# Fluctuations of the multiplicity of produced particles in pA collisions

Stéphane Munier

Centre de physique théorique

Based on a work to appear (Phys Rev D, 2016) with Tseh Liou and A.H. Mueller

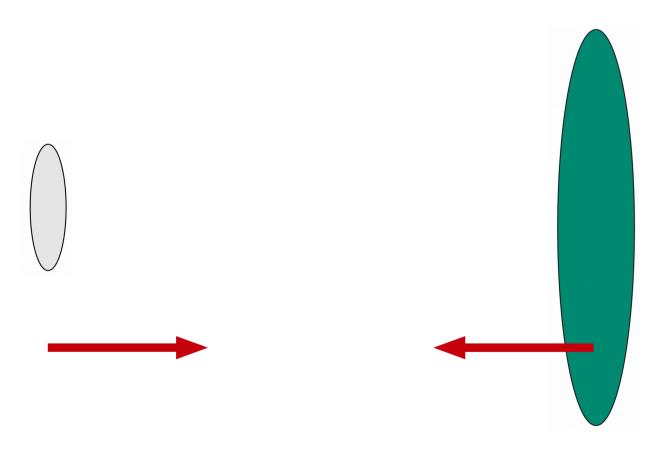






#### Observable

Consider proton-nucleus collisions



What is the distribution of the number of particles observed in the fragmentation region of the proton?

#### Outline

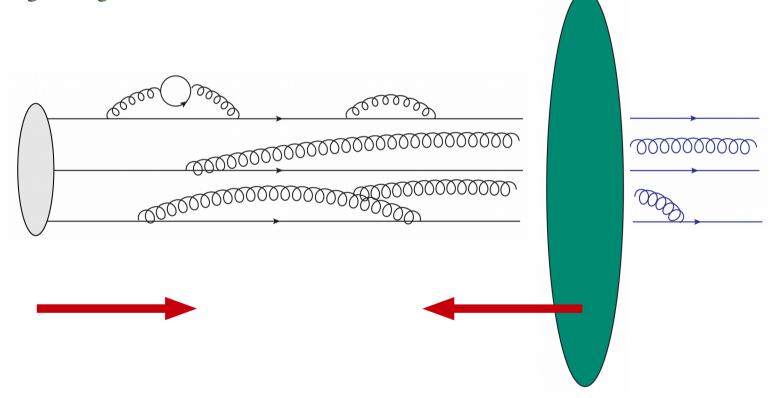
- \* Picture of particle production in pA collisions
- \* How a hadronic state dresses at high energies: the color dipole model
- \* Probability distribution of the particle multiplicity

#### Particle production in a dilute-dense collision

The proton is an initially **dilute** object while the nucleus is a **dense** object, characterized by a saturation scale  $Q_A$ 

At the time of the interaction: The proton is in a given Fock state, essentially made of gluons if

its rapidity is large enough.

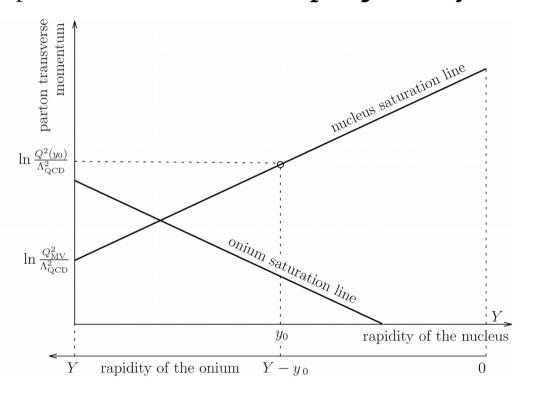


**Final state:** The gluons which have a transverse momentum less than  $Q_A$  are **freed** thanks to multiple scatterings with the nucleus and go the the final state (in the forward region of the proton), while the others essentially do not interact and just recombine.

#### Particle production in a dilute-dense collision

As a model for a dilute system, we consider an onium (quark-antiquark pair in a color-singlet state)

We choose a frame in such that the saturation scale of the nucleus is much larger than that of the proton, and look at central rapidity in that frame.



The multiplicity of produced particles at central rapidity is related to the gluon density in the onium by

$$\frac{dN}{dy}(y \simeq 0) = x G(x, Q_A^2(y_0))$$

where 
$$x = e^{-(Y-y_0)} \frac{Q_A(y_0)}{\text{onium mass}}$$

The multiplicity measured in the onium/proton fragmentation region in an event is the gluon number density at the scale  $Q_{\rm n}$  in the corresponding realization of the QCD evolution.

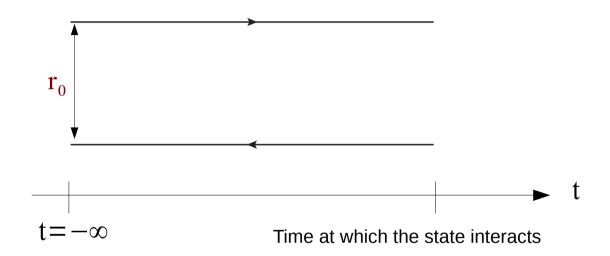
Let's try and understand theoretically the event-by-event fluctuations of this gluon density!

#### Outline

- \* Picture of particle production in pA collisions
- \* How a hadronic state dresses at high energies: the color dipole model
- \* Probability distribution of the particle multiplicity

# QCD calculation: the color dipole model

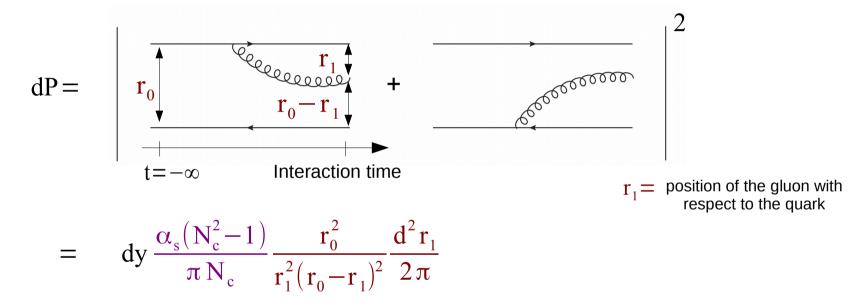
To simplify, we start with a color neutral quark-antiquark pair (=onium) of given transverse size.



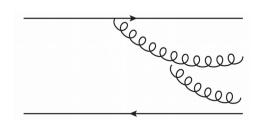
## QCD calculation: the color dipole model

To simplify, we start with a color neutral quark-antiquark pair (=onium) of given transverse size.

Probability of observing 1 gluon fluctuation when one increases the rapidity from 0 to dy:



For finite rapidities, one needs to consider higher-orders:



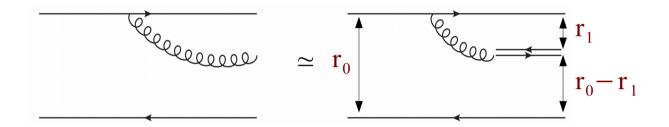


# QCD calculation: the color dipole model



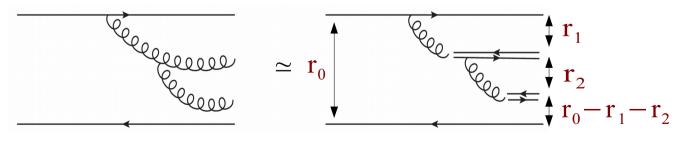
Trick: large number-of-color limit!





Gluon emission is interpreted as a color-dipole splitting, with proba. 
$$dP = dy \frac{\alpha_s N_c}{\pi} \frac{r_0^2}{r_1^2 (r_0 - r_1)^2} \frac{d^2 r_1}{2\pi}$$

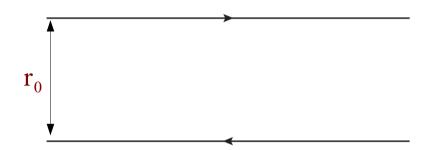
Higher-order fluctuations are generated by a branching process:

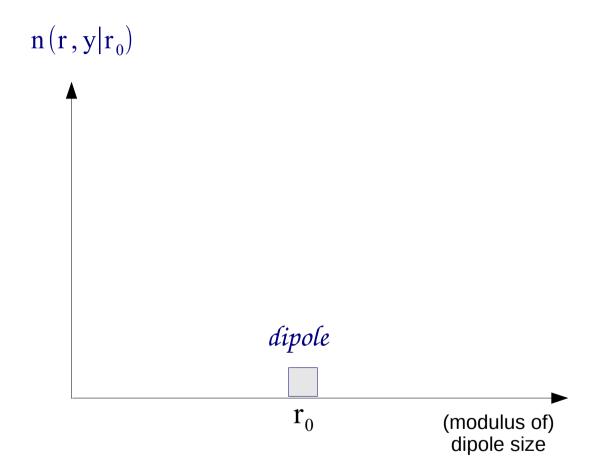


Two successive dipole branchings

 $x G(x,k^2) \sim density of dipoles of size 1/k after evolution over the rapidity <math>y = log 1/x$ 

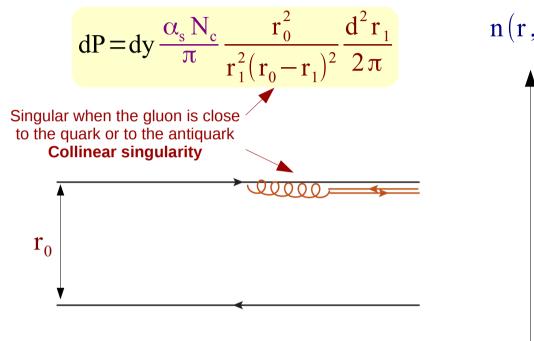
$$dP = dy \frac{\alpha_s N_c}{\pi} \frac{r_0^2}{r_1^2 (r_0 - r_1)^2} \frac{d^2 r_1}{2\pi}$$



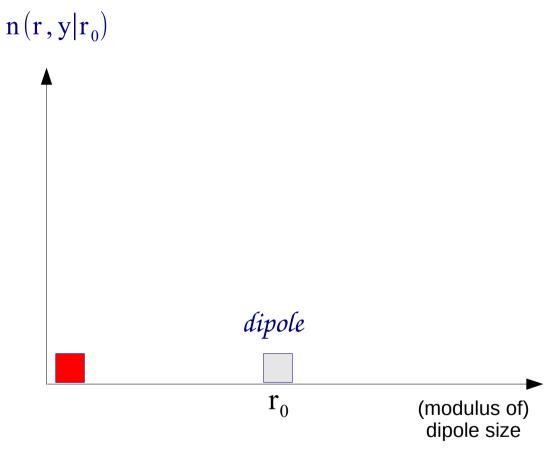


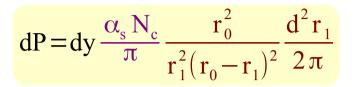
$$y=0$$

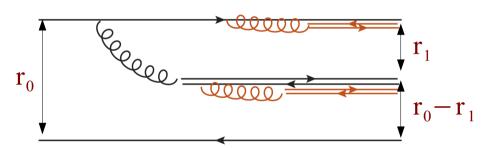
 $y \ge 0$ 



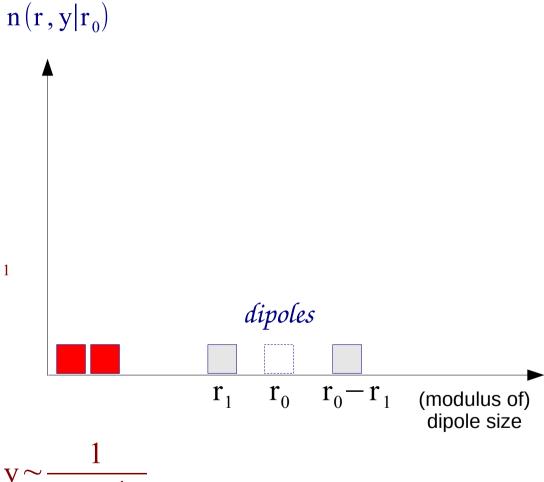
The small-dipole size region gets very easily populated



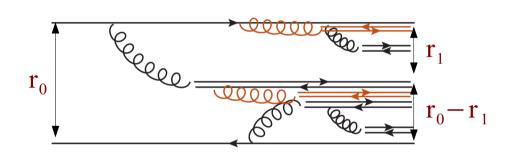




Dipoles of size of the order of the size of the initial dipole (are larger) need sone finite rapidity to get produced

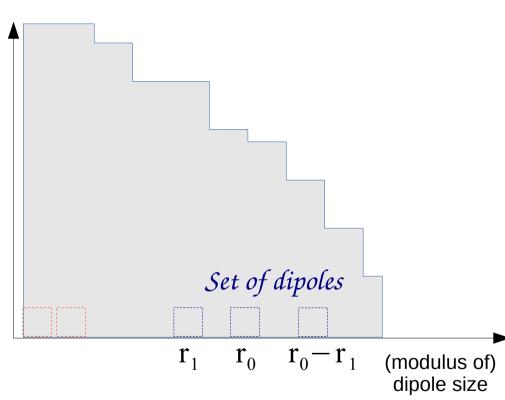


$$dP = dy \frac{\alpha_s N_c}{\pi} \frac{r_0^2}{r_1^2 (r_0 - r_1)^2} \frac{d^2 r_1}{2\pi}$$



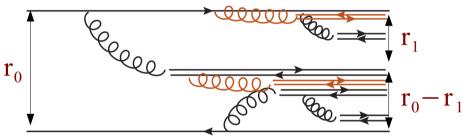
The total number of dipoles grows exponentially with rapidity through a branching process

$$\log n(r, y|r_0)$$

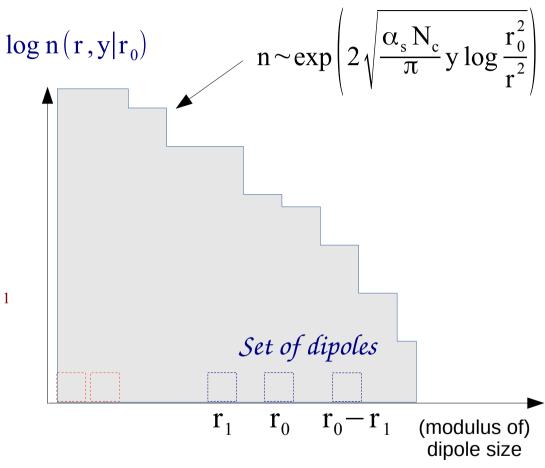


$$y \gg \frac{1}{\alpha_s N_c/\pi}$$

$$dP = dy \frac{\alpha_s N_c}{\pi} \frac{r_0^2}{r_1^2 (r_0 - r_1)^2} \frac{d^2 r_1}{2\pi}$$



The total number of dipoles grows exponentially with rapidity through a branching process



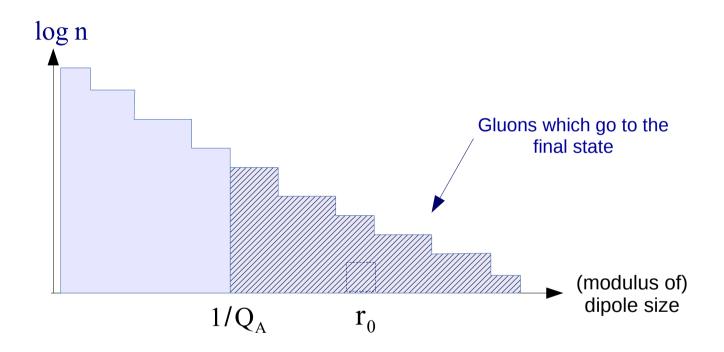
$$y \gg \frac{1}{\alpha_s N_c / \pi}$$

#### Outline

- \* Picture of particle production in pA collisions
- \* How a hadronic state dresses at high energies: the color dipole model
- \* Probability distribution of the particle multiplicity

# $\begin{array}{c} \textbf{\textit{Dipole-nucleus scattering: total multiplicity}} \\ \frac{dN}{dy}(y \! \simeq \! 0) \! = \! x \, G(x, Q_A^2(y_0)) \ \ \text{where} \ \ x \! = \! e^{-(Y-y_0)} \frac{Q_A(y_0)}{\text{onium mass}} \end{array}$

$$\frac{dN}{dy}(y\!\simeq\!0)\!=\!x\,G\!\left(x\,,Q_{A}^{2}\!\left(y_{0}\right)\right)\;\;\text{where}\;\;x\!=\!e^{-(Y-y_{0})}\frac{Q_{A}\!\left(y_{0}\right)}{\text{onium mass}}$$

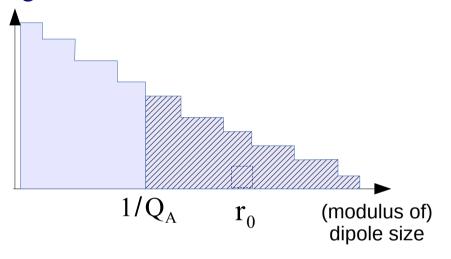


The gluons which go to the final state, i.e. which are freed in the scattering, correspond to dipoles which have a size larger than the inverse saturation scale of the nucleus.

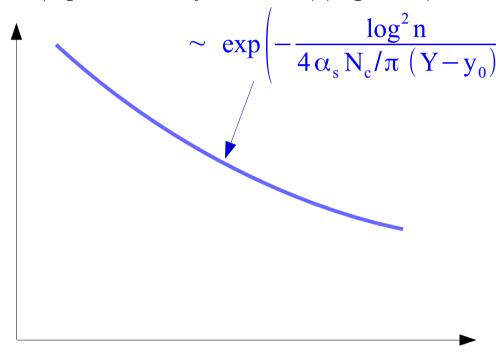
# Dipole-nucleus scattering: total multiplicity

$$\frac{dN}{dy}(y \! \simeq \! 0) \! = \! x \, G\!\left(x \, , Q_A^2\!\left(y_0\right)\right) \quad \text{where} \quad x \! = \! e^{-(Y-y_0)} \frac{Q_A\!\left(y_0\right)}{\text{onium mass}}$$

#### log n Perturbative calculation



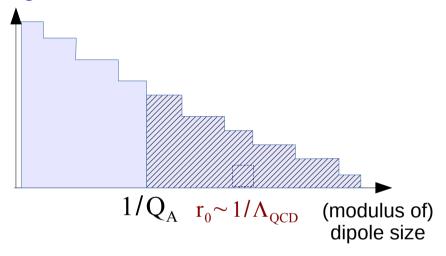
proba(n particles in final state)(log scale)



# Dipole-nucleus scattering: total multiplicity

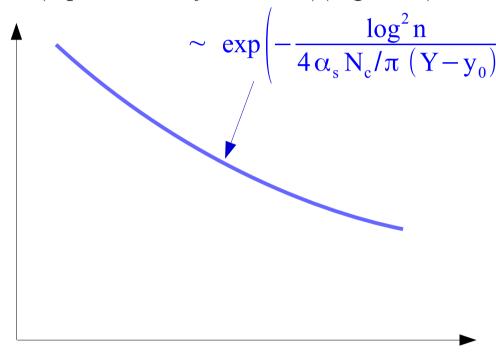
$$\frac{dN}{dy}(y\!\simeq\!0)\!=\!x\,G\!\left(x\,,Q_{A}^{2}\!\left(y_{0}\right)\right)\ \ \text{where}\ \ x\!=\!e^{-(Y-y_{0})}\frac{Q_{A}\!\left(y_{0}\right)}{\text{onium mass}}$$

#### log n Perturbative calculation



But in pA, the "dipole" is a proton!

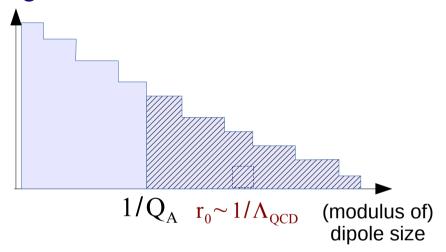
proba(n particles in final state)(log scale)



# Dipole-nucleus scattering: total multiplicity

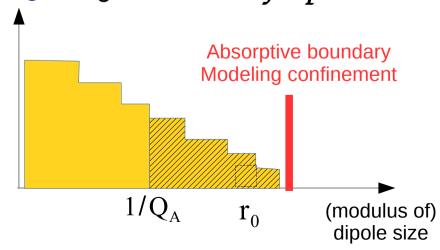
$$\frac{dN}{dy}(y \simeq 0) = x \, G\left(x \,, Q_A^2\left(y_0\right)\right) \quad \text{where} \quad x = e^{-(Y-y_0)} \frac{Q_A(y_0)}{\text{onium mass}}$$

#### log n Perturbative calculation

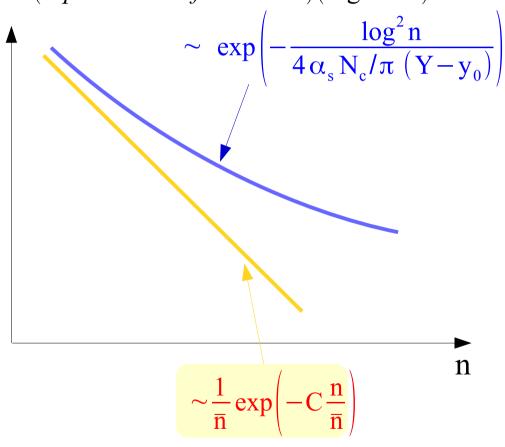


But in pA, the "dipole" is a proton!

#### log n Realistic model for pA



proba(n particles in final state)(log scale)



Consistent with the data?

Work in progress: How much does C depend on the details of confinement?

#### Summary

- At high energies, hadrons look like dense states of gluons (sometimes called "color glass condensates"), very far from the valence picture. This is a property of QCD.
- The (stochastic) evolution of hadronic wave functions towards high energy can be computed in QCD. The color dipole model is a convenient implementation of this evolution.
- Fluctuations of the multiplicity in pA scattering in the proton fragmentation region can be related to the event-by-event fluctuations of the total integrated gluon density in the proton! pA data at the LHC is a great opportunity to study these fluctuations!

#### Outlook

- Better understand the fluctuations in the dipole model+confinement test robustness of our solution
- Try and build a realistic model for phenomenology (proton instead of onium etc...)