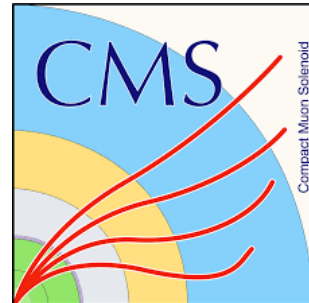


Toward a next generation of simplified DM models

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(representing the work of many)

LHC DM WG Workshop
Sept 19, 2016
CERN



Outline

- Scalar singlet model with mixing (SMM)
- Discovery scenarios
- Other simplified models
- DM mechanisms in SUSY

Primary proposal (and focus of this presentation) is the inclusion of the SMM as recommended benchmark model

The other topics are only briefly reviewed ...

- Propose to revisit these in a later meeting of the LHC DM WG

Proposals derived
from discussions
and studies
involving several
from CMS and the
theory community

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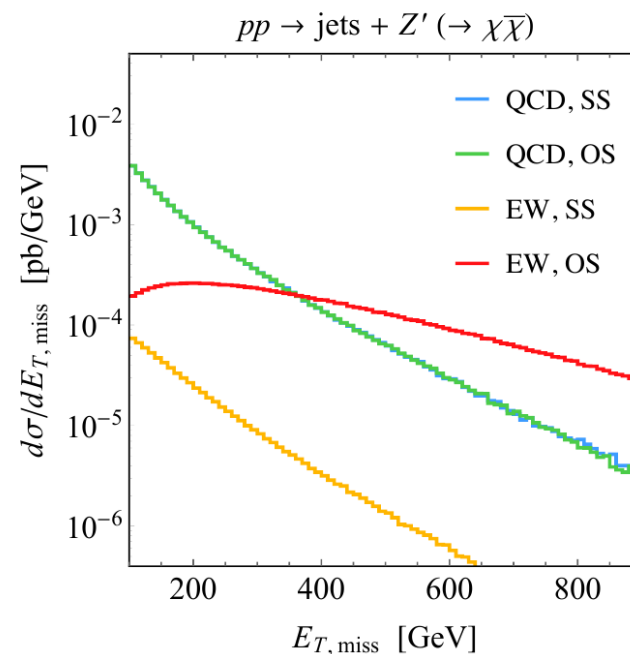
Simplified Models

Motivation

- Capture key features of DM processes @ the LHC
- Really just model “snippets”, but useful :
 - Provides simple framework for probing hallmark MET signatures
 - And a basis for comparison between MET+X channels, with DD/ID, relic density
 - Avoid choosing among numerous UV-completions

Simplification → Caveats

- UV completion not specified
 - Possible bad high s-behaviour: breakdown of perturbative unitarity
 - Possible issues with gauge symmetry: breakdown at lower energy scales
- Can give misleading picture of sensitivity of various search channels
 - Spin-1 Example : mono-W, monojet (+EW)
 - C.f. 1503.07874, 1512.00476, 1603.01267



from 1603.01267

Simplified models should not have pathologies in the channels or regions in which they are applied

- See discussions in 1510.02110, 1604.07975

With current crop of simplified models, one might:

- Restrict couplings
- Add interactions / fields
- Reduce coupling strengths

These treatments developed for mono-W in 1603.01267

For coherent comparison/combination of DM searches, must propagate to all relevant channels

Another approach: use gauge invariant & unitary models

- Avoid pathologies by construction
- Ensure consistency by design

Spin-0 Simplified Models

Some “features” of the current scalar DMF models

- **ttS coupling not gauge invariant**
- **gg → SZ divergent @ high-s**
 - Although not of particular concern at the LHC ...
- **The ttSZ interaction also ?**
 - Probably not a concern either ...
- **Lack some interactions** that are gauge invariant

$$\mathcal{L} \supset -g_t^S \frac{m_t}{v} S \bar{t}t = -g_t^S \frac{m_t}{v} S (\bar{t}_L t_R + \bar{t}_R t_L)$$

$$\sim \frac{\alpha_s}{4\pi} y_t^2 g_t^S \ln^2 \left(\frac{s}{m_t^2} \right)$$

$$S|H|^2, S^2|H|^2, S^3, S^4$$

From U. Haisch

Scalar model with mixing (SMM)

- Example of a simple gauge invariant, unitary model
- **A good basis for an a priori consistent treatment of spin-0 channels**
 - VBF, W/Z associated, mono-H, mono-X ... all with single model!
- **Distinct phenomenology vs current scalar simplified models**

The SMM

Simplest extension of the SM that includes DM

- Fermion DM
- Scalar mediator portal

$$\mathcal{L} \supset -y_{\text{DM}} s \bar{\chi} \chi - \mu s |H|^2$$

Usual mediator-DM coupling, as in DMF models *New portal interaction*

Portal interaction induces scalar/Higgs boson mixing

- Mass eigenstates: h_1 & h_2
- Take h_1 to be observed 125 GeV state

$$\begin{pmatrix} h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} h \\ s \end{pmatrix}$$

h_1, h_2 couplings to DM & SM fermions, W, Z through mixing

$$\mathcal{L} \supset -y_{\text{DM}} (\sin \theta h_1 + \cos \theta h_2) \bar{\chi} \chi$$

$$+ (\cos \theta h_1 - \sin \theta h_2) \left(\frac{2M_W^2}{v} W_\mu^+ W^{-\mu} + \frac{M_Z^2}{v} Z_\mu Z^\mu - \sum_f \frac{m_f}{v} \bar{f} f \right)$$

Model Details

The SMM

- Starting point: singlet model from Khoze *et al.*, 1505.03019
 - ttS loop implemented as an effective operator
 - Missing $h1 \rightarrow h2$, $h2$ and vice versa
 - 1505.03019 also explores more general (2HDM-inspired) mixing scenario
 - We decided to leave this for later
- Migrated to singlet SMM UFO provided by U. Haisch
 - ttS loop fully resolved
 - $h1, h2$ trilinear, some quartic interactions included

“DMF” models

- first pass: DMF ttDM MG model, POWHEG for monojet (as in DMF)
 - Different scale choices in POWHEG/MG \rightarrow difficult to compare
- final pass: use DMSimp for both in MG

SM Higgs

- Also done with DMSimp, provides resolved ttH loop

Kinematic Scenarios

Explore SMM phenomenology in all relevant kinematic regimes, compare with DMF

- A ($m_{h_2} > 2m_\chi > m_{h_1}$) : h2 on shell \rightarrow DM/h1, h1 \rightarrow DM forbidden
- B ($m_{h_1} > 2m_\chi > m_{h_2}$) : h1 on shell \rightarrow DM/h2, h2 \rightarrow DM forbidden
- C ($m_{h_2} > m_{h_1} > 2m_\chi$) : heavy h2, light DM
- D ($m_{h_1} > m_{h_2} > 2m_\chi$) : light h2, lighter DM

Kinematics : generated $p_T(\chi\chi)$ vs mixing angle, coupling

- monojet (1,2 jets) and tt+DM (+ 0,1jet) in SMM & DMF
- Also compare with SM h + 1,2 jet and tth + 0,1 jet

Cross section : vs mixing angle

- tt+DM (+0,1 jet) in SMM & DMF, and SM tth + 0,1 jet
- **NB: Run-I Higgs global fit $\rightarrow \sin\theta < \sim 0.4$**

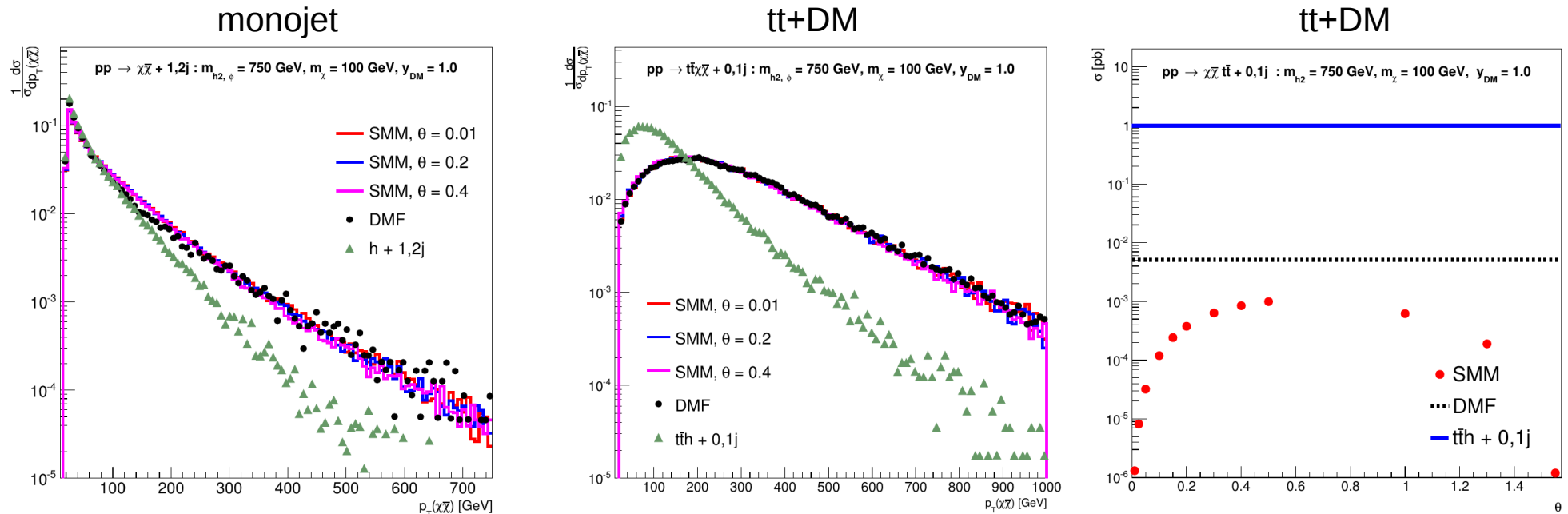
Scenario A : $m_{h_2} > 2m_\chi > m_{h_1}$

Kinematics

- New h_2 scalar mediator can decay into DM
- **SSM kinematics generally map to those of DMF**
 - Disagreement in monojet near $p_T(\chi\chi) \sim 150$ GeV from new diagrams

Cross section

- **Suppressed relative to DMF and SM**



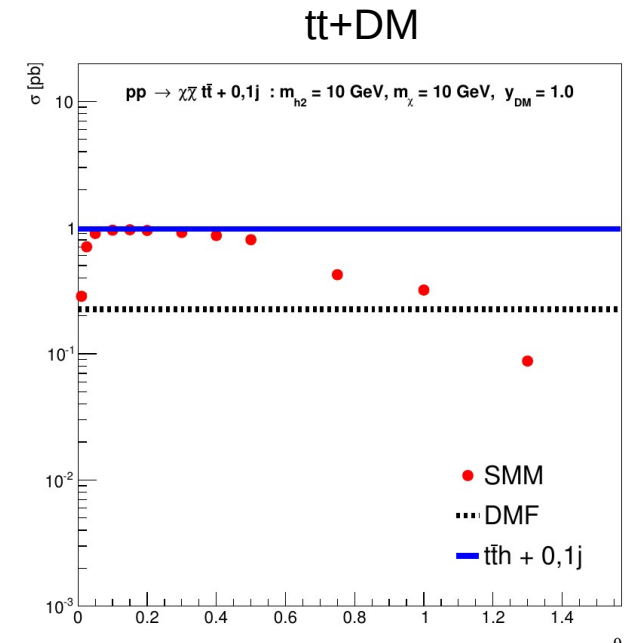
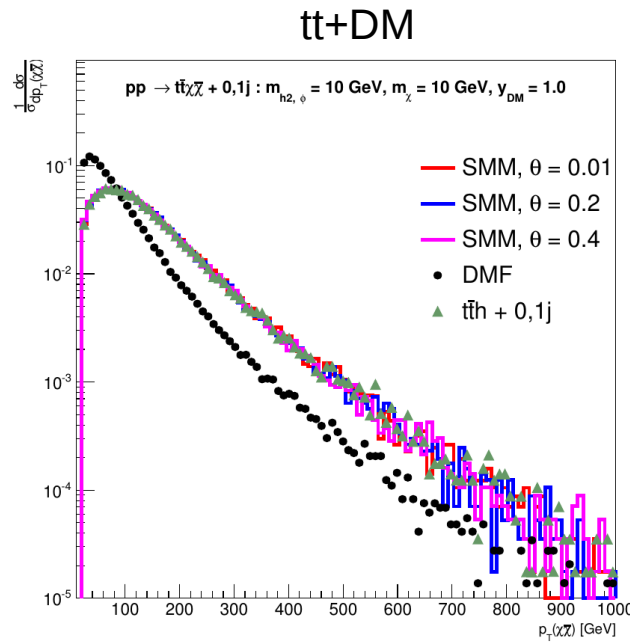
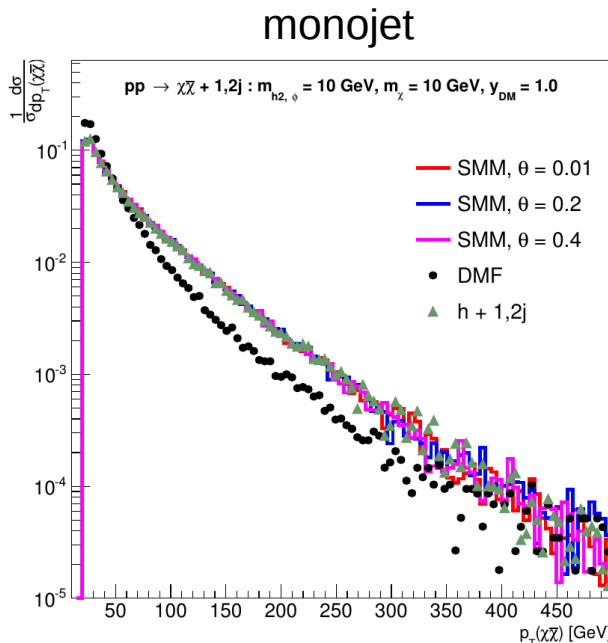
Scenario B : $m_{h_1} > 2m_\chi > m_{h_2}$

Kinematics

- h1 can decay into DM, new h2 mediator can't
- SSM kinematics map directly to those of SM Higgs

Cross section

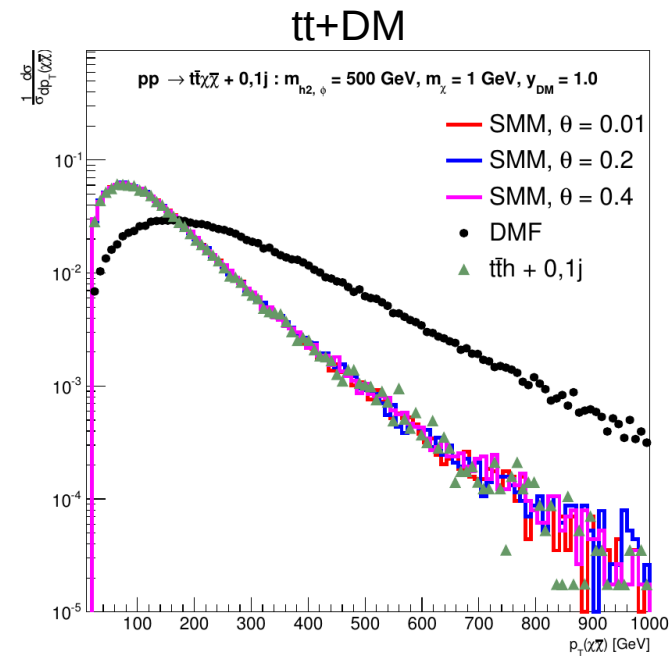
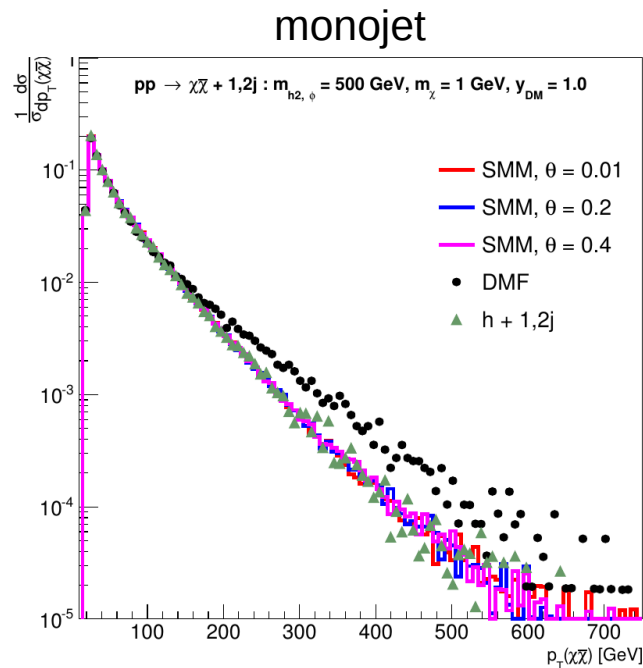
- Approaches SM, higher than in DMF
 - Only off-shell decays in the DMF scenario ...



Scenario C : $m_{h_2} > m_{h_1} > 2m_\chi$

Kinematics

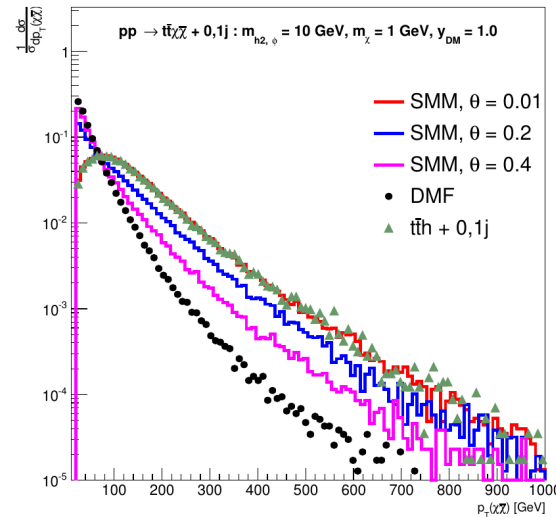
- $h_1 \rightarrow \text{DM}$ resonantly enhanced due to small h_1 width
- SSM kinematics map to those of SM Higgs



Scenario D : $m_{h_1} > m_{h_2} > 2m_\chi$

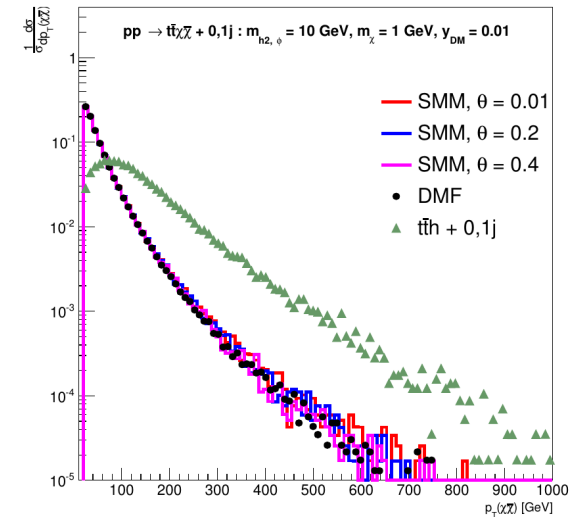
Kinematics

- Impact of mixing clear for nominal ($y_{DM} = 1.0$) couplings
- Both $h_1 \rightarrow DM$ and $h_2 \rightarrow DM$ on-shell
- Both processes important if Γ_{h_1} and Γ_{h_2} comparable
- Kinematics evolve toward DMF as coupling is lowered

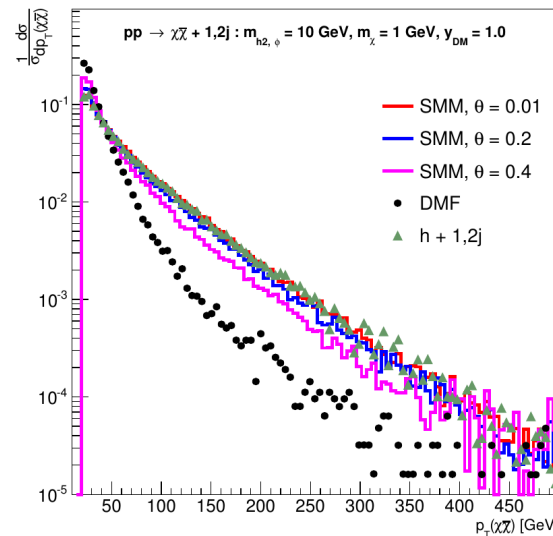


$y_{DM} = 1.0$

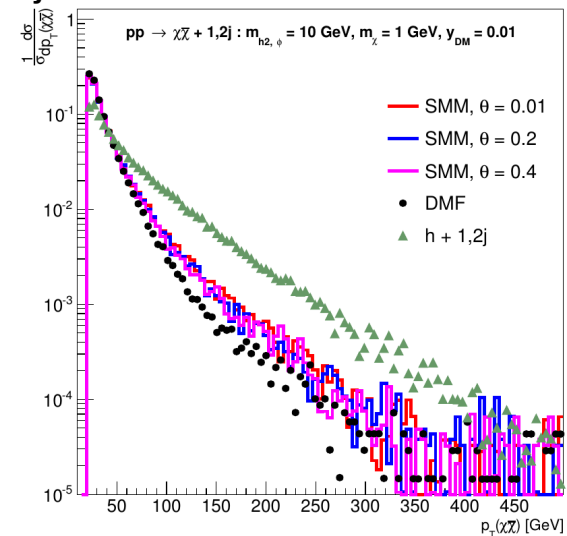
tt+DM



$y_{DM} = 0.01$



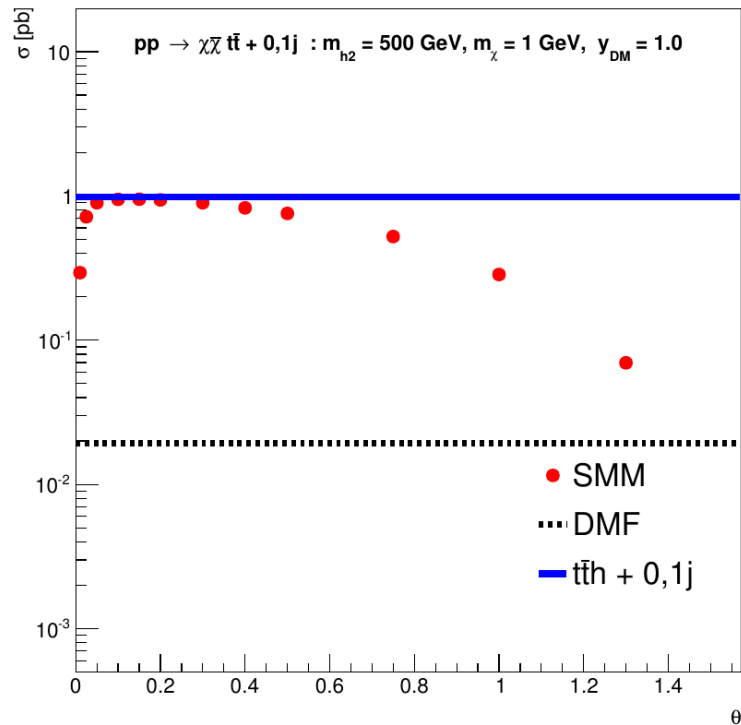
monojet



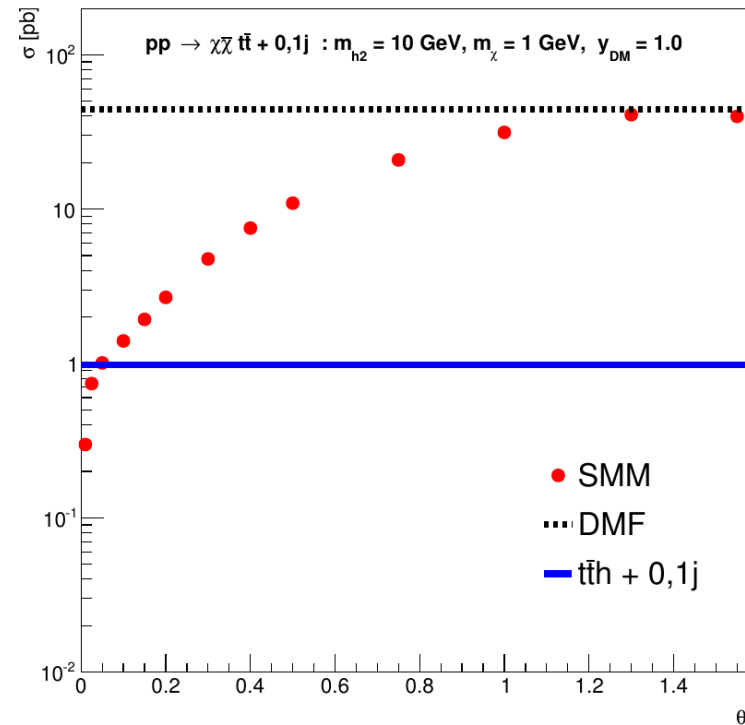
Scenario C/D Cross Sections

Approaches SM Higgs for Scenario C, $> \sigma_{DMF}$

Generally between σ_{SM} and σ_{DMF} for Scenario D



Scenario C

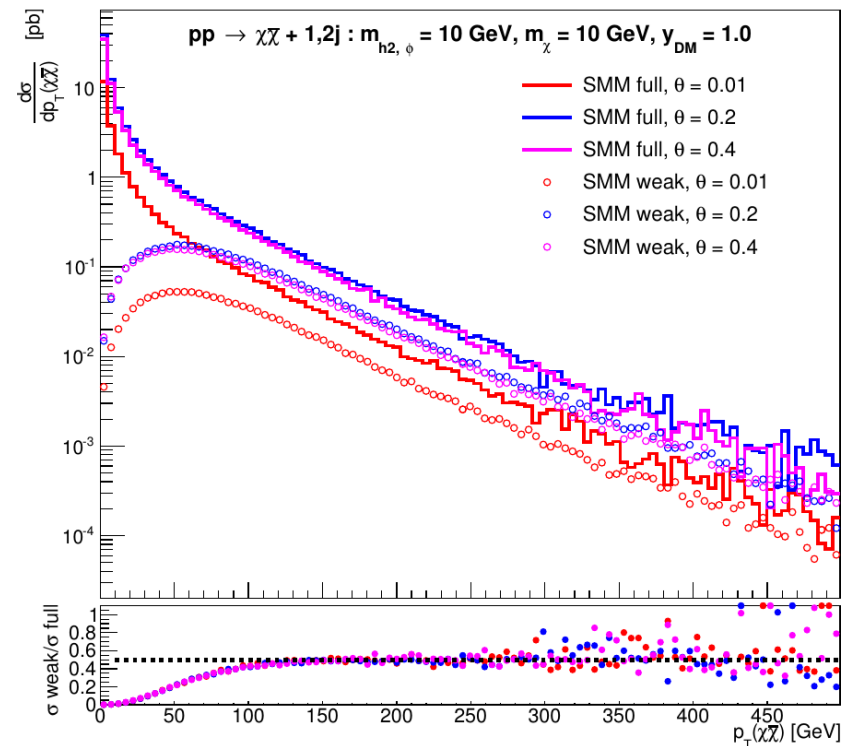
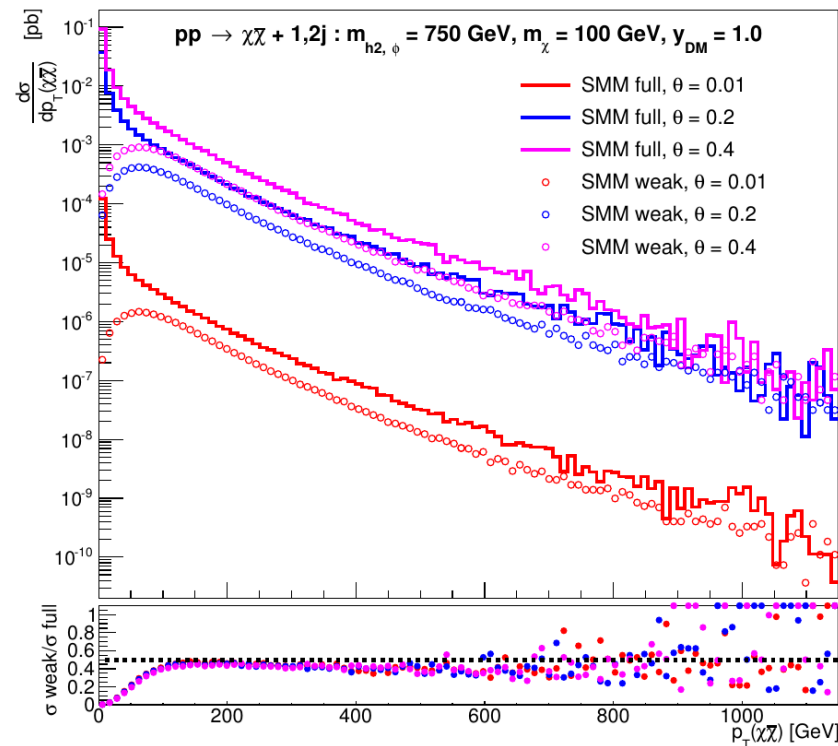


Scenario D

Scenario C/D Differential Cross Sections

Compare full $pp \rightarrow \chi\chi$ monojet cross section vs that of new VB-mediated processes

- Significant contributions from new couplings at high $p_T(\chi\chi)$



Conclusions from SMM

- Higgs mixing is the simplest extension to the current DMF scalar models
- In several regimes, kinematics & cross section differ significantly from those of the DMF models
- SMM displays much richer phenomenology
 - In particular, from new vector boson mediated processes
- The model is gauge invariant and unitary
 - Provides a useful basis for the comparison/combination of scalar-mediated DM searches
 - VBF, W/Z associated, mono-H, mono-X, all with single model!
- ▶ We propose that the SMM be included in the suite of LHC DM WG recommended benchmark models

Conclusions from Other Topics

Discovery scenarios

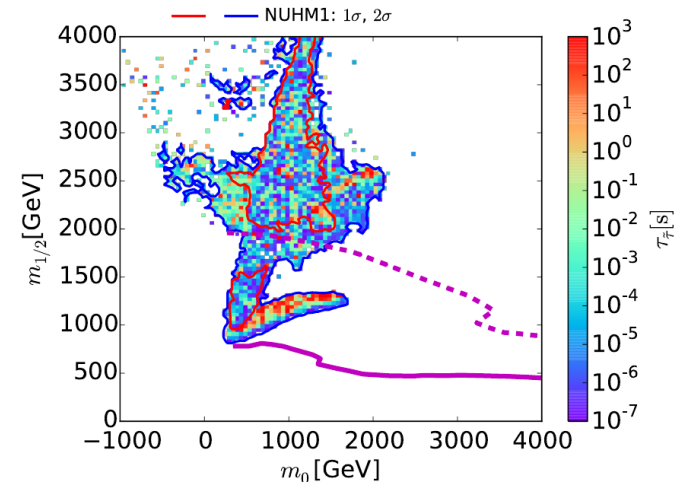
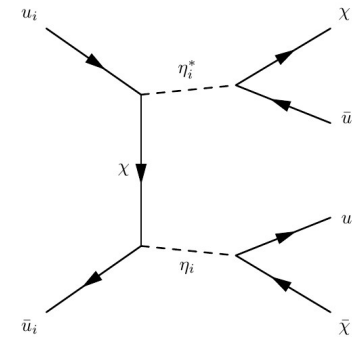
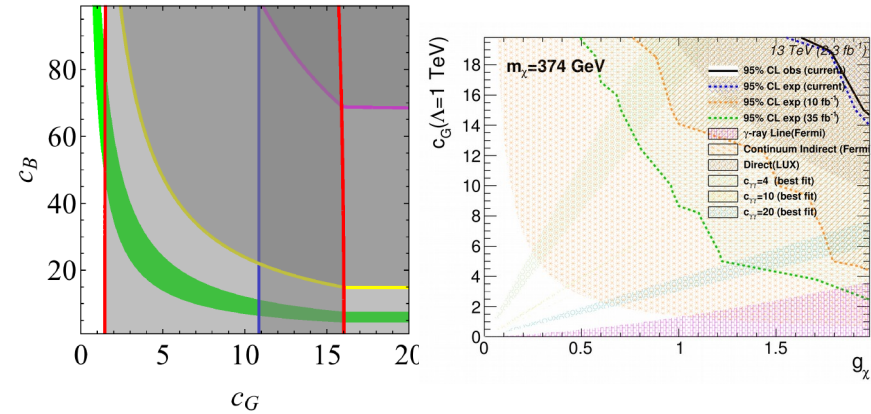
- Use 750 GeV excess to demonstrate complementarity of visible, MET+X, DD/ID
- spin-0 models highlight this particularly
- ▶ Should develop discovery-oriented models that leverage multi-messenger signals

Other simplified models

- t-channel, spin-2 mediator, pseudo-dirac DM
- ▶ Should be a future focus of the DM WG

Lessons from SUSY

- cascade decays → softer MET, more complex topologies
- Displaced vertices, metastable charged particles
- ▶ SDMM extensions needed to capture properties of complete models



**“SOMEBODY
CALL FOR
BACKUP!”**



h1/h2 Widths

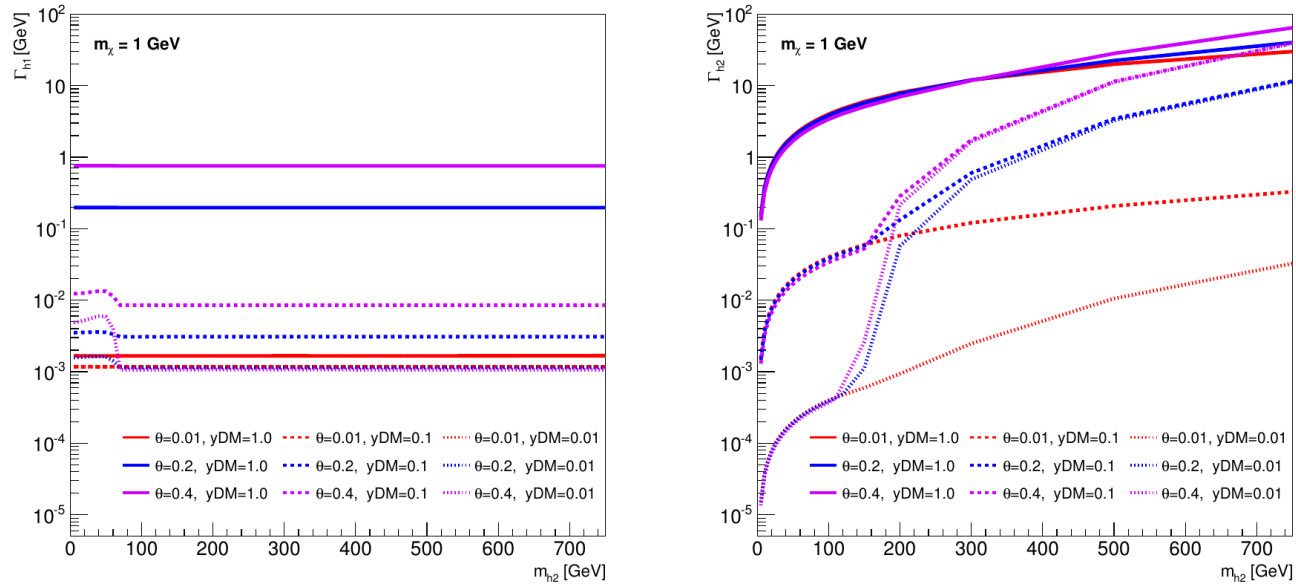


Figure 2. The h_1 (left) and h_2 (right) total decay width in scenarios C and D for several values of y_{DM} and θ . In both panels the DM mass is fixed to 1 GeV. An increase in the total decay width of the h_1 state at low masses and low values of DM coupling y_{DM} is due to the contribution of the $h_1 \rightarrow h_2 h_2$ decay channel. Note that the $\theta = 0.01$ lines for $y_{\text{DM}} = 0.1$ and 0.01 nearly overlap and hence are seen as a single line in the left plot.

2HDM

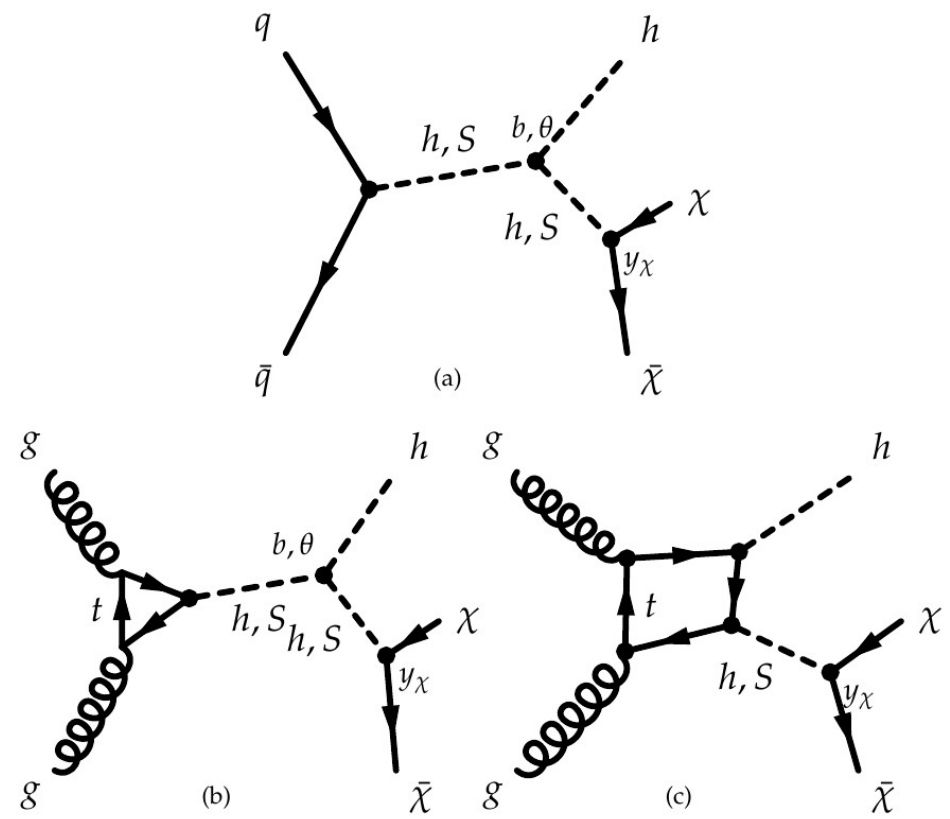
SMM perhaps simplest member of class of gauge invariant / unitary spin-0 DM models

Could also consider the wider class, incorporates 2HDMs

- 2HDM can be less restrictive vs SMM for scalar mediation
- **Could naturally incorporate pseudoscalar mediator**
- Straightforward connection with SUSY
- Included in DMF recommendations for mono-Higgs

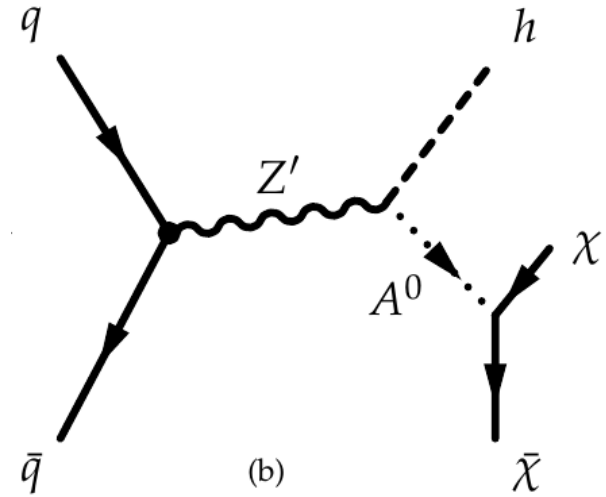
2HDM could be studied more broadly (ie: in terms of general MET+X signatures) within the LHC DM WG

- Eg: 1402.7074, 1502.06000



3.1.2 $\cancel{E}_T + \text{Higgs}$ from a scalar mediator

A real scalar singlet S coupling to DM can be introduced as a portal between SM and the dark sector through the Higgs field. The most general scalar potential is detailed in Ref. [ORMWo7], including terms that break \mathbb{Z}_2 . The \mathbb{Z}_2 symmetry, which causes the new scalar to also be a DM candidate, is not covered in this report, but follows Ref. [Car+14] introducing an additional coupling to DM that breaks \mathbb{Z}_2 and leads to a new invisible decay of S . For this reason, no symmetry is broken and no new interactions arise, so there is no dependence on the vacuum expectation value of S : a shift in the field leads to a redefinition of the model couplings. The new scalar S mixes with the SM Higgs boson, and couples to



3.1.3 *Higgs+ \cancel{E}_T signal from 2HDM model with a Z' and a new pseudoscalar*

In this simplified model [BLW14b], a new Z' resonance decays to a Higgs boson h plus a heavy pseudoscalar state A^0 in the 2HDM framework, which in turn decays to a DM pair. This model is represented in the diagram in Fig. 3.2 (b).

The motivation for coupling the dark matter to the pseudoscalar is that dark matter coupling to a Higgs or Z' boson is generically constrained by other signal channels and direct detection. A reason to consider this model is that it has different kinematics due to the on-shell Z' production, where for heavy Z' masses the \cancel{E}_T and