Exploring the phase structure of minimally doubled fermions

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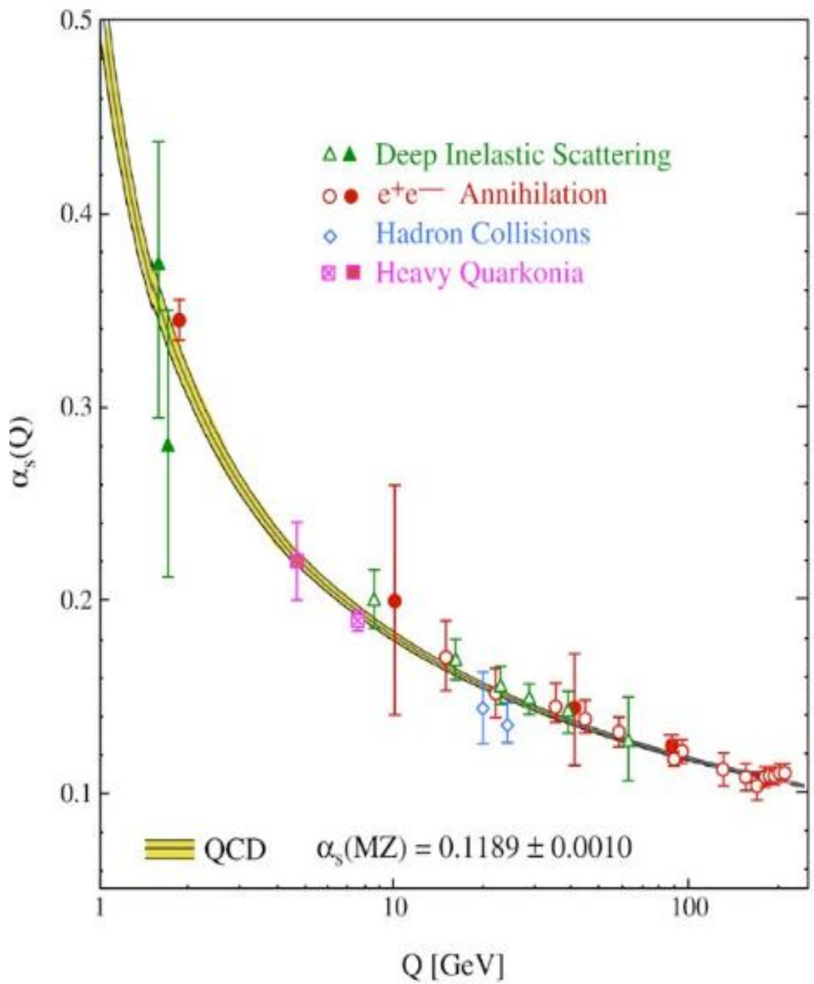


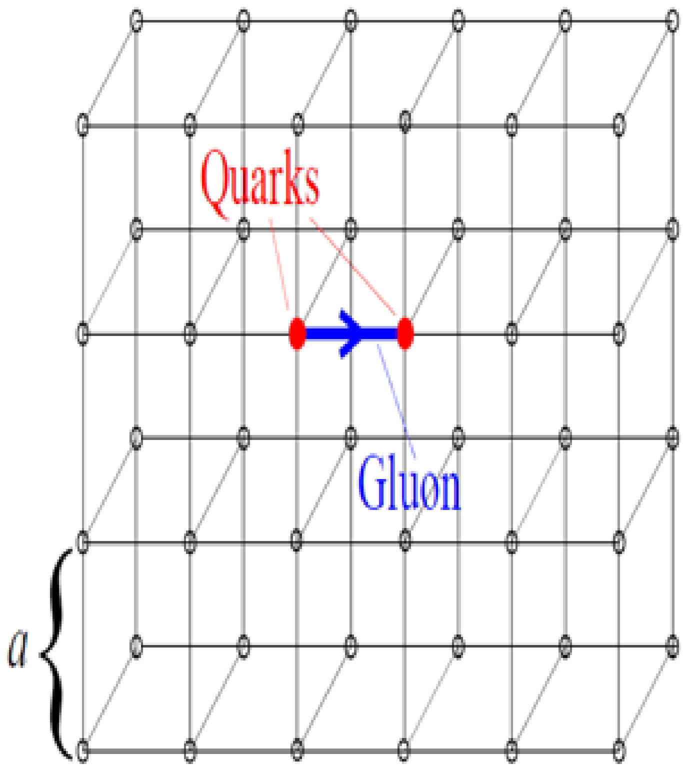
Analytic or perturbative solutions in low-energy QCD are hard or impossible due to the highly nonlinear nature of the strong force and the large coupling constant at low energies.

We can't calculate mesons and baryons masses, even when we know the quark masses and the coupling constant

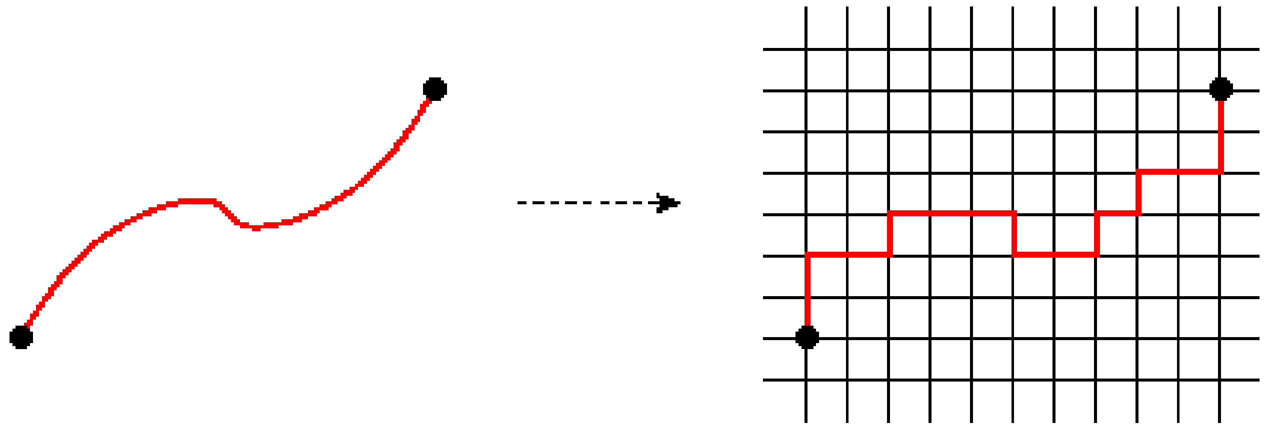
Need for different formulations, simulations

Introduction

- Proposed by Wilson, 1974.
- Non- perturbative low-energy solution of QCD.
- The continuum spacetime discretized on 4d Euclidean space-time lattice.
- Quarks can only exist in lattice points and gluons are the links between them
- Solved by large scale numerical simulations on supercomputers.



Introduction



- From continuum to discretized lattice:

$$\int d^4x \to a^4 \sum_n$$

n - four-vector that labels the lattice site,

a - lattice constant

- Fermions and gauge action are discretized
- Take an appropriate continuum limit ($a \to 0$) to get back the continuum theory.

Introduction

- Fermions action discretization:
  - Doubling problem
  - If no doublers, chiral symmetry is broken

Nielsen – Ninomiya No – go theorem *(Nielsen, Ninomiya, 1981):*

"*A local, real, free fermion lattice action, having chiral and translational invariance, necessarily has fermion doubling* "

Introduction

- Doubling problem : obstacle to simulations

Several ways to bypass No-Go theorem, but….

- Wilson ➡ Broken chiral symmetry
- DW or Overlap ➡ Numerical expensive (Non exact locality)
- Staggered ➡ Rooting procedure (4 tastes)

➢ Another possibility : **Minimally doubled fermions**

i) 2 flavors ← 4 in Staggered

ii) Exact chiral: $U(1)_A \subset SU(2)$ ← Broken in Wilson

iii) Strict locality ← Not strict in DW or Overlap

Symmetries

Translation, Gauge, $U(1)_B$, $\underline{U(1)_A} \subset SU(2)$

P.F.Bedaque, *et.al.*, PLB **662**, 449 (2008)

Common with all Minimal-doubling actions.

- <u>Discrete symmetries</u>

1. Karsten-Wilczek
   <u>CT, P, Cubic, Z_2</u>
2. Creutz-Borici
   <u>CPT, S_4, Z_2</u>
3. Twisted-Ordering
   <u>CPT, Z_4, Z_2</u>

With gauge interaction, broken symmetry leads to anisotropy.

↓

Fine-tuning to cancel redundant operators is required.

Boriçi – Creutz action

Boriçi – Creutz fermionic action with the free Dirac operator (in the momentum space):

$$D(p) = \sum_{\mu} i\gamma_{\mu} \sin p_{\mu} + \sum_{\mu} i\gamma'_{\mu} \cos p_{\mu} - 2i\Gamma$$

This operator has zeros: $p_1 = (0,\ 0,\ 0,\ 0)$ dhe $p_2 = (\pi/2,\ \pi/2,\ \pi/2,\ \pi/2)$.

- There is always a special direction in euclidean space (given by the line that connects these two zeros: hypercubic diagonal)

⬇

- Thus, these actions cannot maintain a full hypercubic symmetry *(P. F. Bedaque et al, 2008).*

⬇

Hypercubic symmetry has to be restored
(Perturbative calculations *Capitani et al, 2010*)

$$D_{BC}(p) = \sum_{\mu} [i\gamma_{\mu} \sin p_{\mu} + i(\Gamma - \gamma_{\mu}) \cos p_{\mu}] + i(c_3 - 2)\Gamma$$

Metodology

- Evaluation of the broken hypercubic symmetry mass

$$\Delta(M^2_{\pi^+}) = (M^{(1,0,0,0)}_{\pi^+})^2 - (M^{(1,1,1,1)}_{\pi^+})^2$$

$$(M^{(1,0,0,0)}_{\pi^+})^2 = (M^{(1,1,1,1)}_{\pi^+})^2 = 0 \quad (M^2_{\pi^+} \xrightarrow[m_{q\to 0}]{} 0\ )$$

- Point splitting (*u(x)*, *d(x)*)
- Charged pion propogator calculation
- Calculation of the charged pion mass from two different directions

- Non perturbative restaturation of broken hypercubic symmetry

- 1 –st Method :: $\Delta(M^2_{\pi^+}) = (M^{(1,0,0,0)}_{\pi^+})^2 - (M^{(1,1,1,1)}_{\pi^+})^2 = 0$

$$M^2_{\pi^+} \xrightarrow[m_{q\to 0}]{} 0$$

- 2-nd Method: Exploring the phase – structure of QCD in T=0K

(*Banks- Casher relation and Lanczos quadrature)*

Evaluation of the broken hypercubic symmetry mass (*preliminary results*)

- Calculation for the charged pion mass from two different directions

Details of the simulations:

- Lattices 8^4, 12^4 dhe 16^4
- Quenched approximation
- Wilson gauge action
- Boriçi – Creutz action
- CGNE inverter
- Five different quark masses

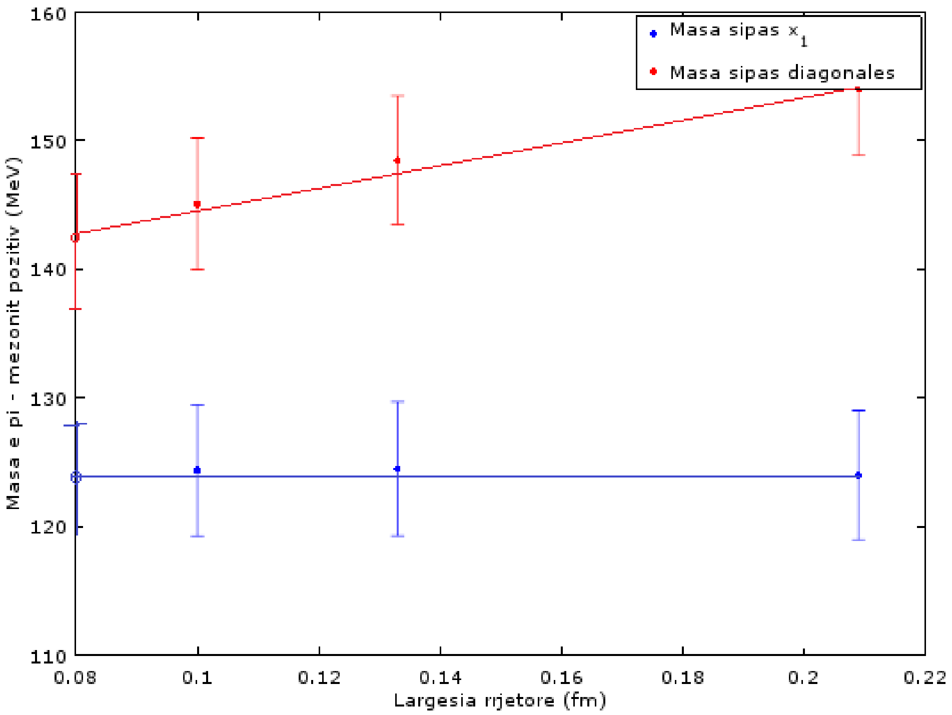
Evaluation of the broken hypercubic symmetry mass (*preliminary results*)

	Lattice constant (fm)	*Lattice constant (MeV)*	*Effective mass*	*Physical Mass (MeV)*	Error*(MeV)*
L=8	*a = 0.209*	a^{-1}*=956.9*	*am=0.161*	*m=154.06*	*± 5,2 MeV*
L=12	*a = 0.133*	a^{-1}*=1500*	*am=0.099*	*m=148.5*	*± 5 MeV*
L=16	*a = 0.1*	a^{-1}*=2000*	*am=0.074*	*m=148*	*± 5,1 MeV*

Results from the hypercubic diagonal mass calculations.

	Lattice constant (fm)	*Lattice constant (MeV)*	*Effective mass*	*Physical Mass (MeV)*	Error*(MeV)*
L=8	*a= 0.209*	a^{-1}*=956.9*	*am=0.130*	*m=124*	*± 5MeV*
L=12	*a= 0.133*	a^{-1}*=1500*	*am=0.083*	*M=124.5*	*± 5,2MeV*
L=16	*a= 0.1*	a^{-1}*=2000V*	*am=0.062*	*M=124.39*	*± 5,1MeV*

Results from one of hypercubic direction mass calculations.



$$m_{\pi^+(diag)} = 141.54 \pm 5 MeV$$

$$m_{\pi^+(x_1)} = 123.89 \pm 5 MeV$$

Banks-Casher relation

- Banks – Casher relation (*T. Banks, A. Casher, 1980*) provides a link between the chiral condensate and the spectral density $\rho(\lambda, m)$

$$\lim_{\lambda \to 0} \lim_{m \to 0} \lim_{V \to \infty} \rho(\lambda, m) = \frac{\Sigma}{\pi} \qquad \Sigma = -\lim_{m \to 0} \lim_{V \to \infty} \langle \bar{\psi}\psi \rangle$$

(*If chiral symmetry is spontaneously broken by a non-zero value of the condensate the density of the quark modes in infinite volume does not vanish at the origin. A non-zero density conversely implies that the symmetry is broken)*

- Instead of the spectral density, the average number ν(M,m) of eigenmodes of the Dirac operator with eigenvalues α ≤ M2 turns out to be a more convenient quantity to consider. Since

$$\nu(\Lambda) = V \int_{-\Lambda}^{\Lambda} d\lambda \rho(\lambda, m)$$

the mode number ultimately carries the same information as the spectral density.

- So we can use: $$\Sigma_{eff} = \frac{\pi}{2} \frac{\nu(\Lambda)}{\Lambda V} \qquad \Lambda = \sqrt{M^2 - m^2}$$

The algorithm for evaluating chiral condensate

Algorithm : Lanczos algorithm for solving $Ax = b$

Let be $A = D^* D$ and $b = z_2$

Set $\beta_0 = 0,\ \rho_1 = 1/||b||_2,\ q_0 = o,\ q_1 = \rho_1 b$

for $i = 1, \dots$ **do**

  $v = Aq_i$

  $\alpha_i = q_i^\dagger v$

  $v := v - q_i\alpha_i - q_{i-1}\beta_{i-1}$

  $\beta_i = ||v||_2$

  $q_{i+1} = v/\beta_i$

  $\rho_{i+1} = -\frac{(\rho_i\alpha_i + \rho_{i-1}\beta_{i-1})}{\beta_i}$

  if $|\frac{\rho_1}{\rho_{i+1}}| < tol$ **then**

    $n = i$

    stop

  end if

end for

Let $A \in \mathbb{C}^{N\times N}$ be a hermitian matrix and $b \in \mathbb{R}^N$ a starting vector. Then the following algorithm computes the Gauss - Lanczos quadrature [17, 16]

Algorithm 1 Algorithm for the Gauss - Lanczos quadrature

Compute α_i and β_i using Lanczos algorithm for $Ax = b$

Set $(T_n)_{i,i} = \alpha_i,\ (T_n)_{i+1,i} = (T_n)_{i,i+1} = \beta_i$ otherwise $(T_n)_{i,j} = 0$

Compute eigenvalues λ_i and eigenvectors v_i of T_n, where $i = 1....n$

Sort eigenvalues and eigenvectors in the increasing order of eigenvalues

Set k as the maximum index which correspond to the cut-off eigenvalue

Set θ_i to the positive square root of the original eigenvalues

Set z_i the first element of eigenvectors v_i where $i = 1....n$

Set $\omega_i = z_i^2$

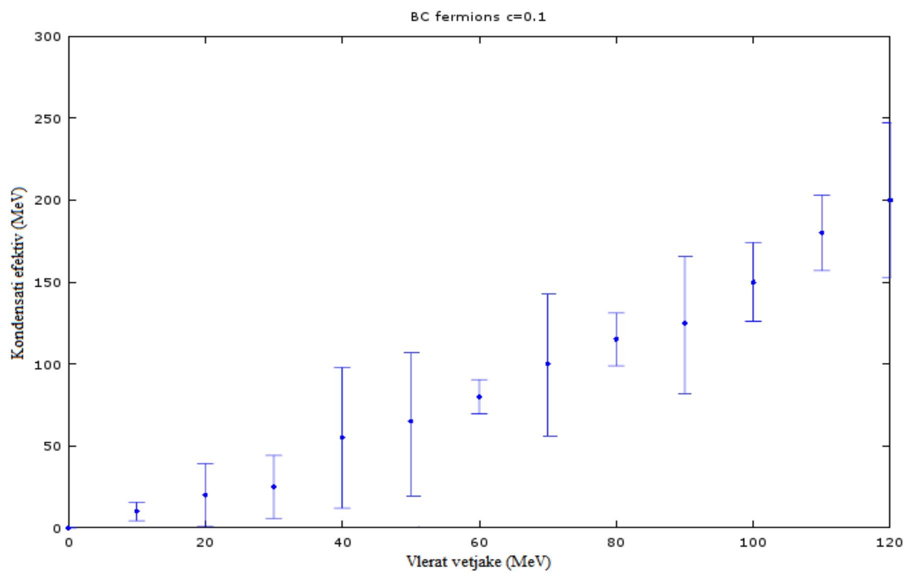
Compute the mode number $\nu_k = \sum_{i=1}^{k} \omega_i$

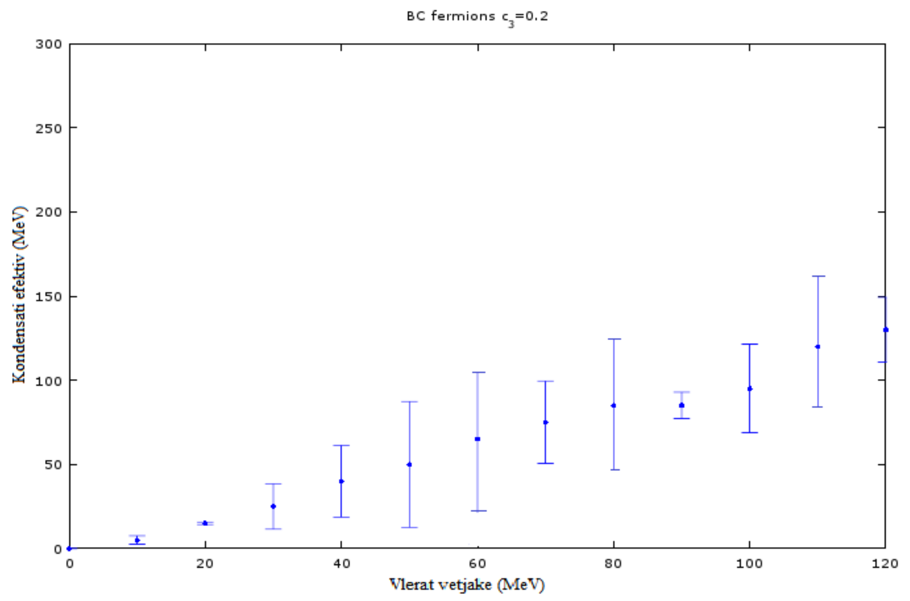
Details of simulations

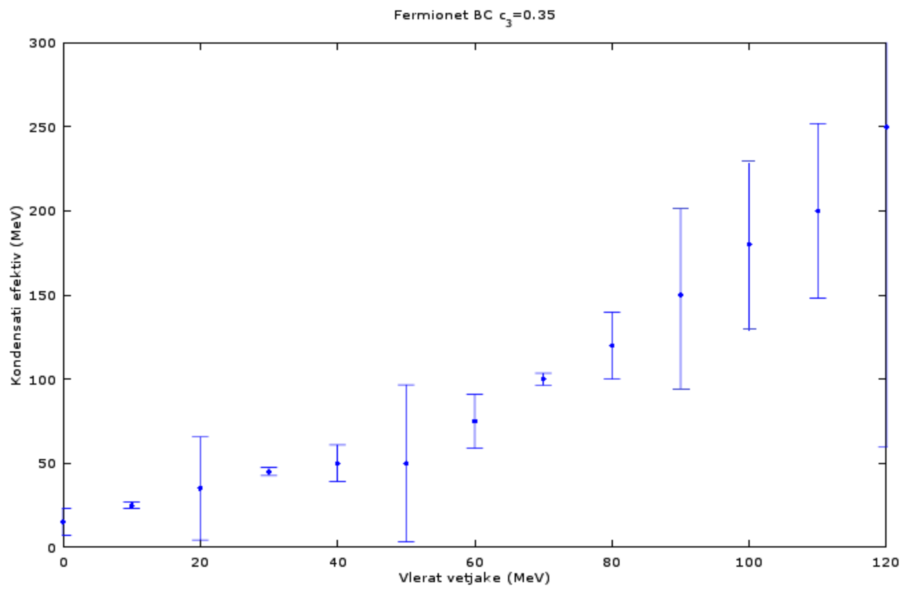
- Calculations of the effective chiral condensate.

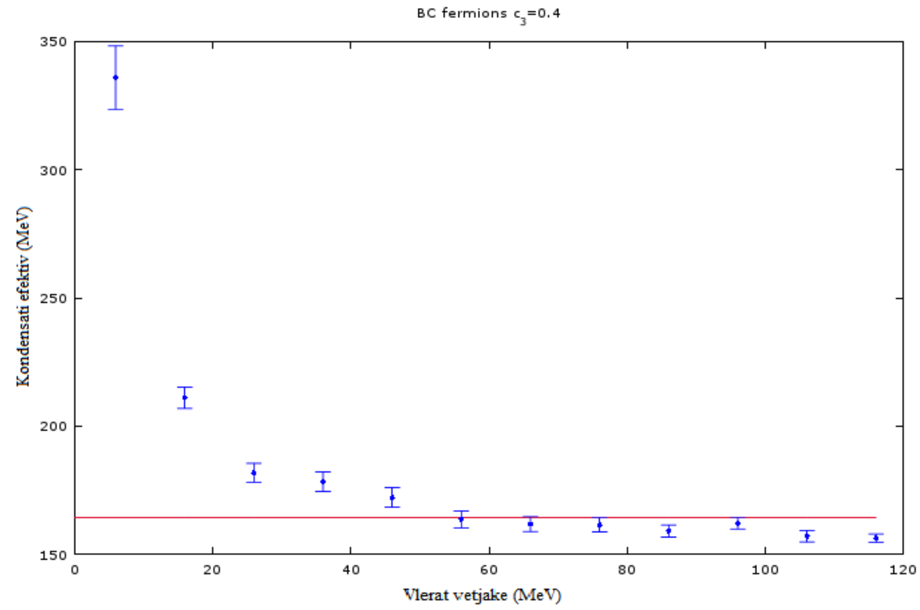
- Lattice 12^4
- Quenched approximation
- Wilson gauge action ($\beta = 6$)
- Boriçi – Creutz action
- Lanczos inverter
- Zero quark mass (BC fermions are chiral fermions)
- Seven different counterterms c3 (0.1, 0.2 0.35, 0.4, 0.6,0.8,1)

Results

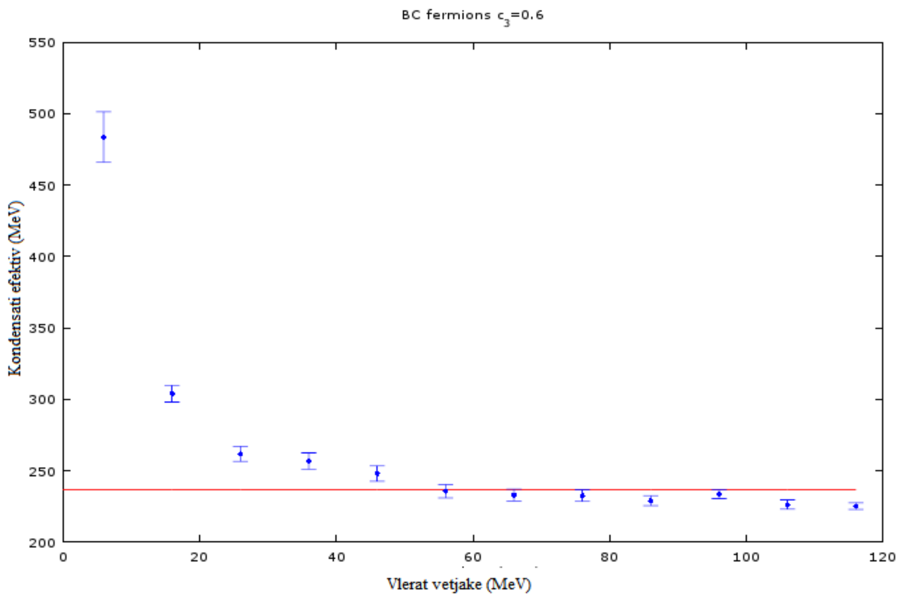

BC fermions c=0.1
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)

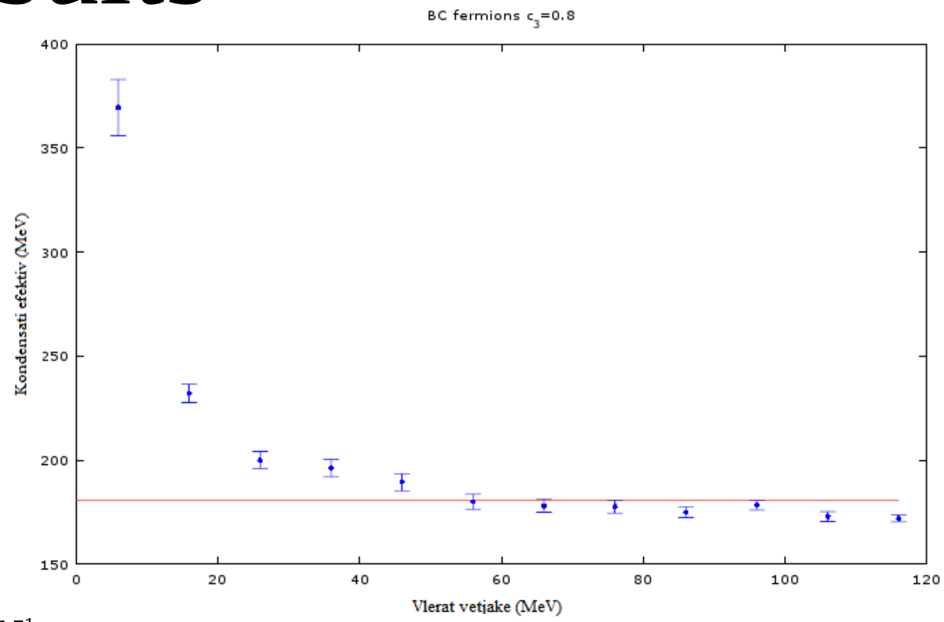

BC fermions c_3=0.2
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)

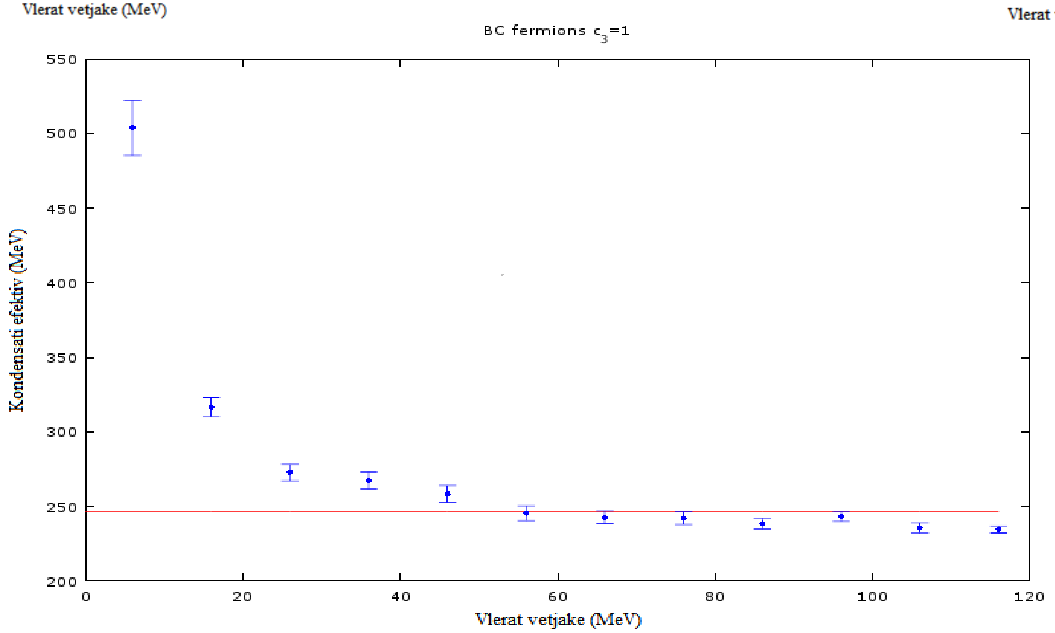

Fermionet BC c_3=0.35
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)


BC fermions c_3=0.4
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)

Results


BC fermions c_3=0.6
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)


BC fermions c_3=0.8
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)


BC fermions c_3=1
Kondensati efektiv (MeV)
Vlerat vetjake (MeV)

Conclusions

- Minimally doubled fermions Boriçi – Creutz present an alternative for the Lattice Quantum Chromodynamics simulations. They preserve an exact chiral symmetry for a degenerate doublet of quarks (chiral symmetry protects mass renormalization). This kind of fermions remains at the same time also strictly local, which means that are fast for simulations.
- BC fermions have a preferred direction in euclidean spacetime → breaking of the hypercubic symmetry
- The calculations of charged pions masses in two different directions show clearly the broken hypercubic symmetry.
- Chiral symmetry and spontanous chiral symmetry breaking is very important in QCD.
- The chiral condensate can be used as an order parameter for BC fermions, and help us to find the proper counterterms that restore partially the broken hypercubic symmetry.
- Using Lanczos quadrature and Banks – Casher relation we can explore the chiral symmetry breaking and find the counterterms for which we have spontanous chiral symmetry breaking
- The value of the counterterm, for which is realised the phase transition of QCD, give us a confirmation that we are working in the proper phase (the one with spontanous symmetry breaking).
- This value depends on the lattice and the coupling constant we use on simulations.
- This work aims to the use of this methodology for further detailed studies of minimally doubled fermions.

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Thank you!