

# A study of the b-fragmentation function in $t\bar{t}$ -production with the ATLAS experiment at the LHC

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# Introduction

Our current understanding of QCD is based on the factorization theorem

$$\frac{d\sigma}{d\mathcal{O}} = \sum_{i,j} \int_0^1 dx_i dx_j \sum_f \int d\Phi_f f_{i/h_1}(x_i, \mu_F^2) f_{j/h_2}(x_j, \mu_F^2) \frac{d\hat{\sigma}_{ij \rightarrow f}}{d\hat{\mathcal{O}}} D_f(\hat{\mathcal{O}} \rightarrow \mathcal{O}, \mu_F^2)$$

In the String (Lund) Model used in Pythia, the fragmentation function is given by (up to a normalization constant)

$$f_B(z) \propto \frac{1}{z} (1-z)^a e^{-bm_T^2/z}$$

The Bowler modification introduces an additional parameter,  $rb$ , which acts as an effective mass for the quark

$$f_B(z) \propto \frac{1}{z^{1+br_b m_q^2}} (1-z)^a e^{-bm_T^2/z}$$

The observables considered which are sensitive to the b-fragmentation function are

$$x_B = 2 \frac{p_b \cdot Q}{Q^2}$$

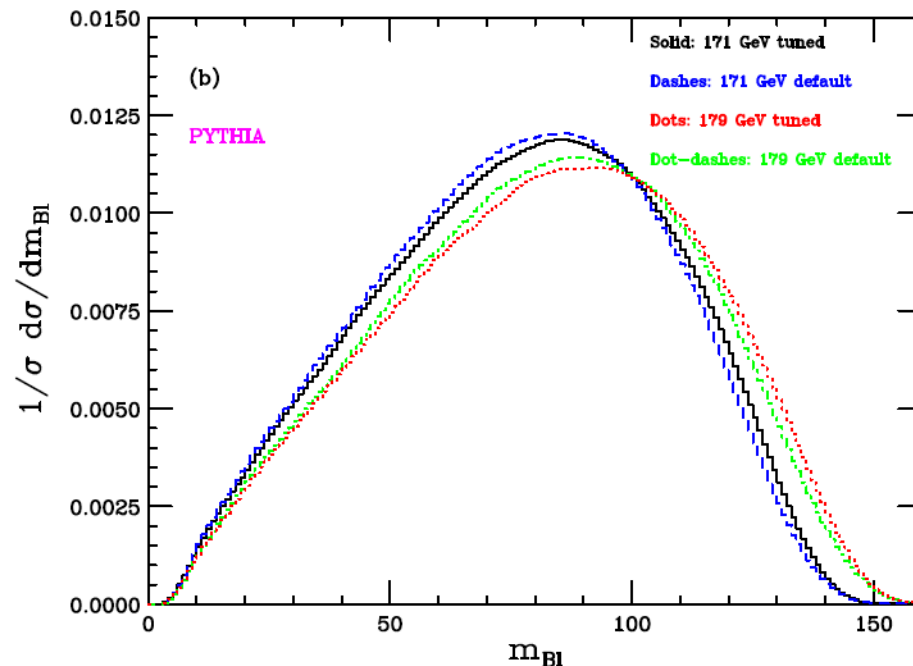
Where  $p_b$  is the 4-momentum of the B-hadron and  $Q$  is the 4-momentum of the Z boson, in  $e^+e^- \rightarrow Z \rightarrow b\bar{b}$  production at LEP.

The b-jet shapes in  $t\bar{t}$  production measured by ATLAS.

$$\rho(r) = \frac{1}{\Delta r} \frac{p_T(r - \Delta r/2, r + \Delta r/2)}{p_T(0, R)}$$

Which is the fraction of transverse energy contained in an annulus of the jet divided by the total energy.

The incomplete knowledge of the b-fragmentation function is an important contributor to the uncertainty in top mass measurements from final states in B-hadron production



B-lepton invariant mass distribution in top decay according to Pythia6, using default and tuned versions for  $m_t = 171$  GeV and 179 GeV (Taken from arxiv:1008.4498)

Three Pythia 8 tunes are considered.

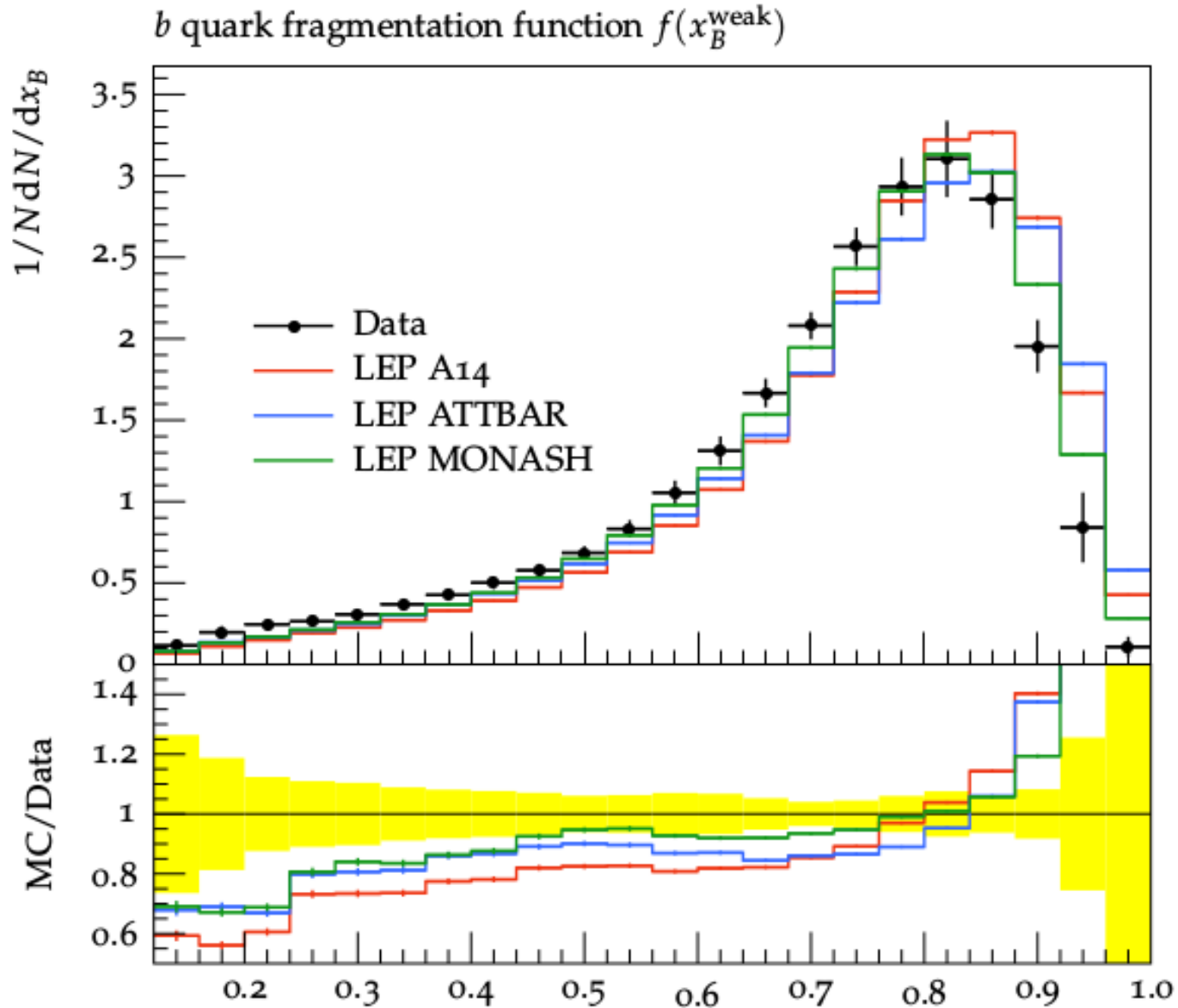
MONASH: The default Pythia 8 settings, the fragmentation functions are tuned to LEP data.

A14: Uses MONASH as baseline and retunes ISR, MPI and FSR to ATLAS data.

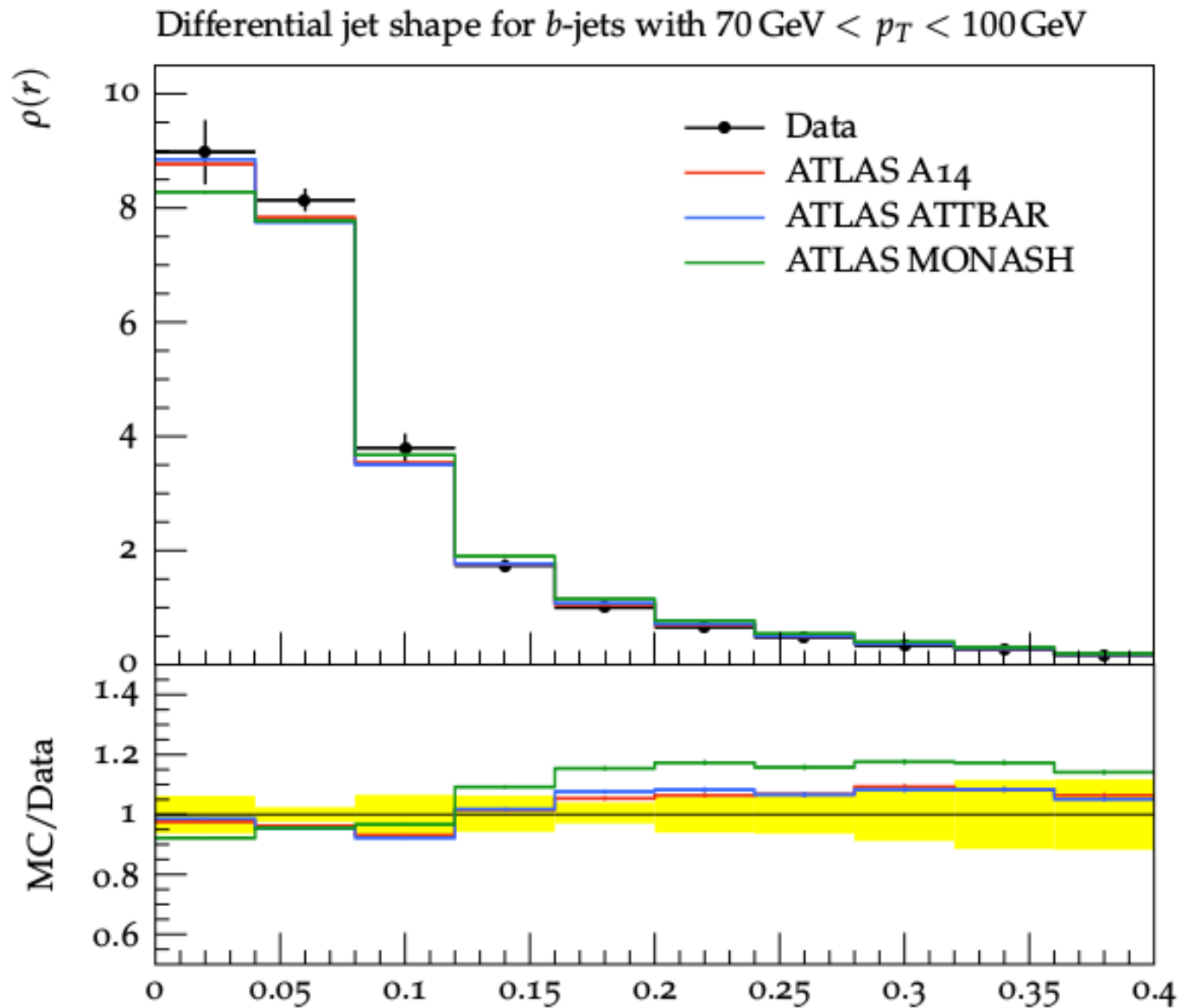
ATTBAR: Uses MONASH as baseline and retunes ISR and FSR to ATLAS top-antitop data. This tune is used as baseline for the extraction of rB.

The main difference between the three tunes lies in the value of  $\alpha_s^{FSR}(m_Z)$ , which is 0.1365 for MONASH, 0.127 for A14 and 0.1374 for ATTBAR.

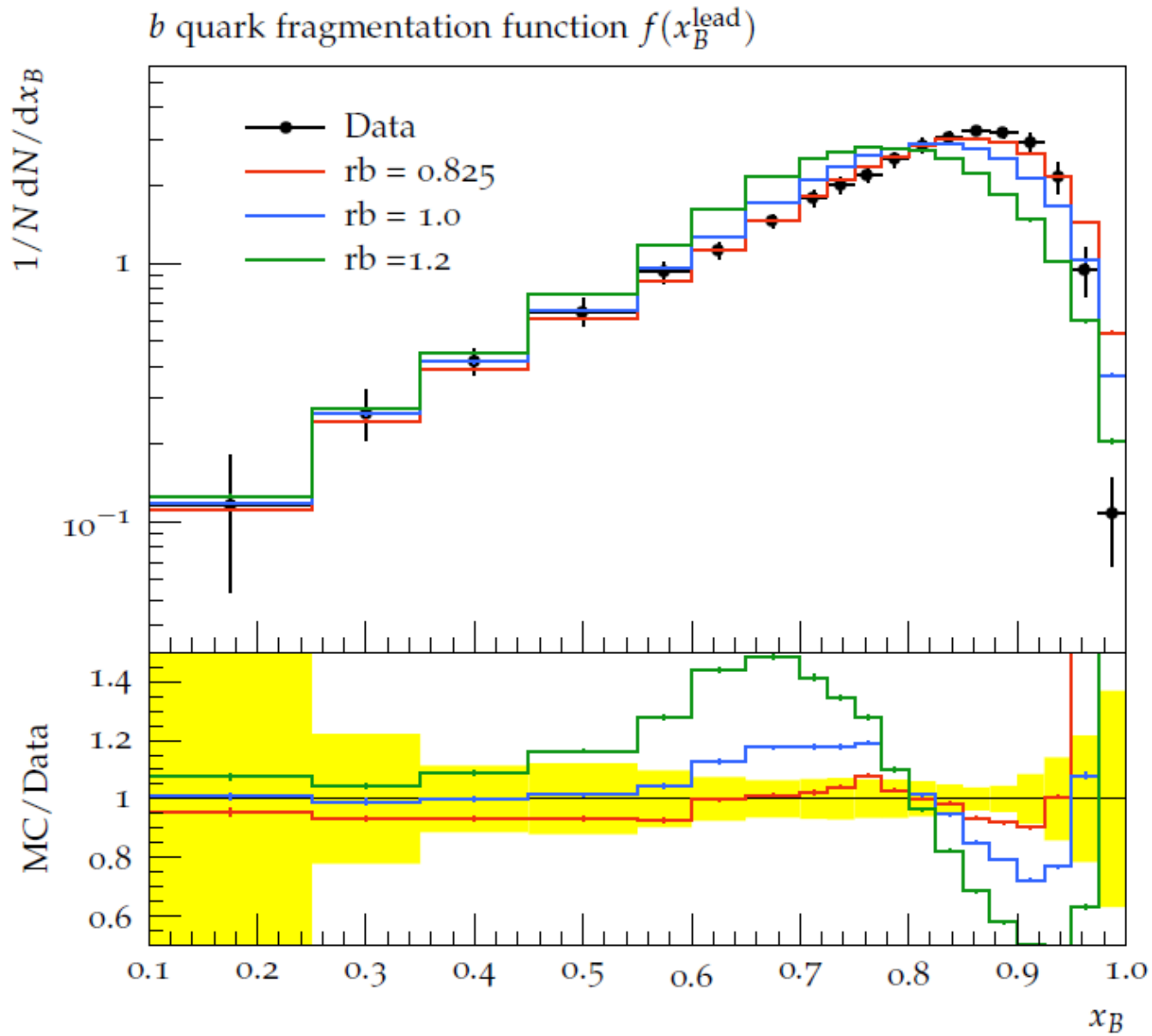
# Comparison of the three tunes for LEP (SLD)



# Comparison of the three tunes for ATLAS

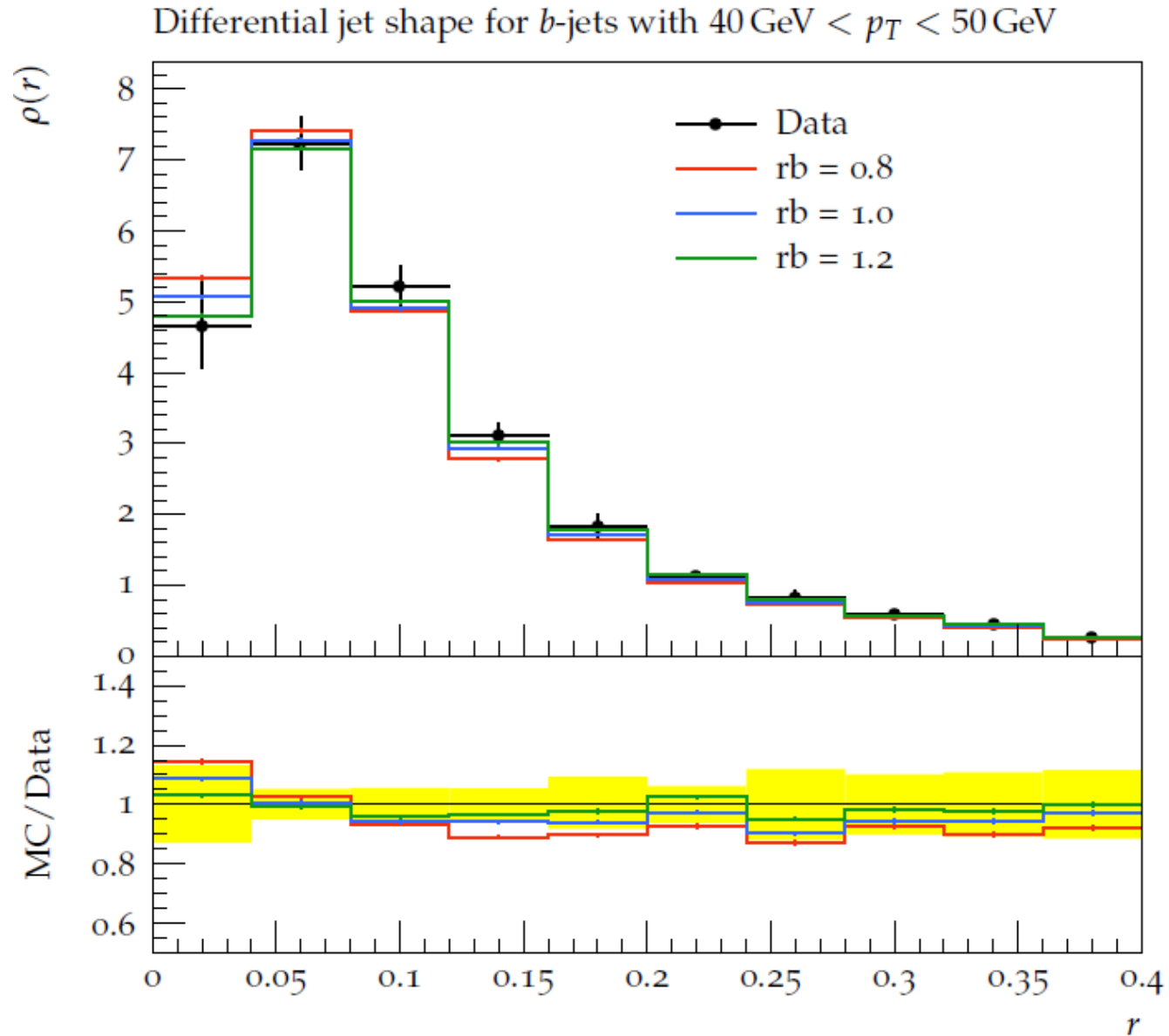


The following plot shows sensitivity to variations of  $r_B$  in ATTBAR for LEP (ALEPH)





The following plot shows sensitivity to variations of  $r_B$  in ATTBAR for ATLAS



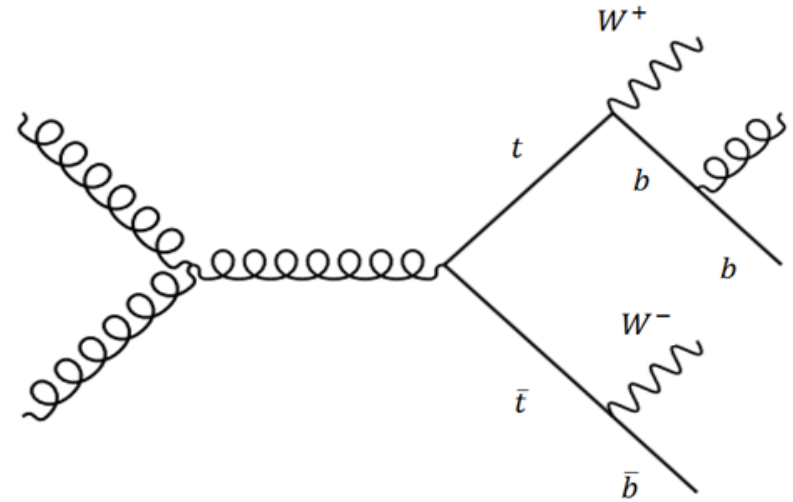
# Methodology

Used the following RIVET routines:

- $t\bar{t}$  b-jet shapes
  - ATLAS\_2013\_I1243871
- LEP b-fragmentation function
  - ALEPH\_2001\_S4656318
  - SLD\_2002\_S4869273
  - DELPHI\_2002\_069\_CONF\_603

The last numbers in the analysis correspond to the Inspire Record. Pythia8.219 was used to generate predictions for several values of  $r_B$ . The ATTBAR tune is used as a baseline. The default value of  $r_B$  is 0.855.

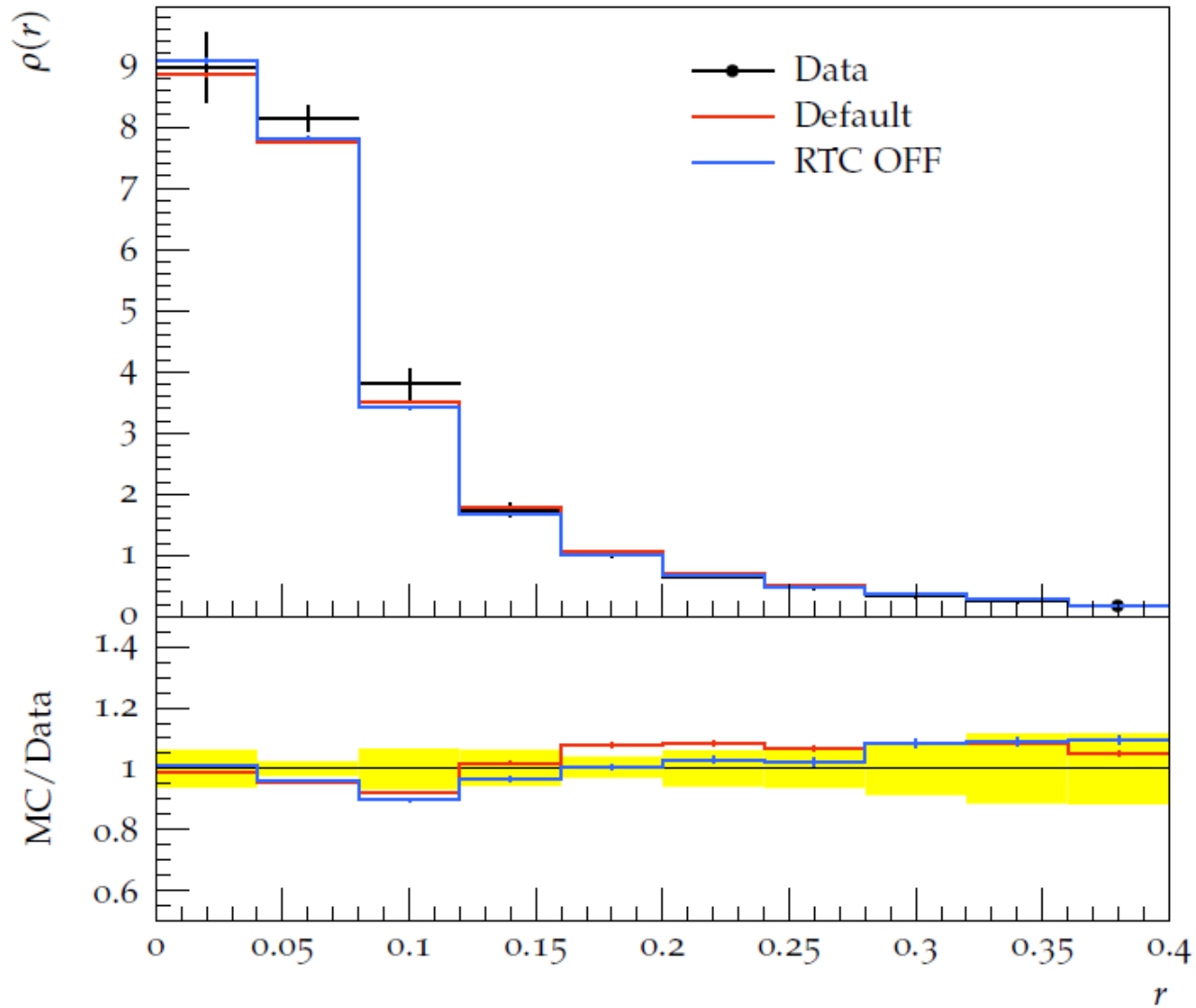
Top quarks decay almost exclusively into b quark and W.



B-jet shapes depend on the details of the FSR recoil strategy, which can be modified with a switch available in Pythia: RecoilToColoured (*on* by default).

*In the decays of coloured resonances, it is not possible to set up dipoles with matched colours. Originally the  $b$  radiator therefore has  $W$  as recoiler, and that choice is unique. Once a gluon has been radiated, however, it is possible either to have the unmatched colour (inherited by the gluon) still recoiling against the  $W$  (off), or else let it recoil against the  $b$  also for this dipole (on)*

Differential jet shape for  $b$ -jets with  $70 \text{ GeV} < p_T < 100 \text{ GeV}$



With RecoilToColoured off the agreement improves at high  $p_T$  and high  $r$

For the uncertainty in the value of  $r_B$ , several runs were made and the plots of  $\chi^2$  vs.  $r_B$  were interpolated with a parabola near the minimum.

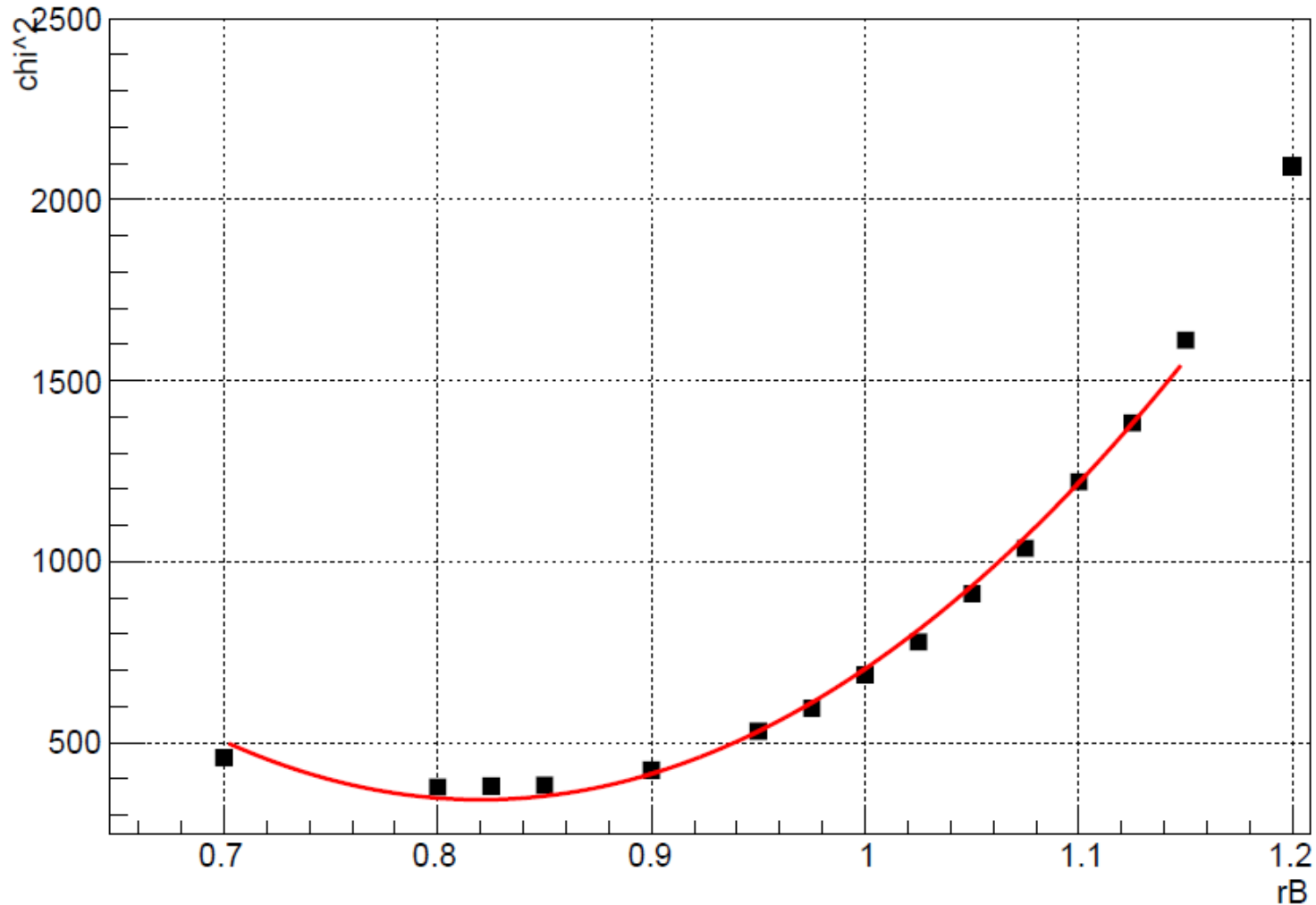
The following formula for  $\chi^2$  was taken for the calculation of the uncertainty (Here, the  $\Delta$  are taken as fully uncorrelated)

$$\chi^2 = \sum_i \frac{(\mu_i - n_i)^2}{\Delta_i^2}$$

Finally, the points where  $\chi_{min}^2 + 1$  cuts the parabola were taken as the uncertainty in  $r_B$

The following plot corresponds to the fit of the averages of the LEP measurements near the minimum value of  $r_B$

chi^2 vs rB

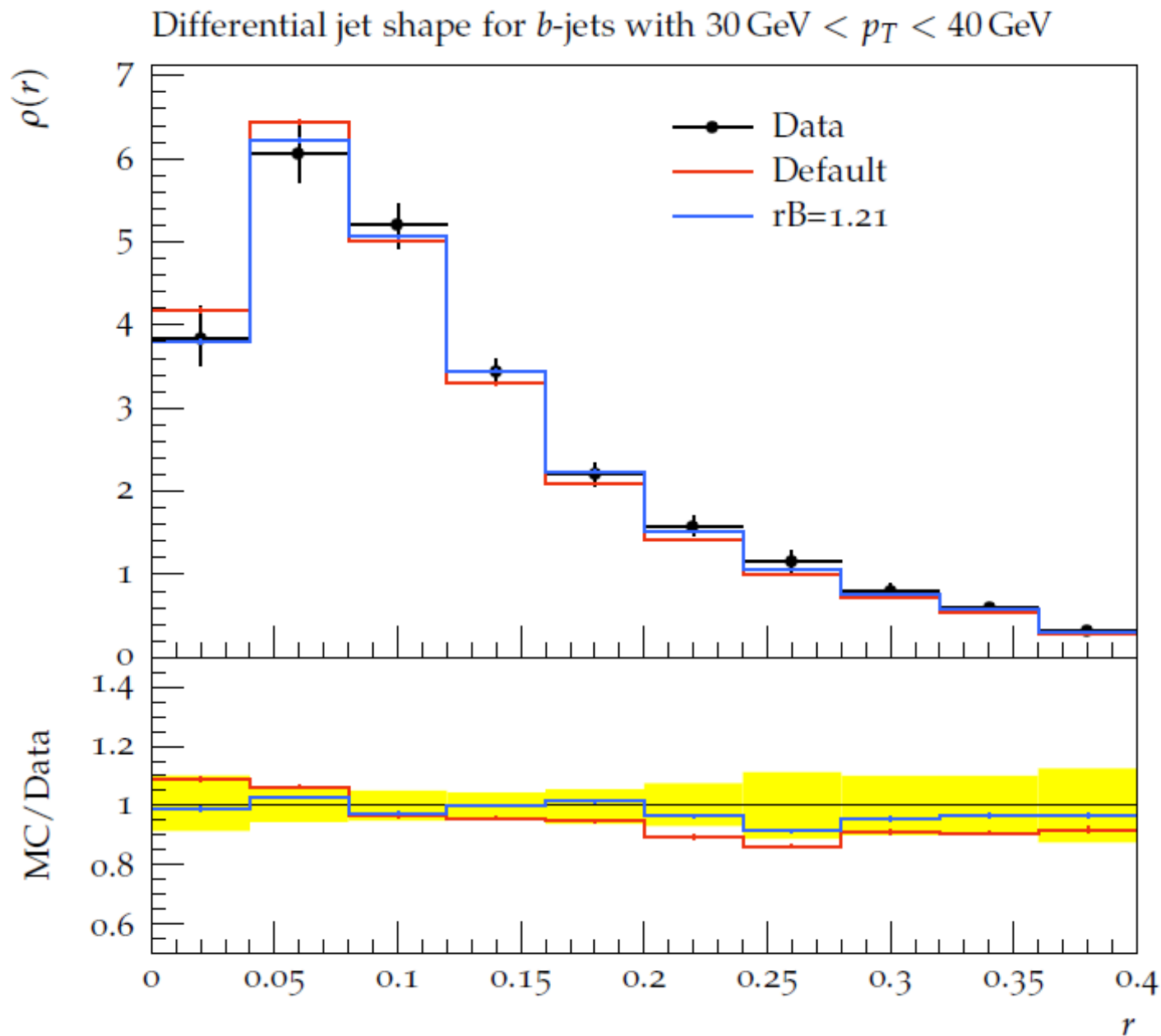


rb extracted from the ATLAS t-tbar b-jet shapes

Pt jet [GeV]	rb	delta rb	X <sup>2</sup> /dof
30 – 40	1,3	0,1	2/9
40 – 50	1,3	0,1	1/9
50 – 70	1,3	0,1	1/9
70 – 100	1,1	0,1	4/9
100 – 150	1,2	0,4	4/9
COMBINATION	1,25	0,06	13/49
COMBINATION (NO RTC)	1,04	0,08	37/49

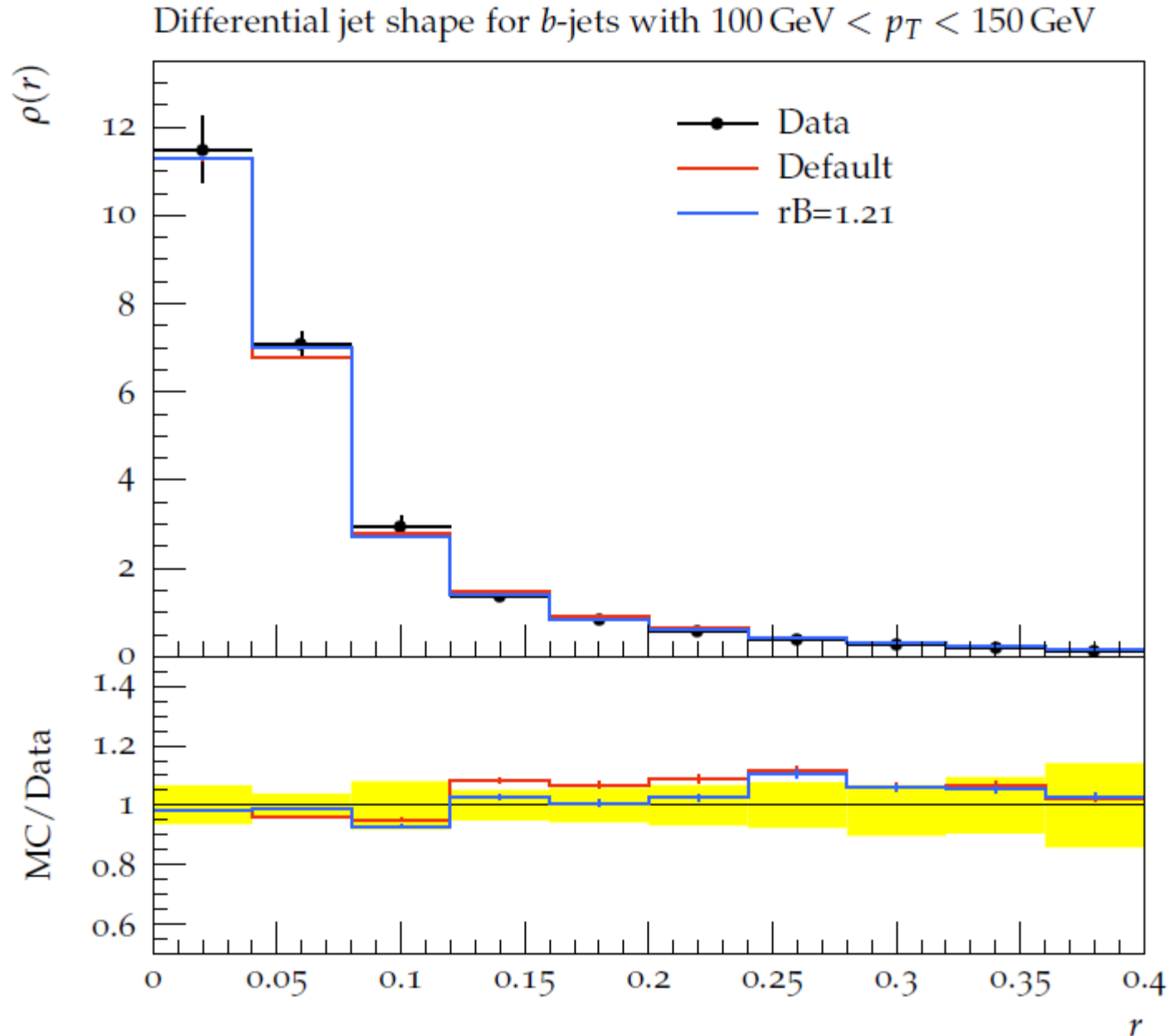
The various pt-jet bins yield compatible results

ATLAS ATTBAR, low  $p_T$   $r_B = 1.21 \pm 0.08$   $\chi^2/\text{dof} = 13/49$  Average for ATLAS





ATLAS ATTBAR, high  $p_T$   $r_B = 1.21 \pm 0.08$   $\chi^2/\text{dof} = 13/49$  Average for ATLAS

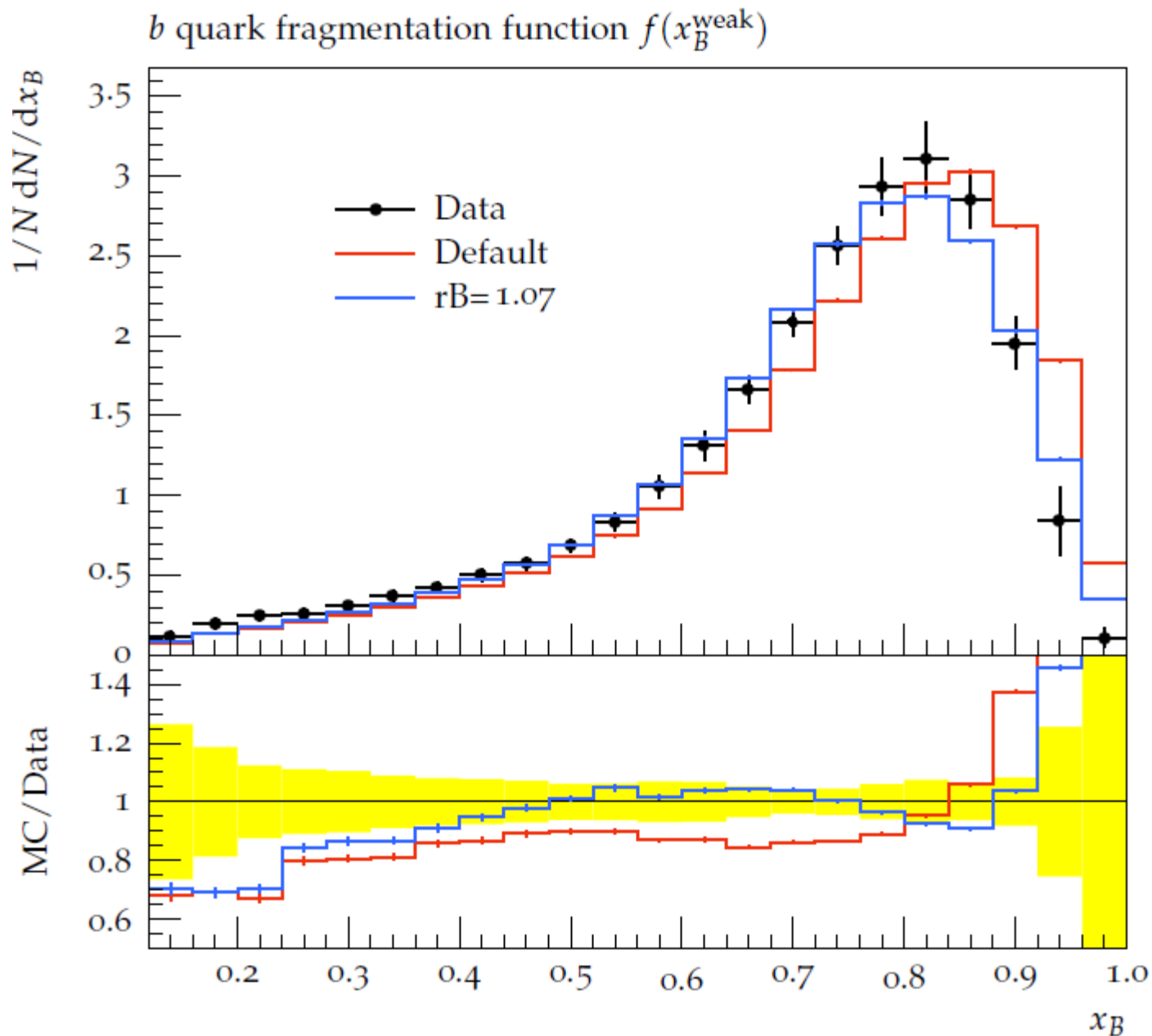


Now we compare the extraction of  $r_b$  from ATLAS  $t$ - $t$ bar data to the same extraction from LEP  $Z \rightarrow b$ - $b$ bar data.

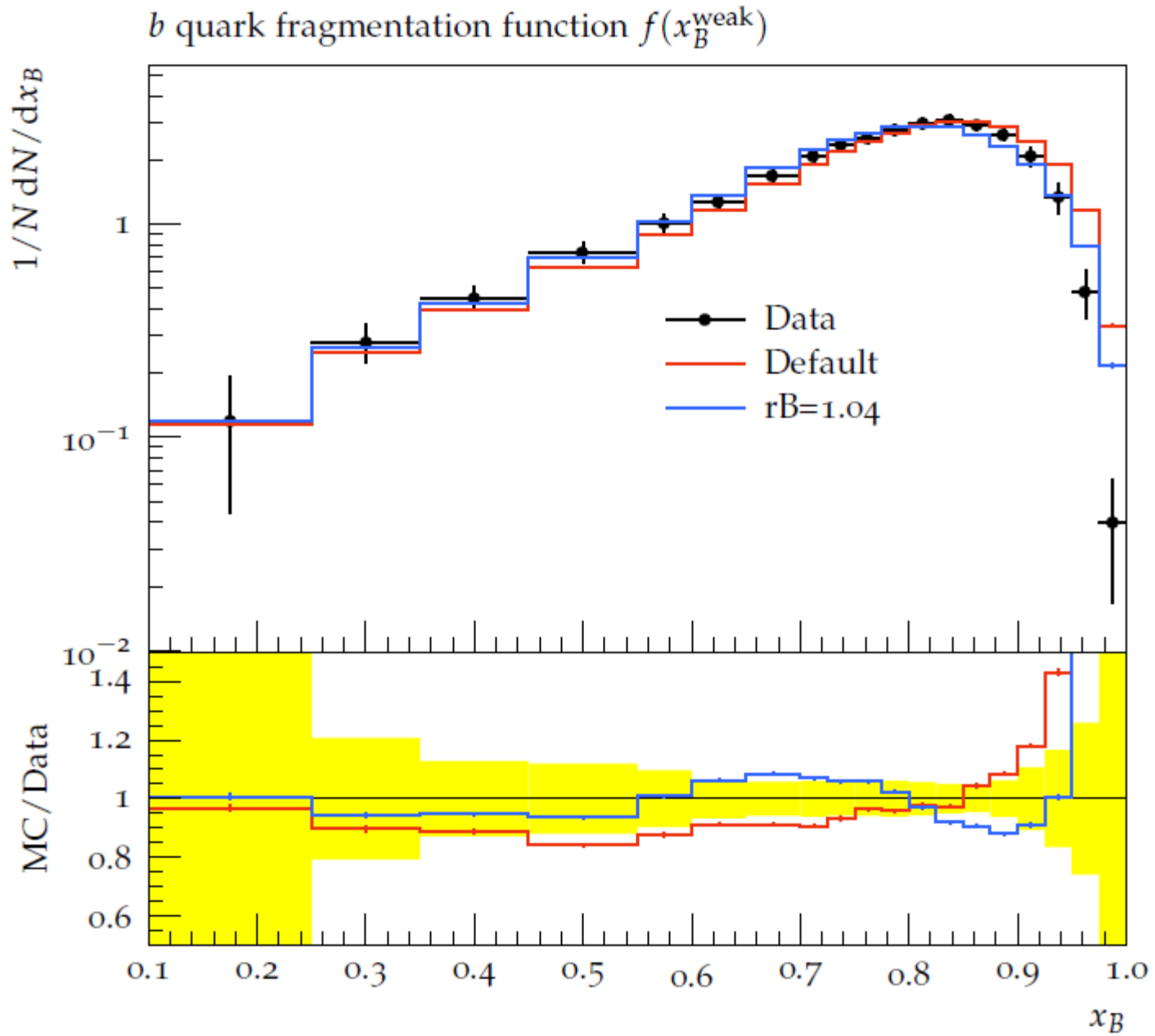
At LEP  $x_b$  was measured from the semi-leptonic decays with ( $x_b$  lead) or without ( $x_b$  weak) reconstructing the missing energy of the neutrino.

	$r_B$	$\Delta r_B$	$\chi^2/\text{dof}$
SLD (weak)	1,08	0,02	39/21
ALEPH (weak)	1,05	0,02	75/18
DELPHI (weak)	1,02	0,02	96/9
ALEPH (lead)	0,92	0,02	106/18
DELPHI (lead)	0,82	0,02	136/9
COMBINED	0,978	0,007	560/79

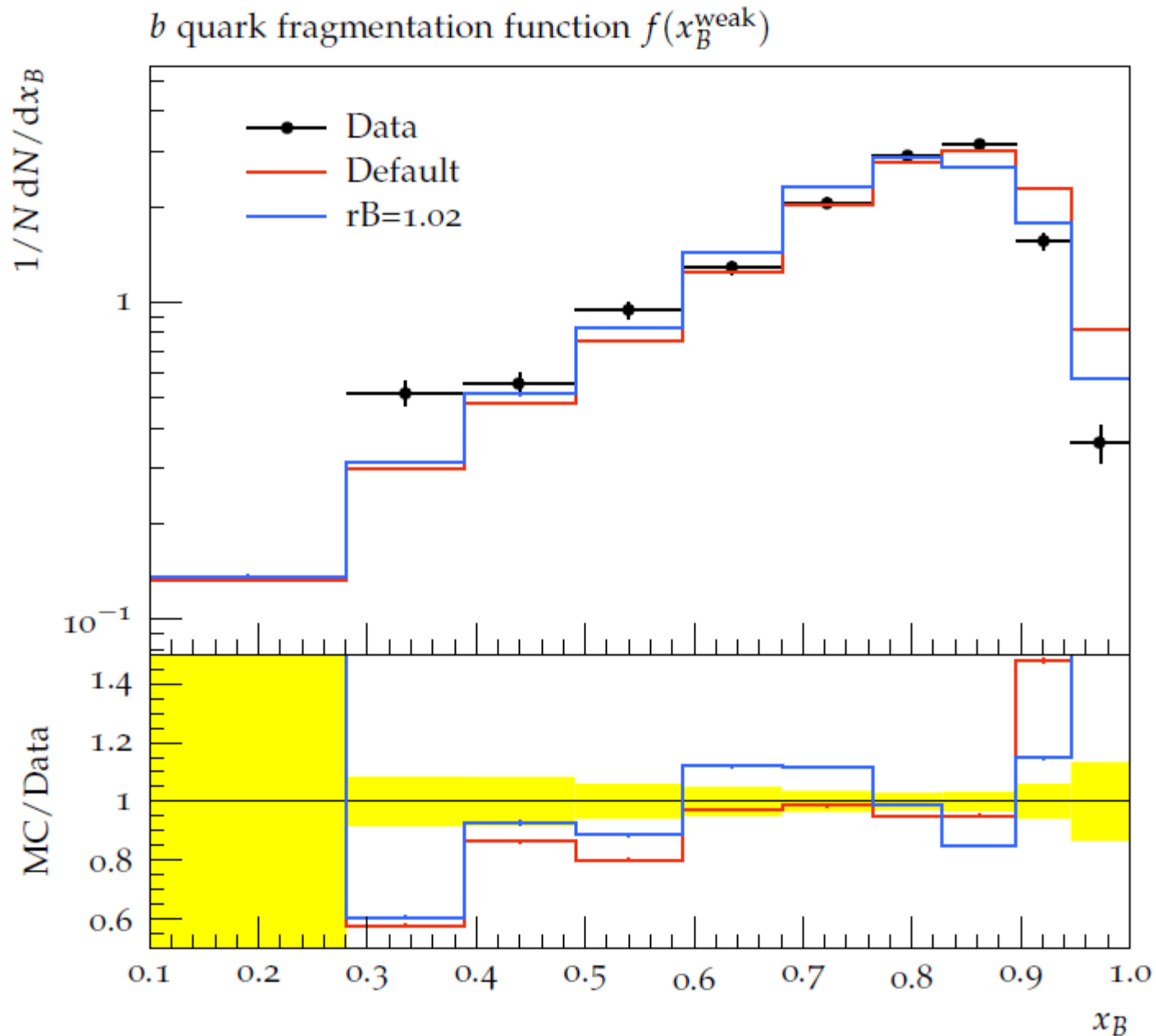
SLD Data for ATTBAR  $r_B = 1.07 \pm 0.02$   $\chi^2/\text{dof} = 39/21$



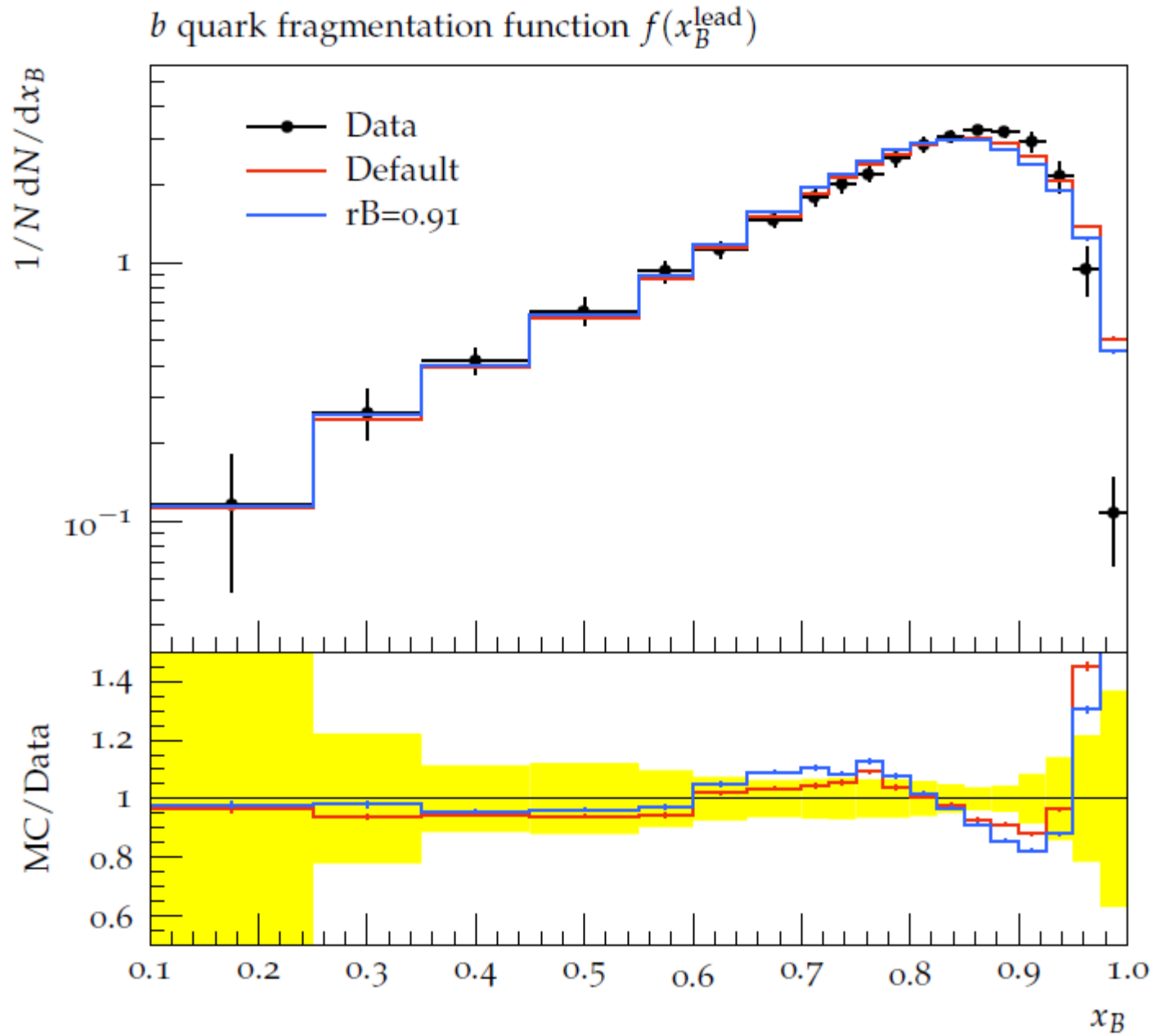
ALEPH Data for ATTBAR  $r_B = 1.04 \pm 0.02$   $\chi^2/\text{dof} = 75/18$



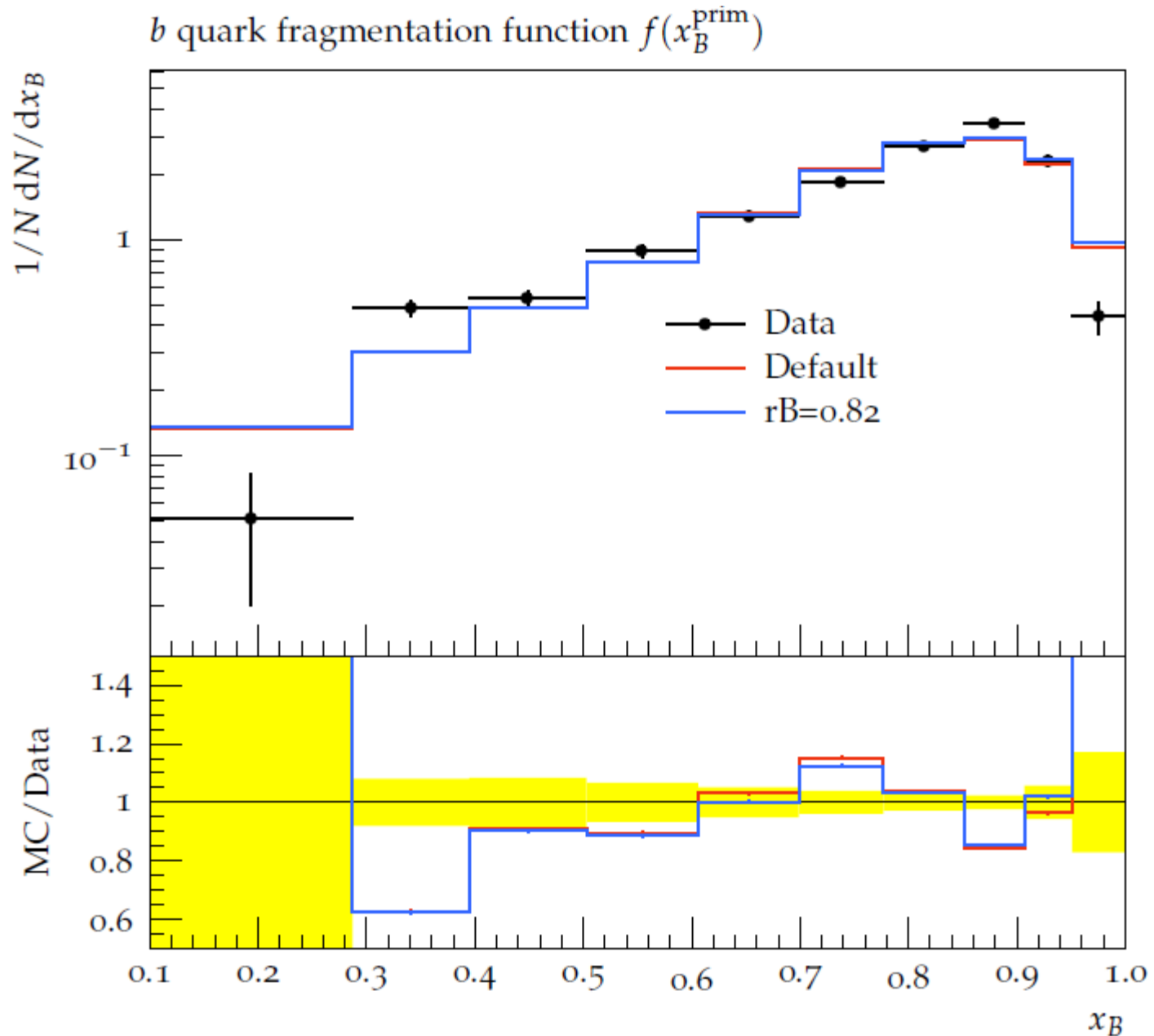
DELPHI Data for ATTBAR  $r_B = 1.02 \pm 0.02$   $\chi^2/\text{dof} = 96/9$



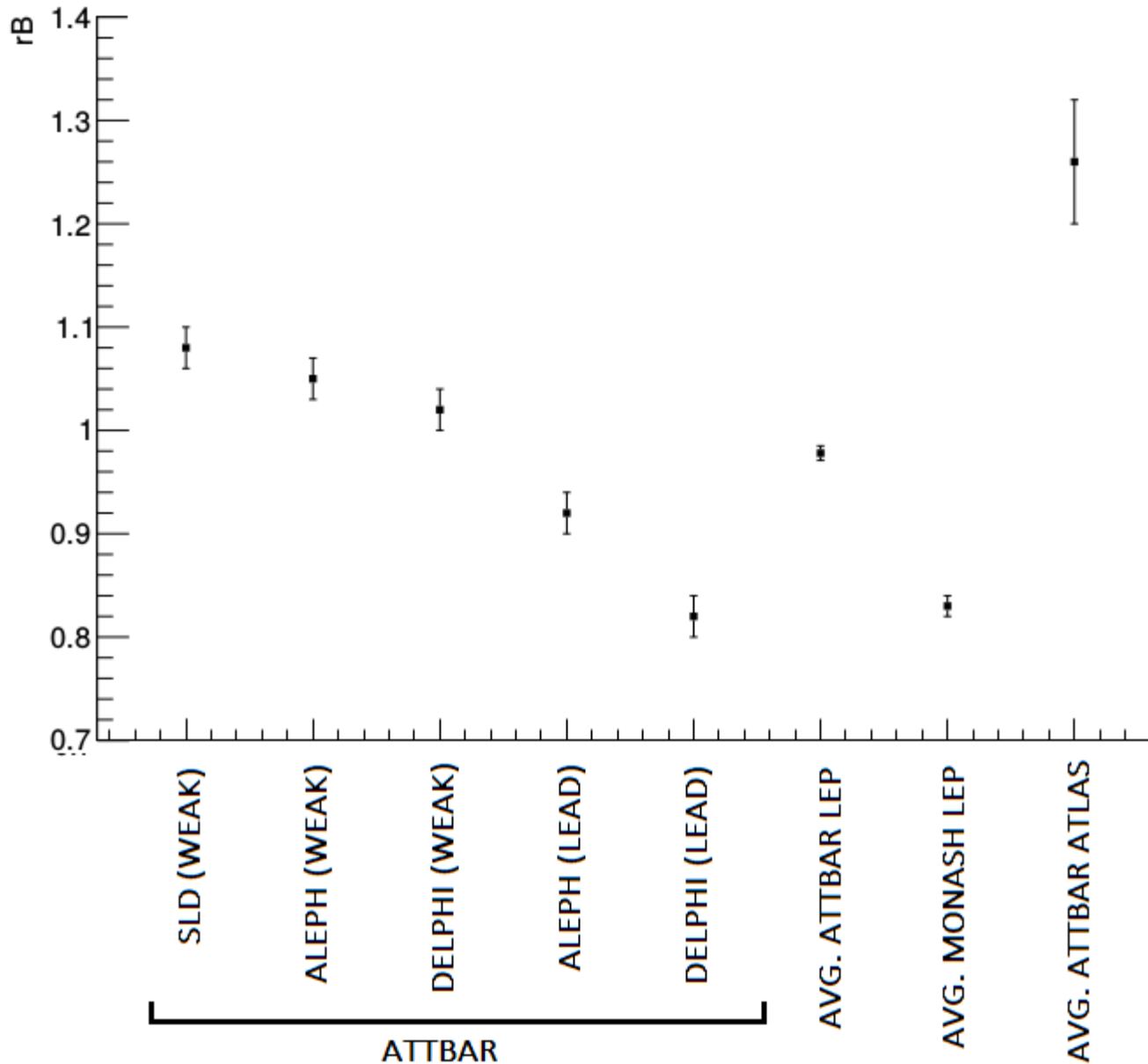
# ALEPH Data for ATTBAR $r_B = 0.91 \pm 0.02$ $\chi^2/\text{dof} = 106/18$



# DELPHI Data for ATTBAR $r_B = 0.82 \pm 0.02$ $\chi^2/\text{dof} = 136/9$

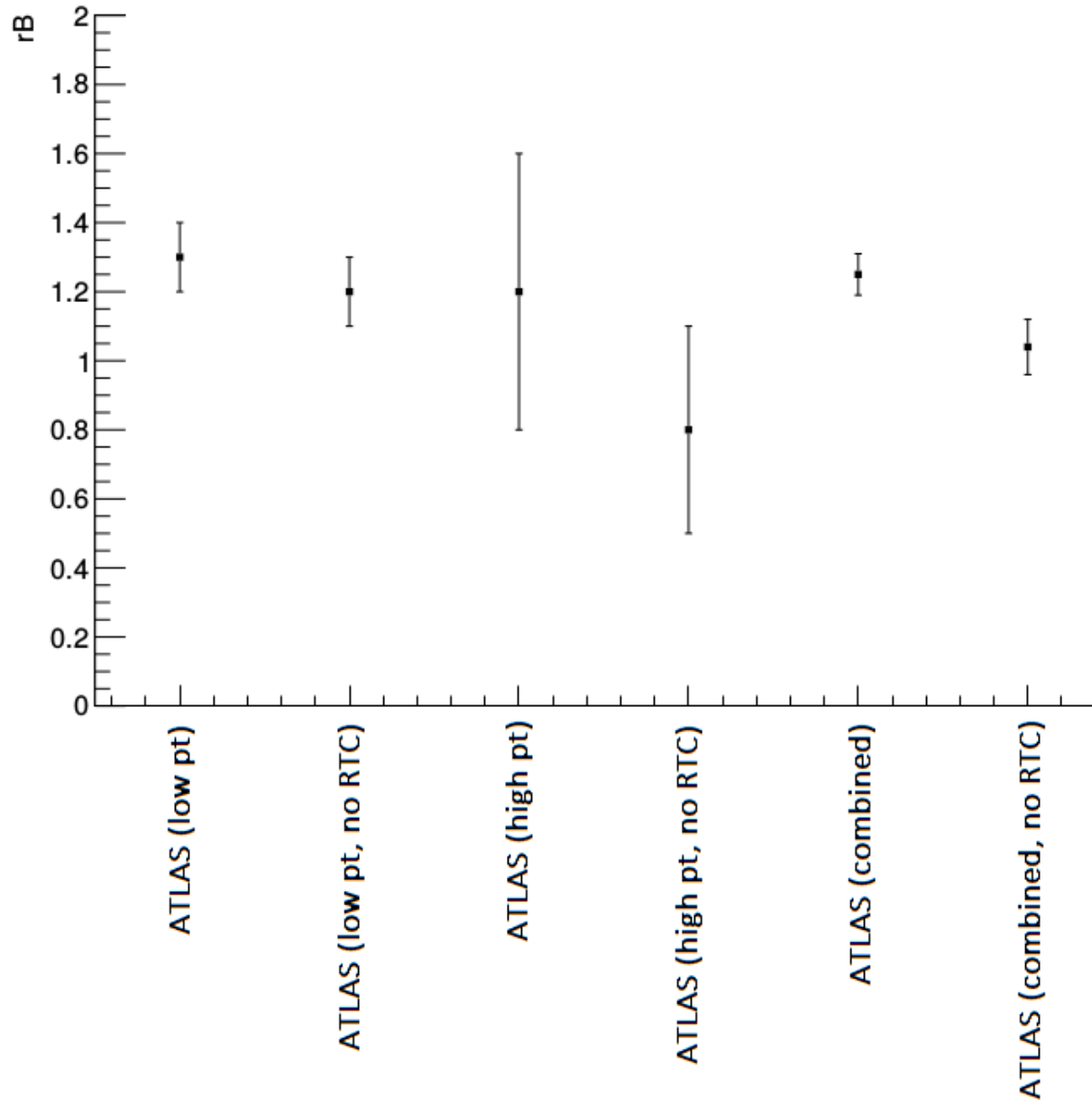


# Comparison of $r_B$ measurements (preliminary)

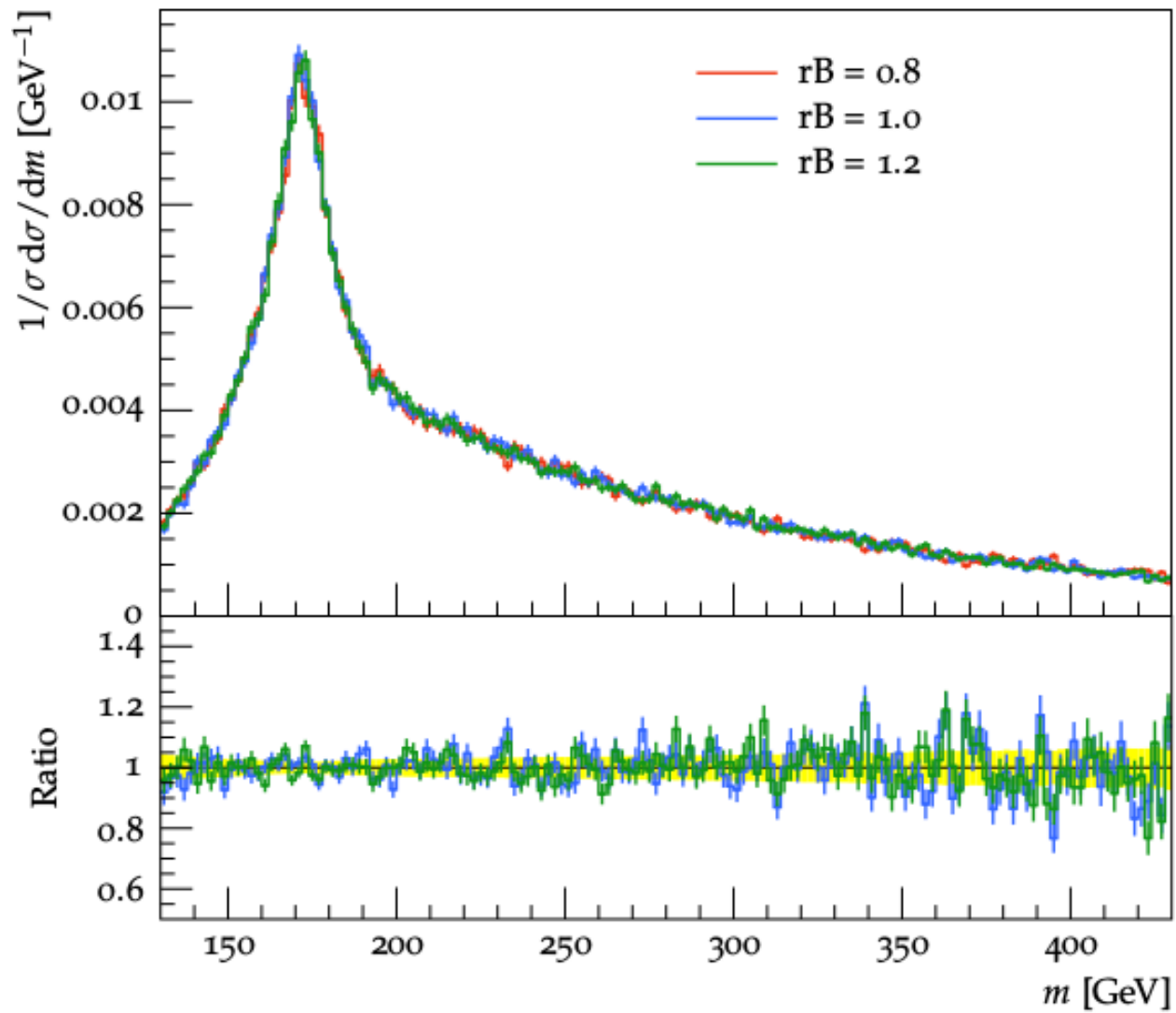




# Comparison of $r_B$ measurements for ATLAS (preliminary)



Mass distribution for reconstructed top



# Summary and Conclusions

- The precise determination of the b-fragmentation function is important to reduce the uncertainty on the top-quark mass.
- The  $r_B$  parameter of the b-fragmentation function in the Bowler modification of the Lund model has been extracted from ATLAS  $t\bar{t}$  b-jet shapes.
- The determination of  $r_B$  is compatible with the same extraction from LEP measurements of the b-fragmentation function.
- Some tension is observed between the  $x_b$  weak and  $x_b$  lead measurements at LEP. The former does not include the reconstructed missing energy of the neutrino.