

# Indirect Detection Constraints on Dark Matter Model Space

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# Fermi Dwarf Analysis

Dwarf Spheroidal Galaxies large amount of DM  
Low Astrophysical Background

photon flux

$$\Phi_\gamma = \frac{1}{4\pi} \sum_f \frac{\langle \sigma v \rangle_f}{2m_\chi^2} \int_{E_{\min}}^{E_{\max}} \left( \frac{dN_\gamma}{dE_\gamma} \right)_f dE_\gamma J.$$

averaged annihilation xsec

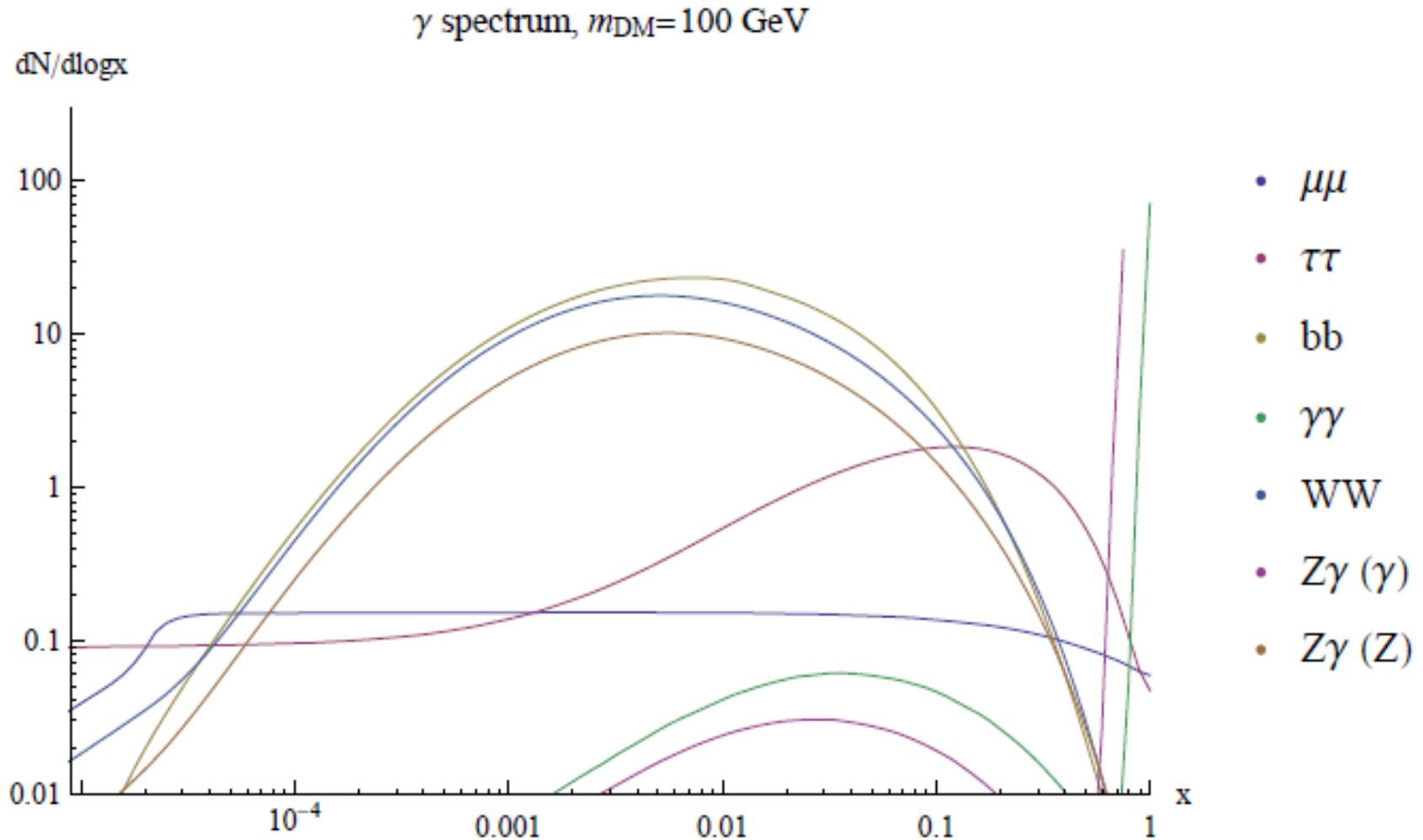
DM mass

Photon energy spectrum

Line of sight integral of DM density

$$J = \int_{\Delta\Omega} \int_{l.o.s} \rho^2(\mathbf{r}) dl d\Omega'.$$

# Spectrum

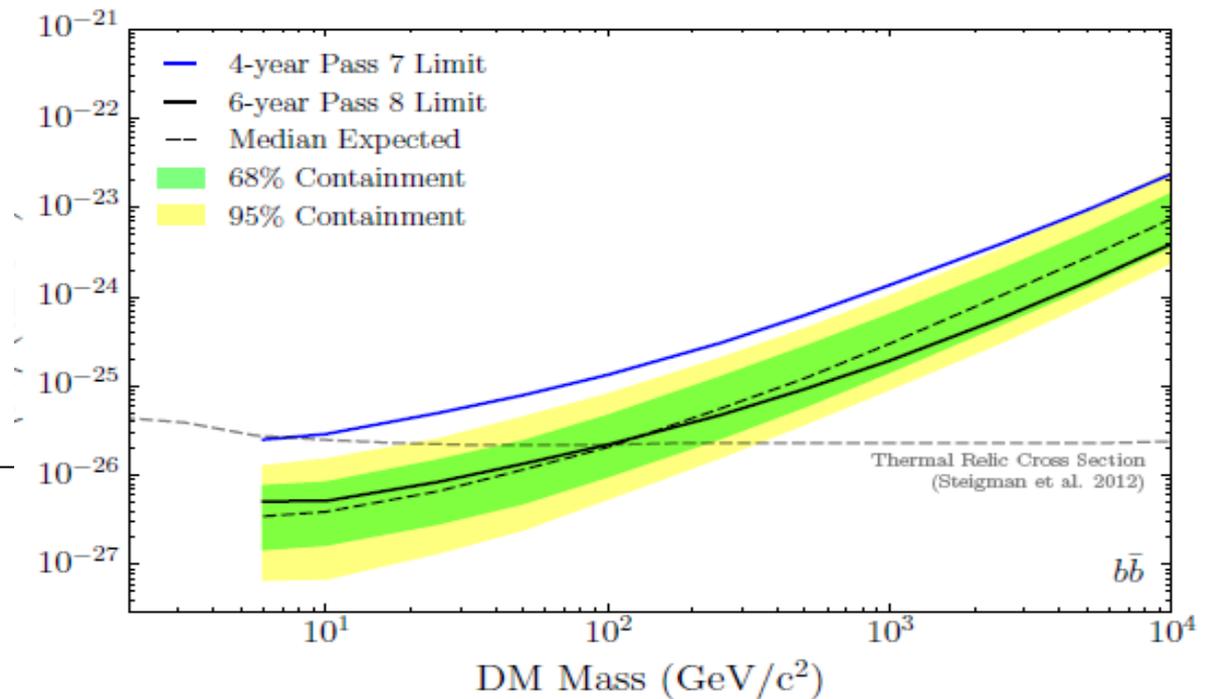


DM annihilates to various SM final states each with a characteristic photon spectrum

# Fermi Analysis combine 15 dwarf's with largest J factors, set 95% c.i. upper bound assuming 100% annihilation into a single channel, e.g. $b\bar{b}$

TABLE I. Properties of Milky Way dSphs.

Name	$\ell^a$ (deg)	$b^a$ (deg)	Distance (kpc)	$\log_{10}(J_{\text{obs}})^b$ ( $\log_{10}[\text{GeV}^2 \text{cm}^{-5}]$ )
Bootes I	358.1	69.6	66	$18.8 \pm 0.22$
Canes Venatici II	113.6	82.7	160	$17.9 \pm 0.25$
Carina	260.1	-22.2	105	$18.1 \pm 0.23$
Coma Berenices	241.9	83.6	44	$19.0 \pm 0.25$
Draco	86.4	34.7	76	$18.8 \pm 0.16$
Fornax	237.1	-65.7	147	$18.2 \pm 0.21$
Hercules	28.7	36.9	132	$18.1 \pm 0.25$
Leo II	220.2	67.2	233	$17.6 \pm 0.18$
Leo IV	265.4	56.5	154	$17.9 \pm 0.28$
Sculptor	287.5	-83.2	86	$18.6 \pm 0.18$
Segue 1	220.5	50.4	23	$19.5 \pm 0.29$
Sextans	243.5	42.3	86	$18.4 \pm 0.27$
Ursa Major II	152.5	37.4	32	$19.3 \pm 0.28$
Ursa Minor	105.0	44.8	76	$18.8 \pm 0.19$
Willman 1	158.6	56.8	38	$19.1 \pm 0.31$
Bootes II <sup>c</sup>	353.7	68.9	42	-
Bootes III	35.4	75.4	47	-
Canes Venatici I	74.3	79.8	218	$17.7 \pm 0.26$
Canis Major	240.0	-8.0	7	-
Leo I	226.0	49.1	254	$17.7 \pm 0.18$
Leo V	261.9	58.5	178	-
Pisces II	79.2	-47.1	182	-
Sagittarius	5.6	-14.2	26	-
Segue 2	149.4	-38.1	35	-
Ursa Major I	159.4	54.4	97	$18.3 \pm 0.24$



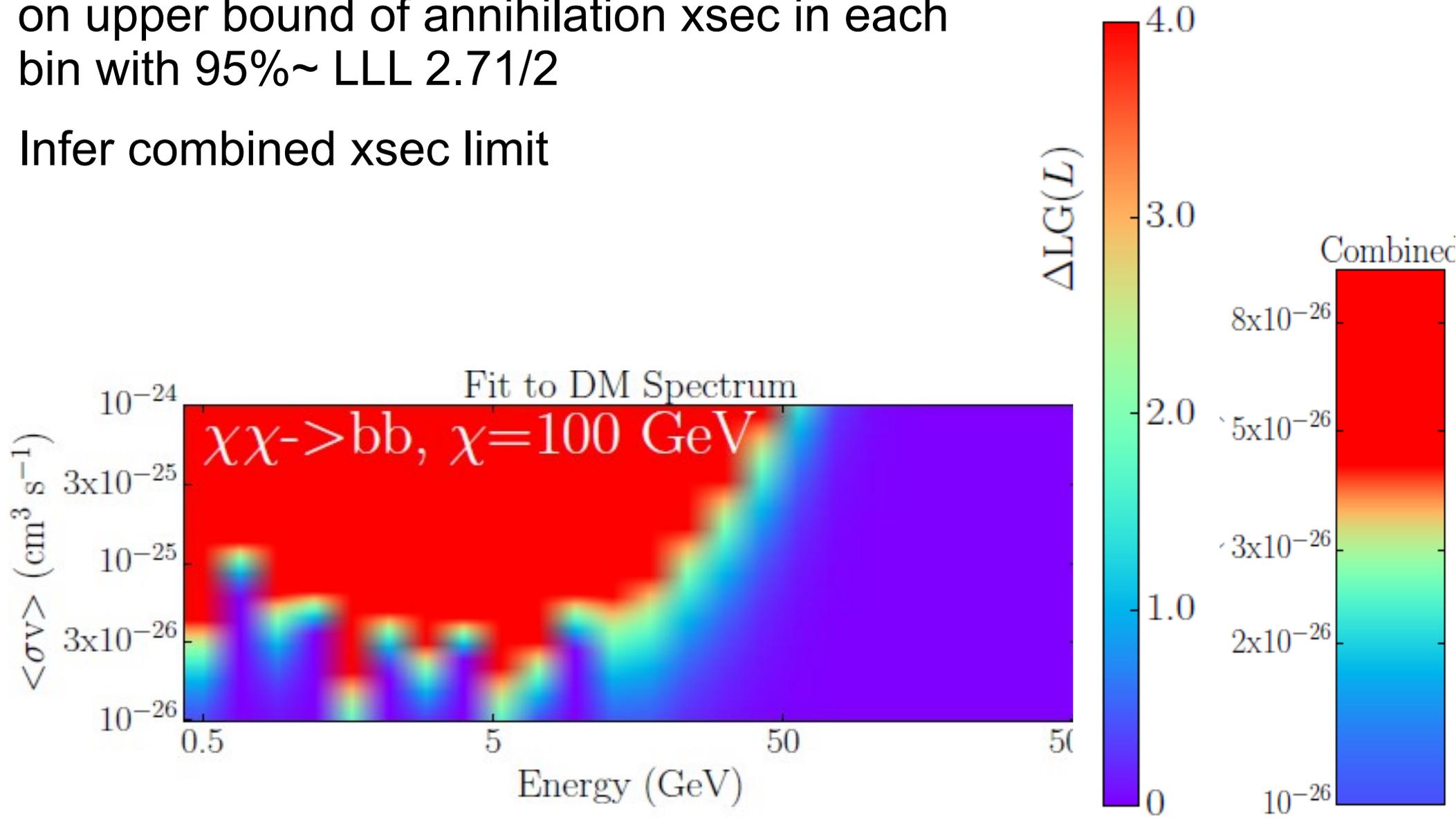
Choose DM mass and annihilation channel

Allow J factor to float with Least Log

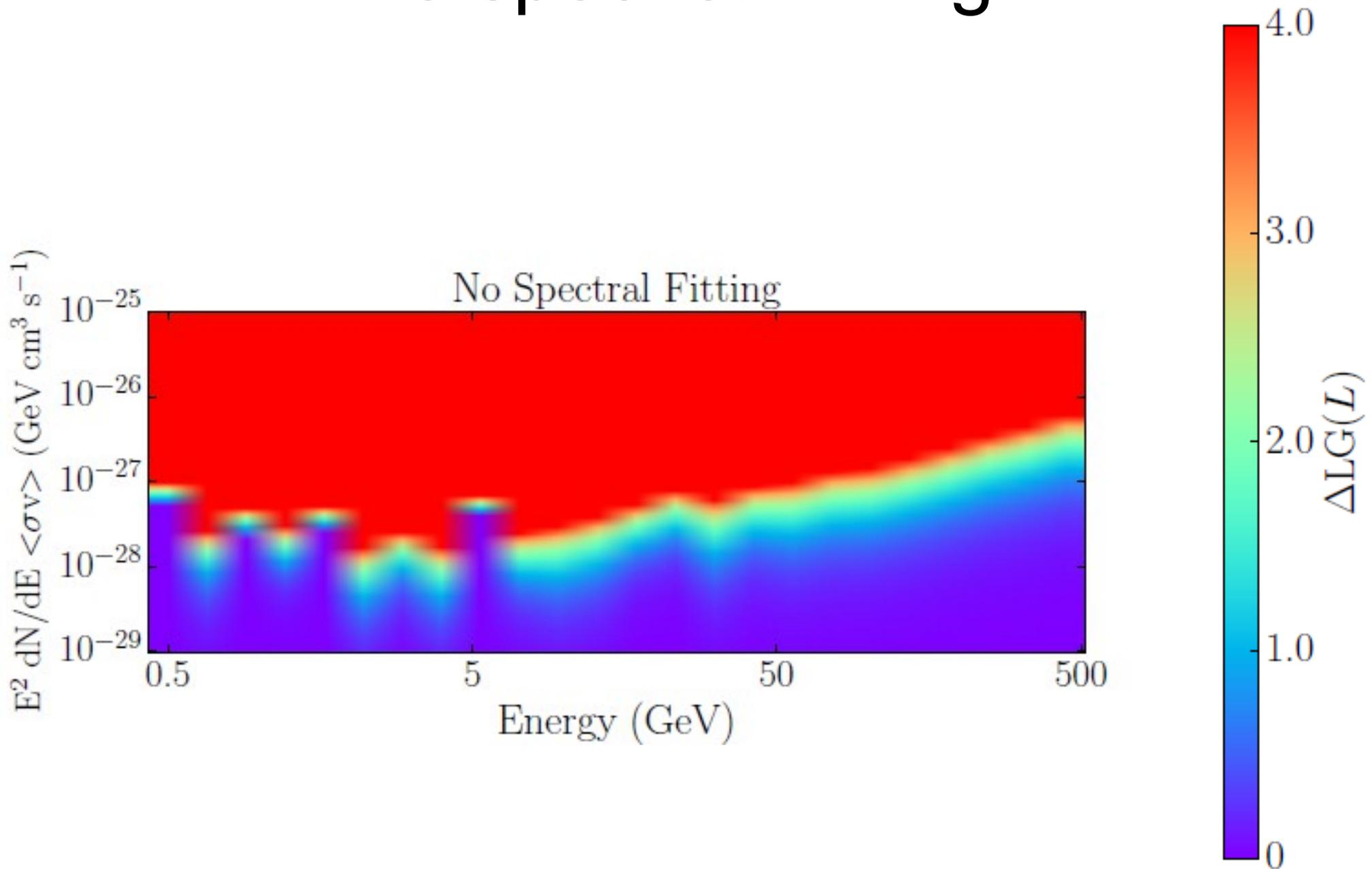
Likelihood cost  $\Delta LG(\mathcal{L}) = (J_{bf} - J_{meas})^2 / (2\sigma_J^2)$

Compare to null hypothesis no DM to set limit on upper bound of annihilation xsec in each bin with 95%~ LLL 2.71/2

Infer combined xsec limit



# No spectral Fitting



# Fermion Portal

Simplest EFT model: 1 operator 1 channel

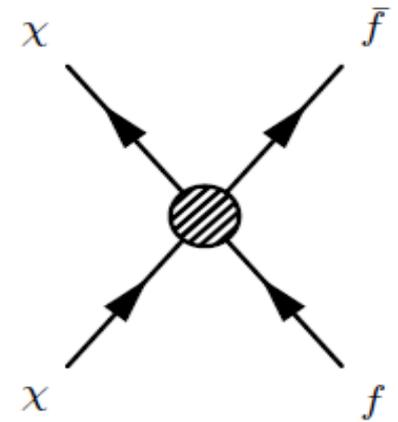
$$\mathcal{L}_f = \frac{\kappa_t}{\Lambda_t^2} \chi \Gamma \bar{\chi} t \Gamma \bar{t} + \frac{\kappa_b}{\Lambda_b^2} \chi \Gamma \bar{\chi} b \Gamma \bar{b} + \frac{\kappa_\tau}{\Lambda_\tau^2} \chi \Gamma \bar{\chi} \tau \Gamma \bar{\tau} + \frac{\kappa_\nu}{\Lambda_\nu^2} \chi \Gamma \bar{\chi} \nu \Gamma \bar{\nu}.$$

Allow visible total annihilation rate below the thermal rate  
Without over-closing the universe

Light DM For now consider annihilation to b,  $\tau$  and invisible channel

$$\langle \sigma v \rangle_{\text{tot}} = \langle \sigma v \rangle_b + \langle \sigma v \rangle_\tau + \langle \sigma v \rangle_\nu.$$

$$\propto a \left( \kappa_b / \Lambda_b^2 \right)^2 + b \left( \kappa_\tau / \Lambda_\tau^2 \right)^2 + c \left( \kappa_\nu / \Lambda_\nu^2 \right)^2$$



First Fix the Annihilation rate as desired

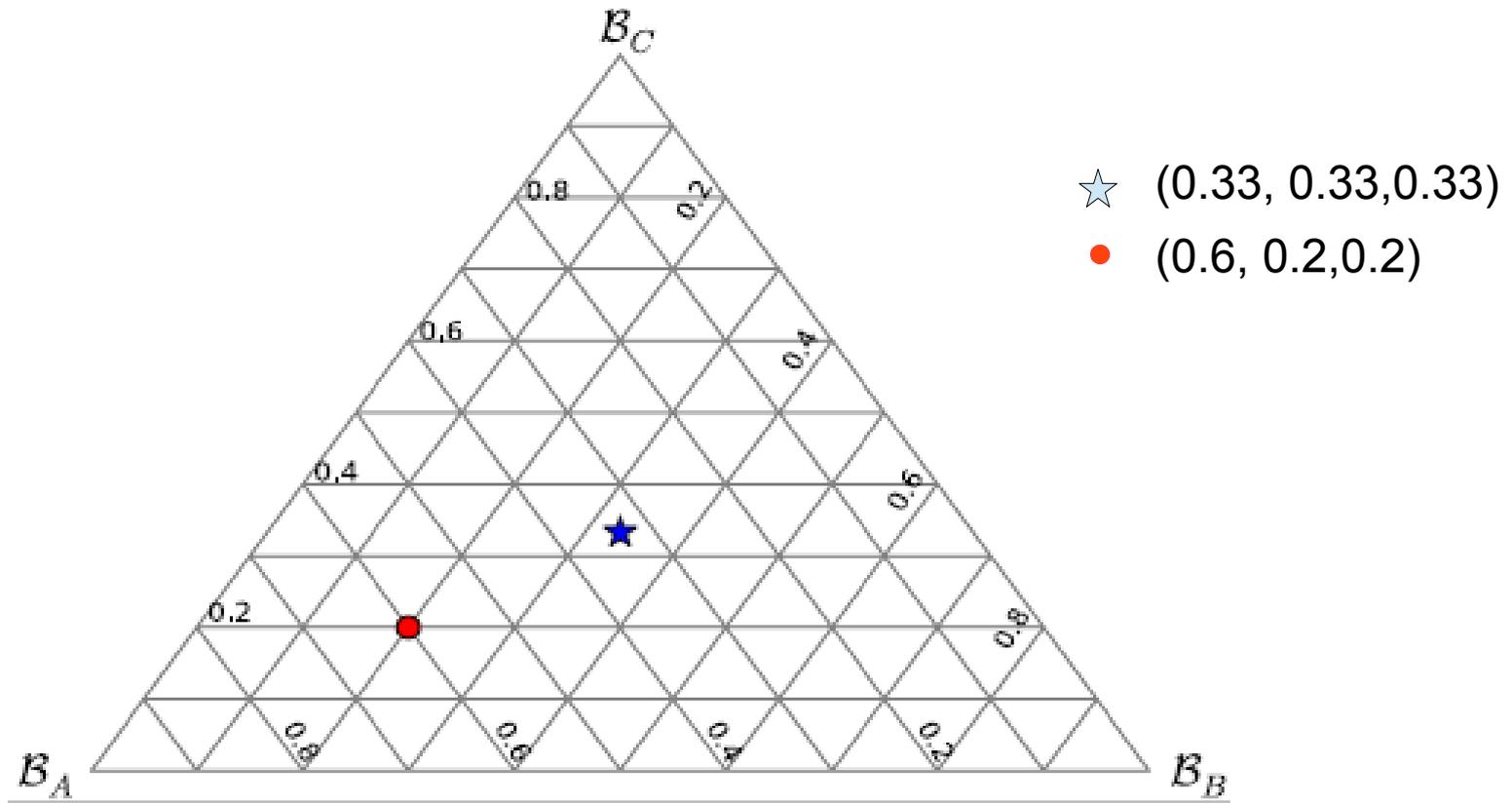
Dividing out by the total rate to define partial rate  $R_i = \langle \sigma v \rangle_i / \langle \sigma v \rangle_{\text{tot}}$

get a constraint between the partial annihilation rates

$$R_1 + R_2 + R_3 + \dots = 1.$$

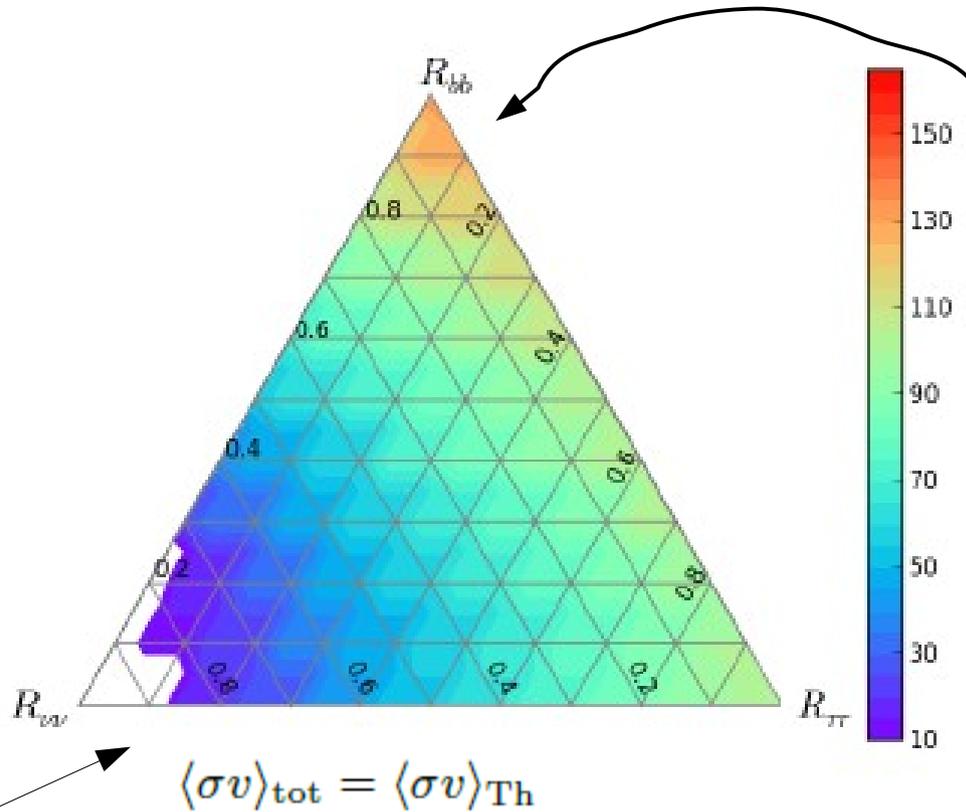
# Three Parameters and 1 constraint may be visualized on 2-D surface as triangle

The partial rates are saturated at the corners of the triangle



$$\Phi_\gamma = \frac{1}{4\pi} \sum_f \frac{\langle \sigma v \rangle_f}{2m_\chi^2} \int_{E_{\min}}^{E_{\max}} \left( \frac{dN_\gamma}{dE_\gamma} \right)_f dE_\gamma J.$$

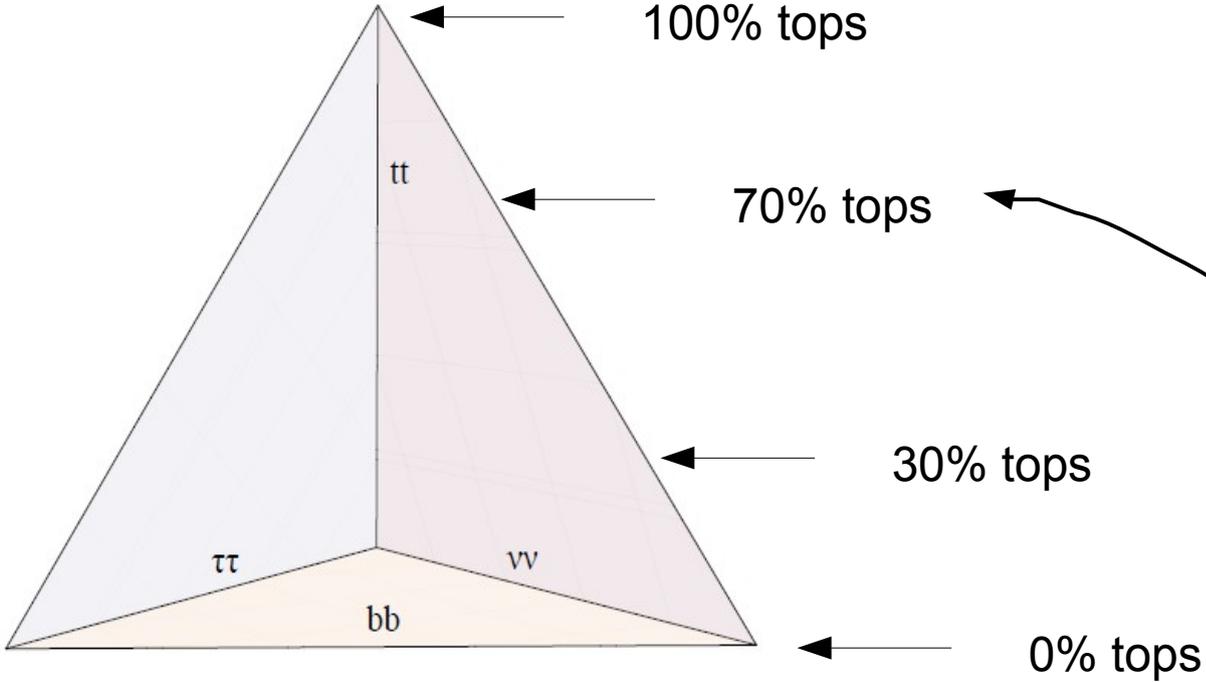
# Pass 8



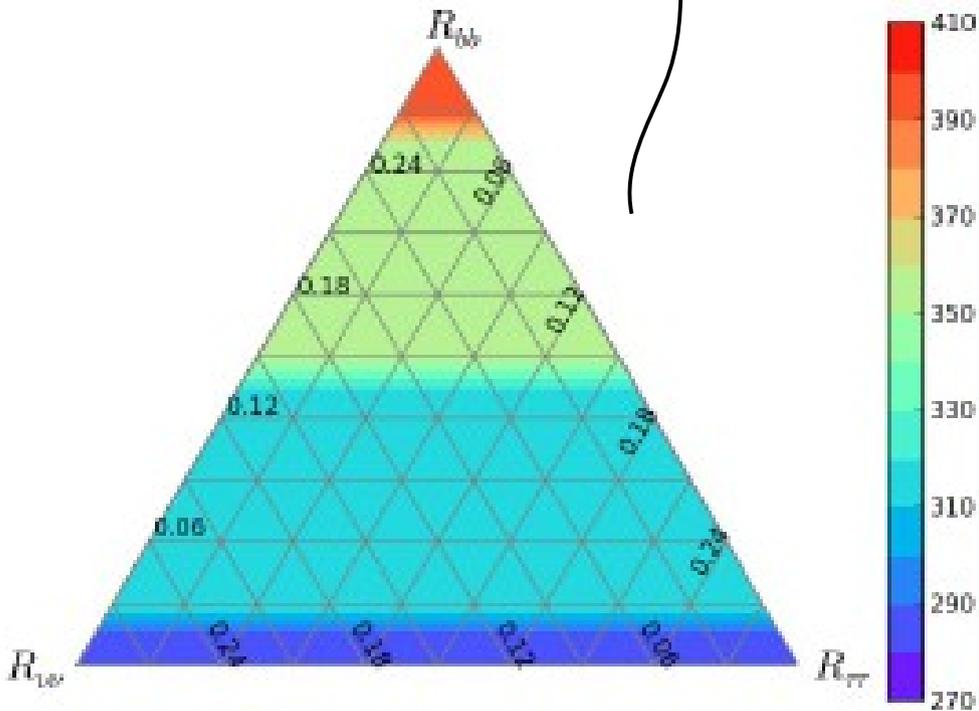
Max exclusion improves by factor  $\sim 4$

order 10s GeV min mass bounds even for visible annihilation rates at 30% of thermal rate

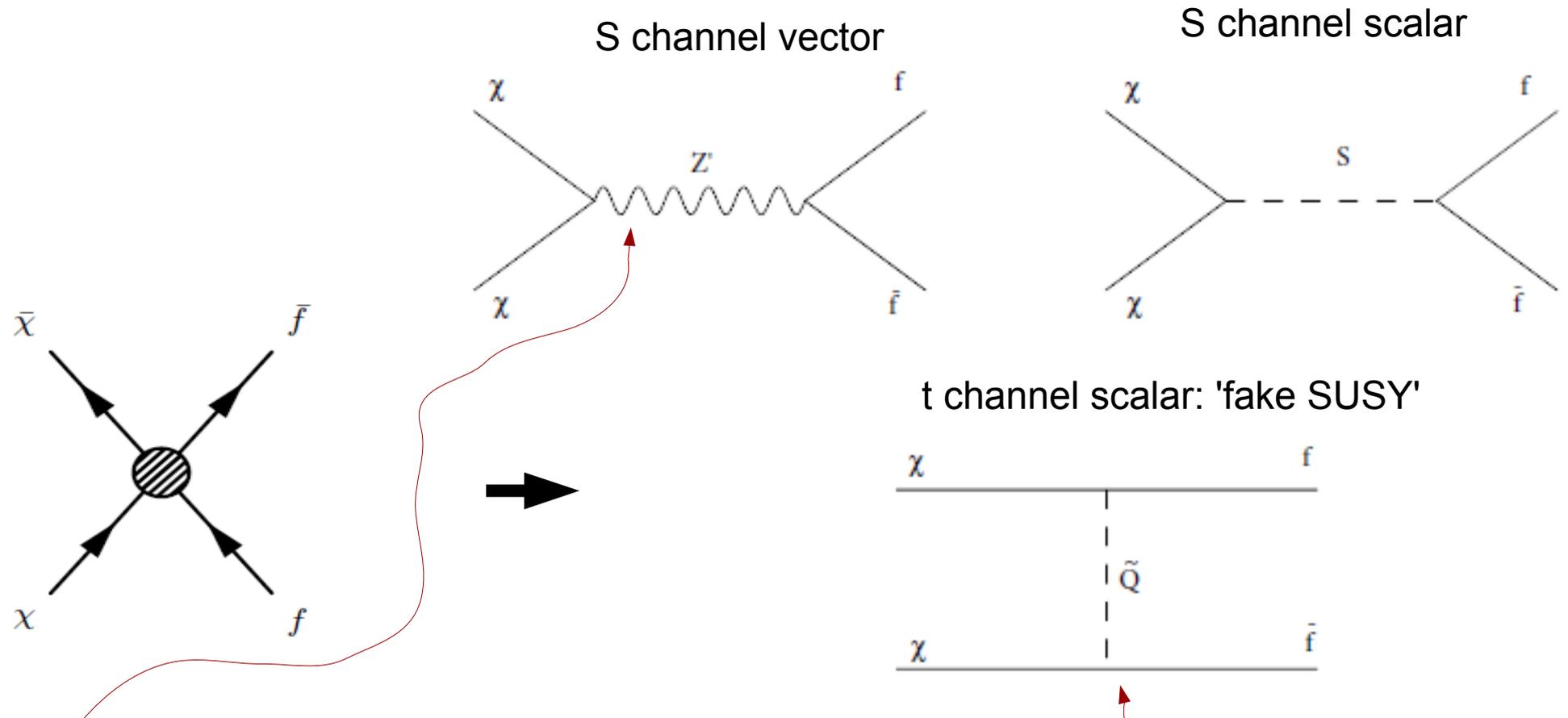
# 4 channels



Shown, slice with 70% ann into tops  
Pass 7 data



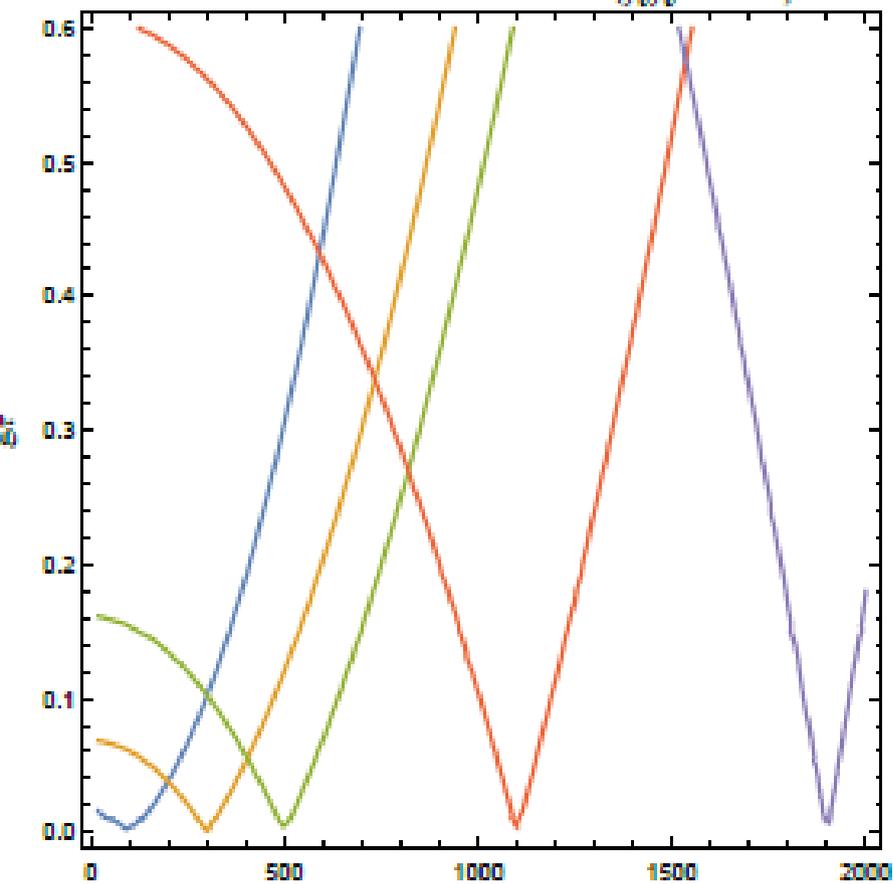
# EFT to Simplified Model



$$\mathcal{L} = V_\mu \bar{\chi} \gamma^\mu (g_\chi^V - g_\chi^A \gamma^5) \chi + \sum_f V_\mu \bar{f} \gamma^\mu (g_f^V - g_f^A \gamma^5) f$$

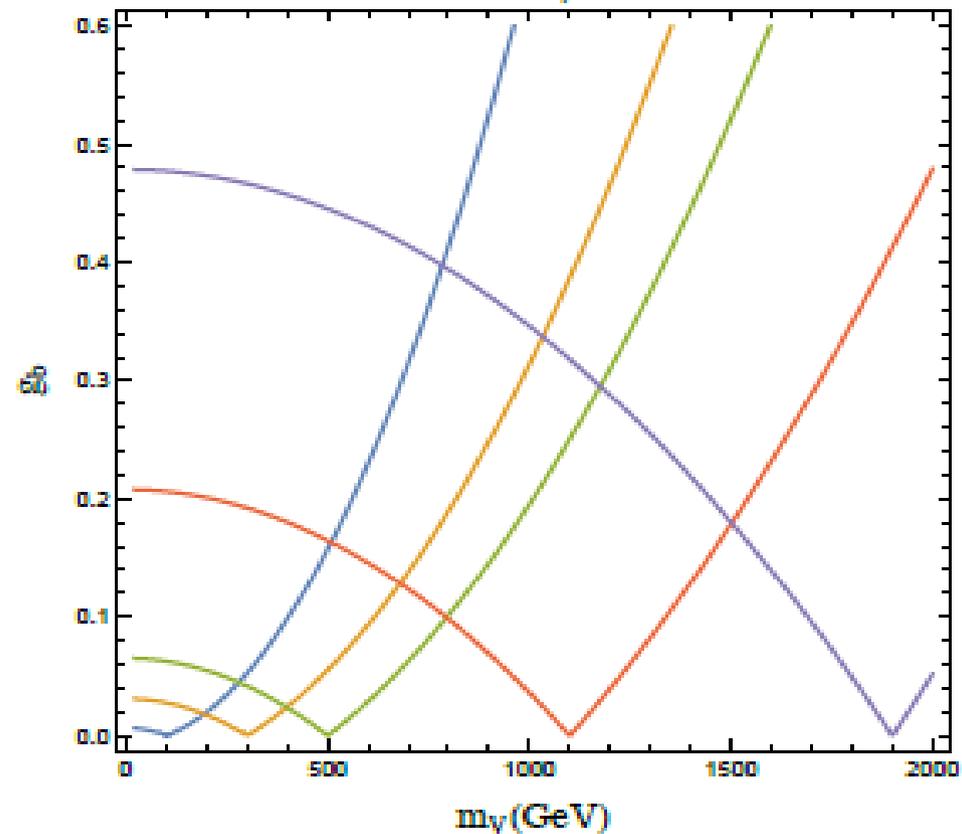
$$\mathcal{L} = \sum_i g_i \phi_i^* \bar{\chi} P_R f_i + h.c.,$$

Vector Mediator Limits ( $\chi\chi \rightarrow \tau\tau$ )



- Blue line:  $m_\chi=50$  GeV,  $\langle\sigma v\rangle=0.56\langle\sigma v\rangle_{\text{Therm}}$
- Orange line:  $m_\chi=150$  GeV,  $\langle\sigma v\rangle=1.4\langle\sigma v\rangle_{\text{Therm}}$
- Green line:  $m_\chi=250$  GeV,  $\langle\sigma v\rangle=2.2\langle\sigma v\rangle_{\text{Therm}}$
- Red line:  $m_\chi=550$  GeV,  $\langle\sigma v\rangle=4.7\langle\sigma v\rangle_{\text{Therm}}$
- Purple line:  $m_\chi=950$  GeV,  $\langle\sigma v\rangle=8.4\langle\sigma v\rangle_{\text{Therm}}$

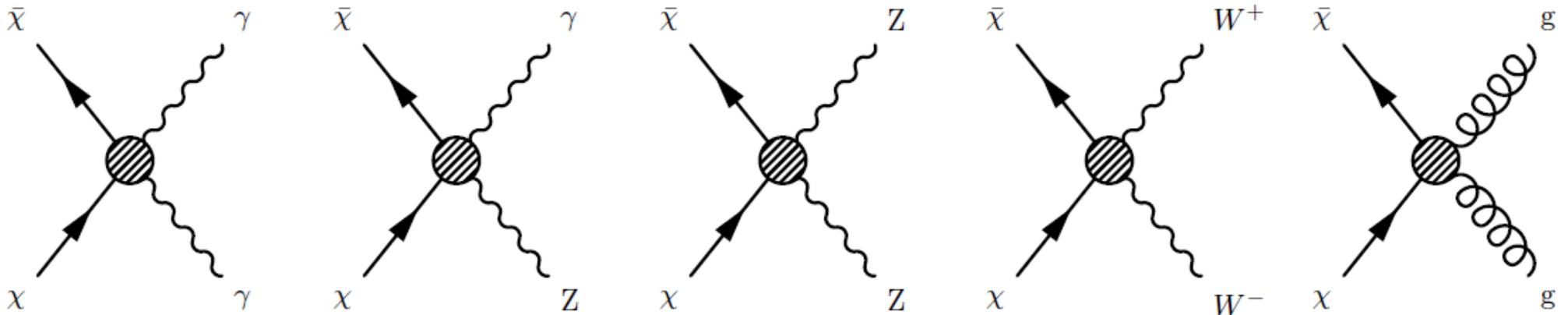
Vector Mediator Limits, 70% bb and 30%  $\tau\tau$



# Models with Interfering Channels

## Gauge Boson Portal Models

$$\mathcal{L} = \frac{\kappa_1}{\Lambda^3} \bar{\chi} \gamma^5 \chi B_{\mu\nu} B^{\mu\nu} + \frac{\kappa_2}{\Lambda^3} \bar{\chi} \gamma^5 \chi W_{\mu\nu}^i W_i^{\mu\nu} + \frac{\kappa_3}{\Lambda^3} \bar{\chi} \gamma^5 \chi G_{\mu\nu}^a G_a^{\mu\nu}$$



$$\langle \sigma v_{rel} \rangle_{WW} = \frac{\kappa_2^2}{4\pi\Lambda^6} \sqrt{1 - \frac{m_W^2}{m_\chi^2}} (16m_\chi^4 - 16m_W^2 m_\chi^2 + 6m_W^4)$$

$$\langle \sigma v_{rel} \rangle_{ZZ} = \frac{(\kappa_1 s_w^2 + \kappa_2 c_w^2)^2}{8\pi\Lambda^6} \sqrt{1 - \frac{m_Z^2}{m_\chi^2}} (16m_\chi^4 - 16m_Z^2 m_\chi^2 + 6m_Z^4)$$

$$\langle \sigma v_{rel} \rangle_{Z\gamma} = \frac{s_w^2 c_w^2 (\kappa_2 - \kappa_1)^2}{16\pi m_\chi^2 \Lambda^6} (4m_\chi^2 - m_Z^2)^3$$

$$\langle \sigma v_{rel} \rangle_{\gamma\gamma} = \frac{4(\kappa_1 c_w^2 + \kappa_2 s_w^2)^2}{2\pi\Lambda^6} m_\chi^4$$

$$\langle \sigma v_{rel} \rangle_{gg} = \frac{16\kappa_3^2}{\pi\Lambda^6} m_\chi^4$$

Defining effective cut-off

$$(k_i = \kappa_i / \Lambda^3)$$

The constraint has the form

$$(ak_1^2 + bk_2^2)^2 + ck_3^2 = \langle \sigma v \rangle_{tot}$$

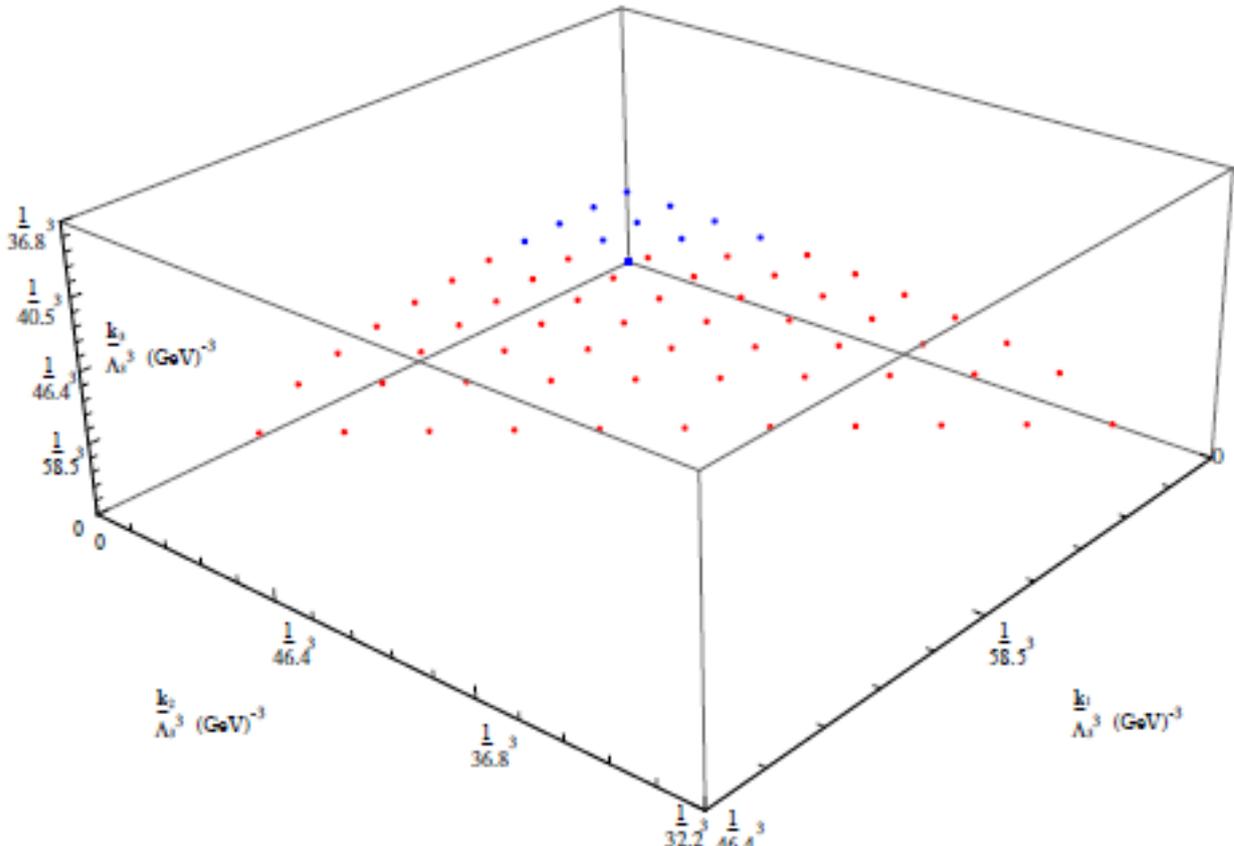
Fixing  $\langle \sigma v \rangle_{tot}$  the coefficients sit on a hypersurface

The constraint no longer factorizes multiple operators contribute to a single channel

A single operator coefficient contributes to multiple channels

To visualize parameter space we can fix total annihilation xsec. We can vary DM mass  
And look at the hypersurface in effective coefficient space where the constraint is satisfied

Each point in parameter space has specific admixture of partial annihilation rates into 5 channels  
We add up the total flux and see if the point is excluded.

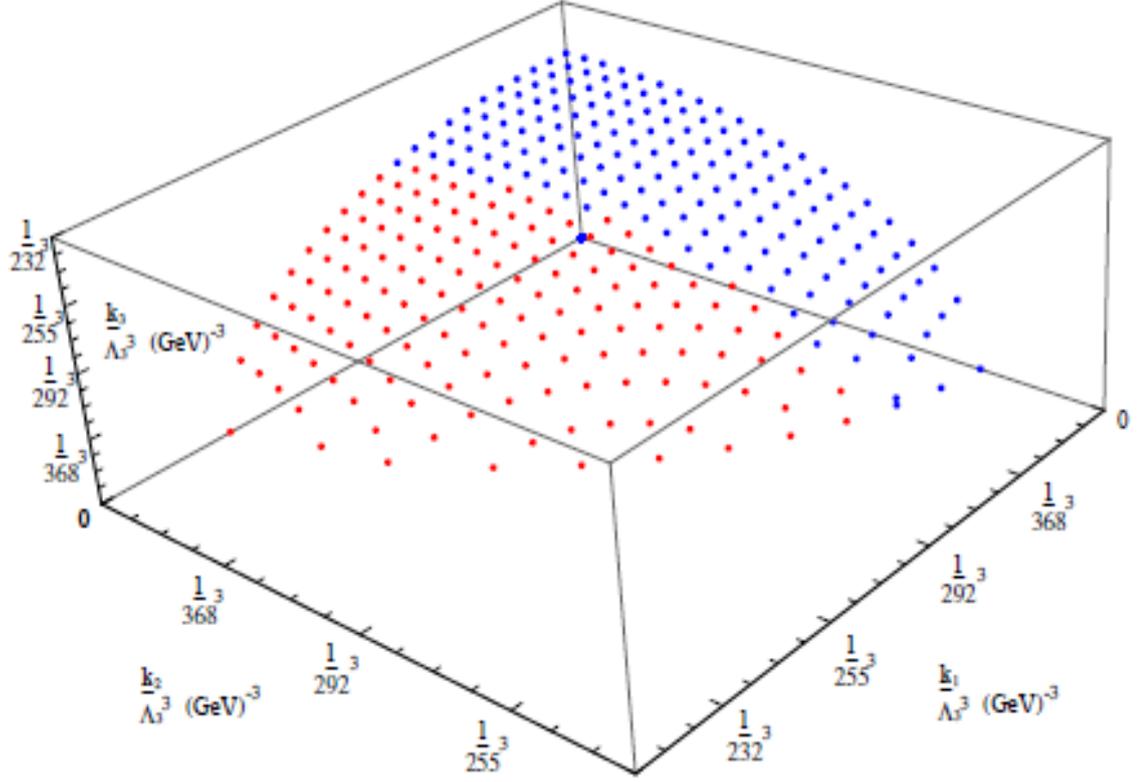


DM mass = 10 GeV  
 Rate constraint is saturated  
 By low effective cut-offs

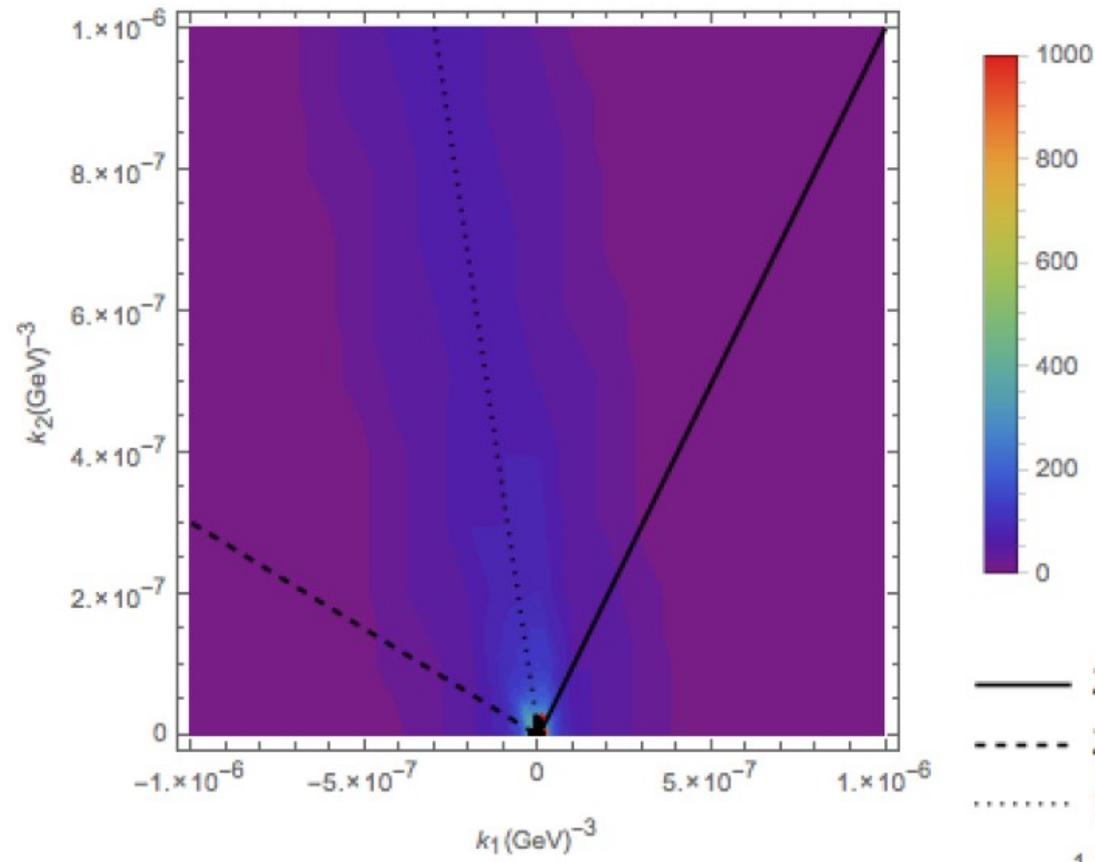
←  
 at low DM mass only g and  $\gamma$   
 channels kinematically accessible

Surfaces in effective cut off space  
 Red points ruled out  
 Total annihilation rate 10x's thermal

DM mass = 150 GeV  
 Rate constraint is saturated  
 By higher effective cut-offs



$$(k_i = \kappa_i/\Lambda^3)$$

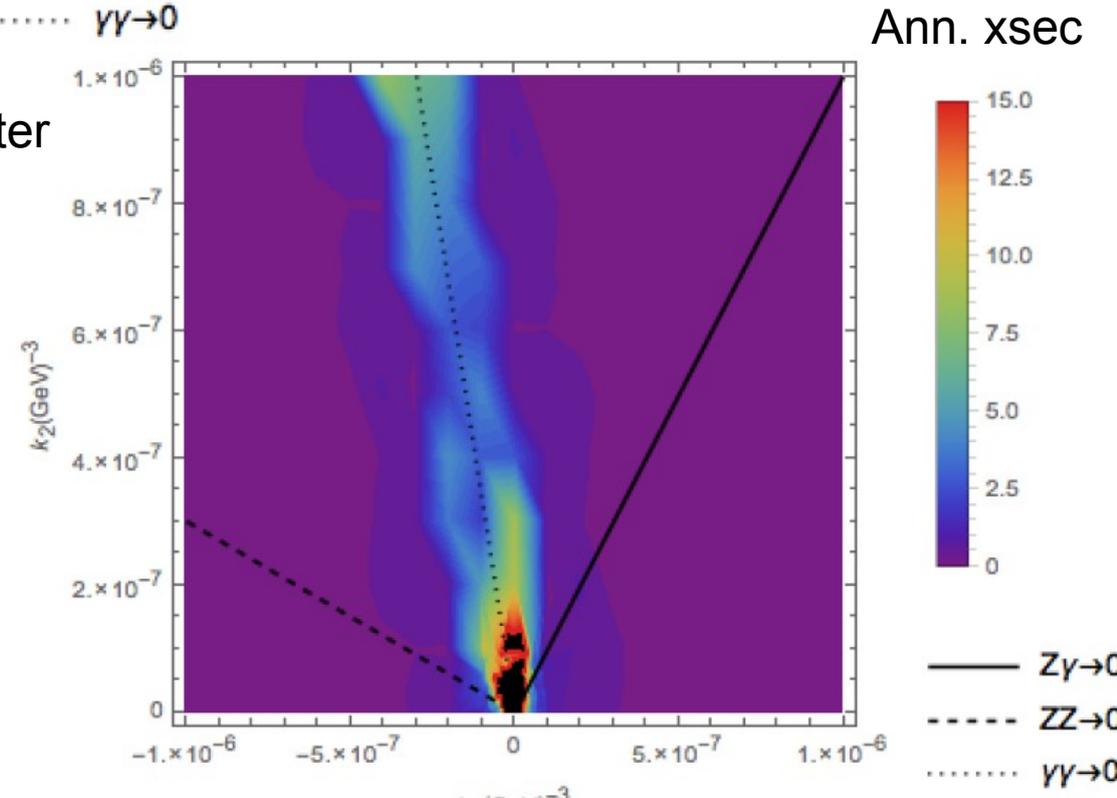


DM mass bound

- $Z\gamma \rightarrow 0$
- - -  $ZZ \rightarrow 0$
- .....  $\gamma\gamma \rightarrow 0$

Let annihilation xsec vary over the parameter Space of effective operator coeff.

Here 0% ann. into gluons



Ann. xsec

- $Z\gamma \rightarrow 0$
- - -  $ZZ \rightarrow 0$
- .....  $\gamma\gamma \rightarrow 0$

# Compare LHC search to Indirect limit

