$
- ▶  $\mathcal{A}_{\text{CP}}^{\text{mix}}$  depends also on  $B_q^0$ - $\bar{B}_q^0$  mixing angle  $\phi_q$

## Original Strategy

- Hadronic parameters are related to those in  $B_d^0 \rightarrow \pi^- \pi^+$  by  $U$ -spin
  - ▶ Extract  $\gamma$  and  $\phi_s$  from CP asymmetries

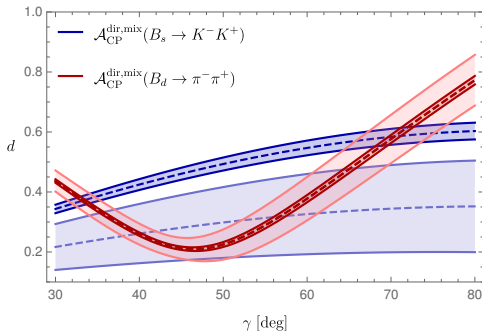
# Original Strategy

R. Fleischer and R. Kneijens [2011]

R. Fleischer [1999,2007]

T. Abe et al. (Belle-II)[2010]

R. Aaij et al. (LHCb)[2013]



- LHCb[2015]  $\gamma = (63.5^{+7.2}_{-6.7})^\circ$
- Future  $\gamma$  of  $\mathcal{O}(1^\circ)$ 
  - ▶ Same uncertainty as future tree determination

- U-spin breaking  $\xi \equiv \frac{d'}{d} = 1 \pm 0.2$  and  $\Delta \equiv \theta' - \theta = (0 \pm 20)^\circ$
- Uncertainty on  $\gamma$  of  $\mathcal{O}(5^\circ)$

# Original Strategy

- CP asymmetries determine the “effective” mixing angle

$$\sin \phi_s^{\text{eff}} = \frac{\mathcal{A}_{\text{CP}}^{\text{mix}}(B_s \rightarrow K^- K^+)}{\sqrt{1 - \mathcal{A}_{\text{CP}}^{\text{dir}}(B_s \rightarrow K^- K^+)^2}}$$

- $B_s^0 - \bar{B}_s^0$  mixing angle determined via  $\sin \phi_s^{\text{eff}} \equiv \sin(\phi_s + \Delta\phi_{KK})$
- Hadronic non-perturbative correction  $\Delta\phi_{KK}$

$$\tan \Delta\phi_{KK} = \frac{2\epsilon d' \cos \theta' \sin \gamma + \epsilon^2 \sin(2\gamma)}{d'^2 + 2\epsilon d' \cos \theta' \cos \gamma + \epsilon^2 \cos(2\gamma)},$$

- $U$ -spin symmetry  $\rightarrow \Delta\phi_{KK}$ 
  - ▶ LHCb[2015]  $\phi_s = -(6.9_{-8.0}^{+9.2})^\circ$
  - ▶ Current PDG average  $\phi_s = -(0.68 \pm 2.2)^\circ$

# Original Strategy

$$\sin \phi_s^{\text{eff}} \equiv \sin(\phi_s + \Delta\phi_{KK})$$

- LHCb upgrade  $\phi_s^{\text{eff}} \sim \mathcal{O}(0.5^\circ)$
- Hadronic non-perturbative correction
  - ▶  $U$ -spin symmetry  $\rightarrow \Delta\phi_{KK} \sim \mathcal{O}(0.3^\circ)$
  - ▶  $U$ -spin breaking of 20% gives  $\Delta\phi_{KK} = -(9.0 \pm 2.6)^\circ$
- Future  $B_s^0 \rightarrow J/\psi\phi$  of  $\mathcal{O}(0.5^\circ)$
- To match future experimental precision  $U$ -spin breaking corrections have to be known at the **few percent level**



# New Strategy

- Minimize use of  $U$ -spin symmetry
- Use  $\gamma$  and  $\phi_d$  as input
- Non-factorizable effects probed by semileptonic ratios

# New Strategy

Non-factorizable effects probed by semileptonic ratios

$$R_\pi \equiv \frac{\Gamma(B_d \rightarrow \pi^- \pi^+)}{|d\Gamma(B_d \rightarrow \pi^- \ell^+ \nu_\ell)/dq^2|_{q^2=m_\pi^2}} = 6\pi^2 X_\pi |V_{ud}|^2 f_\pi^2 |a_{\text{NF}}|^2 r_\pi$$

$$R_K \equiv \frac{\Gamma(B_s \rightarrow K^- K^+)}{|d\Gamma(B_s \rightarrow K^- \ell^+ \nu_\ell)/dq^2|_{q^2=m_K^2}} = 6\pi^2 X_K |V_{us}|^2 f_K^2 |a'_{\text{NF}}|^2 r_K$$

- $r_\pi = 1 - 2d \cos \theta \cos \gamma + d^2 \rightarrow$  from CP asymmetries in  $B_d^0 \rightarrow \pi^- \pi^+$
- $r_K = 1 - 2d'/\epsilon \cos \theta' \cos \gamma + d'^2/\epsilon^2$
- $a_{\text{NF}} \equiv (1 + r_P)(1 + x)a_{\text{NF}}^T$ 
  - ▶  $r_P \equiv P^{(ut)}/T$
  - ▶  $x \equiv \frac{E+PA^{(ut)}}{T+P^{(ut)}}$

# New Strategy

- Determine  $d'$  and  $\theta'$  using  $r_K = r_\pi \frac{R_K}{R_\pi} \left| \frac{V_{ud}}{V_{us}} \right|^2 \left( \frac{f_\pi}{f_K} \right)^2 \frac{X_\pi}{X_K} (\xi_{NF}^a)^2$ 
  - ▶  $\Delta\phi_{KK}$  and  $\phi_s$

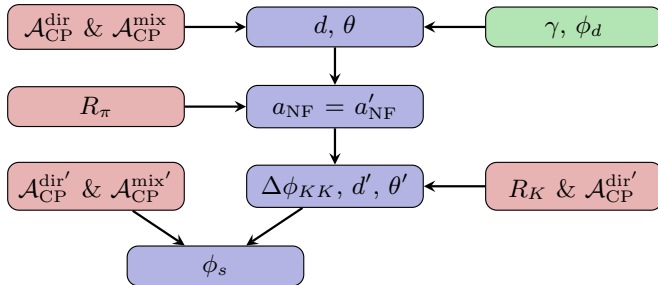
## $U$ -spin parametrization

$$\xi_{NF}^a \equiv \left| \frac{a_{NF}}{a'_{NF}} \right| = \left| \frac{a_{NF}^T}{a_{NF}^{T'}} \right| \left| \frac{1 + r_P}{1 + r'_P} \right| \left| \frac{1 + x}{1 + x'} \right|$$

- Very favorable structure in terms of  $U$ -spin-breaking parameters
  - ▶ Robust structure
  - ▶ Minimal use of  $U$ -spin symmetry

# New Strategy

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- Use  $\gamma$  and  $\phi_d$  as input
- Non-factorizable effects probed by semileptonic ratios



# $U$ -spin-breaking corrections

Beneke, Huber, Li [2010] Gronau *et al.* [1995]  
Fleischer, Jaarsma, and KKV[2016]

$$\xi_{\text{NF}}^a \equiv \left| \frac{a_{\text{NF}}^T}{a_{\text{NF}}^{T'}} \right| \left| \frac{1+r_P}{1+r'_P} \right| \left| \frac{1+x}{1+x'} \right|$$

- Use data to quantify  $U$ -spin-breaking corrections
- QCD factorization  $a_{\text{NF}}^T = 1.000_{-0.069}^{+0.029} + (0.011_{-0.050}^{+0.023})i$
- $a_{\text{NF}}^T/a_{\text{NF}}^{T'} \simeq 1 + \Delta_{\text{NF}}^T \xi_{\text{NF}}^T + \mathcal{O}((\Delta_{\text{NF}}^T)^2)$
- $U$ -spin breaking of 20% gives a correction of  $\mathcal{O}(1\%)$

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- $U$ -spin breaking of 20% gives a correction of  $\mathcal{O}(1\%)$
  
- $r_P = P^{(ut)}/T \simeq 0.2$  from pure Penguin decays
- $\frac{1+r_P}{1+r'_P} \simeq 1 + (1 - \xi_P)r_P + \mathcal{O}(r_P^2)$
- $U$ -spin breaking of  $\xi_P \sim 0.2$  gives a correction of  $\mathcal{O}(4\%)$ .

# Exchange and Penguin Annihilation contributions

Gronau *et al.* [1995]

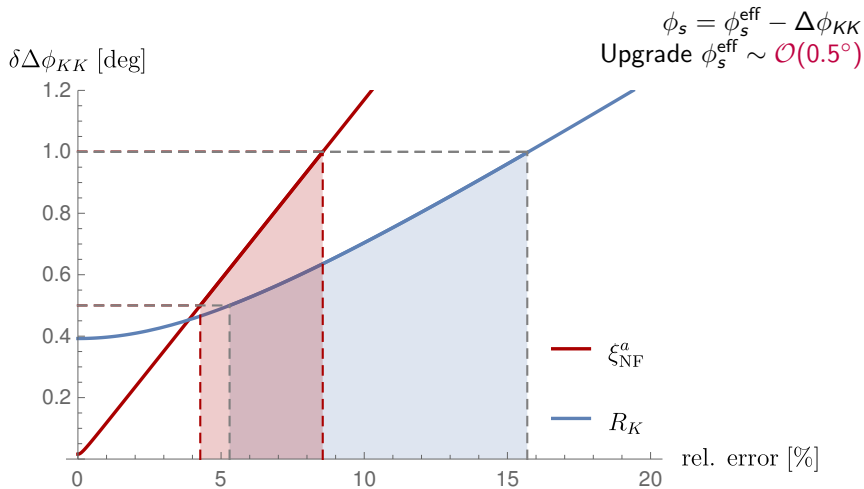
Fleischer, Jaarsma, and KKV[2016]

- $x = \frac{E+PA^{(ut)}}{T+P^{(ut)}}$  expected to be small
  - ▶ Constrained by pure exchange and penguin annihilation topologies  
 $B_d^0 \rightarrow K^+ K^-$  and  $B_s^0 \rightarrow \pi^+ \pi^-$
  - ▶ New ICHEP data  $x \sim 0.15$
- $\frac{1+x}{1+x'} = 1 + (1 - \xi_x)x + \mathcal{O}(x^2)$
- $U$ -spin breaking of 20% gives a correction of  $\mathcal{O}(3\%)$ 
  - ▶ Future data will narrow this down further

## $U$ -spin-breaking corrections

Expected combined correction  $\xi_{NF}^a \sim \mathcal{O}(5\%)$

# Illustration of the Error on $\Delta\phi_{KK}$

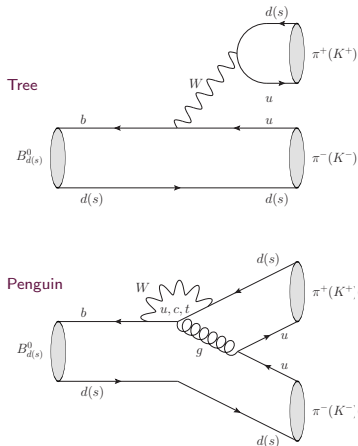


- $0.5^\circ$  precision on  $\Delta\phi_{KK}$  requires  $\mathcal{O}(5\%)$  precision on  $R_K$  and  $\xi_{\text{NF}}^a$



# Picture from Current Data

- $B_s^0 \rightarrow K^- \ell^+ \nu_\ell$  not yet measured
- $B_d^0 \rightarrow \pi^- K^+$  has only Tree and Penguin topologies
  - ▶ Similar to  $B_s^0 \rightarrow K^- K^+$
  - ▶ Ignoring  $E$  and  $PA$  topologies, only spectator quark difference
  - ▶  $\tilde{R}_K \equiv \frac{\Gamma(B_d \rightarrow \pi^- K^+)}{|d\Gamma(B_d \rightarrow \pi^- \ell^+ \nu_\ell)/dq^2|_{q^2=m_K^2}}$
  - ▶ Semileptonic rate cancels largely in the ratio  $R_\pi/\tilde{R}_K$



# Picture from Current Data

- Using  $\gamma = (70 \pm 7)^\circ$

$$d = 0.58 \pm 0.16, \quad \theta = (151.4 \pm 7.6)^\circ,$$
$$\tilde{d}' = 0.50 \pm 0.03, \quad \tilde{\theta}' = (157.2 \pm 2.2)^\circ.$$

- Assuming  $U$ -spin symmetry  $\tilde{d}' = d' \rightarrow \Delta\phi_{KK} = -(10.7 \pm 0.6)^\circ$
- CP asymmetries in  $B_s^0 \rightarrow K^- K^+$  give  $\phi_s^{\text{eff}} = -(17.6 \pm 7.9)^\circ$

$$\phi_s = \phi_s^{\text{eff}} - \Delta\phi_{KK} = -(6.9 \pm 7.9)^\circ.$$

- Very consistent with LHCb determination  $\phi_s = -(6.9_{-8.0}^{+9.2})^\circ$

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## First $U$ -spin test

$$\xi = \tilde{d}'/d = 0.87 \pm 0.20 \quad \Delta = \tilde{\theta}' - \theta = (5.8 \pm 8.3)^\circ$$

# Conclusion

- New strategy to extract mixing angle  $\phi_s$ 
  - ▶ Semileptonic  $B_d^0 \rightarrow \pi^- \ell^+ \nu_\ell$ ,  $B_s^0 \rightarrow K^- \ell^+ \nu_\ell$  decays
  - ▶ Apply  $U$ -spin symmetry to robust quantities
- Theoretical precision of  $\mathcal{O}(0.5^\circ)$  attainable
- Current data show promising picture

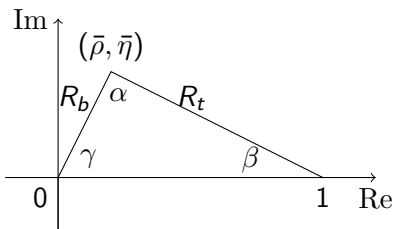
## Outlook

- ▶ Extensive study of Exchange and Penguin Annihilation topologies
- ▶ Analyses of  $B_s^0 \rightarrow K^- \ell^+ \nu_\ell$  strongly advocated
- ▶ New sources of CP violation may be revealed

# CKM Matrix

Wolfenstein [1983]  
Buras et. al [1994]

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$



$$\begin{aligned} \bar{\rho} &\equiv \rho \left(1 - \frac{1}{2}\lambda^2\right) \\ \bar{\eta} &\equiv \eta \left(1 - \frac{1}{2}\lambda^2\right) \end{aligned}$$