

A cathode excitation effect and jets induced breakdown

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Discharge physics in MPGDs was already reviewed during Paris Workshop and WG2 meeting in December

The main focus in these presentations which were made by *P. Fonte and myself* was on a Raether limit and a high rate effect to which in some cases the Raether limit can be also applied

Raether limit for MPGDs:

It was recently discovered* that in micropattern detectors: GEMs, MICROMEAS and others breakdowns appear at the following conditions:

$$A_{\max} n_0 = Q_{\max} = 10^6 - 10^7 \text{ electrons,}$$

where n_0 is the number of primary electrons created in the drift region of the detector

(Q_{\max} depends on the detector geometry and the gas composition)

(*see Y. Ivanchenkov et al., NIM A422, 1999, 300 and V. Peskov et al., IEEE Nucl. Sci. 48, 2001, 1070)

Therefore:

With single primary electrons gains up to 10^6-10^7 in principle are possible

With ^{55}Fe ($n_0 \sim 230$ electrons) the maximum achievable gain is $<10^5$

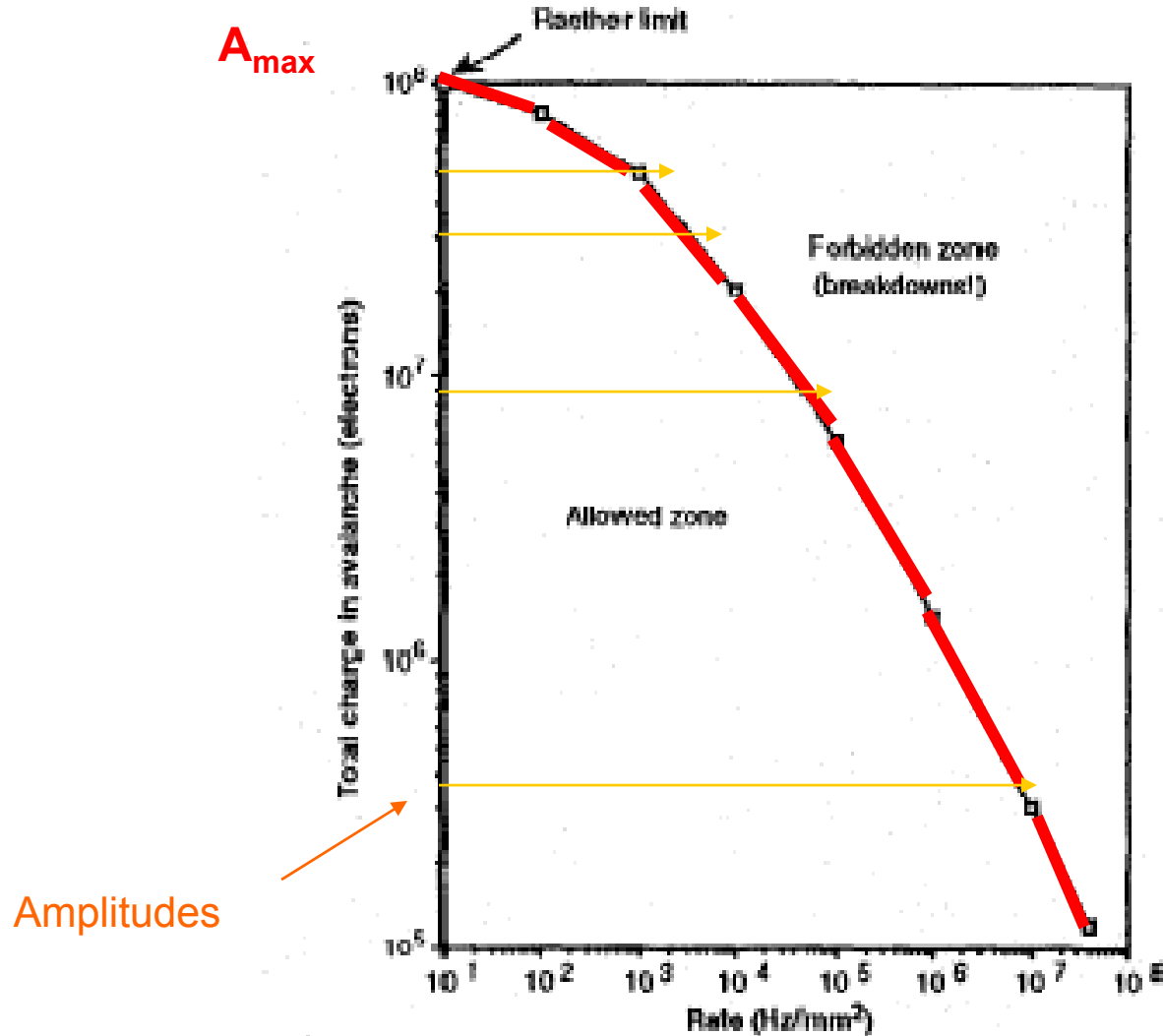
With alphas ($n_0 = 10^5$) the maximum achievable gain <100

Feedbacks related breakdowns can be in the case of MPGDs operating in noble gases or combined with photocathodes)

$A_f \gamma = 1 (A_f \gamma_+ = 1 \text{ or } A_f \gamma_{ph} = 1)$ - "slow" mechanism of discharges

The probabilities γ_+ and γ_{ph} are increasing with the increasing the photocathode QE and it's sensitivity to visible light and with electric field near the cathode

Parallel plate detector (PPAC)

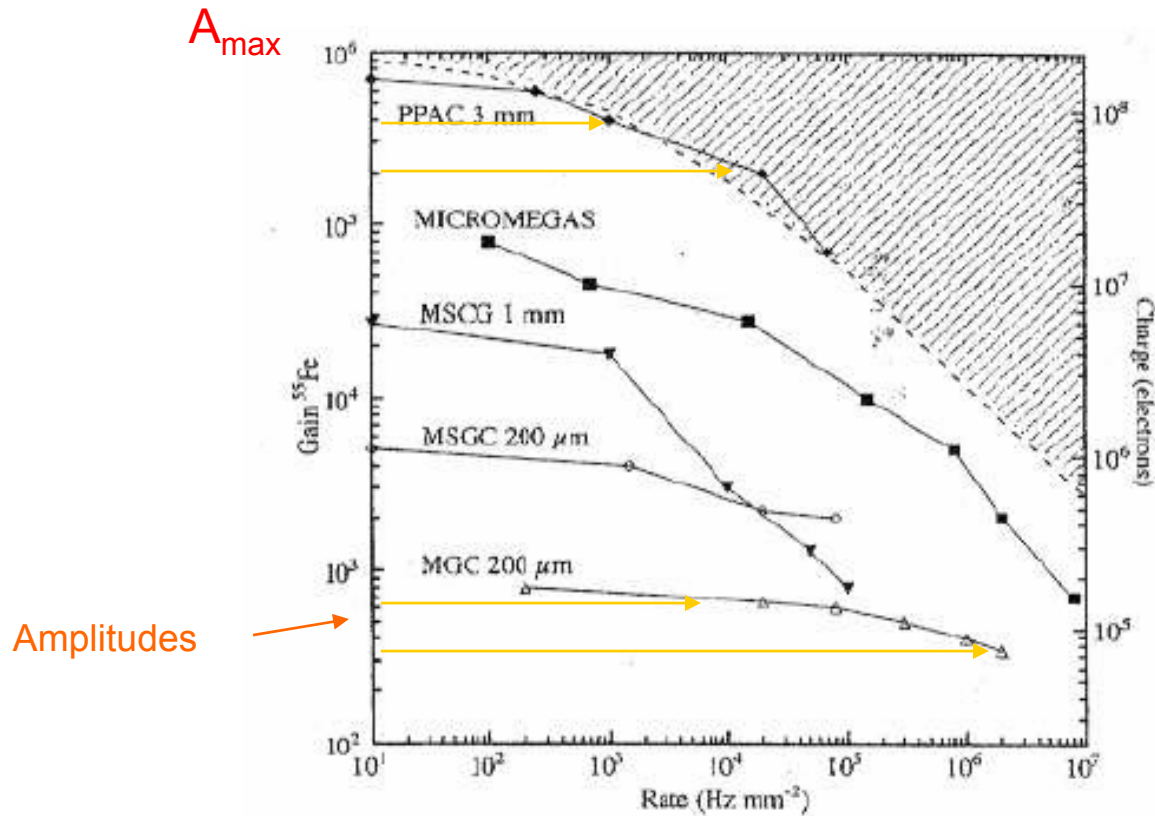


*P. Fonte et al IEEE Nucl. Sci
46, 1999, 321*

Figure 1: General curve reflecting gain limitation with rate for gaseous detectors.

Signal amplitude does not drop with rate, however there is a rate limit for each amplitude

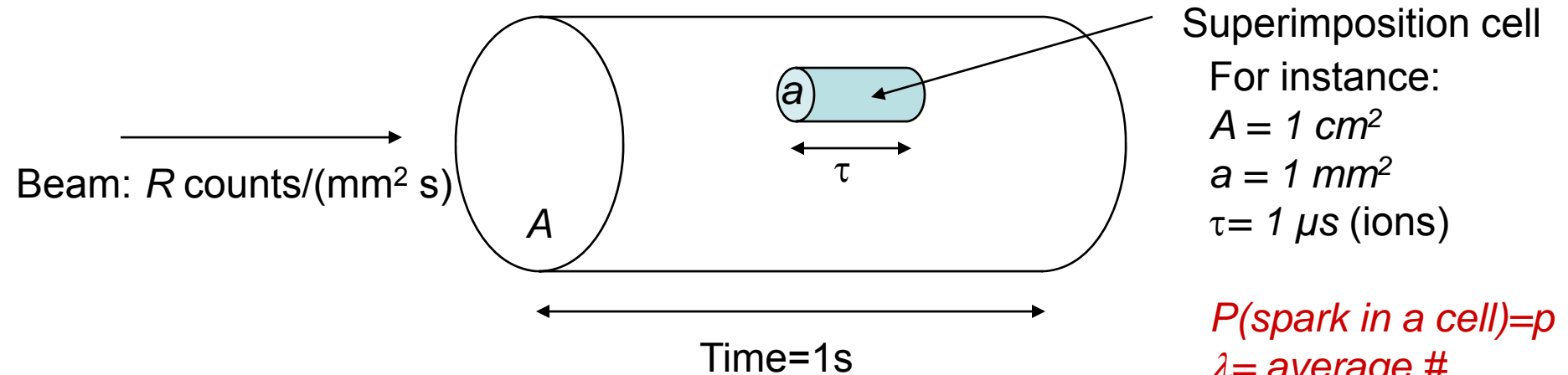
Rate limit of micropattern gaseous detectors



*P. Fonte et al,
NIM A419, 1998, 405*

For each micropattern detector the amplitude remains unchanged with rate, however **the maximum achievable gain** drops with rate

Breakdown statistics via superimposition and Raether limit



There are $N = A/a \times (1\text{s})/\tau$ superimposition cells: $N = 10^8$.

We want to observe a relatively low absolute spark rate $P(\text{spark}) = S \sim 10^{-2} / \text{s}$

$$S = 1 - P(\text{not spark}) = 1 - (1-p)^N \Rightarrow p \approx S/N: p = 10^{-10}.$$

The number of avalanches n in each cell is Poisson-distributed with average $\lambda = Ra\tau$:
 $\lambda = R \times 1 \times 10^{-6}$.

There will be a spark if $nq > Q_R$, q is the average avalanche charge and Q_R the Raether limit.

Then, the required gain reduction owing to superimposition is $1/\tilde{n}$, with \tilde{n} the percentile $1-p$ of the Poisson distribution with average λ .

Rate-induced breakdown? – experimental evidence

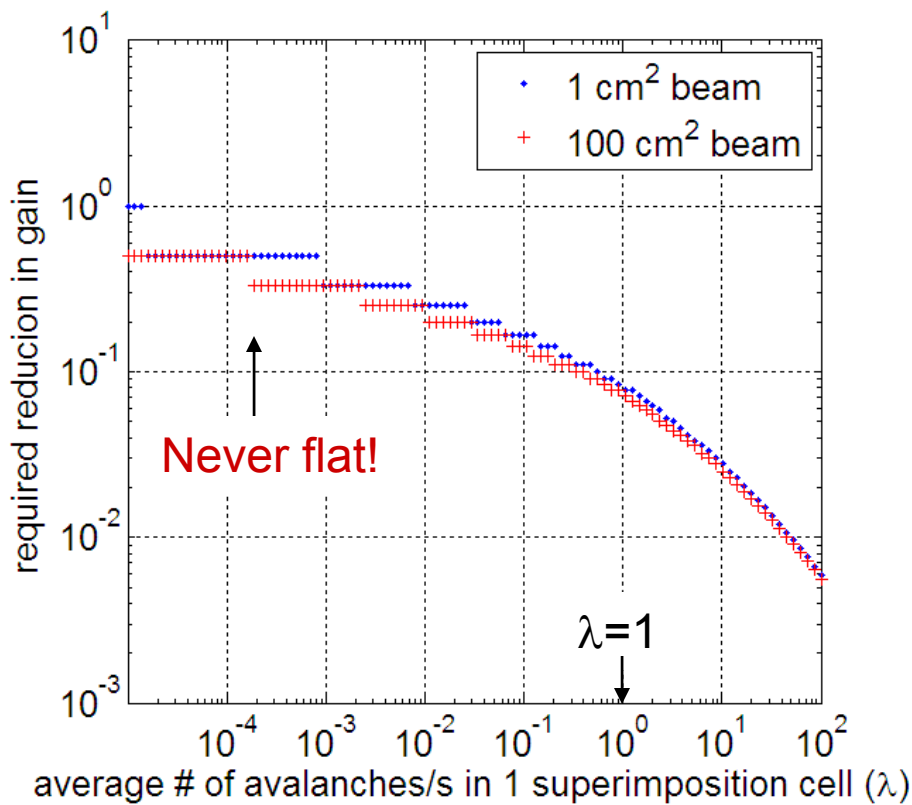
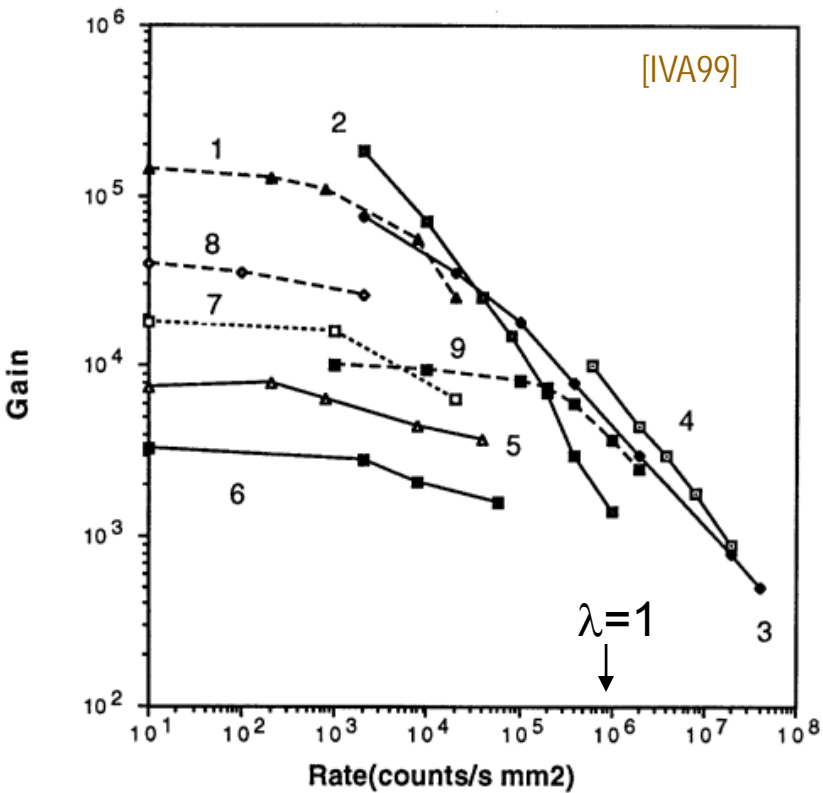


Fig. 1. The maximum achievable gain (curves 1–6), as a function of X-ray flux for various detectors: (1) thick-wire MWPC, (2) PPAC with 3mm gap, (3) PPAC with 0.6mm gap, (4) MICROME GAS (from Ref. [13]), (5) CAT, (6) GEM. (7–9) Space-charge gain limit as a function of rate for other MWPCs: (7) "standard" MWPC, (8) MWPC replotted (from Ref. [14]), (9) thin-gap MWPC (from Ref. [15]).

Mere statistics seem to qualitatively reproduce the data!

BUT...

Puzzle...?

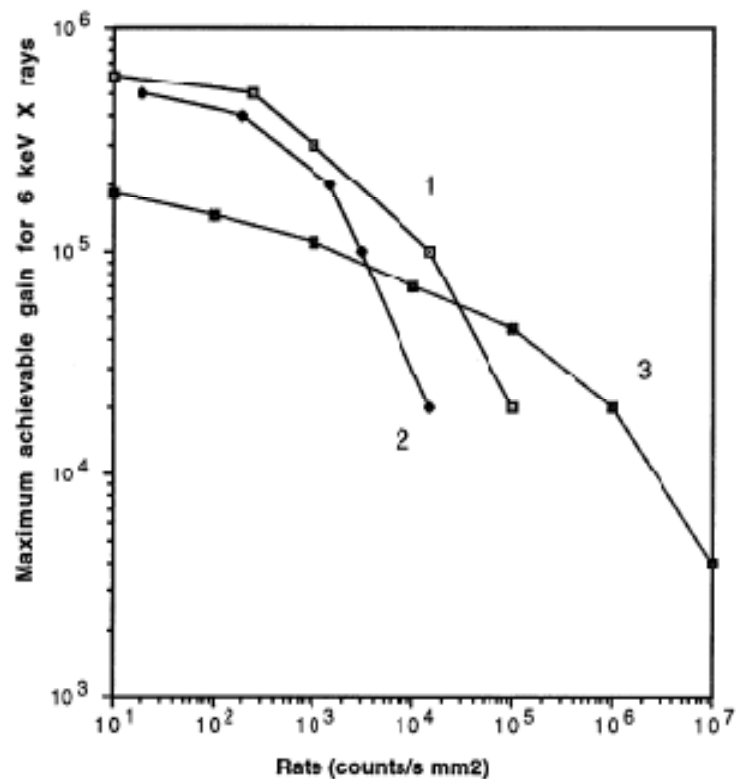
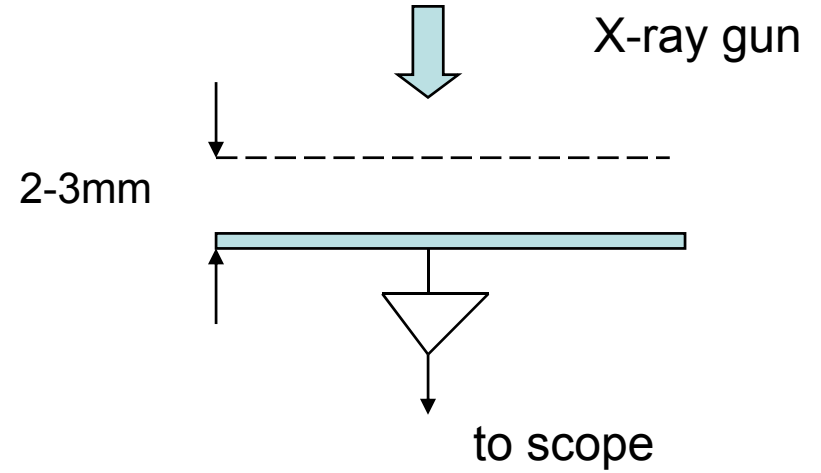
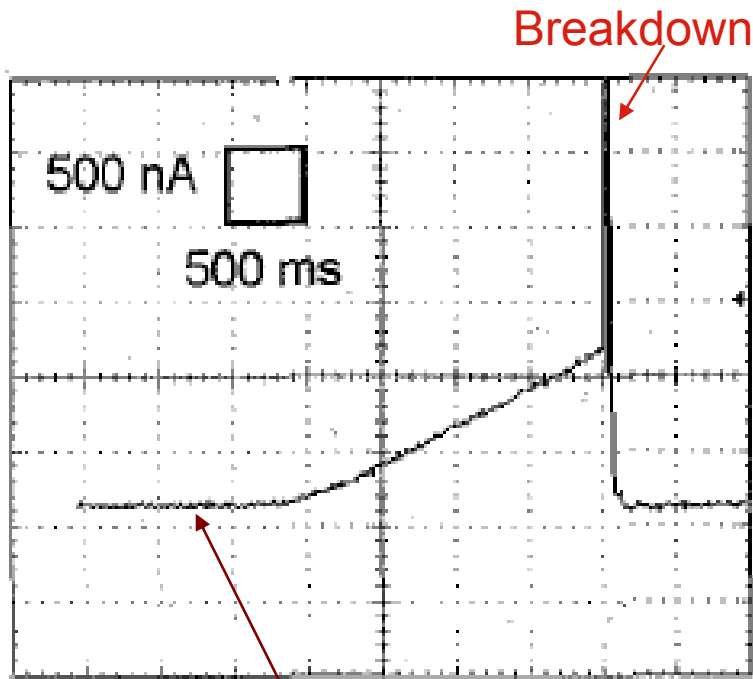


Fig. 3. The maximum achievable gain as function of rate (6 keV X-rays) for the 3 mm-gap PPAC (1,2), and for MICROMEAS (3) (from Ref. [6]). Curve (1) corresponds to a beam diameter of 2 mm and curve (2), to a beam diameter of 20 mm. The gas mixture was Ar + 5% isobutane at 1 atm.

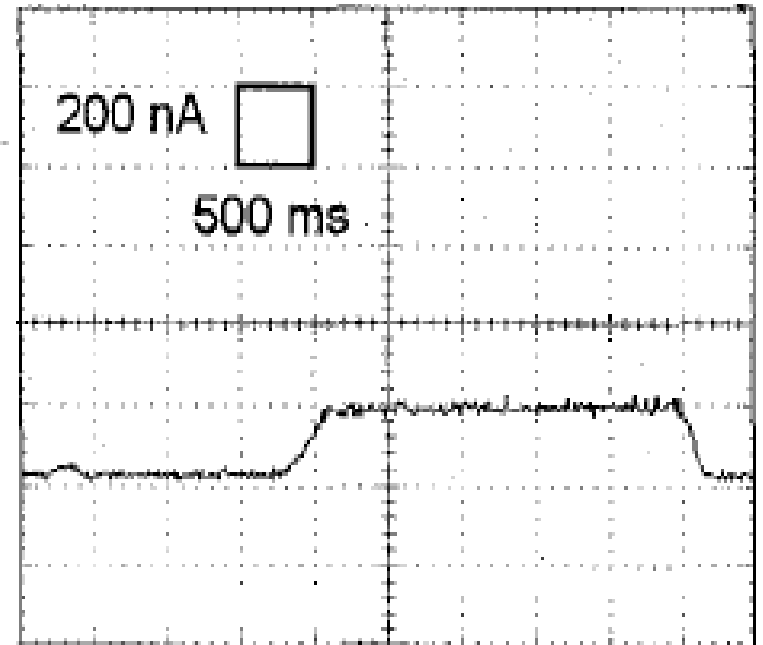
Beside these "classical" mechanisms there were mentioned exotic mechanisms: "preparation", "cathode excitation/memory" effect and electron jets

In this reports I'd like to focus on these exotic breakdown mechanism and demonstrate that they exist in MPGDs

Preparation mechanism

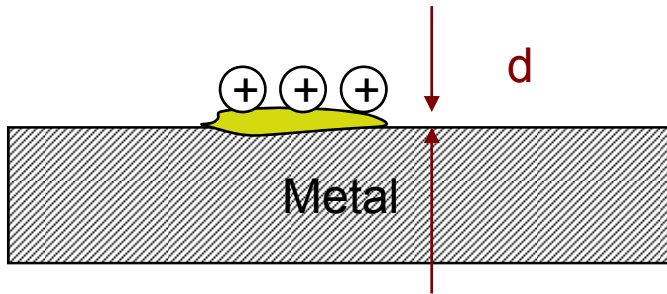


X-ray current



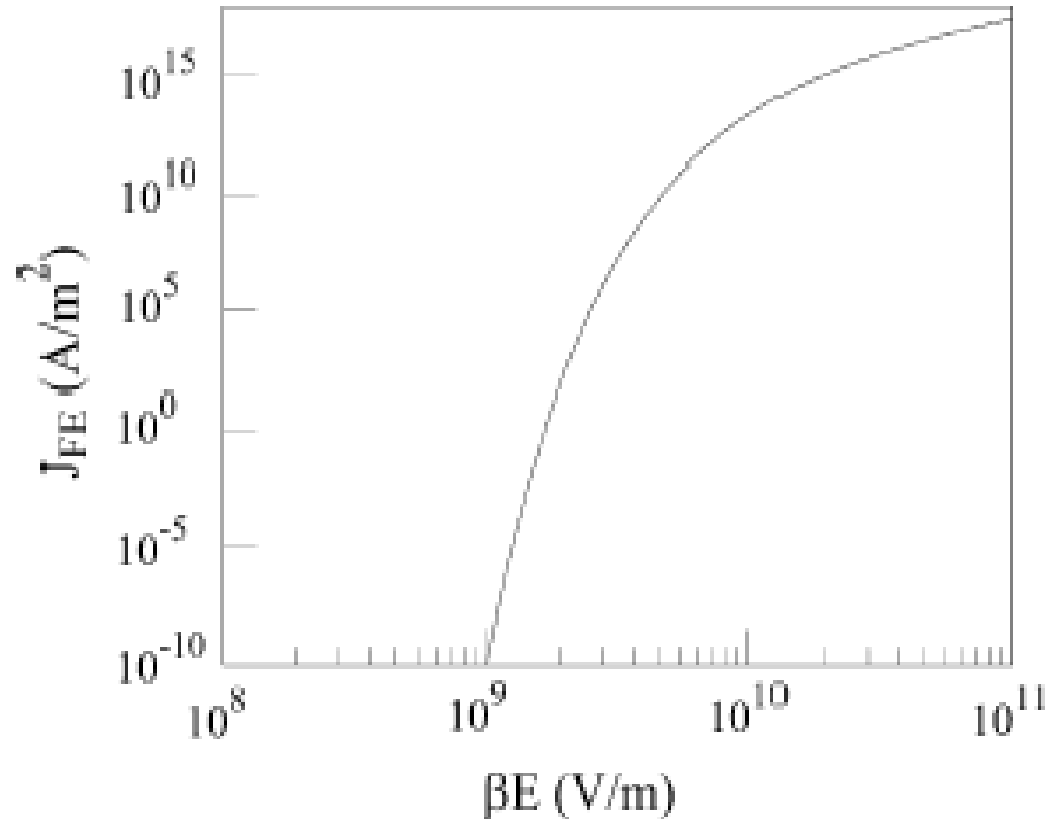
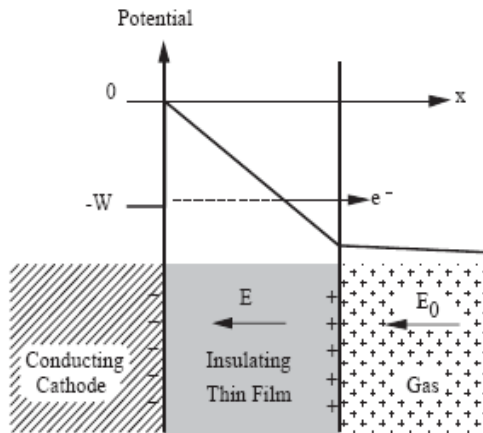
Similar effect is often describes in aging papers (see for example *Aging Workshop, NIM A515, 2003*)

Usual explanation is via a Malter effect..



$$J_{fe} = 5.4 \times 10^{-5} (\beta E)^2 \exp(-5.43 \times 10^{10} / \beta E)$$

$$E = \sigma / d$$



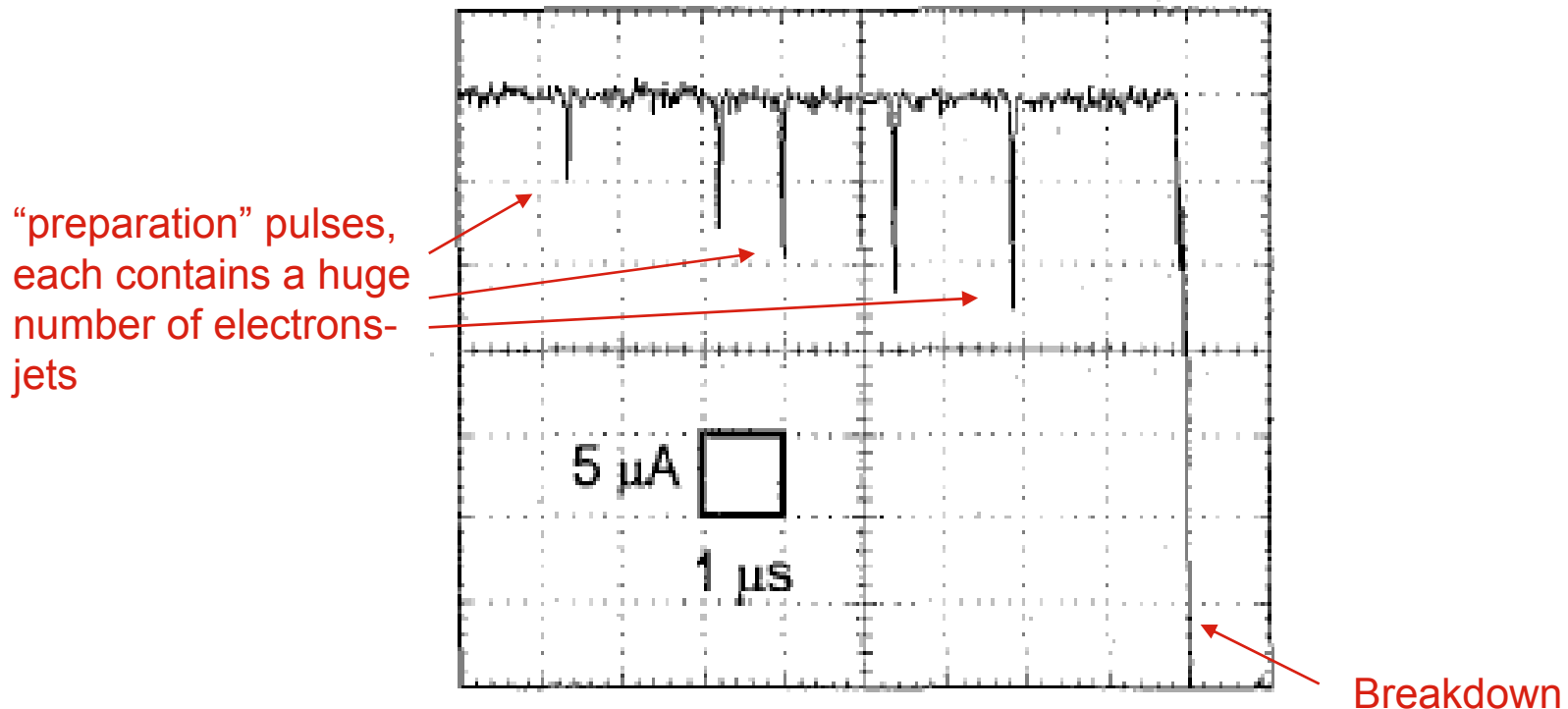
But it is not so simple...

Classical Malter effect predicts single electron emission

(see L. Malter, *Phys. Rev.*, 49, 1936, 478)

However, in most cases a slow current increase is just an integral of high amplitude pulses I. Ivanchenkov et al, *IEEE*, 45, 1998, 258

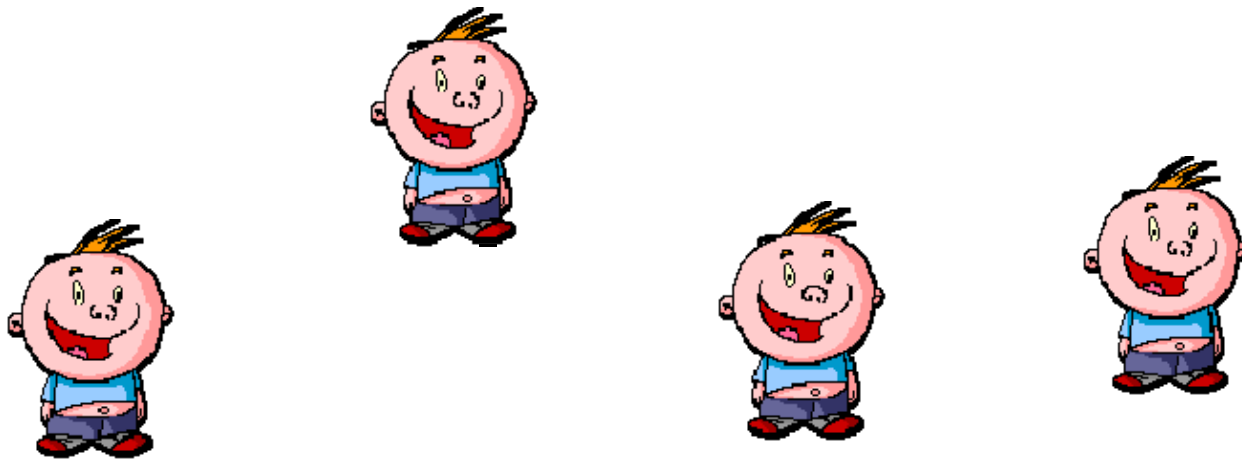
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This strongly contradict to classical Malter effect

More detailed studies reveal that the preparation mechanism may exhibit not only as current pulses but also as a short-term current growth

Cathode excitation



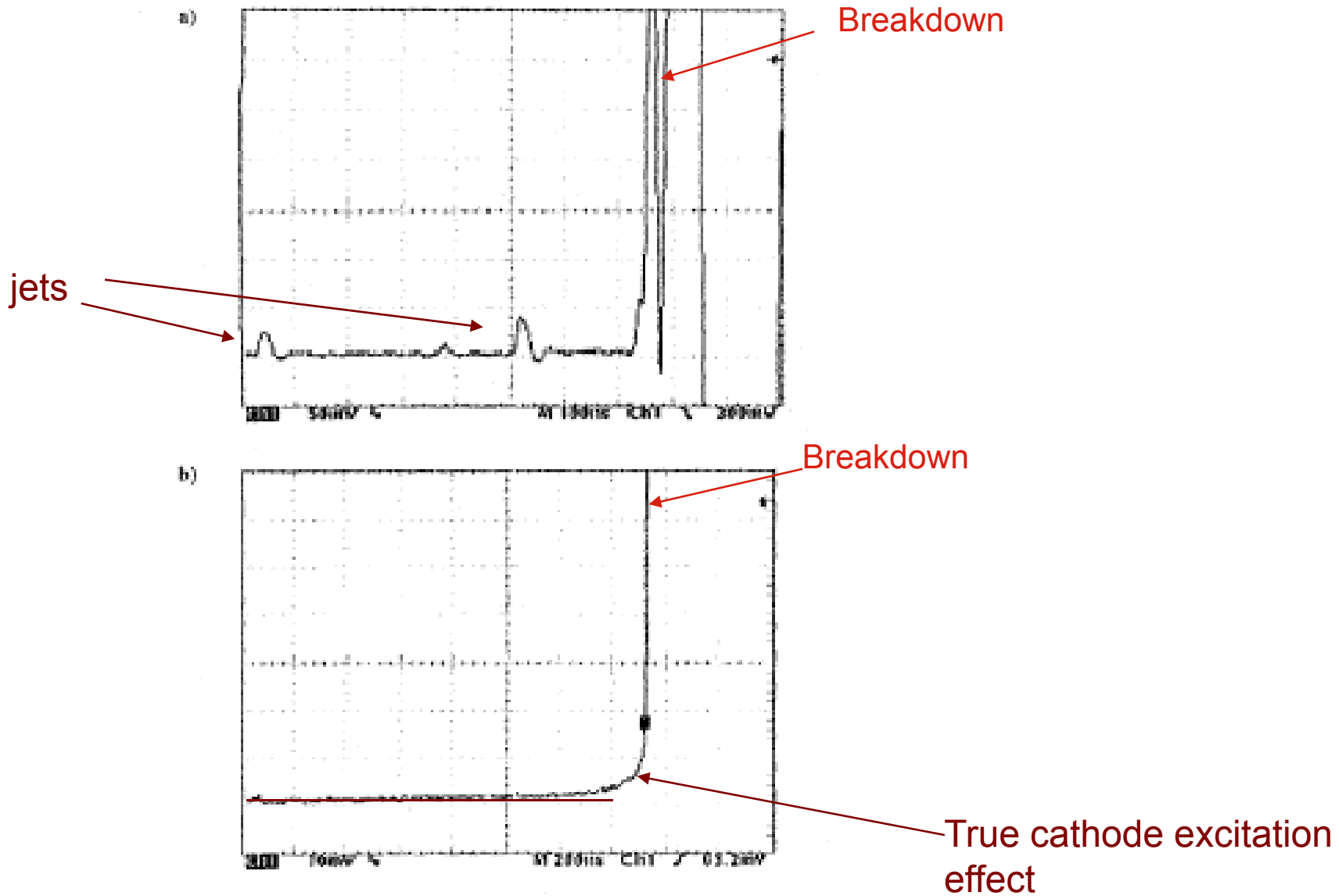


Figure 2a), b): Two typical oscillograms showing a preparation mechanism immediately preceding a high-rate breakdown.

Early studies of the cathode excitation effect

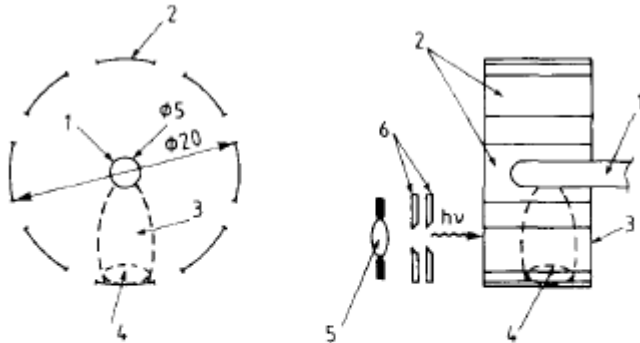


Fig. 8. Sketch of a multisection gas counter for γ_+ measurements. 1 - Cylindrical anode; 2 - multisectional cylindrical cathode; 3 - glow-discharge; 4 - cathode spot of the glow-discharge; 5 - external VUV radiation source; 6 - collimating system for the VUV radiation.

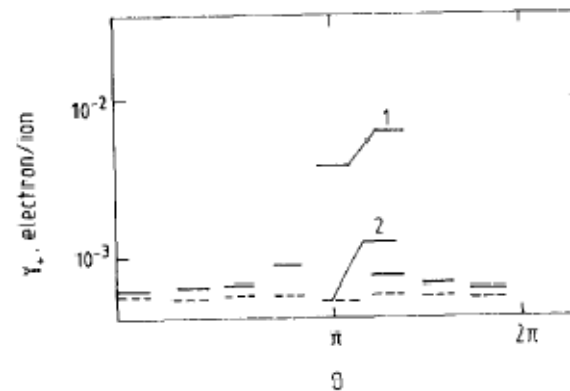
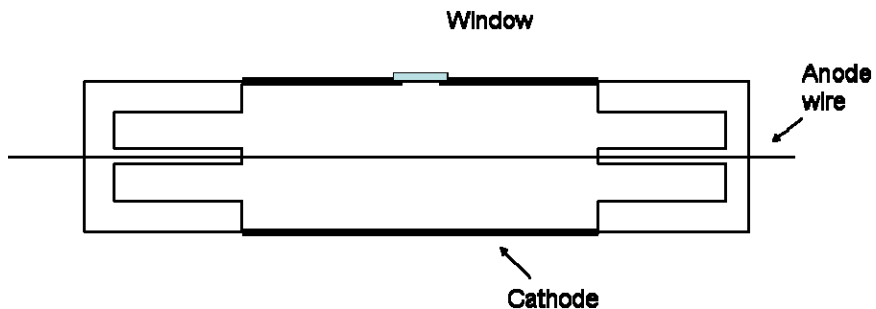


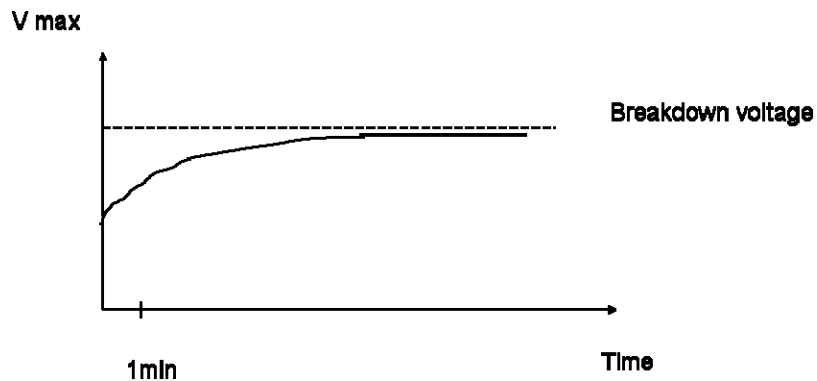
Fig. 9. Values of γ_+ as a function of the angle θ between the upper section (see fig. 8) and investigated one and for different time delays τ after turning off the glow-discharge. 1 - $\tau = 9 \mu\text{s}$; 2 - $\tau = 10 \text{ms}$.

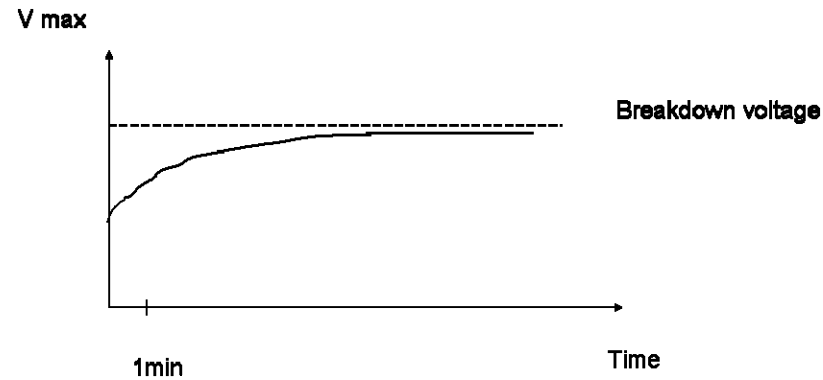
Measurements with single wire counters:



It is well known that
In single wire counters:

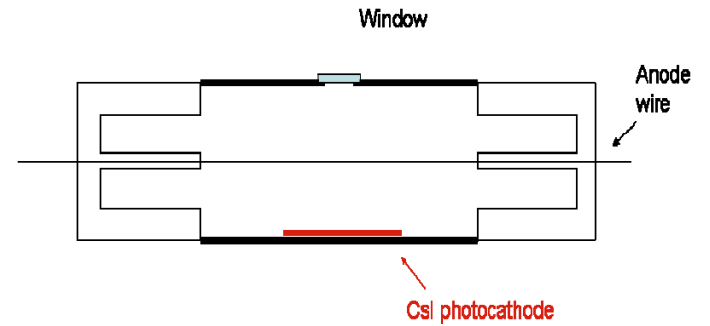
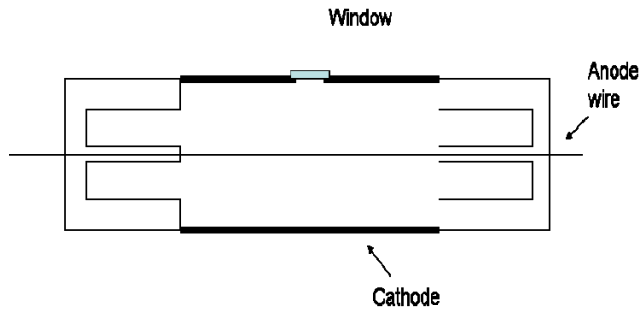
$$A_f \gamma = 1$$



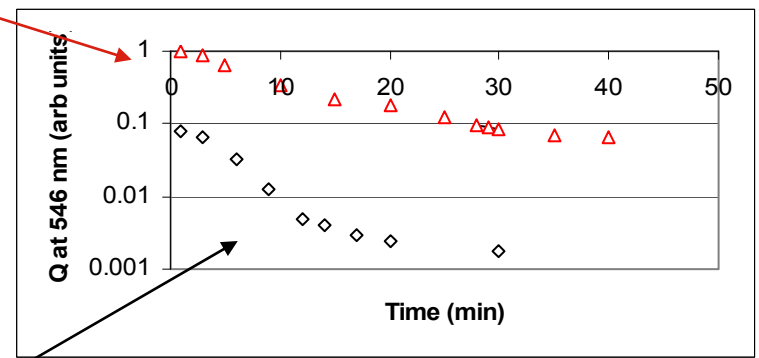
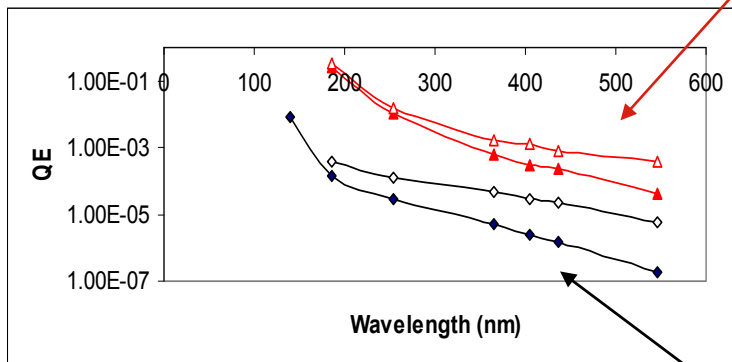


This curve is typical for many gaseous detectors, including MPGDs (check with your experience!)

Changes in QE after intense ion bombardment



CsI



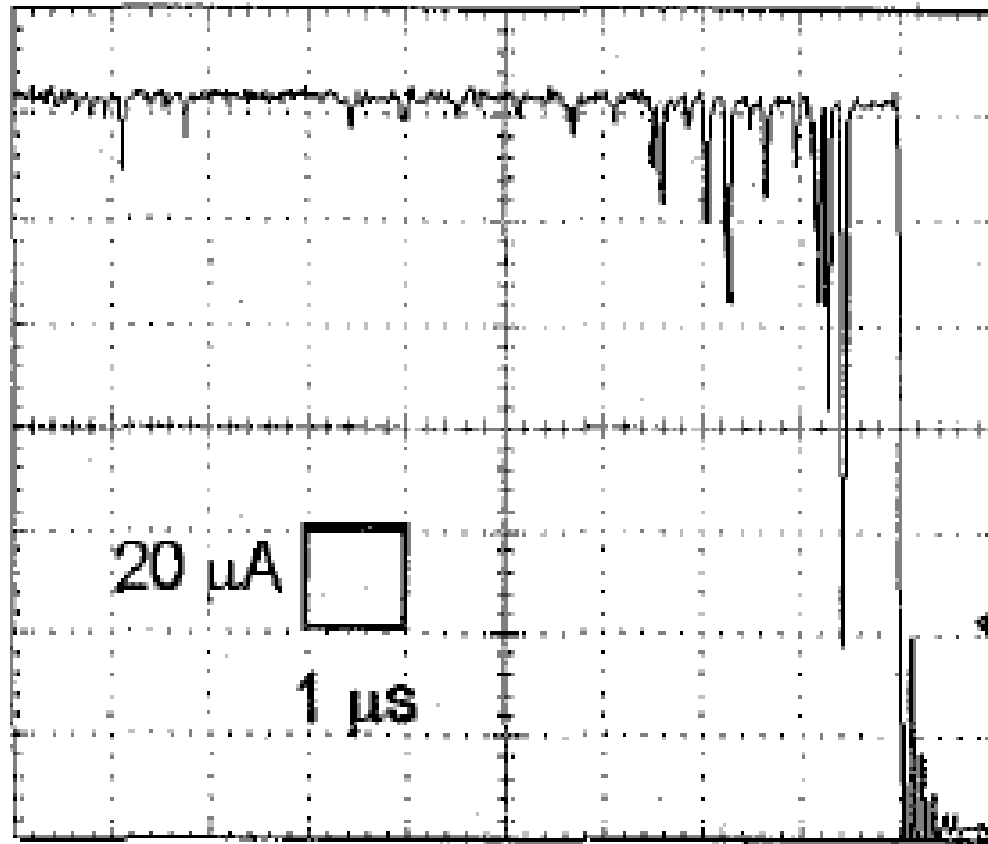
SS

So it was clearly observed that after intense ion bombardment the QE temporally increase as well as γ_{ph} and γ^+

Therefore , the feedback loop $A\gamma=1$ will appear at lower A

Jets

What is the origin of these gigantic pulses?



Explosive field emission

Besides classical field emission calculated by Zommerfeld and others there is another phenomena -explosive field emission

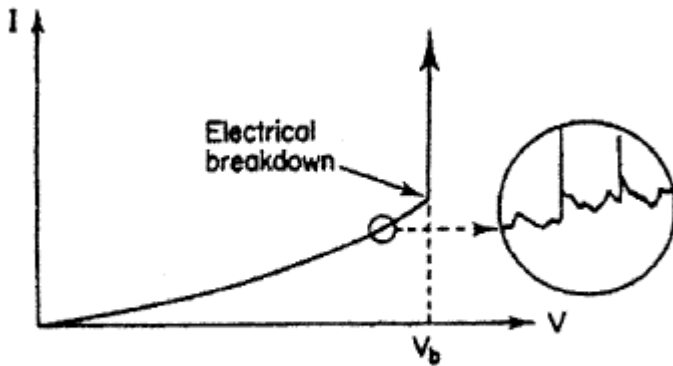
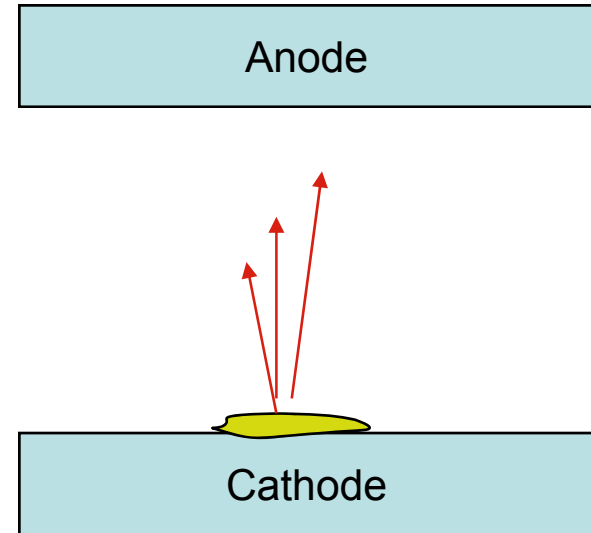
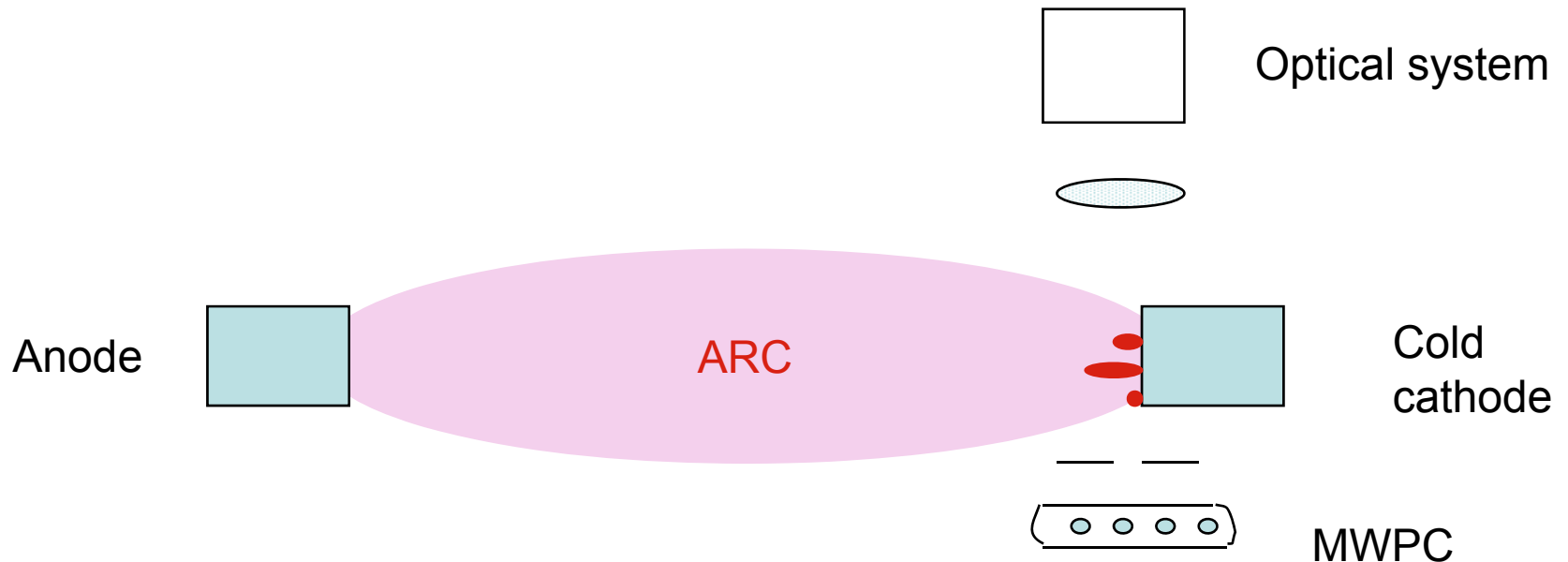


Fig. 14. Current-voltage curve in the case of electrical breakdown in vacuum (from [17]). Enlargement shows pulses due to the explosive field emission.



R. Latham, "High voltage vacuum insulation", new Yoork, 1995

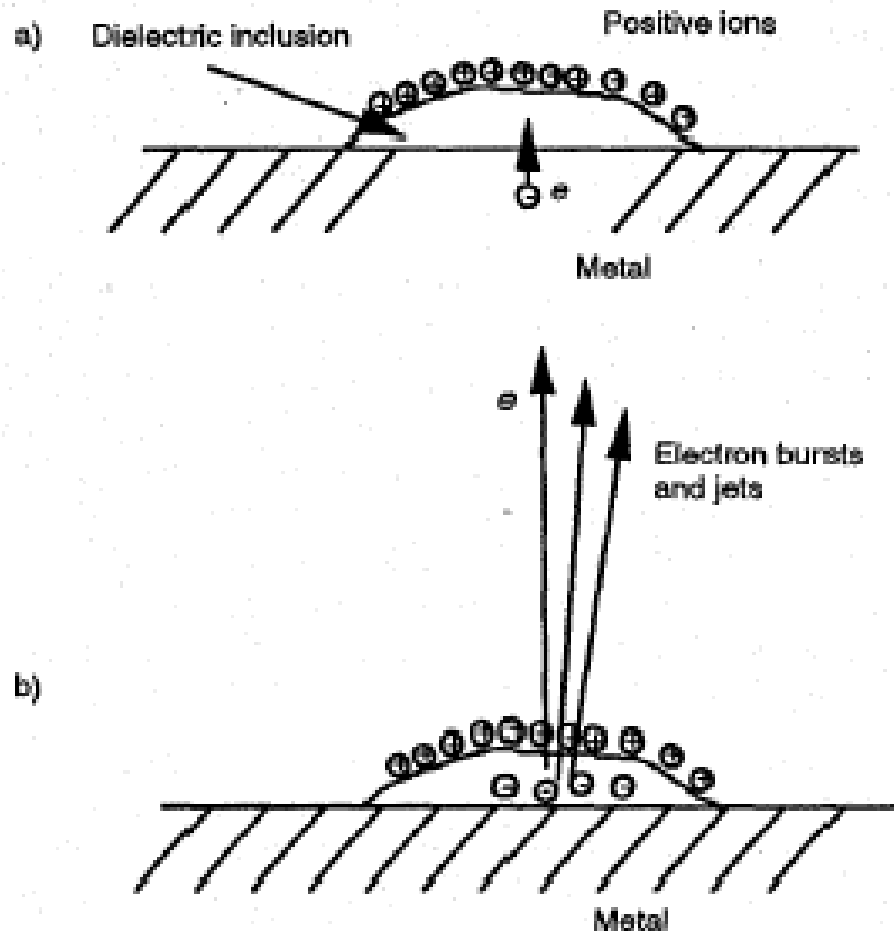
Explosive electron emission was also observed from cold cathodes of some gaseous discharges, for example arcs (*Rachovski phenomena*)



See: G.A. Lubimov, V.I. Rahovski, *Uspek. Phys. Nauk*, 125, 1978, 665,

V. Peskov *Journ, de Physique Coll. C7, suppl#7, 1979,C7-333*

A proposed mode of electrons jets in gaseous detectors:



*P. Fonte et al.,
IEEE Nuc. Sci
46, 1999, 321*

Figure 3: Schematic illustration of a two-step process which leads to emission of jets and bursts from thin dielectric films.

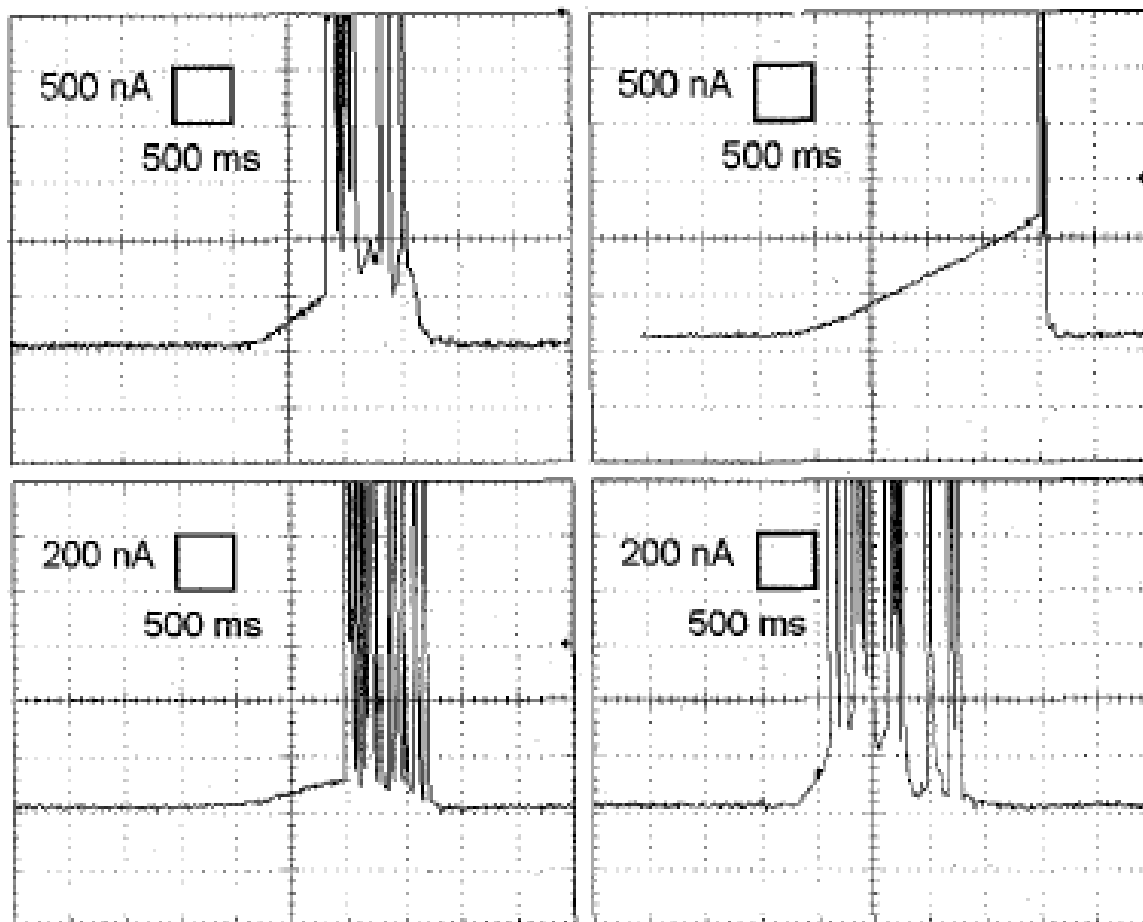


Figure 3 - Gap current as a function of time in the vicinity of a rate-induced spark at low gas gain. A linear increase in current is visible for about one second before breakdown, followed by several large pulses that coincide with a spark.



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Section A

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Model of high-current breakdown from cathode field emission in aged wire chambers[☆]

Adam M. Boyarski*

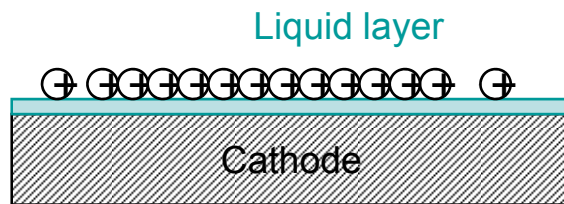
Stanford Linear Accelerator Center, M.S. 95, 2575 Sand Hill Road, Menlo Park, Stanford, CA 94025, USA

Received 18 March 2004; accepted 26 June 2004

Available online 13 August 2004

Some of the results presented in this paper were interpreted via the [jets mechanism](#)

Role of adsorbed layers?



If $E_i < 2\phi$ –no ion recombination

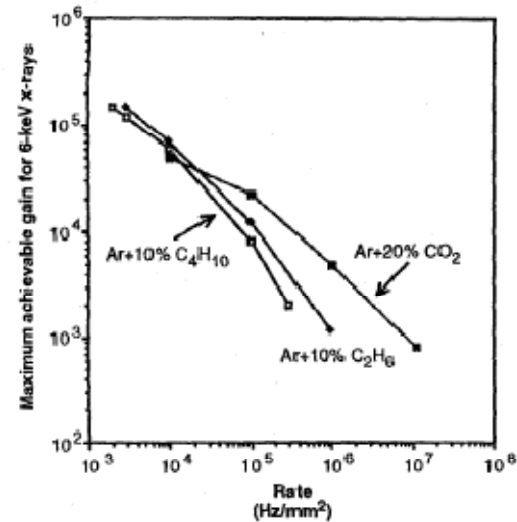


Figure 4: Improvement in the rate characteristics of a PPAC by optimizing the gas mixture.

Observations of jets and cathode
excitation effect in operation of some
gaseous detectors

Examples:

Glass RPC

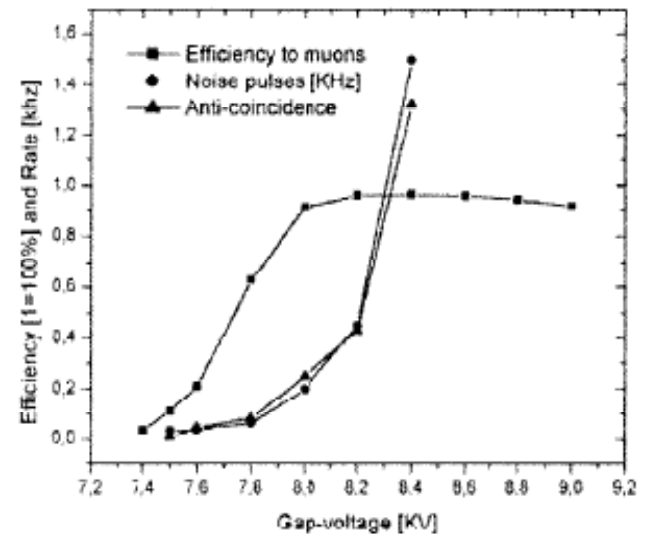
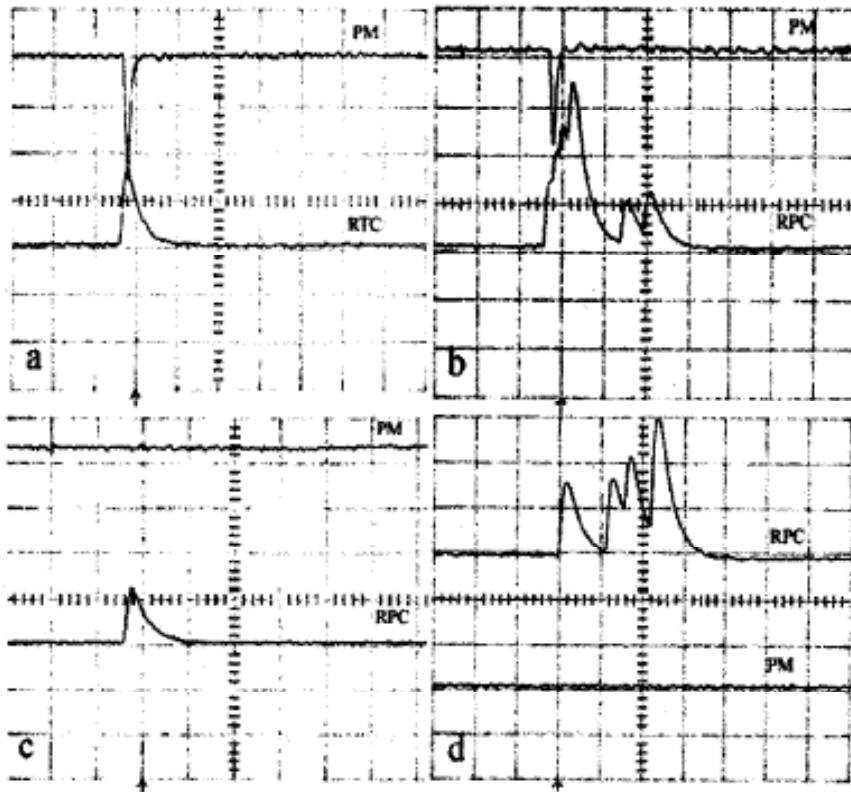


Fig. 5. The efficiency and the rate of noise pulses versus the voltage applied on the RPC [4]. The curve with the triangle symbols correspond to measurements done in anti-coincidence with the signals from the scintillators, respectively. Gas mixture Ar/Isobutane/Freon (R134) in the ration 48/4/48.

High rate Si and GaAs RPCs

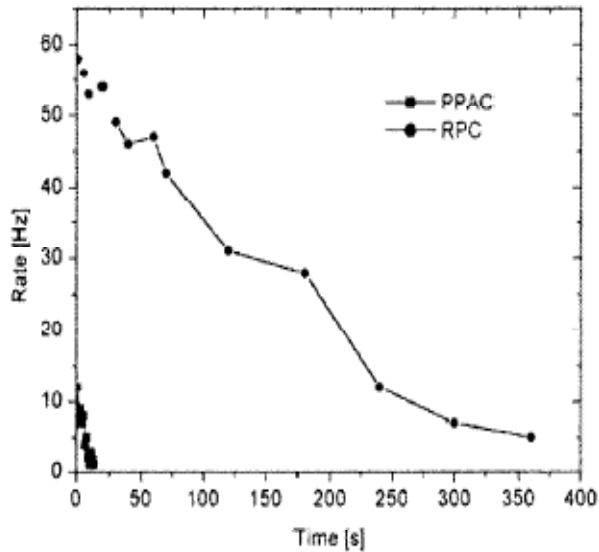


Fig. 11. The rate of the afterpulses for the PPAC (Cu-electrodes) and the RPC (Si). Gas mixture Xe (20%)+Kr (40%)+CO₂ (20%) at 1 atm.

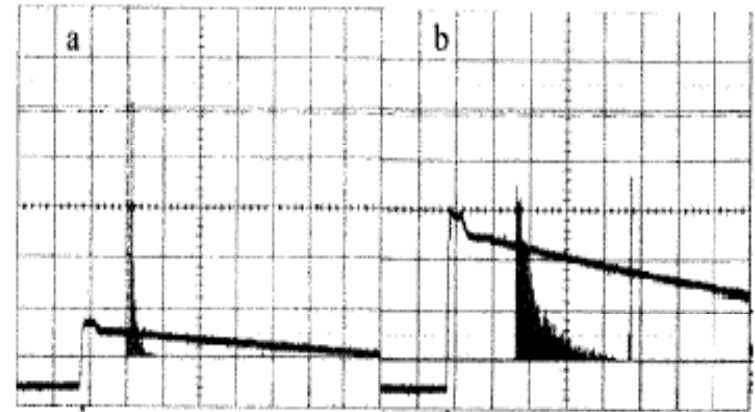


Fig. 6. Pulse-height spectra of signals from RPCs measured in the case of single primary electrons produced from the cathode by (a) UV emission and (b) in the case of noise pulses. The gas mixture Xe (40%)+Kr (40%)+CO₂ (20%) was used (1 atm).

GEM at extreme counting rates

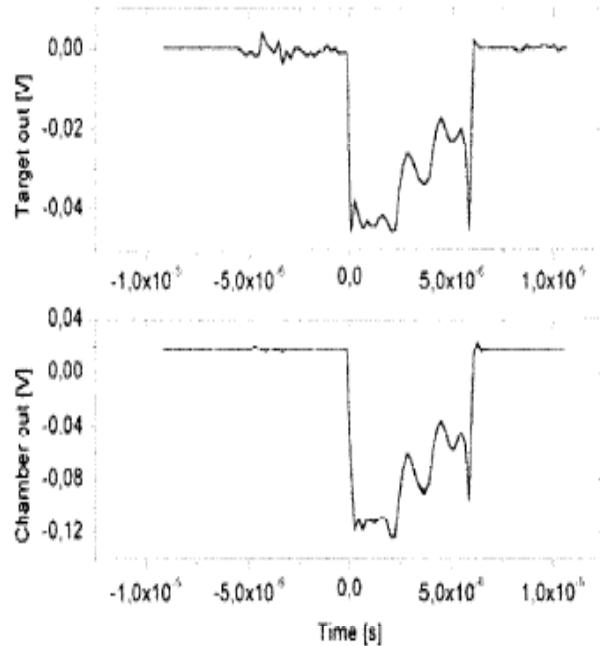


Fig. 12. The current from the GEM (at 350 V) recorded directly on a 50- Ω input of the oscilloscope when the GEM was exposed to a pulsed gamma radiation, producing $\sim 10^7$ counts/mm² on the 2.5 cm \times 2.5 cm GEM area. No other resistors (except the 50- Ω input of the scope) were connected. The upper figure shows the current pulse from a racetrack current monitor. The lower figure shows the corresponding current pulse from the GEM readout. The gas mixture Ar+20%CO₂ was used for the measurement (1 atm).

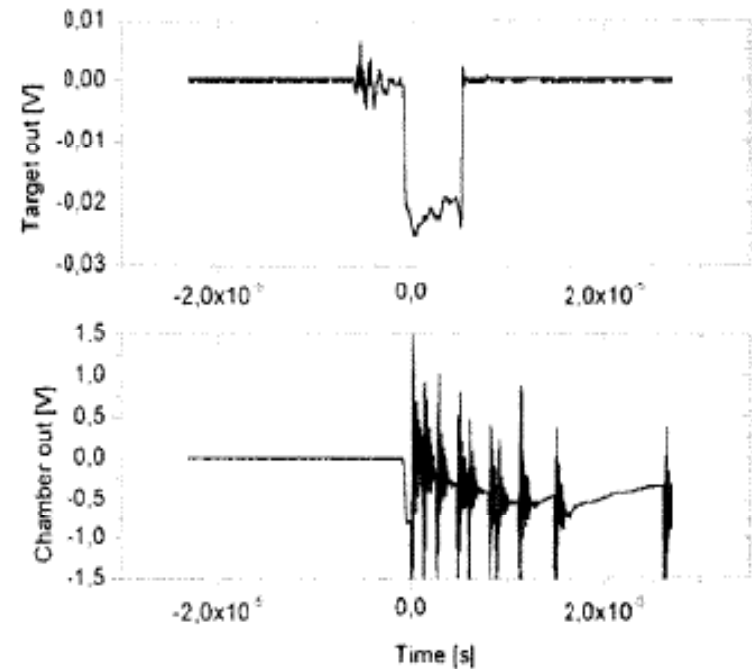
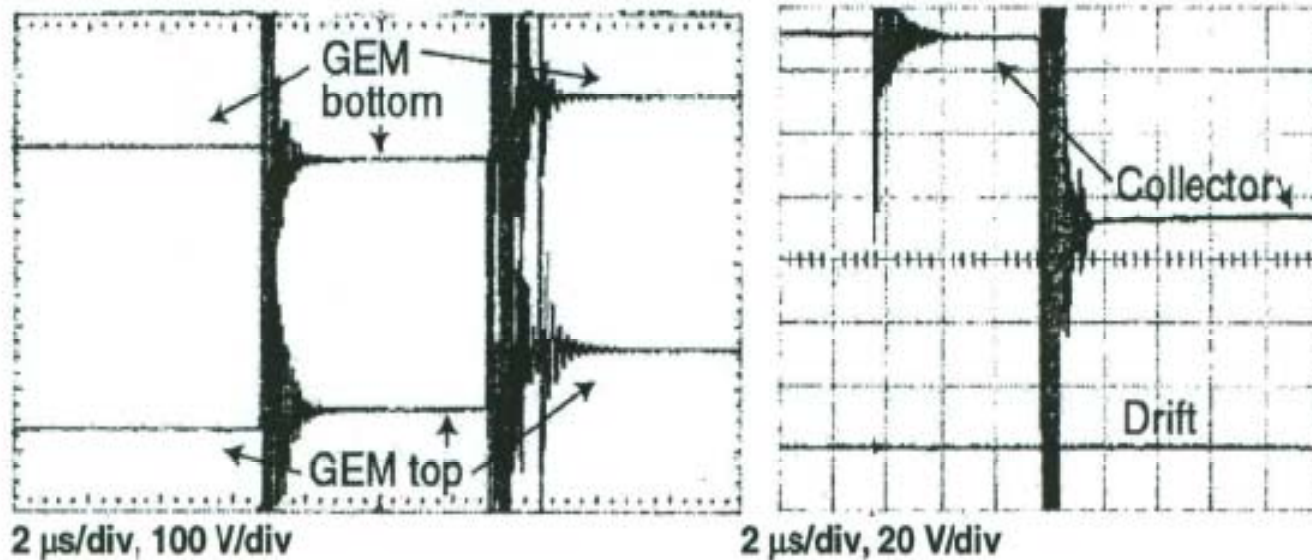


Fig. 13. The same setup as in Fig. 12, but 420 V applied over the GEM electrodes. The upper oscillogram shows the current pulse from the racetrack current monitor, the lower shows the current from the GEM readout. One can clearly see current pulses of large amplitudes, corresponding to a large number of primary electrons $> 10^5$.

Delayed discharge propagation between GEMs



See: V. Peskov. "Discharge propagation between GEMs," this WG-2 meeting

Conclusions

- Besides well established breakdown mechanisms- streamer and feedback related -it was discovered recently another one: “memory/cathode excitation-jets”
- This mechanism mainly show up at high counting rates. For example COMPASS RICH already experiance in memory effect and some tests of RPCs at GIF also revei this effect
- Very often cathode excitation and jets mechanism are mixed
- It will be important to further study these effects , because they can cause problems at future MPGDs applications in LHC experimenst.

