Production of strange particles in jets in p–p, p–Pb and Pb–Pb collisions measured with ALICE

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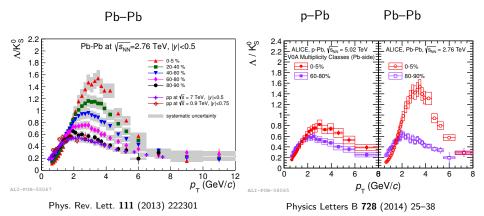






Motivation for PID in jets

- Baryon-to-meson ratio is enhanced in Pb–Pb and p–Pb collisions.
- This phenomenon cannot be explained by fragmentation in vacuum.
- What is the effect of QGP on hadronization mechanism(s) in jets?
- What are the mechanisms (parton recombination)?

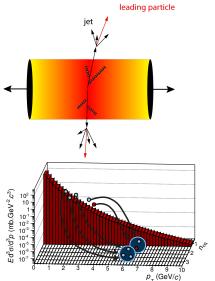


Motivation for PID in jets

We aim to understand the origin(s) of the Λ/K_S^0 enhancement by separating hadrons produced in hard processes (jets) from hadrons produced in soft processes (underlying event).

Is the baryon-to-meson ratio enhanced due to the collective effects in the plasma (parton recombination, radial flow,...) or is it (also) due to a modification of the jet fragmentation in the medium?

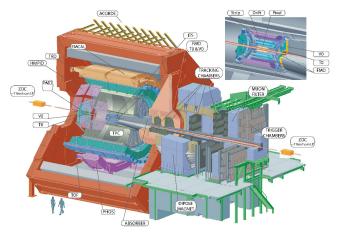
- jet fragmentation
 A high-p_T parton from hard scattering fragments into hadrons.
- parton recombination Multiple partons cluster together to form a hadron.



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ALICE

- ▶ collisions studied: p–p at $\sqrt{s} = 7$ TeV, p–Pb at $\sqrt{s_{\rm NN}} = 5.02$ TeV, Pb–Pb at $\sqrt{s_{\rm NN}} = 2.76$ TeV
- tracking of charged particles by ITS & TPC in magnetic field of 0.5 T
- centrality estimated from the multiplicity of charged particles in the detectors at forward pseudorapidities



Analysis of charged jets

- track selection
 - charged primary particles
 - $p_{\mathrm{T}}^{\mathrm{track}} > 150 \,\mathrm{MeV}/c$
 - uniform in $\phi imes \eta$, $|\eta_{\mathrm{track}}| < 0.9$
- raw-jet reconstruction
 - anti-k_t algorithm
 - resolution parameter R = 0.2, (0.3, 0.4)
- subtraction of average soft background
 - average background density ρ estimated from the median $k_{\rm t}$ cluster
 - $P_{\rm T}^{\rm jet,ch,corr} = p_{\rm T}^{\rm jet,ch,raw} \rho A_{\rm jet,ch}, \qquad ({\rm where} \ A_{\rm jet,ch} \ {\rm is} \ {\rm jet} \ {\rm area})$
- signal-jet selection (good candidates for hard scattering)
 - $p_{T}^{\text{leading track}} > 5 \text{ GeV}/c \text{ (only Pb-Pb)}$
 - $A_{\rm jet,ch} > 0.6\pi R^2$
- further $p_{\rm T}^{\rm jet,ch}$ corrections
 - background anisotropy (intra-event p_T fluctuations)
 - detector response

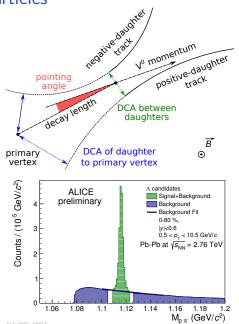
Analysis of neutral strange particles

Strange neutral particles decaying into two charged daughter particles

- $\blacktriangleright \text{ meson } \mathsf{K}^0_S \rightarrow \pi^+ + \pi^- \text{ (BR 69\%)}$
- baryon $\Lambda \rightarrow p + \pi^-$ (BR 64%)

Mother V^0 particle reconstructed using topology of its V-shaped decay.

Combinatorial background suppressed by cuts on decay parameters. Signal yield extracted from the invariant-mass distribution.



Strange particles in jets

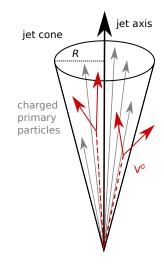
Analysis steps

- V⁰ candidate selection
- candidate-jet matching (V⁰s in jet cones)

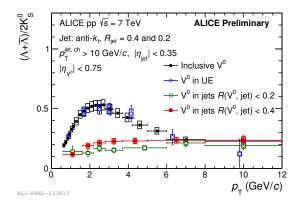
$$\sqrt{(\phi_{\mathsf{V}^0}-\phi_{\mathsf{jet},\mathsf{ch}})^2+(\eta_{\mathsf{V}^0}-\eta_{\mathsf{jet},\mathsf{ch}})^2} < R,$$

 $|\eta_{\rm jet,ch}|^{\rm max} < |\eta_{\rm V^0}|^{\rm max} - R$

- candidate-UE matching (V⁰s in events without selected jets with p_T^{jet,ch} > 5 GeV/c)
- signal extraction (invariant-mass distribution)
- efficiency correction (in jet cones, in UE)
- subtraction of V⁰s in UE
- subtraction of V⁰s coming from decays of jet constituents (Ξ → Λ), i.e. "feed-down" correction

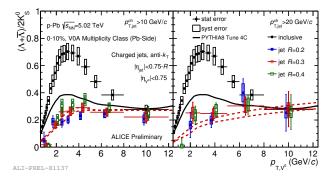


Λ/K_S^0 ratio in charged jets in p-p at $\sqrt{s} = 7 \text{ TeV}$ (127 × 10⁶ minimum-bias collisions)



- The ratio in UE is consistent with the inclusive ratio.
- ► The ratio in jets is clearly different from the inclusive ratio at low and intermediate p_T^{V⁰}.
- ► A slight increase of the ratio in jets with increasing *R*.

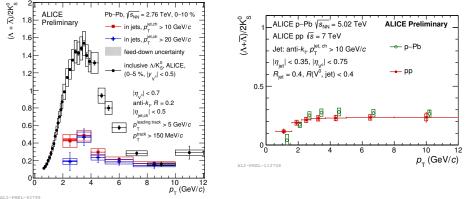
Λ/K_S^0 ratio in charged jets in p–Pb at $\sqrt{s_{NN}}=5.02\,\text{TeV}$ (high-multiplicity collisions, 0–10%)



The ratio in jets

- ▶ is clearly different from the inclusive ratio at low and intermediate $p_T^{V^0}$,
- is different from the inclusive ratio in PYTHIA (black line),
- is similar to the ratios in PYTHIA jets (red dashed lines),
- shows no significant dependence on $p_{\rm T}^{\rm jet,ch}$ and a slight dependence on R.

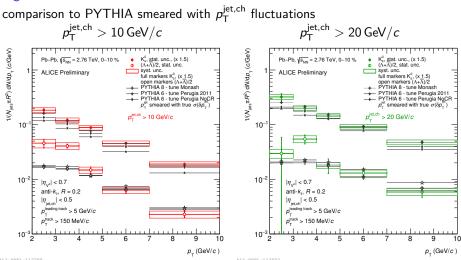
Λ/K_S^0 ratio in charged jets in Pb–Pb at $\sqrt{s_{NN}} = 2.76$ TeV (7.4 \times 10⁶ central collisions, 0–10%)



The ratio in jets

- ▶ is clearly different from the inclusive ratio at low and intermediate $p_{\rm T}^{V^0}$,
- shows no significant dependence on p_T^{jet,ch}
- ▶ is consistent with the ratio in jets in p–Pb and p–p at $p_T^{V^0} > 4 \text{ GeV}/c$.

K_{S}^{0} , Λ spectra in charged jets in Pb–Pb



- Same slopes of spectra from measurement and from PYTHIA.
- Enhancement for Λ at $p_{\rm T}^{\rm V^0} < 4 \,{\rm GeV}/c$.

Summary and outlook

ALICE has performed the first measurement of the Λ/K_S^0 ratio in charged jets in p–p, p–Pb and Pb–Pb collisions at the LHC.

Message

- In every collision system, the Λ/K_S^0 ratio in jets is significantly smaller than the inclusive ratio (and the UE).
- ► The Λ/K_S^0 ratios in jets are consistent within uncertainties in all collision systems for $p_T^{V^0} > 4 \text{ GeV}/c$.
- The dominant source of the enhancement are soft processes associated with collective behaviour.
- A potential modification of jet fragmentation seems to be restricted to the region p^{V⁰}_T < 4 GeV/c and manifest by an enhancement of the Λ yields.

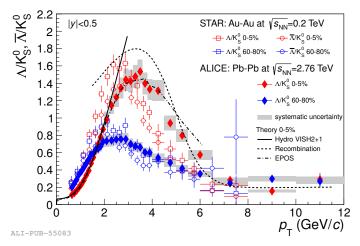
Outlook

- Multiplicity dependence in p-p and p-Pb collisions.
- Better resolution in Pb–Pb collisions in Run 2.

Thank you for your attention.

Backup

Comparison of data with models



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V^0 candidate selection

Г

| Cut variable | Value |
|--|--|
| Daughter tracks | |
| TPC refit | true |
| type of production vertex | not kKink |
| DCA to the primary vertex | $\geq 0.1{ m cm}$ |
| DCA between daughters | $\leq 1\sigma_{TPC}$ |
| $ \eta $ | ≤ 0.8 |
| V ⁰ candidate | |
| reconstruction method | offline |
| cosine of the pointing angle (CPA) | \geq 0.998 |
| radius of the decay vertex | 5–100 cm |
| $ \eta $ | ≤ 0.7 |
| transverse proper lifetime | $\leq 5	au$ |
| Armenteros–Podolanski cut (K ⁰ _S) | $p_{\mathrm{T}}^{\mathrm{Arm.}} \geq 0.2 lpha^{\mathrm{Arm.}} $ |

Jet algorithms

A sequential recombination jet finder is defined according to this general scheme:

1. $\forall i, j$: calculate distances d_{ij} and d_{iB} (NB $k_t \equiv p_T$):

$$d_{ij} = \min\left(k_{t,i}^{2p}, k_{t,j}^{2p}\right) \frac{\Delta_{ij}^2}{R^2}, \quad \Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2, \quad d_{iB} = k_{t,i}^{2p}$$

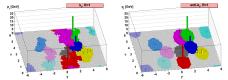
2. Find d_{min}:

$$d_{\min} = \min\left(d_{ij}, d_{iB}\right).$$

- If ∃ i, j : d_{min} = d_{ij}, merge particles i and j into a single particle and combine their momenta.
- If $\exists i : d_{\min} = d_{iB}$, declare particle *i* to be a final jet and remove it from the list.

These steps are repeated until no particles are left.

$$p = \begin{cases} 1 & k_{\rm t} \text{ (background estimation} \\ 0 & \text{Cambridge/Aachen} \\ -1 & \text{anti-}k_{\rm t} \text{ (signal jets)} \end{cases}$$



Matteo Cacciari et al. JHEP 0804 (2008) 063

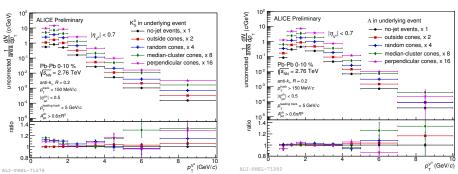
Estimation of V⁰s in the underlying event in Pb–Pb

- no-jet events: V⁰s in events with no selected jets
- outside cones: V⁰s outside jet cones

 K_{s}^{0}

- ▶ random cones: V⁰s in a randomly oriented cone
- median-cluster cones: V^0 s in the cone of the median k_t -cluster
- ▶ perpendicular cones: V⁰s in cones perpendicular to the jet in azimuth

Methods differ in regions, events, statistics, efficiency.

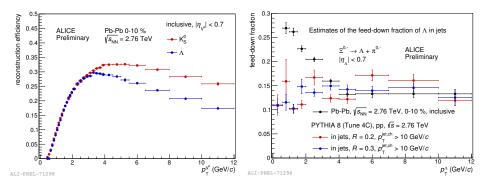


Λ

Reconstruction efficiency, feed-down in Pb-Pb

Reconstruction efficiency depends strongly on $p_T^{V^0}$ and η_{V^0} . Shape of the measured η_{V^0} distribution depends on the selection criteria. Efficiency of inclusive V⁰s is reweighted to get efficiency in jet cones and UE. Feed-down fraction of Λ in jets estimated from:

- inclusive A (Pb–Pb-like),
- jets generated by PYTHIA 8 (p-p-like).



Reweighting of the reconstruction efficiency

- $\blacktriangleright~\epsilon$ reconstruction efficiency of inclusive particles
- ▶ ϵ_s reconstruction efficiency of particles of interest (scaled ϵ)
- ▶ *a_s* yield of associated particles of interest
- ▶ g_s yield of generated particles of interest
- \blacktriangleright *m* uncorrected yield of measured particles (candidates) of interest
- t yield of true (corrected) particles of interest
- P signal purity

Signal extraction in JC, UE (assume that $P_{\text{inclusive}}(p_T^{V^0}, \eta_{V^0})$ is the same as for V⁰s of interest):

$$m(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})=m_{\mathsf{raw}}(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})|_{\mathsf{peak region}}\cdot P_{\mathsf{inclusive}}(p_{\mathsf{T}}^{\mathsf{V}^0},\eta_{\mathsf{V}^0})|_{\mathsf{peak region}}$$

Efficiency calculation:

$$\begin{aligned} a_{s} \equiv m, \quad \sigma_{a_{s}} \equiv 0, \qquad g_{s} = a_{s}/\epsilon \\ \frac{1}{\epsilon_{s}(p_{\mathsf{T}}^{\mathsf{V}^{0}})} = \frac{\sum_{\eta_{\mathsf{V}^{0}_{j}}} g_{s}(\eta_{\mathsf{V}^{0}_{j}}, p_{\mathsf{T}}^{\mathsf{V}^{0}})}{\sum_{\eta_{\mathsf{V}^{0}_{j}}} a_{s}(\eta_{\mathsf{V}^{0}_{j}}, p_{\mathsf{T}}^{\mathsf{V}^{0}})} = \sum_{\eta_{\mathsf{V}^{0}_{j}}} \frac{a_{s}(\eta_{\mathsf{V}^{0}_{j}}, p_{\mathsf{T}}^{\mathsf{V}^{0}})}{\sum_{\eta_{\mathsf{V}^{0}_{j}}} a_{s}(\eta_{\mathsf{V}^{0}_{j}}, p_{\mathsf{T}}^{\mathsf{V}^{0}})} \frac{1}{\epsilon(\eta_{\mathsf{V}^{0}_{j}}, p_{\mathsf{T}}^{\mathsf{V}^{0}})} \end{aligned}$$

Spectra correction:

$$t = m/\epsilon_s$$

Sources of systematic uncertainties

- \blacktriangleright cuts for the selection of V^0 candidates (cuts varied)
- signal extraction (parameters varied)
- estimation of spectra of V⁰s in the underlying event (multiple methods)
- estimation of material budget (from another analysis)
- estimation of feed-down fraction of Λ and $\overline{\Lambda}$ in jets (PYTHIA as alternative)
- correction of p_T^{jet} (embedding of simulated jets in real data)
- detector performance $(p_T^{V^0} resolution)$

Systematic uncertainties in p-Pb

| source | uncertainty |
|--------------------------------|---|
| selection cuts | 2–5 % for K_{S}^{0} , 3–6 % for Λ |
| signal extraction | 6 % (10 %) for $p_{ m T}^{ m jet,ch} > 10~{ m GeV}/c~(20~{ m GeV}/c)$ |
| V ⁰ s in UE | 10 % (2 %) at low (high) $p_{\rm T}^{\rm V^0}$ |
| $p_{\rm T}^{\rm jet,ch}$ scale | 1 % (10 %) at low (high) $\rho_{\rm T}^{\rm V^0}$ |
| feed-down | 5 % |