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## Probing the local structure in Multiferroic $\text{SmCrO}_3$

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Rare-earth orthochromites of the formula  $\text{RCrO}_3$  ( $\text{R}=\text{Dy}, \text{Pr}, \text{Ho}, \text{Yb}, \text{Er}, \text{Y}, \text{Lu}, \text{Sm}$ ) are currently at the center of great controversy regarding ferroelectricity. While dielectric constant anomalies near 400 – 500 K in the heavier rare-earth chromites were associated with non-centrosymmetry, others claim that the polarization observed in these systems is due to the combined effect of the electric applied field, that breaks the symmetry, and exchange-field on the R ion from the Cr sub-lattice. Accordingly to these claims, no spontaneous ferroelectric polar-order exists in these systems and the presence of a magnetic R-ion is essential to induce a metastable ferroelectric state. Contrarily, the appearance of ferroelectricity without direct correlation to the magnetic order, arising from polar octahedral rotations and/or cation displacements, was recently claimed. [1,2]

Clearly, additional efforts are needed to definitely validate these claims. Since these properties might emerge from local structural landscapes that are not well described by long-range average structural methods, the use of local probe studies, such as Perturbed Angular Correlation (PAC) spectroscopy, provide relevant knowledge. In this work the  $\text{SmCrO}_3$  compound was studied. The temperature dependent of the electric field gradient (EFG) on  $\text{SmCrO}_3$  compound was followed, using the  $^{111}\text{Cd}$  PAC probe, in the 16 K < T < 723 K temperature range. A temperature range that spans over the important transition temperatures, namely the reported ferroelectric transition ( $T_{FE} \approx 220$  K), the magnetic ordering of Cr atoms sub-lattice ( $T_N^{\text{Cr}} = 133$  K), the spin reorientation ( $T_{SR} = 34$  K) and magnetic ordering of Sm atoms sub-lattice ( $T_N^{\text{Sm}} = 20$  K). The  $^{111m}\text{Cd}$  implantation and  $^{111}\text{Cd} \rightarrow ^{111}\text{In}$  diffusion was followed by an annealing at high temperatures in air.

At high temperatures,  $T > 300$  K, a frequency triplet corresponding to a single EFG, *i.e.*, one probe local environment, was observed and in this temperature range no significant changes occur in the spectra when the temperature is lowered. However, below 300 K visible changes can be observed in the perturbation function ( $R(t)$ ) data and in the corresponding Fourier transforms. In detail, a second EFG emerges and its relative abundance increases with decreasing temperature. Accordingly, the fits to the  $R(t)$  experimental data were performed considering only one static EFG distribution, which was assumed to be Lorentzian-like, for  $T > 300$  K while two EFG distributions had to be considered to account for the features that emerge below that temperature.

The spectra obtained at high temperatures revealed an EFG characterized by a  $V_{ZZ}^{\text{Sm1}} \approx 76$  V/m<sup>2</sup> and an asymmetry parameter  $\eta \approx 0.2$  in good agreement with similar systems. The second EFG, that emerges at low temperatures, is characterized by a similar fundamental frequency but a higher asymmetry parameter  $\eta \approx 0.6$ . From our data we observed that a distortion of the high temperature local environment start to develop below 300 K within the paramagnetic phase. Although our data might be compatible with the most recent reports, where polar octahedral rotations and/or cation displacements are at the origin of a polar order in the paramagnetic state, remarkably, our results point to a more subtle scenario, where locally an inhomogeneous state emerges. In this new state regular and distorted environments (most probably polar and non polar states) coexist.

[1] <https://doi.org/10.1103/PhysRevB.86.214409>.

[2] <https://doi.org/10.1209/0295-5075/107/47012>.

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