

Littlest Higgs Model and W Pair Production at ILC

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- Alternative models include composit higgs models, models with strongly interacting EW sector, higgsless models, etc.

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- The **EWSB** mechanism leaves its signature in the gauge sector of the theory, for the massive gauge bosons get their longitudinal components from the symmetry breaking sector.
- Proposed high energy linear collider (ILC) is expected to produce a large number of gauge boson pairs, where precision measurements are going to scrutinize their couplings, and other properties in detail.

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Plan

- **The Little Higgs Models**
- **The Process:** $e^+e^- \rightarrow W^+W^-$
- **Summary**

The Little Higgs Models

- The scenario is analogous to the description of **low energy hadronic interactions** by a non-linear realisation of the chiral symmetry, $SU(2)_L \times SU(2)_R$, broken down to the diagonal isospin $SU(2)_I$ with **pions** as pseudo Goldston Bosons.

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- In a specific model, called **Littlest Higgs model** $G \equiv SU(5)$ broken down to $H \equiv SO(5)$ at a scale f . This gives $24 - 10 = 14$ Goldstone Bosons.
- Gauge group $SU(2)_1 \times SU(2)_2 \times U(1)_Y \subset G$, is broken down to $SU(2)_L \times U(1)_Y$, which is identified as the SM gauge group.

The Little Higgs Models

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- $SU(2)_L \times U(1)_Y$ breaking is achieved by the vev, v of the GB's belonging to complex 2 , induced radiatively by Coleman-Weinberg mechanism.
- Quadratically divergent corrections to the higgs mass due to the SM gauge bosons cancel with the new heavy gauge boson contributions.

The Little Higgs Models

- Yukawa couplings:
In order to avoid quadratic divergence due to top quark loop, a pair of Weyl quark \tilde{t} , \tilde{t}^c and a weak singlet quark u_3^{lc} are added, along with the standard left-handed doublet, (t_3, b)
- The top Yukawa coupling terms are given in terms of the mass eigenstates t (the SM top quark with mass m_t) and T (with mass $M_T \geq \sqrt{2} f$) as

$$\mathcal{L}_Y = \lambda_t h t^c t + \lambda_T h T^c t + \frac{\lambda'_T}{2M_T} h h T^c T + h.c.,$$

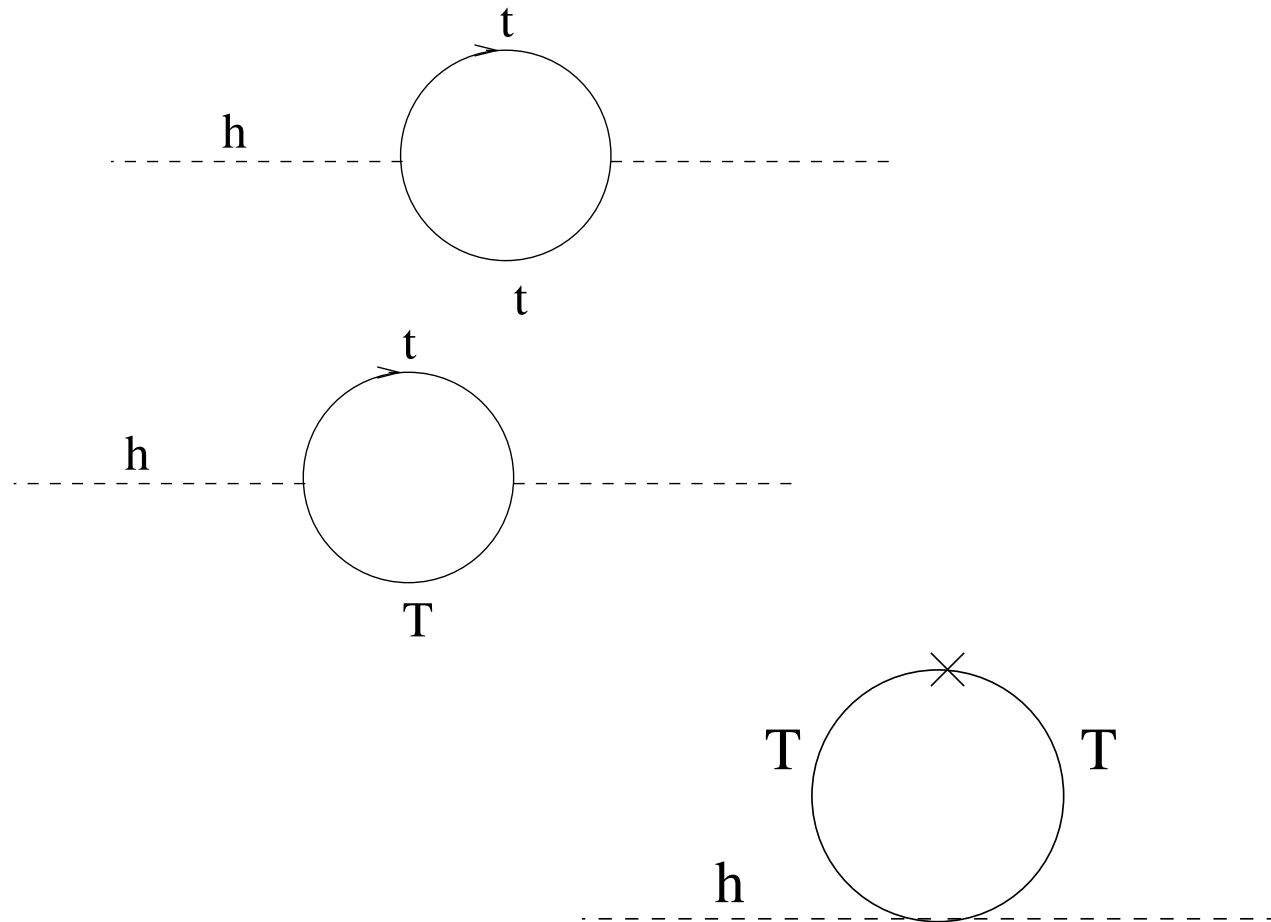
where

$$\lambda_t = m_t/v, \quad \lambda_T = x_\lambda m_t/v, \quad \lambda'_T = (1 + x_\lambda^2) m_t^2/v^2.$$

x_λ is a parameter (of order 1).

The Little Higgs Models

- Thus at one loop the quadratic contributions cancel between the following diagrams:



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- Real production of heavy gauge bosons occur above a TeV. But their virtual effects may be seen at lower energies.

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- Coupling relevant to this case are given in terms of these parameters f and θ as

$$g_{WWZ_H} = \frac{gv^2}{8f^2} \sin 4\theta$$

$$c_{eeZ}^v = \frac{g}{2c_W} \left[\left(-\frac{1}{2} + 2x_W \right) + \frac{v^2}{f^2} \frac{\sin 4\theta}{8} \right]$$

$$c_{eeZ}^a = \frac{g}{2c_W} \left[\frac{1}{2} - \frac{v^2}{f^2} \frac{\sin 4\theta}{16} \cot \theta \right]$$

$$c_{eeZ_H}^v = \frac{-g}{4} \cot \theta$$

$$c_{eeZ_H}^a = \frac{g}{4} \cot \theta$$

The Gauge Sector

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Arkani-Hamed, et.al. *JHEP* **07**(2002) 034;

Han, Logan, Wang, *hep-ph/0506313*;

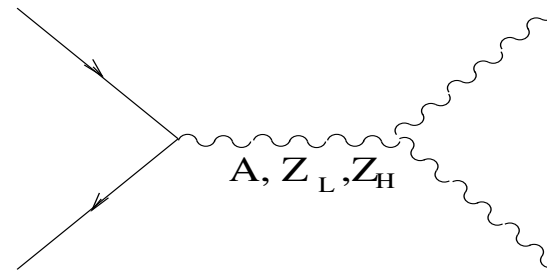
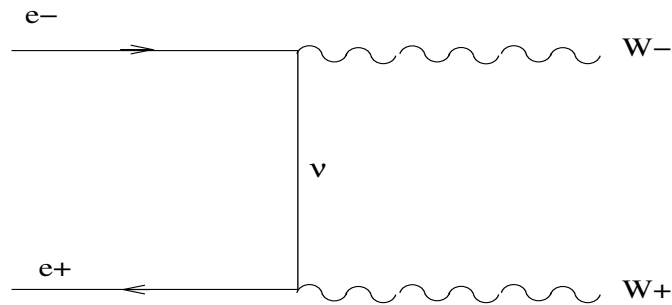
Schmalts, Tucker-Smith, *hep-ph/0502182*;

Conley, Hewet, Phuong Le, *hep-ph/0507198*.

$$e^+e^- \rightarrow W^+W^-$$

We will now consider the effect of this scenario in W pair production at a high energy linear e^+e^- collider.

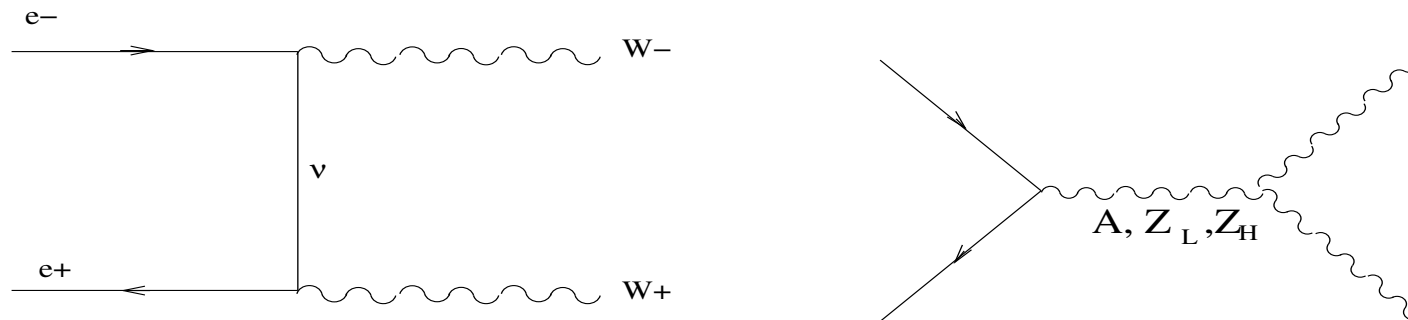
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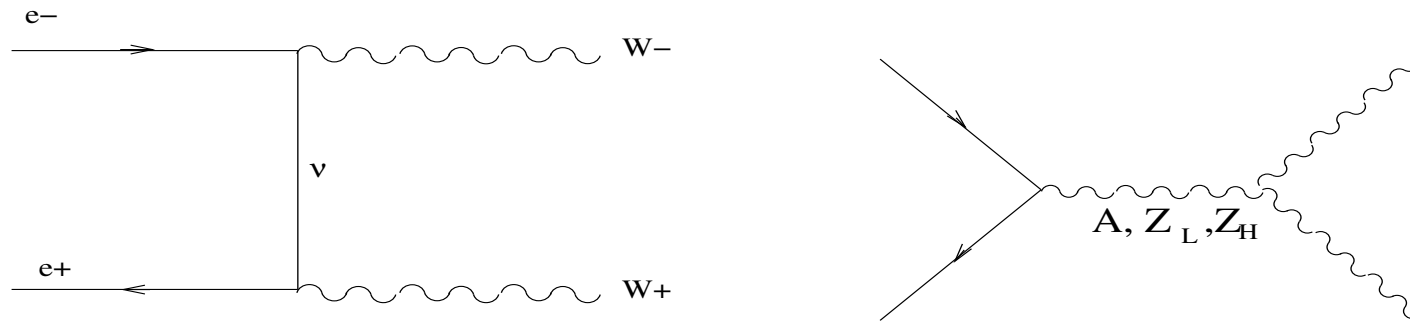


- **ILC** is going to be operated at energies up to a TeV. A Z_H with mass around a TeV will perhaps show its signs there.

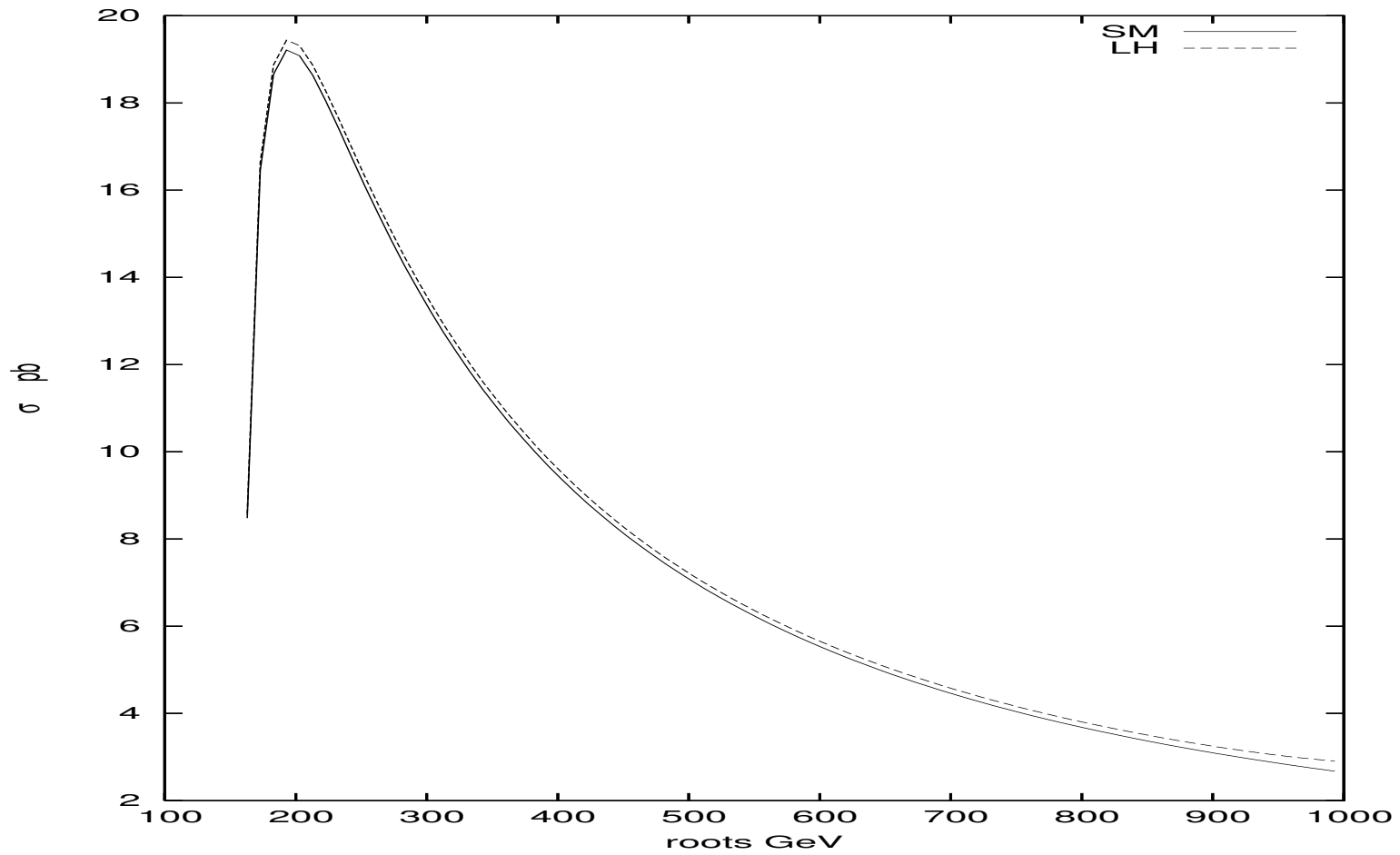
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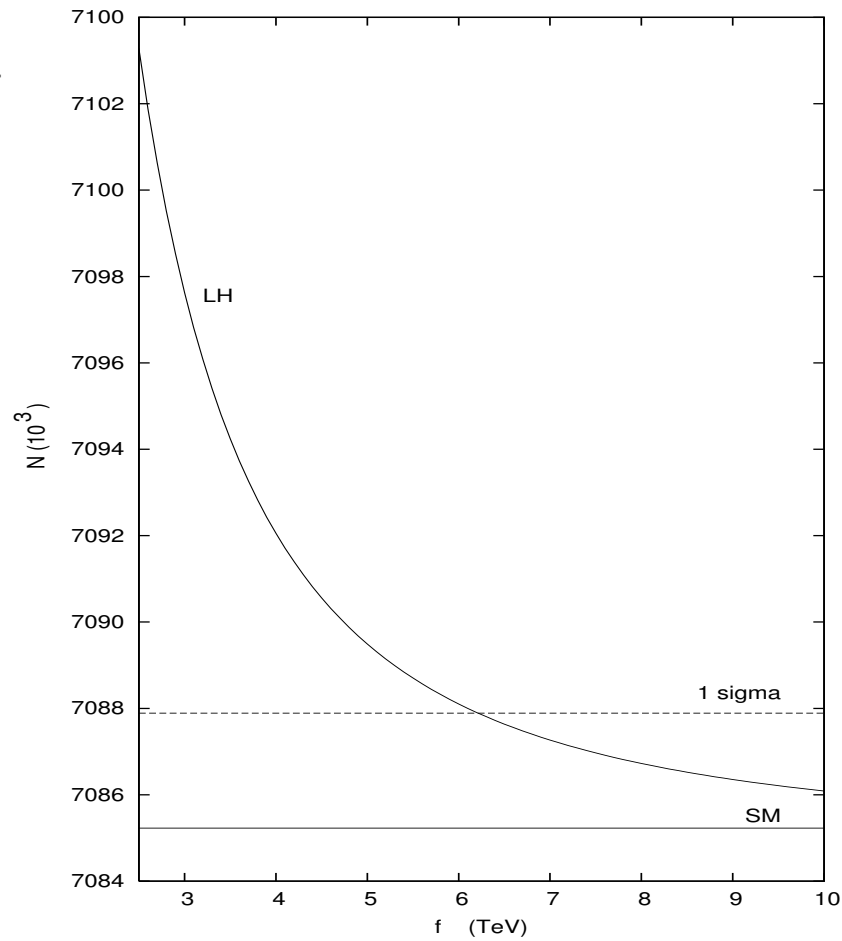
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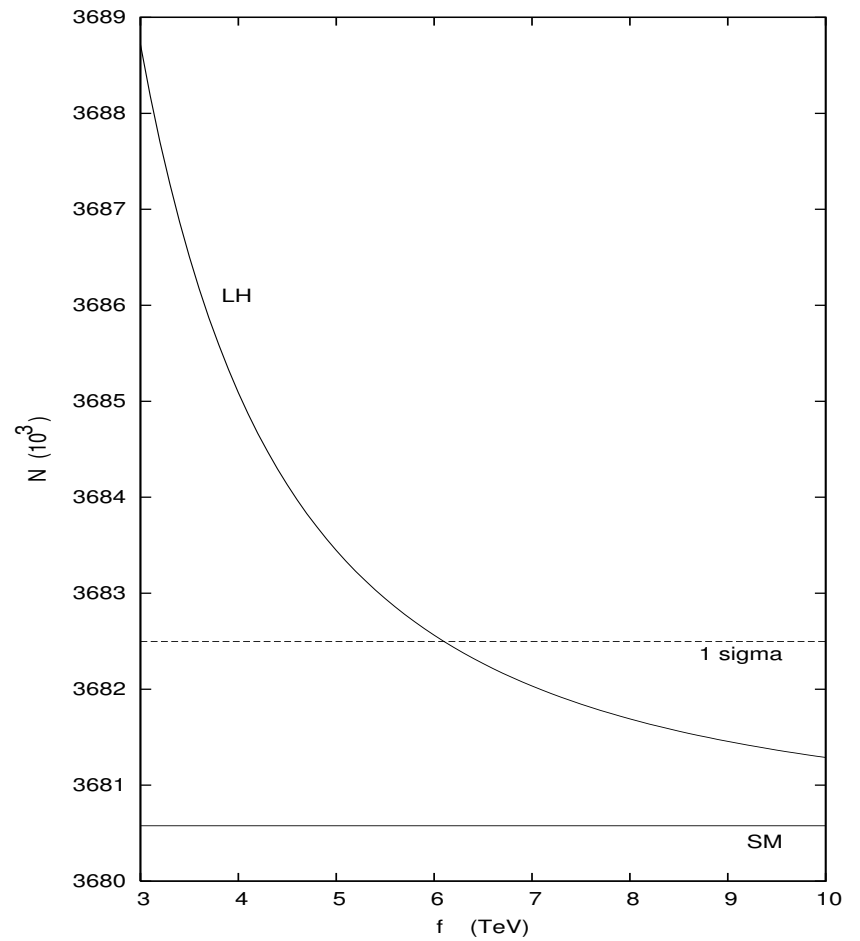
- **ILC** is going to be operated at energies up to a TeV. A Z_H with mass around a TeV will perhaps show its signs there.
- Owing to the large luminosity expected at the ILC, cross section measurement itself might be a good observable.



\sqrt{s} GeV	σ_{SM} pb	σ_{LH} pb
500	7.1	7.2
800	3.7	3.8



$$\sqrt{s} = 500 \text{ GeV}$$



$$\sqrt{s} = 800 \text{ GeV}$$

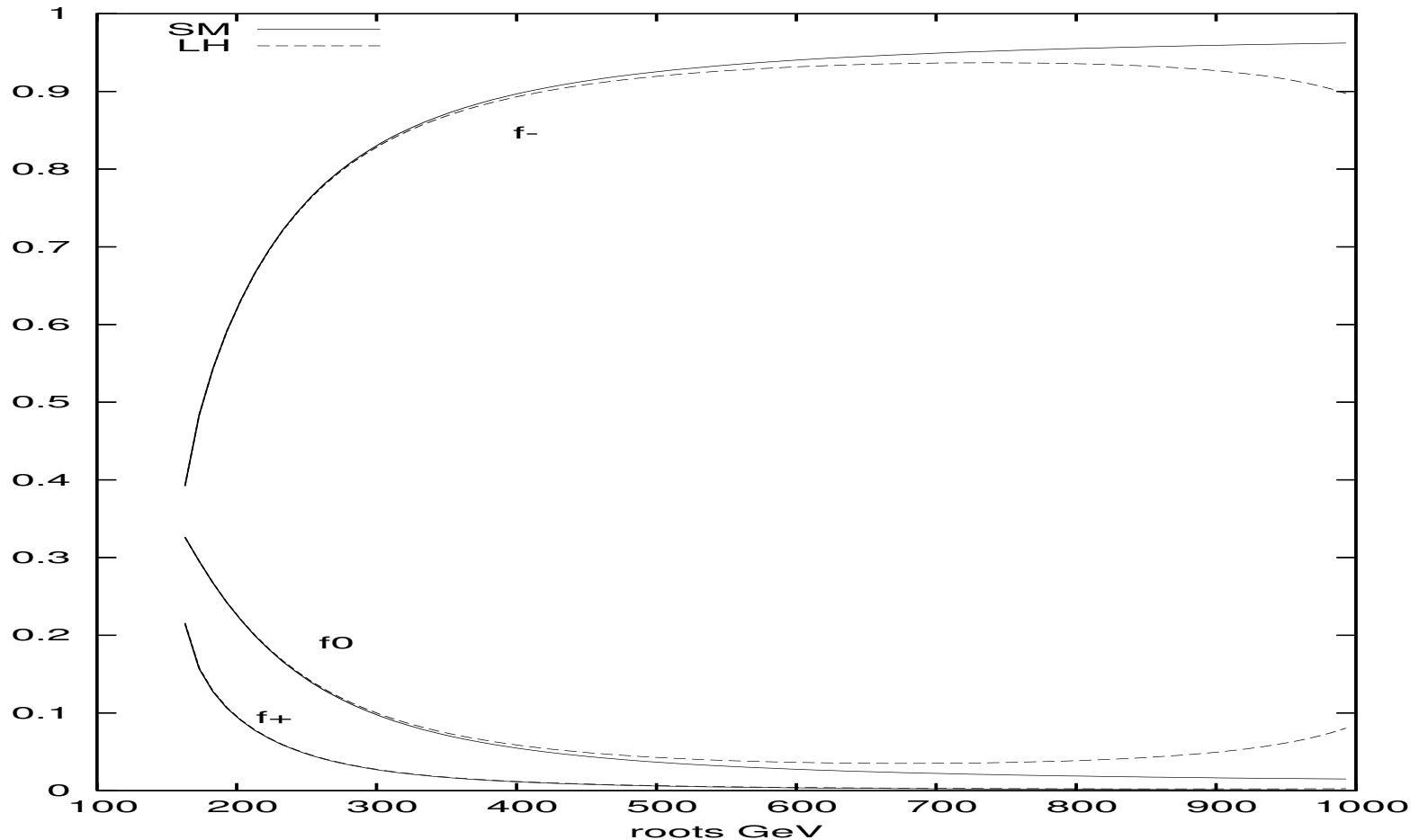
Integrated luminosity: 1 ab^{-1}

1σ limit on f : $\sim 6 \text{ TeV}$

LEP has measured the **fractional cross section of the polarised W 's** very precisely. AT ILC it will be even better.

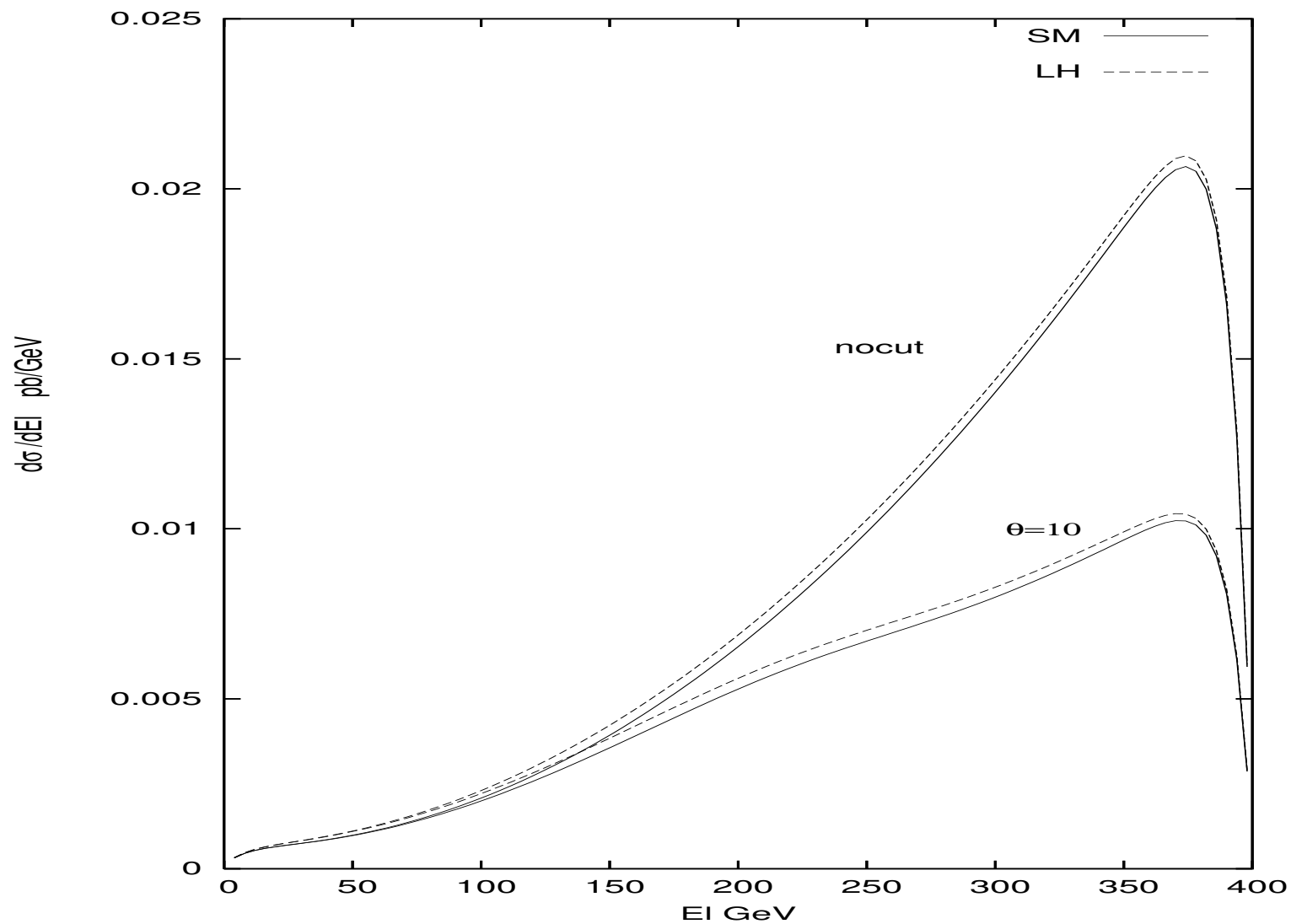
We define

$$f^0 = \frac{1}{\sigma_{unpol}} \sigma(e^+e^- \rightarrow W^+W_L^-), \text{ etc.}$$



\sqrt{s} GeV	σ_{SM} pb		$f_{SM}\%$	$f_{LH}\%$
500	7.1	f^0	3.8	4.4
		f^-	95.6	94.9
		f^+	0.6	0.7
800	7.2	f^0	1.9	4.0
		f^-	97.9	95.8
		f^+	0.2	0.2

Energy distribution of the decay leptons.



Summary

- Little Higgs model with a global $SU(5) \rightarrow SO(5)$ and a local $SU(2)_1 \times SU(2)_2 \times U(1)_Y \rightarrow SU(2)_L \times U(1)_Y$ is a viable alternative to the Standard Higgs mechanism.
- The model has
 - one light neutral higgs with radiatively stable mass.
 - two charged (W_H^\pm) and one neutral (Z_H) gauge bosons with mass $\sim f$, the symmetry breaking scale.
 - one heavy top quark of mass $\sim \sqrt{2}f$.

- Our analysis to see the effect of this model in $e^+e^- \rightarrow W^+W^-$ at a high energy linear collider shows:
 - Considering 1σ deviation at a 500 GeV or 800 GeV ILC with integrated luminosity of 1 ab^{-1} can explore the mass scale up to about 6 TeV .

For parameter values

$$f = 1 \text{ TeV and } \theta = 0.35$$

- cross section deviates by 1.2% at $\sqrt{s} = 500 \text{ GeV}$ and 3.4 % at $\sqrt{s} = 800 \text{ GeV}$
 - polarization fractions show measurable deviations.
 - Beam polarization is expected to show larger effects.
- Studies in progress.